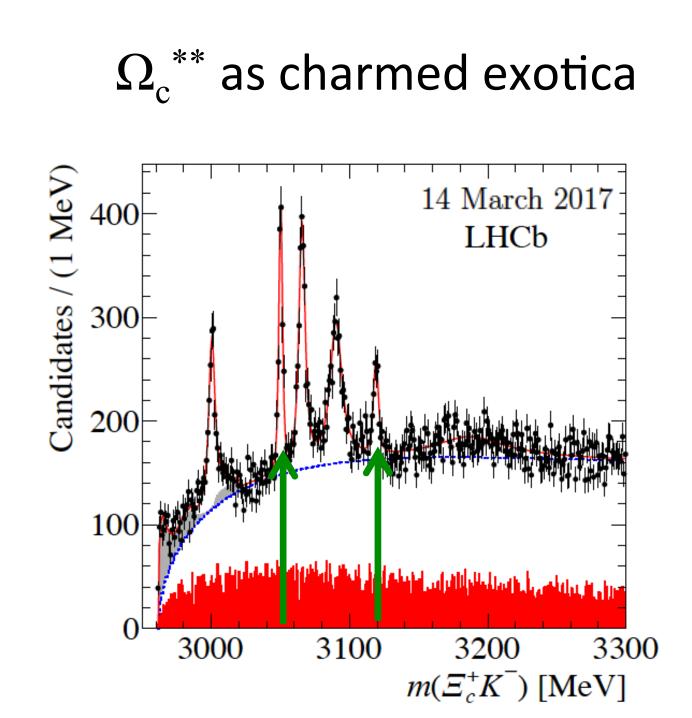
Possibility of the existence of charm exotica

Michał Praszałowicz (Jagiellonian University, Kraków)

in collaboration with M.V. Polyakov (Bochum) H.-C. Kim (Incheon)



Resonance	Mass (MeV)	$\Gamma (MeV)$
$\Omega_{c}(3000)^{0}$	$3000.4 \pm 0.2 \pm 0.1^{+0.3}_{-0.5}$	$4.5\pm0.6\pm0.3$
$\Omega_{c}(3050)^{0}$	$3050.2 \pm 0.1 \pm 0.1^{+0.3}_{-0.5}$	$0.8 \pm 0.2 \pm 0.1$
	70 MeV	$< 1.2\mathrm{MeV}, 95\%~\mathrm{CL}$
$\Omega_{c}(3066)^{0}$	$3065.6 \pm 0.1 \pm 0.3^{+0.3}_{-0.5}$	$3.5\pm0.4\pm0.2$
$\Omega_{c}(3090)^{0}$	$3090.2 \pm 0.3 \pm 0.5^{+0.3}_{-0.5}$	$8.7\pm1.0\pm0.8$
$\Omega_{c}(3119)^{0}$	$3119.1 \pm 0.3 \pm 0.9^{+0.3}_{-0.5}$	$1.1\pm0.8\pm0.4$
		$< 2.6\mathrm{MeV}, 95\%~\mathrm{CL}$
$\Omega_{c}(3188)^{0}$	$3188 \pm 5 \pm 13$	$60 \pm 15 \pm 11$

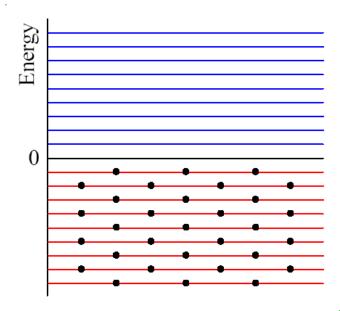
- Chiral Quark Soliton Model intro
- Light sector phenomenology
- Narrow and light pentaquark
- Soliton model for heavy baryons: ground state: 3-bar and 6
- Excitations: regular and exotic
- Test on known ground states and on 3-bar excitations
- Interpretation of the LHCb data
- Consequences and summary

QCD: quarks and gluons integrate out gluons many quark nonlocal interactions Lagrangian chirally symmetric approximation: manyq, nonl. \rightarrow 4q, local Nambu Jona Lasinio model spontaneous chiral symmetry breaking semibosonization: $q\bar{q}q\bar{q} \rightarrow q\bar{q}\pi$

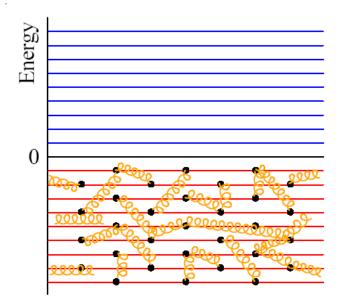
Chiral Quark Model

M. Praszałowicz

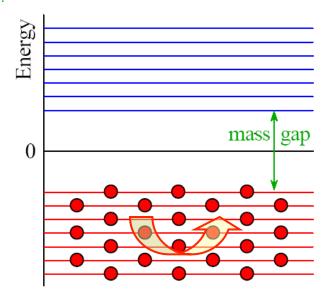
QCD vacuum:



