



Heavy flavour baryons

Mat Charles (UPMC/LPNHE)

QCD spectroscopy and exotic hadrons session,
LHCb implications workshop, 9 Nov 2017

What this talk is about

- I'll cover:
 - **Masses and lifetimes** of singly heavy **ground states** (briefly)
 - **Resonances with one heavy quark**, including
 - Low-lying excited Ξ_b states
 - $\Omega_c \rightarrow \Xi_c$ K resonances
 - b-baryons as a lab for c-baryon spectroscopy
 - Ξ_{cc} past, present & future
 - Ξ_b^0 **oscillation** search
- Other LHCb talks will cover:
 - **Exotic hadrons** (Sebastian Neubert) including pentaquarks
 - **Heavy flavour production** (Michael Winn)
 - **Asymmetries in charm baryon production and decay** (Jolanta Brodzicka)
 - **bc-hadrons at LHCb, including Ξ_{bc}** (Daria Savrina)
 - Heavy baryon decays as probes for CKM physics, $b \rightarrow sl^+l^-$ anomalies, etc (e.g. **Elena Graverini**, **Daniel O'Hanlon**, ...)

**DOON'T
PANIC**

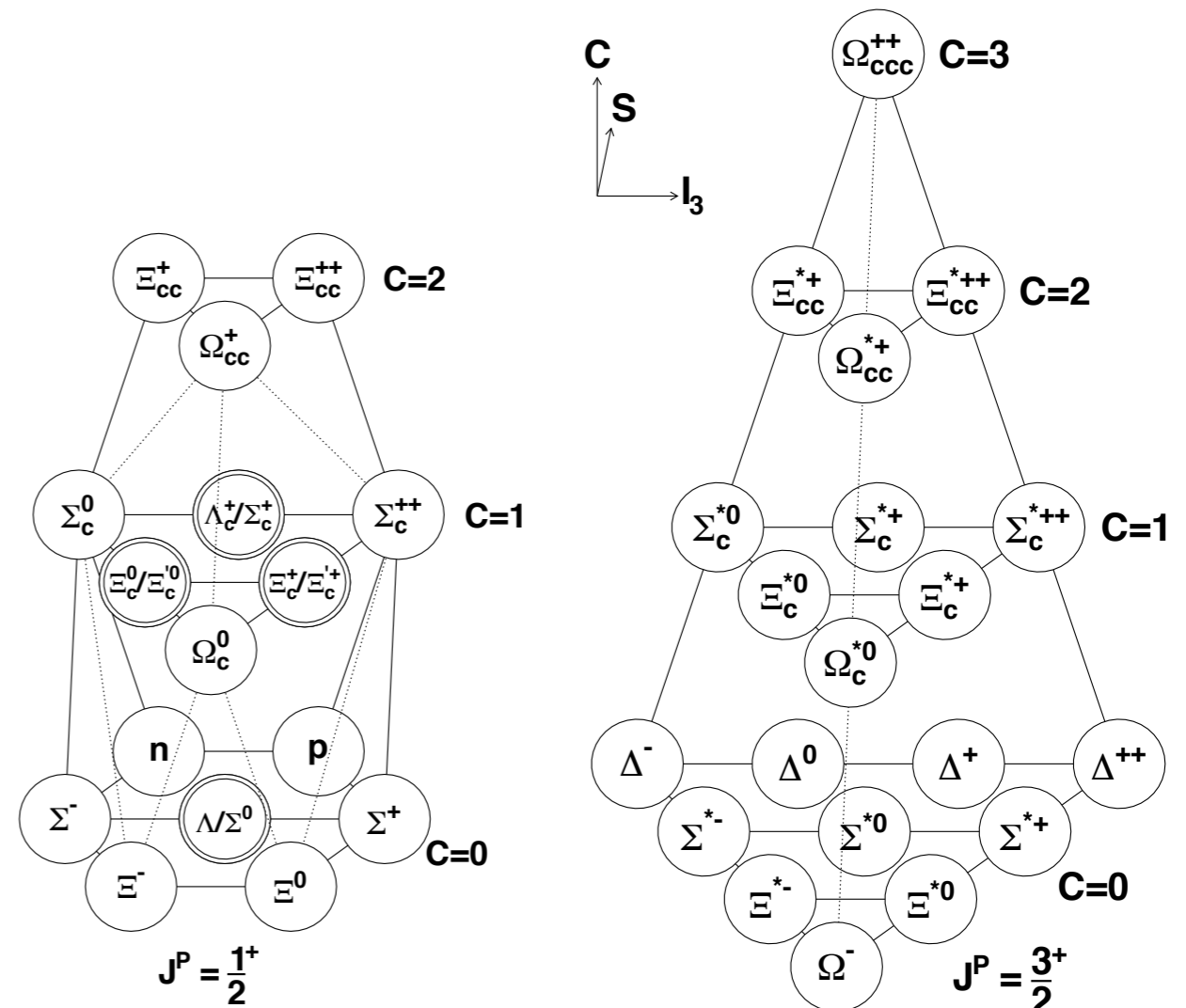
Crib sheet

- Baryon \Rightarrow 3 valence quarks
- If **all three are u or d**:
 - Isospin $1/2 \Rightarrow$ n or p
 - Isospin $3/2 \Rightarrow \Delta$
- If **two are u or d**:
 - Isospin 0 $\Rightarrow \Lambda$
 - Isospin 1 $\Rightarrow \Sigma$
- If **one is u or d**:
 - Isospin $1/2 \Rightarrow \Xi$
- If **none are u or d**:
 - Isospin 0 $\Rightarrow \Omega$
- By default, non-{u/d} quarks are s.
Add subscripts to replace them, e.g.
 - $\Xi^0 = ssu$
 - $\Xi_c^+ = csu$
 - $\Xi_{cc}^{++} = ccu$

"Ground states" ambiguous:

- 1) Lowest states that decay weakly
- 2) All states that are neither orbitally nor radially excited

For second meaning, here are the ground states (with beauty 0):



Why do we do spectroscopy?



In roughly ascending order:

- It's good clean fun[^]W physics
 - ... and sometimes you actually get to see a signal
- It provides **inputs** to pheno calculations in the area
- Empowers **searches for New Physics**
 - Necessary inputs, e.g. charm resonances for $B \rightarrow D^{(*)} \tau \nu$
 - Opens new paths (time until first CPV search with Ξ_{cc} ?)
- Sometimes you find something **unexpected**
 - See: pentaquarks

[Zoltan's talk on Thursday](#)

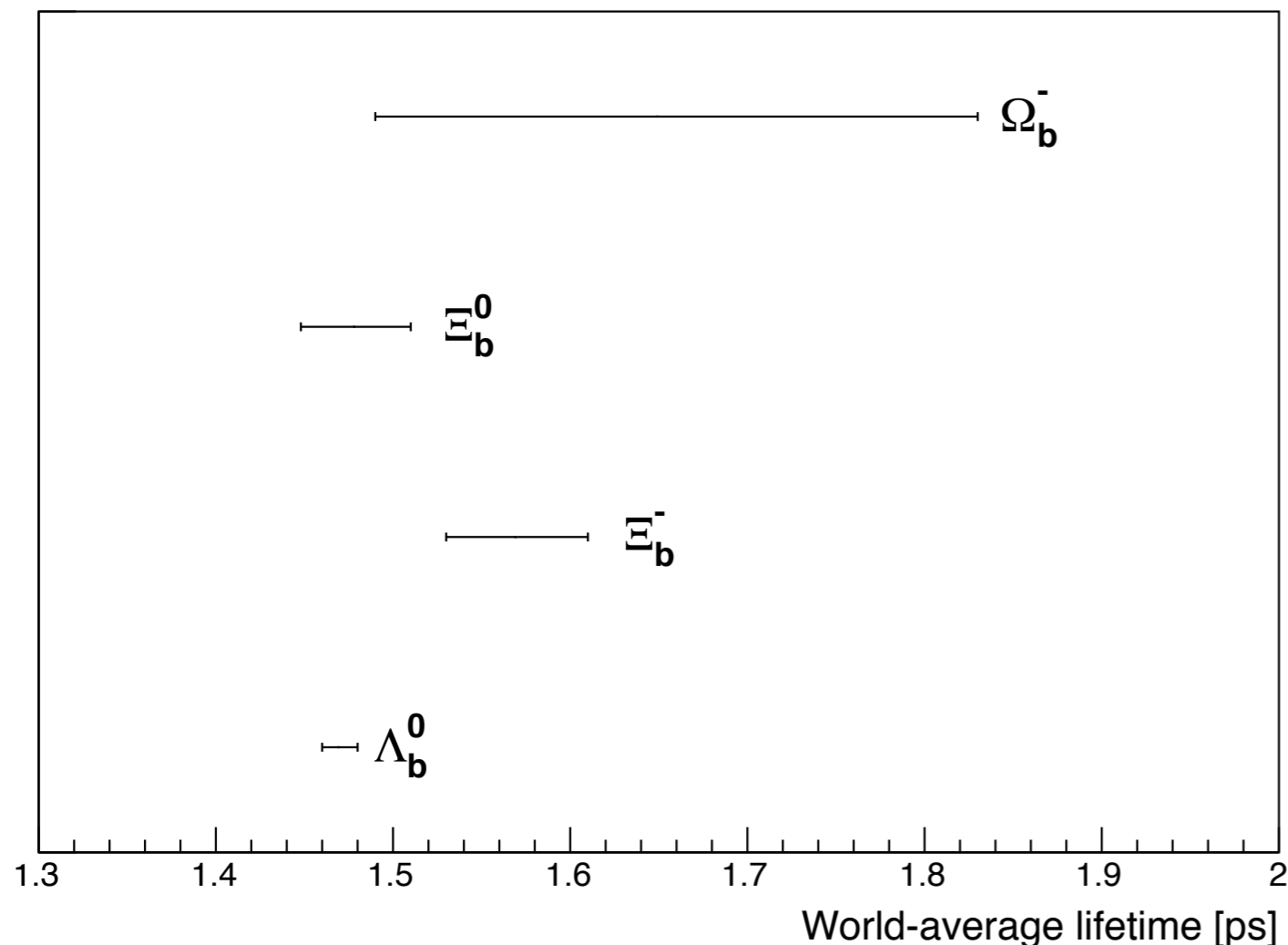
Weakly decaying b-baryons

- Skipping details, but:
- Precise measurements of masses and lifetimes, often using ratio technique.

$$M(\Lambda_b^0) = 5619.65 \pm 0.17 \pm 0.17 \text{ MeV}/c^2$$

$$M(\Xi_b^0) - M(\Lambda_b^0) = 172.44 \pm 0.39 \pm 0.17 \text{ MeV}/c^2$$

$$M(\Xi_b^-) - M(\Lambda_b^0) = 178.36 \pm 0.46 \pm 0.16 \text{ MeV}/c^2$$



See backups for Ω_b^- mass info.

[LHCb-PAPER-2015-060](#) ; [arXiv:1603.06961](#)

[LHCb-PAPER-2014-048](#) ; [arXiv:1409.8568](#)

[LHCb-PAPER-2014-021](#) ; [arXiv:1405.7223](#)

[LHCb-PAPER-2016-008](#) ; [arXiv:1604.01412](#)

[LHCb-PAPER-2012-048](#) ; [arXiv:1302.1072](#)

[LHCb-PAPER-2014-003](#) ; [arXiv:1402.6242](#)

[LHCb-PAPER-2014-048](#) ; [arXiv:1409.8568](#)

[LHCb-PAPER-2014-021](#) ; [arXiv:1405.7223](#)

[LHCb-PAPER-2016-008](#) ; [arXiv:1604.01412](#)

[LHCb-PAPER-2014-010](#) ; [arXiv:1405.1543](#)

LHCb results: weakly decaying c-baryons

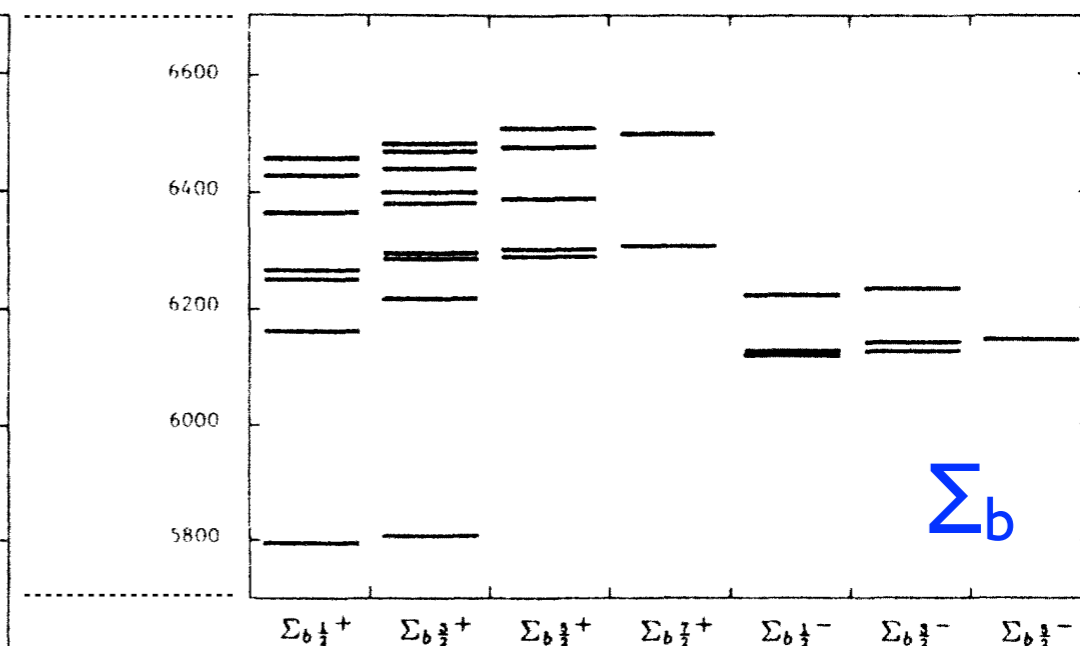
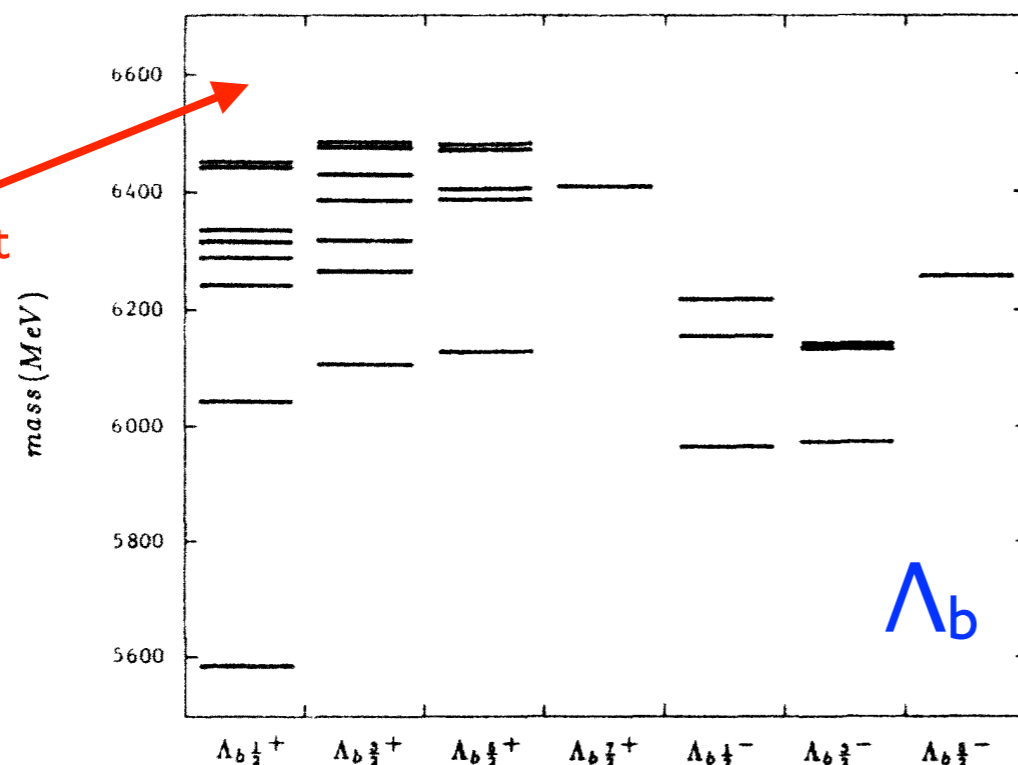


- True, results from B-factories for masses already exist...
- ... but surely LHCb can add something, esp. on lifetimes.

