



Hadron production and spectroscopy at ATLAS

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on behalf of the ATLAS Collaboration

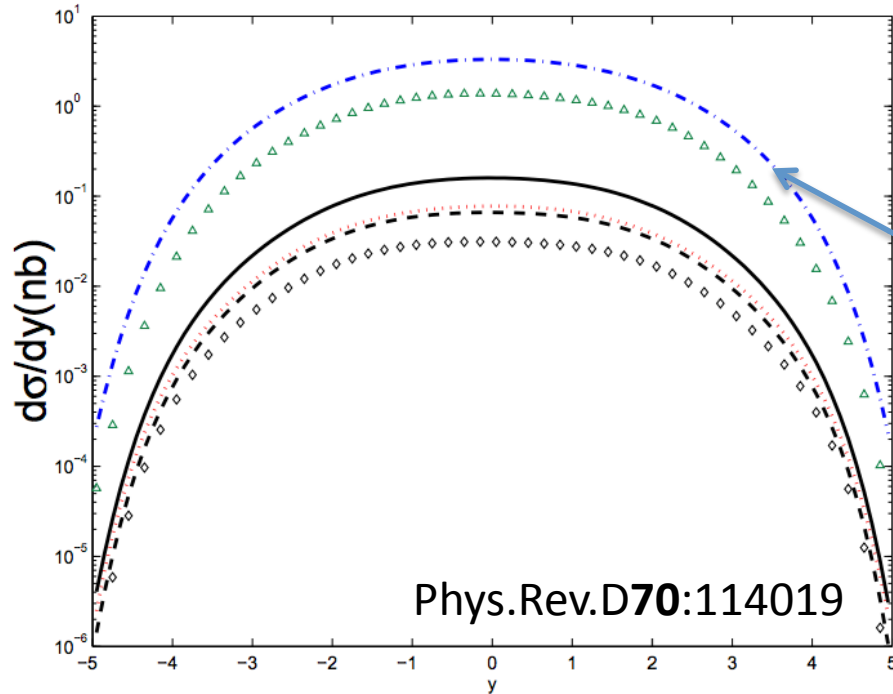
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Outline

- ATLAS has a rich program for studies of hadron production and spectroscopy.
- Selected results covered in this talk:
 - $B_c(2S)$ observation
 - Study of the $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ decays
 - Search for the X_b and other hidden-beauty states
 - Measurement of the branching ratio $\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0) / \Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)$
 - Production measurements of $\psi(2S)$ and $X(3872)$
- <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/BPhysPublicResults>
- Much more to come soon!

B_c(2S) observation

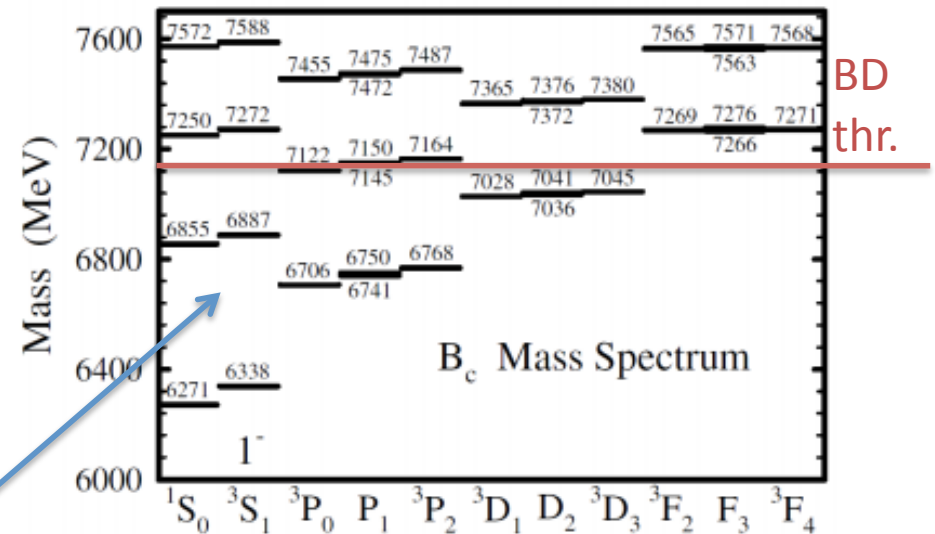
- [Phys. Rev. Lett. 113, 212004 \(2014\)](#)



Ground state: $B_c^\pm \rightarrow J/\Psi(\mu^+\mu^-) \pi^\pm$

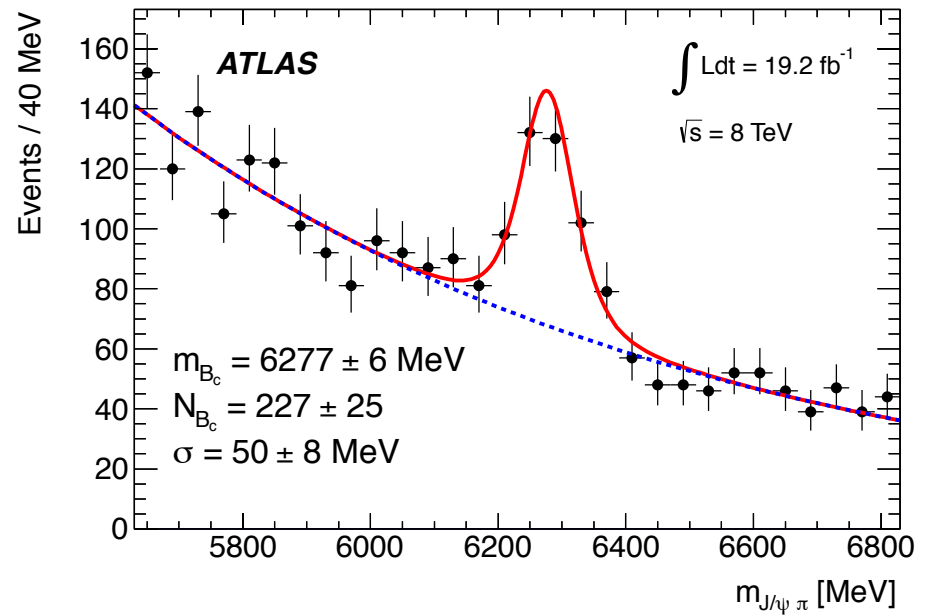
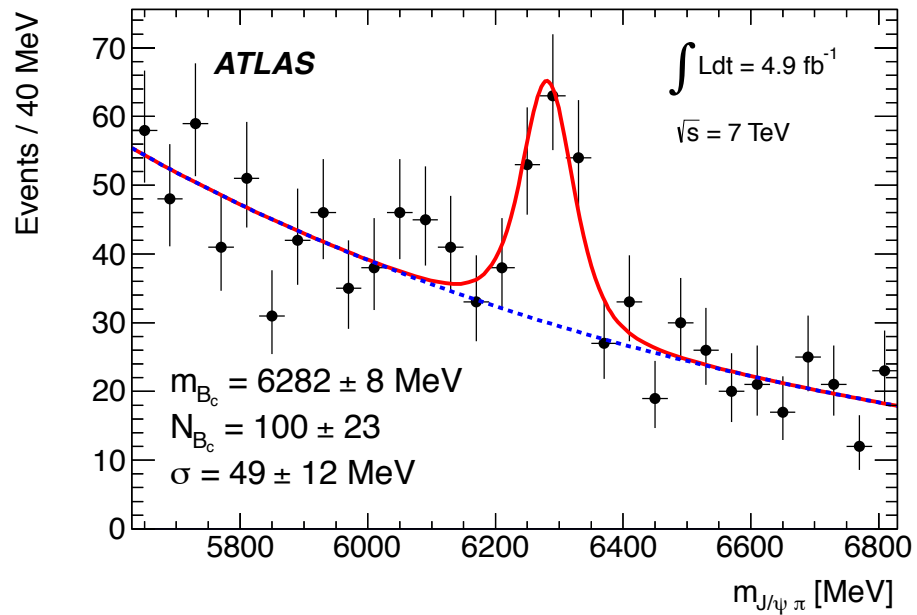
Peak is observed in the mass difference distribution $m(B_c^\pm(2S)) - m(B_c^\pm) - 2m(\pi^\pm)$

- B_c is a unique system: two distinct heavy quarks.
- “Intermediate” between the b \bar{b} and c \bar{c} .
- Many theoretical predictions for the bc-family spectroscopy and production properties.
- Production is predicted to be suppressed in the forward-backward regions (various excited states).



