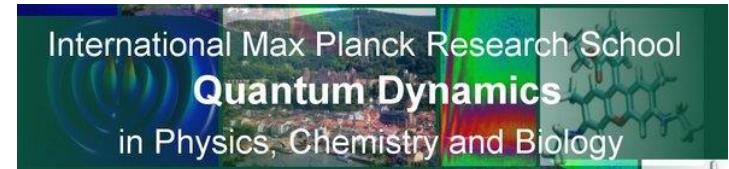




INTERNATIONAL
MAX PLANCK
RESEARCH SCHOOL

PT
FS

FOR PRECISION TESTS
OF FUNDAMENTAL
SYMMETRIES



High-Precision Measurement of the Deuteron's Atomic Mass

International Conference on Precision Physics and
Fundamental Physical Constants
Sascha Rau

L I N
—
T R A P



What is LIONTRAP?



- acronym for Light ION TRAP
- purpose-built high precision Penning-trap mass spectrometer for light ions, located in Mainz, follow up from Mainz g -factor experiment for highly charged ions (electron mass, strong field BS-QED Tests)
- first successful measurement campaign on proton mass with $\delta m/m = 32$ ppt (parts per trillion, 10^{-12})
(F. Heiße et al. PRL 119, 033001, 2017)

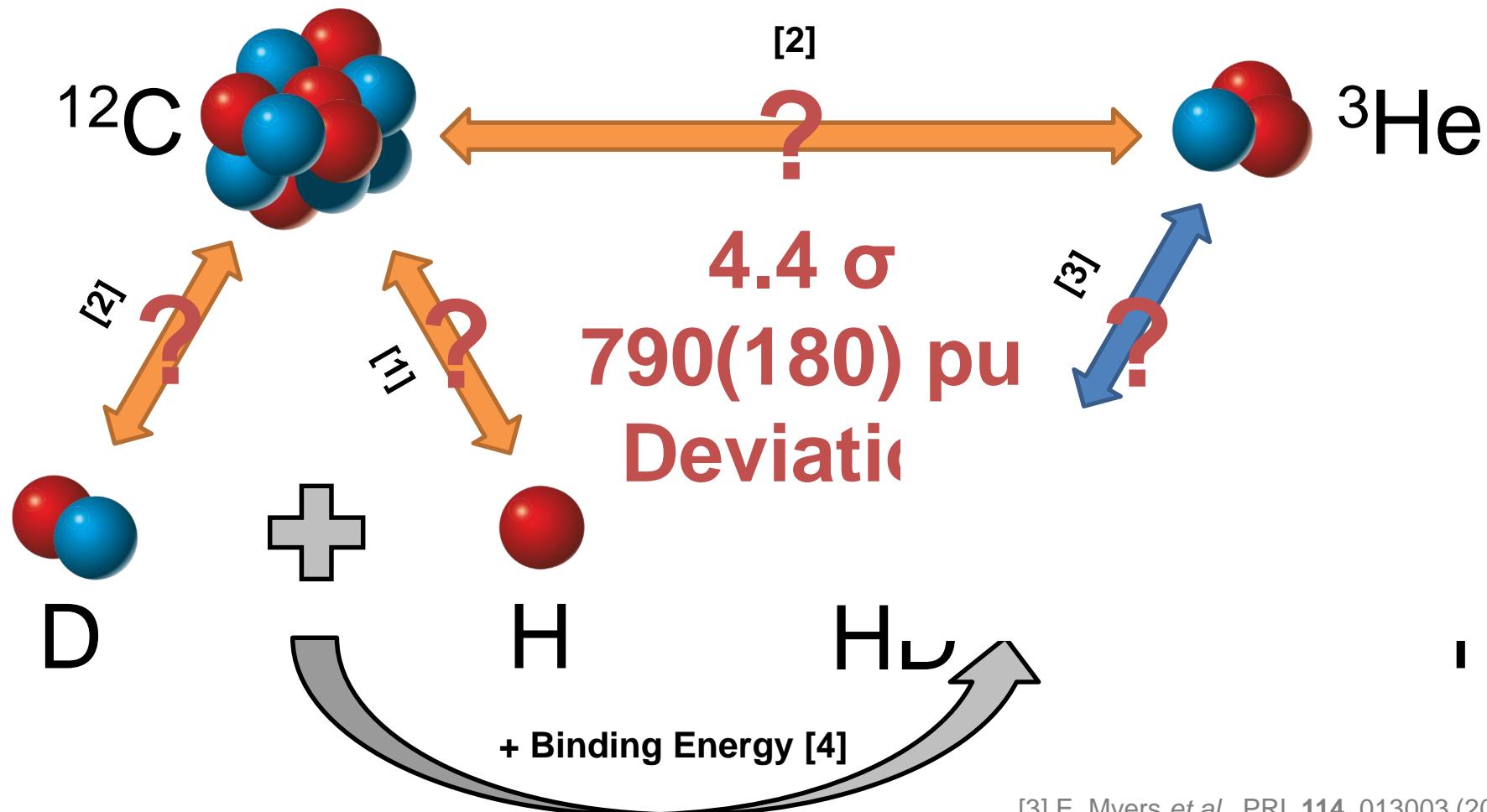
Why are mass measurements on light ions important?

- relevant for precision spectroscopy (m_p/m_e), tests of special relativity (m_d), KATRIN ($m_{^3\text{He}} - m_t$)
- inconsistency in literature calls for independent values
- difficult because of relatively small signals, high relativistic shifts

Puzzle of Light Atomic Masses

R.S. Van Dyck *et al.* @ UW

E. Myers *et al.* @ FSU



[1] R. S. Van Dyck Jr. *et al.*, AIP Conf. Proc. **457**, 101 (1999)

[2] S. L. Zafonte *et al.*, Metrologia **52**, 280 (2015)

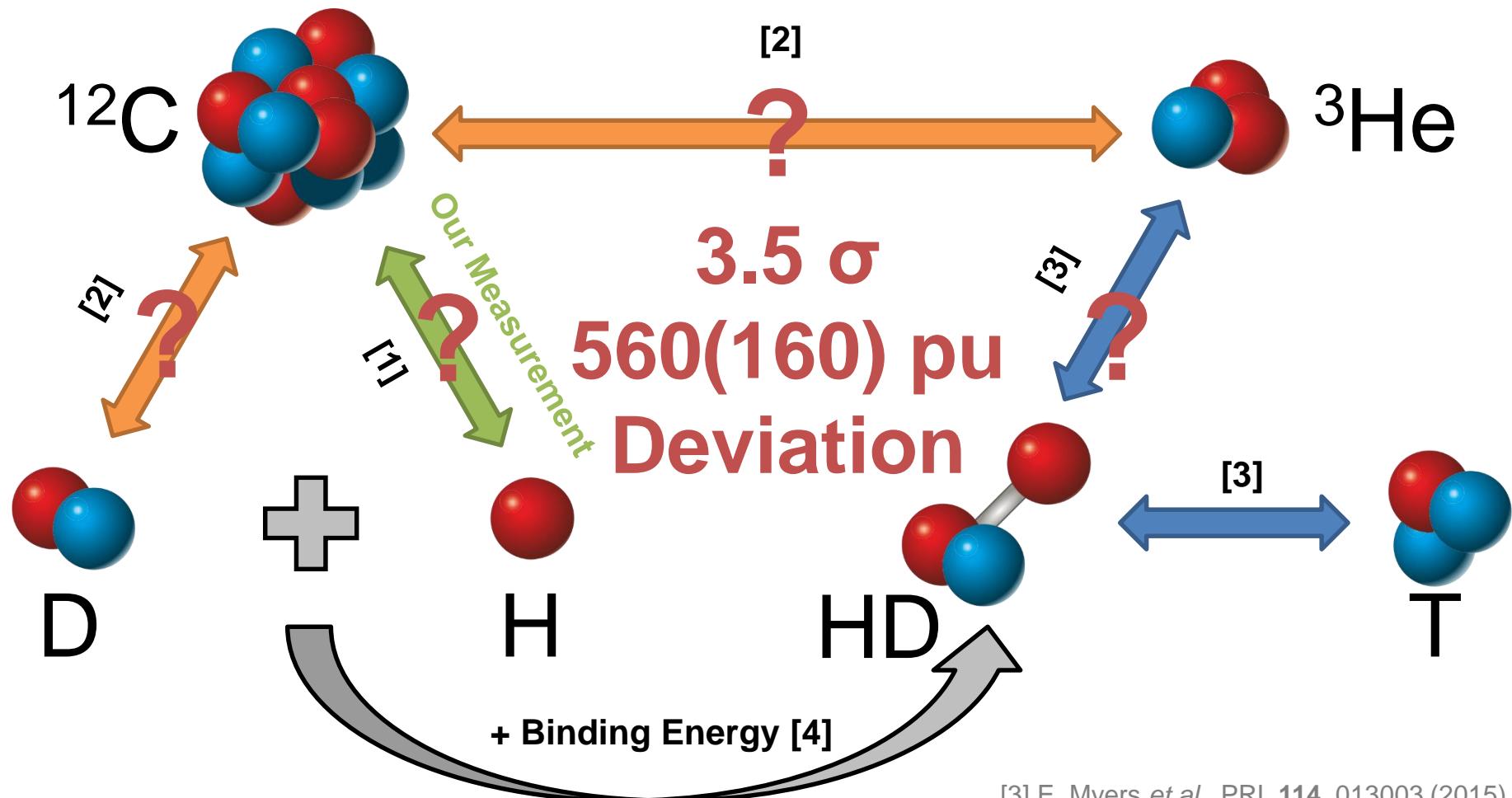
[3] E. Myers *et al.*, PRL **114**, 013003 (2015)

[4] Yan *et al.*, PRA **67**, 062504 (2003)

Puzzle of Light Atomic Masses

R.S. Van Dyck et al. @ UW

E. Myers et al. @ FSU



[1] F. Heiße et al., PRL 119, 033001 (2017)

[2] S. L. Zafonte et al., Metrologia 52, 280 (2015)

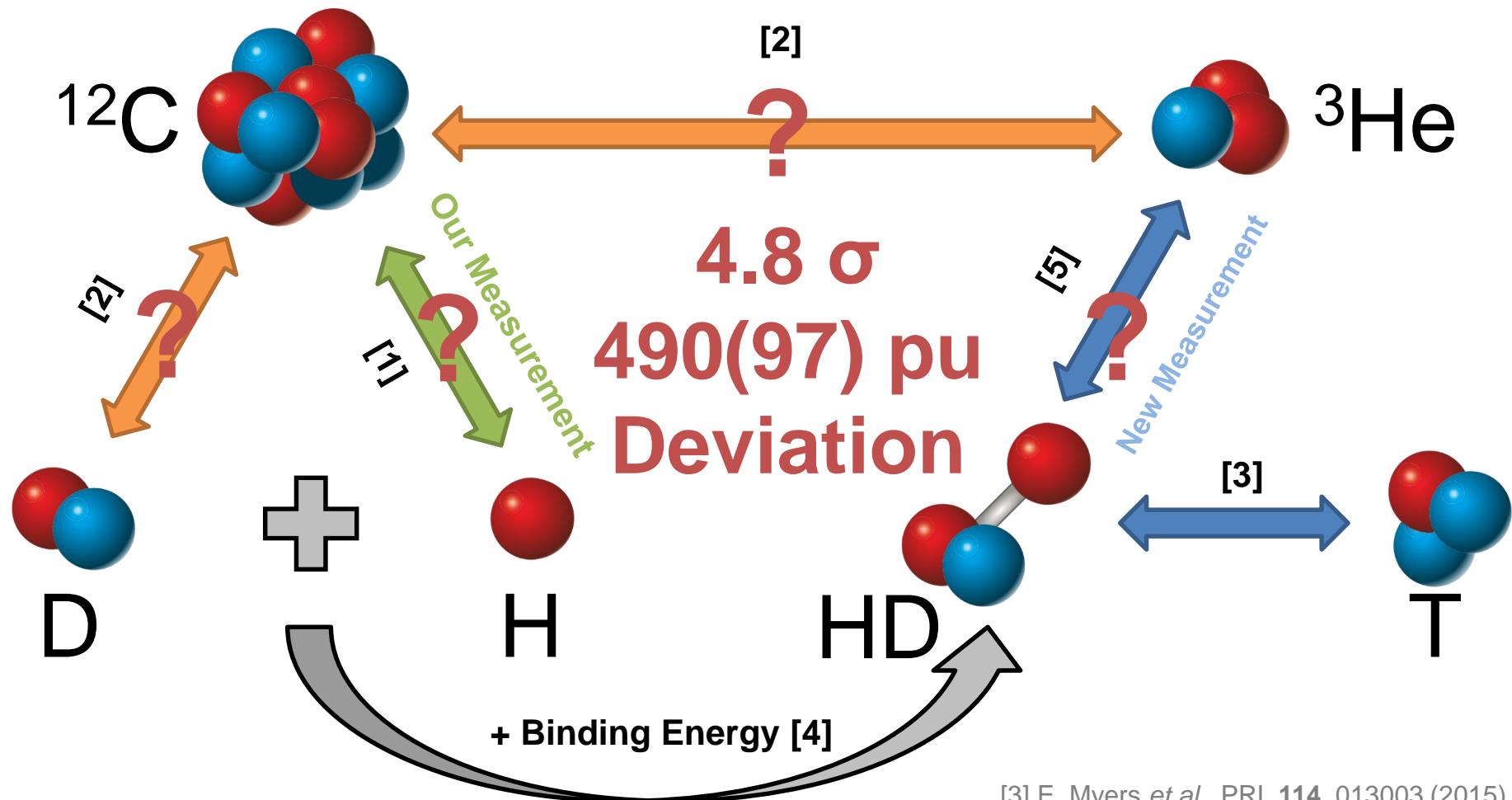
[3] E. Myers et al., PRL 114, 013003 (2015)

[4] Yan et al., PRA 67, 062504 (2003)

Puzzle of Light Atomic Masses

R.S. Van Dyck et al. @ UW

E. Myers et al. @ FSU



[1] F. Heiße et al., PRL 119, 033001 (2017)

[2] S. L. Zafonte et al., Metrologia 52, 280 (2015)

[3] E. Myers et al., PRL 114, 013003 (2015)

[4] Korobov et al., PRL 118, 233001 (2017)

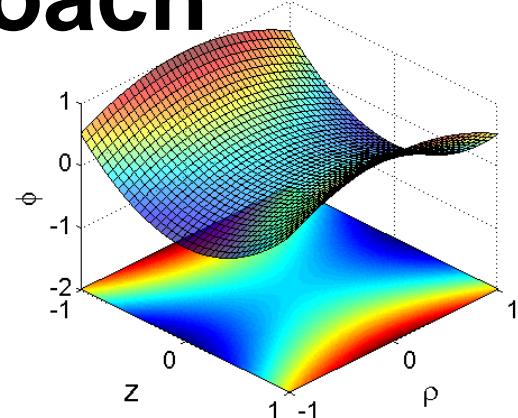
[5] S. Hamzelou et al., PRA 96, 060501 (2017)

