

Nonrelativistic energy levels of D_2

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Collaboration

- Faculty of Physics, University of Warsaw
 - Krzysztof Pachucki
 - Paweł Czachorowski
- Faculty of Chemistry, Adam Mickiewicz University, Poznan
 - Mariusz Puchalski

Energy of an atom or molecule

$\alpha \approx 1/137$ – fine-structure constant

$$E(\alpha) = m\alpha^2 E^{(2)} + m\alpha^4 E^{(4)} + m\alpha^5 E^{(5)} + m\alpha^6 E^{(6)} + \dots$$

- The expansion coefficients can be expressed by expectation values of effective Hamiltonians with **the nonrelativistic wave function** – NRQED

$E^{(2)}$	– nonrelativistic energy	36 749.0910(2)
$E^{(4)}$	– relativistic correction	–0.5276(3)
$E^{(5)}$	– the leading QED correction	–0.1983(2)
$E^{(6)}$	– higher order QED correction	–0.0016(8)
E	– total energy	36 748.3633(9)

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$E^{(5)}$ – the leading QED correction	-0.1983(2)	$1 \cdot 10^{-3}$
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NAPT vs DNVM

NAPT – nonadiabatic perturbation theory

- + all the bound states "at once"
- accuracy limited by cutting the perturbational series

DNVM – direct nonadiabatic variational method

- + "unlimited" accuracy
- state-by-state calculations

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NAPT – nonadiabatic perturbation theory

Pachucki 2009

$$H^{(2)} \Psi(\vec{r}, \vec{R}) = E^{(2)} \Psi(\vec{r}, \vec{R})$$

$$H^{(2)} = H_{\text{el}} + H_{\text{n}}$$

$$\Psi(\vec{r}, \vec{R}) = \phi_{\text{el}}(\vec{r}; \vec{R}) \chi(\vec{R}) + \delta\phi_{\text{na}}(\vec{r}, \vec{R})$$

$$H_{\text{el}} \phi_{\text{el}} = \mathcal{E}_{\text{el}}(\vec{R}) \phi_{\text{el}}$$

$$\langle \delta\phi_{\text{na}} | \phi_{\text{el}} \rangle_{\text{el}} = 0$$

$$\left[(H_{\text{el}} - \mathcal{E}_{\text{el}}) + (\mathcal{E}_{\text{el}} + H_{\text{n}} - E^{(2)}) \right] \left| \phi_{\text{el}}(\vec{r}; \vec{R}) \chi(\vec{R}) + \delta\phi_{\text{na}}(\vec{r}, \vec{R}) \right\rangle = 0$$

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NAPT(2)

Pachucki 2009

Perturbative expansion for the effective nuclear Hamiltonian

$$(\mathcal{E}_{\text{el}} + \mathcal{E}_{\text{a}} + H_{\text{n}} - E^{(2)})|\chi\rangle = -(H_{\text{n}}^{(2)} + H_{\text{n}}^{(3)} + H_{\text{n}}^{(4)} + \dots)|\chi\rangle$$

$$H_{\text{n}}^{(2)} \sim \left(\frac{m_{\text{e}}}{\mu_{\text{n}}}\right)^2$$

$$H_{\text{n}}^{(3)} \sim \left(\frac{m_{\text{e}}}{\mu_{\text{n}}}\right)^2 + \mathcal{O}\left[\left(\frac{m_{\text{e}}}{\mu_{\text{n}}}\right)^3\right]$$

$$H_{\text{n}}^{(4)} = \mathcal{O}\left[\left(\frac{m_{\text{e}}}{\mu_{\text{n}}}\right)^3\right]$$

$$\delta E^{(2)} = E^{(2)} - E_{\text{NAPT}(2)}^{(2)} \sim \left(\frac{m_{\text{e}}}{\mu_{\text{n}}}\right)^3$$

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Perturbative expansion for the effective nuclear Hamiltonian

$$(\mathcal{E}_{\text{el}} + \mathcal{E}_a + H_n - E^{(2)})|\chi\rangle = -(H_n^{(2)} + H_n^{(3)} + H_n^{(4)} + \dots)|\chi\rangle$$