S. Godfrey, PRD 70, 054017(2004)

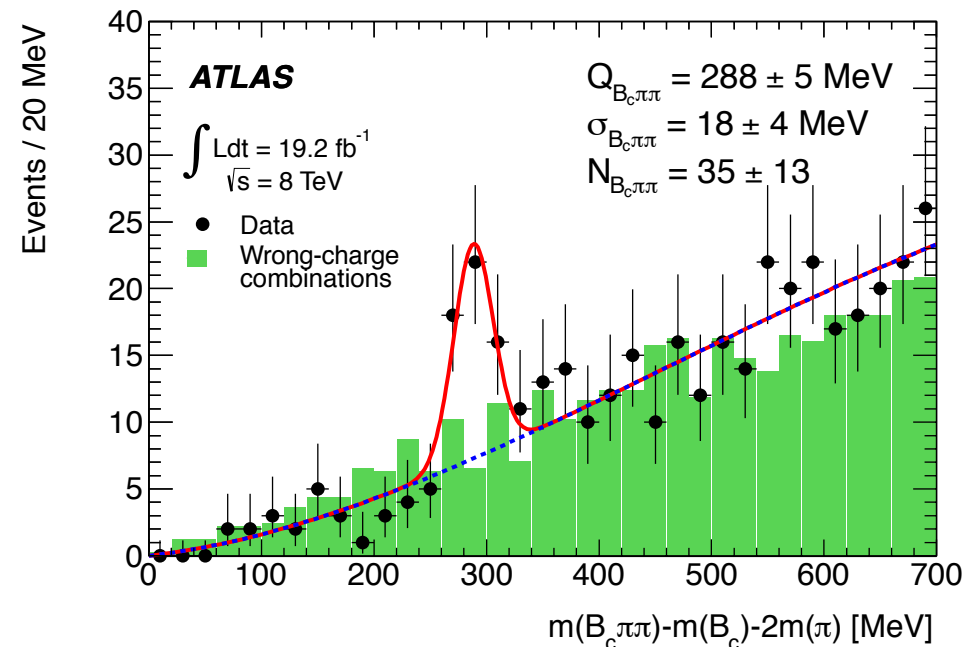
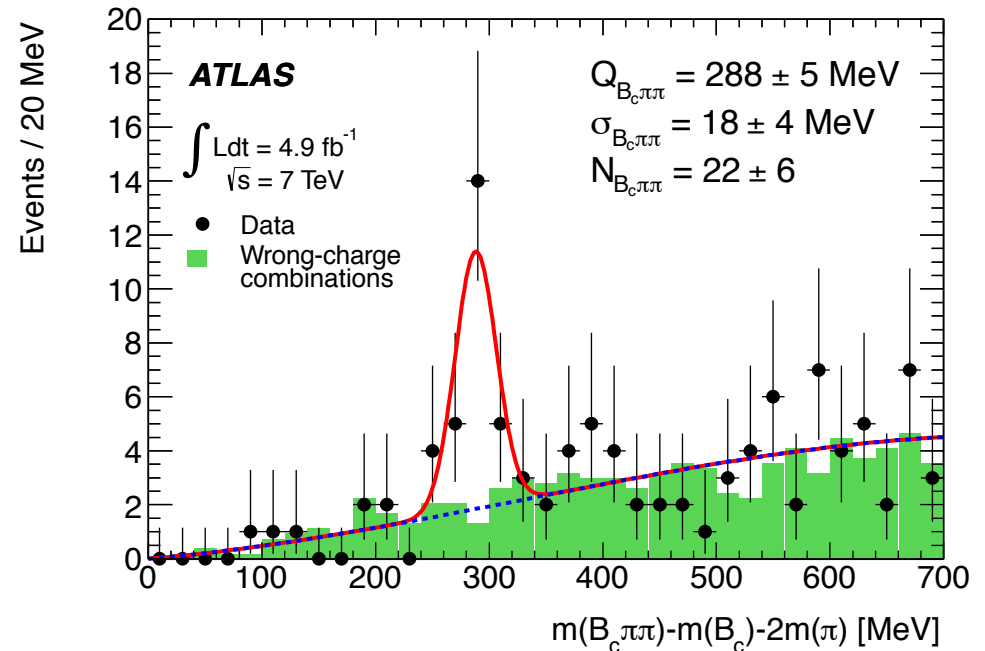
$B_c(2S)$, cont.



- Ground state: $B_c^\pm \rightarrow J/\Psi(\mu^+\mu^-) \pi^\pm$

$B_c(2S)$, cont.

- **First observation** of the excited B_c state, the $B_c(2S)$.
- Mass of the peak $m[B_c(2S)] = 6842 \pm 4_{\text{stat}} \pm 5_{\text{syst}}$ MeV is consistent with the theory [6835-6917], as well as the decay mode.
- The observed yield is within the theoretical expectations.
- **5.2 sigma observation** for the combination of 7 and 8 TeV datasets (LEE included, **local is 5.4**).
- Further studies possible with new data.



Study of the $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ decays

- [Eur. Phys. J. C, 76\(1\), 1 \(2016\)](#).

- $\bar{b}c \rightarrow c\bar{c}\bar{s}$ transition:

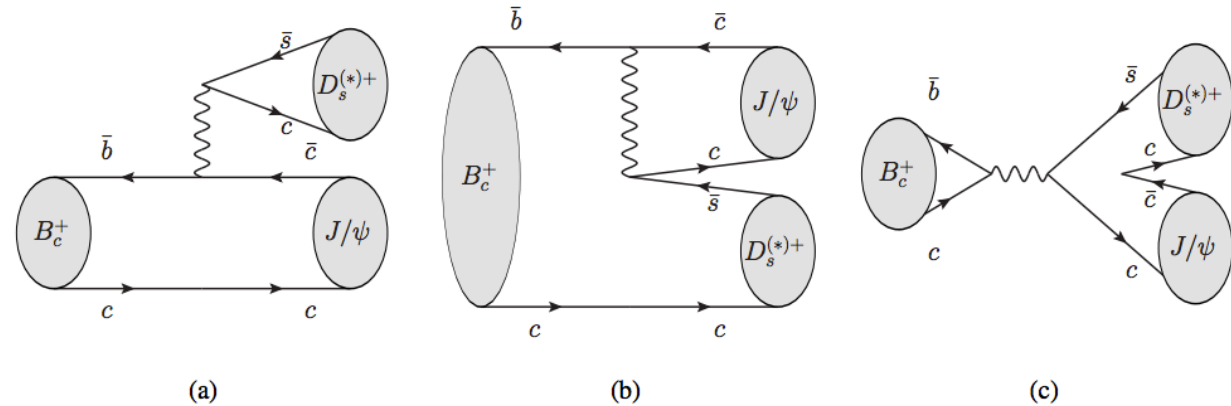
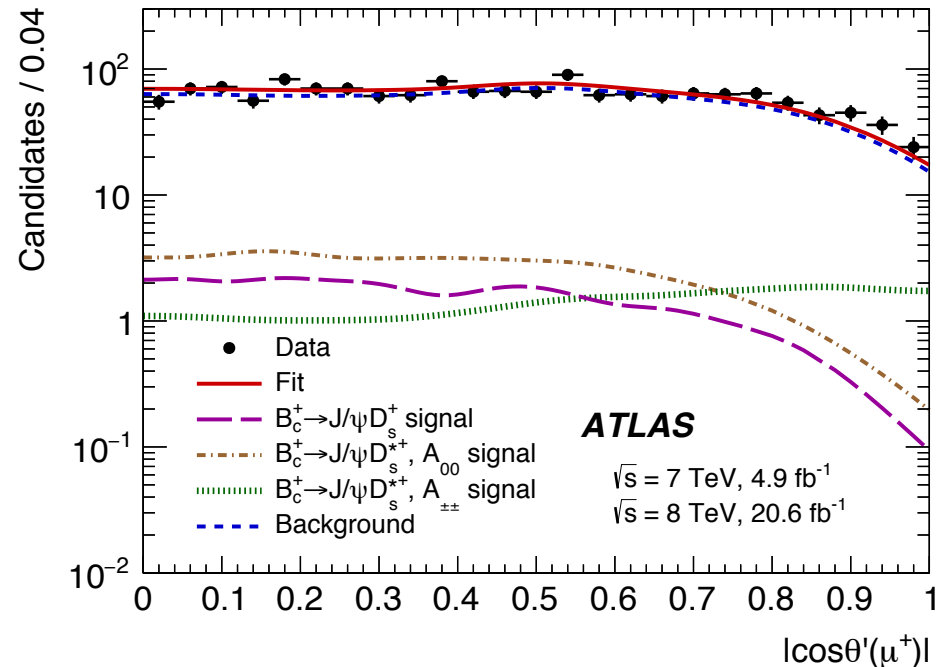
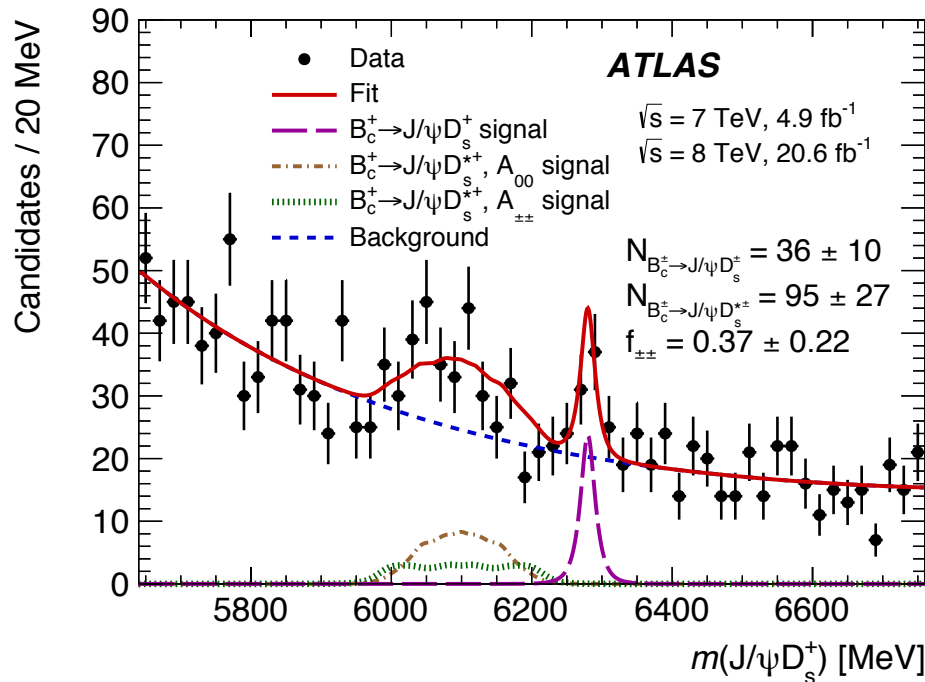


Figure 1: Feynman diagrams for $B_c^+ \rightarrow J/\psi D_s^{(*)+}$ decays: (a) colour-favoured spectator, (b) colour-suppressed spectator, and (c) annihilation topology.

- Measurement of the relative branching fractions of $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ and of the branching fractions relative to $B_c^+ \rightarrow J/\psi \pi^+$ decay.
- $D_s^{*+} \rightarrow D_s^+ [\pi^0/\gamma]_{\text{soft}}$, $D_s^+ \rightarrow \varphi (K^+ K^-) \pi^+$.
- The decay $B_c^+ \rightarrow J/\psi D_s^{*+}$ is a transition of a pseudoscalar meson into a pair of vector states and is thus described by the three helicity amplitudes, A_{++} , A_{--} , and A_{00} .
- The contribution of the A_{++} and A_{--} amplitudes, referred to as the $A_{\pm\pm}$ component, corresponds to the J/ψ and D_s^{*+} transverse polarisation.
- Measurement of the fraction of transverse polarisation, $\Gamma_{\pm\pm}/\Gamma = \Gamma_{\pm\pm}(B_c^+ \rightarrow J/\psi D_s^{*+})/\Gamma(B_c^+ \rightarrow J/\psi D_s^{*+})$. From the naïve spin model should be 2/3.

