

Oblique correction from sextet QCD

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Geneva, July 2010

Outline

- Conformality “for all practical purposes”
- Precision electroweak tests for fun and profit
- The lattice calculation – usual overlap miracles
- QCD games to check techniques
- Oblique correction from sextet QCD
- Conclusions

Work based on big program with B. Svetitsky and Y. Shamir

arXiv:1006.3777

Background

- Lots of interest in BSM systems by lattice community
- TD+Svetitsky + Shamir's system – $SU(3)$ gauge group with $N_f = 2$ flavors of sextet fermions
- Lattice discretization – clover fermions
- Weak coupling phase (in finite volume) is chirally restored, deconfined
- No stable $am_q^{AWI} = 0$ in strong coupling
- Don't know if it has an IR fixed point
- Certainly it is a theory whose coupling runs very slowly
- Let's use the fact that it walks to do a calculation

Physics question: What would the S-parameter look like, for nearly conformal dynamics?

Target: “conformal technicolor” or “unparticle” or “hidden valley” fans

Conformality “for all practical purposes”

- Schrödinger functional coupling runs really slowly
- In any lattice volume, coupling at shortest and longest distance varies by $\sim 10 - 15$ per cent
- This is slow enough that it's hard to see it run
- Slow running is approximately no running
- No running means that the quark mass is the relevant operator, $m = 0$ is critical
- This gives several ways to measure the mass anomalous dimension
 - Finite size scaling analysis of correlation length (TD, November)
 - Schrödinger functional running mass (DSS, June)
- Gives a nice measurement of $\gamma_m(g^2)$ which is
 - Unfortunately, fatally low for TC phenomenology ($\gamma_m = 1$ is desired)
 - Still the world's record for biggest γ_m from a simulation (as of June 2010)

Some pictures:

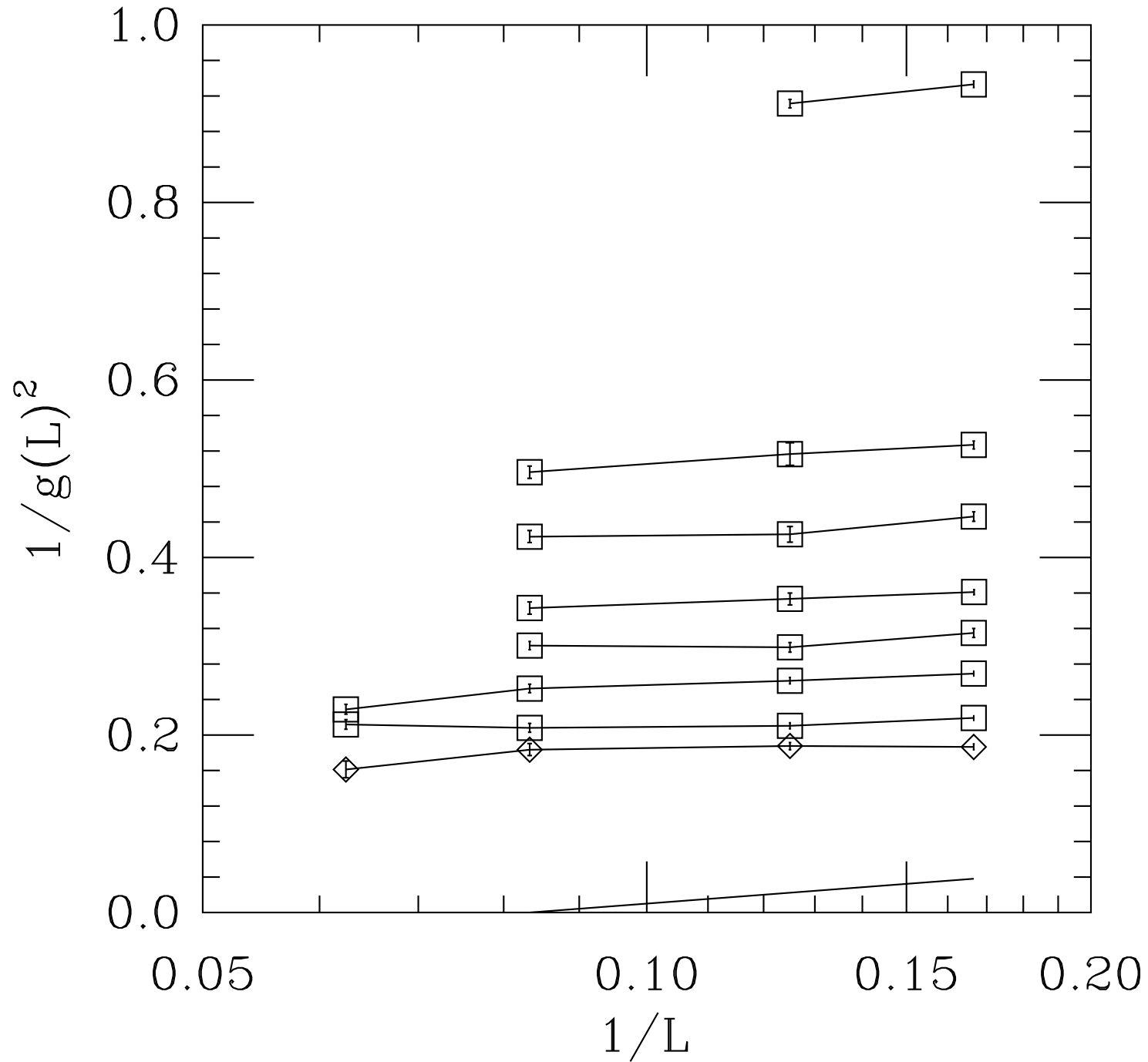


Figure 1: SF coupling $1/g^2$ vs. a/L . Is $1/g^2 = 2b_1/(16\pi^2) \log L + \text{constant}$?

