

Two-Loop Self-Energy Corrections To The g -Factor Of Bound Electrons

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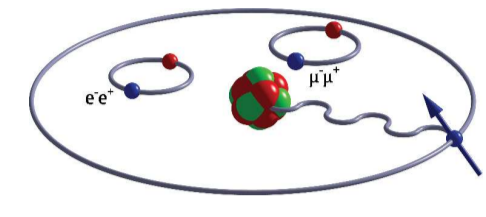
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Zeeman effect and g -factor

Magnetic moment μ of electron

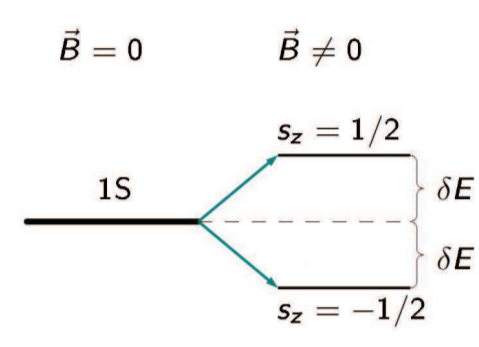
$$\mu = g \frac{e\mathbf{J}}{2m_e}$$



e : electron charge, m_e : electron mass, \mathbf{J} : electron angular momentum, g : electron g -factor [1]

Bound electron: energy shift δE of ground state $|1s\rangle$ with magnetic quantum number m_j in external magnetic field B :

$$\delta E = -\langle 1s | \mu \mathbf{B} | 1s \rangle = -m_j g \frac{e}{2m_e} B$$

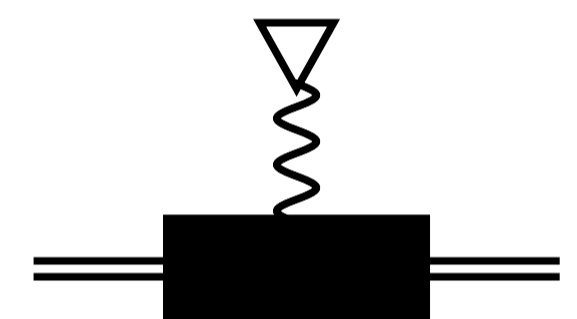


δE computed using the Two-time Green Function method [2].

Applications

- precision tests of QED in the presence of strong electromagnetic background fields [3, 4]
- determination of the electron mass [5, 6]
- determination of the muon mass [7]
- determination of the fine-structure constant [8, 9]
- test of physics beyond the standard model [10]

Theoretical description

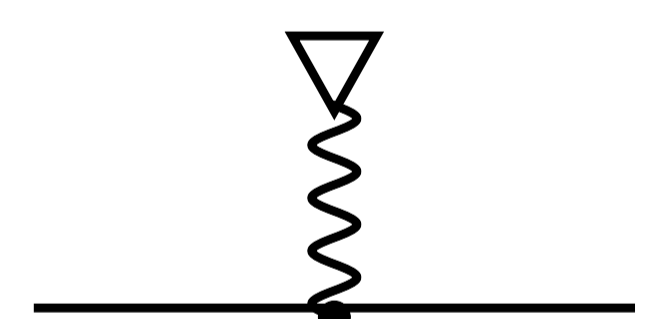


$$g = 2 \left(C^{(0)} + C^{(2)} \left(\frac{\alpha}{\pi} \right) + C^{(4)} \left(\frac{\alpha}{\pi} \right)^2 + \dots \right)$$

$C^{(2n)} = C_{\text{free}}^{(2n)} + \delta C_{\text{bound}}^{(2n)}(Z\alpha) \hat{=}$ sum of all Feynman diagrams with n closed loops [1]

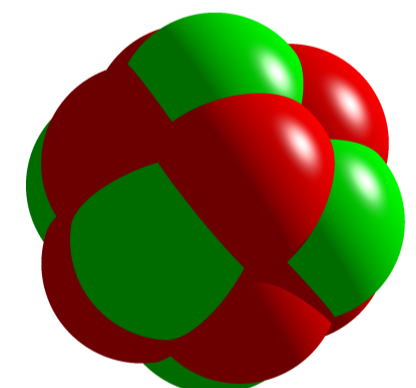
Non-perturbative determination of $C^{(2n)}$: "Furry picture"

