

# CP-violating Bottom Yukawa at NLO

Joachim Brod



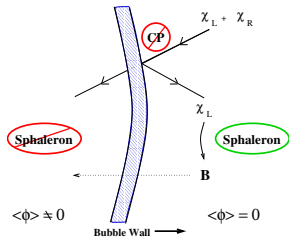
PIKIO Meeting, UKY Lexington  
September 16, 2017

With Emmanuel Stamou – [work in progress](#)

# Motivation – Electroweak Baryogenesis

- Baryogenesis fails within the SM
  - Need strong first-order phase transition
  - Need more CP violation
- A minimal setup for electroweak baryogenesis:

[Huber, Pospelov, Ritz, hep-ph/0610003]



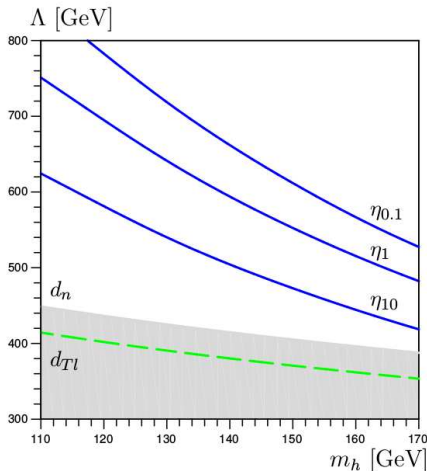
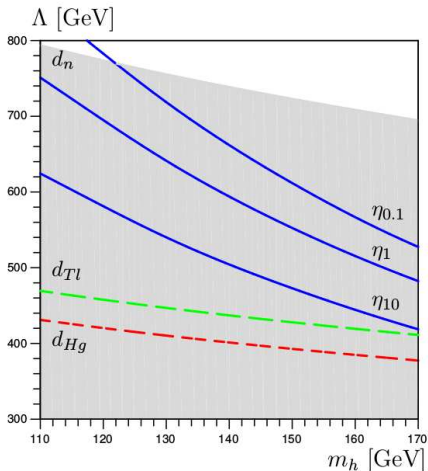
[Image credit: Morrissey et al., 1206.2942]

$$\mathcal{L} = \frac{1}{\Lambda^2} (H^\dagger H)^3 + \frac{Z_t}{\Lambda^2} (H^\dagger H) \bar{Q}_3 H^c t_R$$

- $\Lambda \sim 500 - 800 \text{ GeV}$  gives correct baryon-to-photon ratio  $\eta_b$
- In principle, there are more operators

# Motivation – EDM constraints on baryogenesis

[Huber, Pospelov, Ritz, hep-ph/0610003]



# Modified Bottom Yukawa

- In the SM, Yukawa coupling to fermion  $b$  is

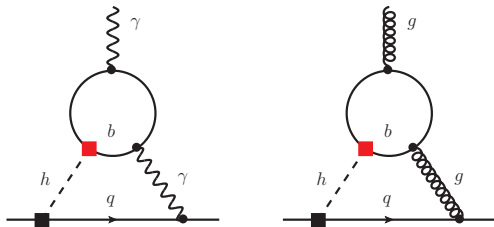
$$\mathcal{L}_Y = -\frac{y_b}{\sqrt{2}}(\bar{b}b)h$$

- We will look at modification

$$\mathcal{L}'_Y = -\frac{y_b}{\sqrt{2}}(\kappa_b \bar{b}b + i\tilde{\kappa}_b \bar{b}\gamma_5 b)h$$

- In the SM,  $\kappa_b = 1$  and  $\tilde{\kappa}_b = 0$
- New contributions will
  - modify Higgs production cross section and decay rates
  - generate EDMs via higher loop effects