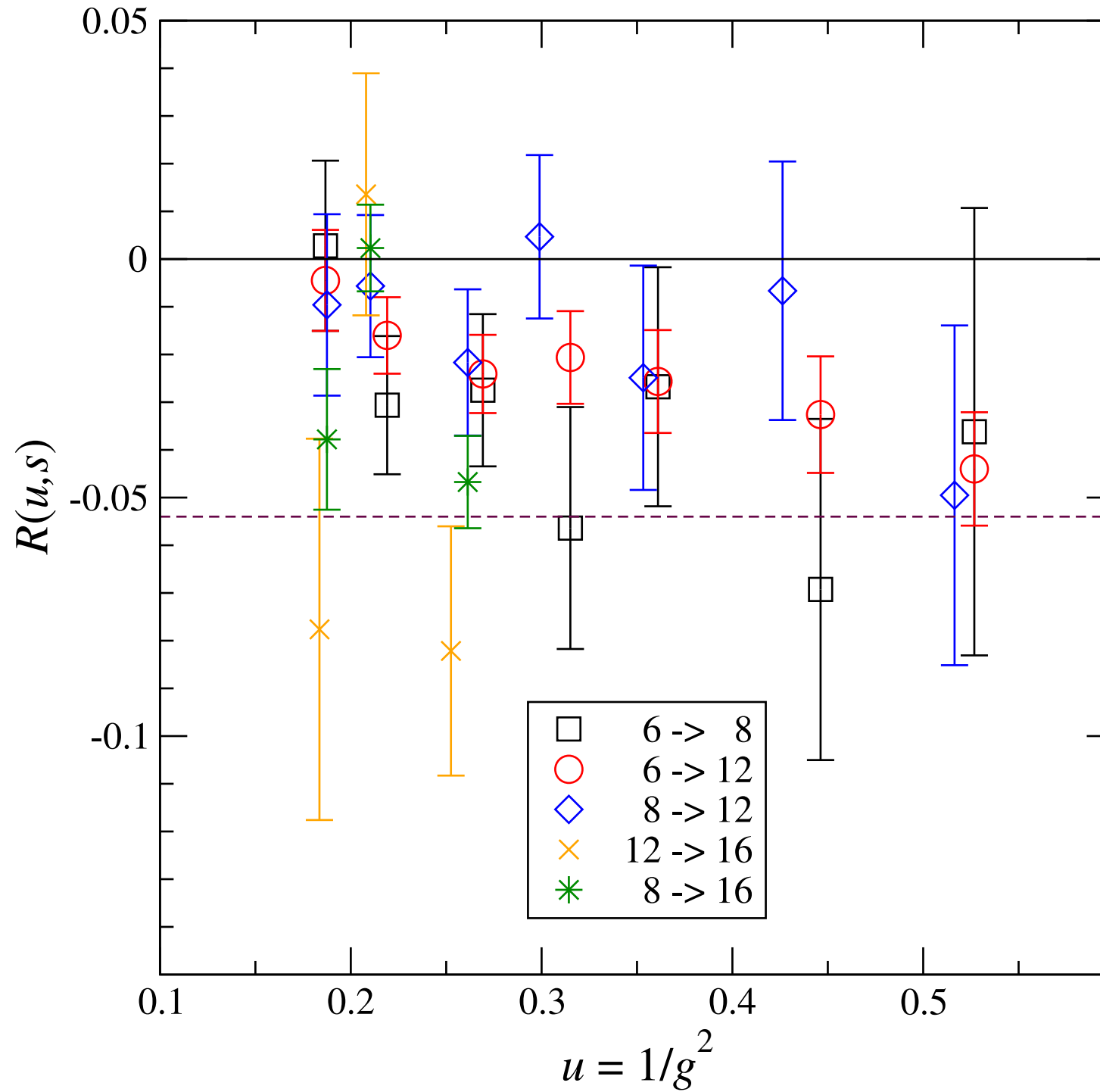


Figure 2: Lattice approximants to the beta function for many scale factors

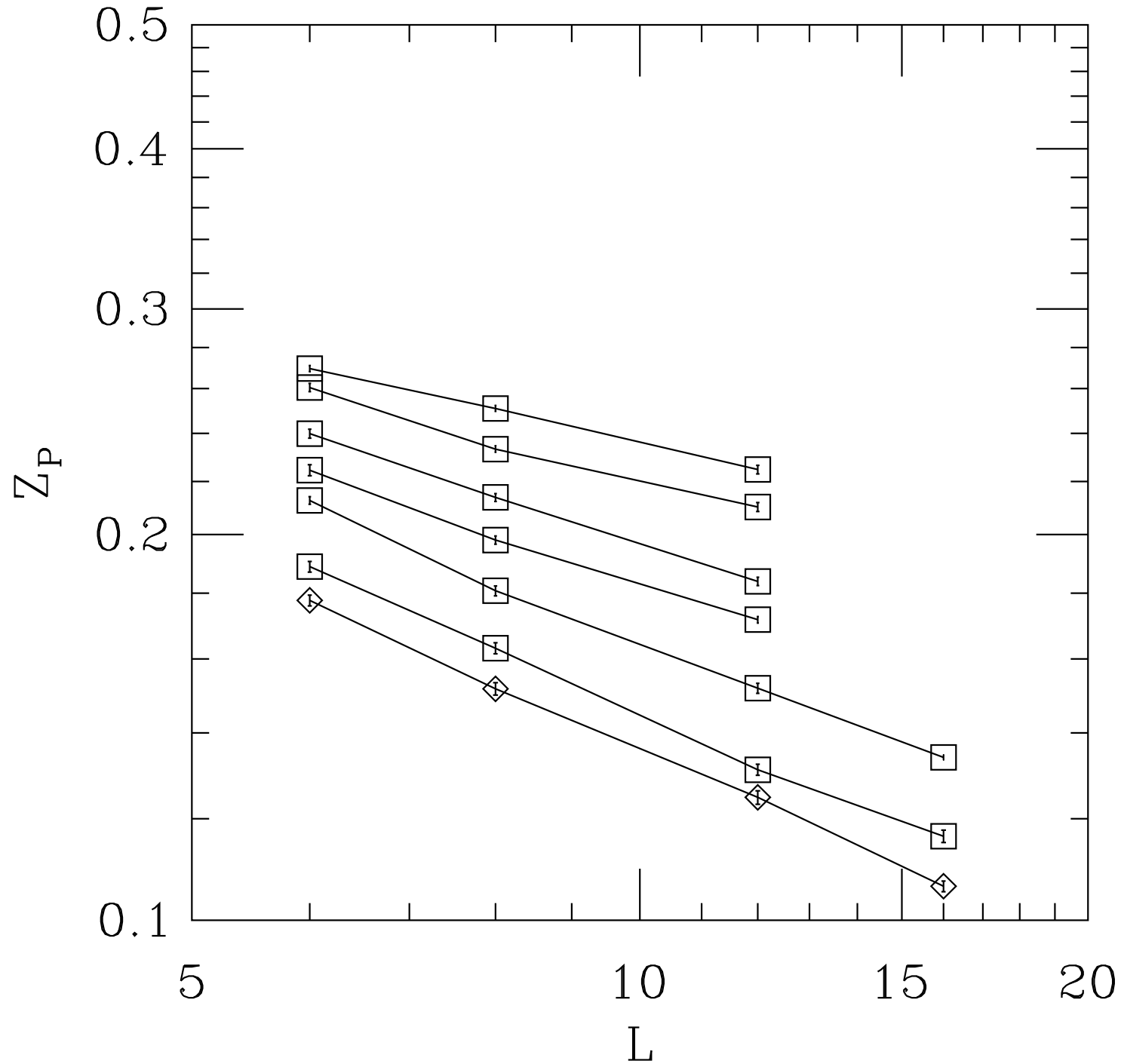
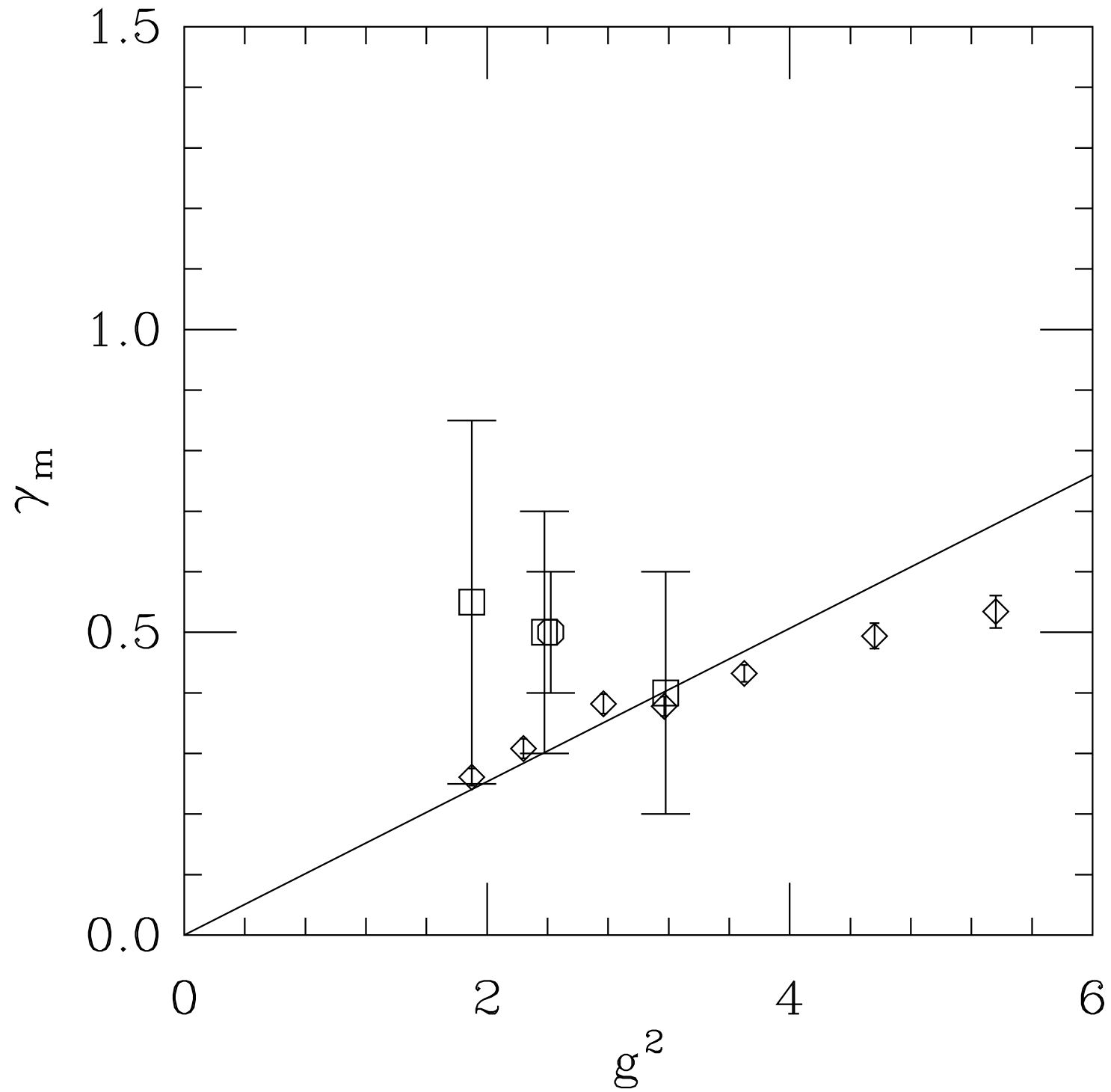


