

EXOTICA IN HADRON SPECTROSCOPY

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4TH PIKIO MEETING

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Ψ

Joint Physics Analysis Center

JPAC is a collaboration between theorists, phenomenologists, and experimentalists to provide phenomenological and data analysis tools for hadron physics

~ 20 active members

> 40 Research Papers (Phys.Rev., Phys.Lett, Eur.J. Phys.) O(10)
ongoing analyses

Regular lecture series on relativistic reaction theory

Summer Workshop on Reaction Theory (2015 & 2017)

<http://www.indiana.edu/~ssrt/>

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National Science Foundation

This project is supported by NSF

$\gamma p \rightarrow \eta p$

We present the model published in [Nys16].
The differential cross section for $\gamma p \rightarrow \eta p$ is computed with Regge amplitudes in the domain $E_\gamma \geq 4 \text{ GeV}$ and $0 \leq -t \leq 1$ (in GeV^2).
We use the CGLN invariant amplitudes A_i defined in [Chew57a].
See the section Formalism for the definition of the variables.
The model and its context is detailed in [Nys16]. We report here only the main features of the model.

Formalism

The differential cross section is a function of 2 kinematic variables. The first is the beam energy in the laboratory frame E_γ (in GeV) or the total energy squared s (in GeV^2). The second is the cosine of the scattering angle in the rest frame $\cos \theta$ or the

Download the [output file](#)
In the file, the columns are:
 $t (\text{GeV}^2)$, $\cos(\theta)$, $D\sigma/dt$ (micro barn/ GeV^2), $D\sigma/d\Omega$ (micro barn), Σ

Observable: differential cross section
Download the [the plot with \$Ox=t\$](#) , [the plot with \$Ox=\cos\$](#) .

gamma p —> eta p

gamma p —> eta p

QUARKS

1st 2nd 3rd

u up 2.3 M $\frac{2}{3}$ $\frac{1}{2}$	c charm 1.27 G $\frac{2}{3}$ $\frac{1}{2}$	t top 173.1 G $\frac{2}{3}$ $\frac{1}{2}$	Mass: eV/c^2 Charge Spin Name	H higgs 126 G 0 0
d down 4.8 M $-\frac{1}{3}$ $\frac{1}{2}$	s strange 95 M $-\frac{1}{3}$ $\frac{1}{2}$	b bottom 4.2 G $-\frac{1}{3}$ $\frac{1}{2}$	g gluon 0 0 1	strong nuclear force
e electron 0.511 M $-\frac{1}{2}$	μ muon 105.7 M $-\frac{1}{2}$	τ tau 1.78 G $-\frac{1}{2}$	γ photon 0 0 1	electromagnetic force
ν_e e neutrino < 2.2 $\frac{1}{2}$	ν_μ μ neutrino 0.17 M $\frac{1}{2}$	ν_τ τ neutrino < 15.5 M $\frac{1}{2}$	W W boson 80.4 G ± 1 1	Z Z boson 91.2 G 0 1

FERMIONS

GAUGE BOSONS

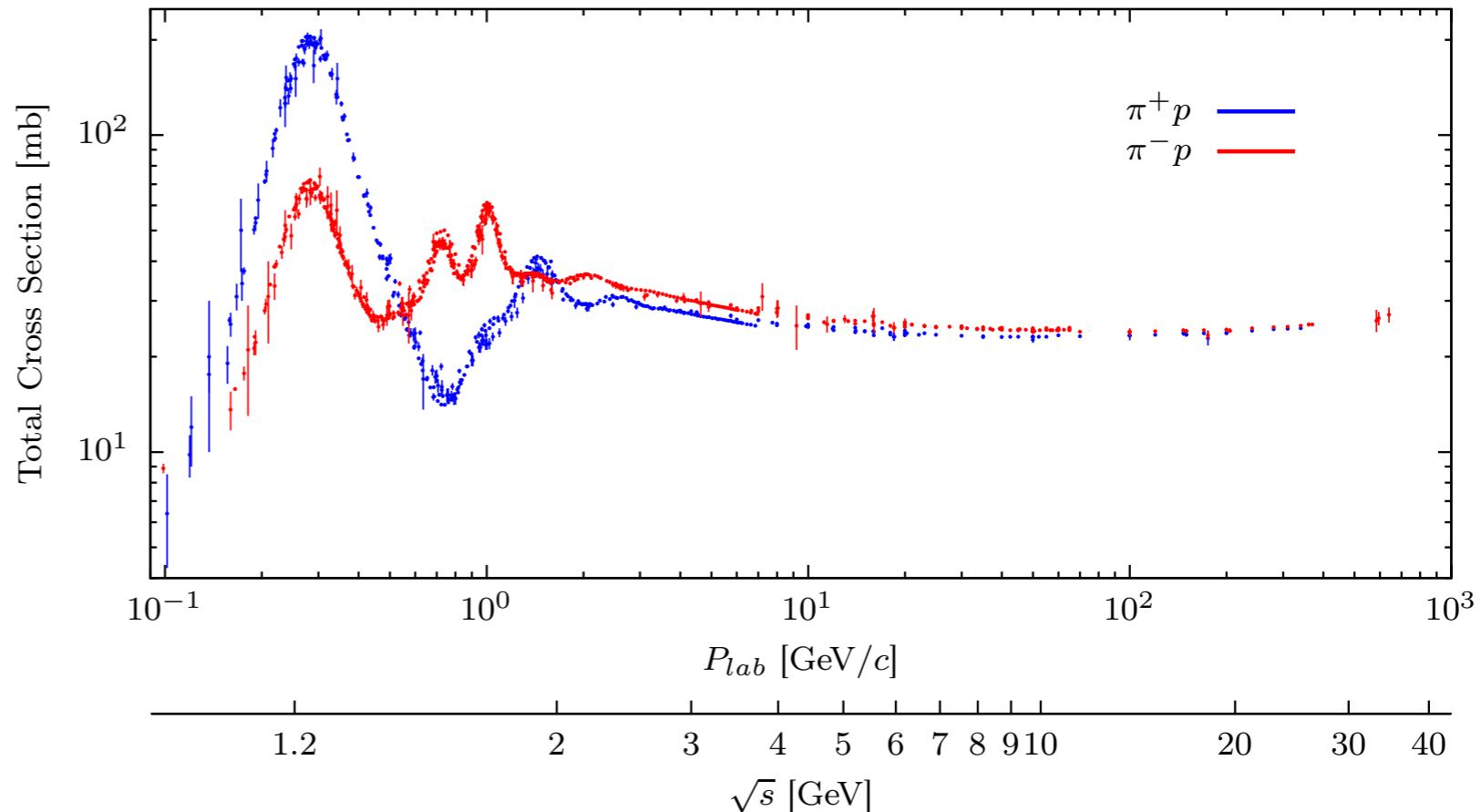
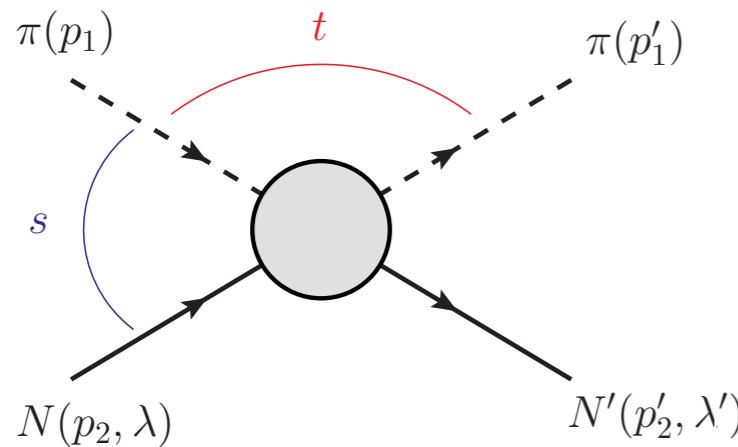
QCD and Hadrons

Theory of quarks and gluons - Quantum Chromodynamics (QCD)

$$\mathcal{L} = \sum_j \frac{1}{2} \bar{\psi}_{q,j} (i \not{D} - m_j) \psi_{q,j} - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a$$

How can such a ‘simple’ theory give rise to the rich structures of hadrons?

$$\pi N \rightarrow \pi N$$



u up	$2.3 M$ $^{2/3}_{1/2}$	c charm	$1.27 G$ $^{2/3}_{1/2}$	t top	$173.1 G$ $^{2/3}_{1/2}$
d down	$4.8 M$ $^{-1/3}_{1/2}$	s strange	$95 M$ $^{-1/3}_{1/2}$	b bottom	$4.2 G$ $^{-1/3}_{1/2}$
					g gluon

Reaction Theory

Use fundamental physics from relativistic reaction theory (unitarity, analyticity, etc.) to constrain hadronic reaction amplitudes

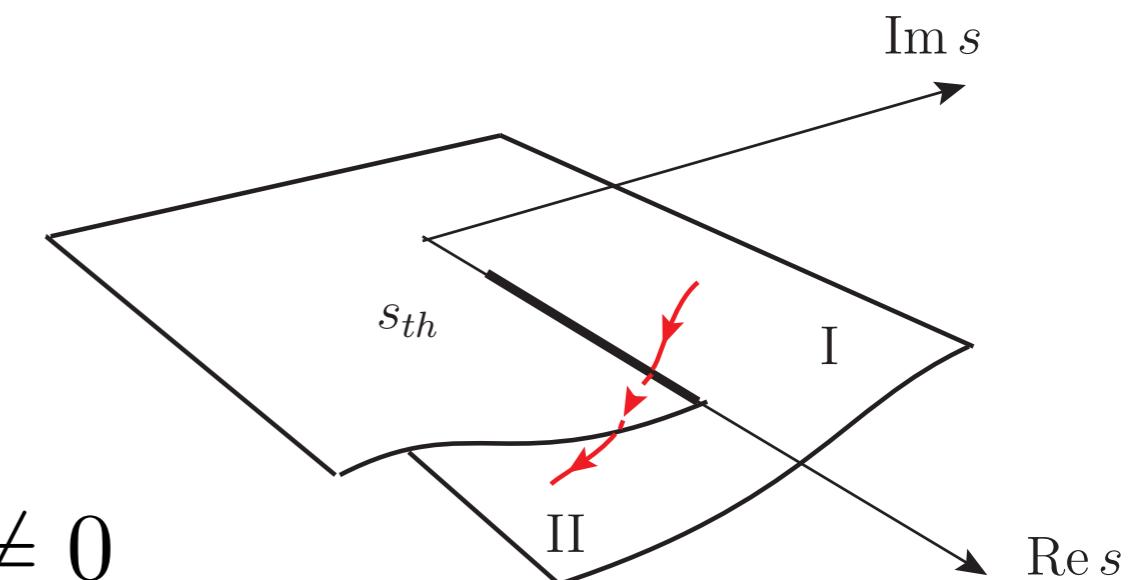
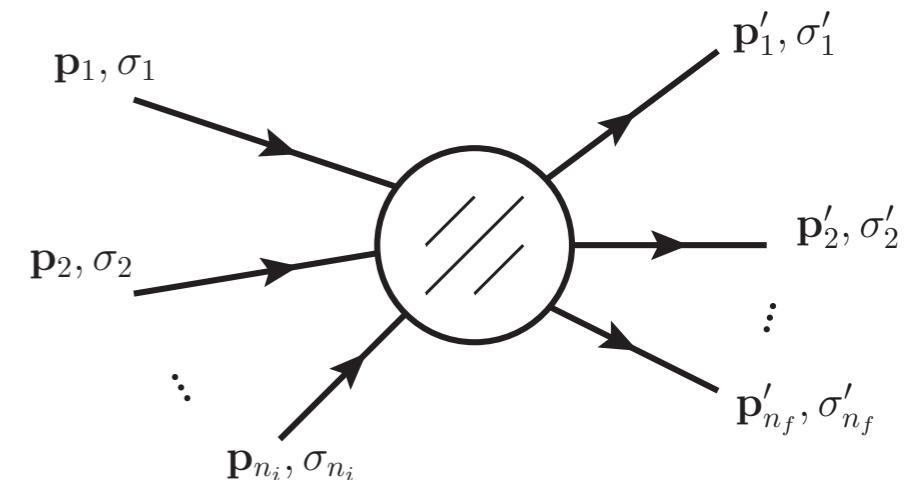
$$\langle \{\mathbf{p}'\sigma'\} | T | \{\mathbf{p}\sigma\} \rangle = (2\pi)^4 \delta^4(P' - P) \mathcal{A}_{\{\sigma', \sigma\}}(\{\mathbf{p}', \mathbf{p}\})$$