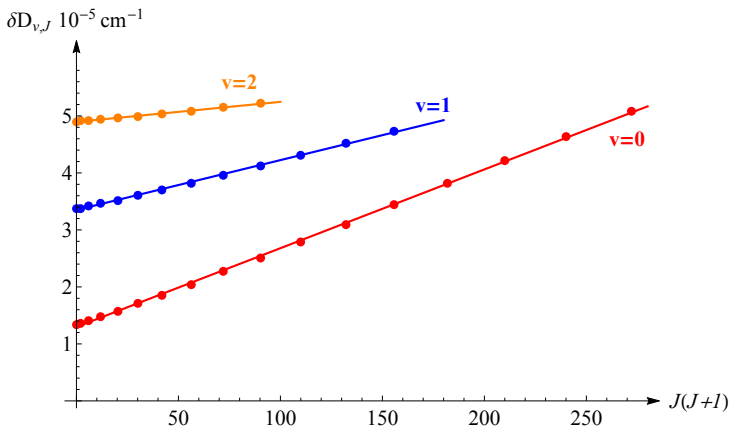
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$$\delta E^{(2)} = E^{(2)} - E_{\text{NAPT}}^{(2)} \quad \text{for } D_2$$



James-Coolidge-type functions

The James-Coolidge-type basis function

$$\psi_{\{n\}}(\vec{r}_1, \vec{r}_2) = (1 \pm \hat{P}_{12})(1 \pm \hat{i})$$

$$(r_{1A} - r_{1B})^{n_2} (r_{2A} - r_{2B})^{n_3} (r_{1A} + r_{1B})^{n_4} (r_{2A} + r_{2B})^{n_5}$$

$$e^{-\beta(r_{1A} + r_{1B} + r_{2A} + r_{2B})} r_{12}^{n_1} R^{-n_1 - n_2 - n_3 - n_4 - n_5 - 3}$$

Closed-form integrals f derived by K. Pachucki^a

$$f(\{n\}; R, \beta) = R^{-n_1 - n_2 - n_3 - n_4 - n_5 - 1}$$

$$\int \frac{d^3 r_1}{4\pi} \int \frac{d^3 r_2}{4\pi} \frac{e^{-\beta r_{1A}}}{r_{1A}} \frac{e^{-\beta r_{1B}}}{r_{1B}} \frac{e^{-\beta r_{2A}}}{r_{2A}} \frac{e^{-\beta r_{2B}}}{r_{2B}} \frac{1}{r_{12}^{1-n_1}}$$

$$(r_{1A} - r_{1B})^{n_2} (r_{2A} - r_{2B})^{n_3} (r_{1A} + r_{1B})^{n_4} (r_{2A} + r_{2B})^{n_5}$$

^aK. Pachucki, Phys. Rev. A 80, 032520 (2009).

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Advertisement

H2SPECTRE

- Fortran code with the NAPT implemented made public!
- Publication with theoretical background and usage
 - in preparation 8(
- AMUR - Adam Mickiewicz University Repository
 - URL will be announced in the publication
 - for now contact Pawel.Czachorowski@fuw.edu.pl

```

komas@e0311:~/Iso_naqed19/h2spectr_4.5$ ./H2spectre.exe -V 2 --napt
D2 Level
(Int. points= 200 , Int. range= 10.0)
-----
D2      Level
0 0
-----
v' = 0  J' = 0

Contribution  Value[cm-1]  Error[cm-1]
E2(NAPT)      36749.09098  8.6E-05
E4            -0.52821  1.2E-06
  E4(B0)      -0.52918
  E4(rec)     -0.00120
  E4(2nd)     -0.00023
E5            -0.19831  1.5E-04
E6            -0.00209  6.2E-06
  E6(B0)     -0.00210
  E6(2nd)     0.00001
E7            0.00010  2.6E-05
Efs          -0.00020  1.2E-07
Total        36749.36227  1.7E-04
-----

```

```

komas@e0311:~/Iso_naqed19/h2spectr_4.5$ ./H2spectre.exe -V 2 --napt --wavl
D2 Transition
(Int. points= 200 , Int. range= 10.0)
-----
D2      Transition
0 1  0 0
-----
v' = 0  J' = 1  -->  v'' = 0  J'' = 0

Contribution  Value[cm-1]  Error[cm-1]
E2(NAPT)      59.7798252  9.3E-07
E4            0.0013106  2.5E-09
  E4(B0)      0.0013098
  E4(rec)     0.0000015
  E4(2nd)    -0.0000007
E5            -0.0005170  3.9E-07
E6            -0.0000045  1.3E-08
  E6(B0)     -0.0000045
  E6(2nd)     0.0000000
E7            0.0000002  6.2E-08
Efs          -0.0000005  2.8E-10
Total        59.7806141  1.0E-06
-----
Wavelength[um] 167.2783084  2.8E-06
-----