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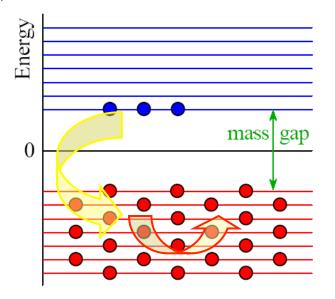
chiral symmetry breaking:



chirally inv. manyquark int.



adding vlence quarks:



chirally inv. manyquark int.

Chiral Quark Soliton Model n: "classical" baryon: Energy $K^P = 0^+$ gap mass

due to hedgehog symmetry of the mean field only grand spin

$$K = T + S$$

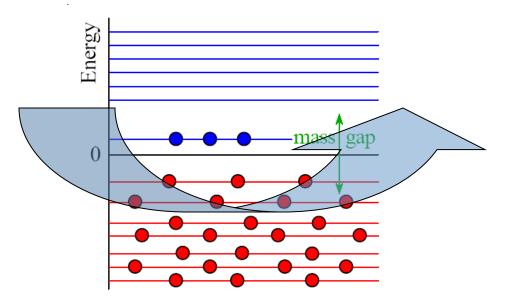
is a *good* quantum number

chirally inv. manyquark int.

soliton configuration no quantum numbers except B



'quantum" baryon:



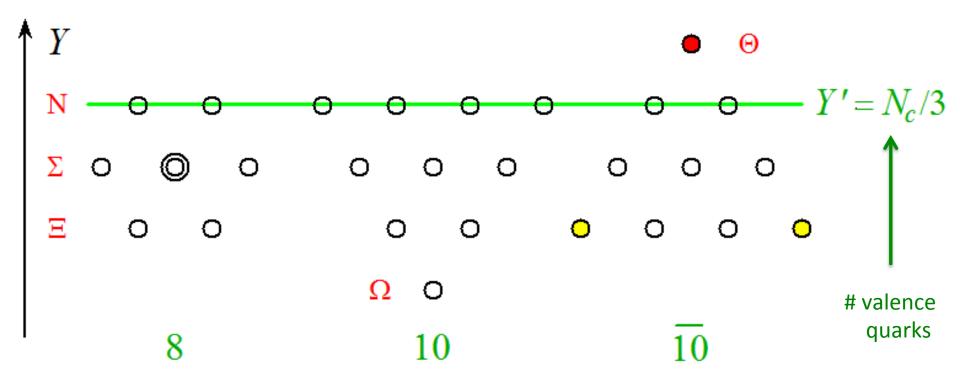
chirally inv. manyquark int.

soliton configuration no quantum numbers except B

rotation generates flavor and spin

Allowed states

- allowed SU(3) representations must contain states with hypercharge $Y' = N_c/3$,
- the isospin T' of the states with $Y' = N_c/3$ couples with the soliton spin J to a singlet: T' + J = 0.



Ch.V. Christov et al. Prog. Part. Nucl. Phys. 37 (1996) 91

Successful Phenomenology

In a "model independent" approach one can get good description of the existing data (including very narrow light pentaquark Θ⁺)

but also one can recover the NRQM result in a special limit

NRQM limit =

= squeezing the soliton to zero size

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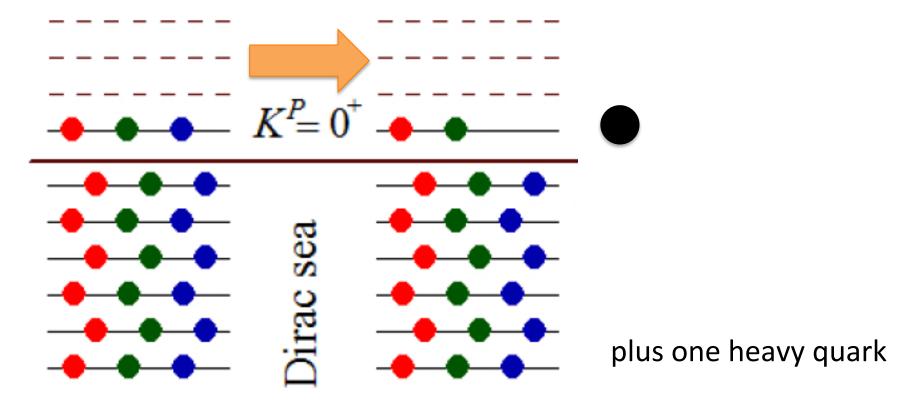
but also one can recover the NRQM result in a special limit