Outline

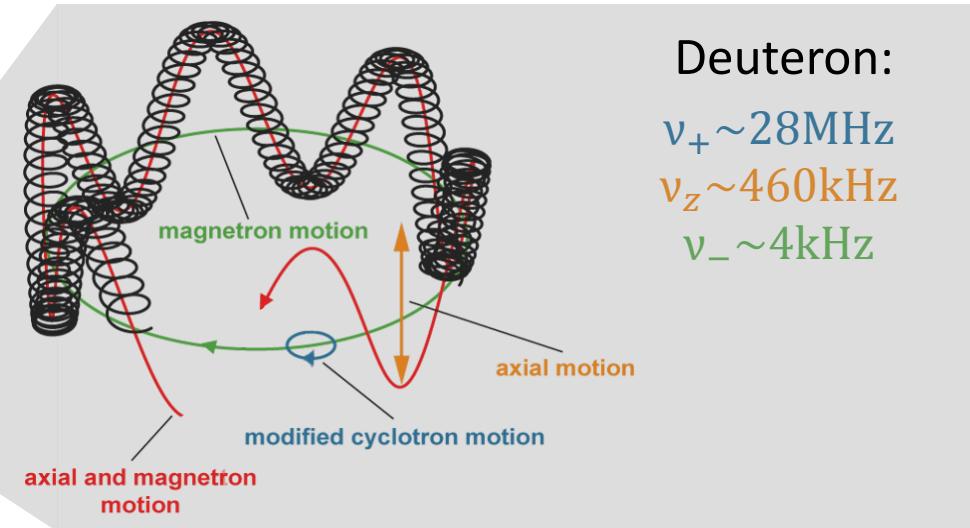
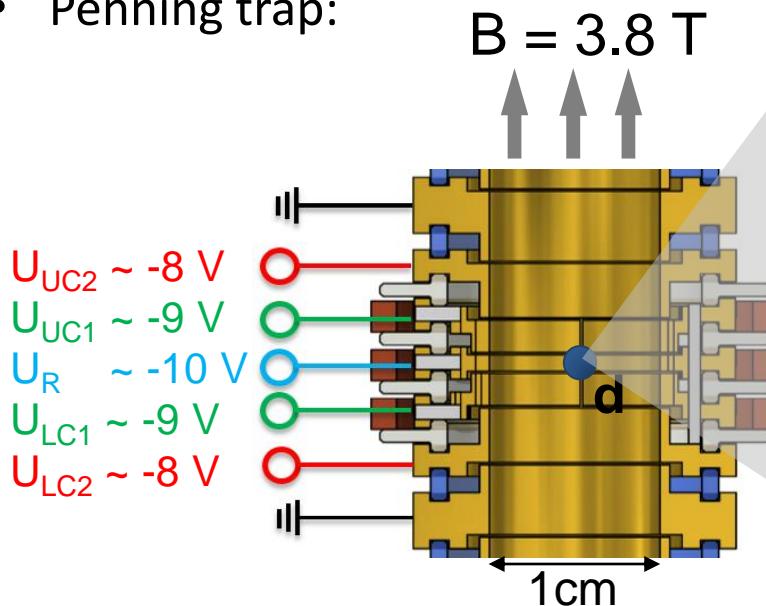
- **Setup and Experimental Techniques**
 - a new ultra-harmonic trap
- Proton's Atomic Mass
F. Heiße *et al.*, PRL **119**, 033001 (2017)
- Deuteron's Atomic Mass

Experimental Approach

- measurement of cyclotron frequency:
 → homogenous static magnetic field $v_c = \frac{1}{2\pi m} \frac{q}{B}$
 → electrostatic quadrupole potential for trapping



- Penning trap:



- cyclotron frequency via invariance theorem [1]:

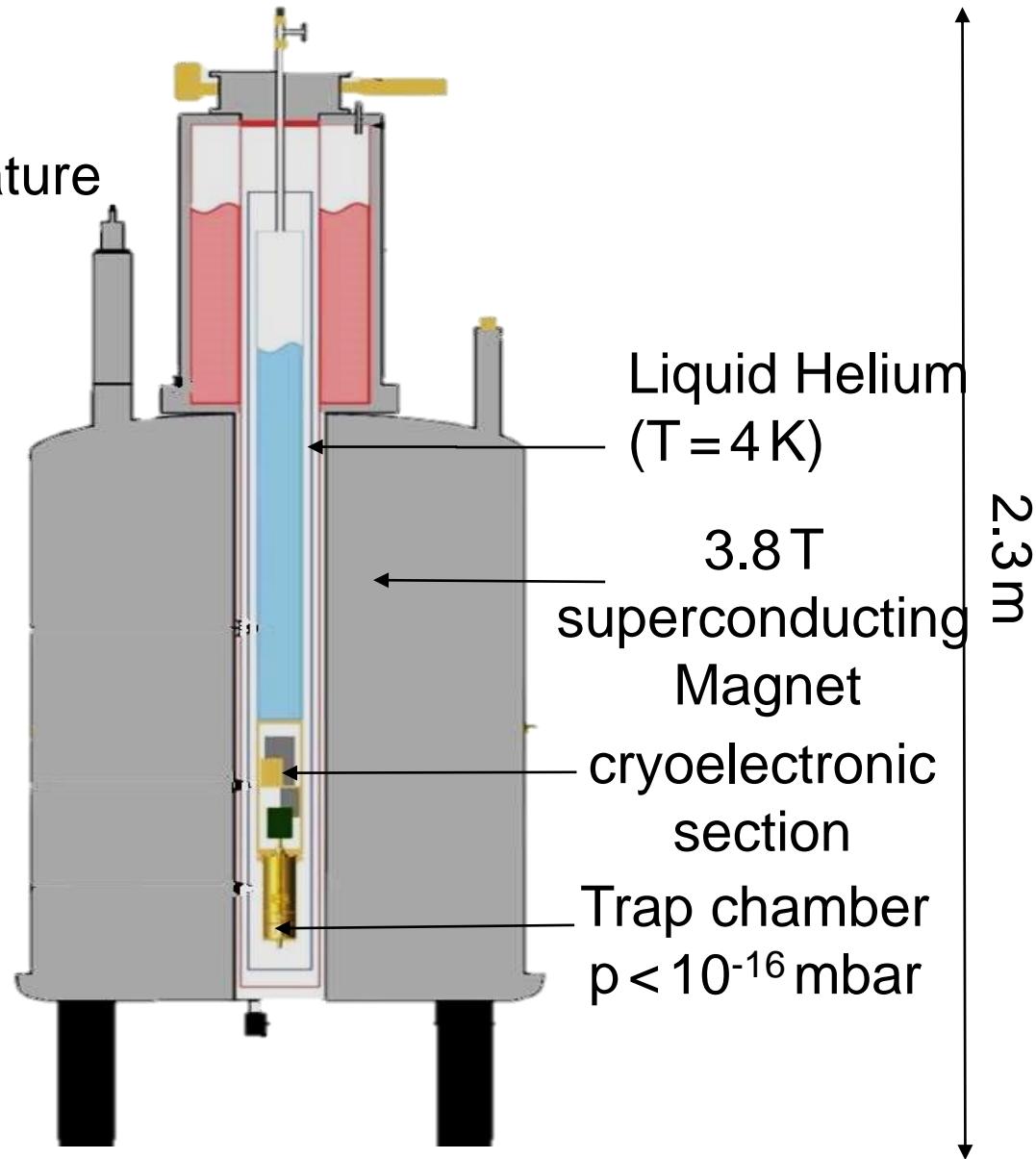
$$v_c = \sqrt{v_+^2 + v_z^2 + v_-^2} = \frac{1}{2\pi m} \frac{q}{B}$$

$$B = \frac{2\pi m}{q(^{12}C)} v_c (^{12}C)$$

[1] L.S. Brown & G. Gabrielse, PRA 25, 2423 (1982)

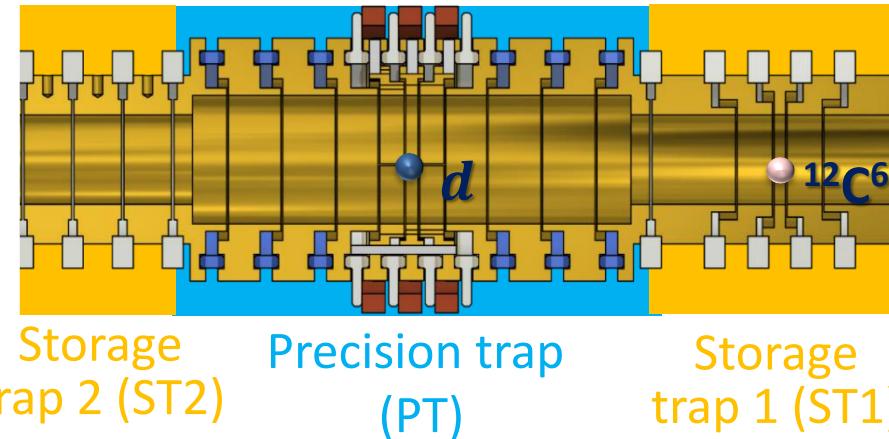
Setup

- offline Experiment
- liquid Helium temperature
- hermetically sealed trap chamber
- in situ ion creation using miniature EBIT/S



Measurement Procedure

$$m_d = \frac{1}{6} \frac{v_c(^{12}\text{C}^{6+})}{v_c(d)} m(^{12}\text{C}^{6+})$$



both ions at the same time within the trap chamber

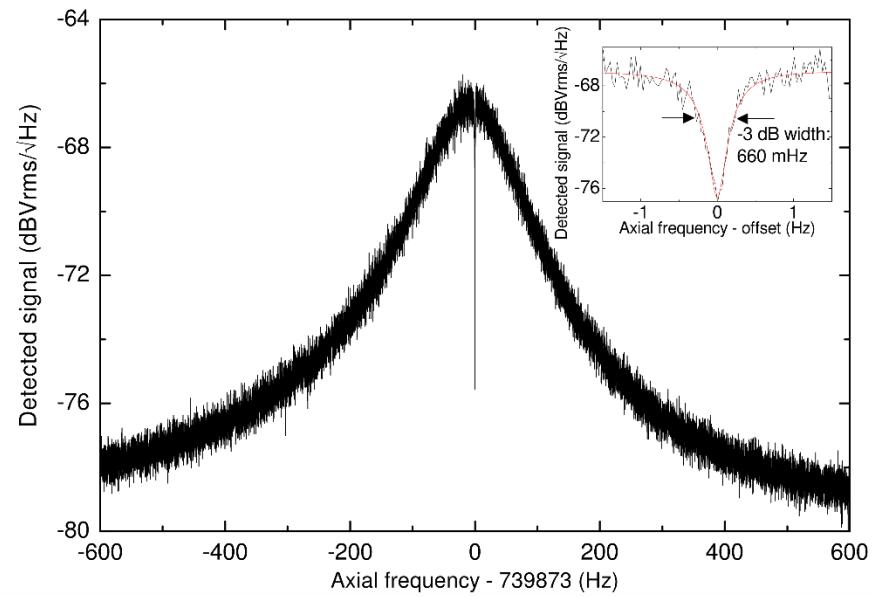
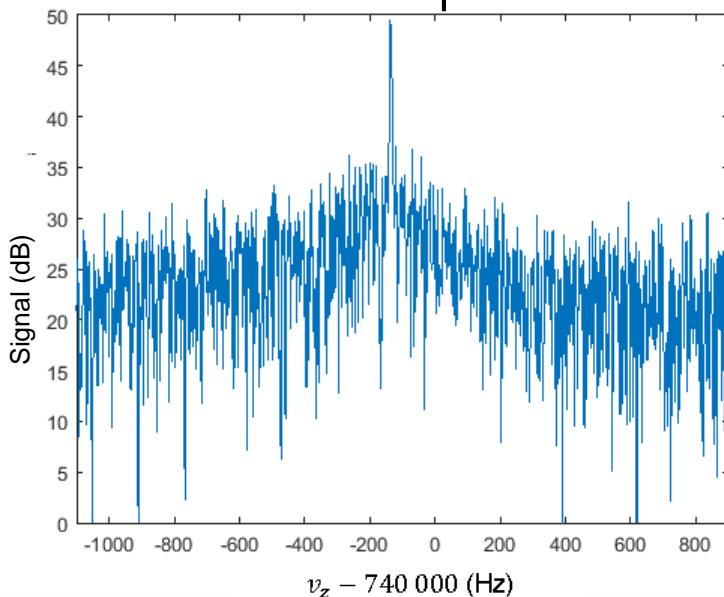
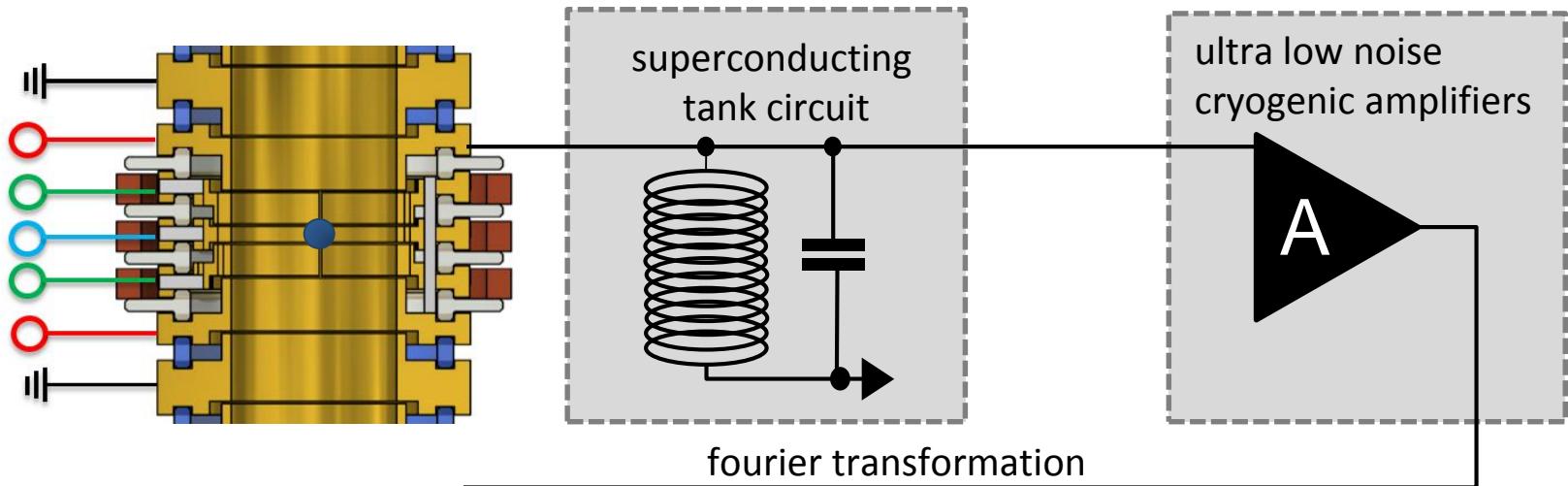
$$\left. \begin{array}{l} \text{(I)} \quad v_c(d) \text{ in PT} \\ \text{(II)} \quad v_c(^{12}\text{C}^{6+}) \text{ in PT} \end{array} \right\} R \equiv \frac{v_c(^{12}\text{C}^{6+})}{v_c(d)}$$

time between v_+ measurements optimized

→ reduction of impact of magnetic field fluctuations compared to former measurements

Eigenfrequency Detection

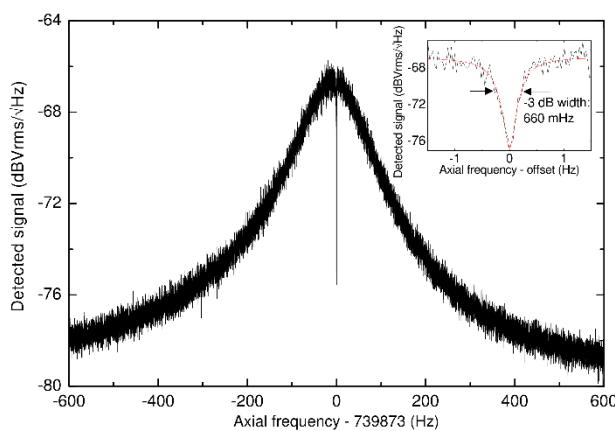
- measurement of induced image currents (\sim fA) on trap electrodes