Study of the $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ decays, results (1)

- 2D unbinned maximum likelihood fit of B_c^+ invariant mass and helicity angle $\theta'(\mu^+)$.



- Measured transverse polarisation fraction $\Gamma_{\pm\pm}/\Gamma = \Gamma_{\pm\pm}(B_c^+ \rightarrow J/\psi D_s^{*+})/\Gamma(B_c^+ \rightarrow J/\psi D_s^{*+})$

$$\Gamma_{\pm\pm}/\Gamma = 0.38 \pm 0.23 \text{ (stat.)} \pm 0.07 \text{ (syst.)}$$

Study of the $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ decays, results (2)

- Ratios of the branching fractions:

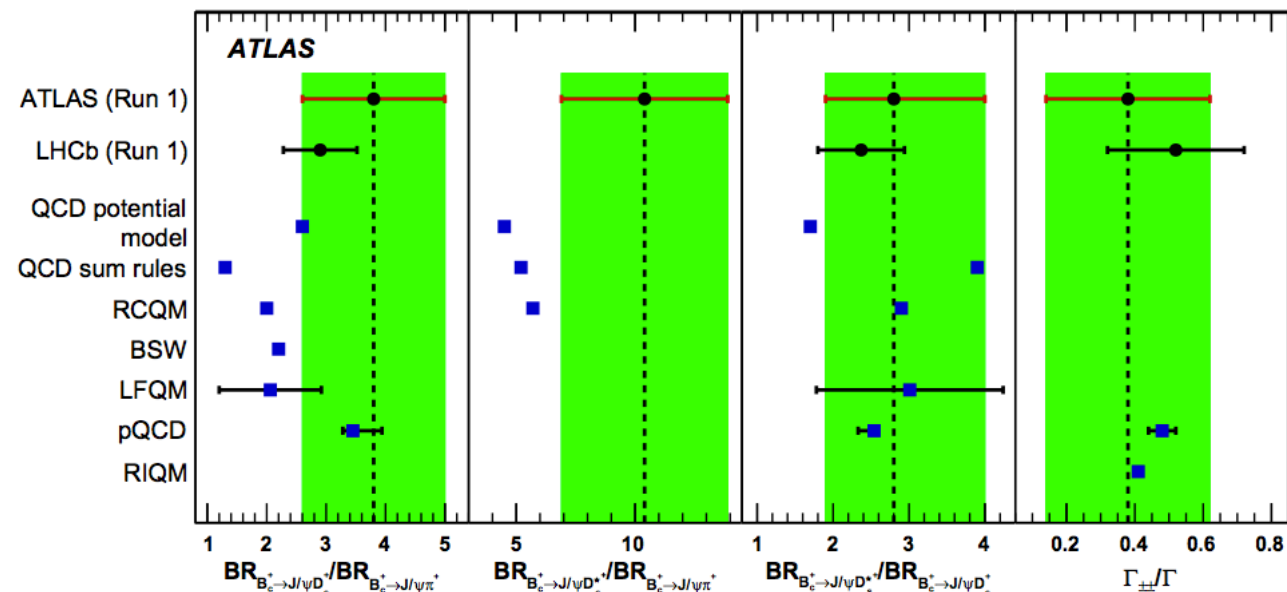
$$\mathcal{R}_{D_s^+/\pi^+} = \frac{\mathcal{B}_{B_c^+ \rightarrow J/\psi D_s^+}}{\mathcal{B}_{B_c^+ \rightarrow J/\psi \pi^+}} = 3.8 \pm 1.1 \text{ (stat.)} \pm 0.4 \text{ (syst.)} \pm 0.2 \text{ (BF)},$$

$$\mathcal{R}_{D_s^{*+}/\pi^+} = \frac{\mathcal{B}_{B_c^+ \rightarrow J/\psi D_s^{*+}}}{\mathcal{B}_{B_c^+ \rightarrow J/\psi \pi^+}} = 10.4 \pm 3.1 \text{ (stat.)} \pm 1.5 \text{ (syst.)} \pm 0.6 \text{ (BF)},$$

$$\mathcal{R}_{D_s^{*+}/D_s^+} = \frac{\mathcal{B}_{B_c^+ \rightarrow J/\psi D_s^{*+}}}{\mathcal{B}_{B_c^+ \rightarrow J/\psi D_s^+}} = 2.8_{-0.8}^{+1.2} \text{ (stat.)} \pm 0.3 \text{ (syst.)}$$

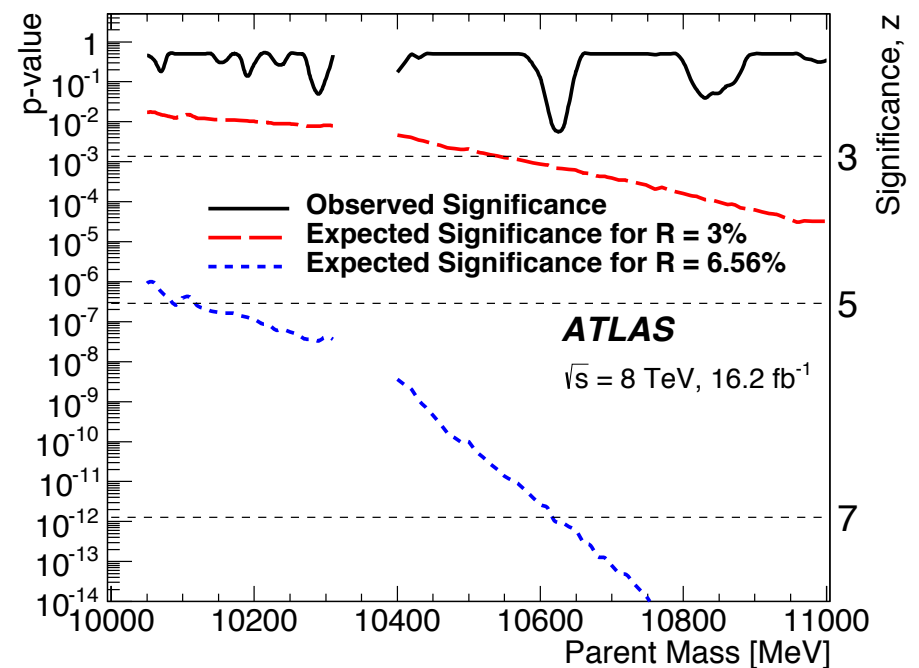
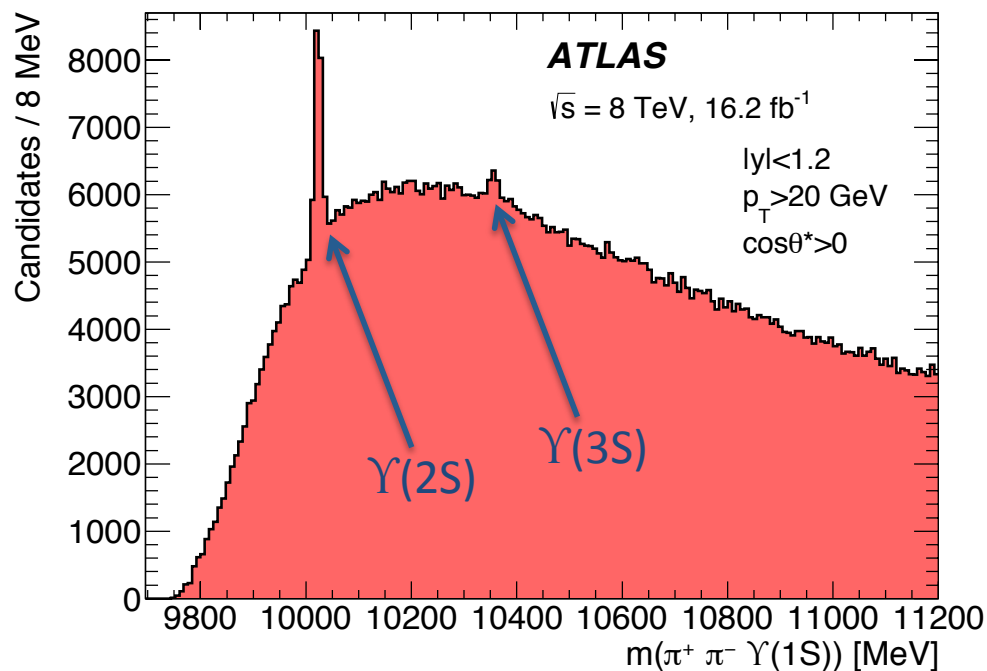
- (BF) corresponds to $\text{Br}(D_s^+ \rightarrow \phi(K^+K^-)\pi^+)$

- Comparison with theory:

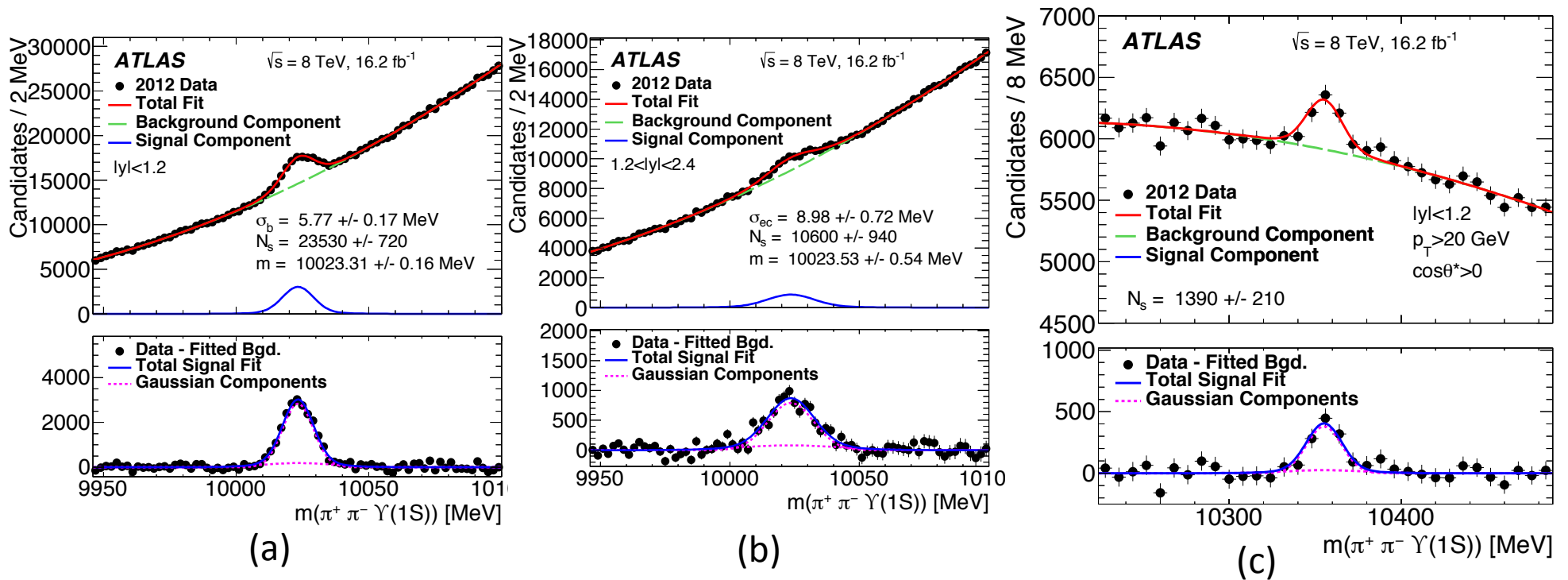


Search for the X_b and other hidden-beauty states at ATLAS (1)