Figure 3: Pseudoscalar renormalization constant Z_P vs. L/a . Is $Z_P \propto L^{-\gamma_m}$?



Precision EW Background

“Oblique correction” to gauge boson vacuum polarization given by current - current correlator

$$\begin{aligned}\Pi_{\mu\nu}(q) &= \int d^4x \exp(iqx) \langle J_\mu^L(x) J_\nu^R(0) \rangle \\ &\equiv (q^2 \delta_{\mu\nu} - q_\mu q_\nu) \Pi_T^{LR}(q^2) + q_\mu q_\nu \Pi_L^{LR}(q^2) \dots\end{aligned}\tag{1}$$

Peskin Takeuchi S-parameter is $S = 16\pi d(q^2 \Pi_T^{LR}(q^2)) / dq^2$ at small q^2 ,

Lattice problems

- Lattice dirt in

$$\Pi_{\mu\nu}(q) = P_{\mu\nu}^T(q) \Pi^T(q) + P_{\mu\nu}^L(q) \Pi_L(q) + \dots\tag{2}$$

- Extra quadratic divergence if lattice currents aren't conserved
- Nonlocal currents have contact terms
- Different Z_V and Z_A

Valence overlap fermions cure the last three!

Wonderful overlap miracles

V and A currents related by Ward identity so

- $Z_V = Z_A$
- quadratic term cancels
- use local currents (actually “improved” ones)

Decompose with $\bar{q}_\mu = (2/a) \sin q_\mu a/2$

$$\Pi_{\mu\nu}(q) = P_{\mu\nu}^T(q)\Pi_T(q) + P_{\mu\nu}^L(q)\Pi_L(q) \quad (3)$$

where

$$P_{\mu\nu}^T(q) = \bar{q}^2 \delta_{\mu\nu} - \bar{q}_\mu \bar{q}_\nu \quad (4)$$

and

$$P_{\mu\nu}^L(q) = \bar{q}_\mu \bar{q}_\nu \quad (5)$$

Check decomposition with JLQCD operator

$$\Delta_J(q) = \sum_{\mu\nu} \bar{q}_\mu \bar{q}_\nu \left(\frac{1}{\bar{q}^2} - \frac{\bar{q}_\nu}{\sum_\lambda \bar{q}_\lambda^3} \right) \Pi_{\mu\nu}^J(q) \quad (6)$$

Force fit to trial decomposition, P^T and P^L are projectors... q by q

QCD fun & games

Did this for quenched and $N_f = 2$ QCD. All looks “normal.”

Pause for physics: resonance dominance (basically large N_c) writes

$$\Pi_T^{LR}(q^2) = \sum_V \frac{f_V^2 M_V^2}{q^2 + M_V^2} - \sum_A \frac{f_A^2 M_A^2}{q^2 + M_A^2} - \frac{f_\pi^2}{q^2}. \quad (7)$$

Weinberg sum rules:

$$\sum_V f_V^2 M_V^2 - \sum_A f_A^2 M_A^2 - f_\pi^2 = 0 \quad (8)$$

$$\sum_V f_V^2 M_V^4 - \sum_A f_A^2 M_A^4 = 0 \quad (9)$$

Usual additional approximation: saturate with lowest resonances, π , ρ , a_1 (5 parameters)

I can

- Measure all 5 in a simulation – this doesn't reproduce WSR's – or data
- Measure f_π , m_ρ , f_ρ , use WSR's to compute f_{a_1} , m_{a_1} – looks like data