Eigenfrequency Measurement Techniques

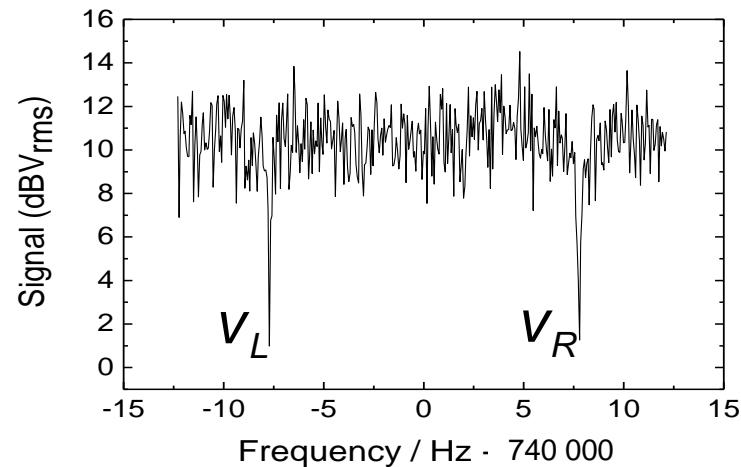
1. Dip Signal

v_z of thermalized ion ($T_z=4.5$ K)
measured as narrow “dip”



2. Double-Dip Signal:

v_+ , v_- coupled to axial motion
via rf-sideband coupling



$$v_+ = v_{rf} - v_z + v_L + v_R \quad v_- = v_{rf} + v_z - v_L - v_R$$

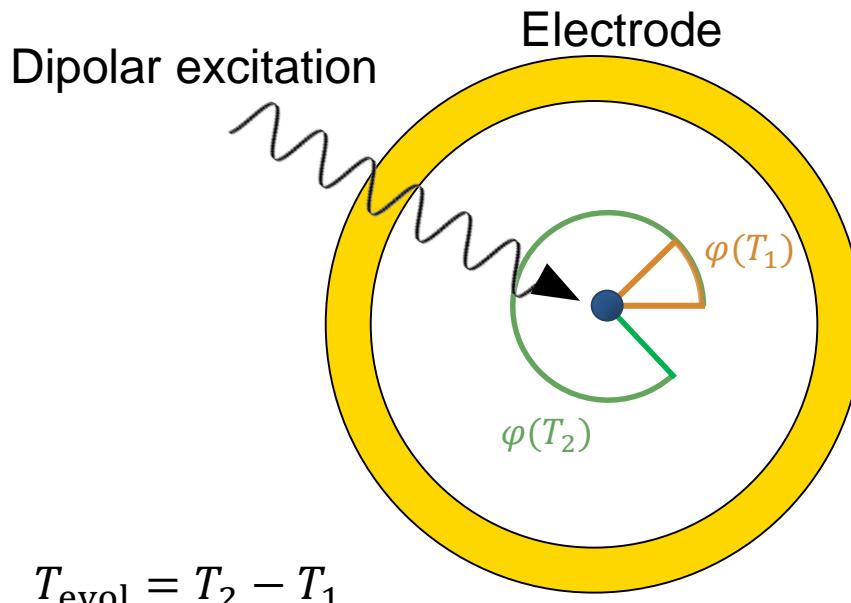
Advantages: - Thermal equilibrium during measurement
→ Small energy dependent systematic shifts

Disadvantages: - $\frac{\delta u_n}{u_n} \propto \frac{1}{\sqrt{t_m}}$ → “slow” (3 min) → impact of field fluctuations
- Lineshape dependent



Phase-Sensitive Detection Method

A phase sensitive measurement scheme allows **coherent detection**.
The complex Fourier spectrum provides besides the peak position in the Fourier spectrum also the instantaneous phase of the ion.

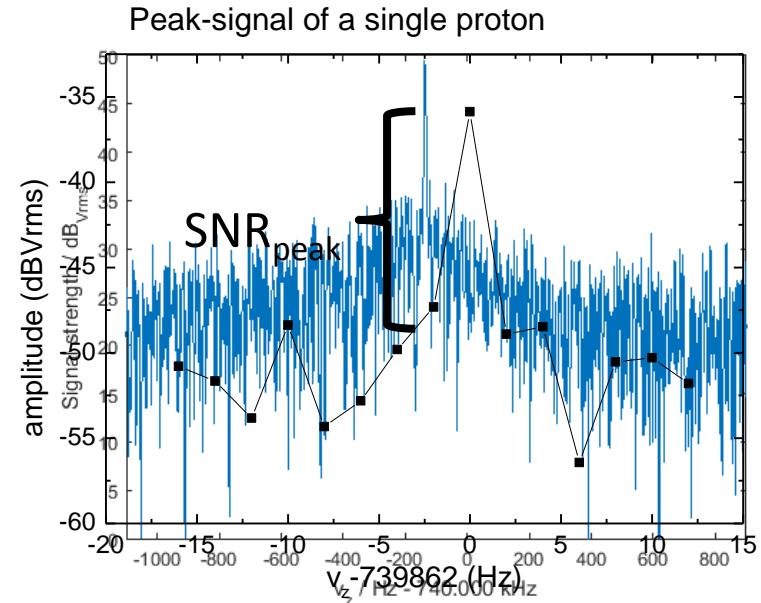


$$T_{\text{evol}} = T_2 - T_1$$

Advantages:

- linear scaling with Measurement Time
- no lineshape model needed

$$\nu = \frac{\varphi(T_2) - \varphi(T_1)}{360^\circ \cdot T_{\text{evol}}}$$



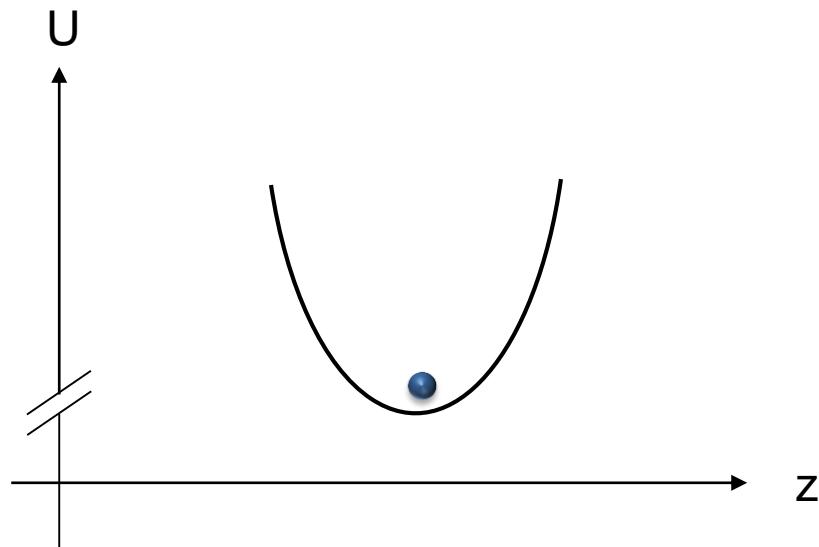
$$\frac{\delta\nu}{\nu} = \frac{\delta\varphi}{360^\circ \cdot \nu \cdot T_{\text{evol}}}, \quad \delta\varphi \triangleq \text{phase jitter}$$

A New Ultra-Harmonic Trap

Ideal world:

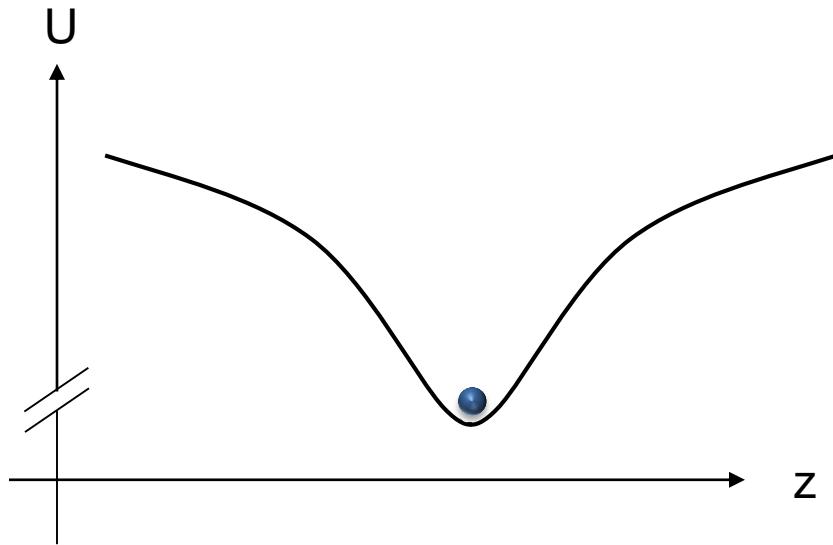
$$U(z) = \frac{1}{2} U_r \left(C_0 + C_2 \frac{z^2}{d_{char}^2} \right)$$

→ v_z independent of energy!



Real world:

→ v_z depends on energy!



A New Ultra-Harmonic Trap

Real world:

$$U(z) = \frac{1}{2} U_r \left(C_0 + C_2 \frac{z^2}{d_{char}^2} + C_4 \cancel{\frac{z^4}{d_{char}^4}} + C_6 \cancel{\frac{z^6}{d_{char}^6}} + C_8 \frac{z^8}{d_{char}^8} + C_{10} \frac{z^{10}}{d_{char}^{10}} + \dots \right) = \frac{1}{2} U_r \sum_{i=0,2,\dots}^{\infty} C_i \frac{z^i}{d_{char}^i}$$

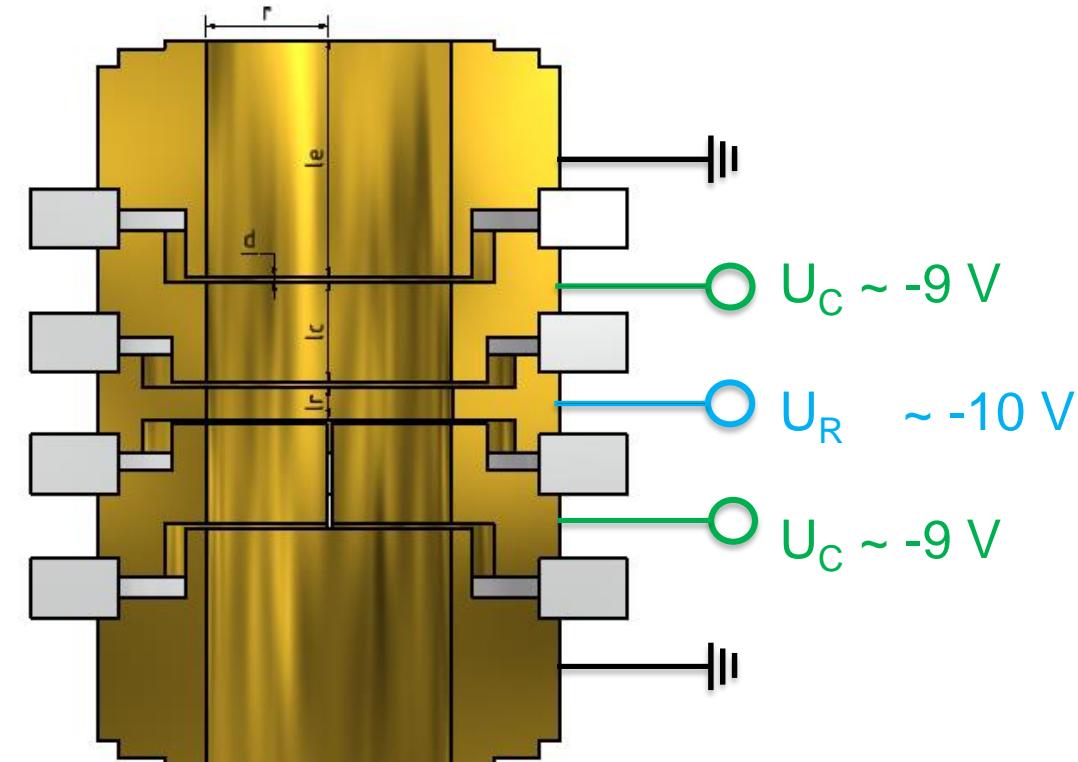
So far:

Use one pair of correction electrodes

Degrees of freedom: L_R , L_C , U_C

Optimized parameters:

$C_4 = C_6 = 0$ (compensated)



A New Ultra-Harmonic Trap

Real world:

$$U(z) = \frac{1}{2} U_r \left(C_0 + C_2 \frac{z^2}{d_{char}^2} + C_4 \cancel{\frac{z^4}{d_{char}^4}} + C_6 \cancel{\frac{z^6}{d_{char}^6}} + C_8 \cancel{\frac{z^8}{d_{char}^8}} + C_{10} \cancel{\frac{z^{10}}{d_{char}^{10}}} + \dots \right) = \frac{1}{2} U_r \sum_{i=0,2,\dots}^{\infty} C_i \frac{z^i}{d_{char}^i}$$

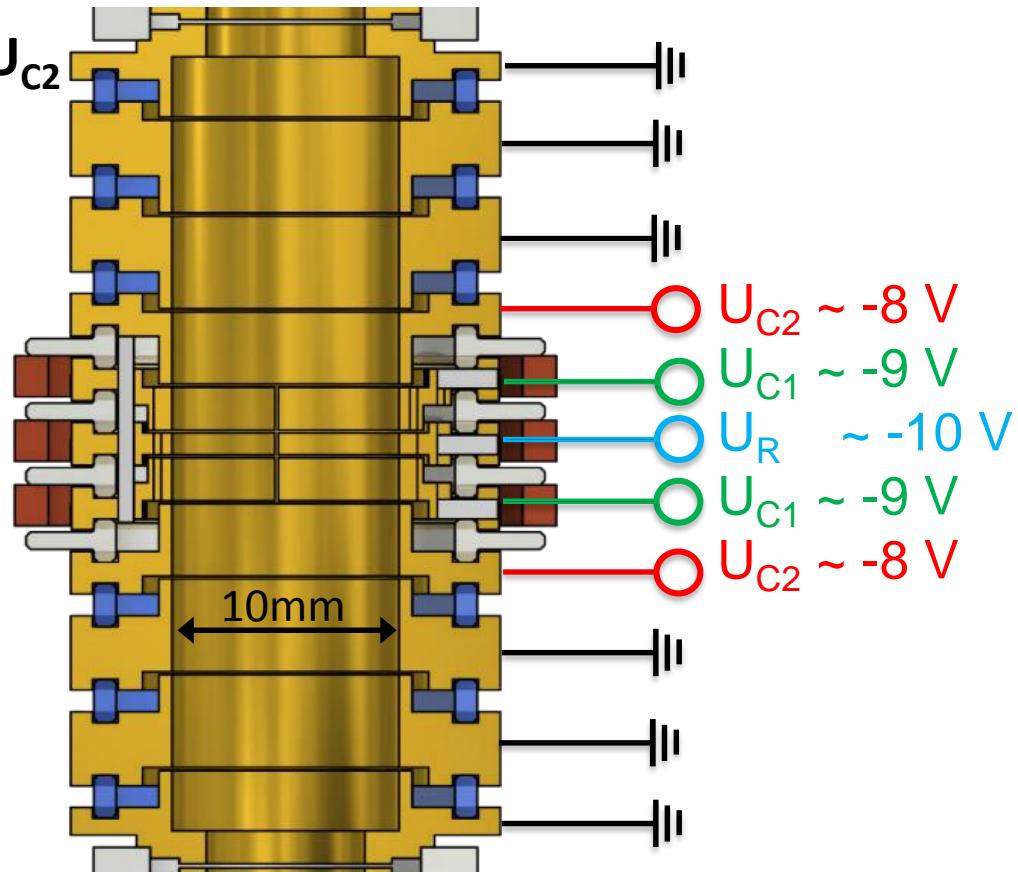
New trap (additional pair of correction electrodes)