# “Naive” Barr-Zee

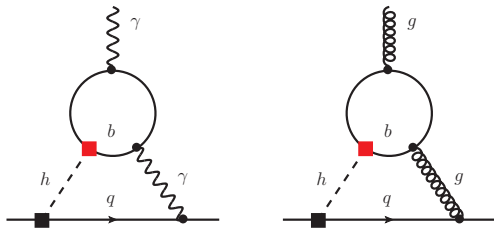


- Generate two operators:
  - EDM ( $d_q$ ):  $\bar{q}\sigma^{\mu\nu}\gamma_5 q F_{\mu\nu}$
  - CEDM ( $\tilde{d}_q$ ):  $\bar{q}\sigma^{\mu\nu} T^a \gamma_5 q G_{\mu\nu}^a$

$$d_q(\mu_W) \simeq -4eQ_q N_c Q_b^2 \frac{\alpha}{(4\pi)^3} \sqrt{2} G_F m_q \kappa_q \tilde{\kappa}_b \frac{m_b^2}{M_h^2} \left( \log^2 \frac{m_b^2}{M_h^2} + \frac{\pi^2}{3} \right),$$

$$\tilde{d}_q(\mu_W) \simeq -2 \frac{\alpha_s}{(4\pi)^3} \sqrt{2} G_F m_q \kappa_q \tilde{\kappa}_b \frac{m_b^2}{M_h^2} \left( \log^2 \frac{m_b^2}{M_h^2} + \frac{\pi^2}{3} \right),$$

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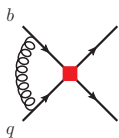
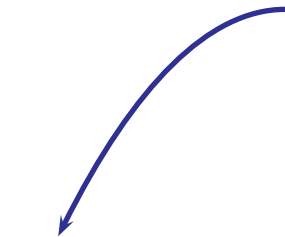
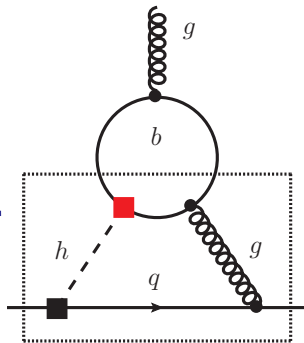
$$\tilde{d}_q(\mu_W) \simeq -2 \frac{\alpha_s}{(4\pi)^3} \sqrt{2} G_F m_q \kappa_q \tilde{\kappa}_b \frac{m_b^2}{M_h^2} \left( \log^2 \frac{m_b^2}{M_h^2} + \frac{\pi^2}{3} \right),$$

- $\alpha_s(M_h)^2 \sim 0.01?$        $\alpha_s(m_b)^2 \sim 0.045?$        $[\alpha_s(2(\text{GeV}))^2 \sim 0.07?]$

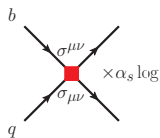
# RG analysis of the $b$ -quark contribution to EDMs

- Factor  $\approx 5$  scale uncertainty in CEDM Wilson coefficient
- Related to different scales in problem:  $\alpha_s \log(M_h/m_b) \sim 1$  is large!
- Use techniques of effective theory and the renormalization group:
  - Sum  $\alpha_s^n \log^n(M_h/m_b)$  to all orders (“LL”)  
[Brod, Haisch, Zupan, 1310.1385]