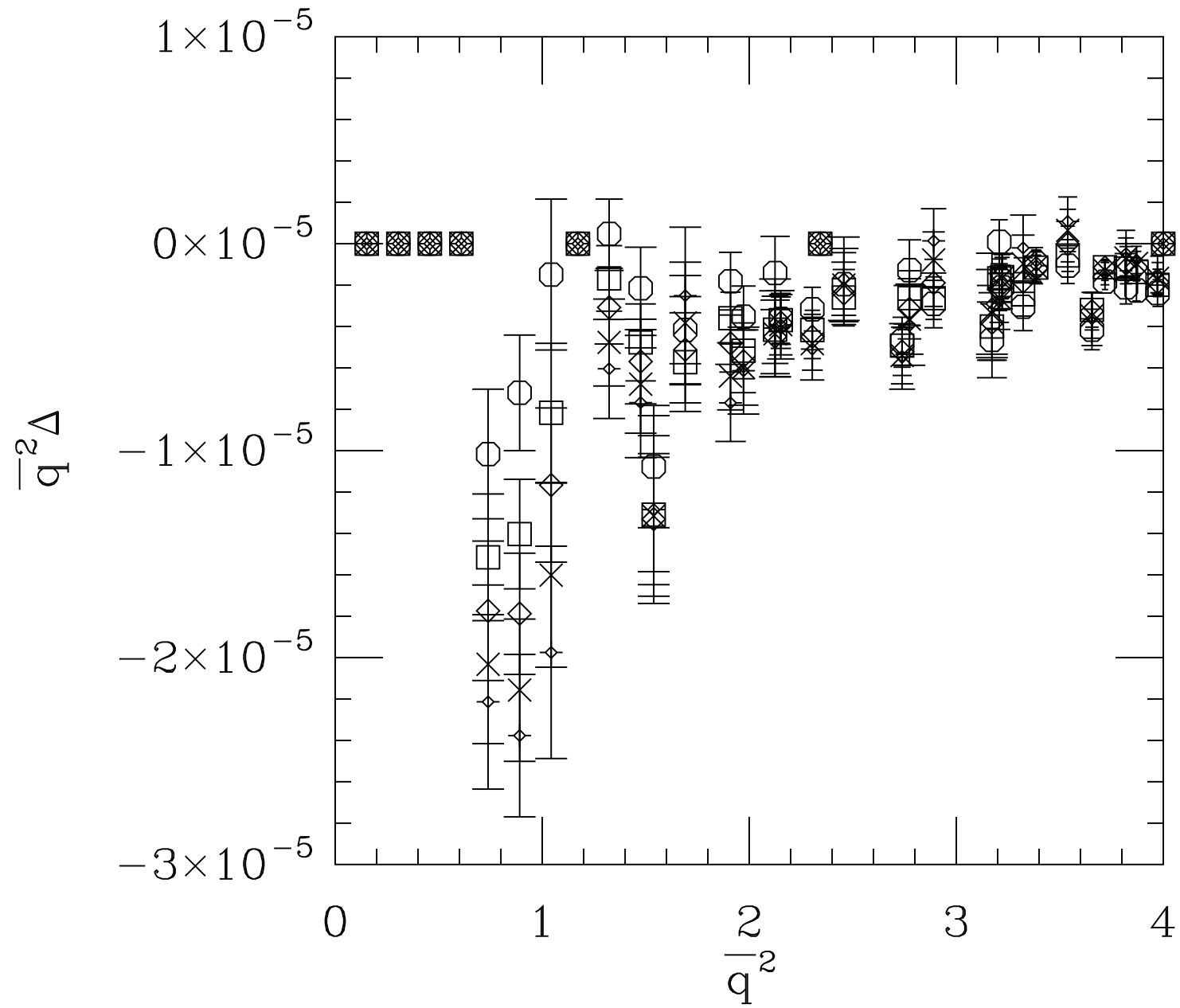


Figure 5: JLQCD parameter for quenched QCD

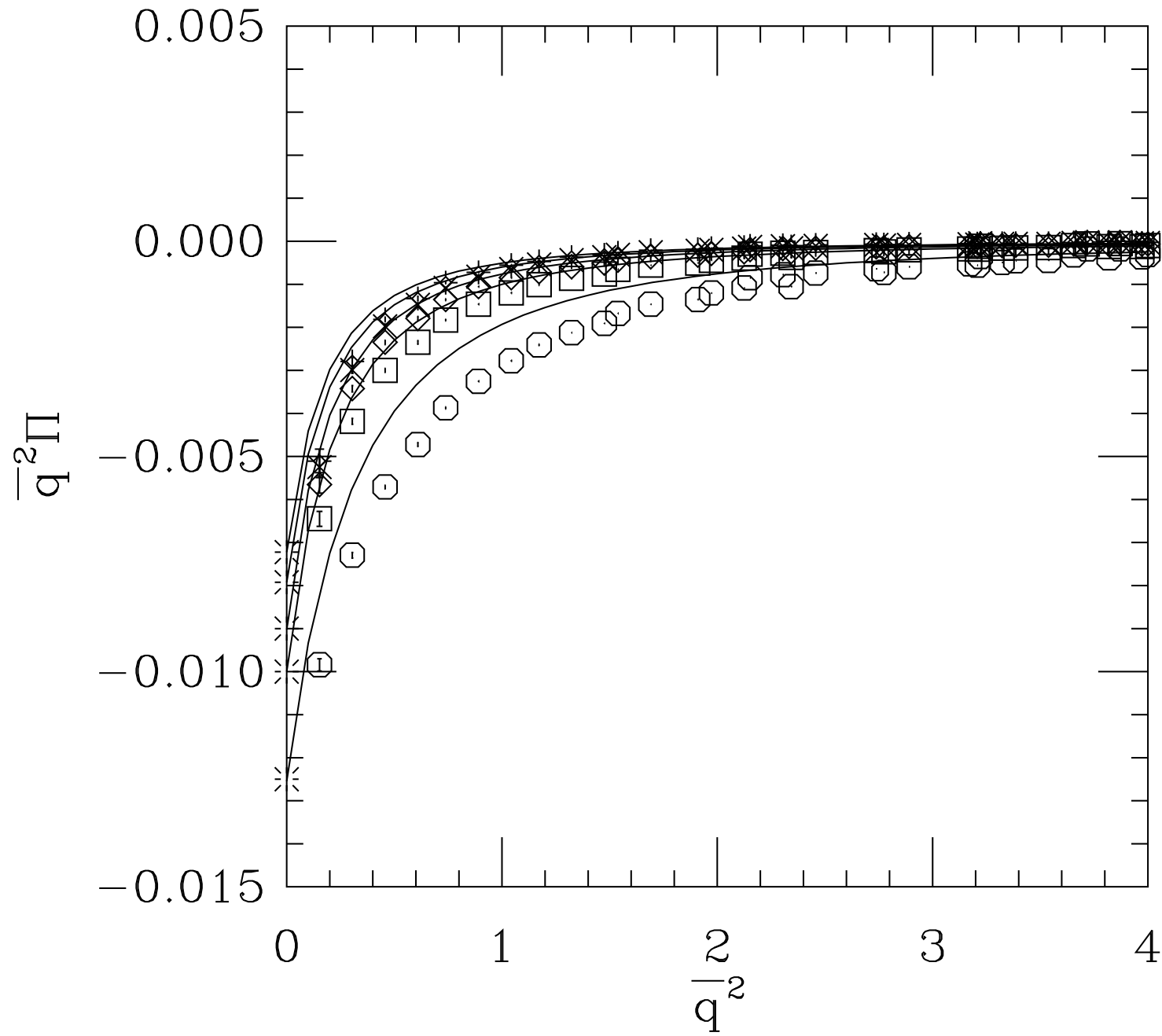


Figure 6: Π_{LR}^T for quenched QCD – lines are low state model

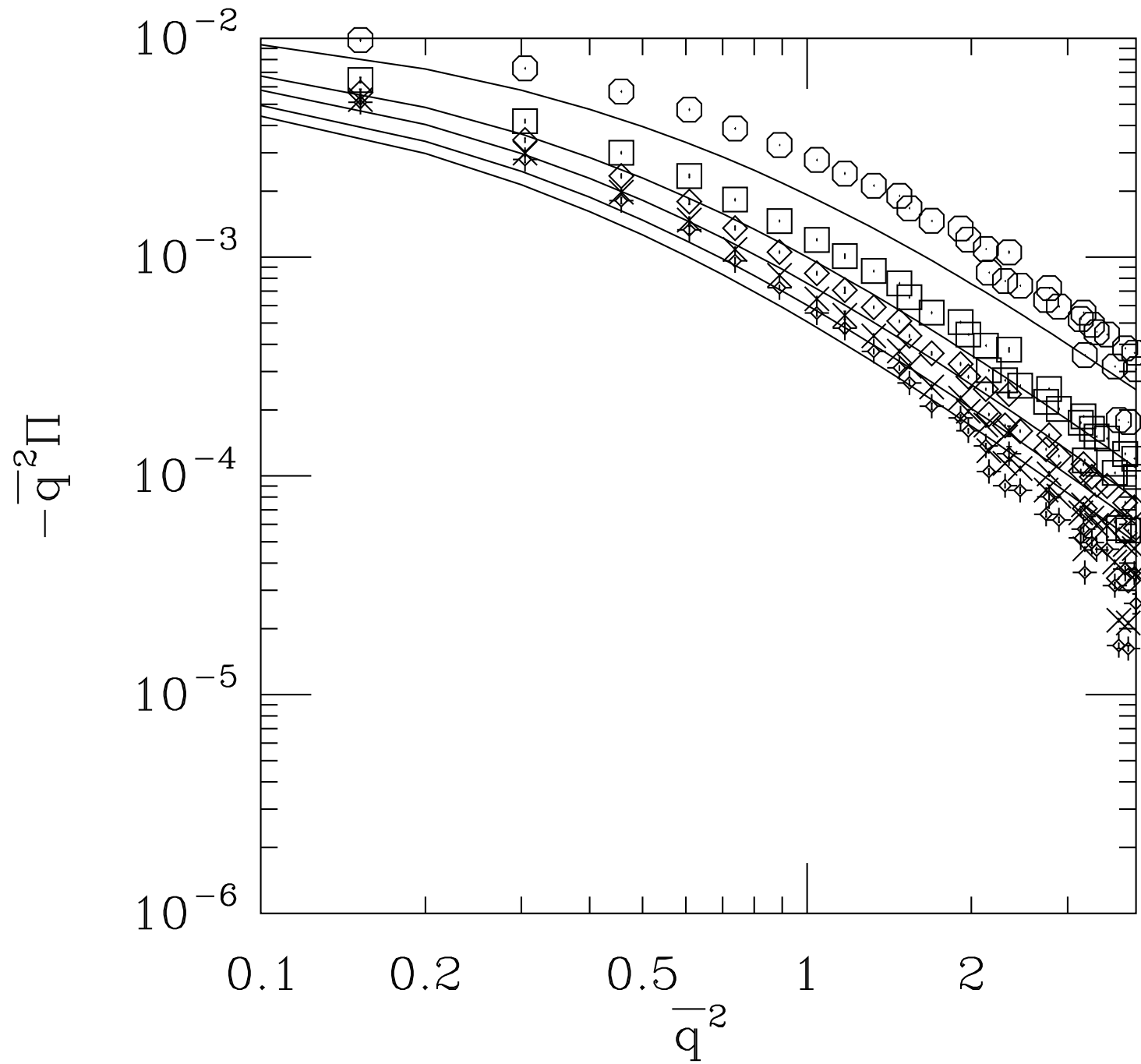


Figure 7: Π_{LR}^T for quenched QCD – log-log plot – lines are low state model

L_{10} and Δm_π^2

With "lowest mass dominance" plus Weinberg sum rules one can predict

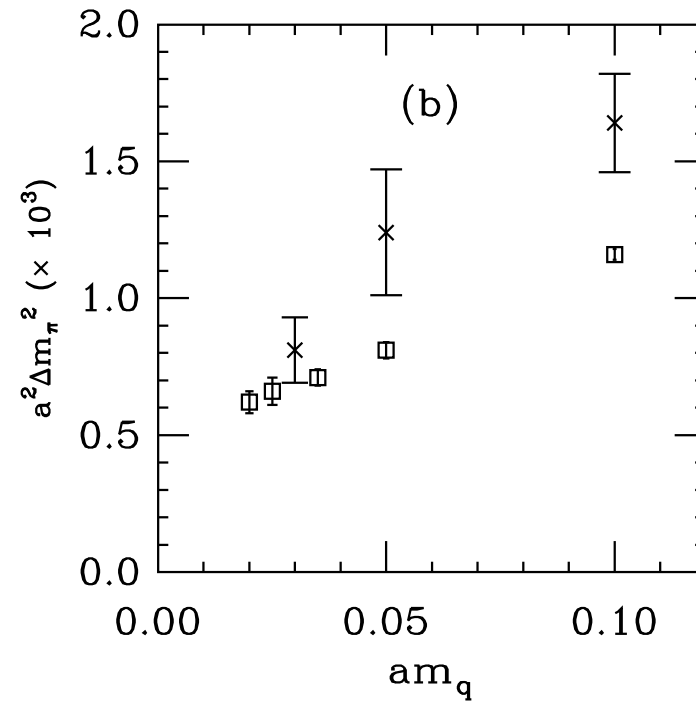