NRQM limit = = squeezing the soliton to zero size

$$g_A^{(3)} = \frac{5}{3}, \quad \Delta \Sigma = 1, \quad \frac{\mu_p}{\mu_n} = -\frac{3}{2} \qquad G_{\overline{10}} = 0$$

Soliton model for heavy baryons

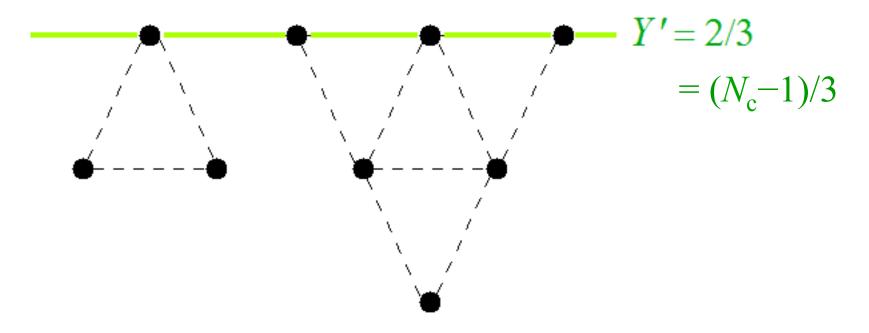
Soliton with N_c – 1 quarks if N_c is large, N_c – 1 is also large and one can use the same mean field arguments



G.S. Yang, H.C. Kim, M.V. Polyakov, MP Phys. Rev. D94 (2016) 071502

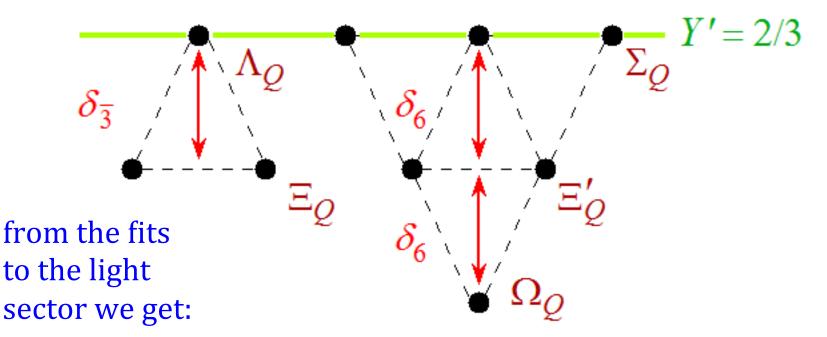
Allowed SU(3) irreps.

 $\bar{3} \ s=0$ 6 s=1



Soliton + heavy quark (splittings)

3 s=0 (1/2) 6 s=1(1/2 3/2)



 $\delta_{\overline{3}} = 203.8 \pm 3.5 \text{ MeV}, \text{ (exp.: 178 MeV)}$ $\delta_{\overline{6}} = 135.2 \pm 3.3 \text{ MeV}, \text{ (exp.: 121 MeV)}$

13%

Soliton + heavy quark (splittings)

3 s=0 (1/2) 6 s=1(1/2 3/2)

 $\delta_{\overline{3}} / \delta_{\overline{6}} / \Sigma_{Q} Y' = 2/3$ $\delta_{\overline{6}} / \Sigma_{Q} / \Sigma_{Q} / \Sigma_{Q}$ from the fits to the light sector we get: $E_{Q} / \Sigma_{Q} /$

13%

- $\delta_{\overline{3}} = -203.8 \pm 3.5 \text{ MeV}, \text{ (exp.: 178 MeV)}$
- $\delta_6 = 135.2 \pm 3.3 \text{ MeV}, \text{ (exp.: 121 MeV)}$

Soliton + heavy quark (splittings)