Leading g -factor diagram



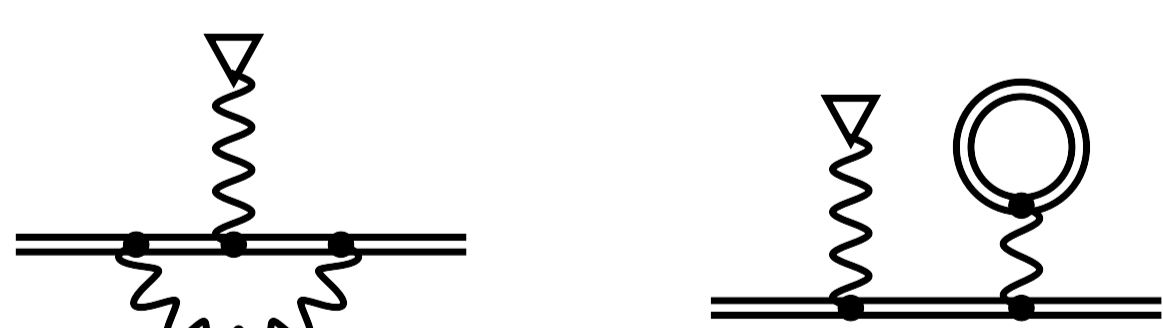
$$\Delta E_{\text{mag}} = -\frac{2}{3} i \int_0^\infty dr r^2 \text{Bier} f(r) g(r)$$

$$g_D = \frac{2}{3} + \frac{4}{3} \sqrt{1 - (Z\alpha)^2}$$



- Finite nuclear size
- Finite nuclear mass (recoil) [11, 12]
- Nuclear deformation [13]
- Nuclear polarization [14]

Feynman diagrams with one loop



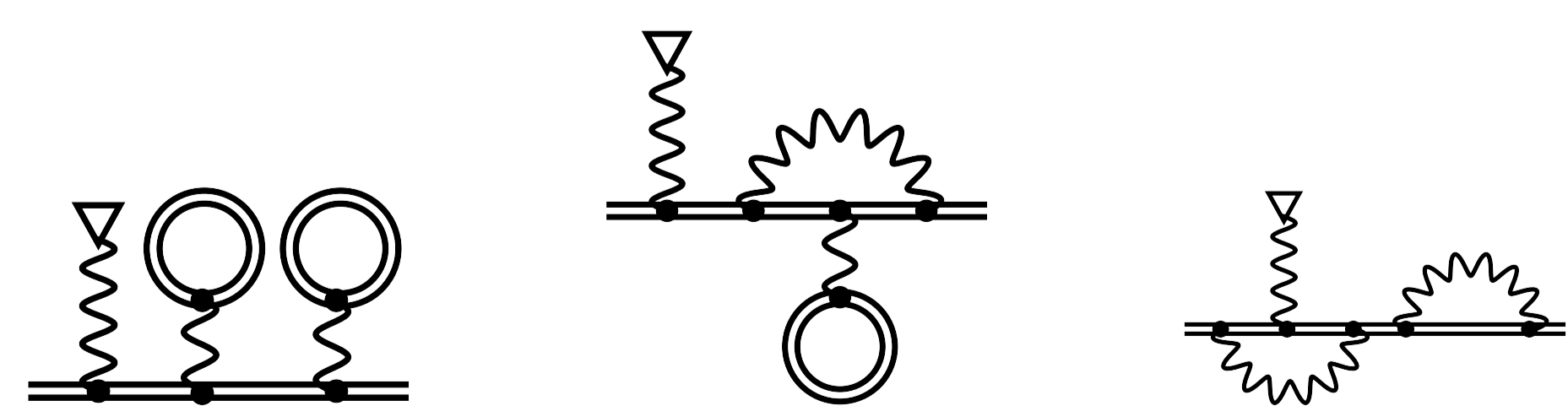
- **Self-energy (SE)**: interaction of the electron with the photon field

$$g_{\text{SE}} = \frac{\alpha}{\pi} \left(1 + \frac{(Z\alpha)^2}{6} \right) + \dots \quad [15, 16]$$

- **Vacuum polarization (VP)**: creation of a virtual charged particle-antiparticle pair

$$g_{\text{VP}} = -\frac{16}{15} \left(\frac{\alpha}{\pi} \right) (Z\alpha)^4 \left(\frac{m_e}{m_{\text{loop}}} \right)^2 + \dots \quad [17, 18]$$

Feynman diagrams with two loops



- Diagrams with two VP loops (**VPVP**) [19, 20]
- Diagrams with one SE and one VP loop (**SEVP**) [20]
- Diagrams with two SE loops (**SESE**) [21]