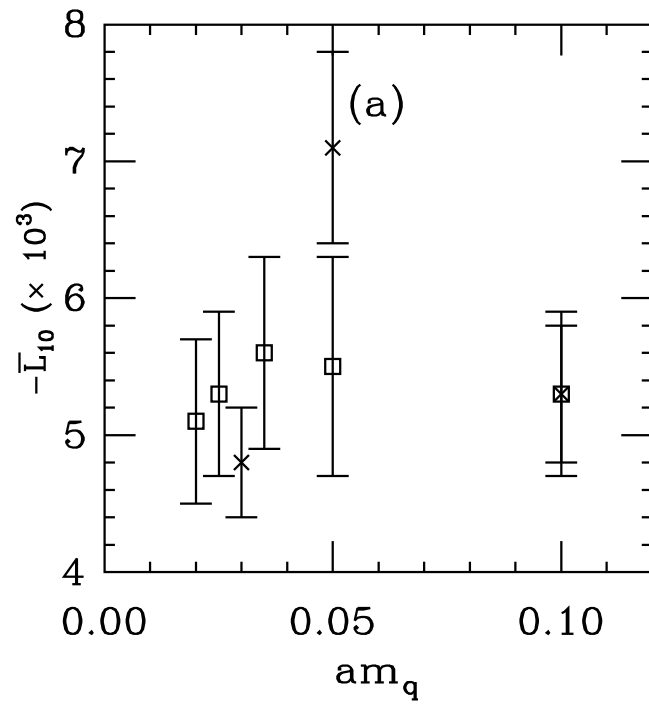
$$L_{10} = f_\rho^2 - f_{a_1}^2 \quad (10)$$

Das, Guralnik, Mathur, Low, Young sum rule – just do the integral

$$\Delta m_\pi^2 = \int dq^2 \Pi_{LR}^T(q^2) = \frac{3\alpha}{4\pi} \frac{1}{\left(\frac{1}{m_\rho^2} - \frac{1}{m_{a_1}^2}\right)} \log \frac{m_{a_1}^2}{m_\rho^2} \quad (11)$$

Works pretty well, $L_{10} \sim 5 \times 10^{-3}$, $\Delta m_\pi^2 = 1100$ MeV.