Singly heavy resonances

- The **low-lying states** are pretty easy to understand.
 - ... especially the "ground states" ($n=L=0$). Even an experimentalist can calculate their masses!
- **Higher excited states** are tougher, because:
 - They are **hard to see** (cross-sections for inclusive production and BFs for exclusive production usually drop rapidly)
 - There are many of them -- **density of states** gets large rapidly
 - They **mix!** States with same J^P and flavour quantum numbers can mix, even if internal numbers differ... and there are a lot of states.

NB states continue up here for $N>2$, just not plotted.



Theory predictions for $N=0,1,2$ from Capstick & Isgur, [Phys.Rev. D34 \(1986\) 280](#)

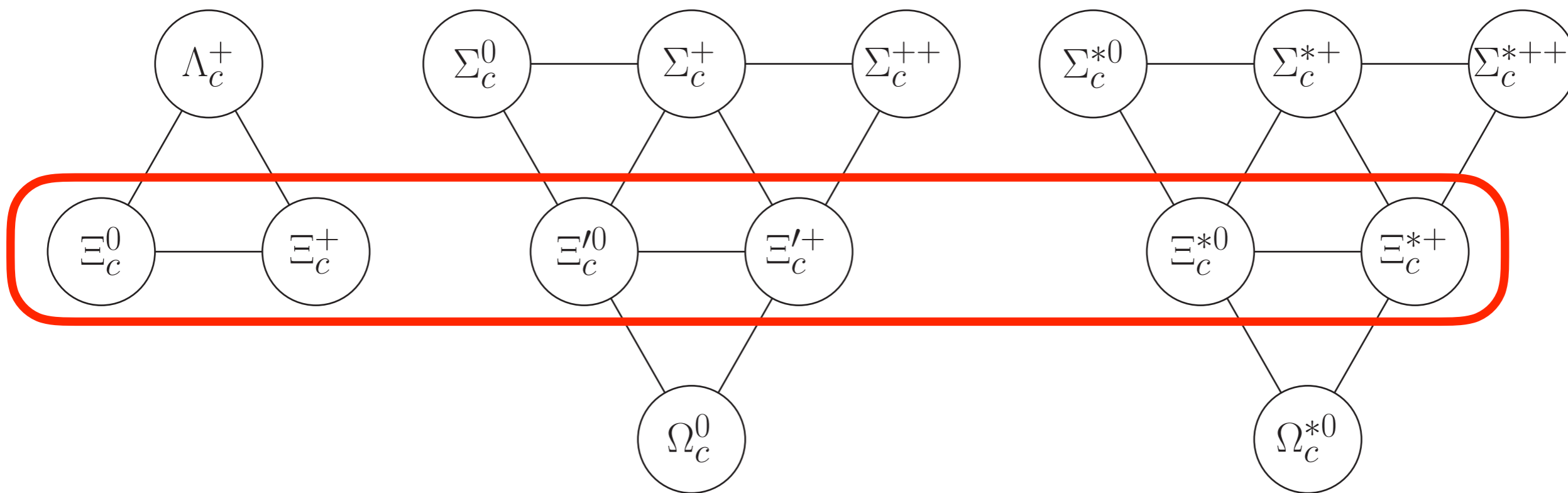
Warm-up: Ξ_b ground states

Diagrams show Ξ_c , but Ξ_b works [almost] the same way

$$j = 0, J^P = \frac{1}{2}^+$$

$$j = 1, J^P = \frac{1}{2}^+$$

$$j = 1, J^P = \frac{3}{2}^+$$

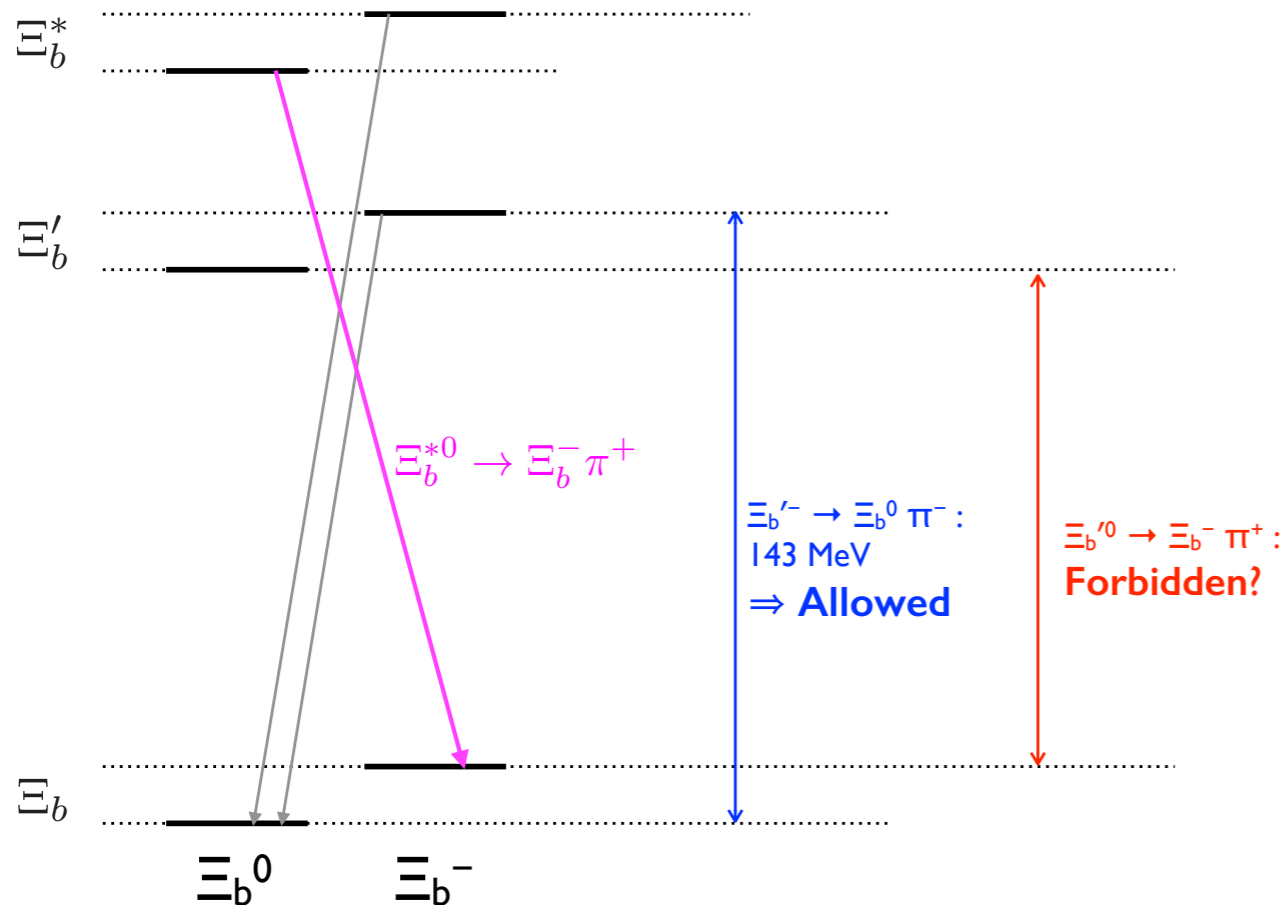
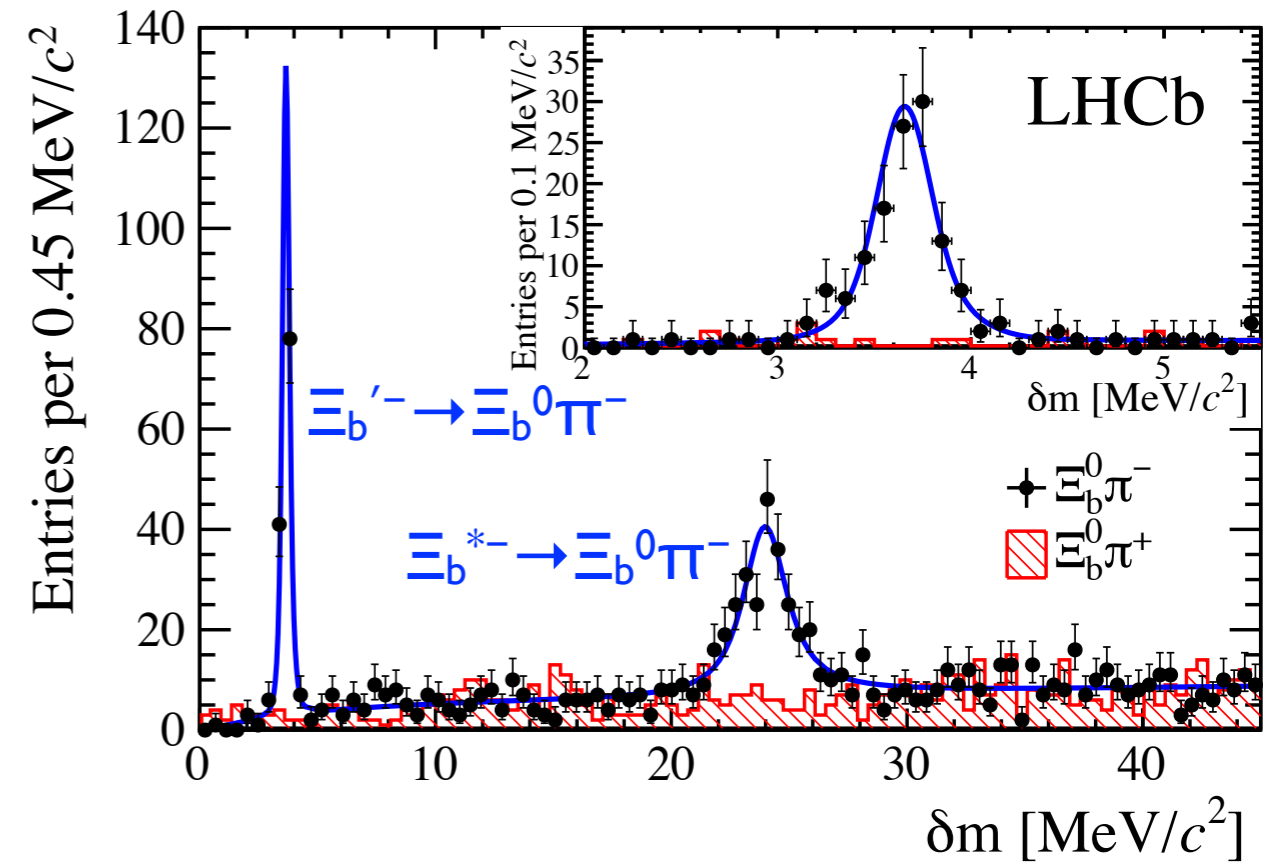
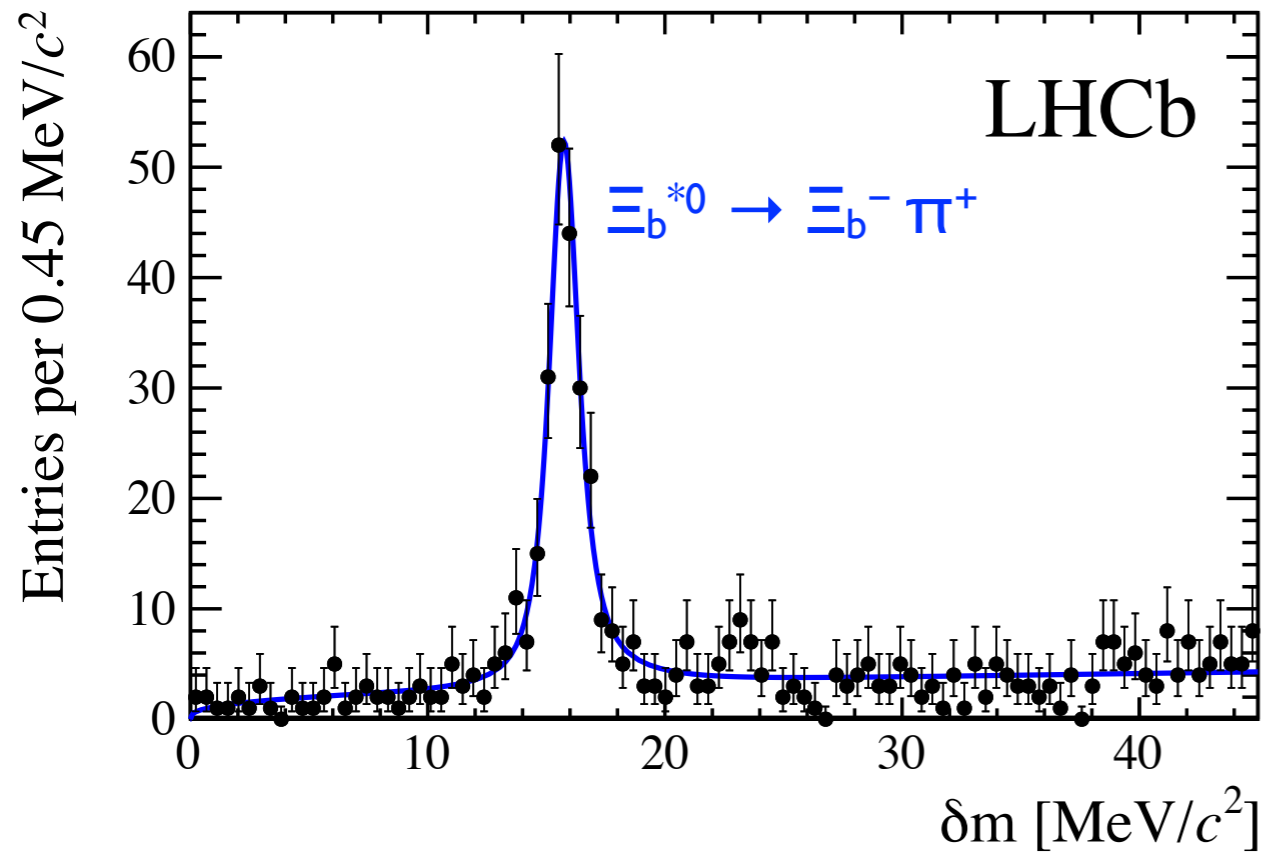


• If you put three quarks with different flavours together*, QM says there are not two but three ways to do it:

- $J = (\frac{1}{2} + \frac{1}{2}) + \frac{1}{2} = 0 + \frac{1}{2} = \frac{1}{2}$, mixed symmetry $\Rightarrow \Xi_b$
- $J = (\frac{1}{2} + \frac{1}{2}) + \frac{1}{2} = 1 + \frac{1}{2} = \frac{1}{2}$, mixed symmetry $\Rightarrow \Xi_b'$
- $J = (\frac{1}{2} + \frac{1}{2}) + \frac{1}{2} = 1 + \frac{1}{2} = \frac{3}{2}$, totally symmetric $\Rightarrow \Xi_b^*$

diquark

Ξ_b measurements at LHCb



5 of the 6 ground states now seen.
 The last one, $\Xi_b'^0$, is (presumably!)
 below threshold for $\Xi_b^- \pi^+$ and so
 hard to see at LHCb.