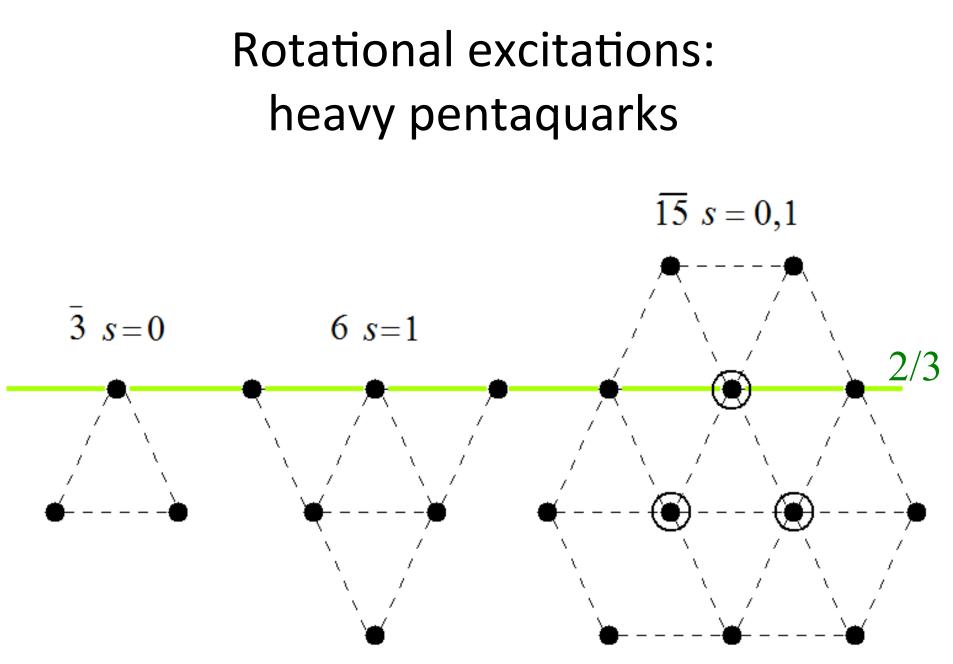
3 s=0 (1/2) 6 s=1(1/2 3/2)

Successful phenomenology fro(also for b-quark baryons) sector we get:

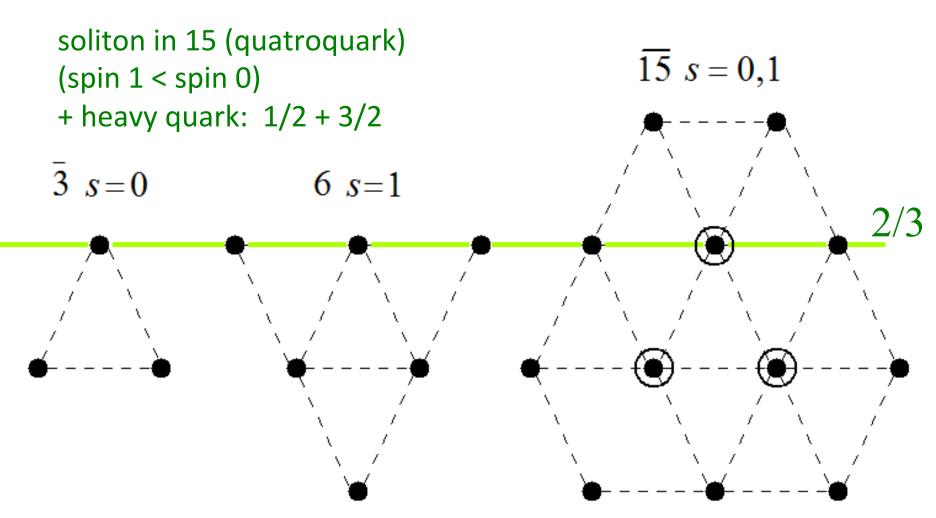
G.S. Yang, H.C. Kim, M.V. Polyakov, MP Phys. Rev. D94 (2016) 071502

13%

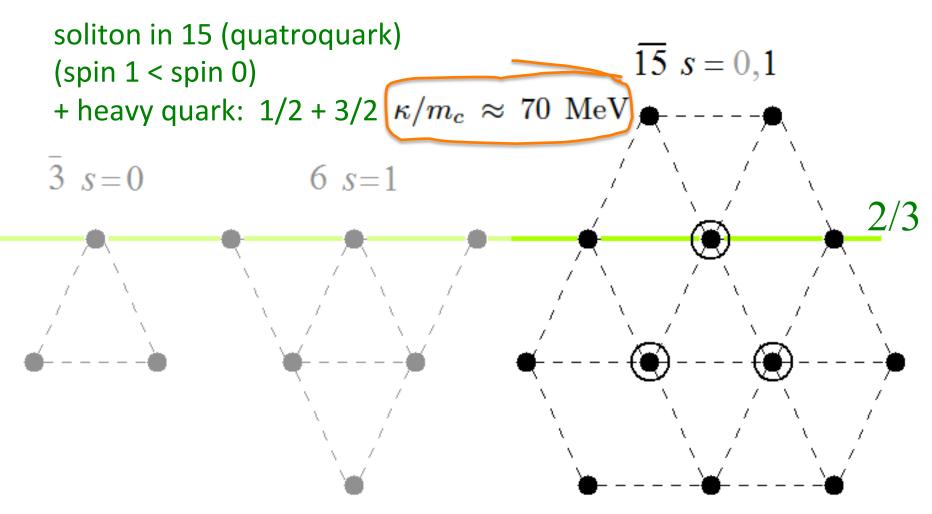
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Rotational excitations: heavy pentaquarks