Probability conservation S -matrix is unitary operator $\implies S^\dagger S = S S^\dagger = \mathbb{1}$

$$S = \mathbb{1} + iT$$

$$\implies T - T^\dagger = iT^\dagger T$$

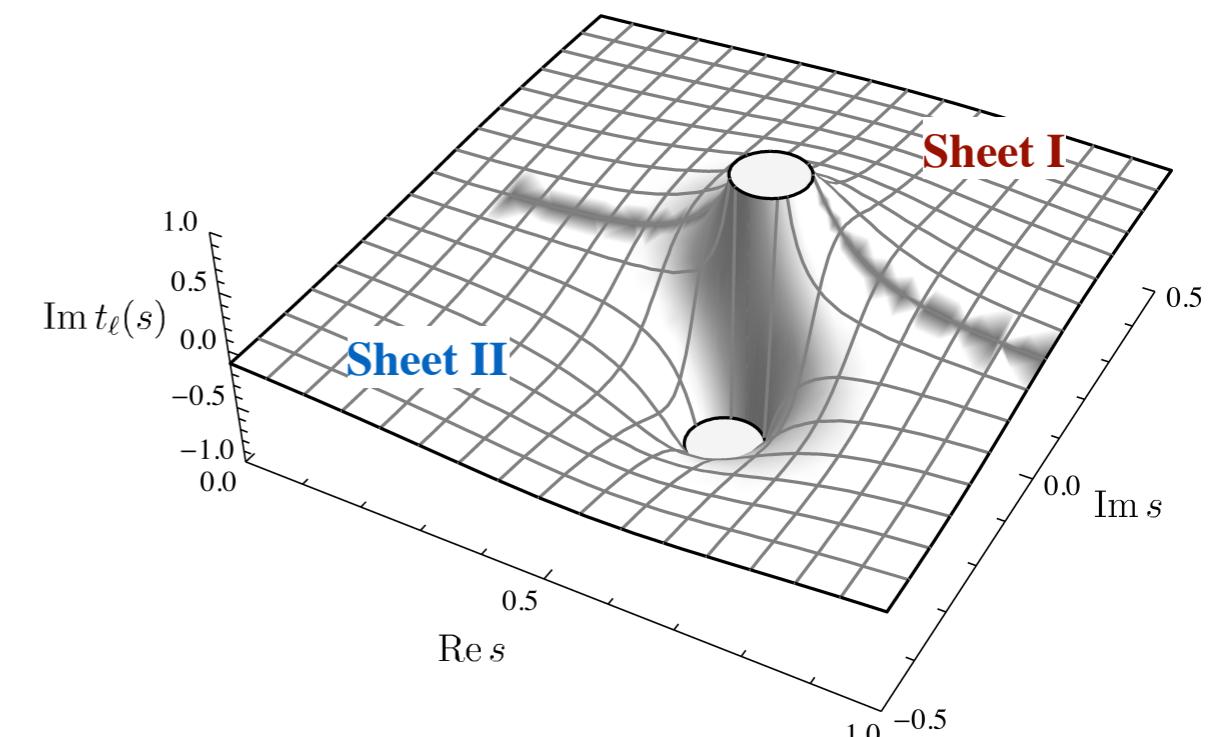
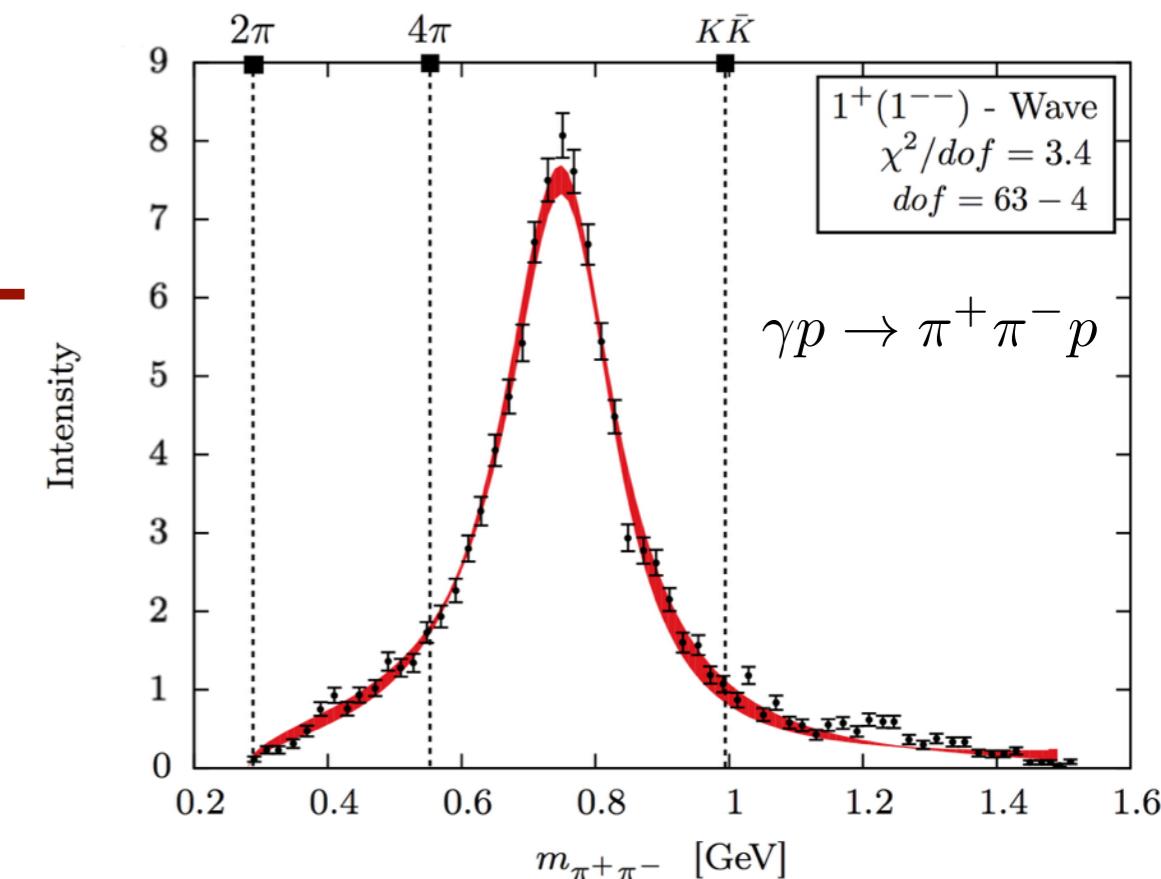
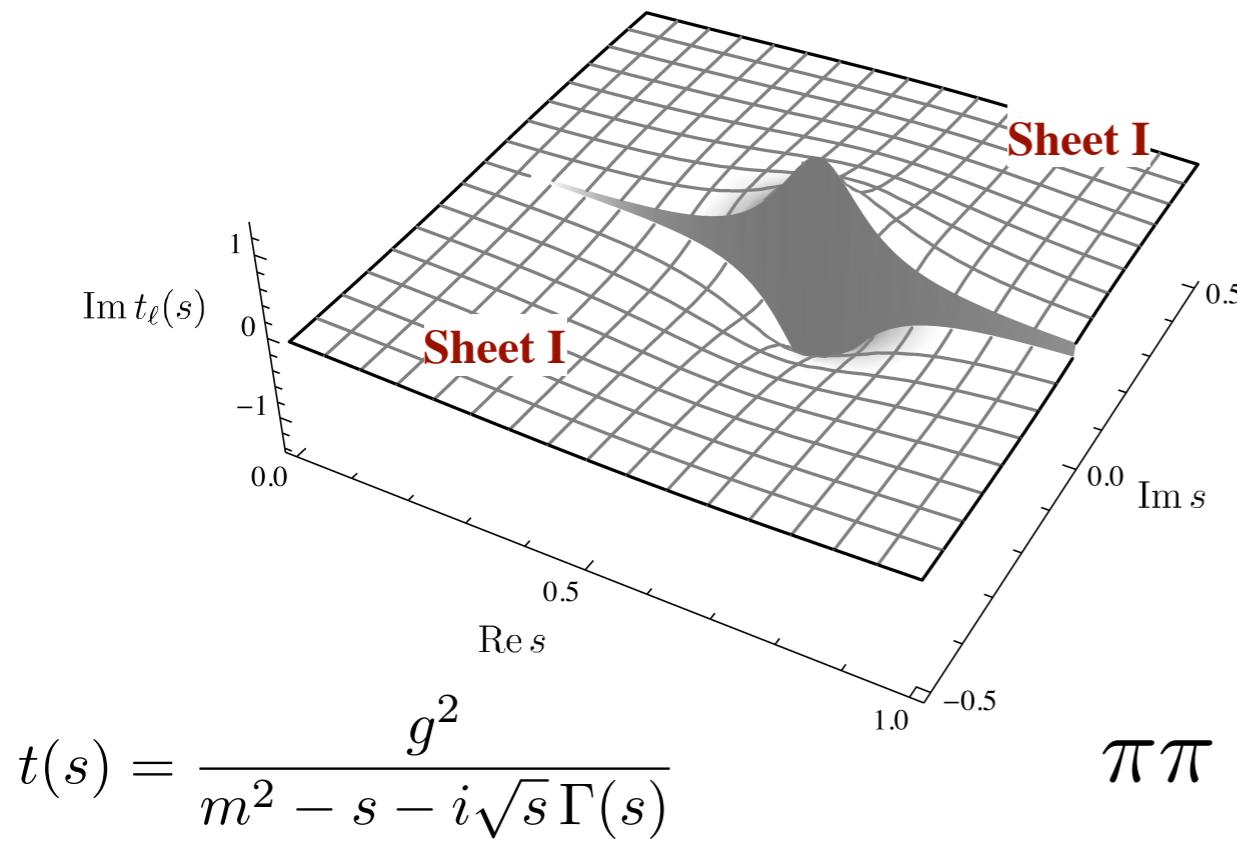
$$\implies \mathcal{A}(s + i\epsilon, t) - \mathcal{A}(s - i\epsilon, t) \neq 0$$