Ξ_b^{*0} first observed by CMS: [PRL 108, 252002 \(2012\)](#)

[LHCb-PAPER-2016-010](#) ; [arXiv:1604.03896](#)

[LHCb-PAPER-2014-061](#) ; [arXiv:1411.4849](#)

Ξ_b measurements at LHCb

- **Very precise** measurements of masses and widths
- ... especially for the state closest to threshold.

$$m(\Xi_b^{\prime-}) - m(\Xi_b^0) - m(\pi^-) = 3.653 \pm 0.018 \pm 0.006 \text{ MeV}/c^2,$$

$$m(\Xi_b^{*-}) - m(\Xi_b^0) - m(\pi^-) = 23.96 \pm 0.12 \pm 0.06 \text{ MeV}/c^2,$$

$$\Gamma(\Xi_b^{*-}) = 1.65 \pm 0.31 \pm 0.10 \text{ MeV},$$

$$\Gamma(\Xi_b^{\prime-}) < 0.08 \text{ MeV at } 95\% \text{ confidence level.}$$

$$m(\Xi_b^{*0}) - m(\Xi_b^-) - m(\pi^+) = 15.727 \pm 0.068 \text{ (stat)} \pm 0.023 \text{ (syst)} \text{ MeV}/c^2,$$

$$\Gamma(\Xi_b^{*0}) = 0.90 \pm 0.16 \text{ (stat)} \pm 0.08 \text{ (syst)} \text{ MeV}.$$

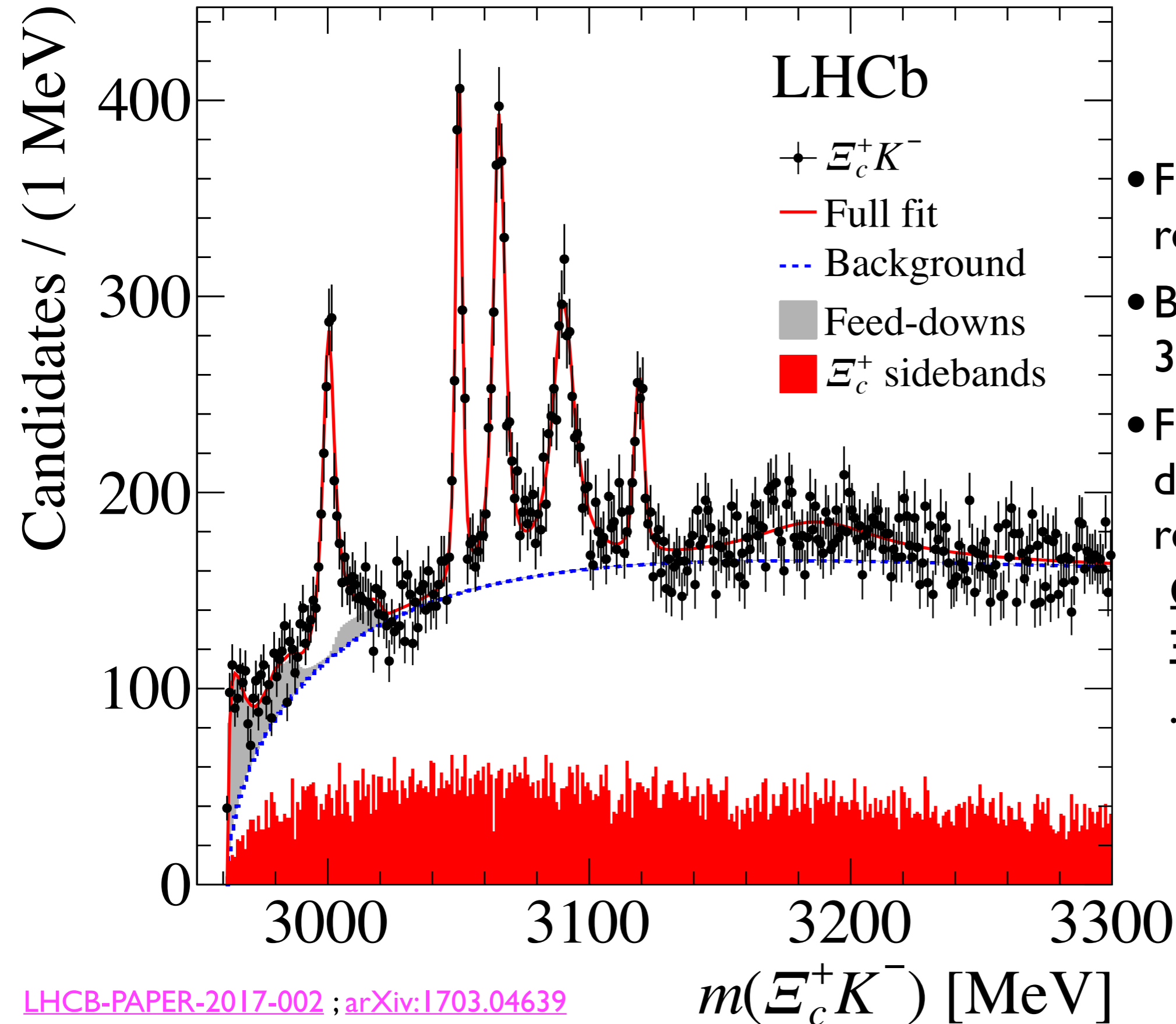
J consistent with expectations (but little information)

Ratios of $(\sigma \times \text{BF})$ also measured.

Low-lying excited $\Omega_c^0 \rightarrow \Xi_c^+ K^-$

- Switch now from a study of Ξ_b ground states to low-lying excited Ω_c states.
- Reminder: "excited" here means:
 - orbital excitation: $L > 0$, e.g. 1P, 1D, ... with parity = $(-1)^L$
 - radial excitation, e.g. 2S, 3S, ...
 - or both (2P, 3D, etc etc etc)
- Note that more obvious mode $\Omega_c^0 \pi^0$ violates isospin.

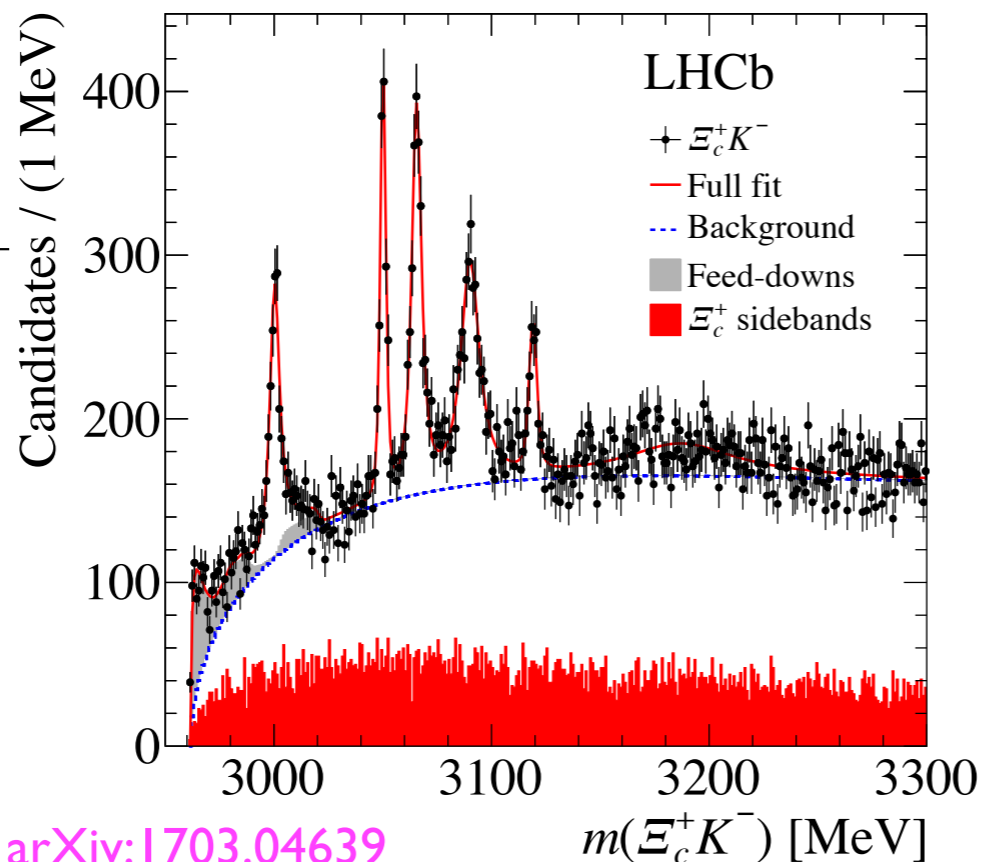
Low-lying excited $\Omega_c^0 \rightarrow \Xi_c^+ K^-$



- Five clear, narrow resonances
- Broad excess around 3190 MeV
- Fit also includes feed-down from same resonances:
 $\Omega_c^0 \rightarrow \Xi_c'^+ K^-$,
 $\Xi_c'^+ \rightarrow \Xi_c^+ \gamma$
 ... and photon missed

Low-lying excited $\Omega_c^0 \rightarrow \Xi_c^+ K^-$

Resonance	Mass (MeV)	Γ (MeV)	Yield	N_σ
$\Omega_c(3000)^0$	$3000.4 \pm 0.2 \pm 0.1^{+0.3}_{-0.5}$	$4.5 \pm 0.6 \pm 0.3$	$1300 \pm 100 \pm 80$	20.4
$\Omega_c(3050)^0$	$3050.2 \pm 0.1 \pm 0.1^{+0.3}_{-0.5}$	$0.8 \pm 0.2 \pm 0.1$ $< 1.2 \text{ MeV, 95\% CL}$	$970 \pm 60 \pm 20$	20.4
$\Omega_c(3066)^0$	$3065.6 \pm 0.1 \pm 0.3^{+0.3}_{-0.5}$	$3.5 \pm 0.4 \pm 0.2$	$1740 \pm 100 \pm 50$	23.9
$\Omega_c(3090)^0$	$3090.2 \pm 0.3 \pm 0.5^{+0.3}_{-0.5}$	$8.7 \pm 1.0 \pm 0.8$	$2000 \pm 140 \pm 130$	21.1
$\Omega_c(3119)^0$	$3119.1 \pm 0.3 \pm 0.9^{+0.3}_{-0.5}$	$1.1 \pm 0.8 \pm 0.4$ $< 2.6 \text{ MeV, 95\% CL}$	$480 \pm 70 \pm 30$	10.4
$\Omega_c(3188)^0$	$3188 \pm 5 \pm 13$	$60 \pm 15 \pm 11$	$1670 \pm 450 \pm 360$	
$\Omega_c(3066)_{\text{fd}}^0$			$700 \pm 40 \pm 140$	
$\Omega_c(3090)_{\text{fd}}^0$			$220 \pm 60 \pm 90$	
$\Omega_c(3119)_{\text{fd}}^0$			$190 \pm 70 \pm 20$	



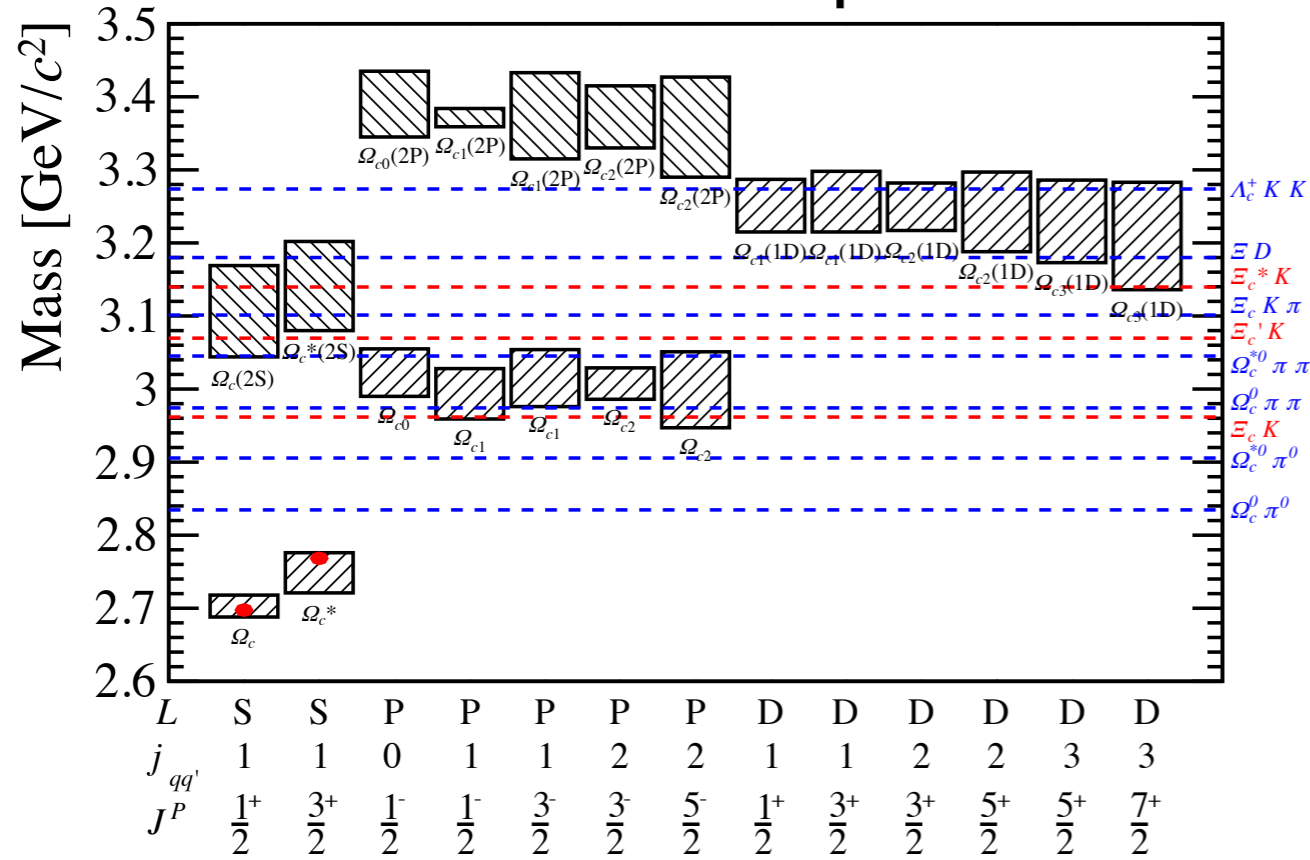
Two states are very narrow;
possible exotic interpretation
(see [talk](#) by Sebastian Neubert;
also [Tuesday's workshop](#)).

e.g. Montaña et al: [arXiv:1709.08737](#)