Rotational excitations: heavy pentaquarks



Decay constants

$\overline{\mathbf{15}}_1 o \overline{\mathbf{3}}_0$

$\overline{\mathbf{15}}_1 o \mathbf{6}_1$

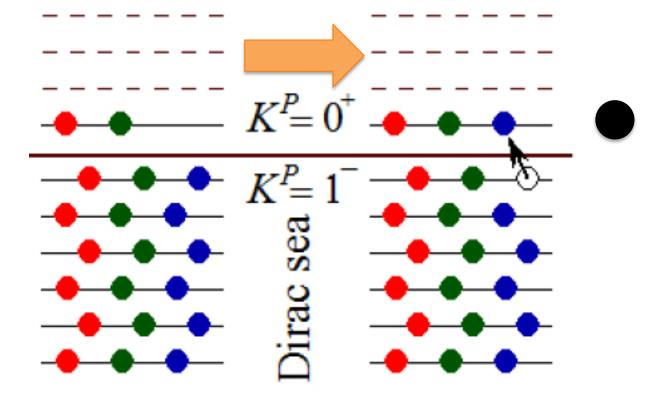
Decay constants

$\mathbf{15}_1 ightarrow \mathbf{3}_0$

${f 15}_1 o {f 6}_1$ in NRQM limit: $G_{f 6}=0$

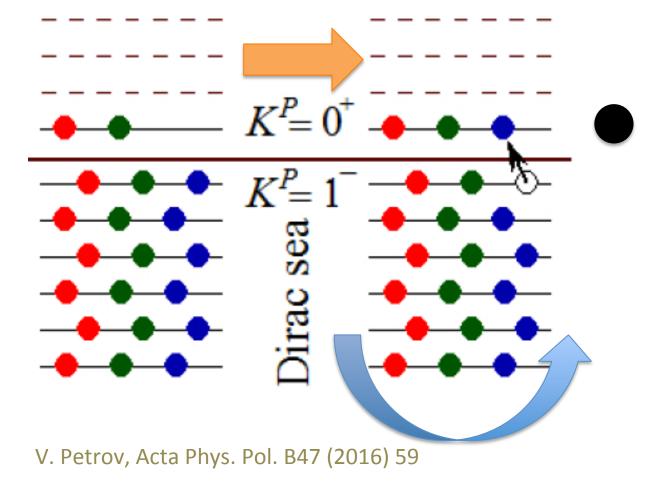
Expectations: some decays will be suppressed

Quark excitations: non-exotic heavy baryons



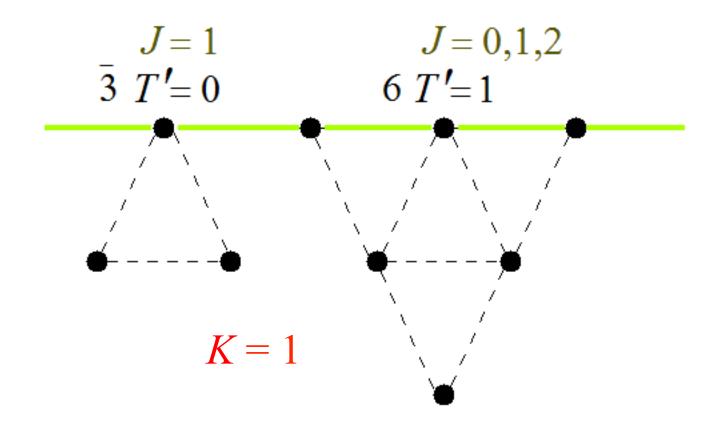
V. Petrov, Acta Phys. Pol. B47 (2016) 59

Quark excitations: non-exotic heavy baryons

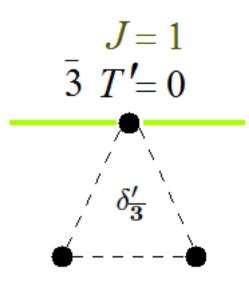


One K=1 quark excited solitons

• the isospin T' of the states with $Y' = (N_c - 1)/3$ couples with the soliton spin J as follows: T' + J = K, where K is the grand spin of the excited level.