Reaction Theory

Causality implies amplitudes are analytic functions of kinematic variables (energy)

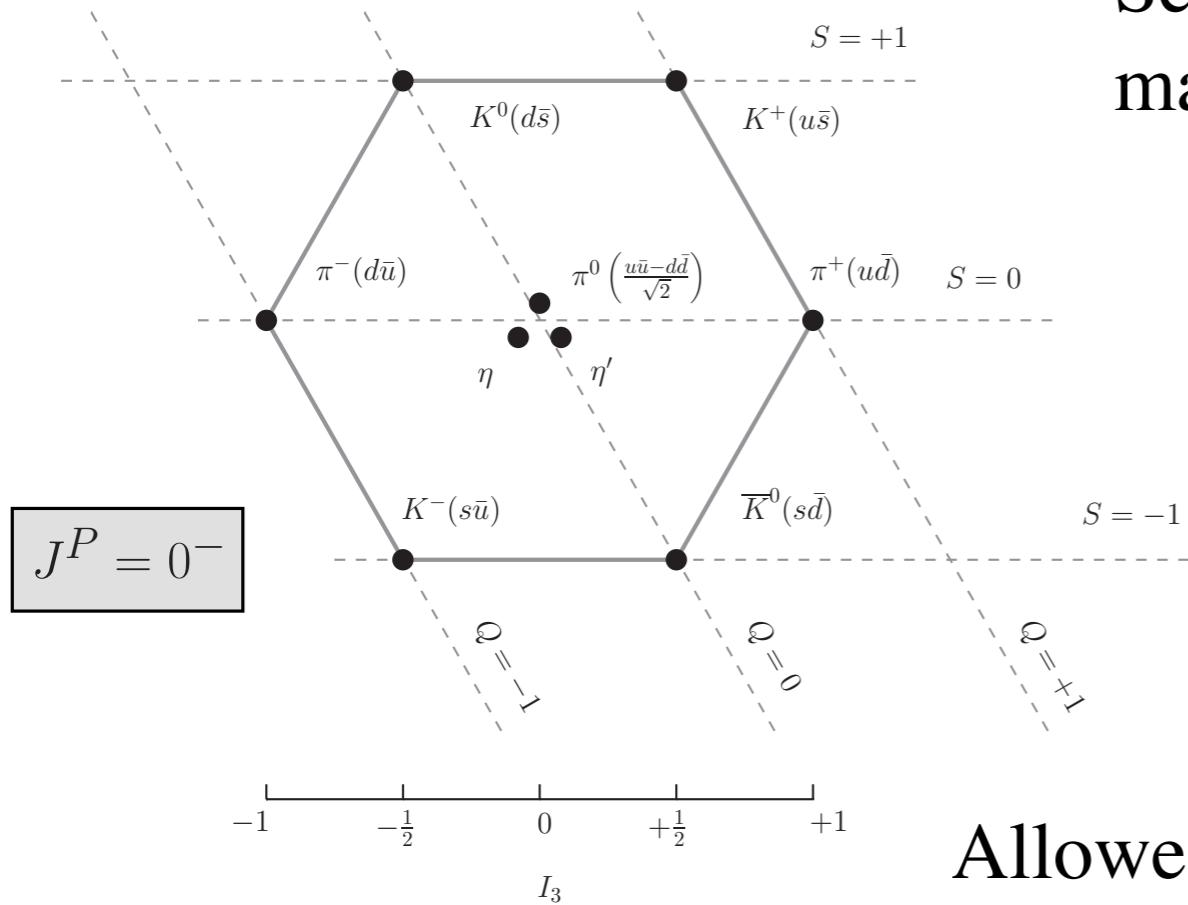
Resonance poles lie underneath unitarity cuts on unphysical Riemann sheets



Spectroscopy

The constituent quark model provides a classification scheme for hadrons

Some quantum numbers are not included



Search for exotic quantum numbers is goal of many experiments (GlueX, COMPASS, etc.)

$$S = 0, 1$$

$$L = 0, 1, 2, \dots$$

$$|L - S| \leq J \leq |L + S|$$

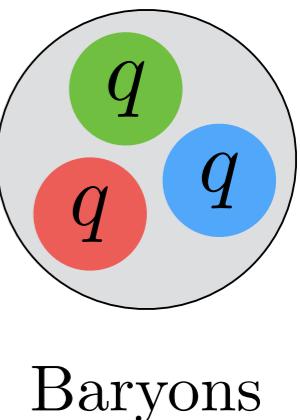
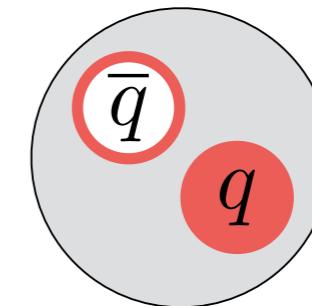
$$P = (-1)^{L+1}$$

$$C = (-1)^{L+S}$$

Allowed : $0^{-+}, 0^{++}, 1^{--}, 1^{+-}, 2^{++}, \dots$

Forbidden : $0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, \dots$

Mesons

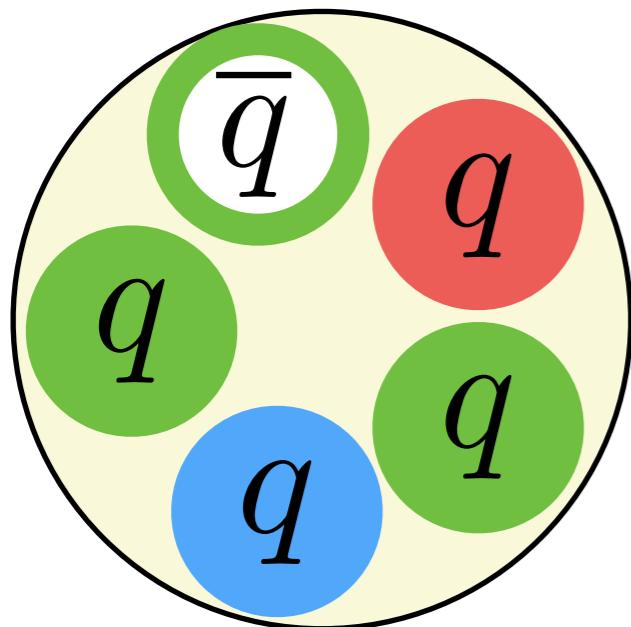


Baryons

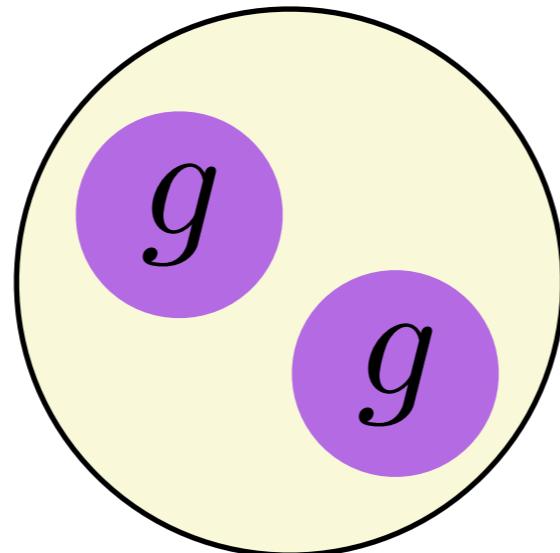
Exotic Hadrons

Not forbidden by QCD

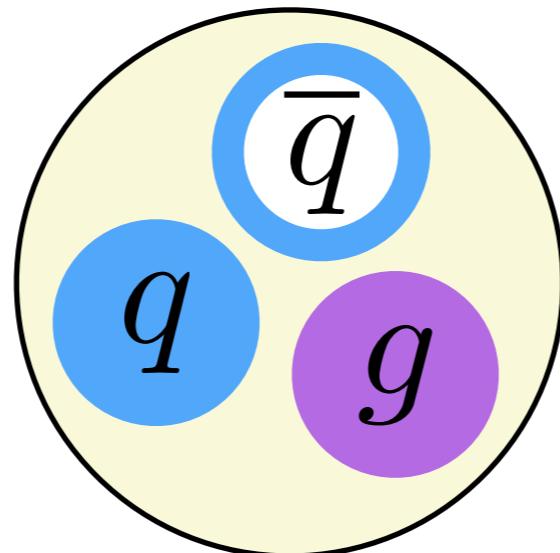
Do these states exist?
Have we seen them?