This is just fooling around, to show $\Pi_{LR}^T(q^2)$ looks ordinary



S-parameter for our conformal theory

Ran valence sextet overlap fermions on a set of 16^4 dynamical sextet clovers at one parameter value

$$(am_q^{AWI} = 0.04, \gamma_m = 0.35)$$

What I found:

- It doesn't look anything like QCD (but why should it?)
- $\Pi_{LR}^T(q^2)$ vanishes at large q^2 , vanishes as $m^2 \rightarrow 0$
- Saturation with lowest states fails!
 - Open any strong coupling BSM review
 - Everyone assumes low state saturation
 - Even true for papers with “conformal” in the title
- Finite differences for S-parameter – power law scaling

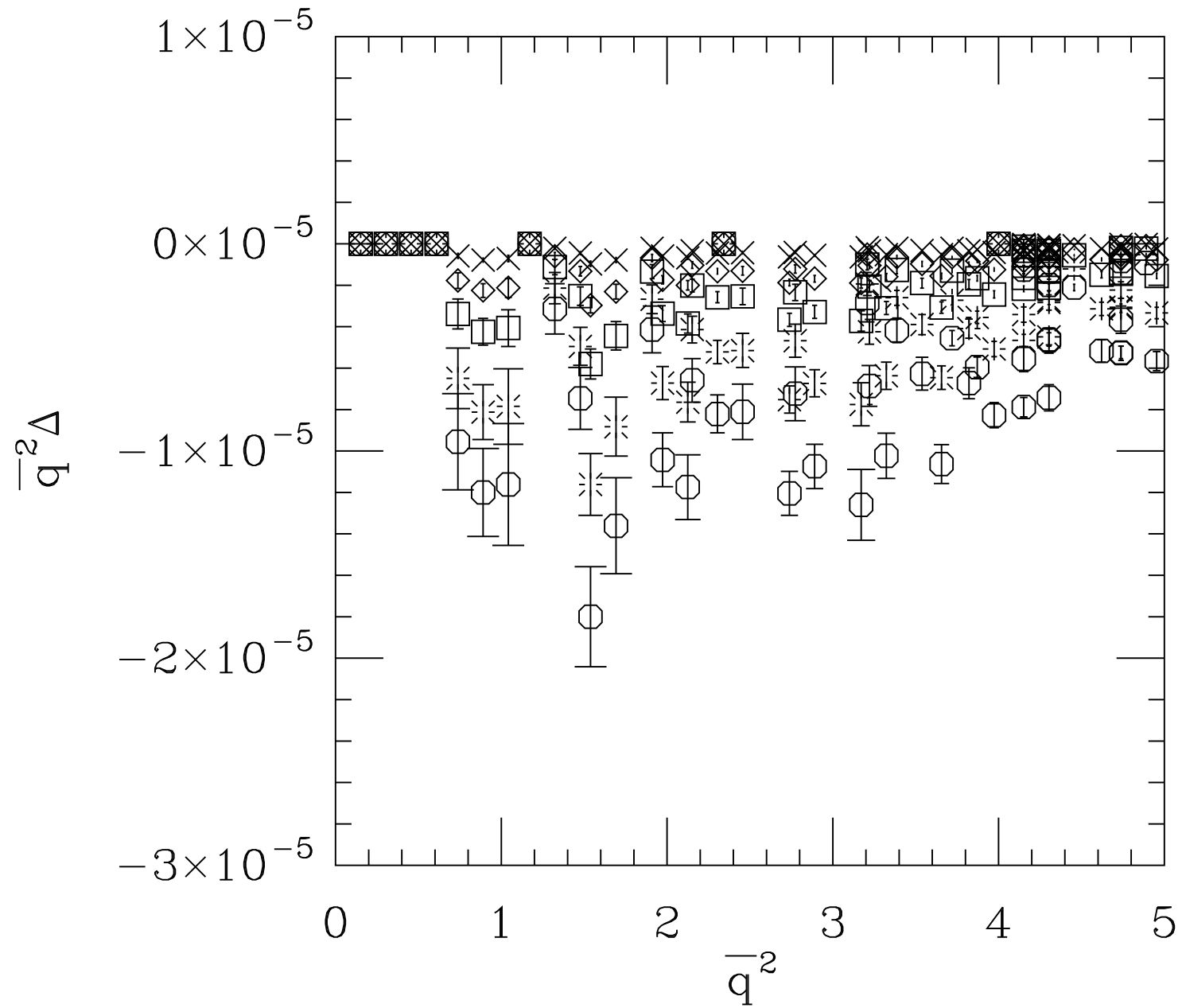


Figure 8: JLQCD parameter for sextet QCD

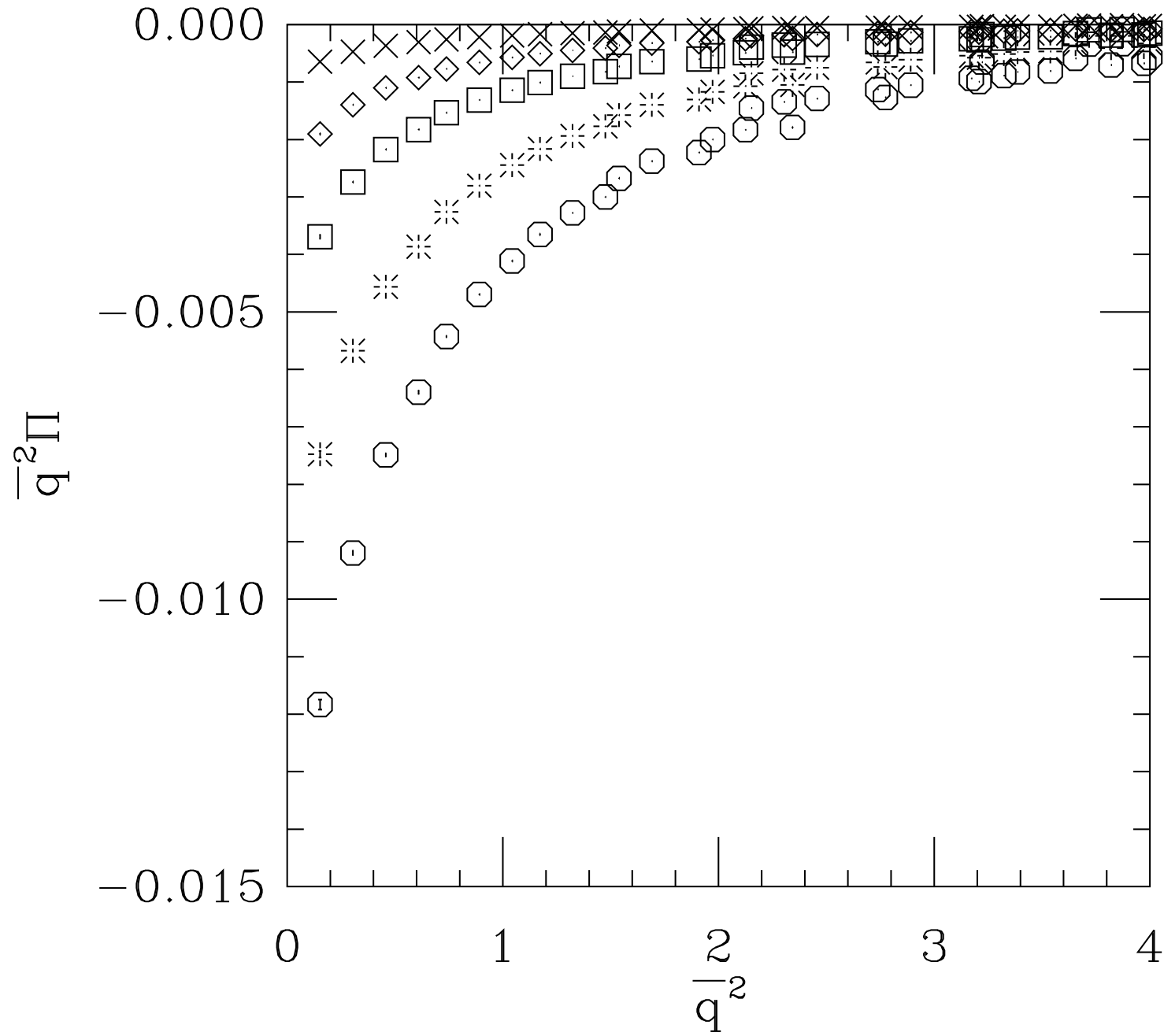


Figure 9: Π_{LR}^T for sextet QCD

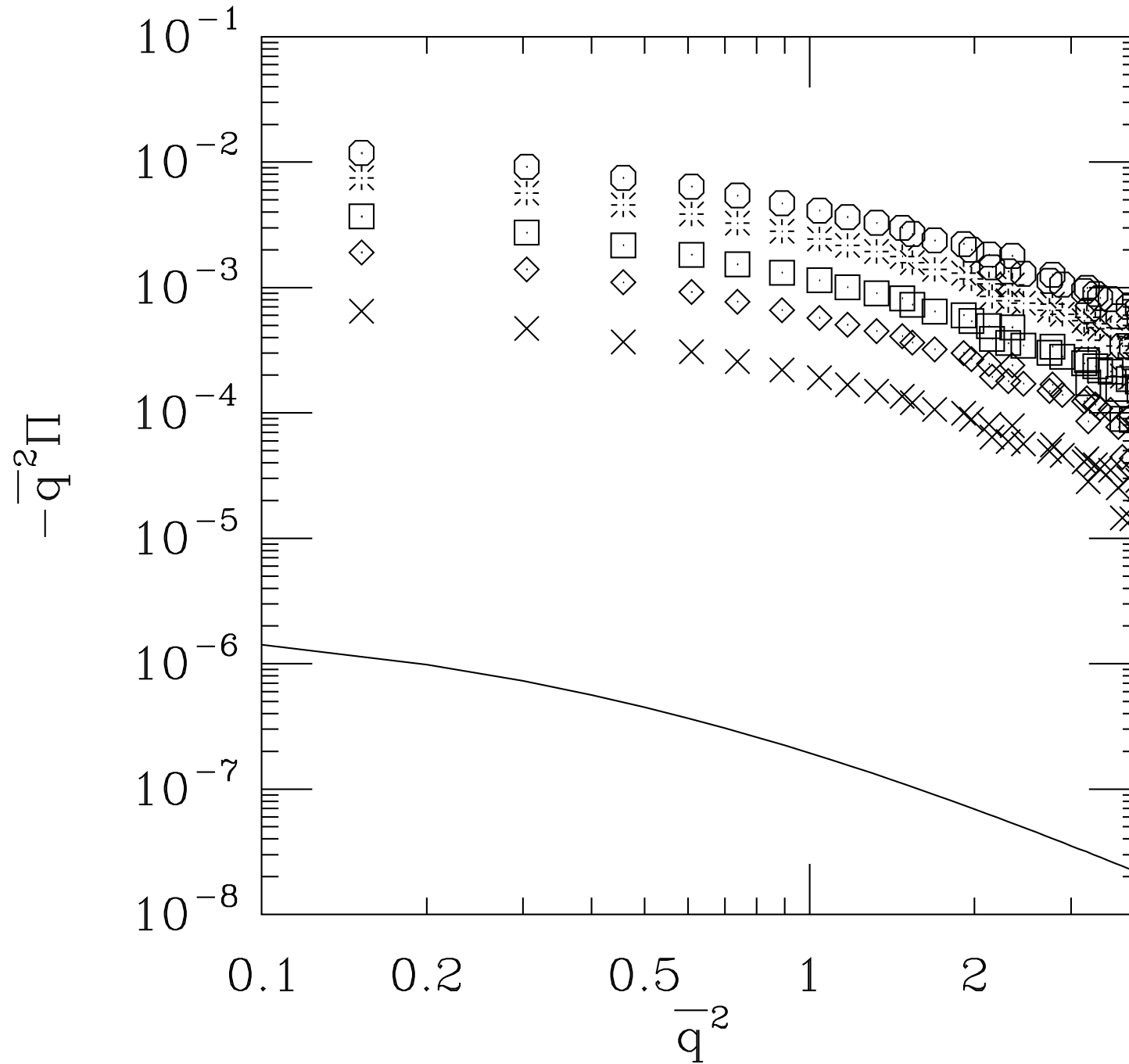


Figure 10: Π_{LR}^T for sextet QCD – log-log plot plus low state formula at one mass