Degrees of freedom: L_R , L_{C1} , U_{C1} , L_{C2} , U_{C2}

Optimized parameters:

$$C_4 = C_6 = C_8 = C_{10} = 0$$

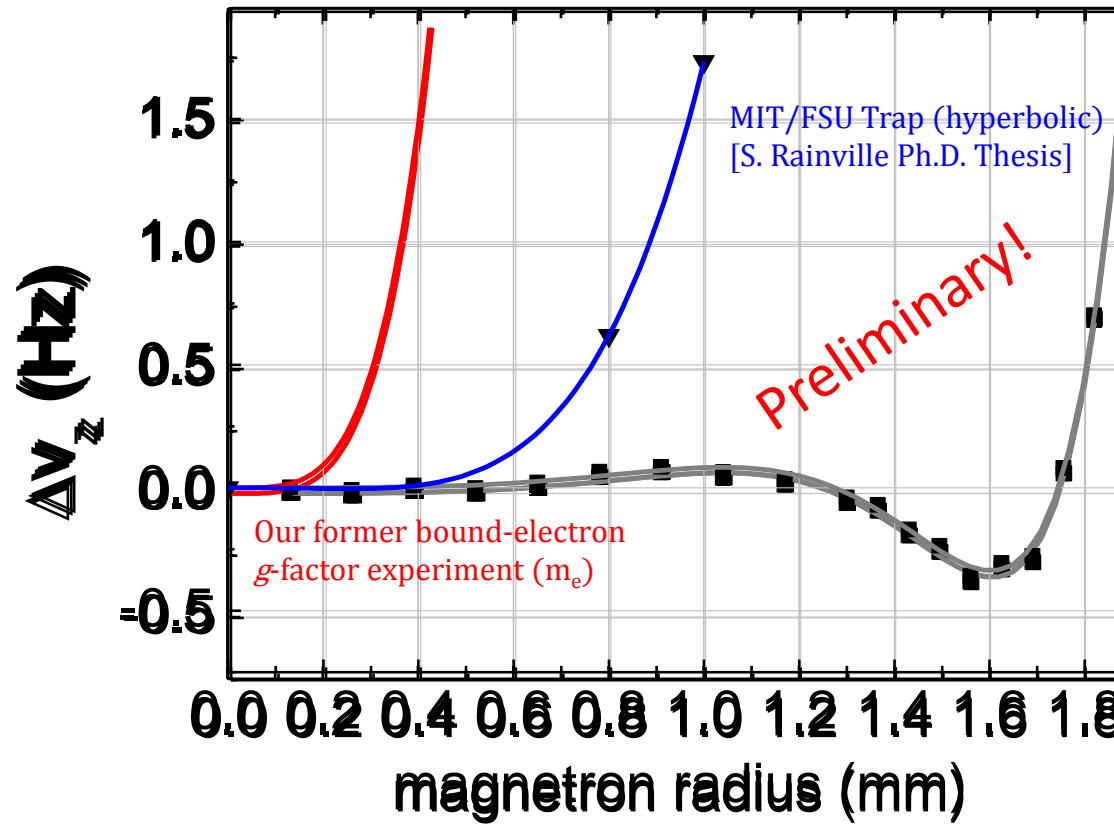
(doubly compensated)



A New Ultra-Harmonic Trap - Performance

Study shifts of axial freq. due to magnetron bursts:

$$\frac{\Delta\nu_z}{\nu_z} = -\frac{3}{2} \frac{C_4}{d_{char}^2 C_2} r_-^2 + \frac{15}{8} \frac{C_6}{d_{char}^4 C_2} r_-^4 - \frac{35}{8} \frac{C_8}{d_{char}^6 C_2} r_-^6 + \dots$$

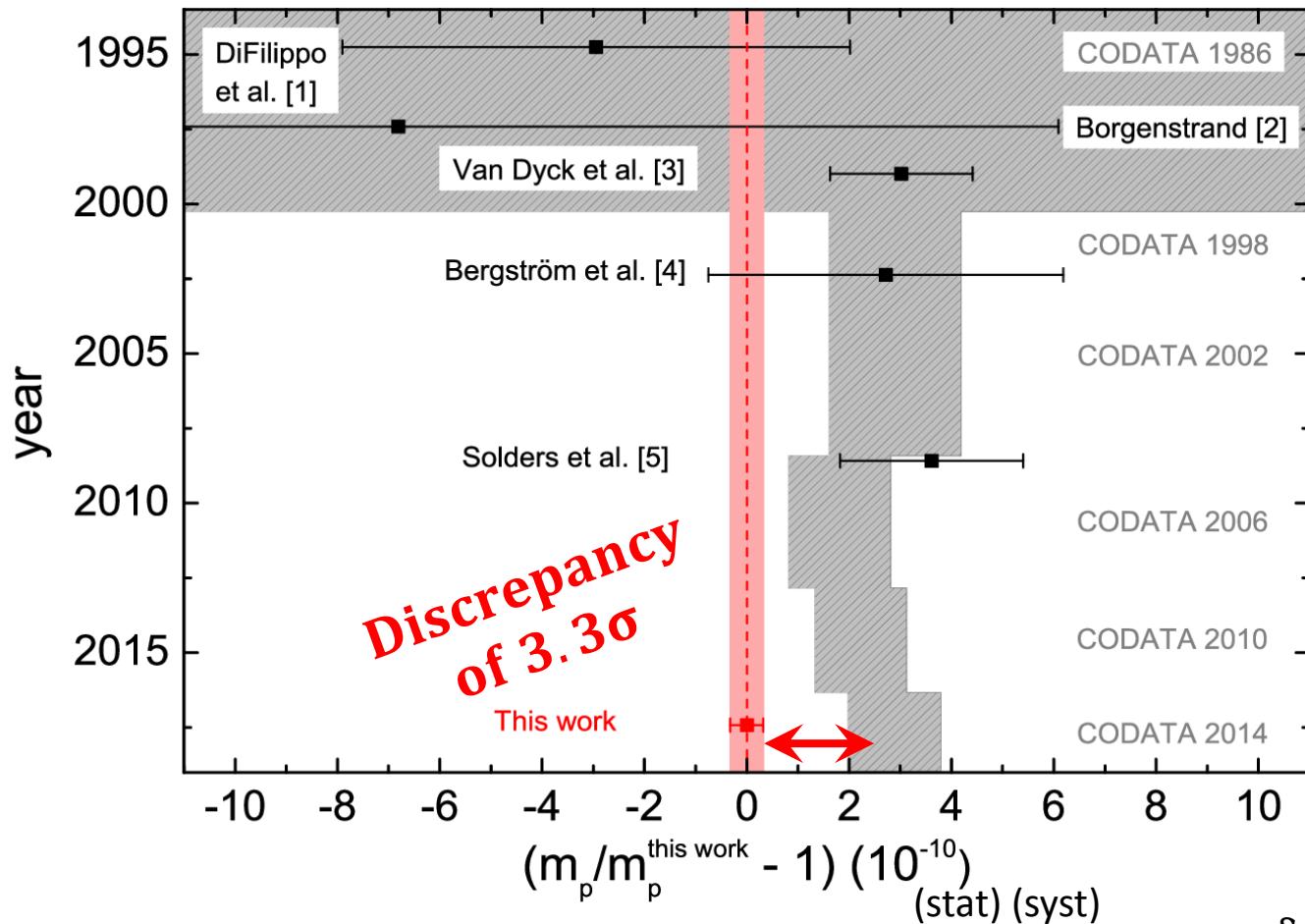


F. Heiße *et al.*, PRA, accepted

Outline

- Setup and Experimental Techniques
 - a new ultra-harmonic trap
- **Proton's Atomic Mass**
F. Heiße *et al.*, PRL 119, 033001 (2017)
- Deuteron's Atomic Mass

Results



$$m_p = 1.007\ 276\ 466\ 583\ (15)(29) \text{ u}$$

$$\frac{\delta m_p}{m_p} = 3.2 \cdot 10^{-11}$$

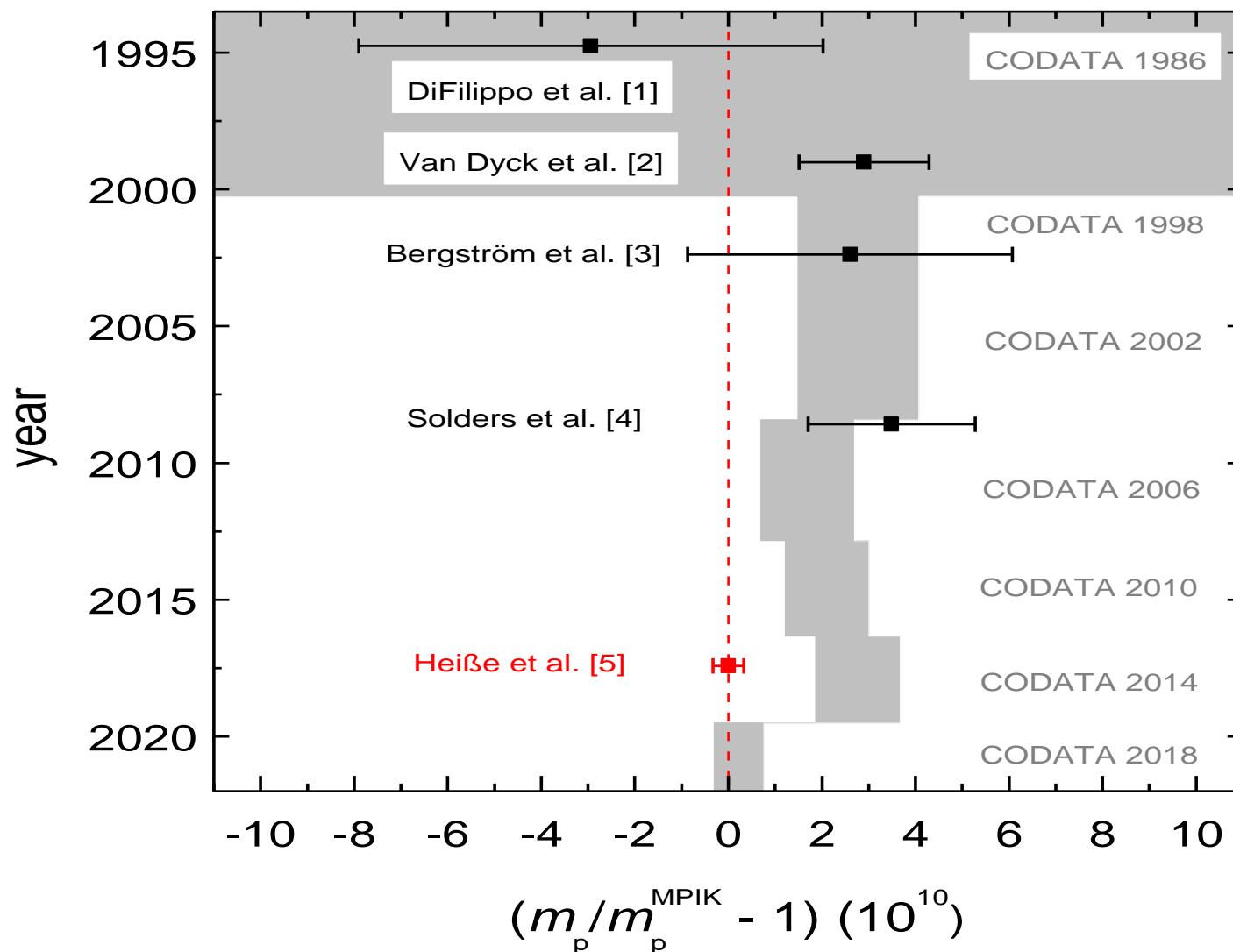
→ Improvement by a factor of 3

[1] PRL **73**, 1481 (1994). [2] Ph.D. thesis, Stockholm University (1997).

[3] AIP Conf. Proc. **457**, 101 (1999). [4] Phys. Scr. **66**, 201 (2002). [5] PRA **78**, 2514 (2008).

[PRL **119**, 033001 (2017)]

Results

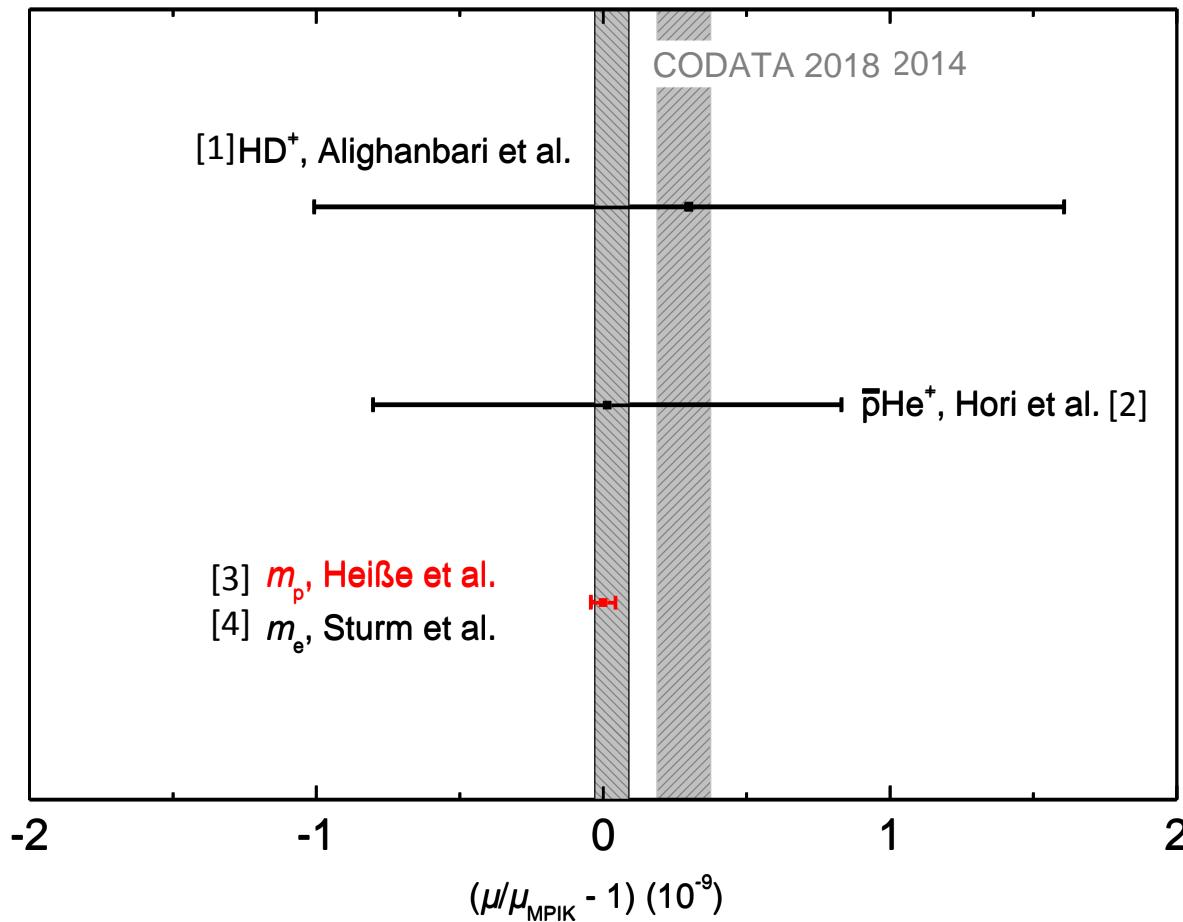


[1] PRL **73**, 1481 (1994). [2] AIP Conf. Proc. **457**, 101 (1999).