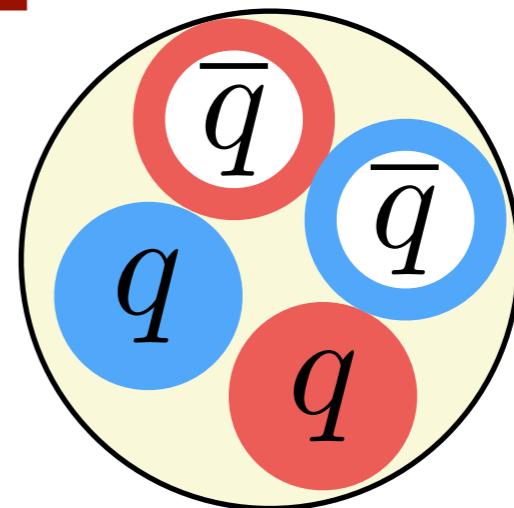
Pentaquarks



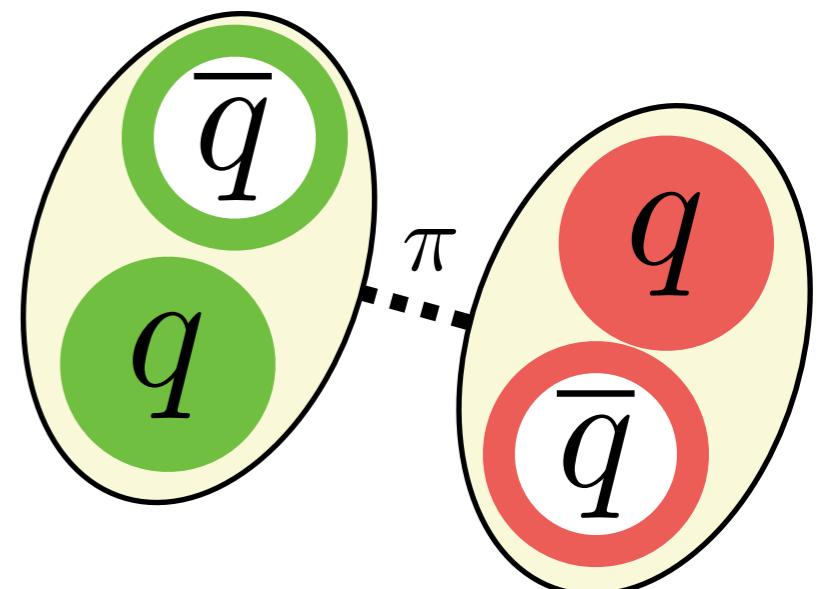
Glueballs



Hybrids



Tetraquarks

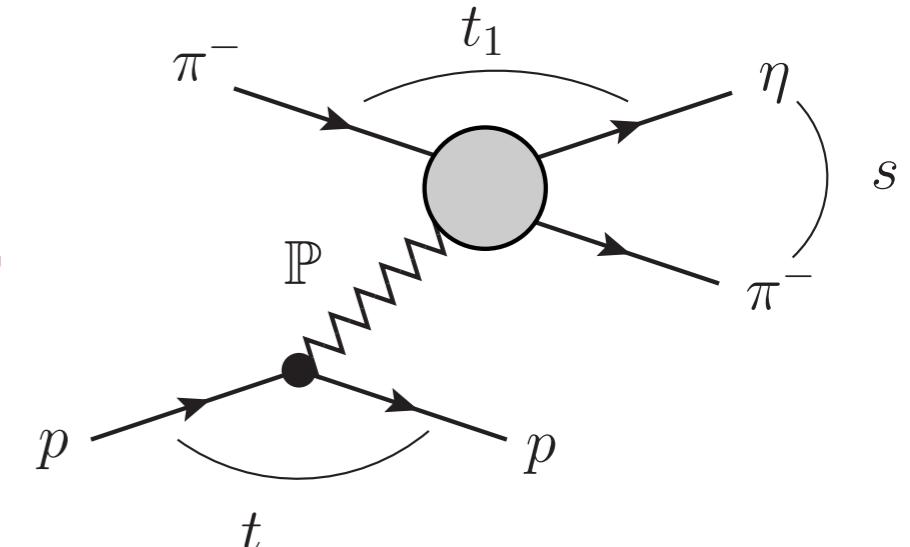
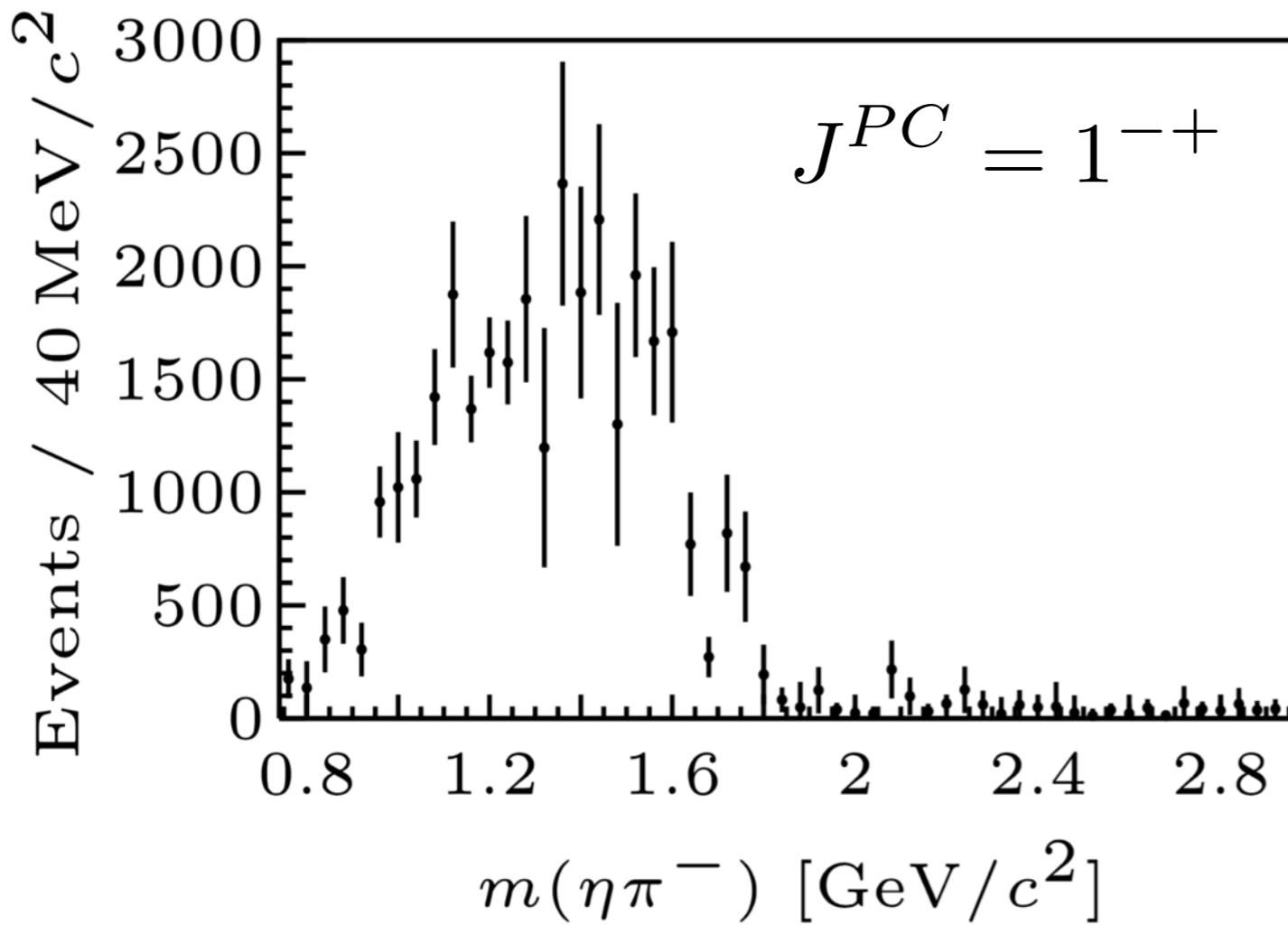


Mesonic-Molecules

Exotics at COMPASS

$$\pi p \rightarrow \eta^{(\prime)} \pi p$$

JPAC working with COMPASS
to analyze this data



JPAC developed relativistic
reaction model to compare
with data

Before analyzing this data,
apply reaction model to
simpler system

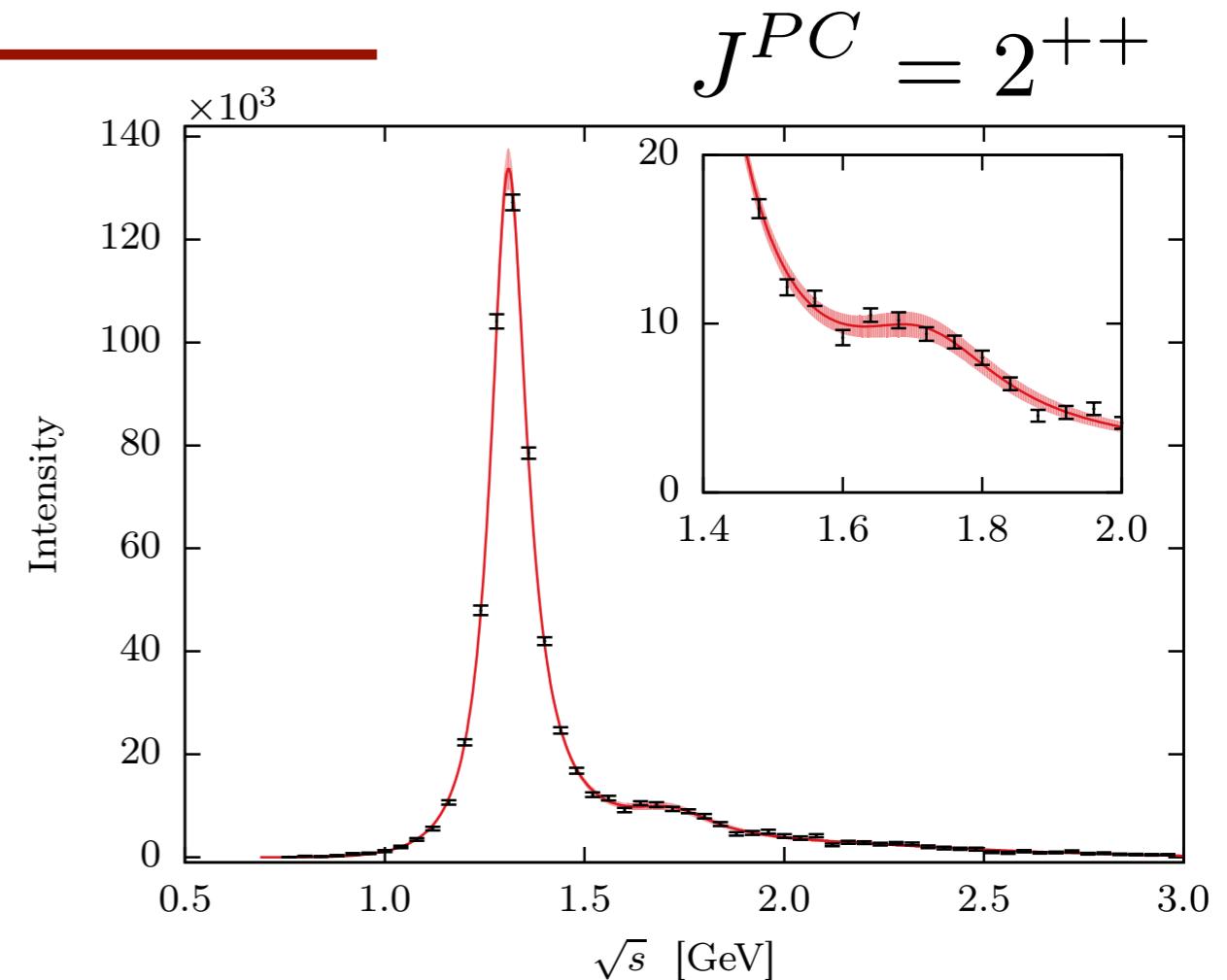
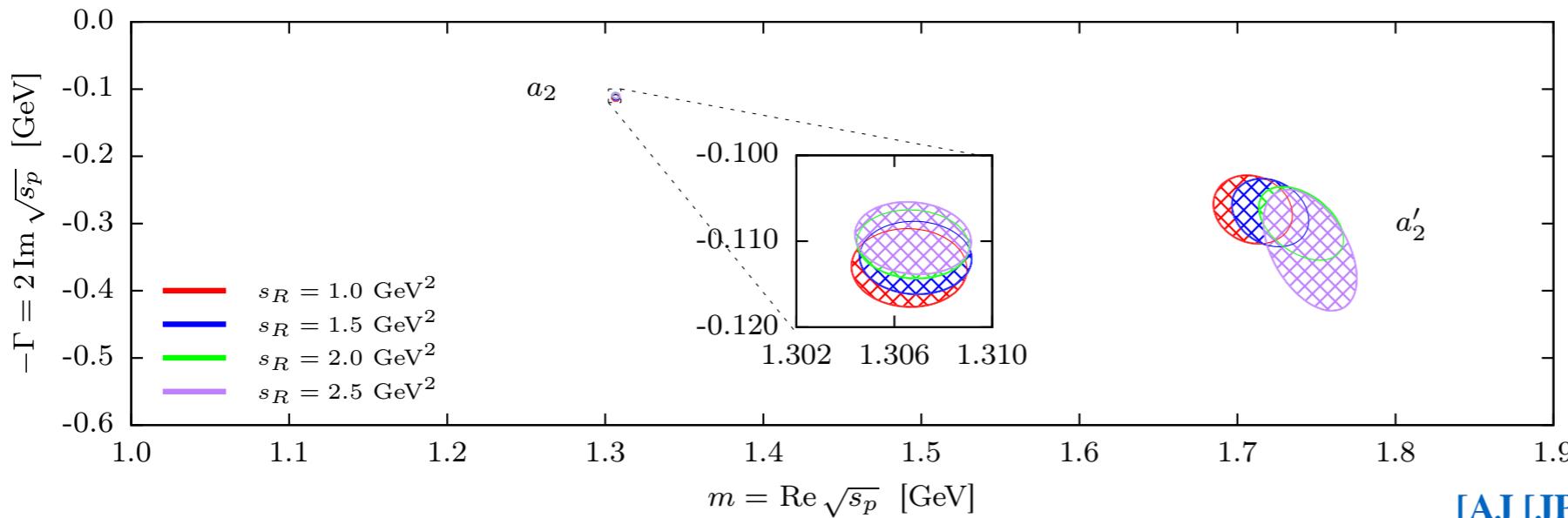
[C. Adolph [COMPASS], Phys. Lett. B 740, 303 (2015)]