- [Physics Letters B 740 \(2015\), pp. 199–217](#)
- 16.2 fb⁻¹ of 8 TeV ATLAS data.
- Search in the decay channel $X_b \rightarrow \pi^+ \pi^- \Upsilon(1S) (\rightarrow \mu^+ \mu^-)$.
- Analysis is performed in eight bins of rapidity, transverse momentum, and the angle (in the rest frame of the parent state) between the dipion system and the laboratory-frame momentum of the parent.

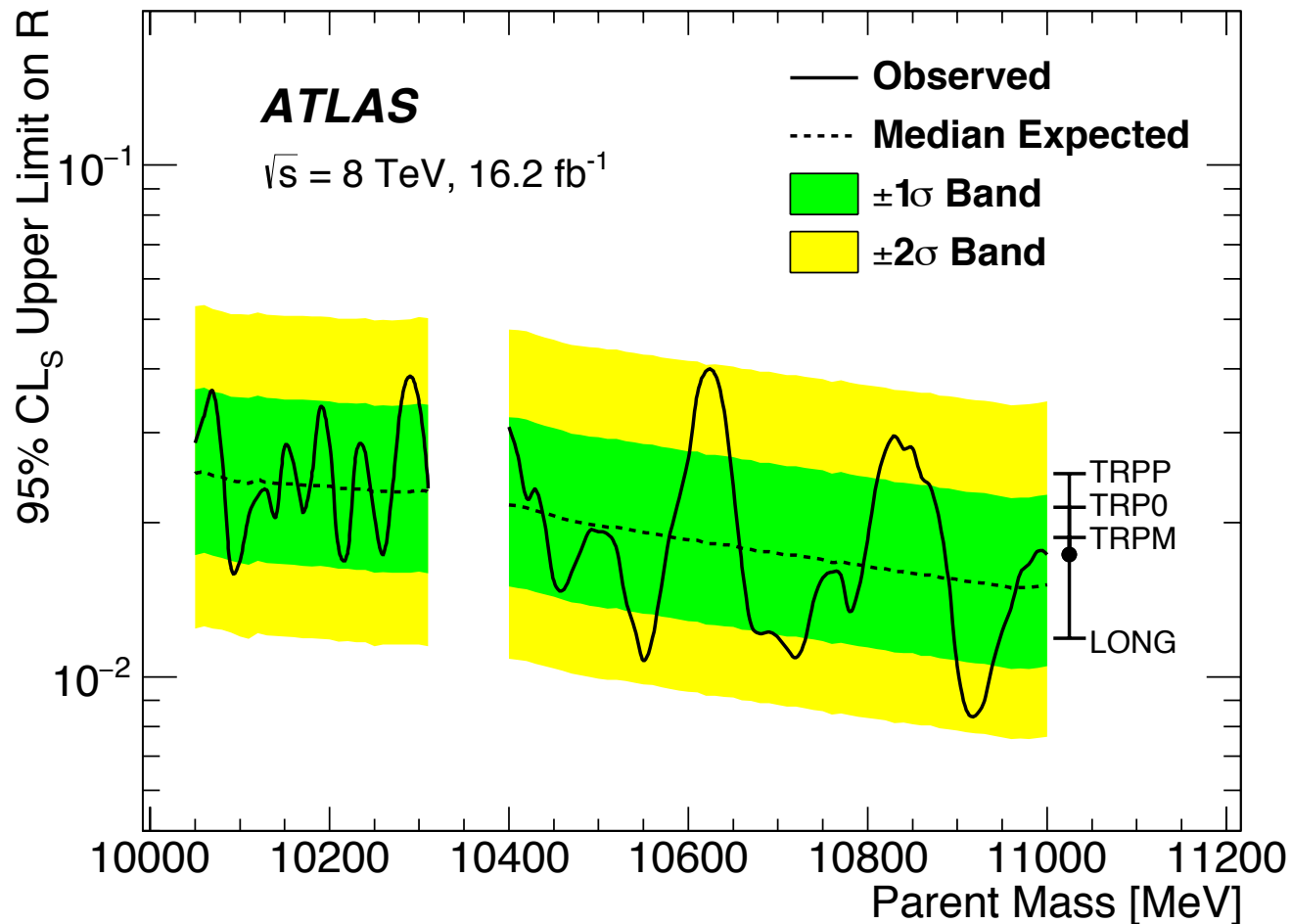


Search for the X_b and other hidden-beauty states at ATLAS (2)



- The signal shape, distribution of signal among analysis bins, and efficiency, are first calibrated on the $\Upsilon(2S)$ peak in data \rightarrow that it turns out only one parameter: the division of signal between barrel and endcap that needs to be adjusted.
- After this, the simultaneous fit to the 8 bins is validated on the $\Upsilon(3S)$.
- Both the expected overall yield and the division among bins are reproduced well.

Search for the X_b and other hidden-beauty states at ATLAS (3)



- The bar on the right shows typical shifts under alternative X_b spin-alignment scenarios.
- Upper limits are recalculated under longitudinal ('LONG') and three transverse ('TRPP', 'TRPO', 'TRPM') spin-alignment scenarios.

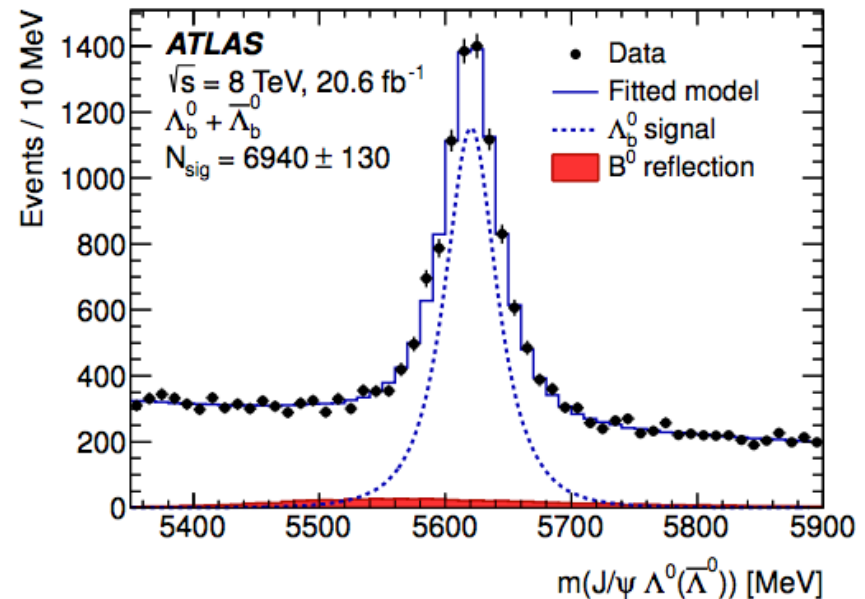
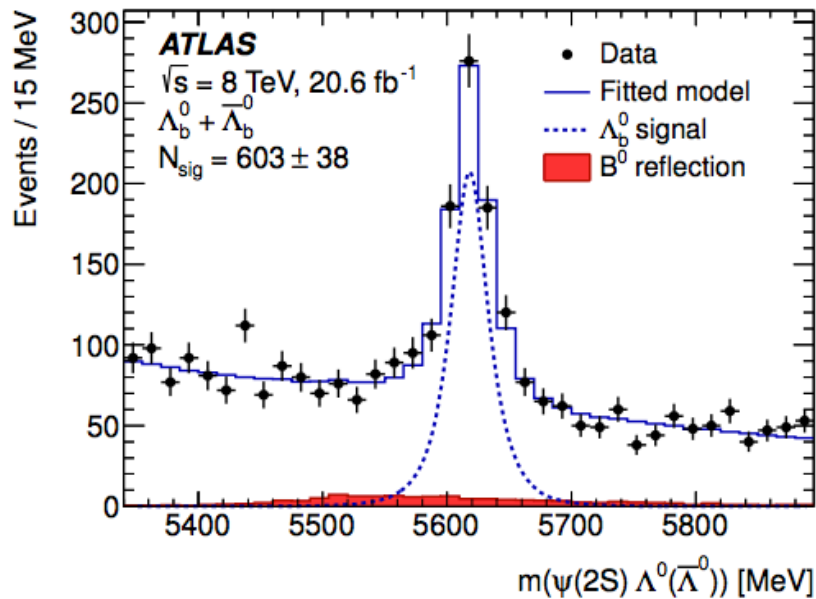
- No evidence for new narrow states is found for masses 10.05–10.31 GeV and 10.40–11.00 GeV.
- Upper limits are also set on the ratio

$$R = \frac{[\sigma(pp \rightarrow X_b) B(X_b \rightarrow \pi^+\pi^-\Upsilon(1S))]}{[\sigma(pp \rightarrow \Upsilon(2S)) B(\Upsilon(2S) \rightarrow \pi^+\pi^-\Upsilon(1S))]},$$
 with results ranging from 0.8% to 4.0% depending on the X_b mass.