[3] Phys. Scr. **66**, 201 (2002). [4] PRA **78**, 2514 (2008).

[5] PRL **119**, 033001 (2017)

Proton/Electron Mass Ratio



[1] S. Alighanbari *et al.* Nat. Phys. **14**, 555 (2018)

[2] M. Hori *et al.* Science **354**, 610 (2016)

[3] F. Heiße *et al.*, Phys. Rev. Lett. **119**, 033001 (2017)

[4] S. Sturm *et al.* Nature **506**, 467 (2014)

Systematic Uncertainties

Effect	Relative Uncertainty (10^{-12})
Residual Magnetostatic Inhomogeneity	27.5
Special Relativity	7.1
Image Charge	4.6
Quadratic Sum	≈ 29

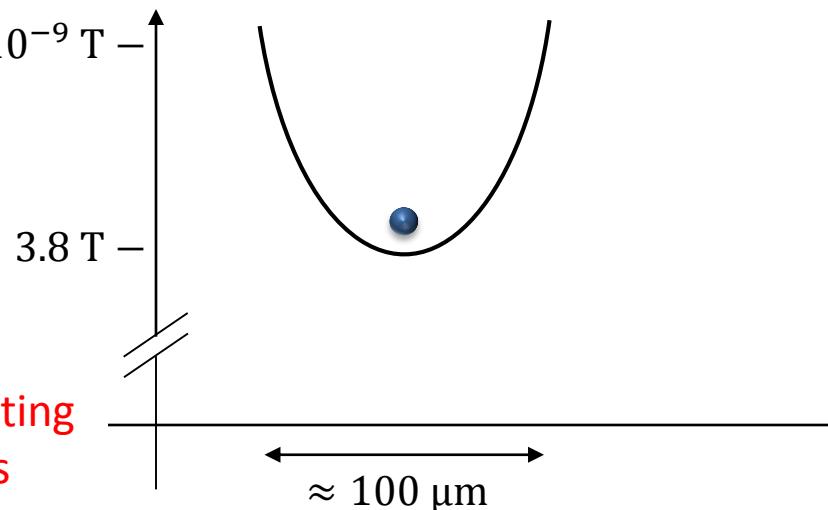
Magnetic inhomogeneity:

$$B(z) = B_0 + B_1 z + B_2 z^2 + \dots$$

$$\langle B \rangle = B_0 + B_2 \langle z_0^2 \rangle / 2$$

$$B_2 = -0.27(2) \text{ T/m}^2$$
$$B_1 = 0.92(1) \text{ mT/m}$$

\rightarrow Superconducting
B1/B2-shim coils

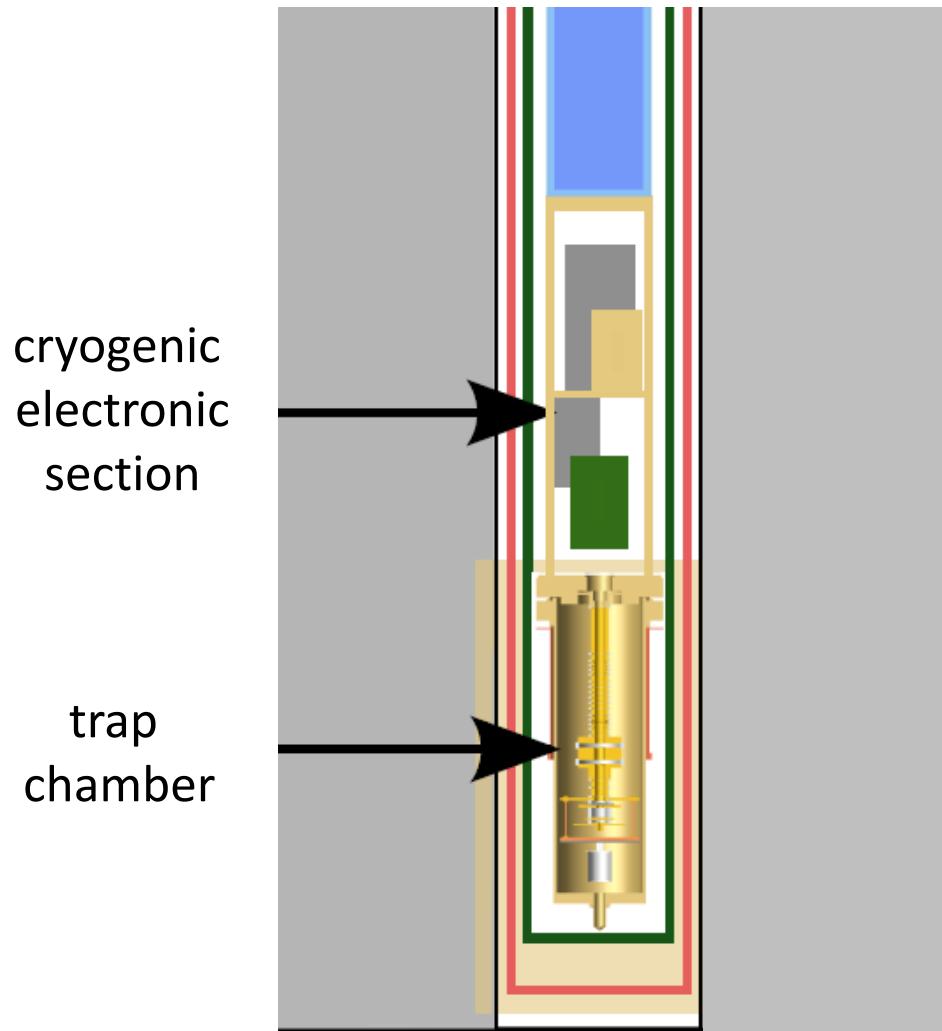


Outline

- Setup and Experimental Techniques
 - a new ultra-harmonic trap
- Proton's Atomic Mass
F. Heiße *et al.*, PRL **119**, 033001 (2017)
- Deuteron's Atomic Mass

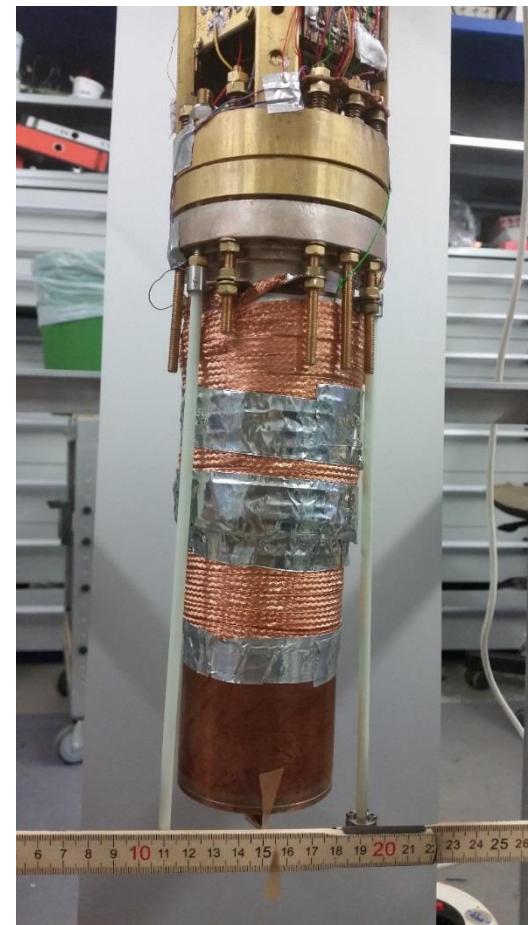
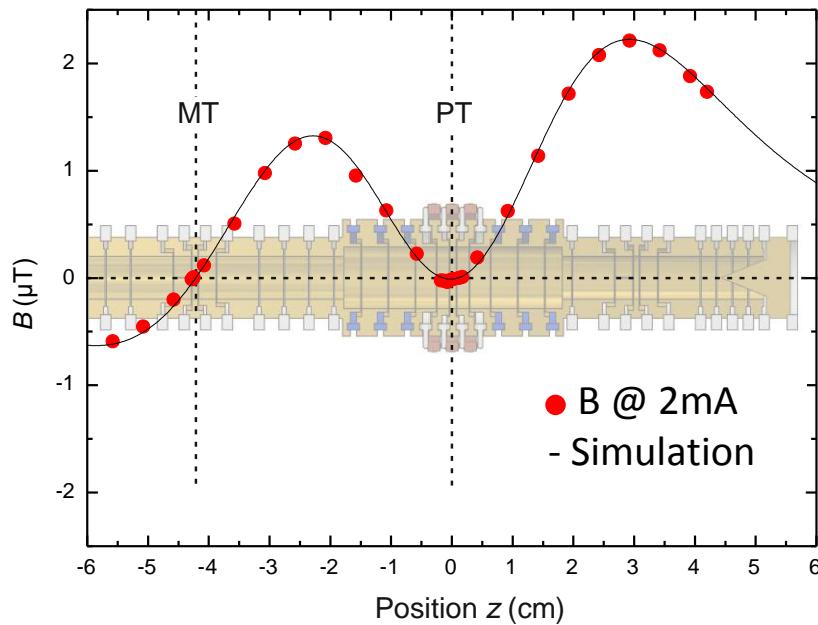
Solenoidal Superconducting B_2 Shim Coil

- placed around trap chamber:

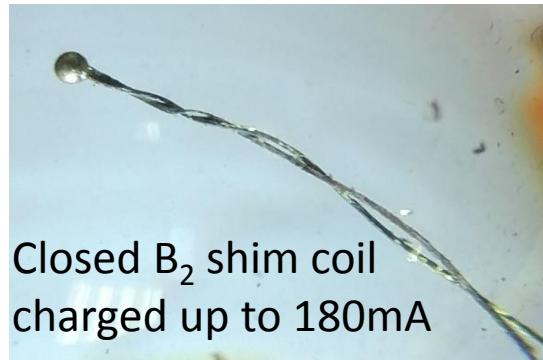


Solenoidal Superconducting B_2 Shim Coil

measurement in a testsetup

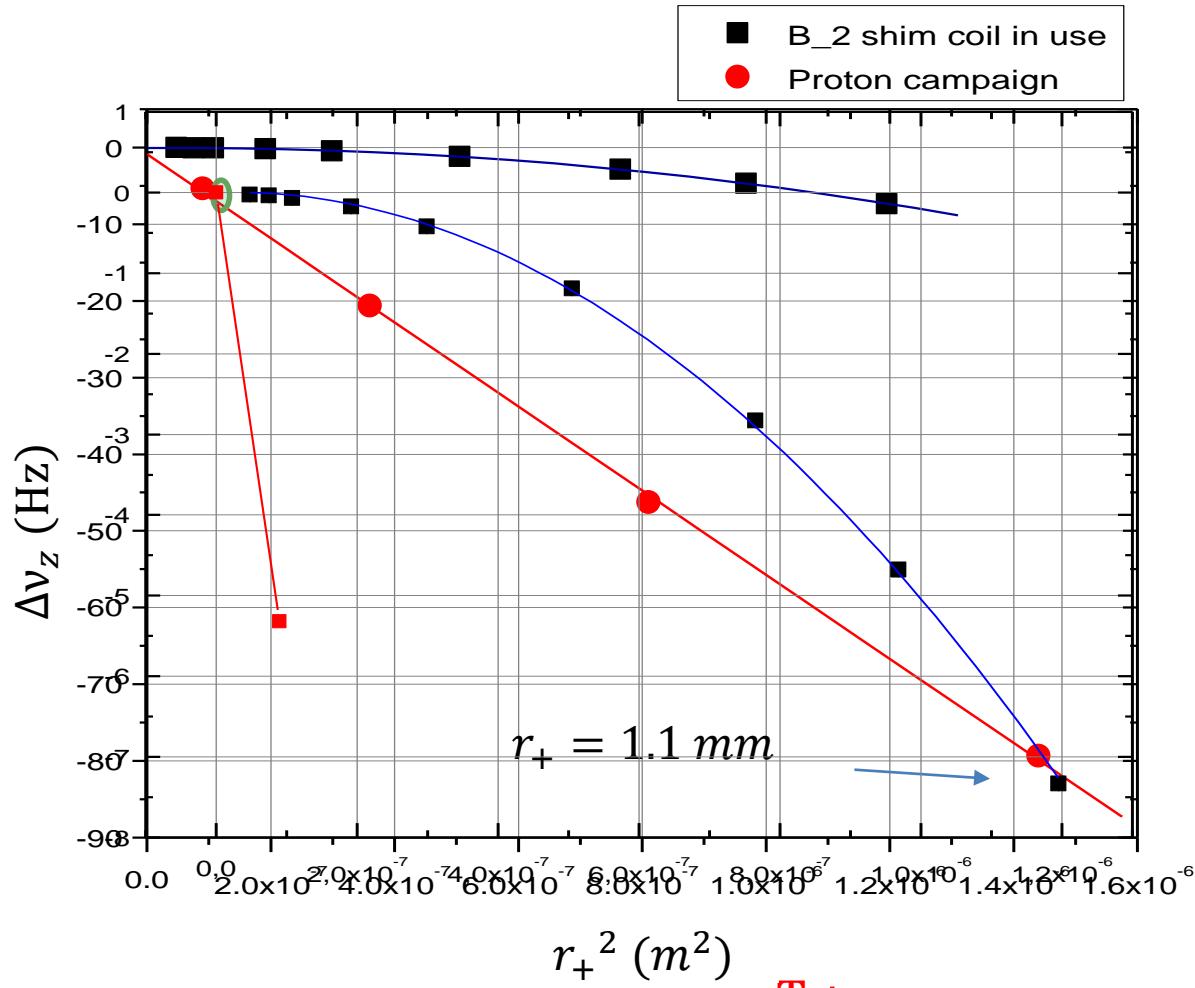


Superconducting Connection:



B₂ Shim Coil Performance

axial frequency shift as function of energy in ν_+ -mode $\sim r_+^2$



Proton mass campaign: $B_2 = -0.27(2) \mu\text{T}/\text{mm}^2$
now: $B_2 = -0.00127(40) \mu\text{T}/\text{mm}^2$

⇒ Factor 200 improvement

Estimated Error Budget Deuteron

Effect	Rel. Unc. (10^{-12}) Pmass	Rel. Unc. (10^{-12}) Deuteron (predicted)
Residual Magnetostatic Inhomogeneity	27.5	0.2
Special Relativity	7.1	2.9
Image Charge	4.6	4.1
Quadratic Sum	≈ 29	≈ 5.0