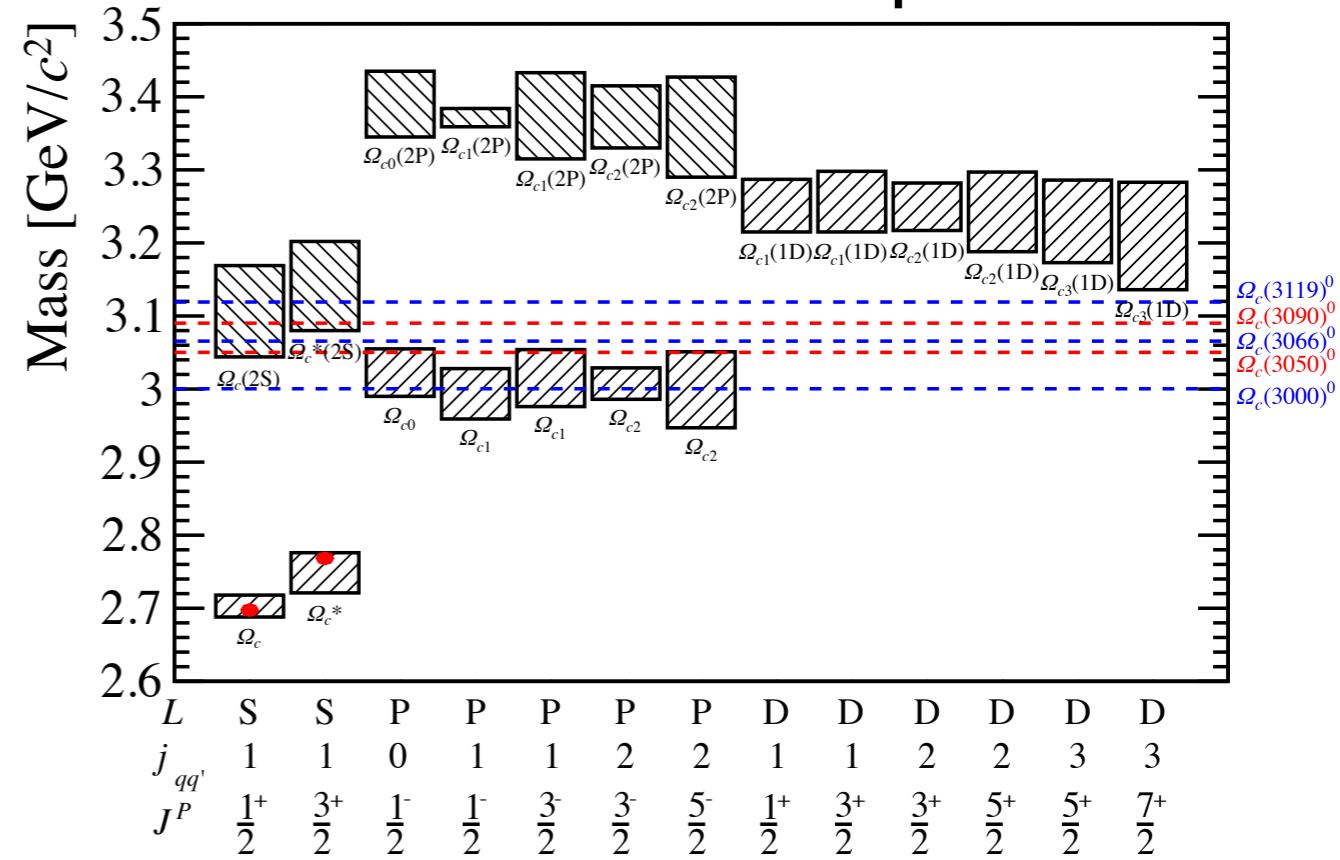
e.g. Debastiani et al: [arXiv:1710.04231](#)

Low-lying excited $\Omega_c^0 \rightarrow \Xi_c^+ K^-$

Thresholds vs some predictions



Observations vs some predictions

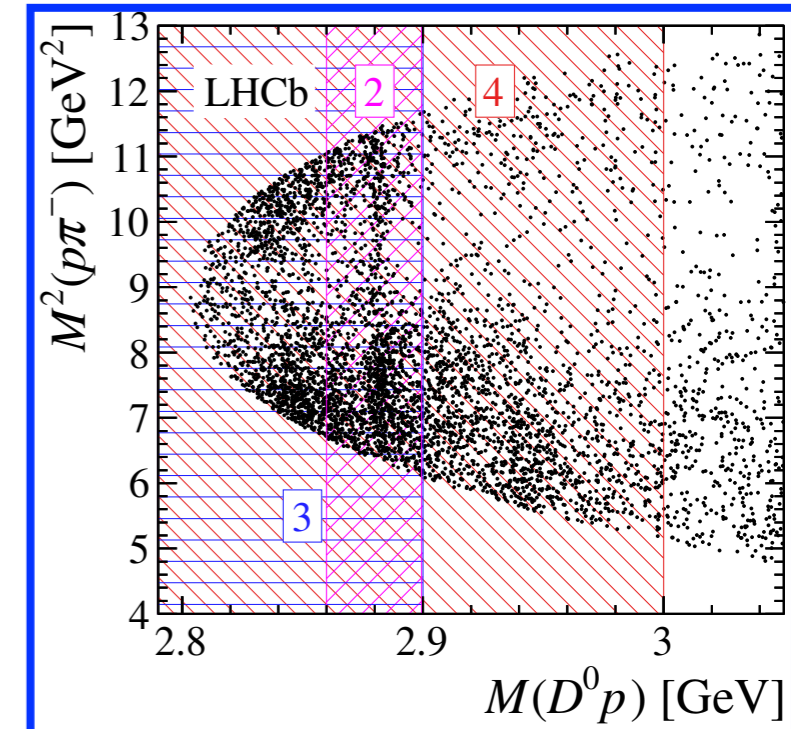
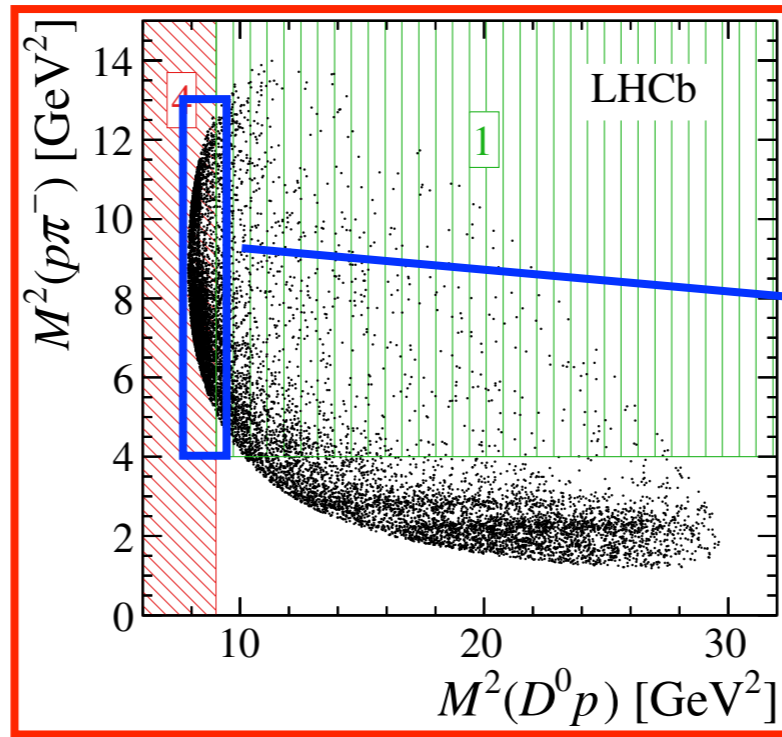
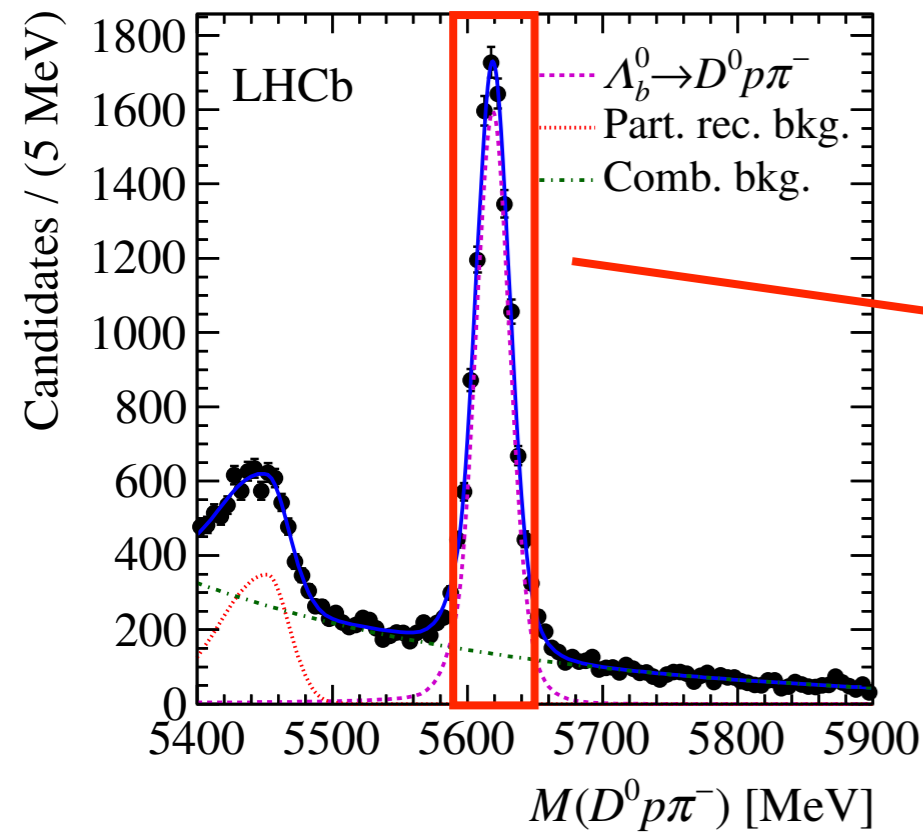


- Without measurements of quantum numbers, matching the five peaks to states is tricky
 - Obvious first guess: $\frac{1}{2}^-$, $\frac{1}{2}^-$, $\frac{3}{2}^-$, $\frac{3}{2}^-$, $\frac{5}{2}^-$
 - If first two states are below threshold, peaks could be $\frac{3}{2}^-$, $\frac{5}{2}^-$, $\frac{5}{2}^-$, $\frac{1}{2}^+$, $\frac{3}{2}^+$
 - Or exotic interpretation, or...
- Aside: $\Xi_c^{*+} K^-$ would be a nice mode for a follow-up search.

Exclusive decays

- Interpretation is difficult without knowing QNs
- So try to **work in a system where we can measure them!**
- Decays of charmed hadrons used to study light hadrons...
 - ... e.g. spin of Ω^- measured for first time in $\Xi_c^0 \rightarrow \Omega^- K^+$
(BABAR, [Phys.Rev.Lett.97:112001,2006](#))
- ... and likewise we **can use beauty decays to study charm.**
- Recent example: $(D^0 p)$ in $\Lambda_b^0 \rightarrow D^0 p \pi^-$
- I can't do justice to this detailed amplitude analysis in just a couple of slides, but let me give you a taste...

$D^0 p$ amplitude in $\Lambda_b^0 \rightarrow D^0 p \pi^-$



Caution: x axis changes from M^2 to M

• Series of amplitude fits

- ... beginning with easier bits (non-res region 1, $\Lambda_c(2880)$ -dominated region 2)
- ... then adding near-threshold resonances (region 3) and $\Lambda_c(2940)$ in region 4.
- N^* -dominated region not used in fit (simplifies life -- but lose a bit of info)
- Known $\Lambda_c(2880)$ and $\Lambda_c(2940)$ studied
 - $\Lambda_c(2880)$: $J = 5/2$ strongly preferred, consistent with Belle
 - $\Lambda_c(2940)$: $J^P = 3/2^-$ preferred, but other values not fully ruled out
- Threshold enhancement, consistent with $3/2^+$ resonance: $\Lambda_c(2860)$

$D^0 p$ amplitude in $\Lambda_b^0 \rightarrow D^0 p \pi^-$

$$m(\Lambda_c(2860)^+) = 2856.1_{-1.7}^{+2.0}(\text{stat}) \pm 0.5(\text{syst})_{-5.6}^{+1.1}(\text{model}) \text{ MeV},$$

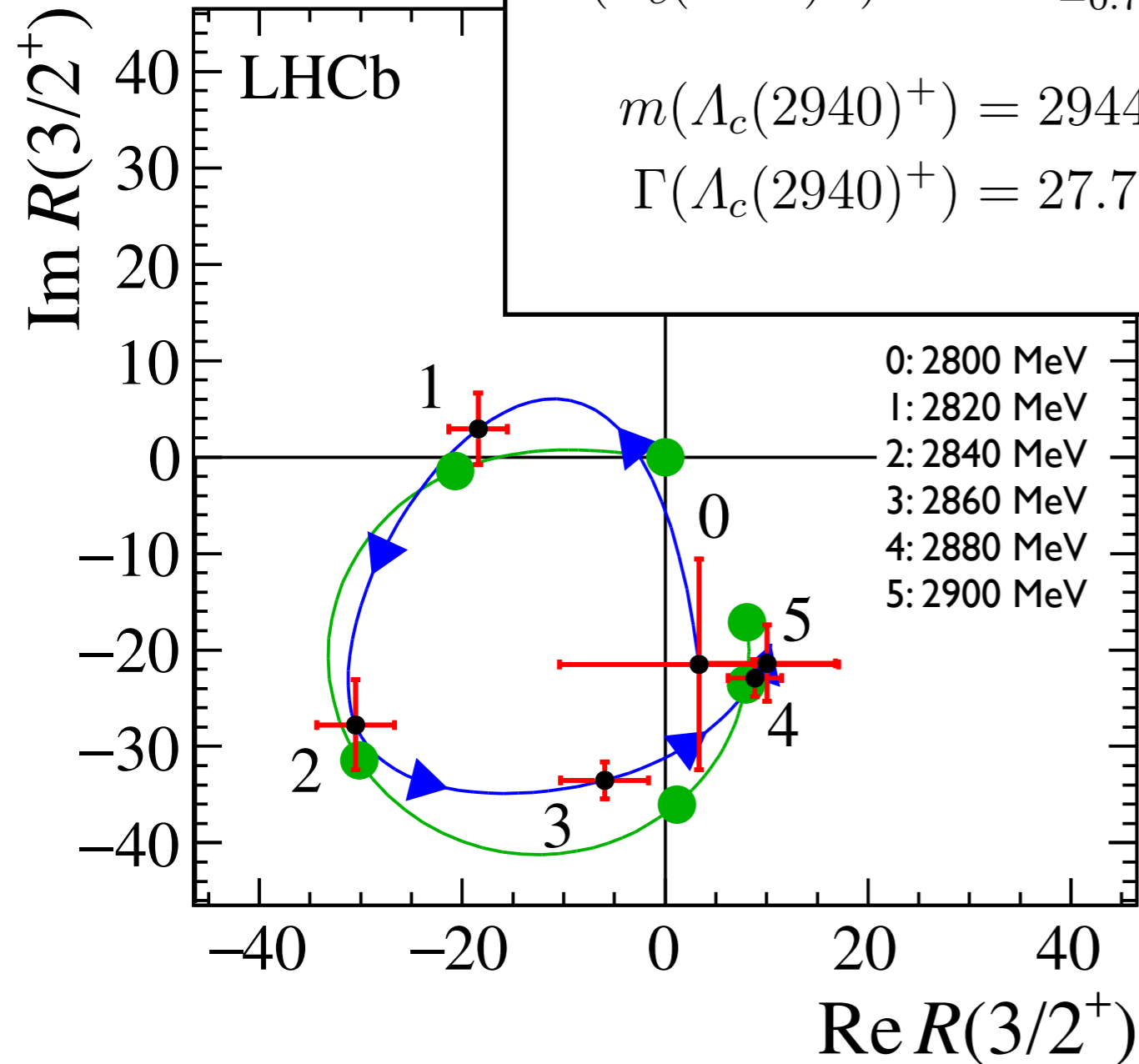
$$\Gamma(\Lambda_c(2860)^+) = 67.6_{-8.1}^{+10.1}(\text{stat}) \pm 1.4(\text{syst})_{-20.0}^{+5.9}(\text{model}) \text{ MeV}$$

$$m(\Lambda_c(2880)^+) = 2881.75 \pm 0.29(\text{stat}) \pm 0.07(\text{syst})_{-0.20}^{+0.14}(\text{model}) \text{ MeV},$$

$$\Gamma(\Lambda_c(2880)^+) = 5.43_{-0.71}^{+0.77}(\text{stat}) \pm 0.29(\text{syst})_{-0.00}^{+0.75}(\text{model}) \text{ MeV}.$$

$$m(\Lambda_c(2940)^+) = 2944.8_{-2.5}^{+3.5}(\text{stat}) \pm 0.4(\text{syst})_{-4.6}^{+0.1}(\text{model}) \text{ MeV},$$

$$\Gamma(\Lambda_c(2940)^+) = 27.7_{-6.0}^{+8.2}(\text{stat}) \pm 0.9(\text{syst})_{-10.4}^{+5.2}(\text{model}) \text{ MeV}.$$



Argand diagram for $J^P=3/2^+$ partial wave in the $\Lambda_c(2860)$ region.