Measurement of the branching ratio

$$\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0) / \Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)$$

- [Physics Letters B 751 \(2015\) pp. 63-80.](#)
- 20.6 fb⁻¹ of 8 TeV data.
- First observation of $\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0$ decay.



Measurement of the branching ratio

$$\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0) / \Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)$$

- Fiducial range: $p_T(\Lambda_b^0) > 10 \text{ GeV}$, $|\eta(\Lambda_b^0)| < 2.1$
- The relative branching ratio is calculated as

$$\frac{\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0)}{\Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)} = \frac{N_{\text{cor}}(\Lambda_b^0 \rightarrow \psi(\mu^+\mu^-)\Lambda^0)}{N_{\text{cor}}(\Lambda_b^0 \rightarrow J/\psi(\mu^+\mu^-)\Lambda^0)} \cdot \frac{\mathcal{B}(J/\psi \rightarrow \ell^+\ell^-)}{\mathcal{B}(\psi(2S) \rightarrow \ell^+\ell^-)}$$

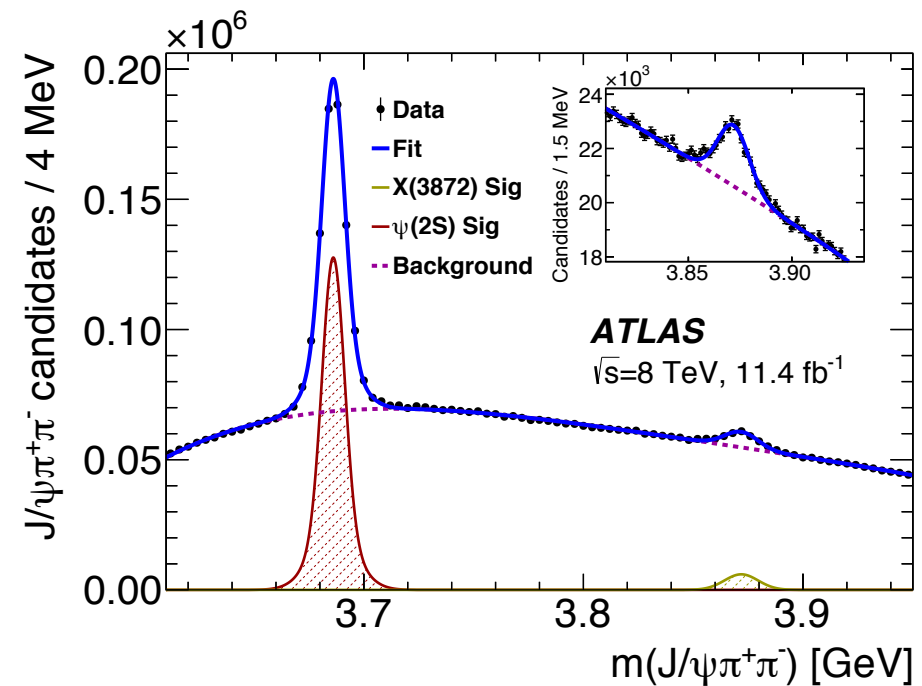
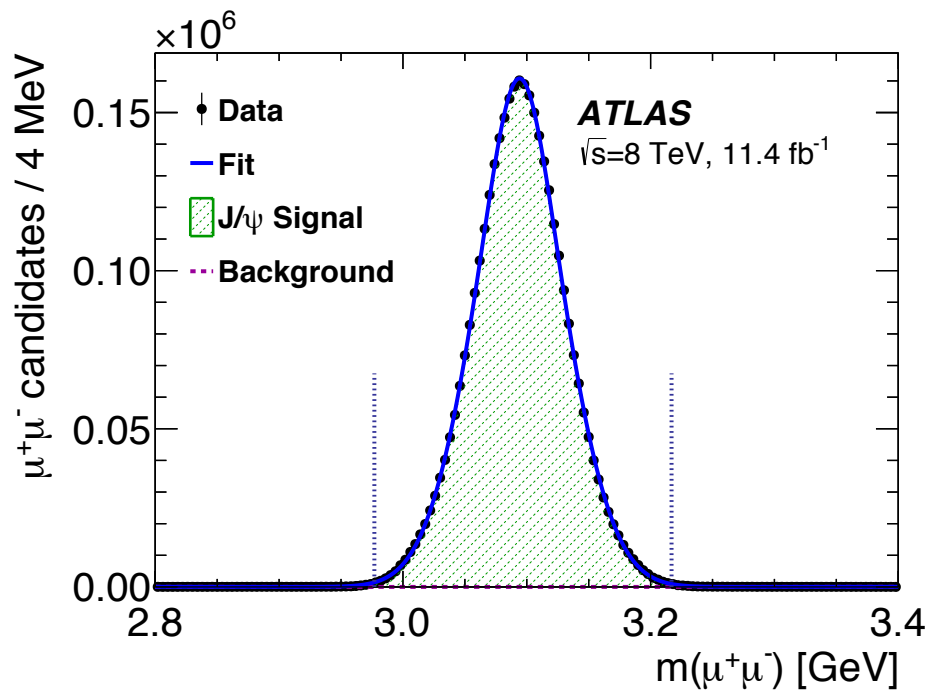
- Where the N_{cor} is the MC-corrected yield.
- Results in:

$$\frac{\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0)}{\Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)} = 0.501 \pm 0.033(\text{stat}) \pm 0.016(\text{syst}) \pm 0.011(\mathcal{B})$$

- Which is lower than the theoretical prediction of 0.8 ± 0.1 from PRD **88** (2013) 114018 and PRD **92** (2015) 114008.

Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (1)

- [JHEP01 \(2017\) 117](#)
- Decay mode: $J/\psi \pi^+ \pi^-$, 11.4 fb⁻¹ of 8 TeV ATLAS data.
- Rapidity range $|y| < 0.75$, p_T range of $J/\psi \pi^+ \pi^- = (10 - 70)$ GeV.
- MC simulation is used for studies of selection and reconstruction efficiencies.

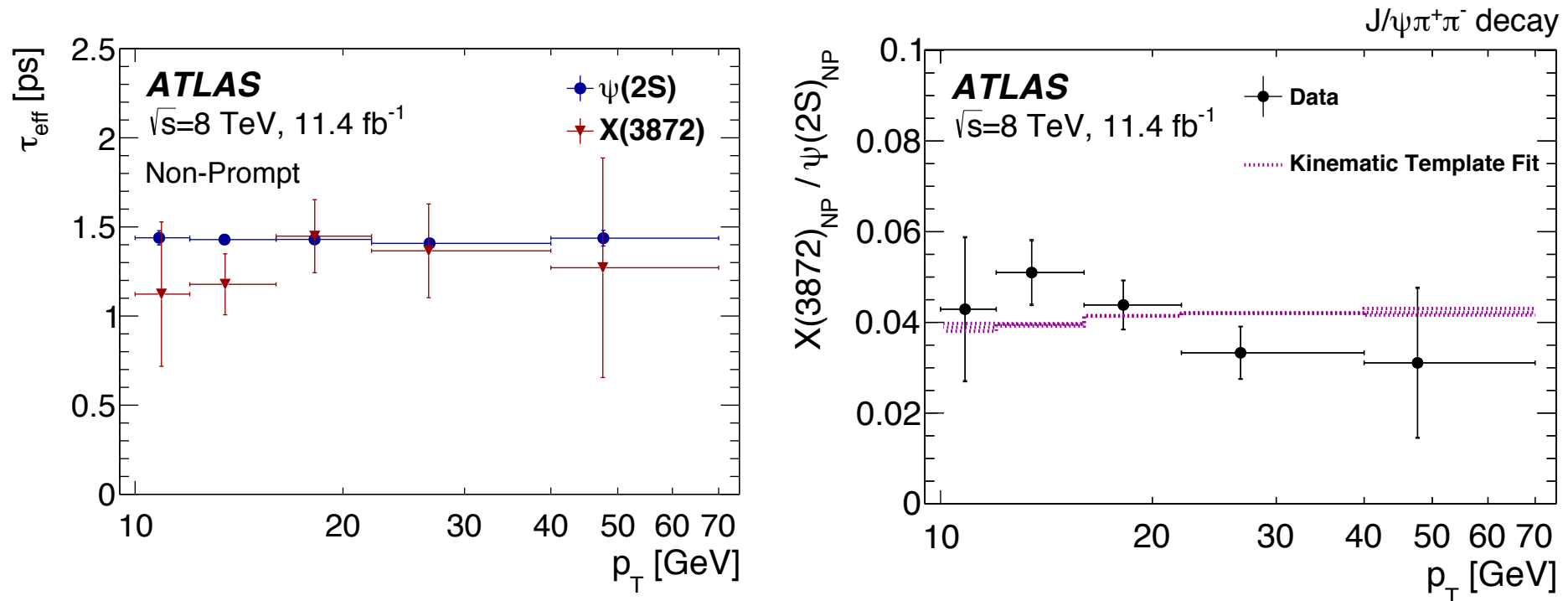


Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (2)

- The production cross sections of the $\psi(2S)$ and $X(3872)$ states are measured in five bins of $J/\psi\pi^+\pi^- p_T$, with bin boundaries
 - (10, 12, 16, 22, 40, 70) GeV.
- The cross sections measured are obtained under the assumption of
 - no spin alignment, but
 - appropriate sets of correction factors for a number of extreme spin alignment scenarios are calculated and presented in the Appendix of the corresponding Conference Note.
- In order to separate prompt production of the $\psi(2S)$ and $X(3872)$ states and the non-prompt production from the decays of long-lived particles such as b-hadrons, the data sample in each p_T bin was divided into intervals of pseudo-proper lifetime $\tau=L_{xy}m/p_T$.
- Four intervals of $\tau(J/\psi\pi^+\pi^-)$ were defined:
 - $-0.3 \text{ ps} < \tau < 0.025 \text{ ps}$
 - $0.025 \text{ ps} < \tau < 0.3 \text{ ps}$
 - $0.3 \text{ ps} < \tau < 1.5 \text{ ps}$
 - $1.5 \text{ ps} < \tau < 15.0 \text{ ps}$

Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (3)

- Measured effective pseudo-proper lifetimes for non-prompt $X(3872)$ and $\psi(2S)$, and the ratio of non-prompt $X(3872)$ and $\psi(2S)$ production.