- current best value for the deuteron's atomic mass:

$$\delta m/m = 20 \times 10^{-12}$$

S. L. Zafonte and R. S. Van Dyck Jr., Metrologia 52, 280 (2015)

Schuh *et al.*, PRA, accepted

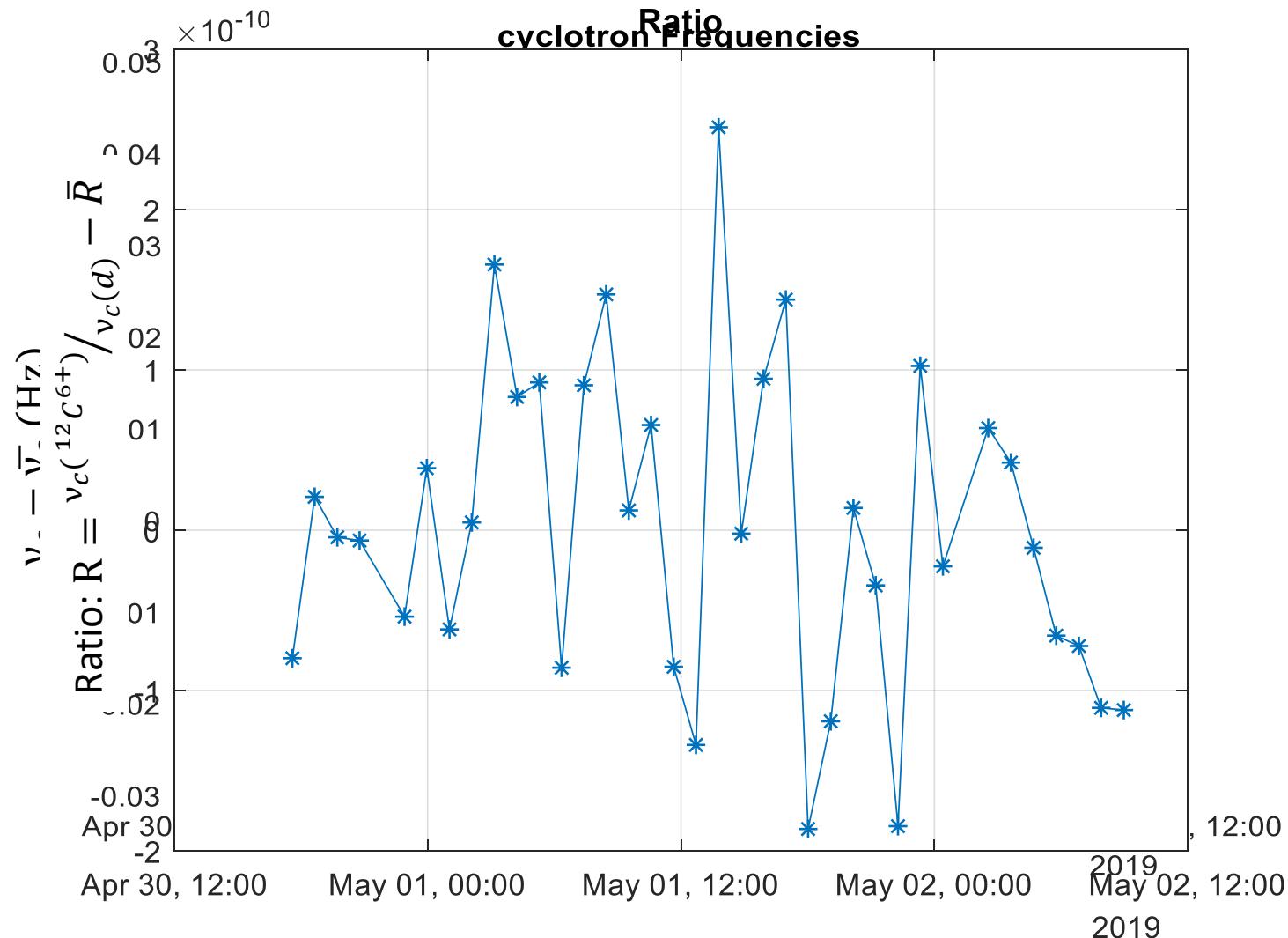
- current best mass ratio: $R = m(^{29}\text{Si}^+)/m(^{28}\text{SiH}^+)$

$$\delta R/R = 6.5 \times 10^{-12}$$

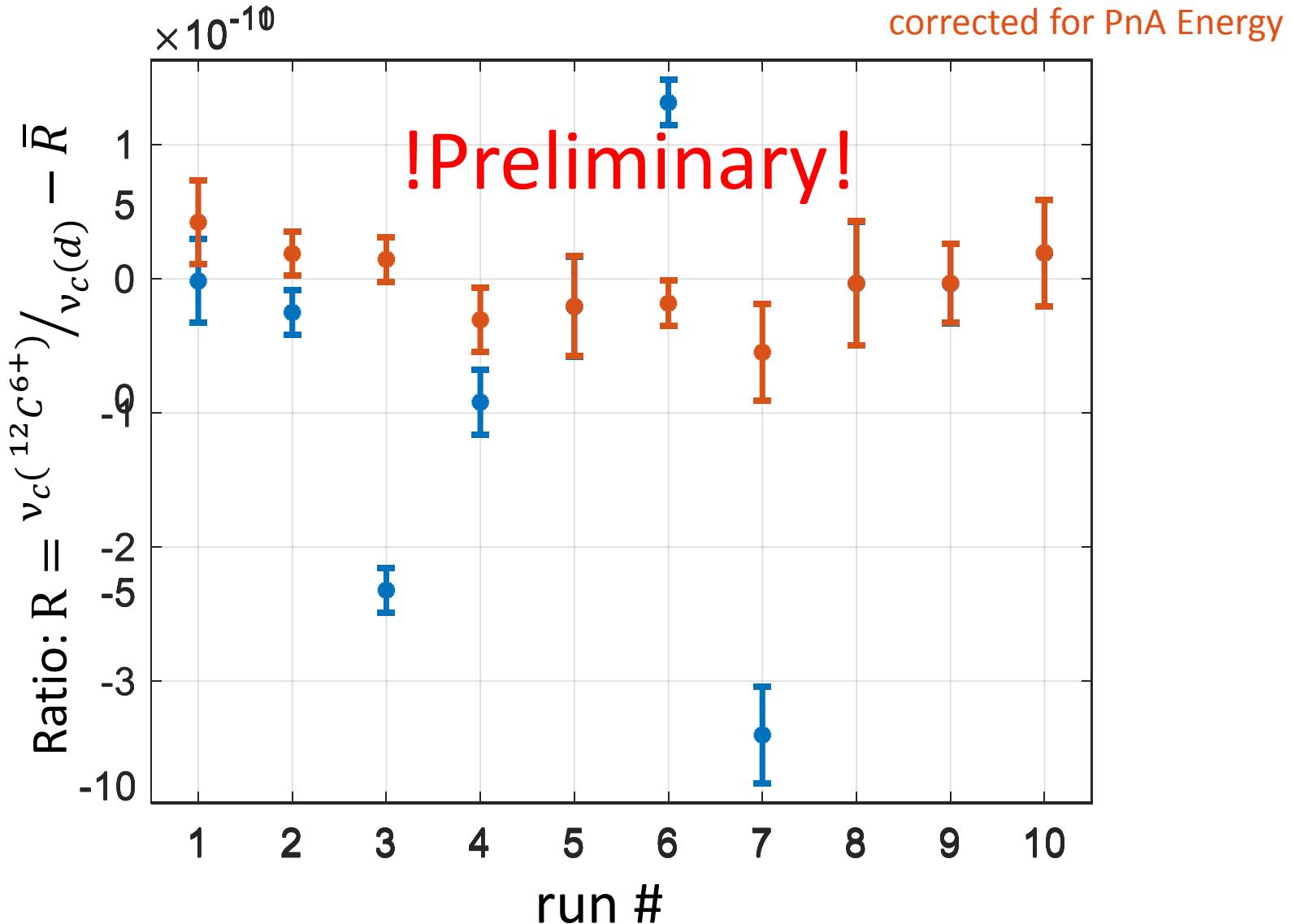
*S. Rainville *et al.*, Nature 438, 1096 (2005)*

Preliminary Results

example run from 30.04.2019 to 02.05.2019



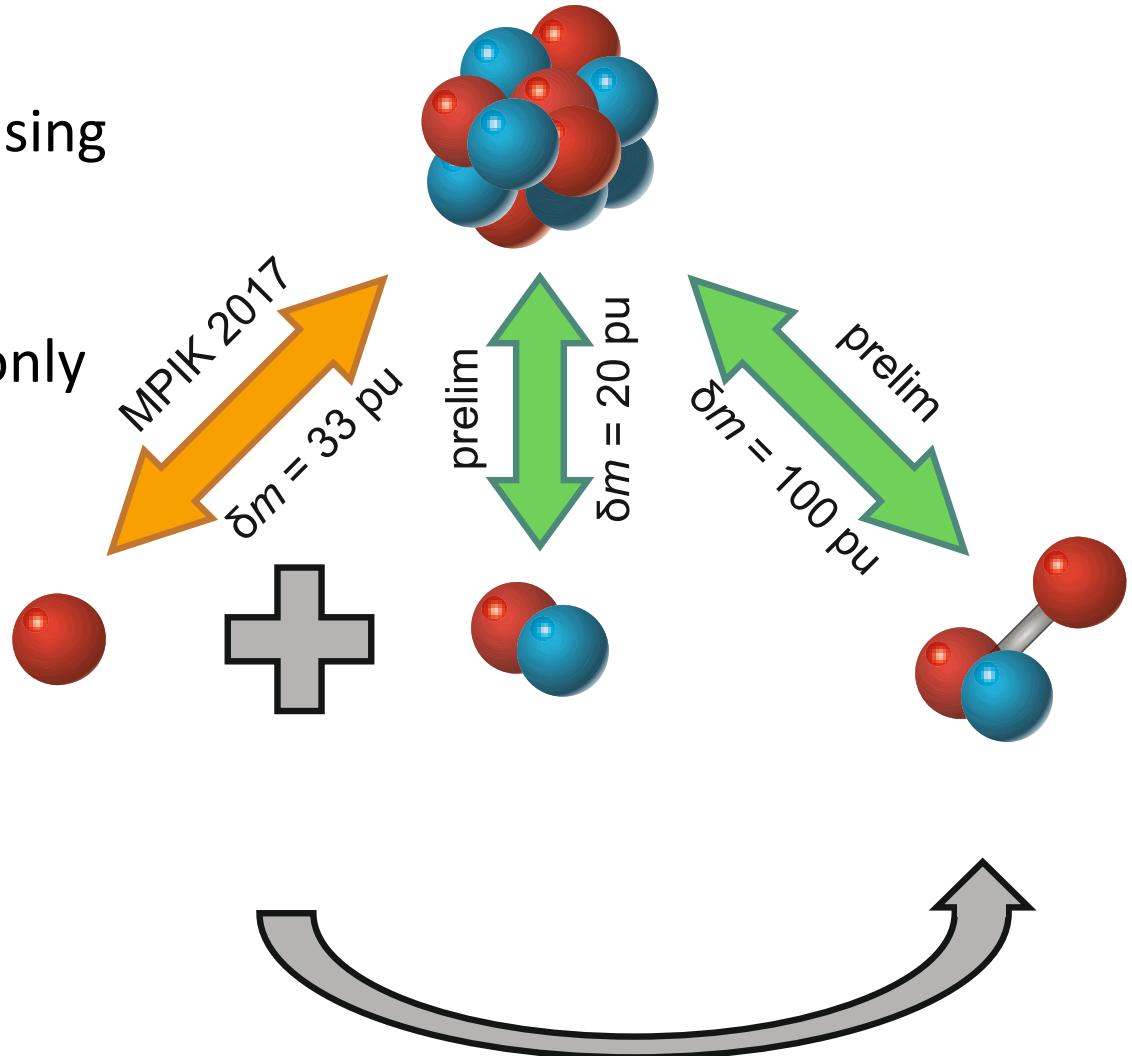
Preliminary Results



Zafonte and Van Dyck Jr., Metrologia 52, 280 (2015)

The mass of HD⁺

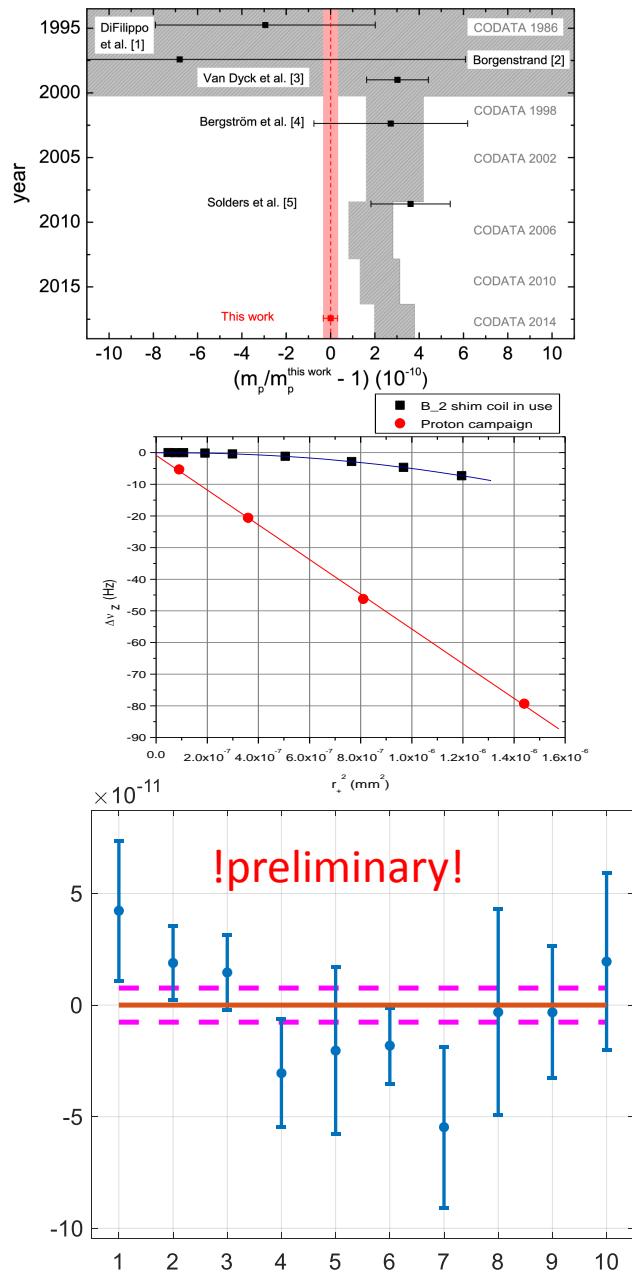
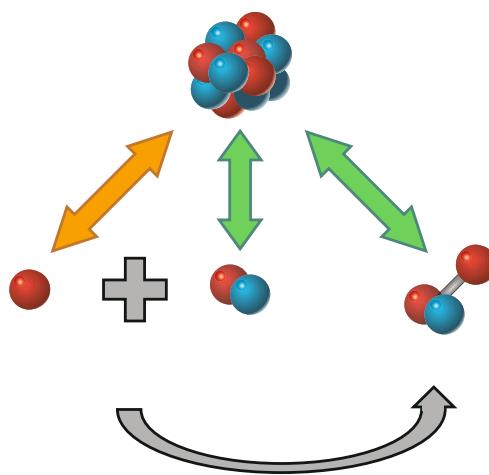
- consistency check using HD⁺
- closed circle using only our values



HD atomic & molecular Binding Energy:
Korobov et al., Phys. Rev. Lett. 118, 233001 (2017)

Summary

- proton mass with 33 ppt precision
- leading order systematic effect made negligible with B_2 shim coil
- estimated systematic uncertainty for deuteron to be around 5×10^{-12}
- ongoing measurement, statistical error around 8×10^{-12}
- strong consistency check using HD^+



Acknowledgement

Collaboration: Sascha Rau, Fabian Heiße^{1,2}, Florian Köhler-Langes¹, Wolfgang Quint², Sven Sturm¹ and Klaus Blaum¹

¹*Division of Stored and Cooled Ions at MPIK, Heidelberg*

²*Atomic Physics Division at GSI Helmholtzzentrum, Darmstadt*



Target: Raphael Haas^{2,3,4}, Dennis Renisch^{3,4} and
Christoph Düllmann^{2,3,4}