- Known $\Lambda_c(2880)$ and $\Lambda_c(2940)$ studied
 - $\Lambda_c(2880)$: $J = 5/2$ strongly preferred, consistent with Belle
 - $\Lambda_c(2940)$: $J^P = 3/2^-$ preferred, but other values not fully ruled out
- $\Lambda_c(2860)$ consistent with $3/2^+$ resonance:

En 2002, en utilisant l'expérience FELEX de l'accélérateur « Fermilab », dans l'Illinois, des chercheurs avaient annoncé avoir trouvé une particule similaire.

Cependant, cette détection était juste en dessous du seuil de signification.

SCIENCE

Demystifying Science: What is Xicc++ ?

JULY 09, 2017 04:02 IST

UPDATED: JULY 09, 2017 00:42 IST

SHARE ARTICLE



PRINT



Package: xicc (0.2-3) [universe]

set the ICC colour profile for an X display

Other Packages Related to xicc

● depends ● recommends ● suggests ● enhances

- **libc6** (>= 2.0) [i386]
Embedded GNU C Library: Shared libraries
also a virtual package provided by **libc6-udeb**
- **libc6** (>= 2.17) [arm64, ppc64el]
- **libc6** (>= 2.2.5) [amd64]
- **libc6** (>= 2.4) [armhf]

It's a new kind of quark, or a subatomic particle, that is a basic building block of matter. Something like the Xicc++ was known to exist in theory since decades but was finally 'observed' by physicists working at the Large Hadron Collider facility in Geneva and made public last week.



GAME ▾

STORY ▾

NEWS

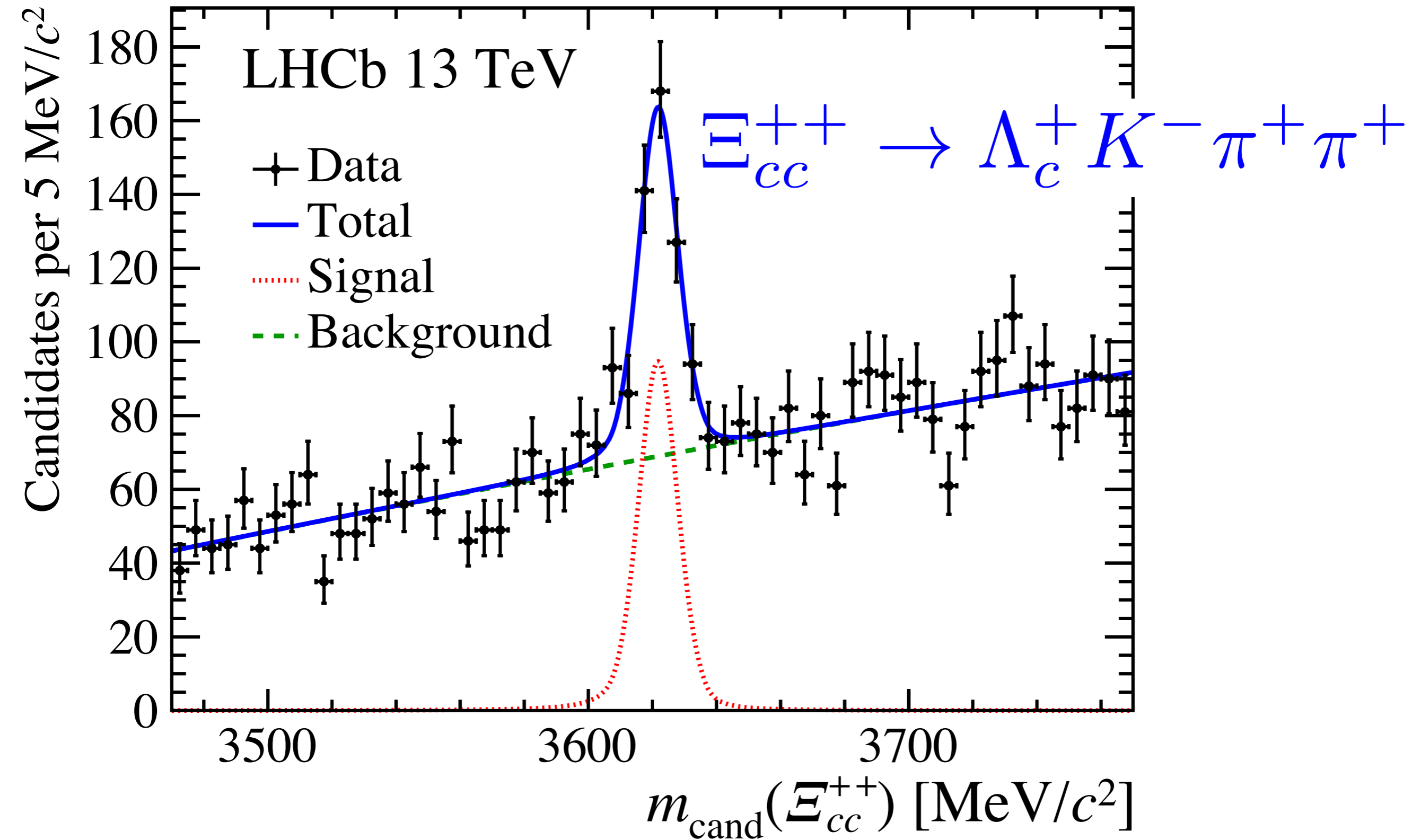


Xicc

15620 ✕ 850 ILVL

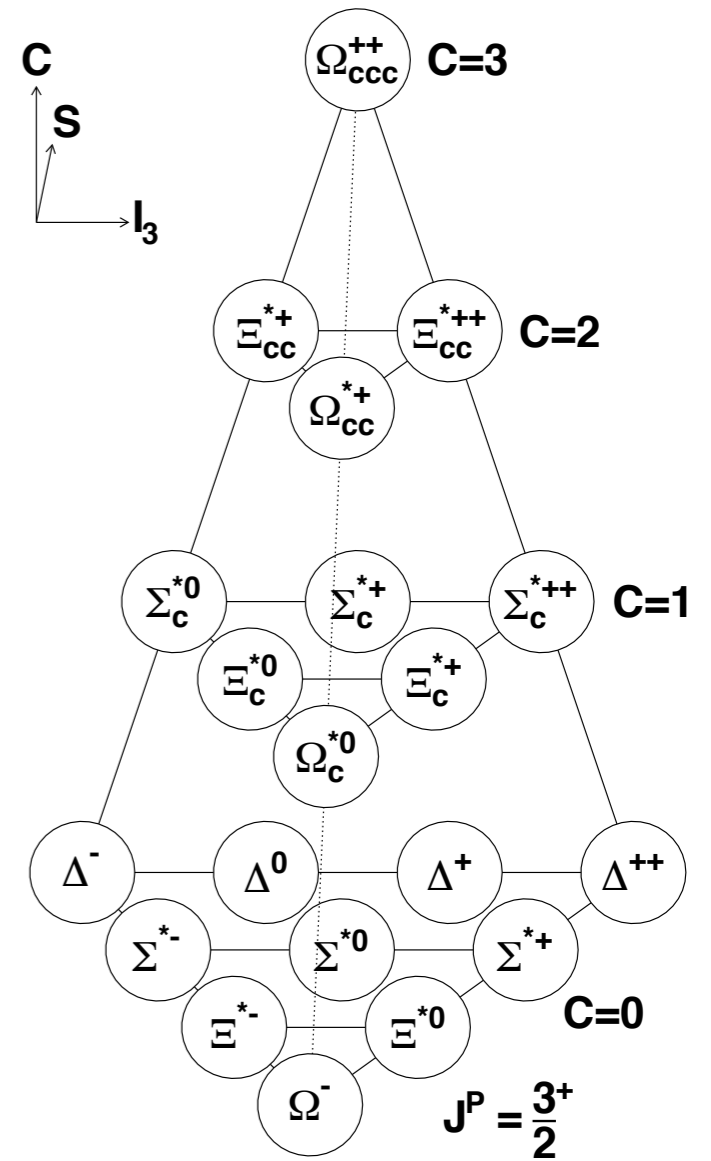
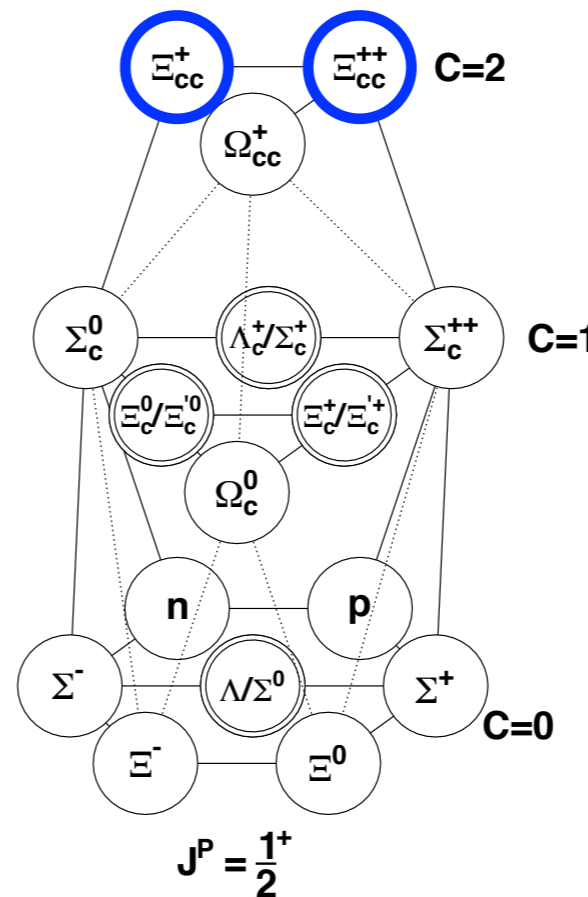
110 Worgen Marksmanship Hunter Stormrage

Doubly charmed baryon!



One year earlier...

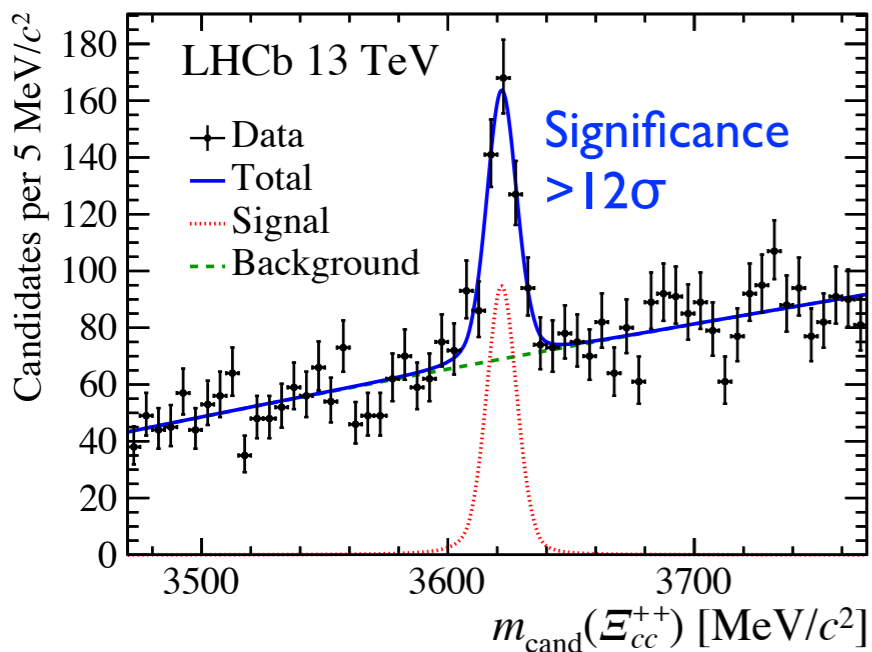
- Ξ_{cc} predictions:
 - Isodoublet of weakly decaying states: Ξ_{cc}^{++} (ccu) and Ξ_{cc}^{+} (ccd)
 - Mass around 3500-3700 MeV
 - Lifetimes: $\tau(\Xi_{cc}^{++}) \sim$ few hundred fs; $\tau(\Xi_{cc}^{++})/\tau(\Xi_{cc}^{+}) \sim 3$ to 4
- Ξ_{cc} observation(s) claimed by SELEX, but never reproduced.
 - Various oddities with SELEX result
- Longstanding experimental puzzle!



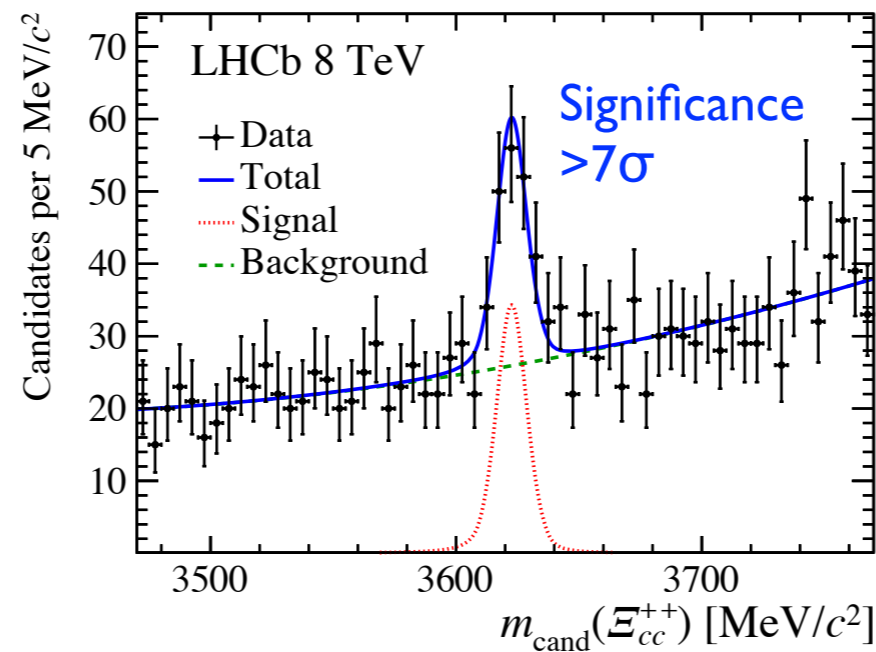
LHCb observation of Ξ_{cc}^{++}

- LHCb observed Ξ_{cc}^{++} with 2016 sample (main dataset) and measured $m = 3621.40 \pm 0.72$ (stat) ± 0.27 (syst) ± 0.14 (Λ_c^+) MeV/c^2
- Observation confirmed with 2012 sample (crosscheck).
- Lifetime significantly different from zero: observation persists when requiring $(t / \sigma_t) > 5$
 - σ_t varies event by event, but typical resolution $\sim 40\text{-}50\text{fs}$