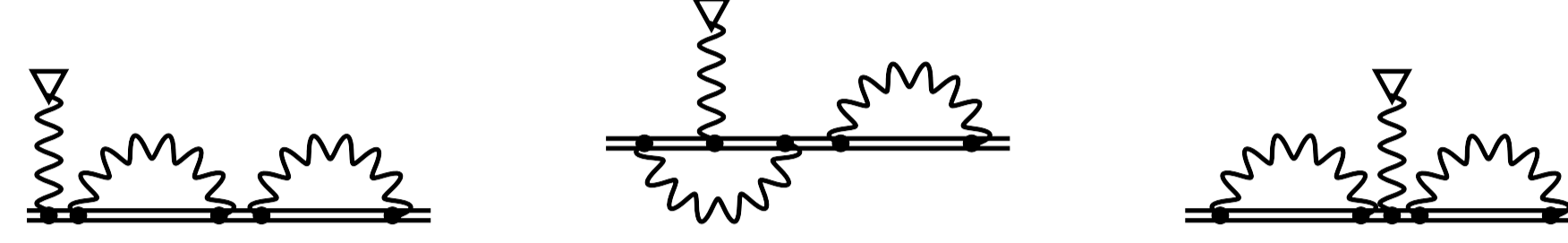
Magnetic loop VP diagrams: see poster by V. Debierre

Two-loop SESE corrections

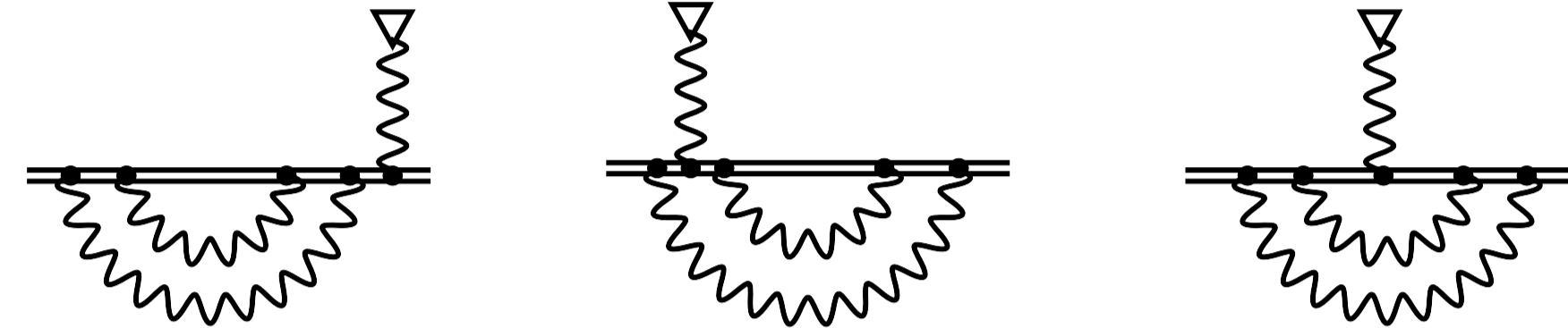
Theoretical uncertainty of g -factor dominated by **uncalculated higher-order SESE** correction

SESE diagrams

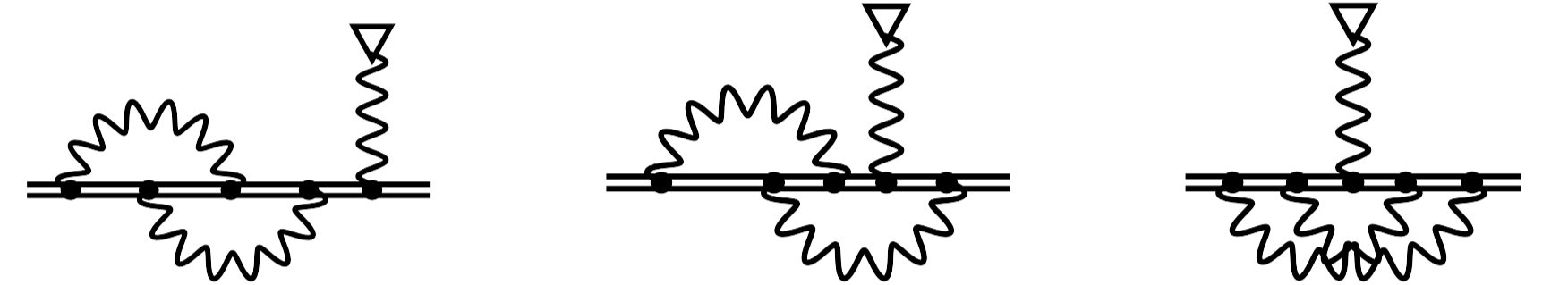
Loop after loop (LAL)