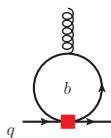
# RG in a nutshell



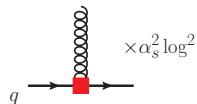
$\Rightarrow$



$\Rightarrow$

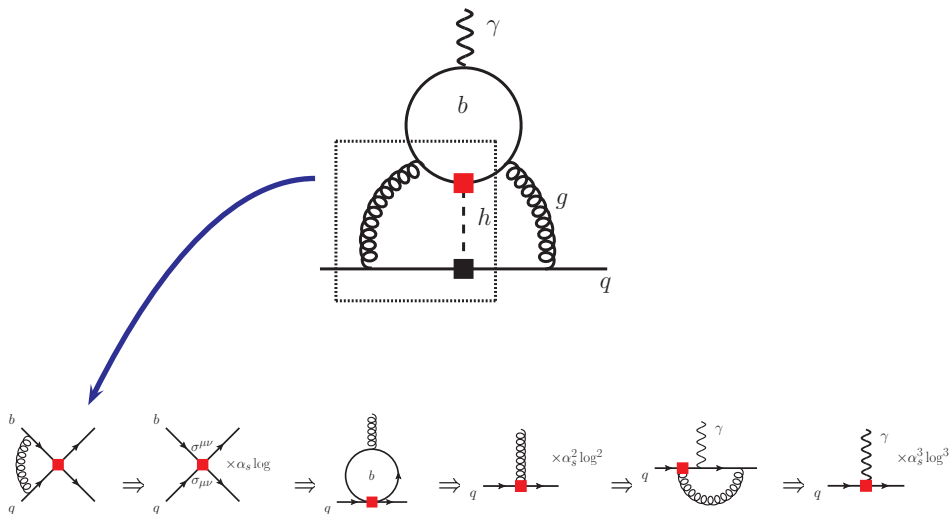


$\Rightarrow$



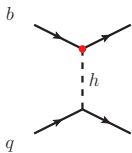


# More RG in a nutshell

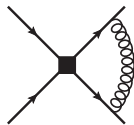


- This contribution dominates over two-loop Barr-Zee by a factor of  $\approx 10!$

# Leading-logarithmic results

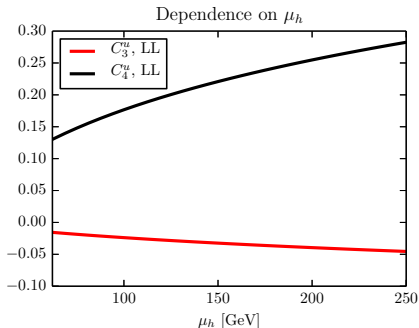


- Tree-level matching



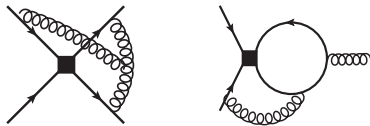
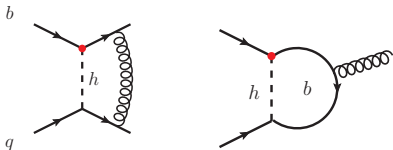
- One-loop running

[Hisano et al., 1205.2212, Misiak et al., hep-ph/9409454]



- LL RG sums  $\alpha_s^n \log^n$  to all orders
- Still factor 2 uncertainty after LL resummation
- $\Rightarrow$  need NLO analysis

# NLO calculation



- One-loop matching:

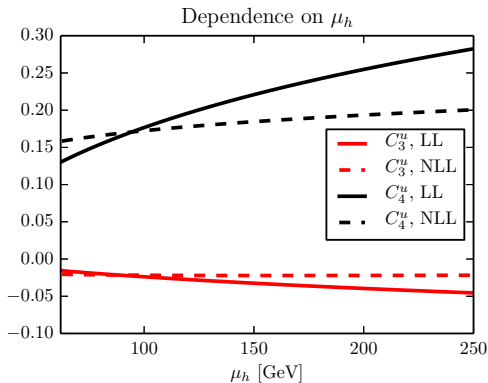
- Cancels linear  $\log \mu$  dependence in LL running
- Finite part is scheme dependent

- Two-loop running:

- Sums  $\alpha_s^{n+1} \log^n$  to all orders
- Cancels scheme dependence of one-loop initial conditions
- Partial results available  
[Misiak et al., hep-ph/9409454, Buras et al. hep-ph/0005183, Degrossi et al. hep-ph/0510137]

# Next-to-leading-logarithmic results

- PARTIAL RESULT excluding full NLO running (scheme dependent!)



# Contribution to the neutron EDM

- Hadronic matrix elements:

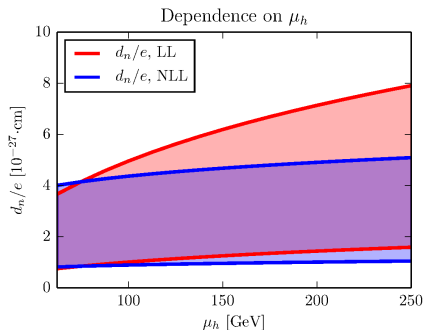
- qEDM  $\rightarrow$  lattice:  $g_T^u = -0.233(28)$ ,  $g_T^d = 0.774(66)$  ( $\overline{\text{MS}}$  @ 2 GeV)

[Battacharya et al., 1506.04196, 1506.06411]

- qCEDM: ChPT and NDA

[E.g. Pospelov & Ritz, hep-ph/0504231]

- Exp. bound:  $|d_n/e| < 2.9 \times 10^{-26}$  cm (90% CL) [Baker et al., hep-ex/0602020]



$$\frac{d_n}{e} = (1.1 \pm 0.55)(\tilde{d}_d + 0.5\tilde{d}_u) - \left( \frac{g_T^u}{e} d_u + \frac{g_T^d}{e} d_d \right)$$