```


DNVM – direct nonadiabatic variational method

Four-body nonadiabatic Schrödinger equation solved variationally

The trial wave function

$$\Psi(\vec{r}_1, \vec{r}_2, \vec{R}_A, \vec{R}_B) = \sum_{k=1}^K c_k \hat{S} \psi_{\{k\}}(\vec{r}_1, \vec{r}_2, \vec{R}_A, \vec{R}_B)$$

Four-particle basis of exponential functions (*naJC*)

$$\psi_{\{k\}} = \exp[-\alpha r_{AB} - \beta (r_{1A} + r_{1B} + r_{2A} + r_{2B})] \\ \times r_{AB}^{k_0} r_{12}^{k_1} (r_{1A} - r_{1B})^{k_2} (r_{2A} - r_{2B})^{k_3} (r_{1A} + r_{1B})^{k_4} (r_{2A} + r_{2B})^{k_5}$$

Schrödinger equation in a matrix form

$$(\mathbb{H} - E_K \mathbb{N}) \mathbf{c} = 0$$

$$E_K \xrightarrow{K \rightarrow \infty} E^{(2)}$$

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Schrödinger equation in a matrix form

$$(\mathbb{H} - E_K \mathbb{N}) c = 0$$

$$E_K \xrightarrow{K \rightarrow \infty} E^{(2)}$$

Nonadiabatic wave function for an arbitrary J

$$\Psi^{J,M} = \Psi_{\Sigma}^{J,M} + \Psi_{\Pi}^{J,M} + \Psi_{\Delta}^{J,M} + \dots$$

where

$$\Psi_{\Sigma}^{J,M} = Y_J^M \Phi_{\Sigma}^J, \quad \text{for } J \geq 0$$

$$\Psi_{\Pi}^{J,M} = \rho^i (\nabla_R^i Y_J^M) \Phi_{\Pi}^J, \quad \text{for } J \geq 1$$

$$\Psi_{\Delta}^{J,M} = (\rho^i \rho'^j)^{(2)} \left(\nabla_R^i \nabla_R^j Y_J^M \right) \Phi_{\Delta}^J, \quad \text{for } J \geq 2$$

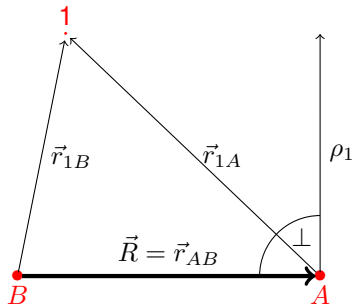
using Einstein notation and the following symbols

$$(\rho^i \rho'^j)^{(2)} \equiv \frac{1}{2} \left(\rho^i \rho'^j + \rho^j \rho'^i - (\delta^{ij} - n^i n^j) \vec{\rho} \cdot \vec{\rho}' \right)$$

$$\vec{\rho}, \vec{\rho}' \equiv \vec{\rho}_1 \text{ or } \vec{\rho}_2 \quad n^i \equiv R^i / R$$

$$\rho_a^i = (\delta^{ij} - n^i n^j) r_{aA}^j$$

Nonadiabatic wave function, cont.



$$\Phi_{\Lambda}^J = \sum_{\{k\}} c_{\{k\}} (1 + \mathcal{P}_{12}) \Phi_{\Lambda\{k\}}^J$$

primitive basis function

$$\Phi_{\{k\}} = e^{-\alpha R - \beta(\zeta_1 + \zeta_2)} R^{k_0} r_{12}^{k_1} \eta_1^{k_2} \eta_2^{k_3} \zeta_1^{k_4} \zeta_2^{k_5}$$

$$\zeta_a = r_{aA} + r_{aB}, \quad \eta_a = r_{aA} - r_{aB}, \quad \vec{R} = \vec{r}_{AB}$$