- Kinematic template fit was calculated as a ratio of the simulated p_T distributions of non-prompt $X(3872)$ and non-prompt $\psi(2S)$, assuming that the same mix of the parent b-hadrons contributes for both signals.
- The shape of the template reflects the kinematics of the decay of a b-hadron into $\psi(2S)$ or $X(3872)$, with the width of the band showing the range of variation for extreme values of the invariant mass of the recoiling hadronic system.

Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (4)

- The fit of the measured ratio to that template allows to determine the ratio of the average branching fractions:

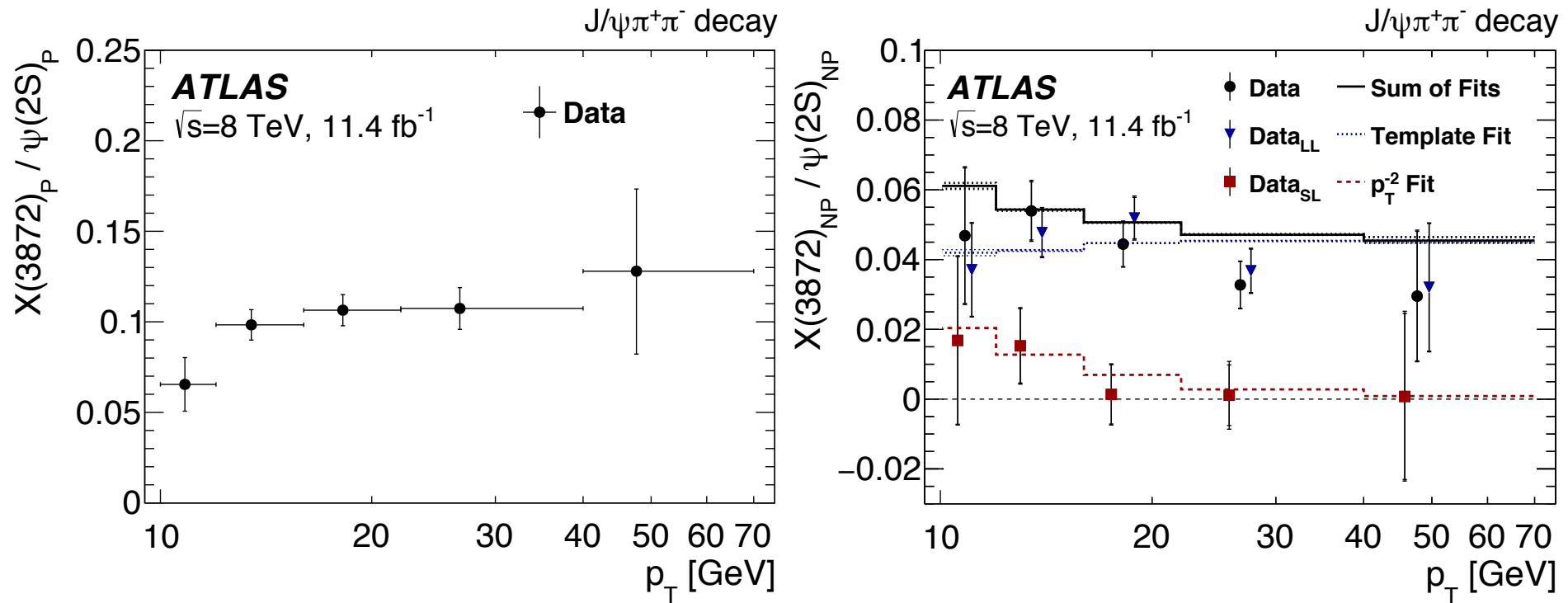
$$R_B^{1L} = \frac{Br(B \rightarrow X(3872))Br(X(3872) \rightarrow J/\psi\pi^+\pi^-)}{Br(B \rightarrow \psi(2S))Br(\psi(2S) \rightarrow J/\psi\pi^+\pi^-)} = (3.95 \pm 0.32(\text{stat}) \pm 0.08(\text{sys}))\%$$

- The somewhat falling trend shown on the right figure on the previous slide does not completely agree with the shape of the template, also **possibly suggesting the presence of an additional contribution** to the non-prompt $X(3872)$ yield in the low- p_T bins.
- The short-lived component is understood to be due to B_c decays, and the value of the pseudo-proper-lifetime parameter of the short-lived component is based on the value expected for B_c .
- An alternative fit model hence was implemented in the analysis, which allows for **two non-prompt contributions** with distinctly different effective lifetimes.

$$R_B^{2L} = \frac{Br(B \rightarrow X(3872))Br(X(3872) \rightarrow J/\psi\pi^+\pi^-)}{Br(B \rightarrow \psi(2S))Br(\psi(2S) \rightarrow J/\psi\pi^+\pi^-)} = (3.57 \pm 0.33(\text{stat}) \pm 0.11(\text{sys}))\%$$

Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (5)

- Ratio of cross section times branching fraction between $X(3872)$ and $\psi(2S)$ for prompt (left) and non-prompt (right) production. For the non-prompt production, the total ratio is separated into short-lived and long-lived components for the $X(3872)$.



Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (6)

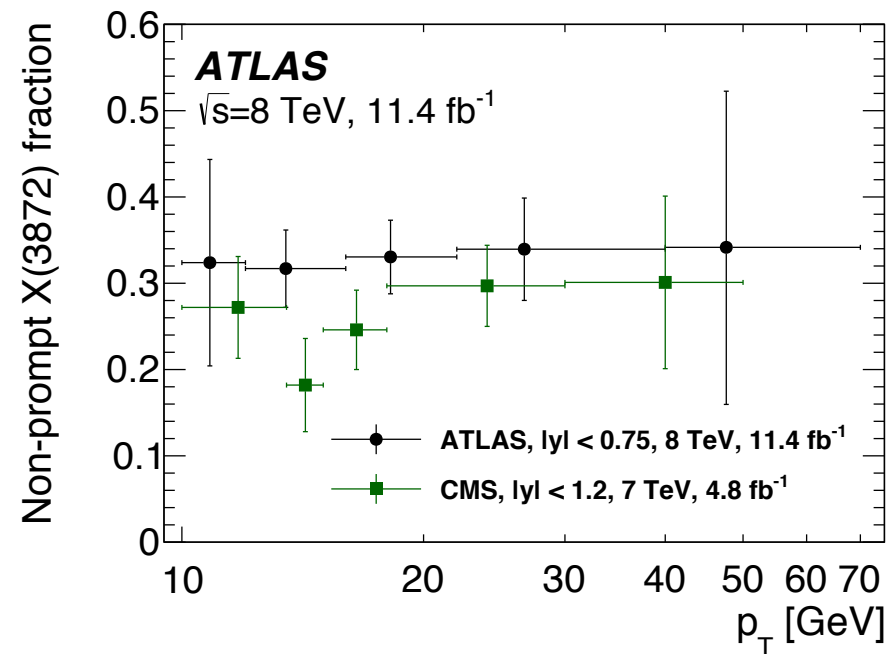
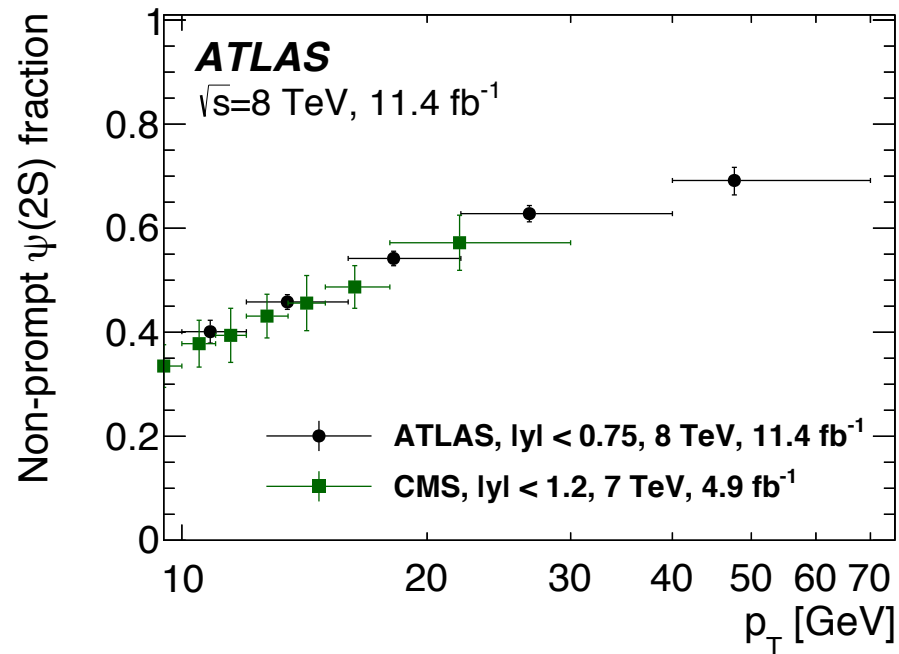
- The above is used to determine the fraction of non-prompt $X(3872)$ from short-lived sources, integrated over the p_T range ($p_T > 10$ GeV) covered in the analysis:

$$\frac{\sigma(pp \rightarrow B_c) Br(B_c \rightarrow X(3872))}{\sigma(pp \rightarrow \text{non-prompt } X(3872))} = (25 \pm 13(\text{stat}) \pm 2(\text{sys}) \pm 5(\text{spin}))\%$$

- The invariant mass distributions of the di-pion system in $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$ and $X(3872) \rightarrow J/\psi\pi^+\pi^-$ decays are also measured. The results do not favour a phase space distribution in either decay, and favour $X(3872) \rightarrow J/\psi\rho^0$.