# Summary

- EDMs yield strong constraints on new sources of CP violation
- Large scale dependence hints at large NLO corrections
- Need to complete NLO calculation
- Study theory uncertainty in more detail
- Future: Systematic analysis of all Yukawas

# Appendix

# EDMs as probes of CP violation

- $T^2$  acts on any one-particle state as

$$T^2 \psi_{p,\sigma} = (-1)^{2j} \psi_{p,\sigma}$$

- For odd number of non-interacting spin-1/2 particles, have

$$T^2 \psi = -\psi$$

- Remains true if interactions are T-invariant (e.g. static electric fields)
- Since Hamiltonian  $H$  commutes with  $T$ , both  $\psi$  and  $T\psi$  are eigenstates
- They cannot be the same state, since  $\psi = \zeta T\psi$  implies

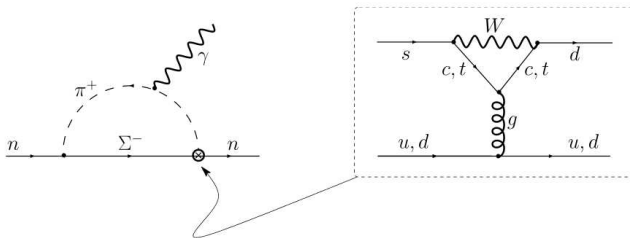
$$T^2 \psi = T(\zeta T\psi) = \zeta^* T\psi = |\zeta|^2 \psi = \psi$$

- An EDM would entirely remove this degeneracy in a static electric field
- Thus, EDMs are forbidden by  $T$  ( $CP$ ) invariance



# Sources of CP violation

- QCD is CP invariant...
  - ... apart from possible  $\theta$  term  $\propto \epsilon_{\mu\nu\alpha\beta} G^{\mu\nu} G^{\alpha\beta}$
  - Neglect for the purpose of this talk
- Microscopic origin of CP violation:
  - Weak interactions
  - New Physics
- E.g. neutron EDM: SM contribution is tiny,  $d_n^{SM} \sim 10^{-32} e \text{ cm}$   
[Khriplovich & Zhitnitsky, PLB 109 (1982) 490]



# EDM experiments, bounds

- Measure different EDMs
  - Elementary: neutron, proton, deuteron
  - Atomic: mercury, radium, xenon
  - Molecular: ThO (mainly electron)
- Current bounds and prospects [ $e\text{ cm}$ ]:

	$d_e$	$d_n$	$d_{p,D}$
current	$8.7 \times 10^{-29}$	$2.9 \times 10^{-26}$	–
expected	$5.0 \times 10^{-30}$	$1.0 \times 10^{-28}$	$1.0 \times 10^{-29}$
	$d_{\text{Hg}}$	$d_{\text{Xe}}$	$d_{\text{Ra}}$
current	$2.6 \times 10^{-29}$	$5.5 \times 10^{-27}$	$4.2 \times 10^{-22}$
expected	$1.0 \times 10^{-29}$	$5.0 \times 10^{-29}$	$1.0 \times 10^{-27}$

# Hadronic matrix elements

- At low scales, three types of operators contribute:
  - qEDM:  $\bar{q}\sigma^{\mu\nu}\gamma_5qF_{\mu\nu}$
  - qCEDM:  $\bar{q}\sigma^{\mu\nu}T^a\gamma_5qG_{\mu\nu}^a$
  - Weinberg:  $f^{abc}\epsilon_{\mu\nu\alpha\beta}G_{\alpha\beta}^aG_{\mu\rho}^bG_{\nu}^{c,\rho}$
- Hadronic matrix elements:
  - qEDM  $\rightarrow$  lattice  
[Battacharya et al., 1506.04196, 1506.06411]
  - qCEDM: ChPT and NDA, strangeness  
[E.g. Pospelov & Ritz, hep-ph/0504231]
  - Weinberg: No systematic calculation exists, even sign unknown  
[NDA: Weinberg PRL 63 (1989) 2333, Sum rules: Demir et al. hep-ph/0208257]