Four-body integrals

$$G = \int \frac{d^3 R}{4\pi} \int \frac{d^3 r_{1A}}{4\pi} \int \frac{d^3 r_{2A}}{4\pi} e^{-u_1 R - w_1 r_{12} - y \eta_1 - x \eta_2 - u \zeta_1 - w \zeta_2} \\ \times R^{n_0} r_{12}^{n_1} \eta_1^{n_2} \eta_2^{n_3} \zeta_1^{n_4} \zeta_2^{n_5} / \mathcal{D}$$

where

$$\mathcal{D} = R r_{12} r_{1A} r_{1B} r_{2A} r_{2B} = \frac{1}{16} R r_{12} (\zeta_1 + \eta_1) (\zeta_1 - \eta_1) (\zeta_2 + \eta_2) (\zeta_2 - \eta_2)$$

Master integral

$$g = \int \frac{d^3 R}{4\pi} \int \frac{d^3 r_{1A}}{4\pi} \int \frac{d^3 r_{2A}}{4\pi} \frac{e^{-u_1 R - w_1 r_{12} - y \eta_1 - x \eta_2 - u \zeta_1 - w \zeta_2}}{\mathcal{D}}$$

- D. M. Fromm and R. N. Hill, Phys. Rev. A 36, 1013 (1987)
 E. Remiddi, Phys. Rev. A, 44, 5492 (1991)
 F. E. Harris, Phys. Rev. A 55, 1820 (1997)

Master integral

Pachucki method

Pachucki differential equations¹

$$\sigma \frac{\partial g^{(n)}}{\partial a} + \frac{1}{2} \frac{\partial \sigma}{\partial a} g^{(n)} + P_a = 0$$

a – a nonlinear parameter u_1, w_1, y, x, u, w

σ – 6th degree polynomial in these parameters

P_a – a combination of logarithmic functions

For $naJC$ basis functions

$$g(w_1 = y = x = 0, w = u)$$

$$= \frac{1}{16 u^4} \left[\frac{-u^3 \ln(2u)}{t(t+2u)} + \frac{u^3 \ln(4u)}{(t-2u)(t+2u)} - \frac{2u^4 \ln(t+2u)}{t(t-2u)(t+2u)} \right]$$

¹ K. Pachucki, Phys. Rev. A **80**, 032520 (2009)

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Technical issues

- $\sim 10^5 \times 10^5$ dense matrices buildup and diagonalization
- quad-double precision (212 bit)
- up to 700 GB of RAM
- parallelization of diagonalization routine

Convergence in DNVM

Convergence of $E_{0,0}^{(2)}$ (a.u.) and $D_{0,0}$ (cm^{-1}) for the ground level of D_2

Ω	K	$E_{0,0}^{(2)}$	$D_{0,0}$
9	33 733	-1.167 168 809 240 572	36 749.090 980 266
10	49 959	-1.167 168 809 276 440	36 749.090 988 138
11	71 736	-1.167 168 809 282 247	36 749.090 989 411
12	100 772	-1.167 168 809 283 632	36 749.090 989 716
13	138 348	-1.167 168 809 283 929	36 749.090 989 780
14	186 660	-1.167 168 809 284 013	36 749.090 989 800
∞	∞	-1.167 168 809 284 05(4)	36 749.090 989 81(1)

^aS. Bubin, M. Stanke and L. Adamowicz, J. Chem. Phys., 135, 074110 (2011).

^bM. Puchalski, A. Spyszkiwicz, J. Komasa and K. Pachucki, Phys. Rev. Lett., 121, 073001 (2018).

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11	71 736	-1.167 168 809 282 247	36 749.090 989 411
12	100 772	-1.167 168 809 283 632	36 749.090 989 716
13	138 348	-1.167 168 809 283 929	36 749.090 989 780
14	186 660	-1.167 168 809 284 013	36 749.090 989 800
∞	∞	-1.167 168 809 284 05(4)	36 749.090 989 81(1)
	Ref. ^a	-1.167 168 809 21(5)	36 749.090 97(1)
	Ref. ^b	-1.167 168 809 27(4)	36 749.090 98(1)

^aS. Bubin, M. Stanke and L. Adamowicz, J. Chem. Phys., **135**, 074110 (2011).

^bM. Puchalski, A. Spyszkiwicz, J. Komasa and K. Pachucki, Phys. Rev. Lett., **121**, 073001 (2018).