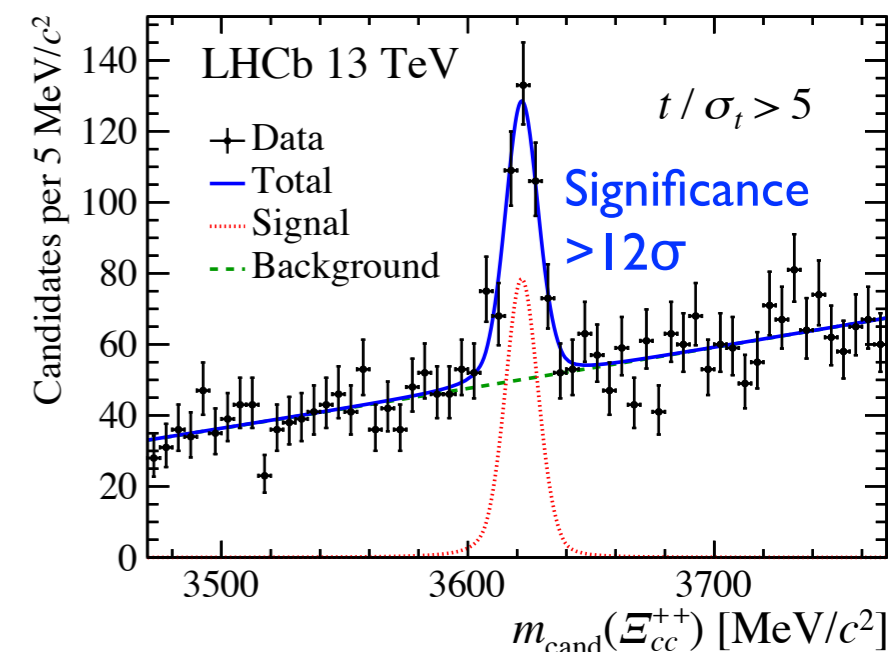
2016 data



2012 data

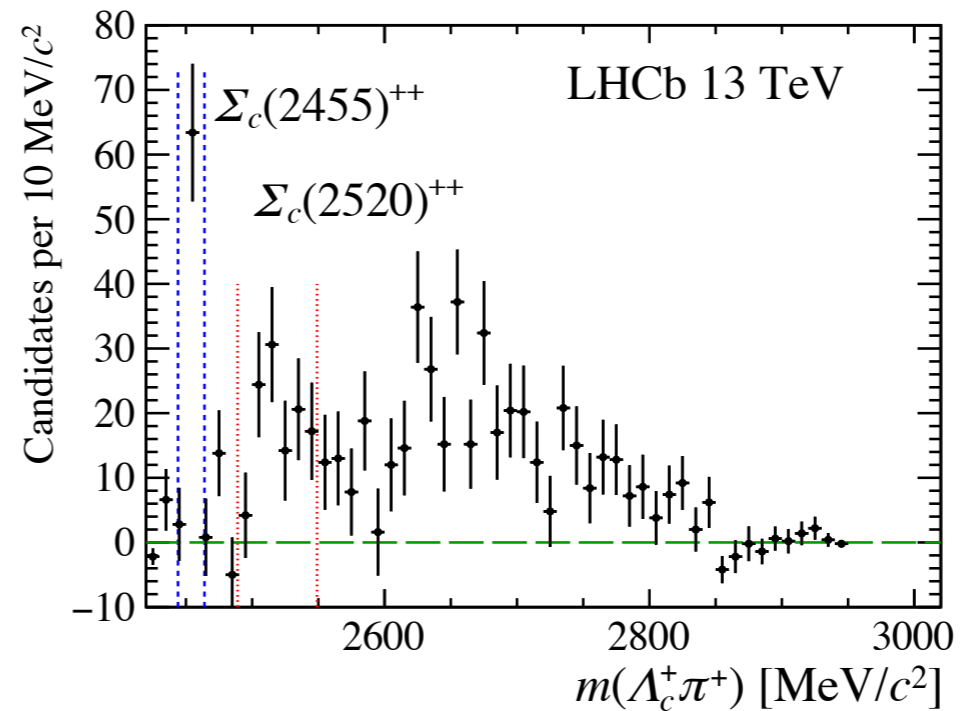
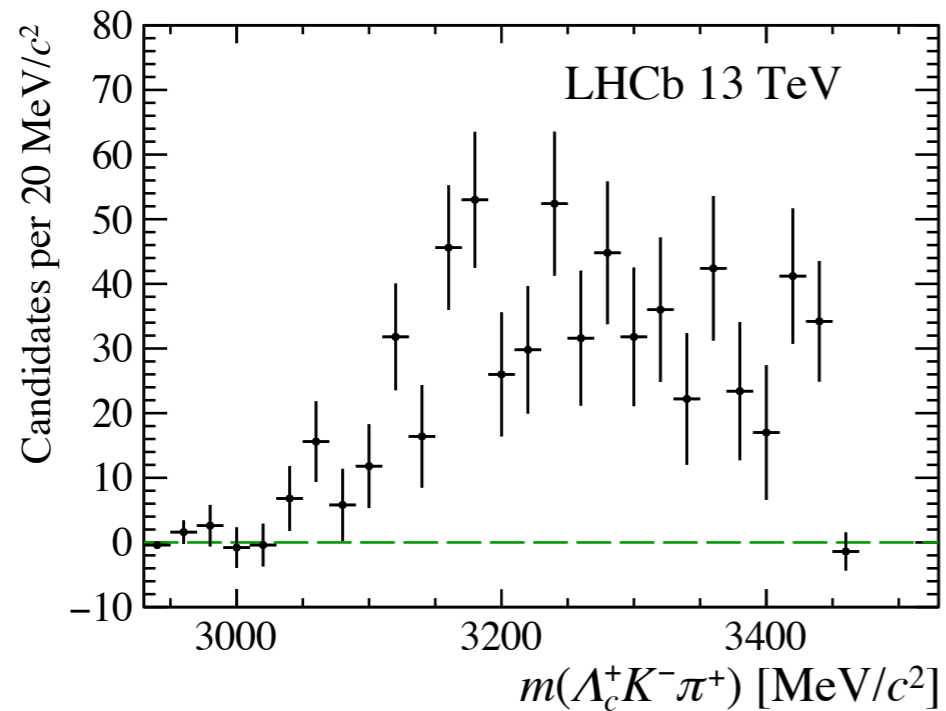
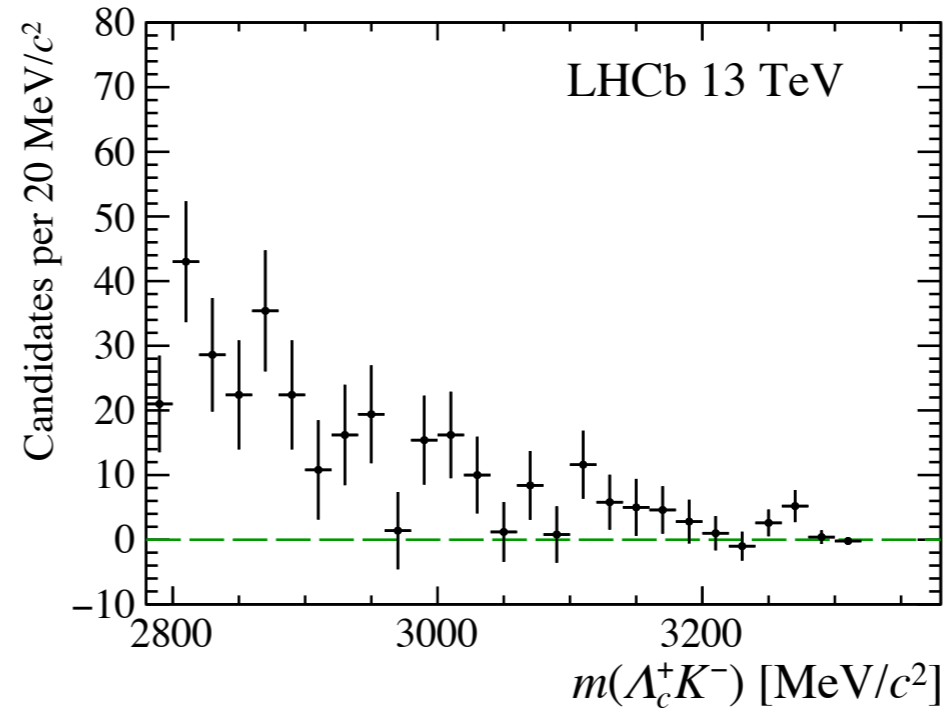
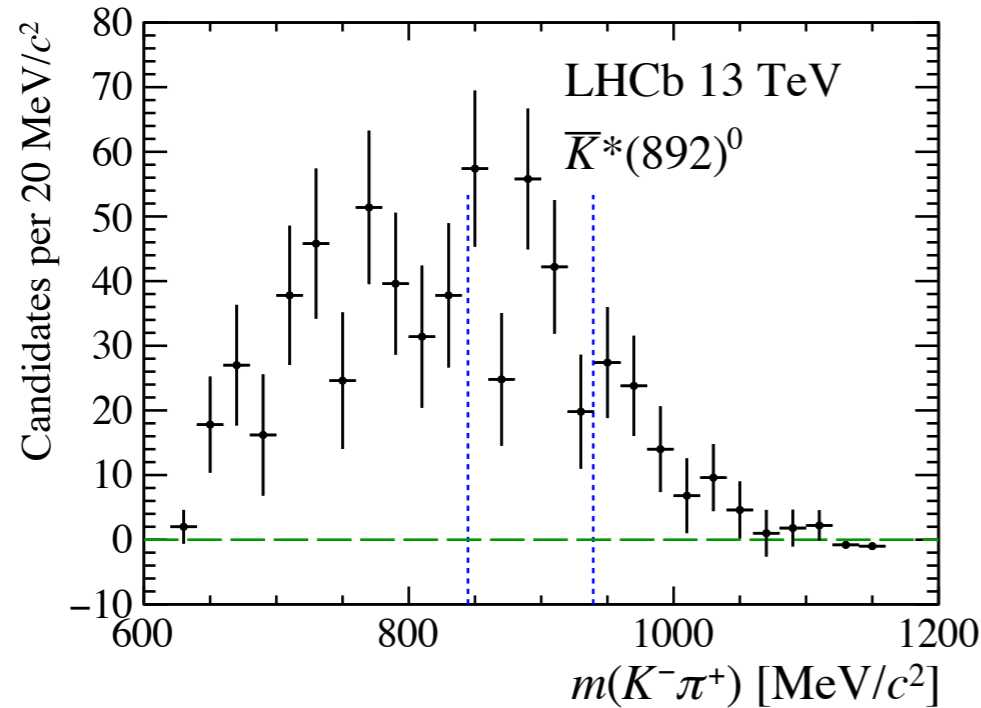


With lifetime cut



Resonant structure?

Background-subtracted $\Xi_{cc}^{++} \rightarrow \Lambda_c^+ K^- \pi^+ \pi^+$



Picture is fuzzy. Clearly some resonant contributions, e.g. $\Sigma_c(2455)^{++}$, but no single resonance dominates.

LHCb Ξ_{cc}^{++} vs SELEX Ξ_{cc}^+

- Masses differ by 103 ± 2 MeV (3621 vs 3519 MeV).
- Clearly not compatible with being an isodoublet.
- Both claims are for weakly decaying states \Rightarrow not compatible with one being an excitation of the other.
- Reported lifetime of SELEX Ξ_{cc}^+ very short (< 33 fs).

In a very narrow, legal sense, the results aren't mutually exclusive, but in practice either (a) one is wrong*, or (b) the constituent quark model is badly broken.

But enough history -- what about the future?

[LHCb-PAPER-2017-018](#) ; [arXiv:1707.01621](#)

* SELEX peak could also be a narrow, strongly decaying Ξ_c resonance that they misidentified... but this interpretation doesn't explain nonobservation of $\Xi_c(2970)$, $\Xi_c(3055)$, etc.

What's next for Ξ_{cc}^{++} ?

- Ξ_{cc}^{++} lifetime measurement! Crucial to the interpretation.
 - Will also give us a clue about Ξ_{cc}^+ lifetime.
- Ξ_{cc}^{++} production cross-section veil of ignorance

- More Ξ_{cc}^{++} decay modes! Ratios of branching fractions.
- Constraints on spin, parity
- Ξ_{cc}^+ search
 - Tougher, due to shorter expected lifetime.
 - We failed to find $\Xi_{cc}^+ \rightarrow \Lambda_c^+ K^- \pi^+$ with 0.65 fb^{-1} earlier... but we have better triggers, more modes, and much more data now.
- Ξ_{cc}^+ properties, production relative to Ξ_{cc}^{++} , etc
- Ω_{cc}^+ search (tougher still)
- Excited states
 - Lowest expected to decay electromagnetically -- tough for LHCb
 - Add pions, kaon... see what's out there!

Déjà vu: Why do we do spectroscopy?

In roughly ascending order:

- It's good clean fun^W physics
 - ... and sometimes you actually get to see a signal
- It provides inputs to pheno calculations in the area
- **Empowers searches for New Physics**
 - Necessary inputs, e.g. charm resonances for $B \rightarrow D^{(*)} \tau \nu$
 - **Opens new paths** (time until first CPV search with Ξ_{cc} ?)
- Sometimes you find something unexpected
 - See: pentaquarks

Let's have a concrete example!

Search for Ξ_b^0 oscillations

- Basic idea: search for **baryon-antibaryon oscillations** of Ξ_b^0 (bsu), tagged by resonances: $\Xi_b^{(*\prime)-} \rightarrow \Xi_b^0 \pi^-$
- Baryon oscillations similar to meson mixing, but **without DCS/** interference terms. [B-field effects negligible for Ξ_b^0 at LHCb]

Ξ_b^0 mixing:

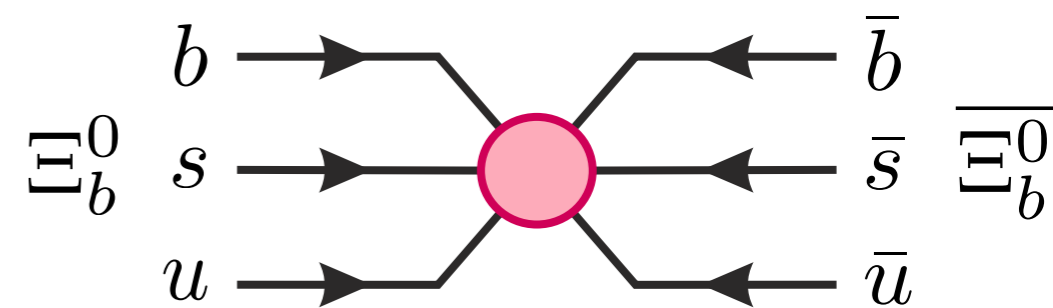
$$\frac{\Gamma(\Xi_b^0 \rightarrow \bar{f})}{\Gamma(\Xi_b^0 \rightarrow f)} \simeq \left| \frac{q}{p} \right|^2 [(\Delta\Gamma/4)^2 + (\Delta m/2)^2] t^2 = \alpha t^2$$

pure quadratic in t

D^0 mixing:

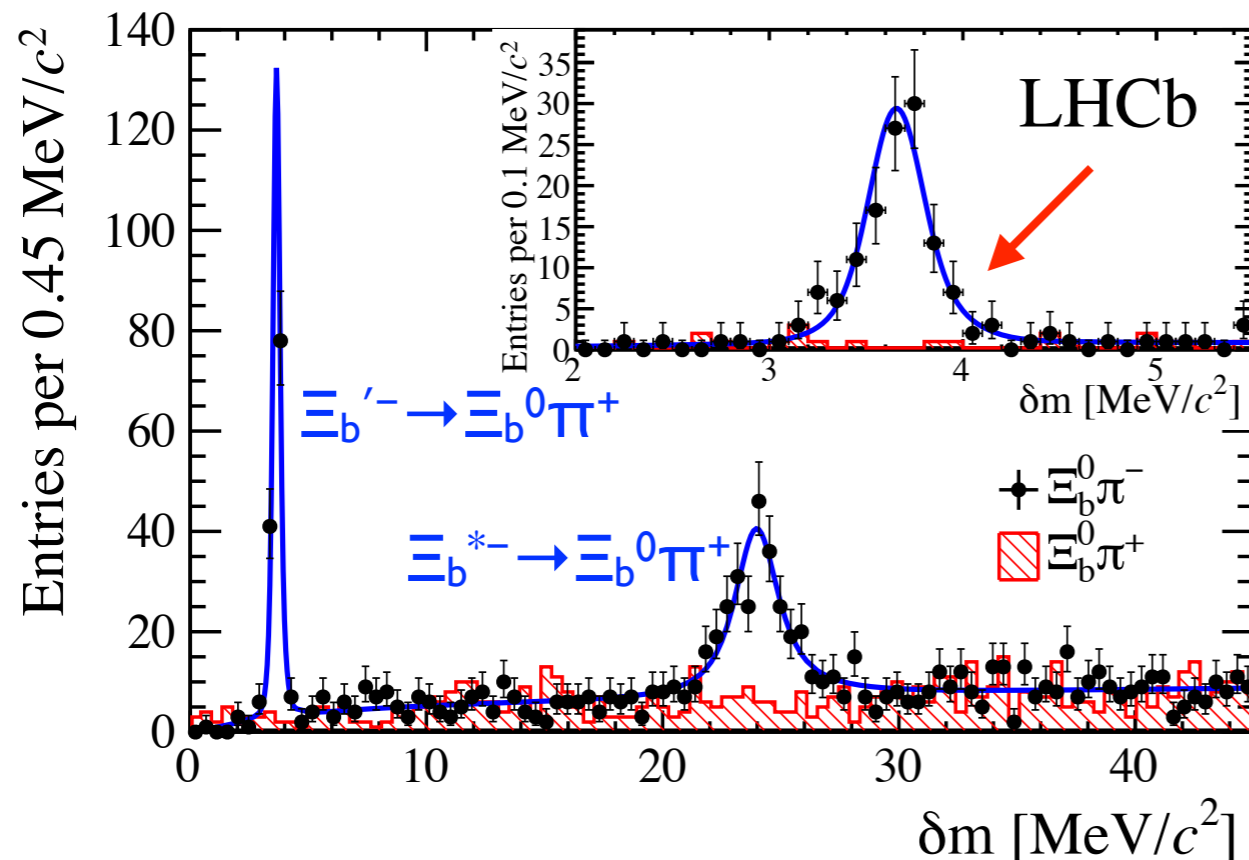
$$\frac{\Gamma(D^0 \rightarrow K^+ \pi^-)}{\Gamma(D^0 \rightarrow K^- \pi^+)} \simeq R_D + \sqrt{R_D} y' \Gamma t + (x'^2 + y'^2) \Gamma^2 t^2$$