³*Institut for nuclear chemistry at JGU, Mainz*

⁴*Helmholtz Institute Mainz*



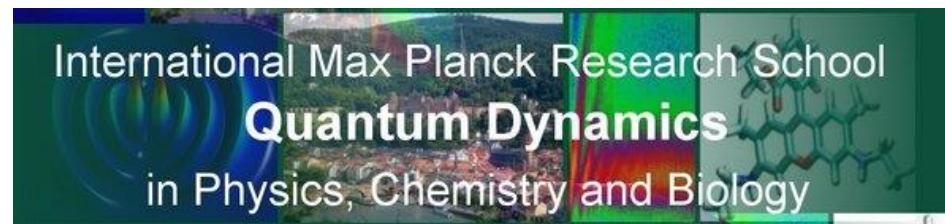
Funding:



INTERNATIONAL
MAX PLANCK
RESEARCH SCHOOL

PT
FS

FOR PRECISION TESTS
OF FUNDAMENTAL
SYMMETRIES



MAX-PLANCK-GESELLSCHAFT

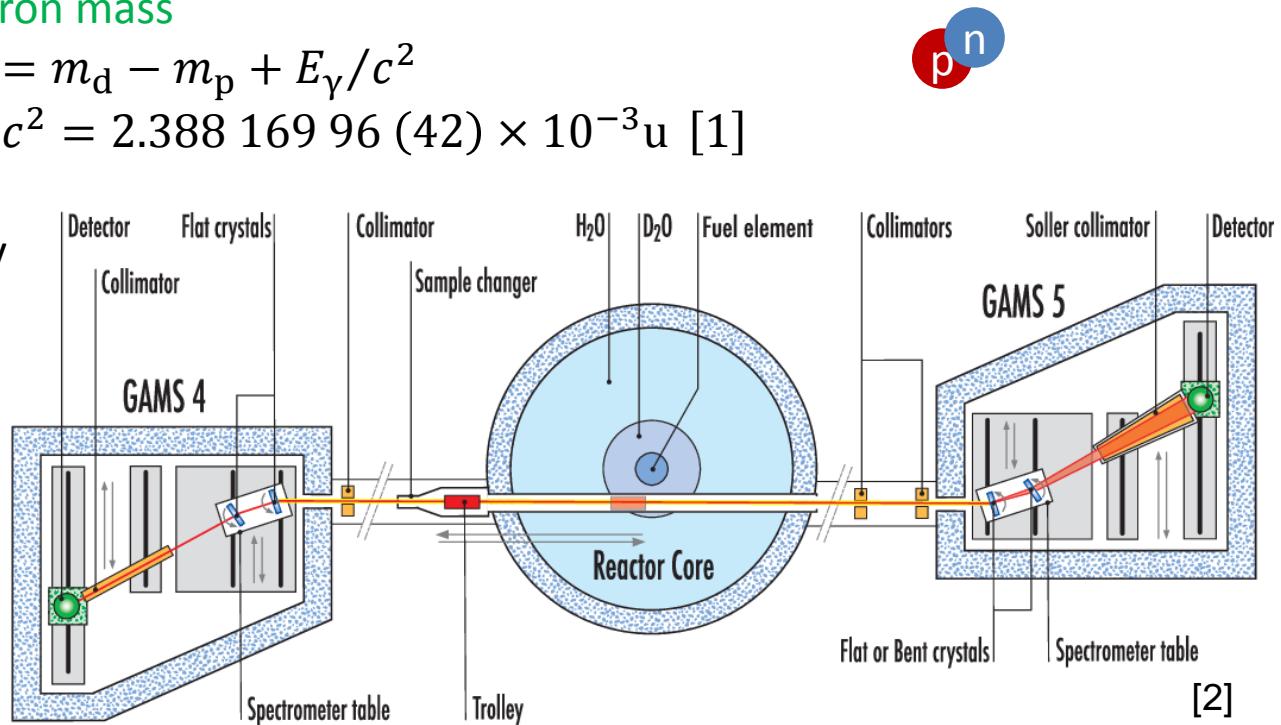


Why deuteron?

- Deuteron mass might help with puzzle of light ion masses
- allows for easier measurement of the Image Charge Shift (q/m doublet)
- one can extract the neutron mass

$$m_n = m_d - m_p + E_\gamma/c^2$$
$$E_\gamma/c^2 = 2.388\ 169\ 96(42) \times 10^{-3} \text{u} \ [1]$$

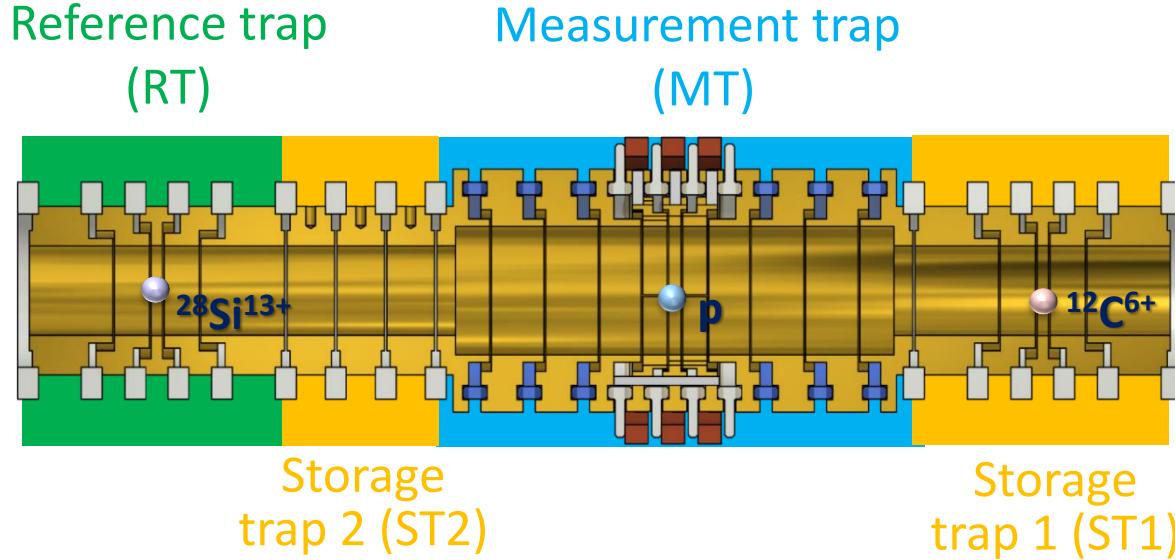
- improved binding energy needed, currently measured at the ILL in France



[1] M.S. Dewey et al. Physical Review C 73, 044303 (2006)

[2] <https://www.ill.eu/instruments-support/ „PN3-GAMS“>

Reference Trap



Simultaneous comparison of cyclotron frequencies to reference ion

(I) $v_c(^{28}\text{Si}^{13+})$ and $v_c(\text{p})$

$$R_1 = \frac{v_c(\text{p})}{v_c(^{28}\text{Si}^{13+})}$$

(II) $v_c(^{28}\text{Si}^{13+})$ and $v_c(^{12}\text{C}^{6+})$ $R_2 = \frac{v_c(^{12}\text{C}^{6+})}{v_c(^{28}\text{Si}^{13+})}$

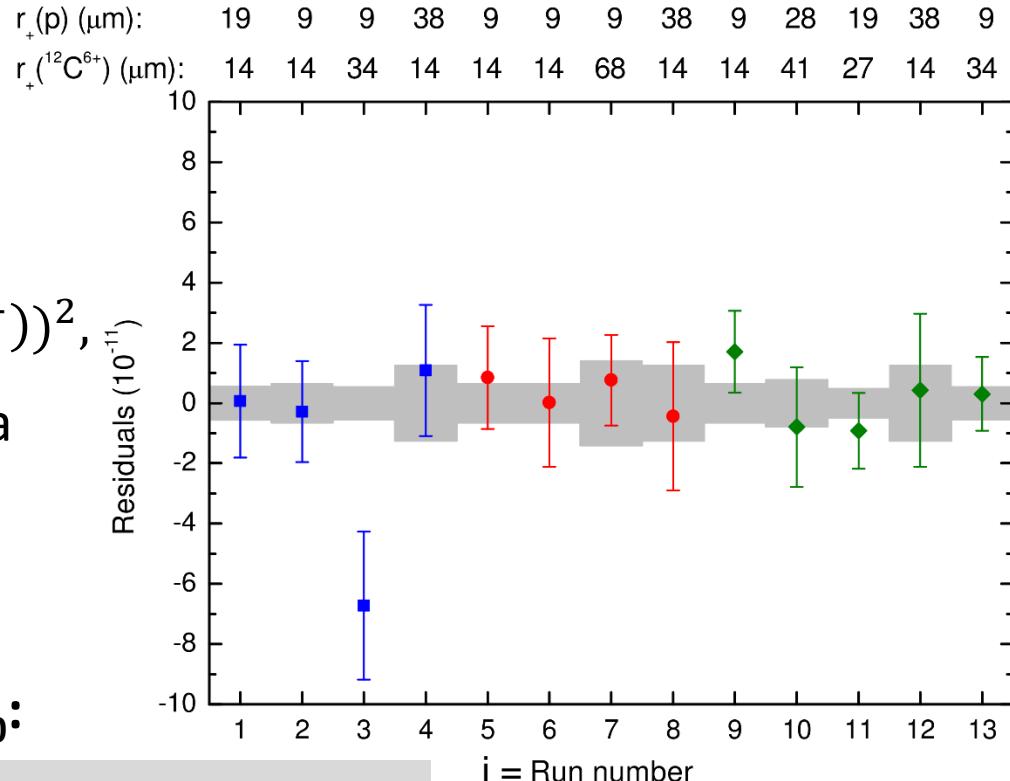
$$\boxed{\frac{R_1}{R_2} = \frac{v_c(\text{p})}{v_c(^{12}\text{C}^{6+})}}$$

→ cancel impact of common mode magnetic field fluctuations

Data

- Several 100's of measurement cycles have been performed with:

- three different p / $^{12}\text{C}^{6+}$ pairs
- at different $r_+(p) = \kappa \cdot A(p)$ and $r_+(\text{C}^{6+}) = \kappa' \cdot A(\text{C}^{6+})$.



- Apply a global planar fit:

$$R_i = R_0 + a (A_i(p))^2 + b (A_i(\text{C}^{6+}))^2,$$

where i is the run number and R_0 , a and b are fit parameters.

- Extrapolated frequency ratio R_0 :

$$R_0 = R_{\text{stat}} = \frac{\nu_c(\text{C}^{6+})}{\nu_c(p)} = 0.503\ 776\ 367\ 634\ 1 (77).$$

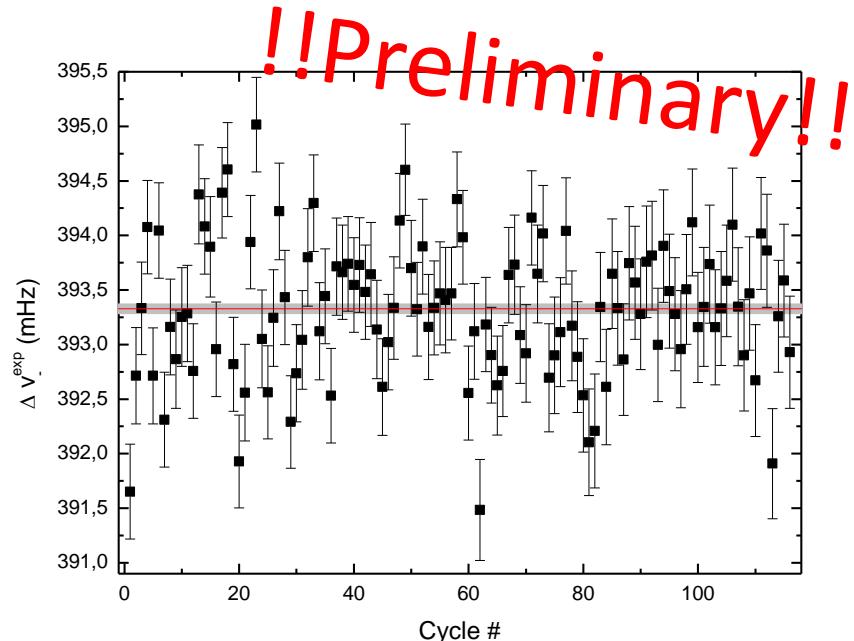
$$\frac{\delta R_0}{R_0} = 1.5 \cdot 10^{-11}$$



The Image Charge Shift

$$\Delta v_- = v_-(\text{Proton}) - v_-(^{12}\text{C}^{6+})$$

Effect	Shifts Δv_- [mHz]
ICS	2.394
B_2	0.076
Tilt	0.162
q/m missmatch	390.714
Total	393.346



BUT:

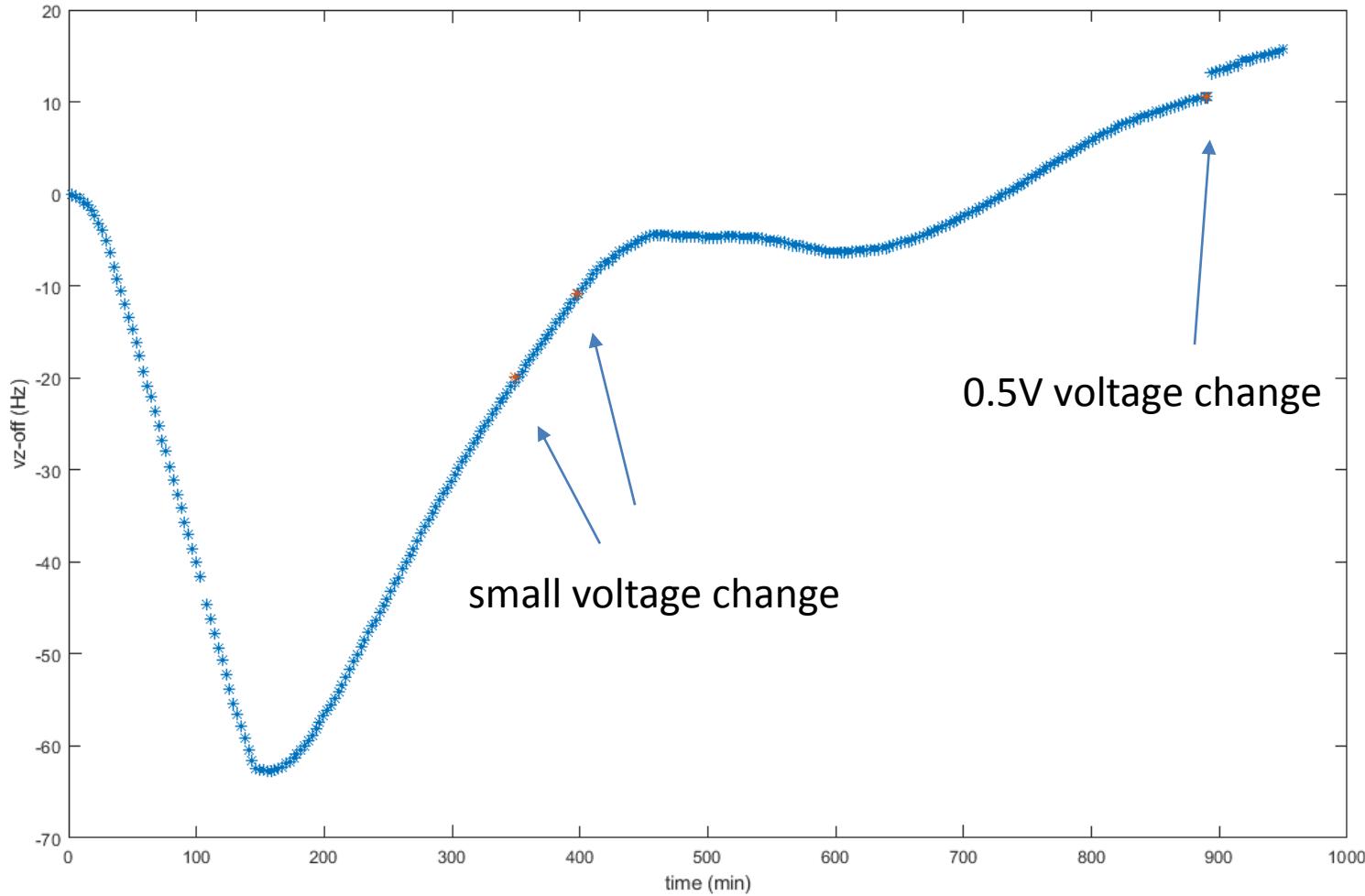
- Tilt between magnetic and electric field (0.5°)
- Variations of this tilt are in the axial mode indistinguishable from voltage fluctuations

$$\begin{aligned}\delta v_-^{\text{tilt fluc.}}(p) &\approx 1.4 \cdot 10^{-5} \cdot \delta v_z^{\text{tilt fluc.}}(p) \\ \delta v_-^{\text{volt. fluc.}}(p) &\approx 1.3 \cdot 10^{-2} \cdot \delta v_z^{\text{volt. fluc.}}(p)\end{aligned}$$