Physical constants uncertainty

Constant	Rel. uncertainty	$\delta D_0/\text{cm}^{-1}$
Deuteron-electron mass ratio	$3.5 \cdot 10^{-11}$	$2.7 \cdot 10^{-8}$
Rydberg constant	$1.9 \cdot 10^{-12}$	$1.4 \cdot 10^{-7}$

CODATA 2014

$$m_d/m_e = 3\,670.482\,967\,85(13)$$

CODATA 2018

$$m_d/m_e = 3\,670.482\,967\,88(13)$$

Derivative over nuclear mass

PCCP 21, 10272 (2019)

(v, J)	$E_{v,J}^{(2)}/\text{a.u.}$	$D_{v,J}/\text{cm}^{-1}$	$\frac{\partial E_{v,J}^{(2)}}{\partial \mu_a} \cdot 10^6$
(0, 0)	- 1.167 168 809 284 05(4)	36 749.090 989 81(1)	
(0, 1)	- 1.166 896 432 359 76(4)	36 689.311 164 76(1)	
(0, 2)	- 1.166 352 930 288 83(4)	36 570.026 248 10(1)	
(0, 3)	- 1.165 540 786 627 44(4)	36 391.781 317 40(1)	
(0, 4)	- 1.164 463 677 472 61(4)	36 155.383 182 70(1)	
(0, 5)	- 1.163 126 414 451 1(2)	35 861.887 874 01(4)	
(0, 6)	- 1.161 534 872 009 1(2)	35 512.584 683 27(4)	

Derivative over nuclear mass

PCCP 21, 10272 (2019)

(v, J)	$E_{v,J}^{(2)}/\text{a.u.}$	$D_{v,J}/\text{cm}^{-1}$	$\frac{\partial E_{v,J}^{(2)}}{\partial \mu_d} \cdot 10^6$
(0, 0)	- 1.167 168 809 284 05(4)	36 749.090 989 81(1)	- 1.02
(0, 1)	- 1.166 896 432 359 76(4)	36 689.311 164 76(1)	- 1.10
(0, 2)	- 1.166 352 930 288 83(4)	36 570.026 248 10(1)	- 1.24
(0, 3)	- 1.165 540 786 627 44(4)	36 391.781 317 40(1)	- 1.46
(0, 4)	- 1.164 463 677 472 61(4)	36 155.383 182 70(1)	- 1.75
(0, 5)	- 1.163 126 414 451 1(2)	35 861.887 874 01(4)	- 2.10
(0, 6)	- 1.161 534 872 009 1(2)	35 512.584 683 27(4)	- 2.52

Comparison with experimental data for D_2

Dissociation energy of the ground level

Contrib.	$D_{0,0}$
$E^{(2)}$	36 749.090 989 81(1)
$E^{(4)}$	-0.528 206 05(9)
$E^{(5)}$	-0.198 256(3)
$E^{(6)}$	-0.002 087(6)
$E^{(7)}$	0.000 103(26)
E_{FS}	-0.000 202
E	36 748.362 342(26)
Exper. ^a	36 748.362 86(68)
Diff.	0.000 57(68)

^aJ. Liu, E. J. Salumbides, U. Hollenstein, J. C. J. Koelemeij, K. S. E. Eikema, W. Ubachs, F. Merkt, J. Chem. Phys., **130**, 174306 (2009).

Comparison with experimental data for D₂

Selected transition energies, $(v'', J'') \leftarrow (v', J')$

Contrib.	$(0, 2) \leftarrow (1, 2)$
$E^{(2)}$	2 987.291 387 53(1)
$E^{(4)}$	0.017 498
$E^{(5)}$	- 0.015 33(8)
$E^{(6)}$	- 0.000 137
$E^{(7)}$	0.000 007(2)
E_{FS}	- 0.000 015
E	2 987.293 41(9)
Exper.	2 987.293 52(15) ^a
Diff.	0.000 11(17)

^aM. Niu, E. Salumbides, G. Dickenson, K. Eikema, W. Ubachs, *J. Mol. Spectrosc.*, **300**, 44 (2014).

^bP. Wcislo, et al., *Journal of Quantitative Spectroscopy and Radiative Transfer*, **213**, 41 (2018).

^cD. Mondelain, S. Kassi, T. Sala, D. Romanini, D. Gatti, A. Campargue, *J. Mol. Spectrosc.*, **326**, 5 (2016).

^dP. Maddaloni, P. Malara, E. De Tommasi, M. De Rosa, I. Ricciardi, G. Gagliardi, F. Tamassia, G. Di Lonardo, P. De Natale, *J. Chem. Phys.*, **133**, 154317 (2010).