the S-parameter

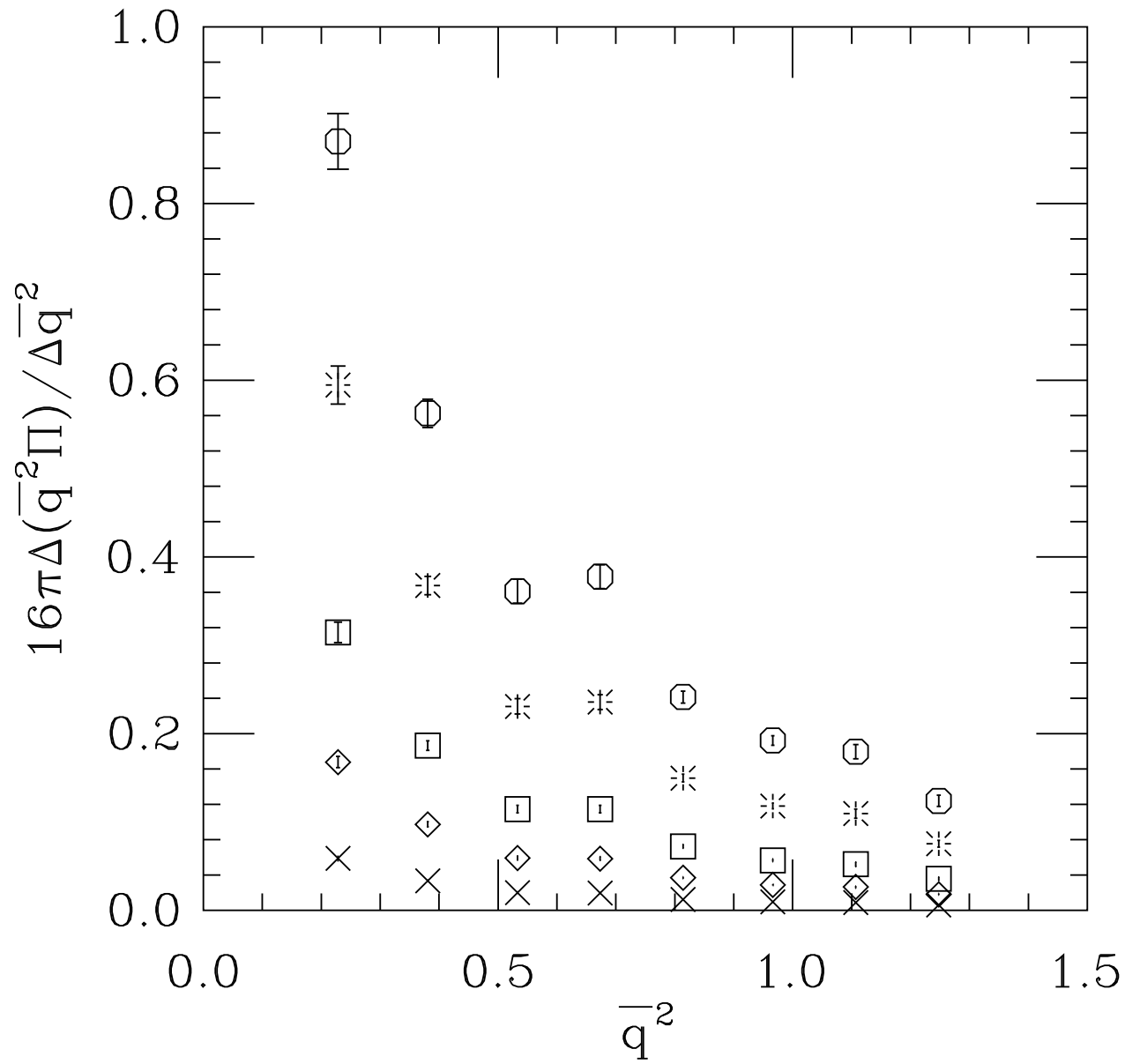


Figure 11: Numerically differentiated “S–parameter” for sextet QCD – log-log plot

the S-parameter vs q^2/m^2

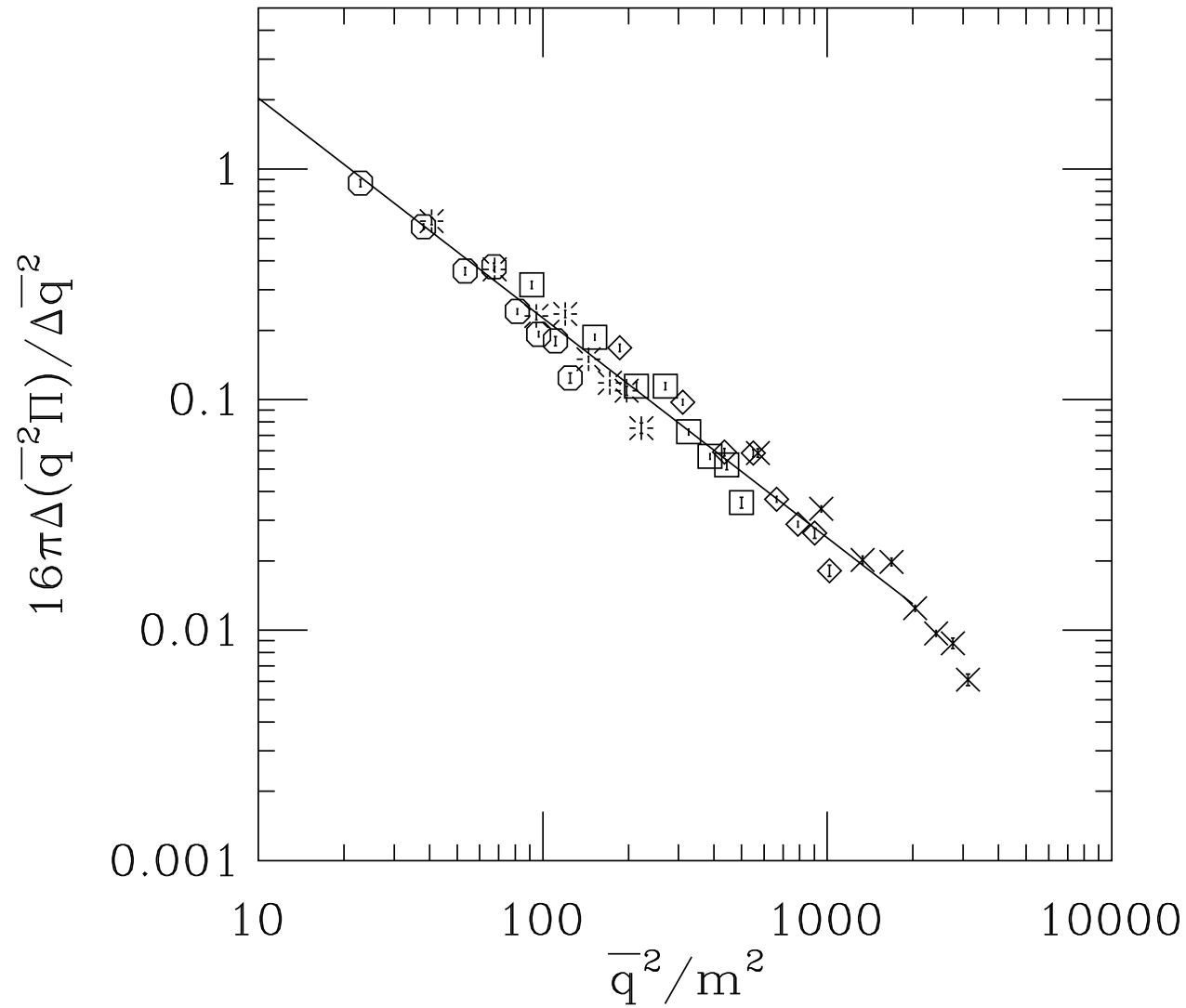


Figure 12: Numerically differentiated “S–parameter” for sextet QCD – log-log plot

Conclusions

- A fun calculation – and things like this will be needed for lattice BSM pheno
- Obviously, not the most wonderful system, but you use what you have
- Not much literature on precision EW for (near) conformal theories
 - $\Pi_T^{LR}(q^2) \rightarrow 0$ as $m_q \rightarrow 0$ – no surprise
 - Power law scaling – maybe no surprise if you believe in conformality
 - “Low resonance dominance” failed. I was surprised!
- Same techniques will work for any TC candidate – if anyone can find one!