Nested loop (N)



Overlapping loop (O)



LAL, Irreducible and LAL, Reducible

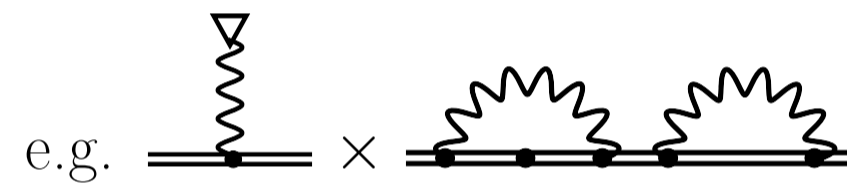
- Electron propagator between SE loop and magnetic interaction:

$$G_D = \sum_n \frac{|n\rangle\langle n|}{E_n - E_{1s}}$$

\Rightarrow Split into "irreducible" ($E_n \neq E_{1s}$) and "reducible" ($E_n = E_{1s}$)

- In LAL, up to two such electron propagators \Rightarrow Regroup LAL contributions into

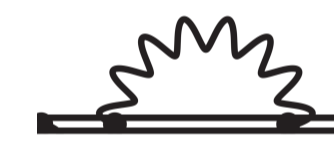
LAL, irred



LAL, red



Computation of LAL, irred: generalization of one-loop QED, with $|1s\rangle \rightarrow |\delta_{\Sigma} 1s\rangle$



$$|\delta_{\Sigma} 1s\rangle = \sum_{n, n \neq 1s} \frac{|n\rangle\langle n| \Sigma |1s\rangle}{E_{1s} - E_n}$$

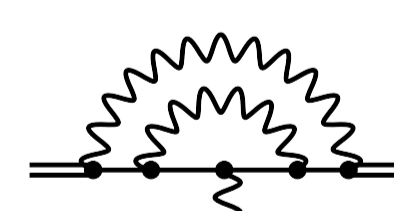
Renormalization

- **LAL, irred**: UV and IR finite
 - **LAL, red**: UV and IR divergent
 - **N**: UV and IR divergent
 - **O**: UV and IR divergent
- \Rightarrow divergences cancel

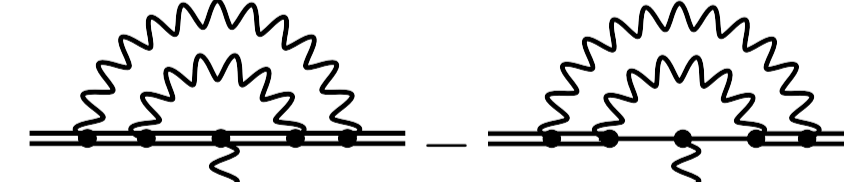
To deal with UV divergences, the N and O diagrams are divided into three terms [22].

- **F-term**: diagrams with only free internal electron lines (UV divergences)
- **M-term**: diagrams with bound internal electron lines (no UV divergences)
- **P-term**: diagrams with bound internal electron lines and a UV divergent subdiagram

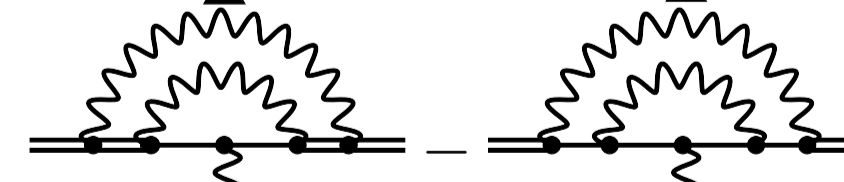
Example **F-term**:



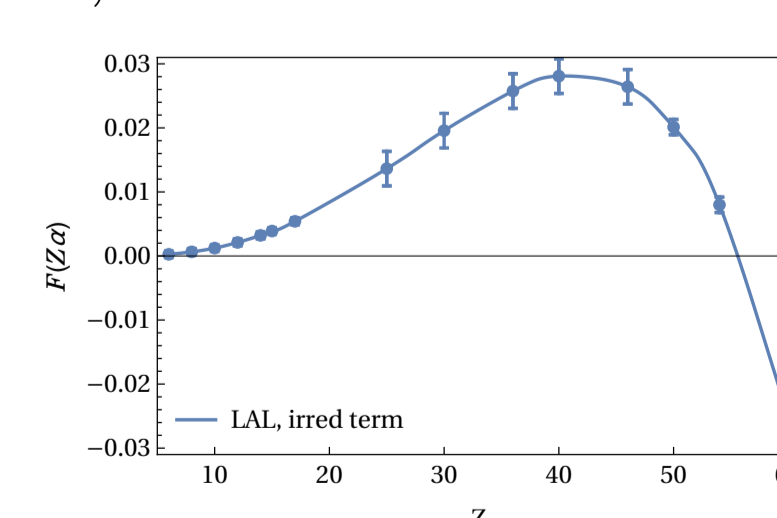
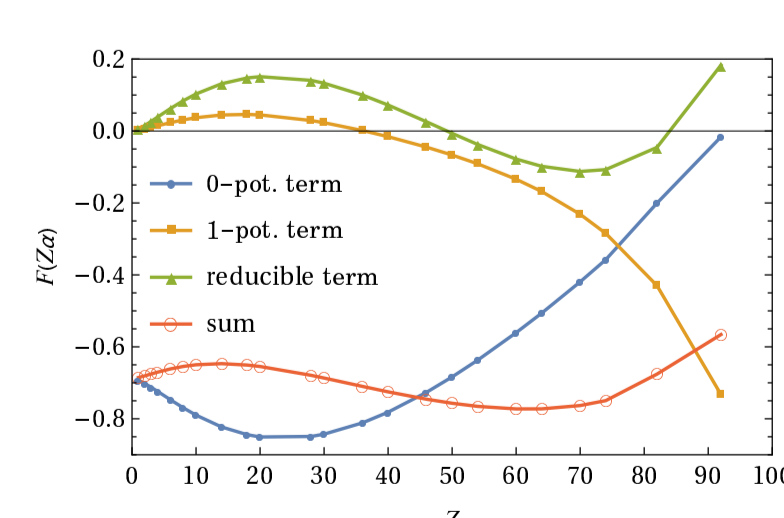
M-term:



P-term:



Numerical results for F-term & LAL, irred



$$\Delta g = \left(\frac{\alpha}{\pi} \right)^2 F(Z\alpha)$$

- **Zero-potential**: two-loop diagrams with free internal electron lines, no interaction with the nuclear potential
- **One-potential**: two-loop diagrams with free internal electron lines, one interaction with the nuclear potential
- **Reducible**: products of two one-loop functions ($\hat{=}$ part of LAL, red)

Our numerical results for the F-term converge to the free-electron value $F(0) = -0.68833\dots$ for low Z . [21, 15]

Two-loop SESE – outlook

IR divergences, methods

- Summation of diagrams whose IR divergences cancel (g -factor, one-loop SE [16])
- Subtraction terms to cancel IR divergences (Lamb shift, two-loop SESE [22])

Numerical challenges

- **M-term**: $\int d\omega_1 \int d\omega_2 \int dr_1 \dots \int dr_5 f(\dots)$
- $\sum_{\kappa_1, \kappa_2} g_{\kappa_1, \kappa_2}$ (infinite summations)
- **P-term**: Numerical Fourier transform
- $\sum_{\kappa_1} g_{\kappa_1}$ (infinite summation)

Consistency checks

- Comparison with perturbative determination of $C^{(4)}$ [15, 23]
- Comparison with SESE correction to Lamb shift [22]
- **Further ideas would be highly appreciated...**

Possibilities after complete two-loop calculation

- Improved theoretical accuracy of bound-electron g -factor for high $Z \Rightarrow$ Comparison with experimental g -factors (ALPHA-TRAP/ARTEMIS)
- Improvement of accuracy of α

Access to the muon mass

$$\left. \begin{array}{l} \text{Larmor frequency: } \omega_L = g \frac{e}{2m_\mu} B \\ \text{Cyclotron frequency: } \omega_{\text{cycl}} = \frac{Q}{m_{\text{ion}}} B \end{array} \right\} \Rightarrow m_\mu = \frac{g e \omega_{\text{cycl}}}{2Q \omega_L} m_{\text{ion}}$$

- $\frac{Q}{e}$ known exactly
 - m_{ion} known very precisely
 - ω_{cycl} from experiment
 - g -factor from theory
 - **Muonic ${}^4\text{He}^+$** : Small uncertainty of nuclear effects
 - Sum of all considered terms:
- $g = 2.002\,195\,193\,4(20)_{\text{calc}}(50)_{\text{uncalc}}$
- \rightarrow **possibility of improvement of muon mass accuracy by one order of magnitude** [7]
- \rightarrow **alternative access to the controversial free-muon g factor with the subtraction of binding effects from theory** [7]

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