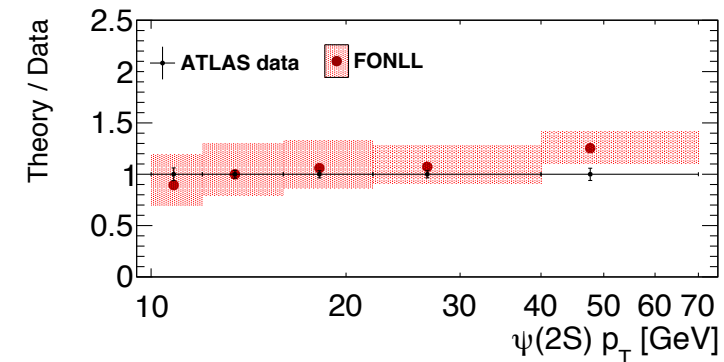
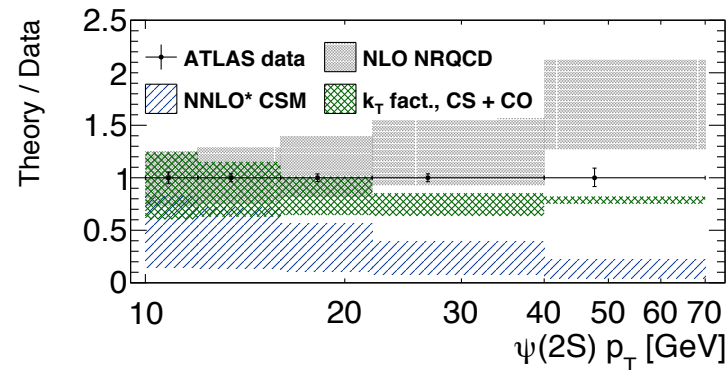
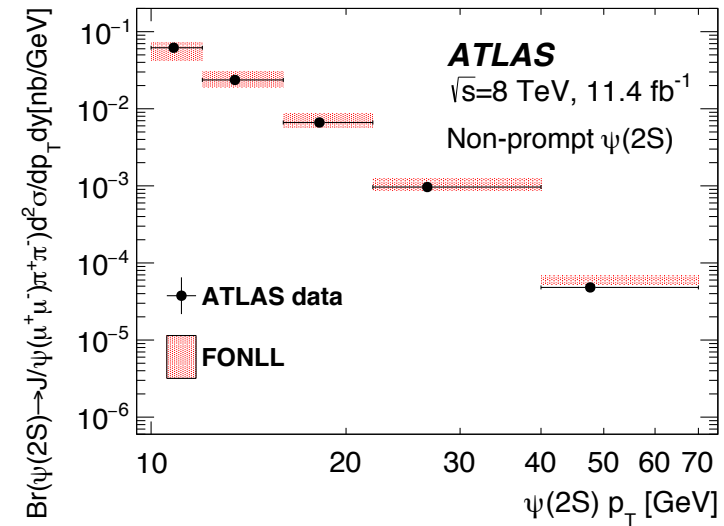
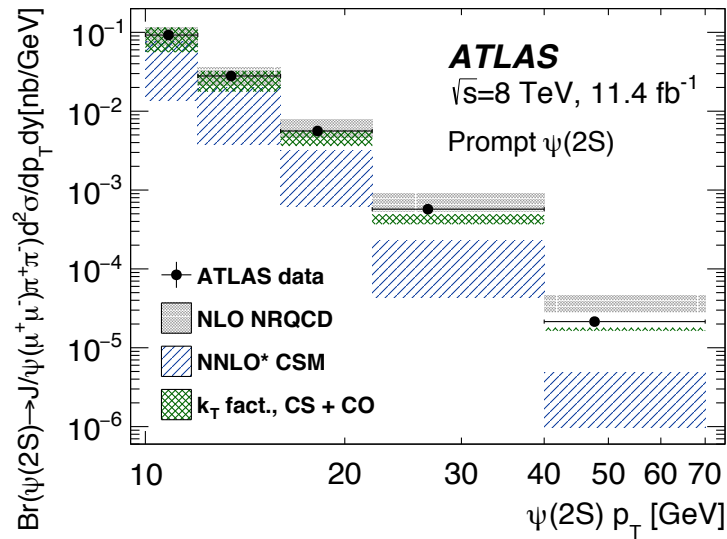
Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (7)

- Measured non-prompt fractions for $\psi(2S)$ (left) and $X(3872)$ (right) production, compared to CMS results at $\sqrt{s} = 7$ TeV.



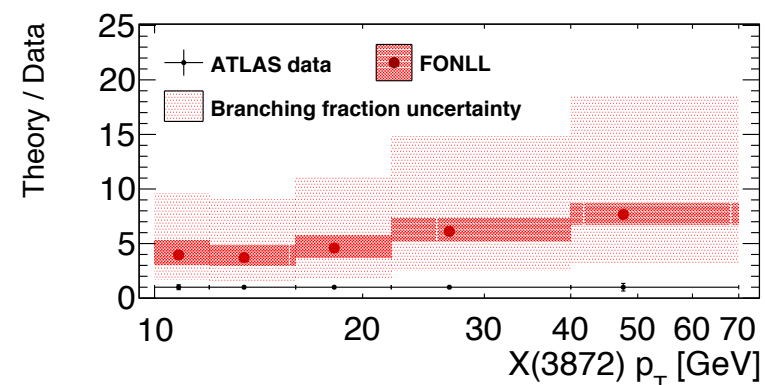
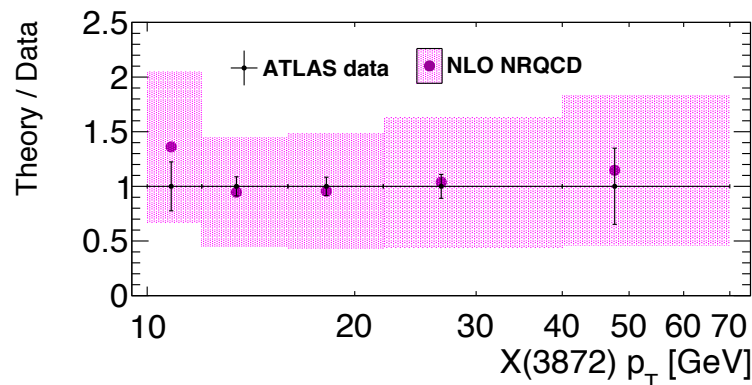
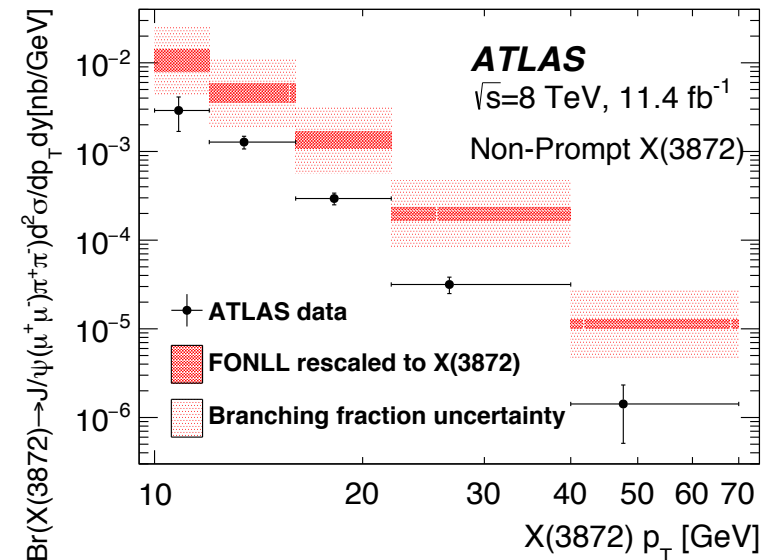
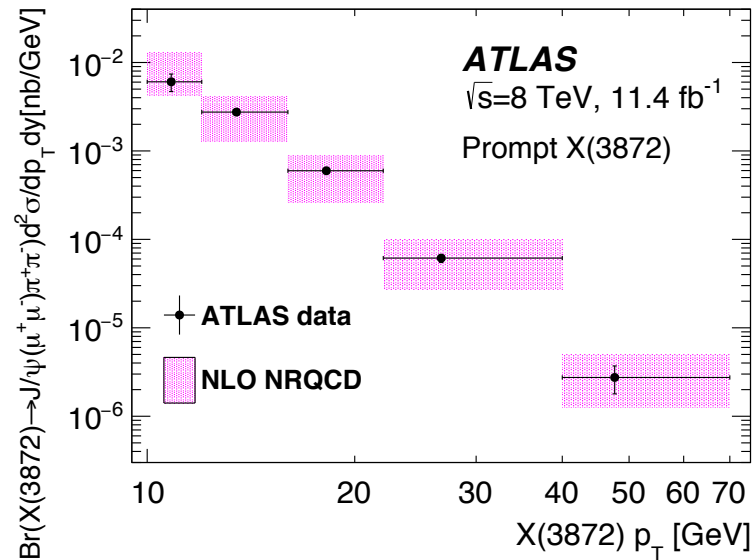
Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (8)

- The measured differential cross section (times the product of the relevant branching fractions) for prompt production of $\psi(2S)$ is presented on the left, and the non-prompt production is presented on the right. Compared to various theoretical models.



Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (9)

- The measured differential cross section (times the product of the relevant branching fractions) for prompt production of $X(3872)$ is presented on the left, and the non-prompt production is presented on the right. The $X(3872)$ is modeled as a mixture of a $\chi_{c1}(2P)$ and a $\bar{D}^0 D^{*0}$ molecular state.