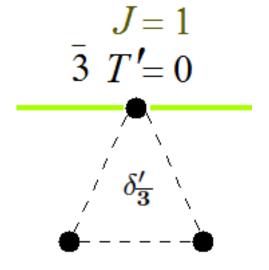
3bar excited heavy baryons



add heavy quark total spin 1/2 and 3/2

 $\delta_{\overline{\mathbf{3}}}' = \delta_{\overline{\mathbf{3}}} = -180 \text{ MeV}$

3bar excited heavy baryons



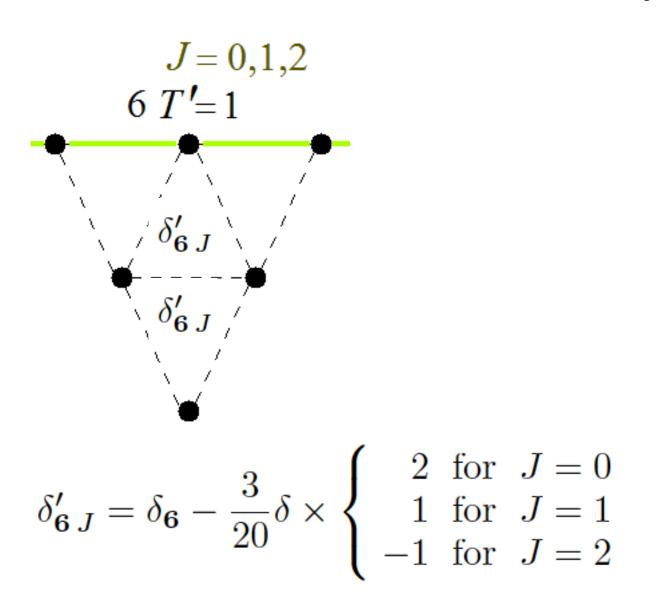
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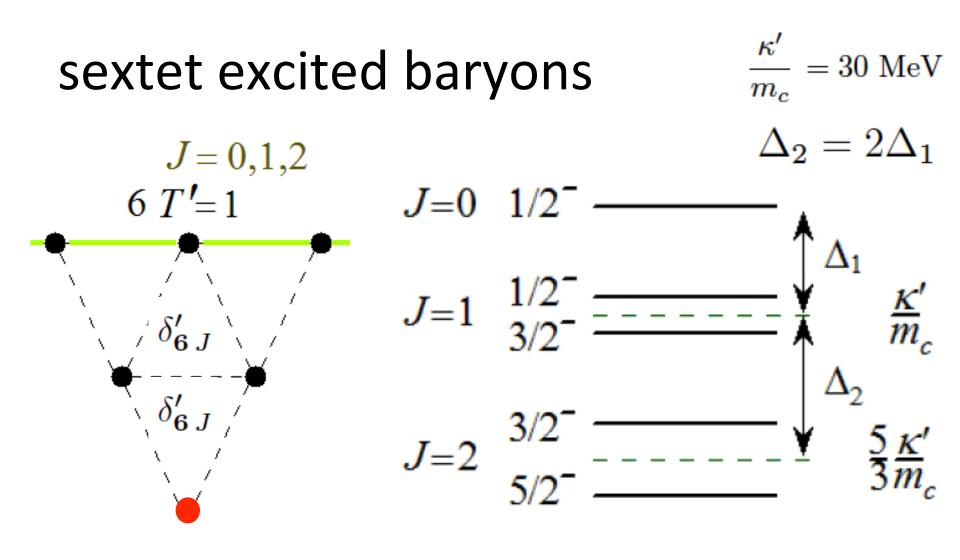
 $\delta_{\overline{\mathbf{3}}}' = \delta_{\overline{\mathbf{3}}} = -180 \text{ MeV}$

experimentally: $\Lambda_{c}(2628)$ $\Lambda_{c}(2592)$ 190 MeV 198 MeV $\Xi_{c}(2818)$ $\Xi_{c}(2790)$ $\frac{\kappa'}{\kappa} = 30 \text{ MeV}$ $(1/2)^{-}$ m_c $H_{\rm hf} = \frac{2}{3} \frac{\kappa}{m_{\rm O}} \boldsymbol{J} \cdot \boldsymbol{J}_Q$