- **Physics motivation:** flavour-diagonal six-fermion BSM vertices allow such BNV without excessive proton decay. (Ξ_b^0 is special: one valence quark from each generation)



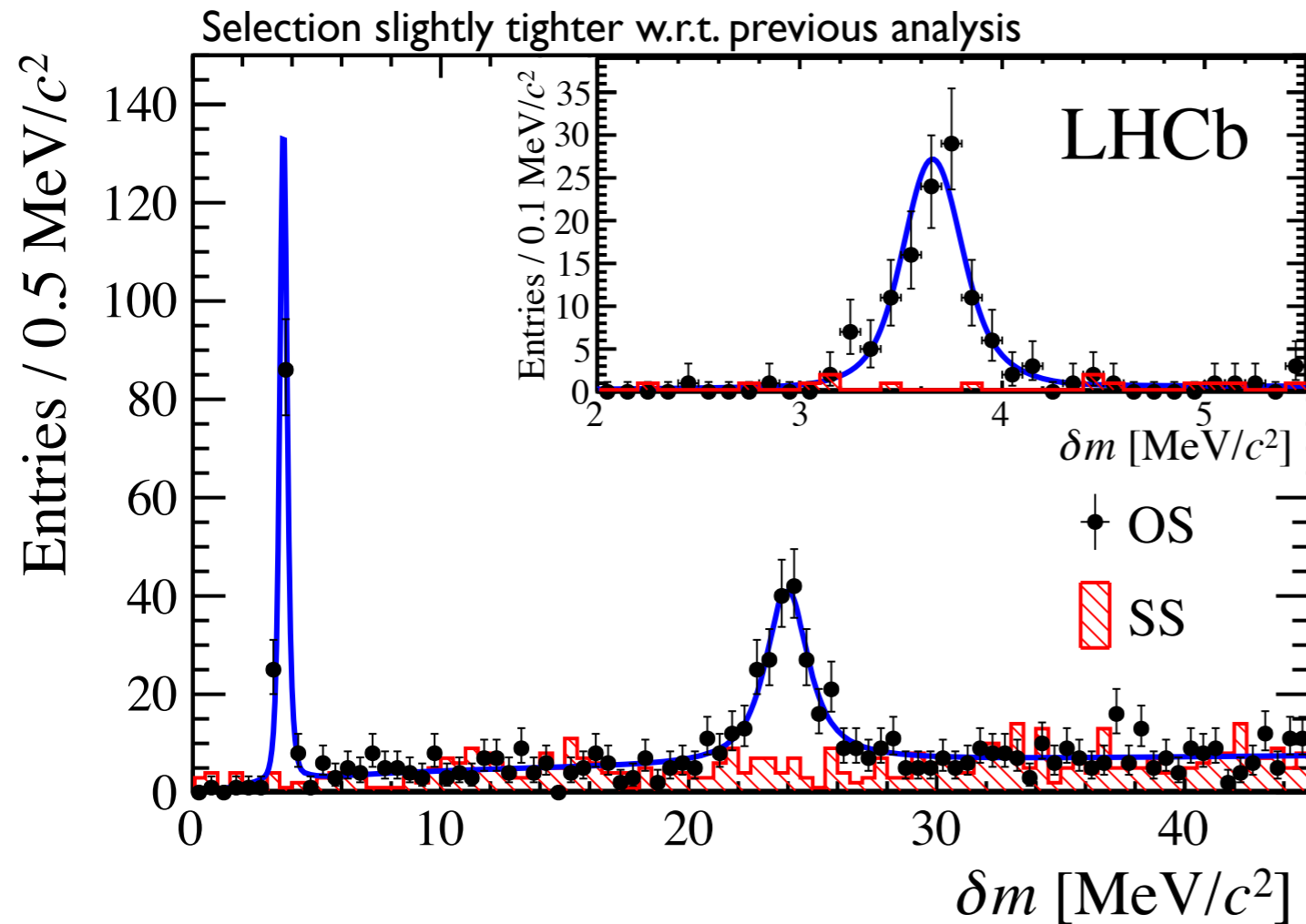
Search for Ξ_b^0 oscillations

- To make the analysis go, you need to do three things:
 - Measure the Ξ_b^0 **decay time** -- easy at LHCb
 - Identify the Ξ_b^0 **baryon number at decay** -- easy from final state
 - Identify (**tag**) the Ξ_b^0 **baryon number at production** -- hmm!
- LHCb has tagging tools... but tagging power is $O(\text{few } \%)$.
 - Typical mistag rate: forty-some percent. This is fine if looking at B^0/B_s oscillations with high stats, but fatal for a very rare BSM process.
- **But we have a solution! Remember this?**



With these freshly discovered resonances, we have a small but very clean sample of Ξ_b^0 whose flavour at production is known.

Search for Ξ_b^0 oscillations

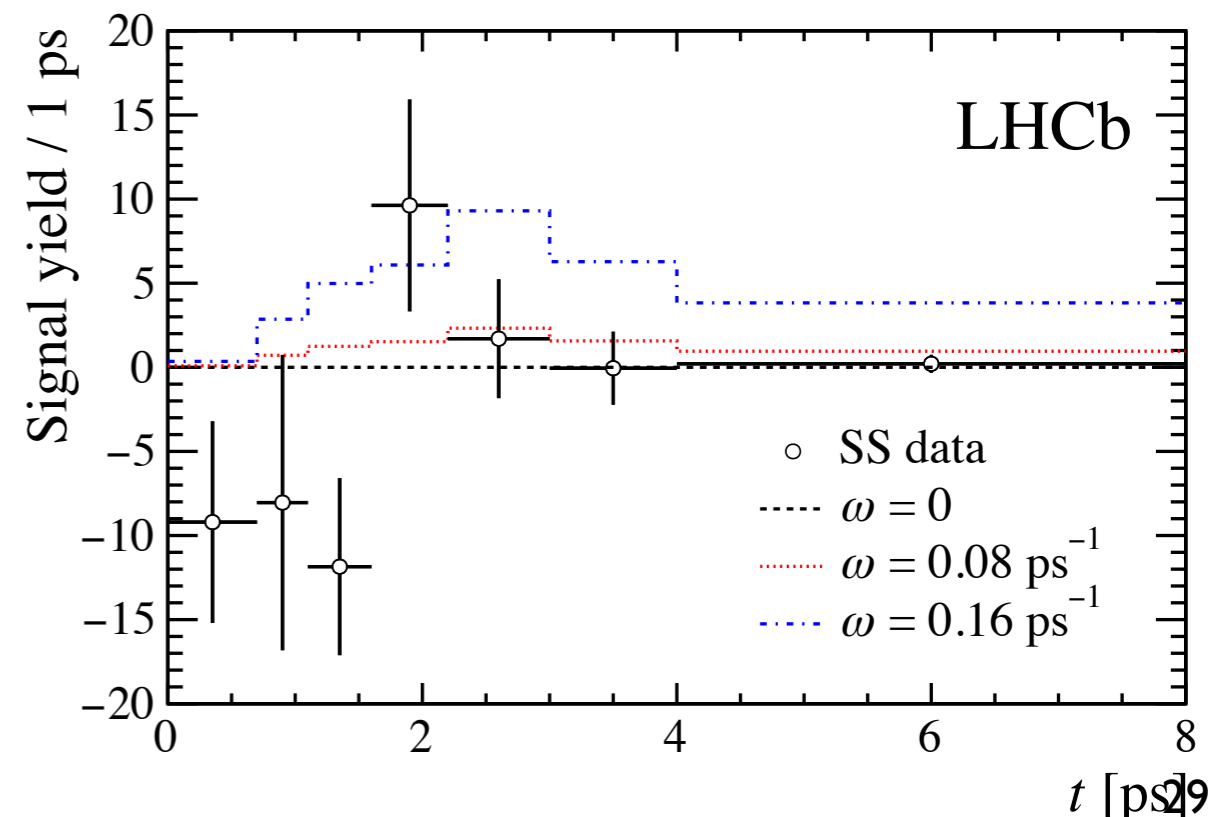


$$R(t) \equiv \frac{P_{X \rightarrow \bar{X}}(t)}{P_{X \rightarrow X}(t)} = \tan^2(t/\tau_{\text{mix}}) \simeq \frac{t^2}{\tau_{\text{mix}}^2} \equiv (\omega t)^2$$

- Data => $\omega < 0.08 \text{ ps}^{-1}$ at 95% CL
- Corresponds to oscillation period $(2\pi/\omega) > 80 \text{ ps}$ at 95% CL
- c.f. Ξ_b^0 lifetime $\sim 1.5 \text{ ps}$
- Can improve with Run 2, and with semileptonic decays of Ξ_b^0

Absence of a clear signal in the same-sign (BNV) channel.

Quantitative analysis to set upper limit on oscillation rate.



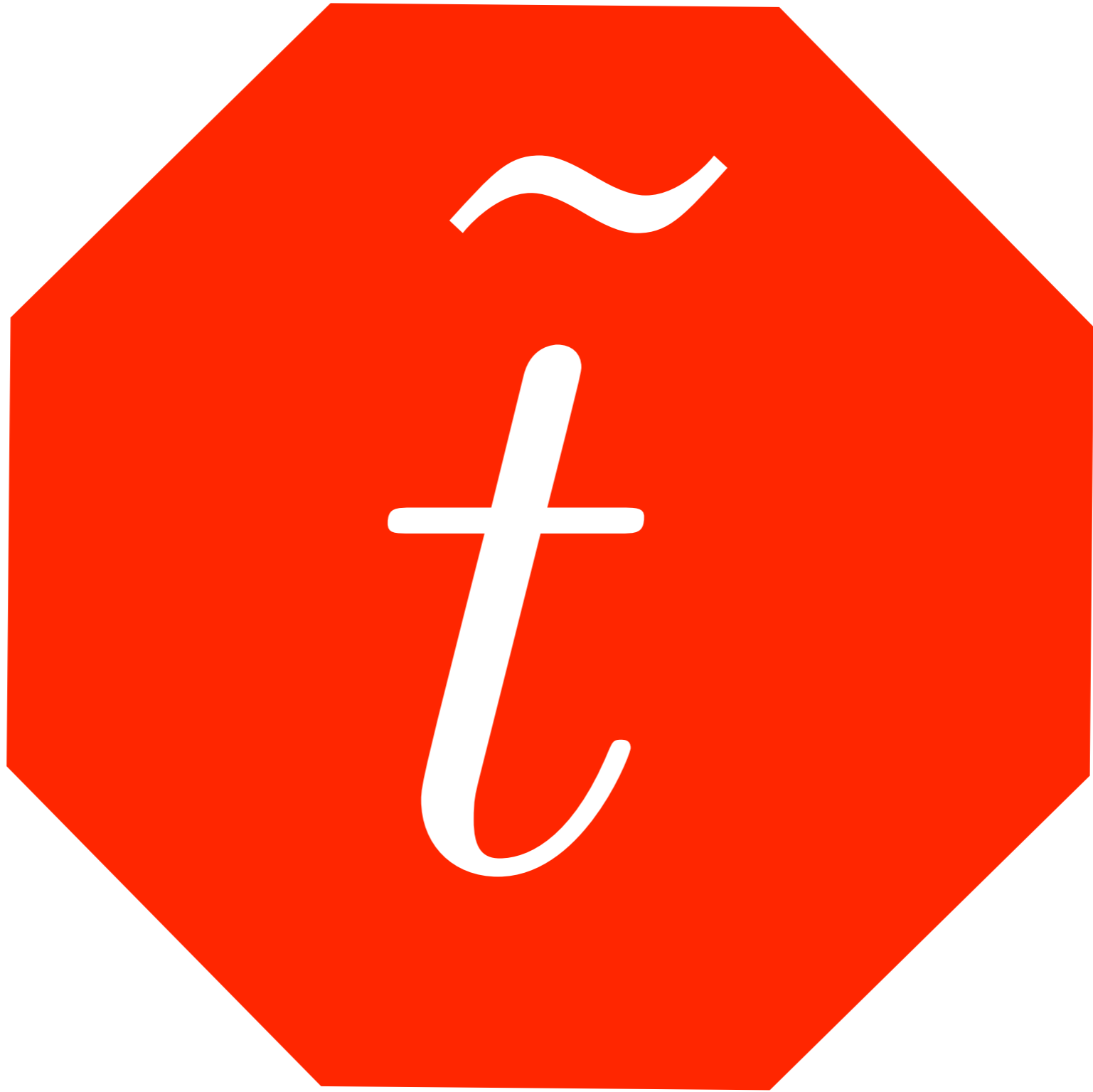
Penultimate words

- LHCb has done some really cool work in spectroscopy.
 - Hat tip especially to the folks who worked on $Z(4430)^-$ and pentaquark analyses.
- But these things don't come for free.
- As we upgrade the detector and move to a real-time trigger, we will get bigger and bigger samples of things we look for.
 - Ξ_{cc} is the poster child here -- analysis used [Turbo \(real-time\) output](#).
Not a coincidence that 2016 sample was more sensitive.
- ... but we will have **zero** sensitivity to things we don't plan for.
- Therefore: [we need to plan carefully ahead of Run 3](#).
 - Maxim: the technical challenges charm has today, beauty will have a couple of years later.
- Also therefore: if you have a clever theory about new particles, make sure we have a trigger that can find them!

Last words

- Lots of good physics going on
 - Some makes headlines and some does not, but it all helps us understand flavour physics better.
- "Yesterday's discovery is today's calibration and tomorrow's background."
- Looking forward to many new backgrounds!