Exotics at COMPASS

$$\pi p \rightarrow \eta^{(\prime)} \pi p$$

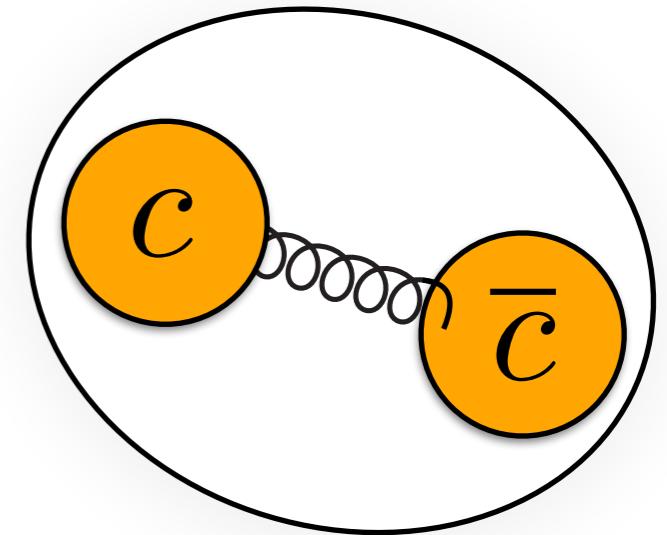
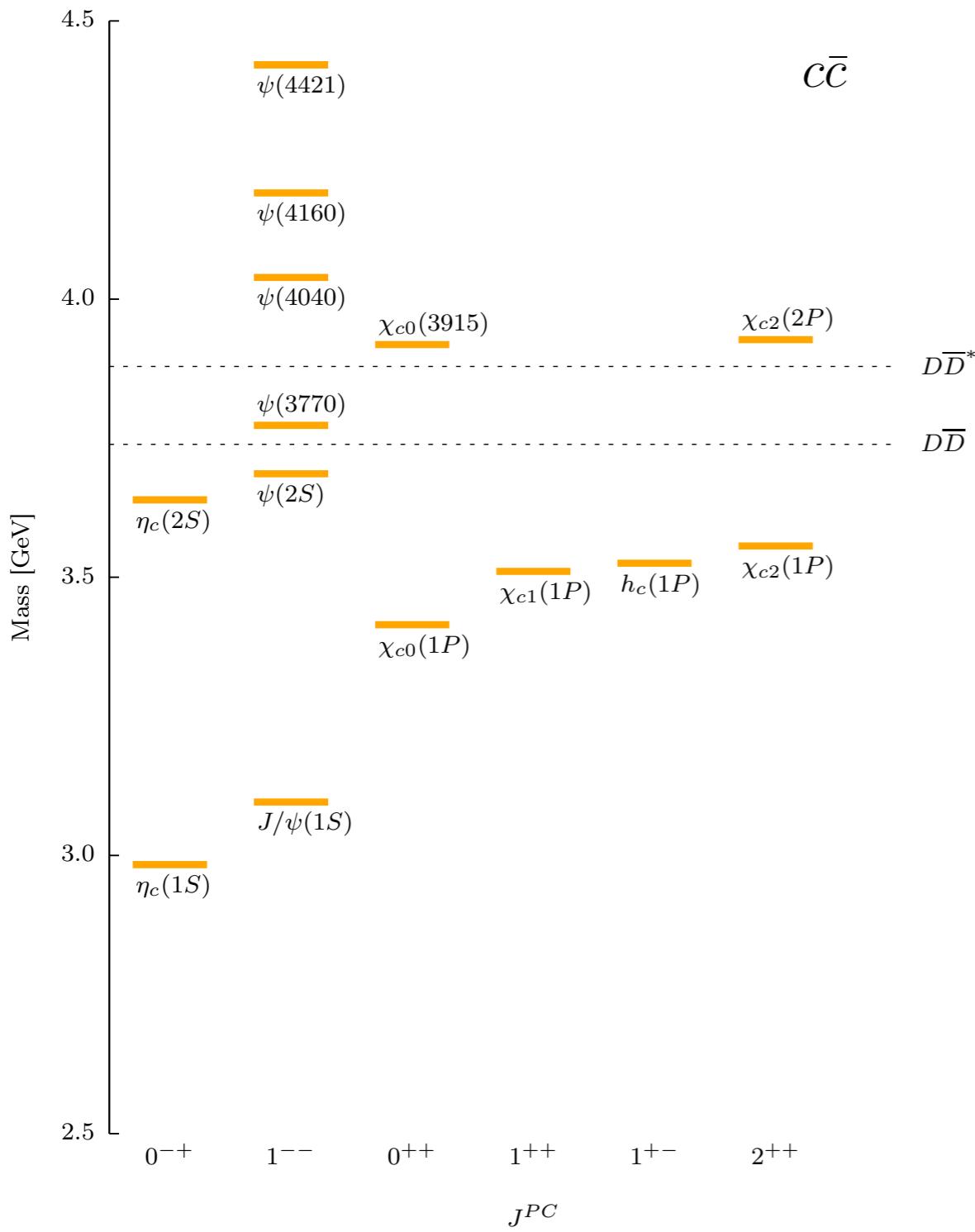
First joint publication between
JPAC and COMPASS out,
laying the ground work of this
analysis

Extract pole positions of a_2, a'_2



[AJ [JPAC & COMPASS], arXiv:1707.02848]

Charmonium



Quark models give qualitative features of heavy quarkonia ($c\bar{c}$, $b\bar{b}$)