Comparison with experimental data for D₂

Selected transition energies, $(v'', J'') \leftarrow (v', J')$

Contrib.	$(0, 2) \leftarrow (1, 2)$	$(0, 2) \leftarrow (2, 4)$
$E^{(2)}$	2 987.291 387 53(1)	6 241.120 920 12(1)
$E^{(4)}$	0.017 498	0.040 174
$E^{(5)}$	- 0.015 33(8)	- 0.033 167(18)
$E^{(6)}$	- 0.000 137	- 0.000 297(1)
$E^{(7)}$	0.000 007(2)	0.000 016(4)
E_{FS}	- 0.000 015	- 0.000 031
E	2 987.293 41(9)	6 241.127 615(19)
Exper.	2 987.293 52(15) ^a	6 241.127 647(11) ^a 6 241.127 655(13) ^b 6 241.127 637(17) ^c
Diff.	0.000 11(17)	0.000 040(23)

^aM. Niu, E. Salumbides, G. Dickenson, K. Eikema, W. Ubachs, J. Mol. Spectrosc., **300**, 44 (2014).

^bP. Wcisło, et al., Journal of Quantitative Spectroscopy and Radiative Transfer, **213**, 41 (2018).

^cD. Mondelain, S. Kassi, T. Sala, D. Romanini, D. Gatti, A. Campargue, J. Mol. Spectrosc., **326**, 5 (2016).

^dP. Maddaloni, P. Malara, E. De Tommasi, M. De Rosa, I. Ricciardi, G. Gagliardi, F. Tamassia, G. Di Lonardo, P. De Natale, J. Chem. Phys., **133**, 154317 (2010).

Comparison with experimental data for D_2

Selected transition energies, $(v'', J'') \leftarrow (v', J')$

Contrib.	$(0, 2) \leftarrow (1, 2)$	$(0, 2) \leftarrow (2, 4)$	$(0, 1) \leftarrow (1, 3)$
$E^{(2)}$	2 987.291 387 53(1)	6 241.120 920 12(1)	3 278.517 313 65(1)
$E^{(4)}$	0.017 498	0.040 174	0.023 746(1)
$E^{(5)}$	- 0.015 33(8)	- 0.033 167(18)	- 0.017 82(1)
$E^{(6)}$	- 0.000 137	- 0.000 297(1)	- 0.000 160(1)
$E^{(7)}$	0.000 007(2)	0.000 016(4)	0.000 009(5)
E_{FS}	- 0.000 015	- 0.000 031	- 0.000 017
E	2 987.293 41(9)	6 241.127 615(19)	3 278.523 06(1)
Exper.	2 987.293 52(15) ^a	6 241.127 647(11) ^a 6 241.127 655(13) ^b 6 241.127 637(17) ^c	3 278.522 0(2) ^d
Diff.	0.000 11(17)	0.000 040(23)	0.001 1(2)

^aM. Niu, E. Salumbides, G. Dickenson, K. Eikema, W. Ubachs, J. Mol. Spectrosc., **300**, 44 (2014).

^bP. Wcisło, et al., Journal of Quantitative Spectroscopy and Radiative Transfer, **213**, 41 (2018).

^cD. Mondelain, S. Kassı, T. Sala, D. Romanini, D. Gatti, A. Campargue, J. Mol. Spectrosc., **326**, 5 (2016).

^dP. Maddaloni, P. Malara, E. De Tommasi, M. De Rosa, I. Ricciardi, G. Gagliardi, F. Tamassia, G. Di Lonardo, P. De Natale, J. Chem. Phys., **133**, 154317 (2010).

Conclusion

$$E(\alpha) = m\alpha^2 E^{(2)} + m\alpha^4 E^{(4)} + m\alpha^5 E^{(5)} + m\alpha^6 E^{(6)} + m\alpha^7 E^{(7)} + \dots$$

- **Variational solutions to four-body Schrödinger equation enable 10^{-8} cm^{-1} (10^{-13}) accuracy on nonrelativistic D_0 and transition energy**
- Outlook:

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- Outlook:
 - The DNVM wave function will be applied to evaluate relativistic corrections
 - $m\alpha^5 E^{(5)}$ QED term and the missing $m\alpha^7 E^{(7)}$ correction limit the accuracy of the current theoretical predictions

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