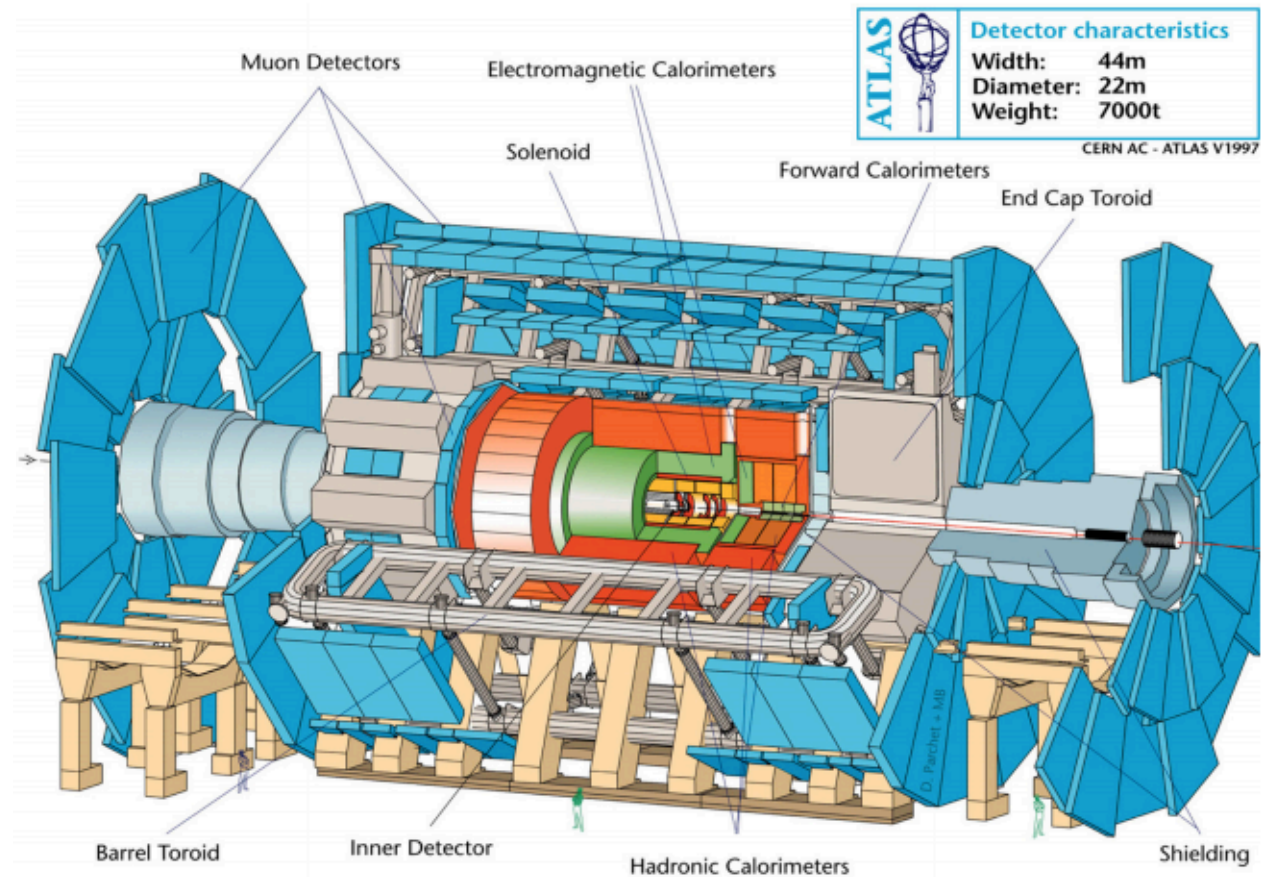
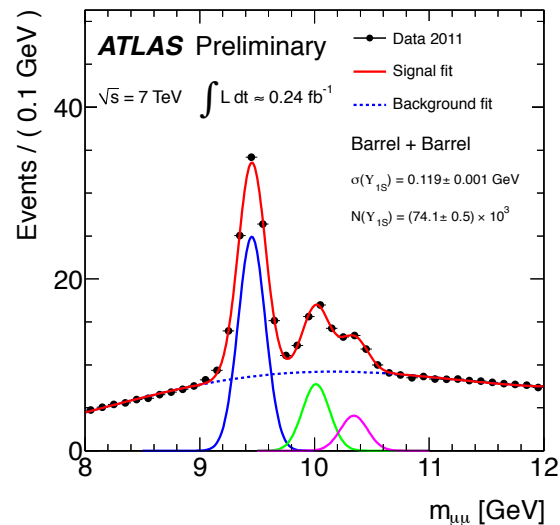
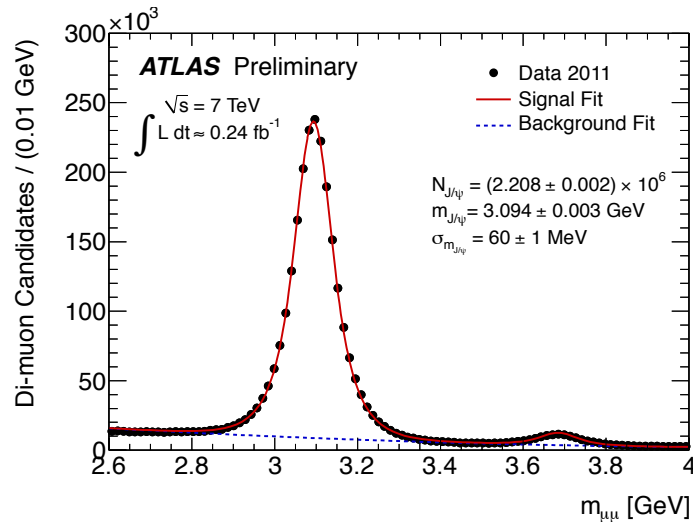
Conclusions

- $B_c(2S)$ has been observed for the first time. Mass and decay mode are consistent with theoretical expectations.
- $B_c^+ \rightarrow J/\psi D_s^+$ and $B_c^+ \rightarrow J/\psi D_s^{*+}$ decays have been studied. Four relative branching fractions were measured, theoretical predictions are within 2σ .
- X_b has been sought in $X_b \rightarrow Y(1S)\pi\pi$. No signal has been observed and upper limits have been set.
- Decay $\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0$ has been observed for the first time. Measured ratio $\Gamma(\Lambda_b^0 \rightarrow \psi(2S)\Lambda^0)/\Gamma(\Lambda_b^0 \rightarrow J/\psi\Lambda^0)$ exceeds the only available theoretical prediction.
- $X(3872) \rightarrow J/\psi\pi^+\pi^-$ decay has been studied. It was found that the molecular tetraquark component is not needed in description of the prompt production. The indication of enhanced contribution of B_c to the non-prompt production has been revealed.

BACKUP

ATLAS detector

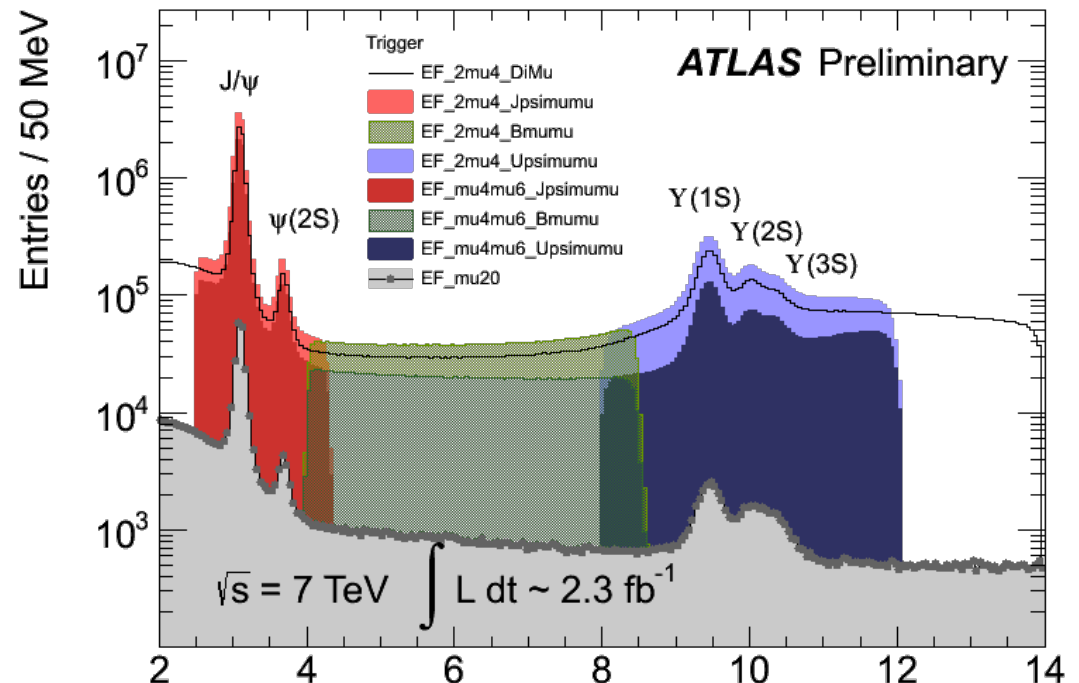
- Subsystems essential for B-physics: Inner detector and Muon spectrometer.
- Inner detector: tracking, momentum and vertexing, $|\eta| < 2.5$, d_0 resolution $\sim 10\mu\text{m}$.
- Muon spectrometer: trigger and muon identification, $|\eta| < 2.7$.
- J/ψ mass resolution: 60 ± 1 MeV, $\Upsilon(1S)$: 119 ± 1 MeV.



Trigger and datasets

- B-physics starts with single or di-muon triggers with various thresholds:

- $p_T(\mu) > 6$ GeV
- $p_T(\mu) > 18$ GeV
- $p_T(\mu_1) > 4$ GeV & $p_T(\mu_2) > 4$ GeV
- $p_T(\mu_1) > 6$ GeV & $p_T(\mu_2) > 4$ GeV
- $p_T(\mu_1) > 6$ GeV & $p_T(\mu_2) > 6$ GeV



- Di-muon mass range: $m(\mu\mu) \in [2.5; 4.3]$ GeV (final states containing J/ψ) and $m(\mu\mu) \in [4.0; 8.5]$ GeV (B to μ transitions).
- No displaced vertex selection requirements: advantage for lifetime measurements.
- Datasets: 4.9 fb^{-1} @ 7 TeV and 20.6 fb^{-1} @ 8 TeV.

Production measurements of $\psi(2S)$ and $X(3872)$ at ATLAS (10)

- (left) Normalised differential decay width of $\psi(2S) \rightarrow J/\psi(\rightarrow\mu^+\mu^-)\pi^+\pi^-$ in bins of dipion invariant mass over the range $0.280 \text{ GeV} < m(\pi\pi) < 0.595 \text{ GeV}$, fitted with the Voloshin–Zakharov model. Also shown is the normalised $m(\pi\pi)$ phase-space distribution (red shaded histogram).
- (right) Normalised differential decay width of $X(3872) \rightarrow J/\psi(\rightarrow\mu^+\mu^-)\pi^+\pi^-$ in bins of dipion invariant mass over the range $0.28\text{GeV} < m(\pi\pi) < 0.79\text{GeV}$. Also shown is the MC prediction for the decay $X(3872) \rightarrow J/\psi(\rightarrow\mu^+\mu^-)\rho^0(\rightarrow \pi^+\pi^-)$ (blue histogram) and the normalised distribution of $m(\pi\pi)$ phase-space (red shaded histogram).