All seemed to be generally understood until 2003

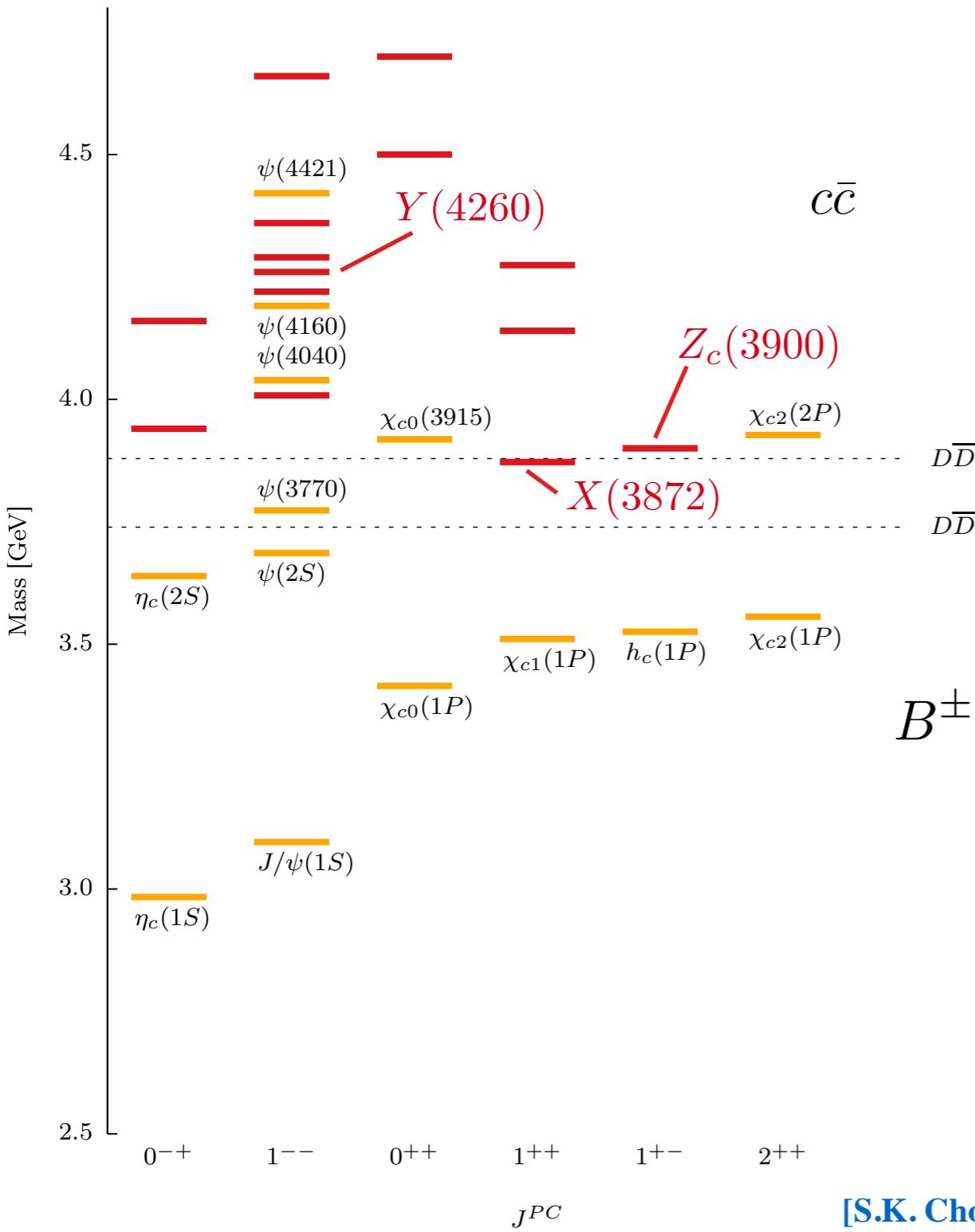
one-gluon exchange

$$V(r) = -\frac{C_F \alpha_s}{r} + \sigma r$$

linear confinement

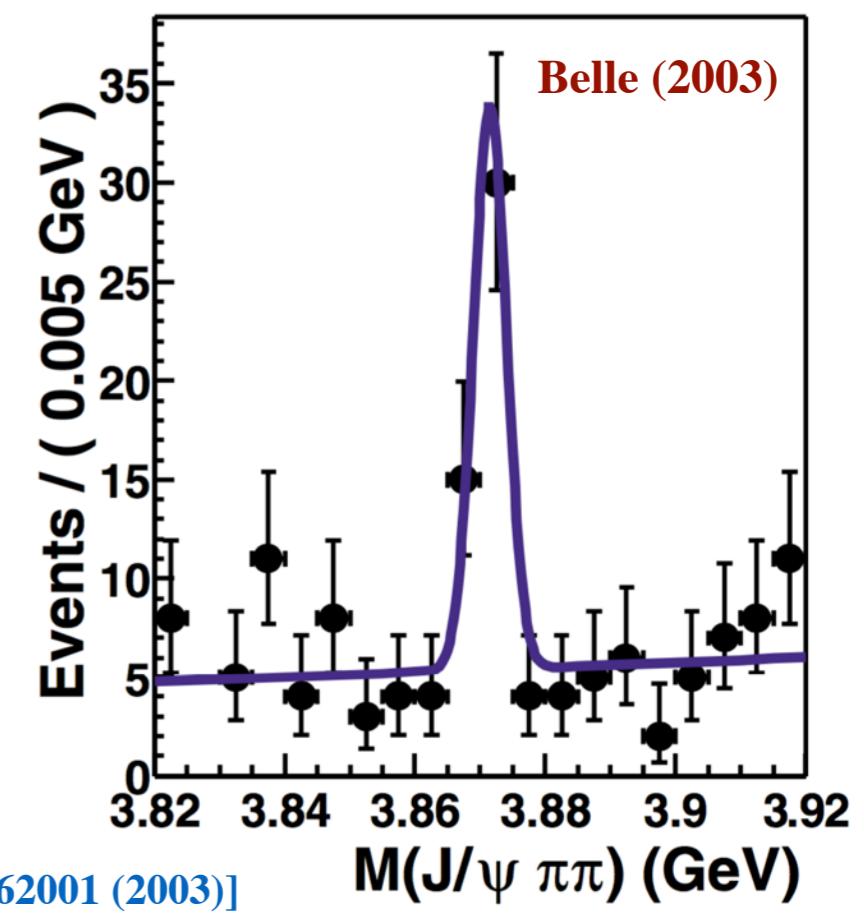
$$\alpha_s(M_Q) \sim 0.3$$

Exotics in Charmonium

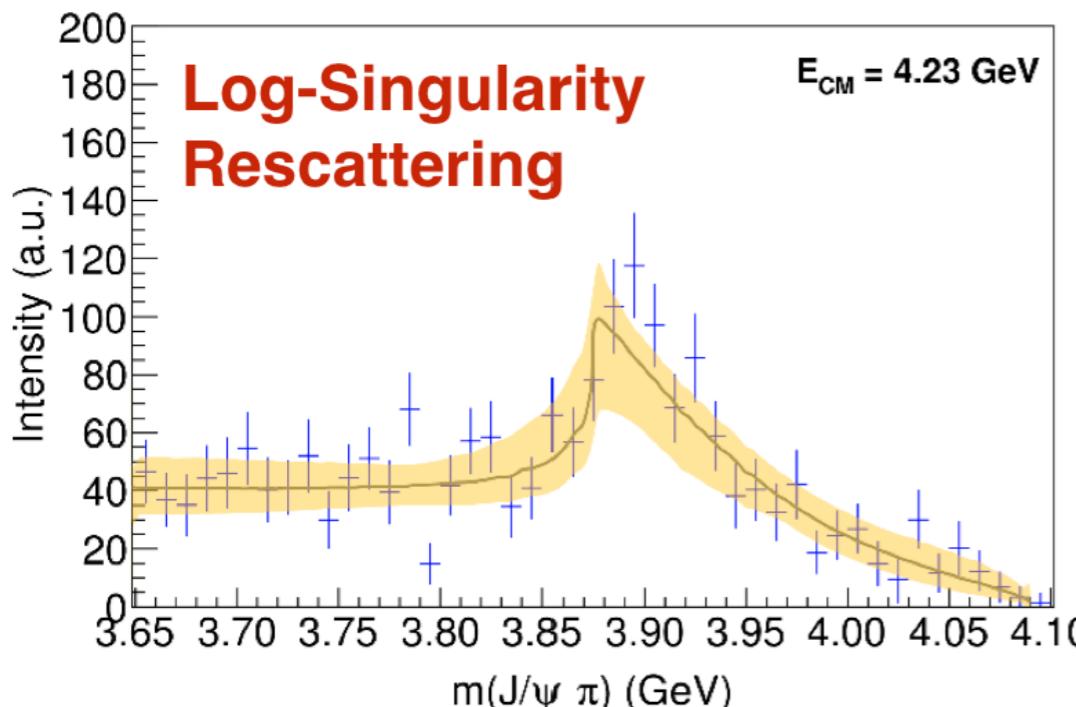
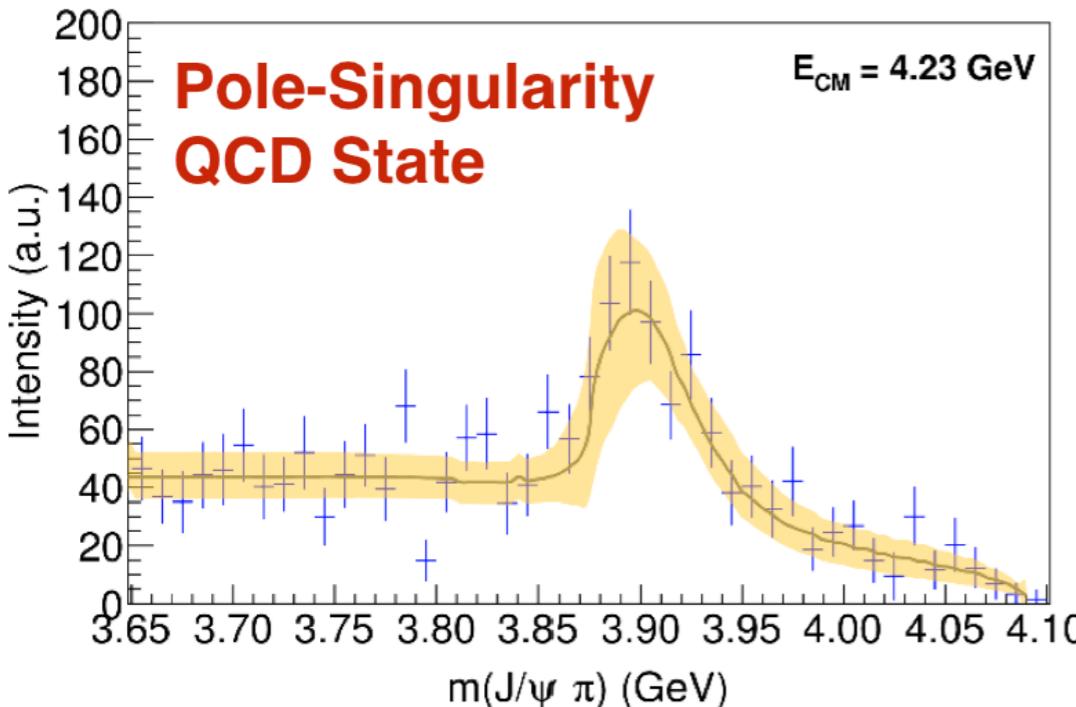


[S.K. Choi [Belle], Phys. Rev. Lett. 91, 262001 (2003)]

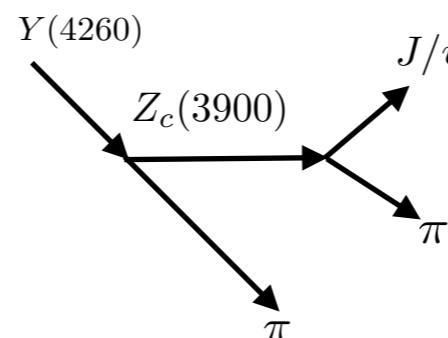
In 2003, the state $X(3872)$ was discovered by Belle
 Began the XYZ era of hadron spectroscopy
 Unconventional quark assignments (tetraquark-like)
 What is the nature of these states?



$Z_c(3900)$



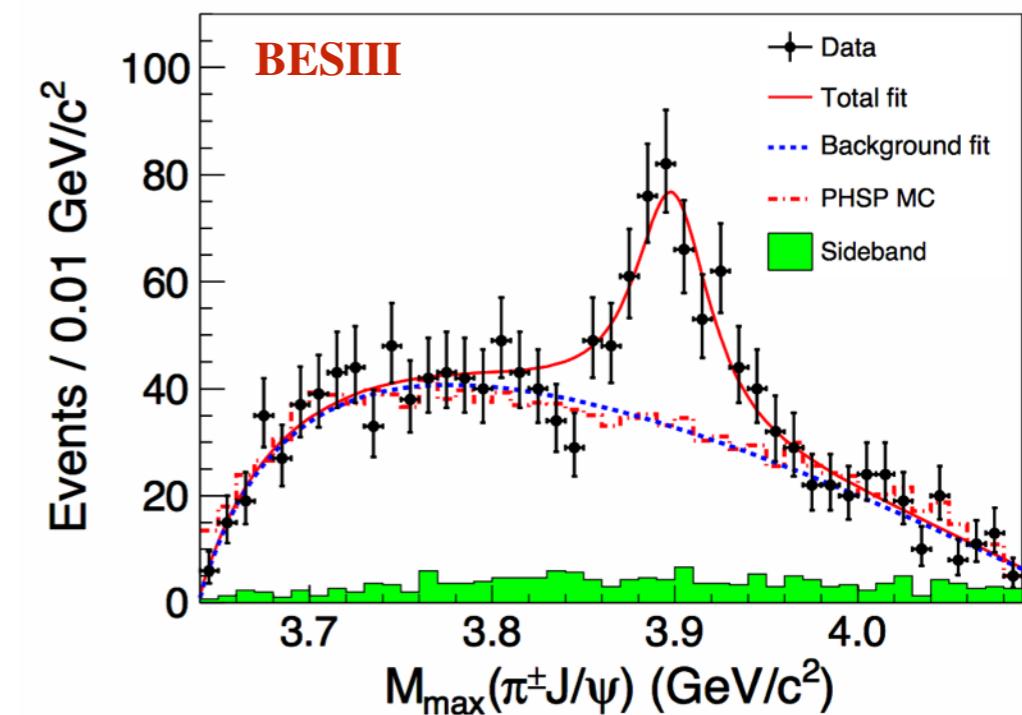
[A.Pilloni, [JPAC] Phys. Lett. B 772, 200 (2017)]



Found in decay of

$$Y(4260) \rightarrow J/\psi \pi \pi$$

Test QCD-pole models and non-QCD pole models



[Liu, BESIII & Belle Collaboration, arXiv:1311.0762v1]