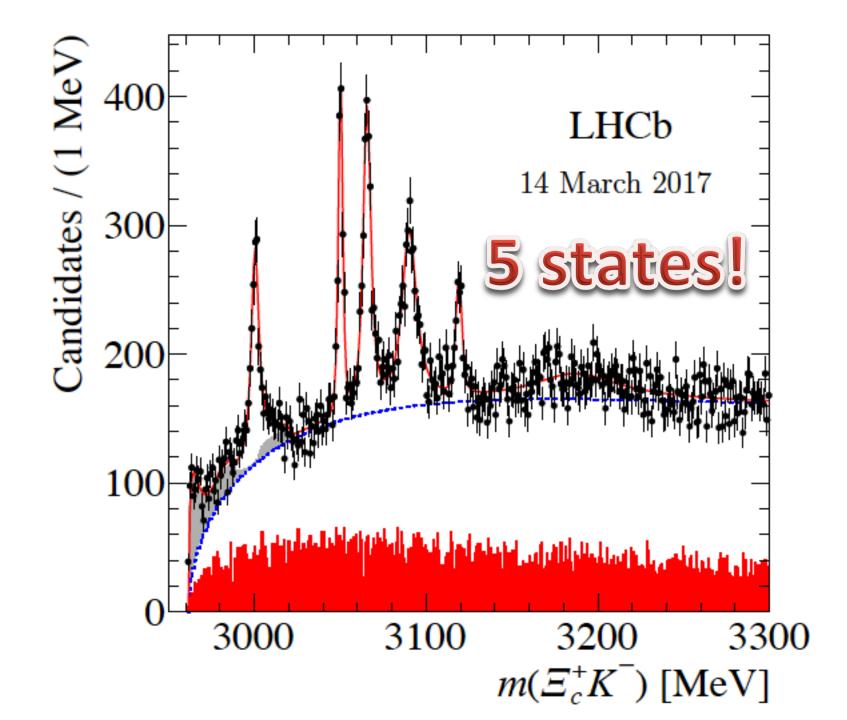
hyprfine splitting different from the ground state

sextet excited baryons





excited Omega_Q spectrum, 5 states



Scenario 1:

all LHCb Omega's are sextet states

J	S^P	$M \; [MeV]$	$\kappa'/m_c \; [\text{MeV}]$	$\Delta_J [{ m MeV}]$
0	$\frac{1}{2}^{-}$	3000	_	_
1	$\frac{1}{2}^{-}$	3050	16	61
	$\frac{3}{2}^{-}$	3066		
2	$\frac{3}{2}^{-}$	3090	17	47
	$\frac{5}{2}$ -	3119		

violates constraints: $\frac{\kappa'}{m_c} = 30 \text{ MeV} \quad \Delta_2 = 2\Delta_1$

Scenario 1:

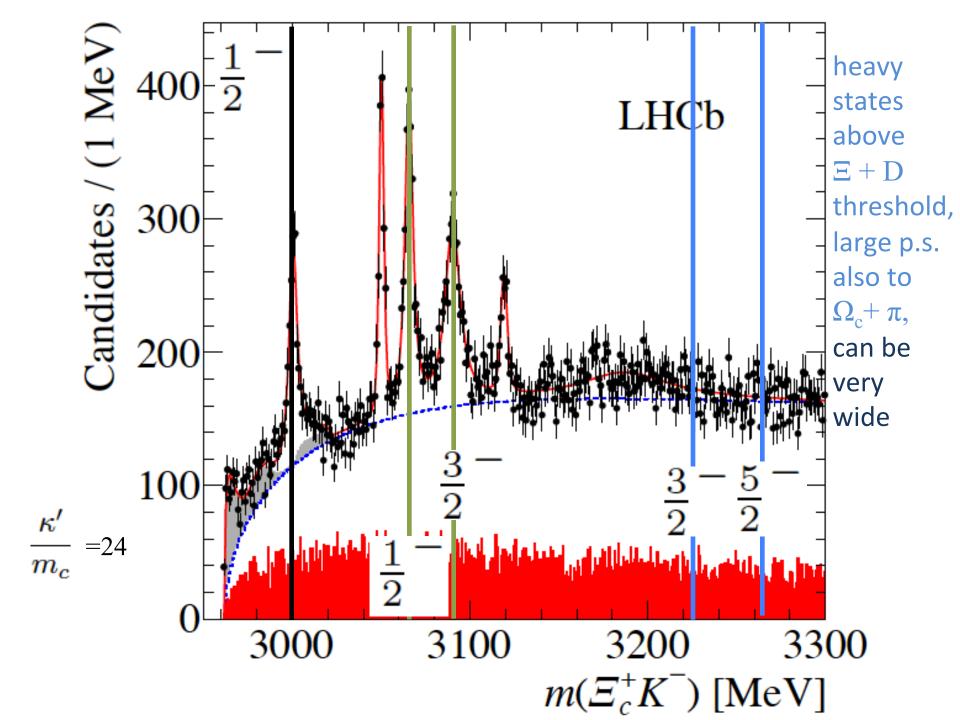
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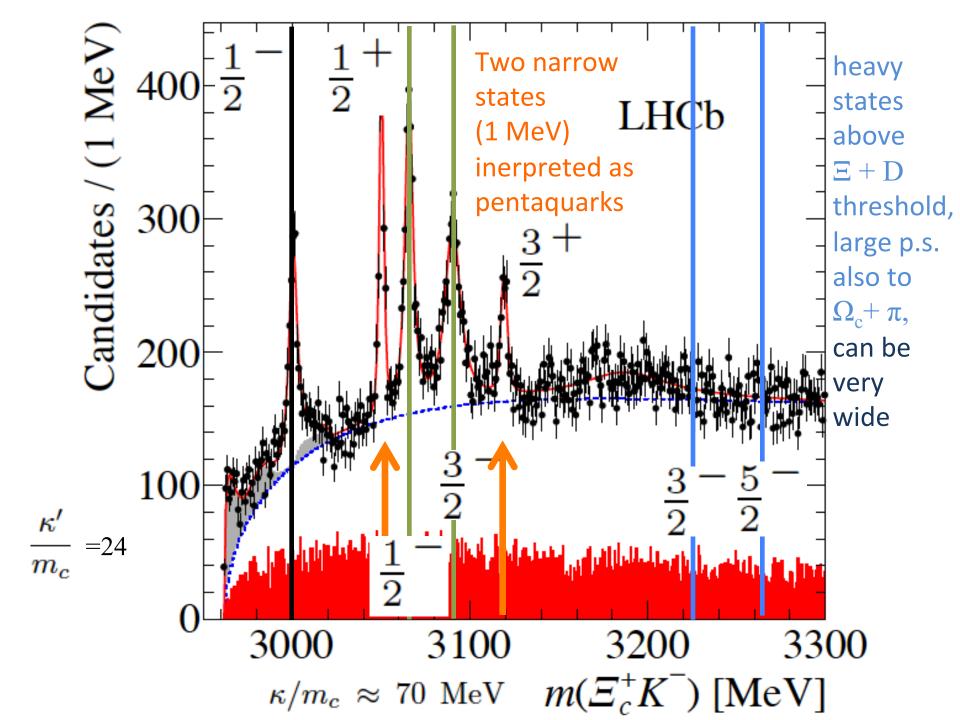
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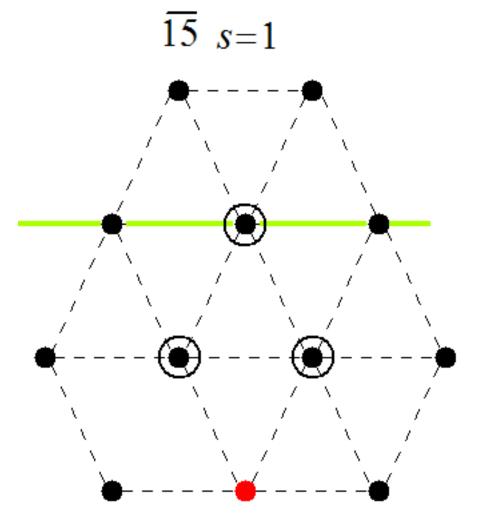
similar problem in the quark models

Scenario 2 force sextet constraints





Consequences



$\Xi_c^0 + K^- \text{ or } \Xi_c^+ + \overline{K}^0$

Omega's form isospin triplet, easy to check experimentally

Consequences

$$15 \ s=1$$

rich structure -

- many new states,

also in the case of b baryons

$$\Xi_c^0 + K^- \text{ or } \Xi_c^+ + \bar{K}^0$$

Omega's form isospin triplet, easy to check experimentally

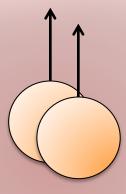
Conclusions

- soliton models **ARE** quark models
- successful phenomenology in the light baryon sector
- in soliton models pentaquarks are **naturally light**
- in NR limit **no decay** of antidecuplet to ocet (!)
- heavy baryons can be desribed in terms of N_c-1 quark soliton
- two types of excitations:
 - rotations: 3-bar, 6 (regular) and 15-bar (exotic)
 - **quark** excitations: 3-bar, 6 (regular)
- **two** of the LHCb Omega_c states may be interpreted as **5q**
- isospin partners make the model easy to falsify or to confirm
- similar structure for **bottom**

Further developements



Further developements



heavy tetraquark