- Fit gives $\delta(\Delta v_-) = 43 \mu\text{Hz}$ with Red $\chi^2 = 2$

⇒ confirmation limited to 5% level ($120 \mu\text{Hz}$)

Improving axial stability



- leakage currents in the order of a few fA, resistance to ground $\sim 10^{15} \Omega$
- transistors DC resistance at least in same magnitude, acts as very small capacity (17fF)

Improving axial stability

- improve voltage stability by using a filter with a switchable time constant
- use „freeze-out“ of silicon transistor (BF545) happening at $\sim 20\text{ K}$:



Characteristics:

Conducting at: 13 mW for 20 s

Freezing out: 3 min

$R_{\text{cold}} > 500\text{ G}\Omega$

- Fluctuations in angle lead to instabilities:

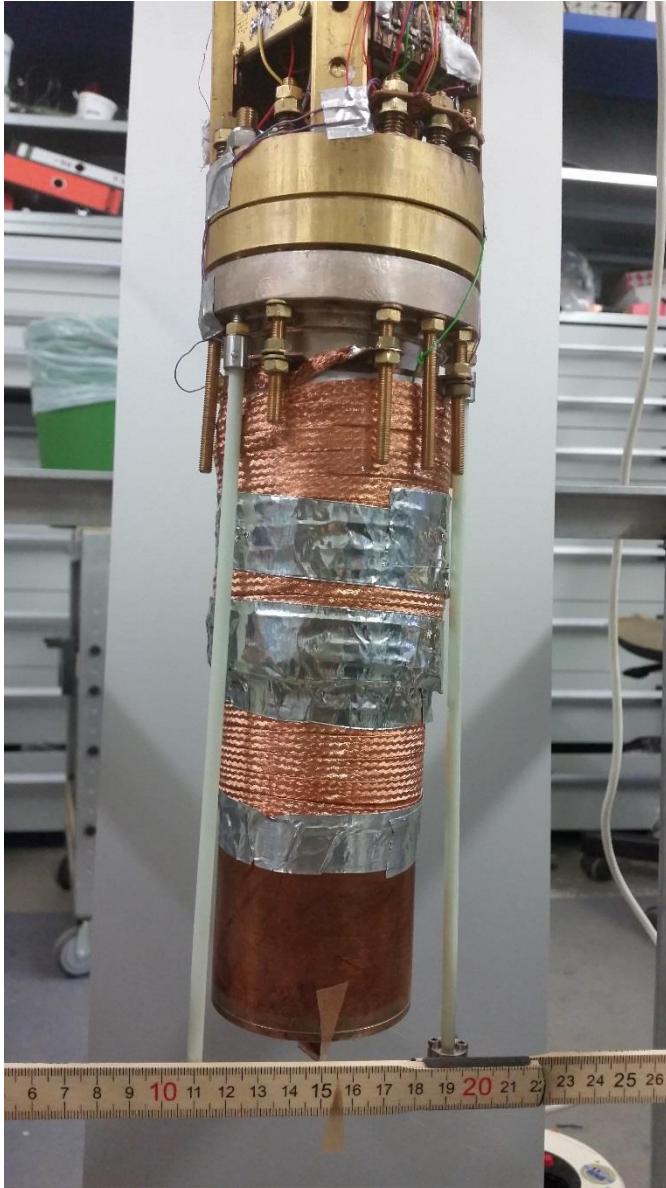
$$\bar{v_z} \approx v_z \sqrt{1 - \frac{3}{2} \sin^2 \theta}$$

$$\theta = \theta_0 + \delta\theta$$

$$\delta\bar{v_z} \approx -\frac{3}{2} v_z \theta_0 \delta\theta$$

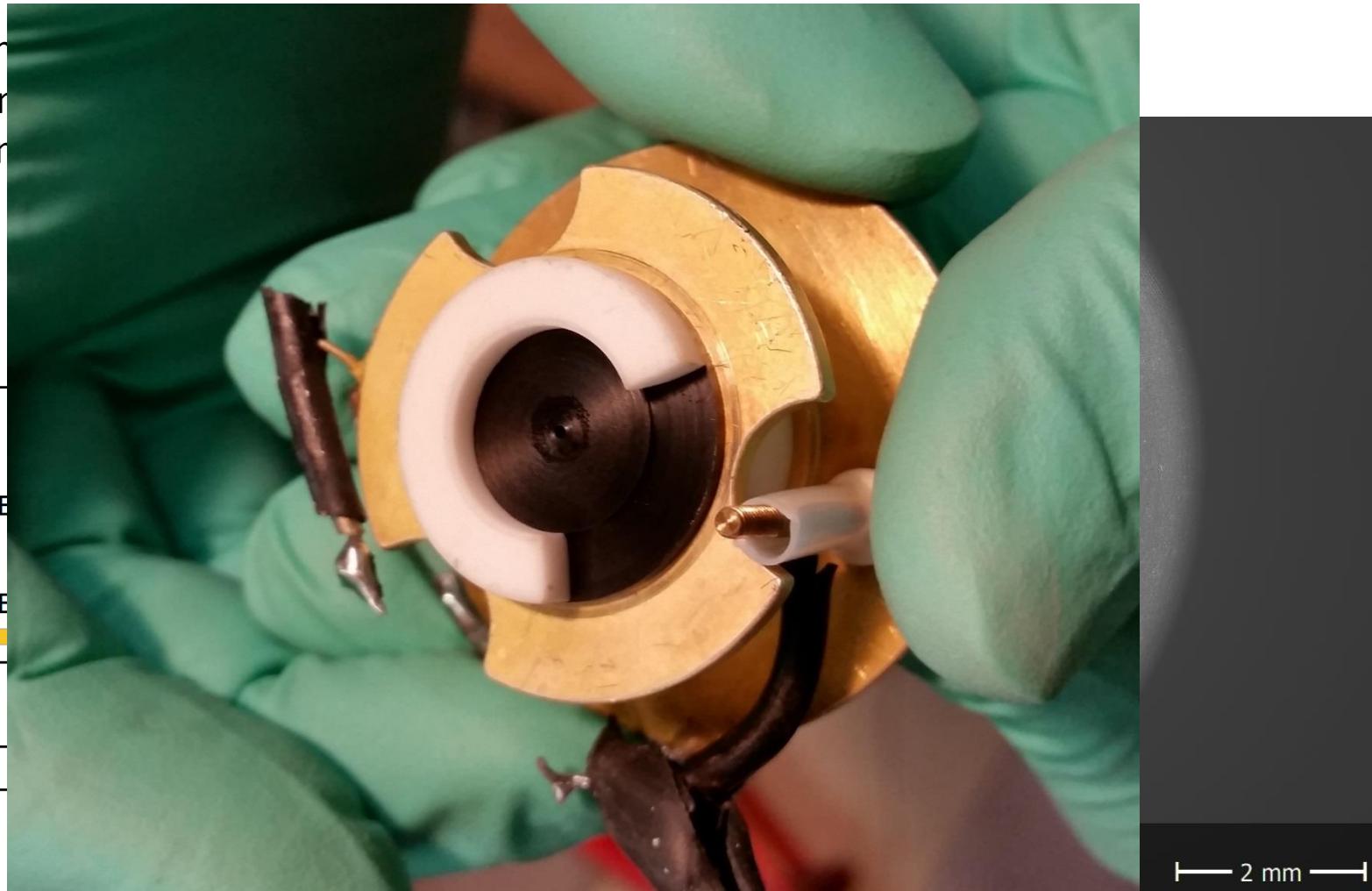
⇒ Mechanical Feedthrough to adjust angle to zero in situ

Mechanical Feedthrough



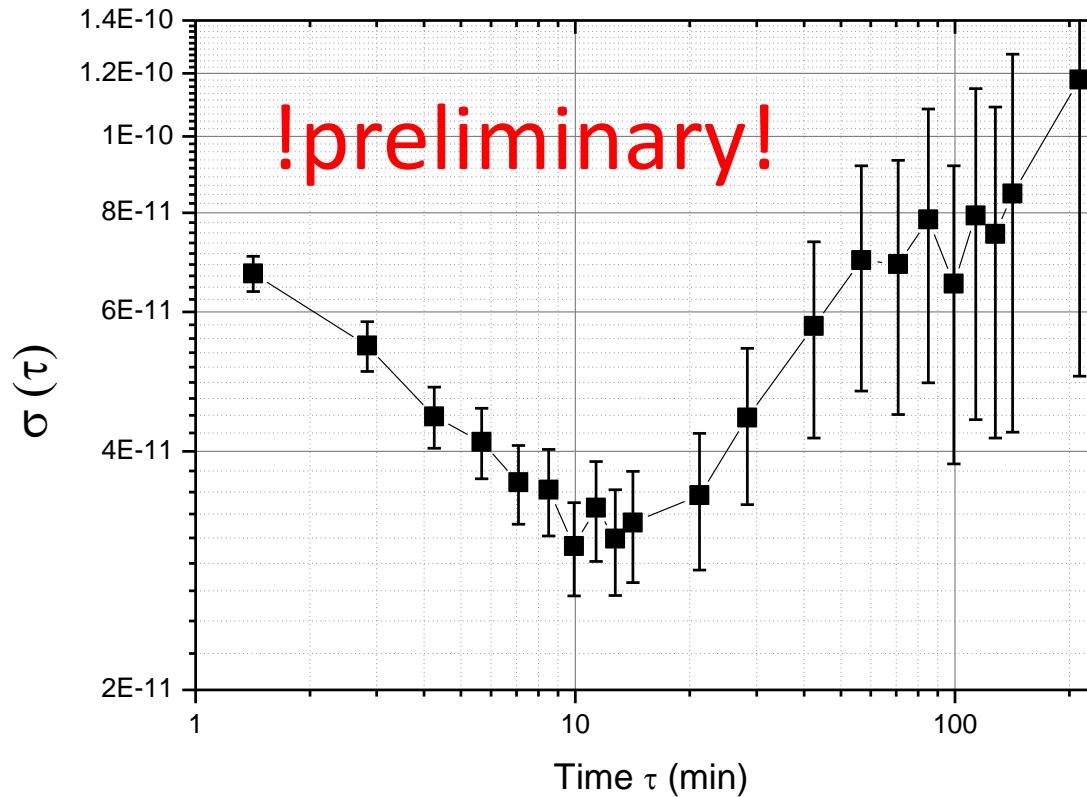
Target Preparation

- use target printer to apply deuterated molecules (Thymidine- $\alpha,\alpha,\alpha,6-d_4$) to surface of standard mEBIS target *R. Haas et al., Nucl. Instr. Meth. Phys. Res. A 874*
- small number of deuterium atoms while remaining in vacuum



Improved Statistics

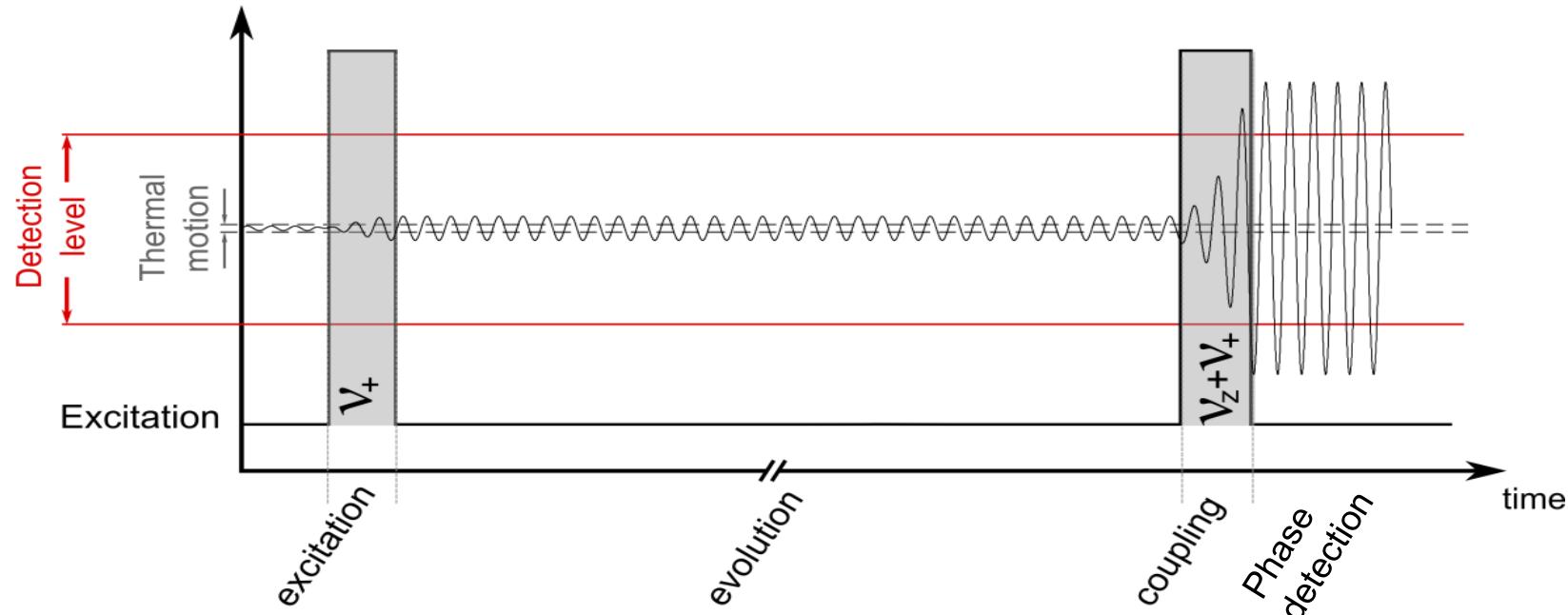
- implemented pressure and temperature stabilization
- Allan deviation with consecutive v_+ measurements with a single deuteron promises statistics in the few 10^{-11} per cycle



Pulse and Amplification – Method

- Phase sensitive detection method for v_+

Reduced cyclotron amplitude



- Rapid measurement time (~ 10 s instead of ~ 3 min for Double-Dip)
→ **Reduction of impact of B-field fluctuations**
- Small radial kinetic energies during phase evolution
→ **Small magnetic and relativistic shifts**