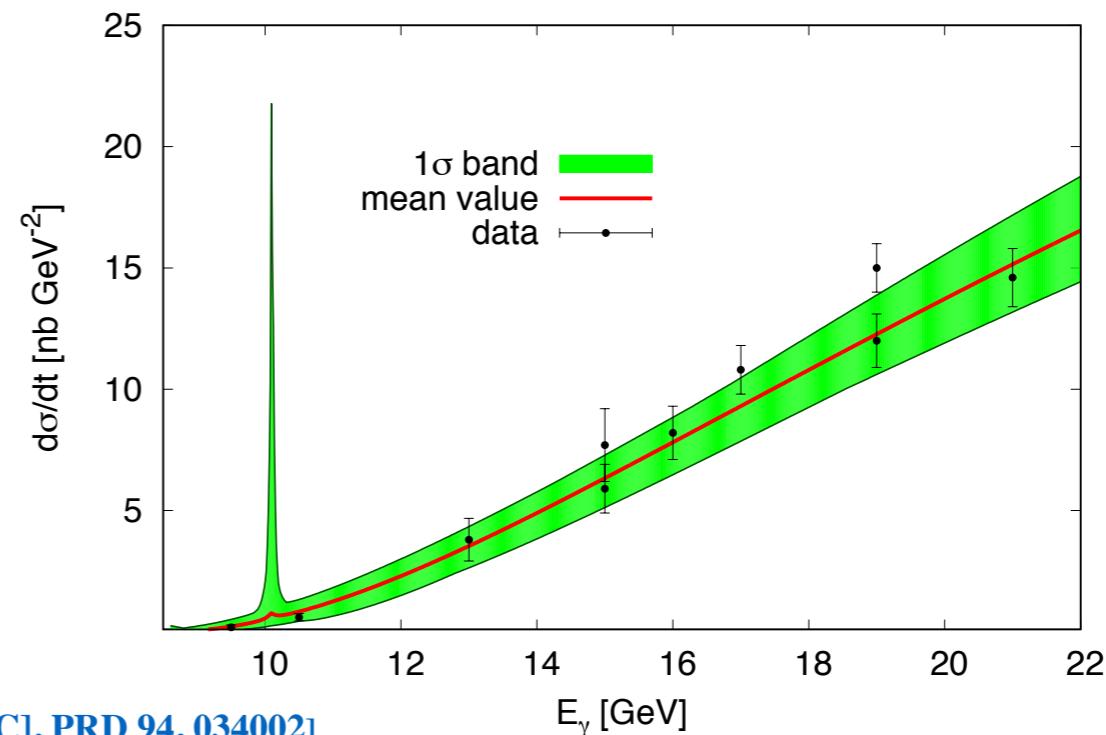
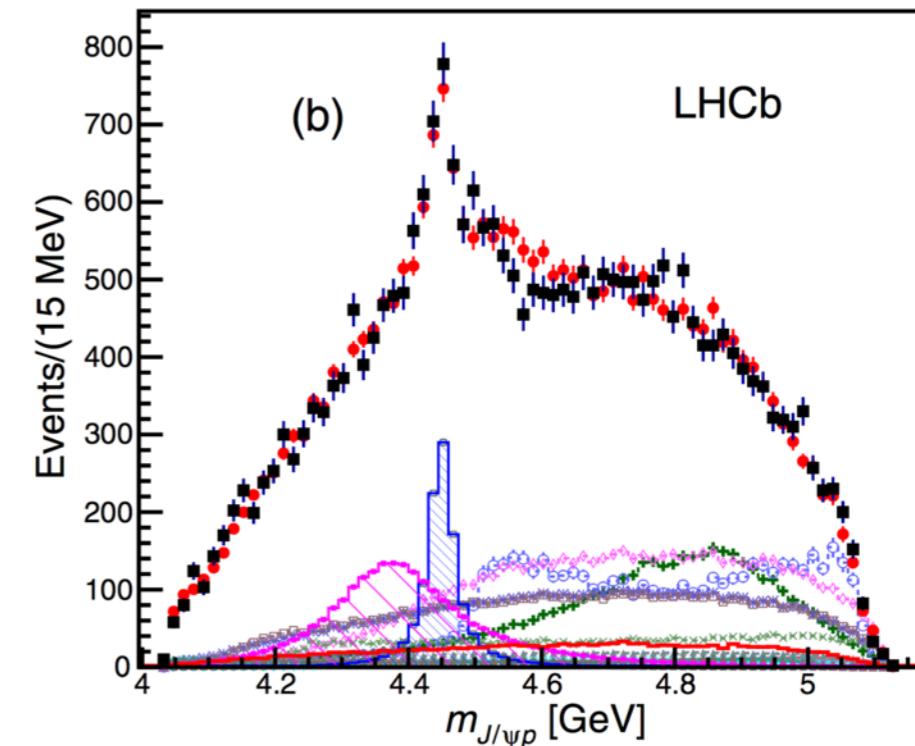
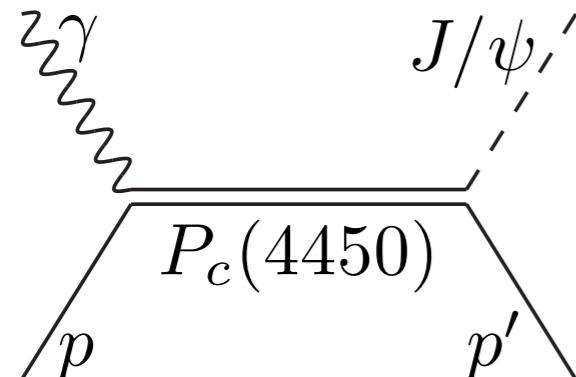
LHCb Pentaquark

[R. AAIJ [LHCb], PRL 115, 072001]

LHCb claim discovery of pentaquarks $P_c(4450)$ and $P_c(4380)$ in $\Lambda_b \rightarrow K J/\psi p$

Proposal to search for $P_c(4450)$ in photoproduction at JLab

Combined JPAC-LHCb analysis on Λ_b decay using JPAC hyperon spectrum



[A.H. BLIN [JPAC], PRD 94, 034002]

Conclusion

Hadron Spectroscopy is an increasingly active field

Many states with unknown origin have been discovered, with more discovered yearly — Exotica breed new life into low-energy QCD

Active work between theorist and experimentalist to construct reliable models to extract physics — JPAC is leader in this effort

General principles of reaction theory can be used to constrain models for experimental analysis — and can be used to discriminate interpretations

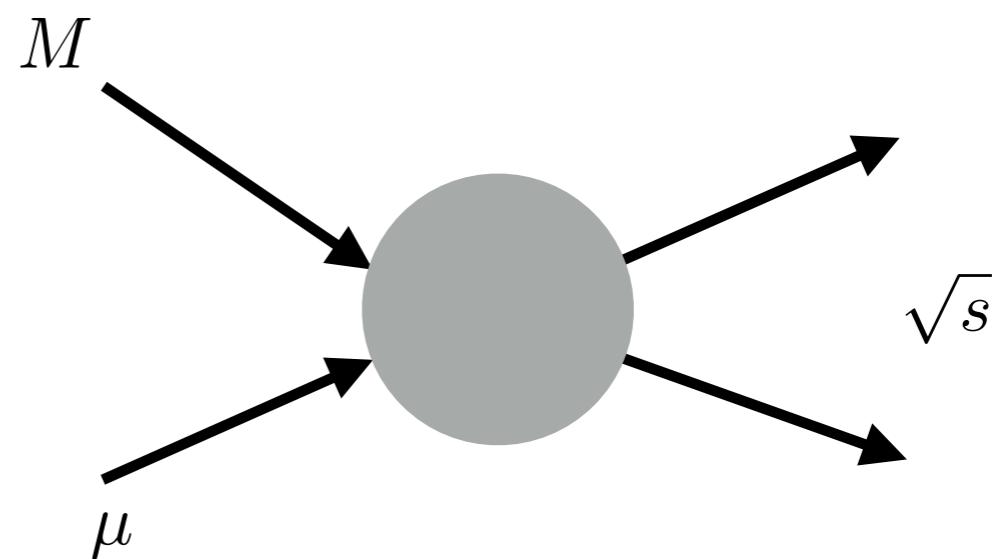
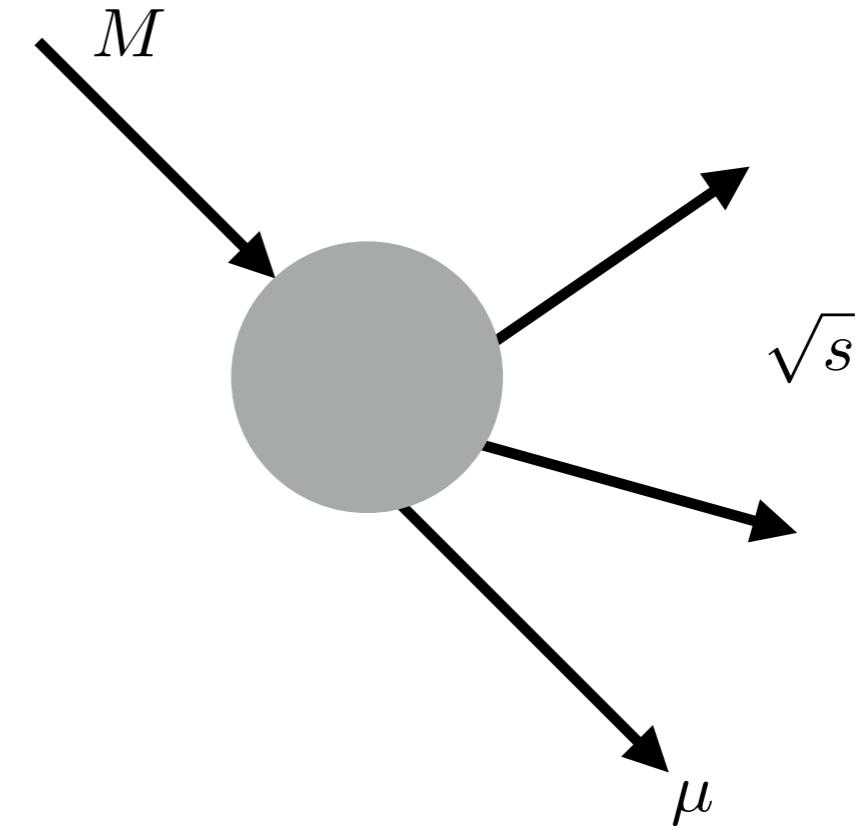
Back-Up

Triangle Mechanism

Look at 3-Body decay of particle M

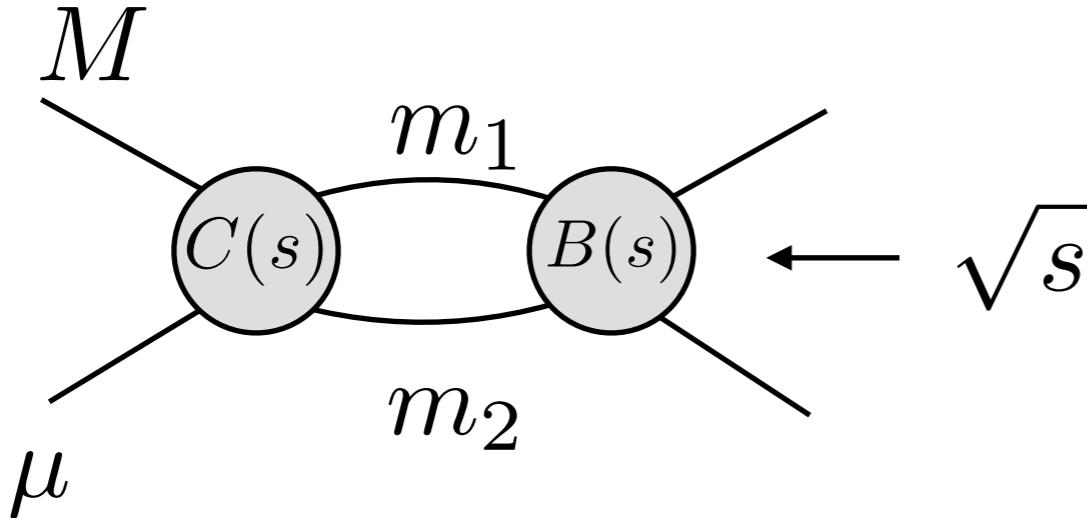
Masses can be such that non-resonant peaks appear in invariant mass

For these examples, assumption that all particles are scalar



From Crossing

Triangle Mechanism

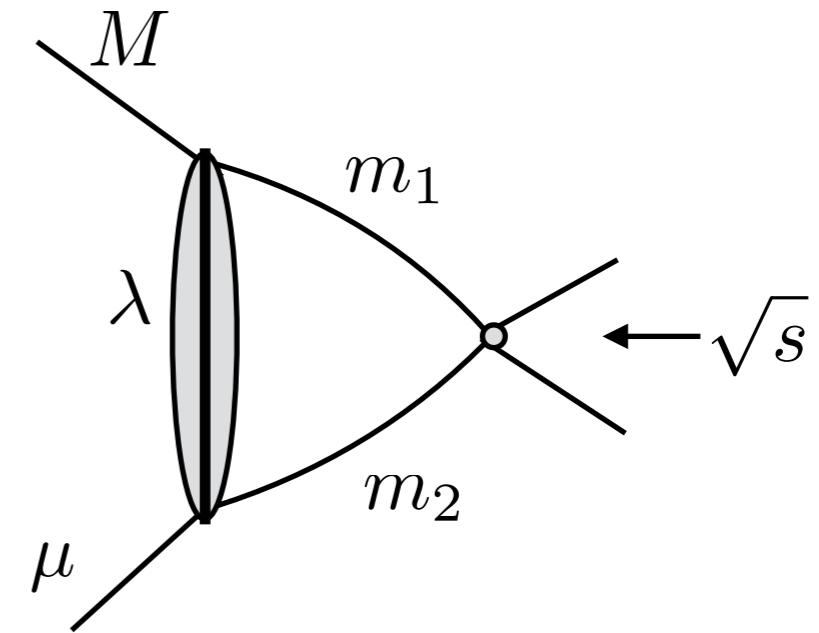


$$\text{Im } T(s) = B^*(s)\rho(s)C(s)\Theta(s - s_t)$$

$$T(s) = \frac{1}{\pi} \int ds' \frac{B^*(s')\rho(s')C(s')}{s' - s}$$

2-Body intermediate phase space

$$\rho(s)$$



For simplicity, $B(s)=\text{CONST}$

After partial wave projection

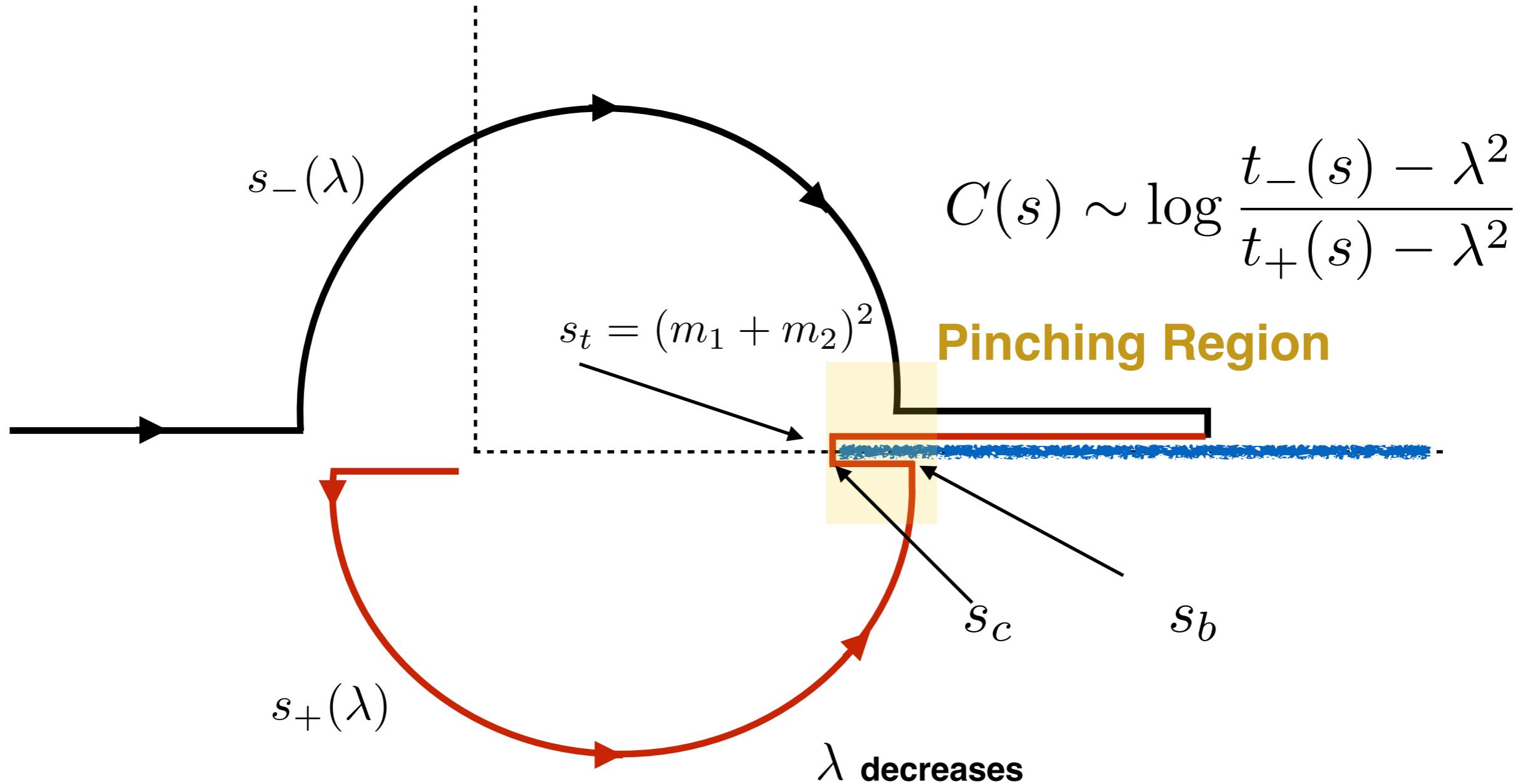
$$C(s) \sim \log \frac{t_-(s) - \lambda^2}{t_+(s) - \lambda^2}$$

Look for singularities of $C(s)$

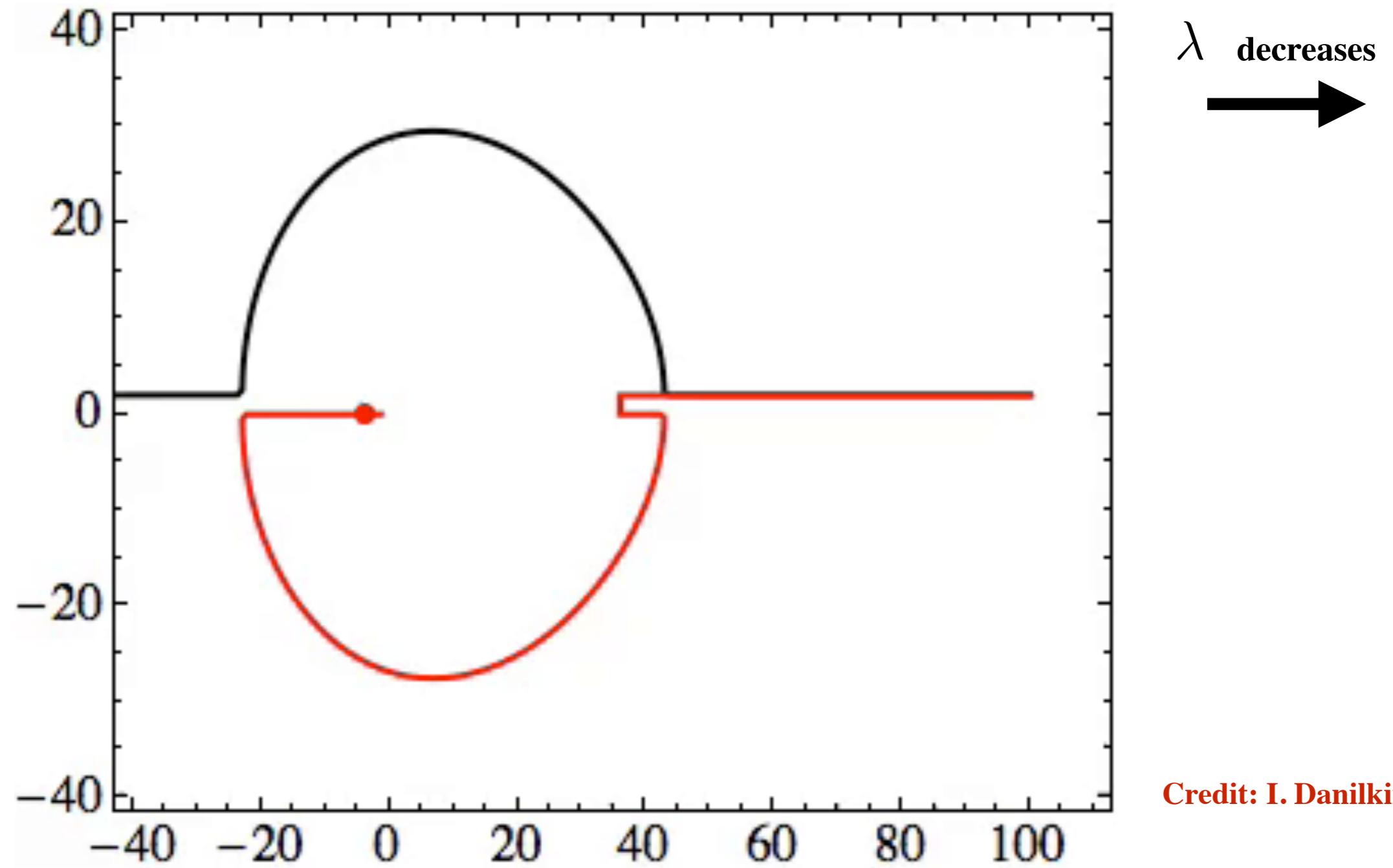
[A. Szczepaniak, arXiv:1501.01691v3]

Triangle Branching Points

- Find 2 solutions $s_{+,-}$ as a function of λ for singular $C(s)$



Branch Cut Movement



Triangle Mechanism

- Peaks appear in physical region when
$$\lambda_c < \lambda < \lambda_b$$
$$s_c < s_{peak} < s_b$$
- These are associated with the pinching effect, not poles, in the amplitude
- Two examples: Z(3900) & Pc(4450)