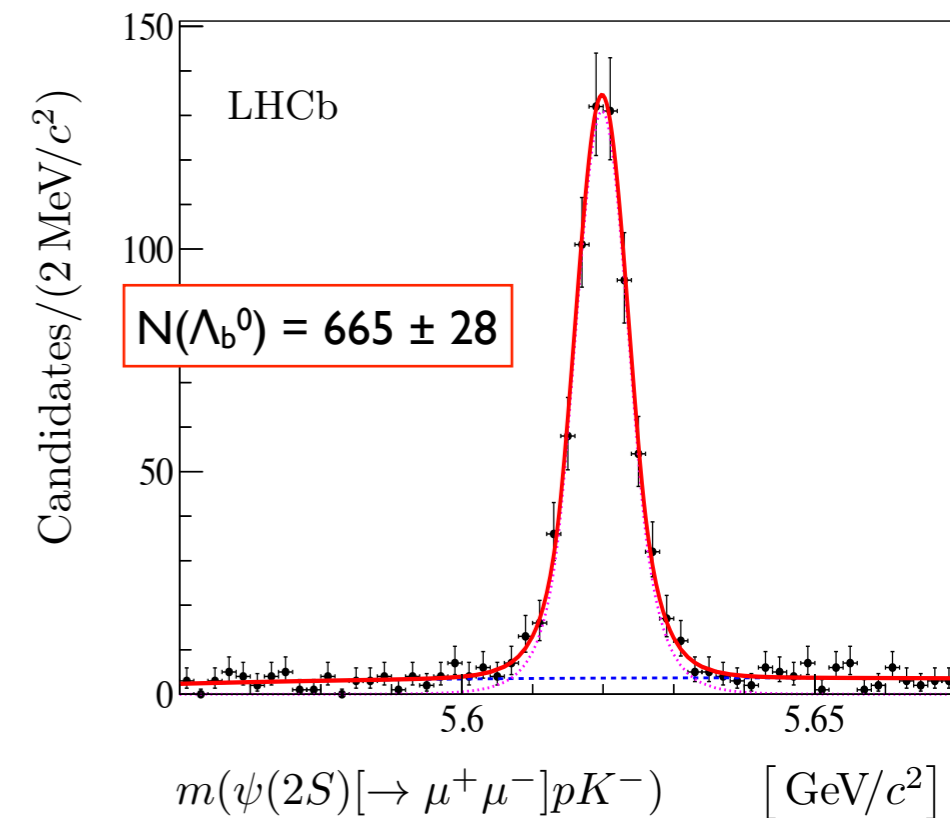
LHCb links & references

- [LHCB-PAPER-2017-023](#) ; [arXiv:1708.05808](#) : Search for baryon-number-violating Ξ_b^0 oscillations
- [LHCB-PAPER-2017-018](#) ; [arXiv:1707.01621](#) : Observation of the doubly charmed baryon Ξ_{cc}^{++}
- [LHCB-PAPER-2017-002](#) ; [arXiv:1703.04639](#) : Observation of five new narrow Ω_c^0 states decaying to $\Xi_c^+ K^-$
- [LHCB-PAPER-2016-061](#) ; [arXiv:1701.07873](#) : Study of the $D^0 p$ amplitude in $\Lambda_b^0 \rightarrow D^0 p \pi^-$ decays
- [LHCB-PAPER-2016-010](#) ; [arXiv:1604.03896](#) : Measurement of the properties of the Ξ_b^{*0} baryon
- [LHCB-PAPER-2016-008](#) ; [arXiv:1604.01412](#) : Measurement of the mass and lifetime of the Ω_b^- baryon
- [LHCB-PAPER-2015-060](#) ; [arXiv:1603.06961](#) : Observation of $\Lambda_b^0 \rightarrow \psi(2S) p K^-$ and $\Lambda_b^0 \rightarrow J/\psi \pi^+ \pi^- p K^-$ decays and a measurement of the Λ_b^0 baryon mass
- [LHCB-PAPER-2015-047](#) ; [arXiv:1510.03829](#) : Evidence for the strangeness-changing weak decay $\Xi_b^- \rightarrow \Lambda_b^0 \pi^-$
- [LHCB-PAPER-2014-061](#) ; [arXiv:1411.4849](#) : Observation of two new Ξ_b^- baryon resonances
- [LHCB-PAPER-2014-048](#) ; [arXiv:1409.8568](#) : Precision measurement of the mass and lifetime of the Ξ_b^- baryon
- [LHCB-PAPER-2014-021](#) ; [arXiv:1405.7223](#) : Precision measurement of the mass and lifetime of the Ξ_b^0 baryon
- [LHCb-PAPER-2014-010](#) ; [arXiv:1405.1543](#) : Measurement of the Ξ_b^- and Ω_b^- baryon lifetimes
- [LHCB-PAPER-2014-003](#) ; [arXiv:1402.6242](#) : Precision measurement of the ratio of the Λ_b^0 to \bar{B}^0 lifetimes
- [LHCB-PAPER-2014-002](#) ; [arXiv:1403.3606](#) : Study of beauty hadron decays into pairs of charm hadrons
- [LHCB-PAPER-2013-056](#) ; [arXiv:1311.4823](#) : Studies of beauty baryon decays to $D^0 p h^-$ and $\Lambda_c^+ h^-$ final states
- [LHCB-PAPER-2013-049](#) ; [arXiv:1310.2538](#) : Search for the doubly charmed baryon Ξ_{cc}^+
- [LHCB-PAPER-2012-048](#) ; [arXiv:1302.1072](#) : Measurement of the Λ_b^0 , Ξ_b^- and Ω_b^- baryon masses
- [LHCB-PAPER-2012-012](#) ; [arXiv:1205.3452](#) : Observation of excited Λ_b^0 baryons
- [LHCB-PAPER-2011-035](#) ; [arXiv:1112.4896](#) : Measurement of b-hadron masses

b-baryon ground state masses

- Masses all relative to Λ_b^0
- World average Λ_b^0 mass dominated by an LHCb study of $\Lambda_b^0 \rightarrow \mu^+ \mu^- p K^- [\pi^+ \pi^-]$

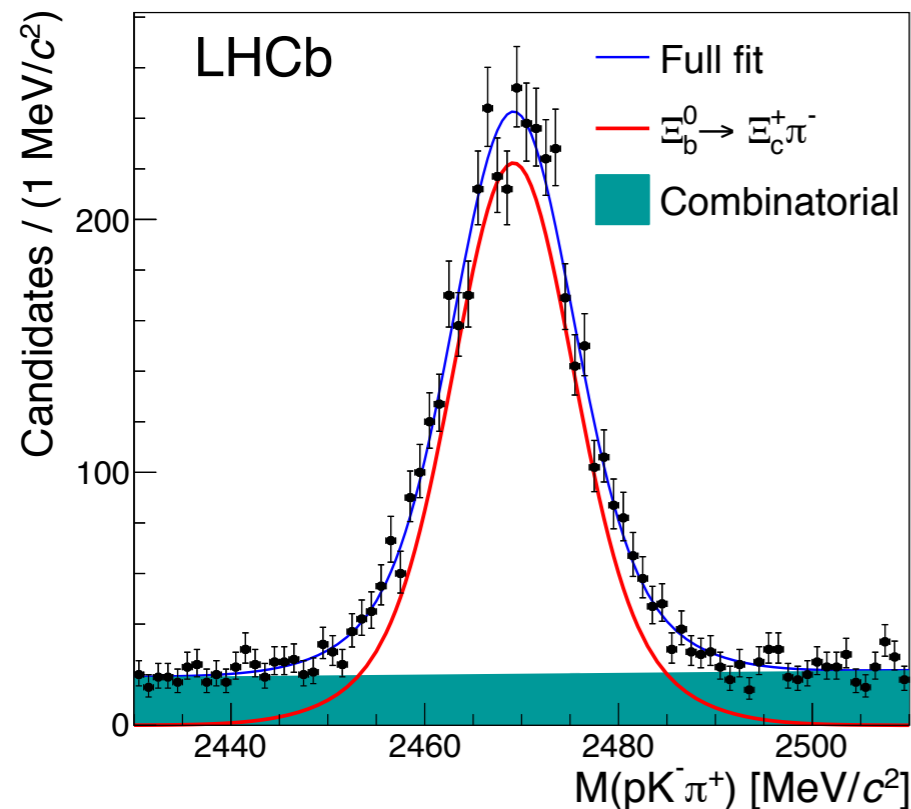
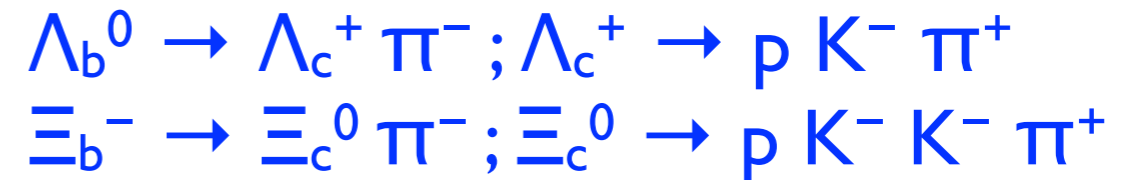
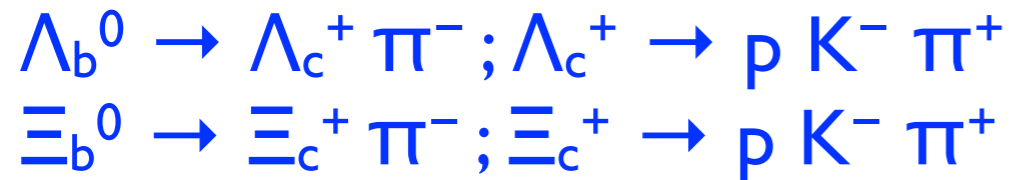
Channel	$M(\Lambda_b^0)$ [MeV/c ²]
$\Lambda_b^0 \rightarrow J/\psi p K^-$	$5619.62 \pm 0.04 \pm 0.34$
$\Lambda_b^0 \rightarrow \psi(2S)[\rightarrow \mu^+ \mu^-] p K^-$	$5619.84 \pm 0.18 \pm 0.19$
$\Lambda_b^0 \rightarrow \psi(2S)[\rightarrow J/\psi \pi^+ \pi^-] p K^-$	$5619.38 \pm 0.33 \pm 0.18$
$\Lambda_b^0 \rightarrow J/\psi \pi^+ \pi^- p K^-$ excluding $\psi(2S)$	$5619.08 \pm 0.30 \pm 0.27$
$M(\Lambda_b^0) = 5619.65 \pm 0.17 \pm 0.17 \text{ MeV}/c^2$	



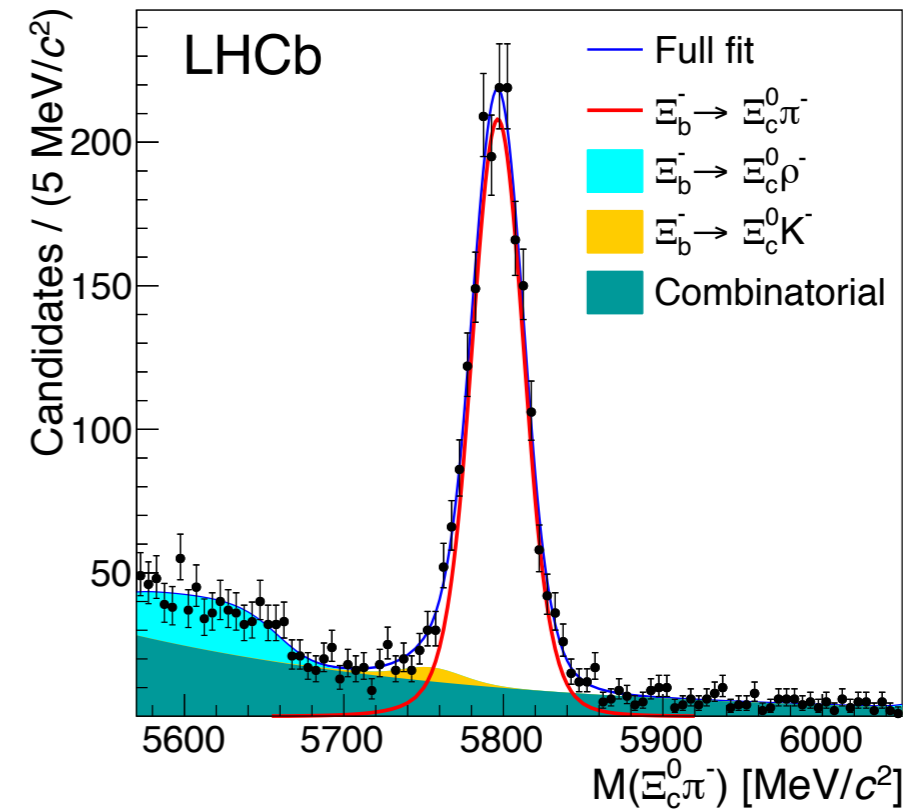
- Trade-off between stat & sys errors!
- Large q-value => large phase space => more stats
- Small q-value => closer to threshold => less to measure
- Winner is a near-threshold mode. This will be a pattern!

b-baryon ground state masses

- Masses all relative to Λ_b^0
- Masses of Ξ_b^- and Ξ_b^0 dominated by a pair of papers:



$$M(\Xi_b^0) - M(\Lambda_b^0) = 172.44 \pm 0.39 \pm 0.17 \text{ MeV}/c^2$$



$$M(\Xi_b^-) - M(\Lambda_b^0) = 178.36 \pm 0.46 \pm 0.16 \text{ MeV}/c^2$$

In passing, also obtain:

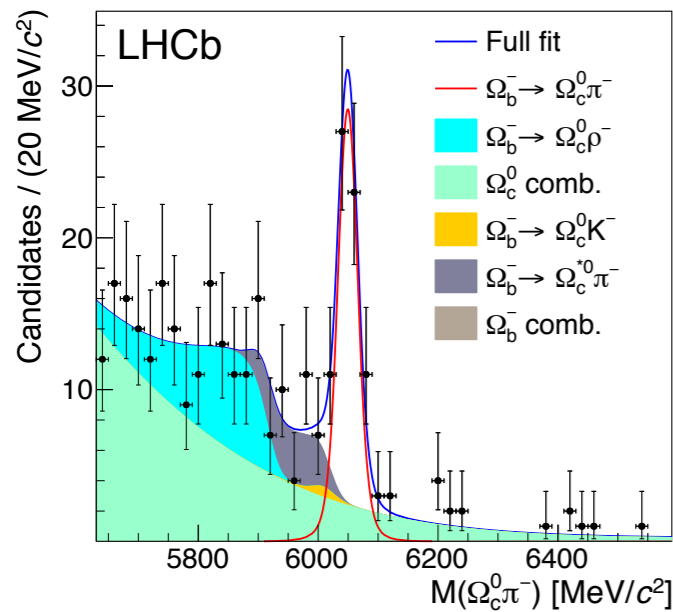
$$M(\Xi_c^+) - M(\Lambda_c^+) = 181.51 \pm 0.14 \pm 0.10 \text{ MeV}/c^2$$

[LHCb-PAPER-2014-048](#) ; [arXiv:1409.8568](#)

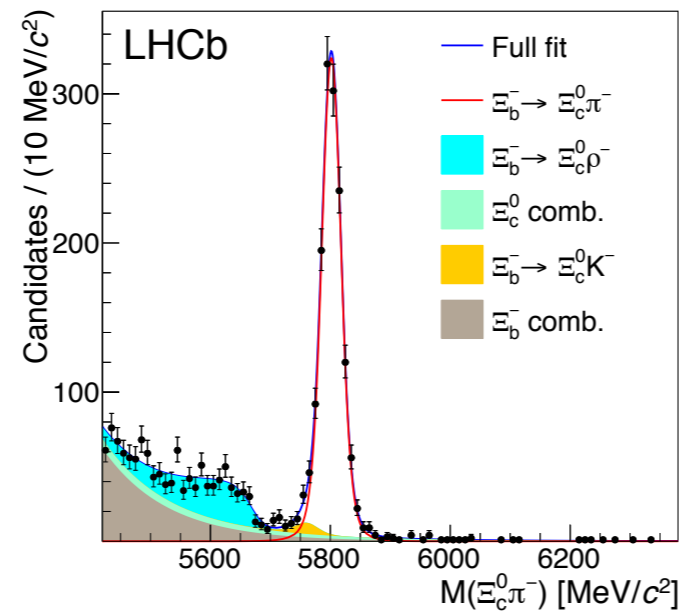
[LHCb-PAPER-2014-021](#) ; [arXiv:1405.7223](#)

b-baryon ground state masses

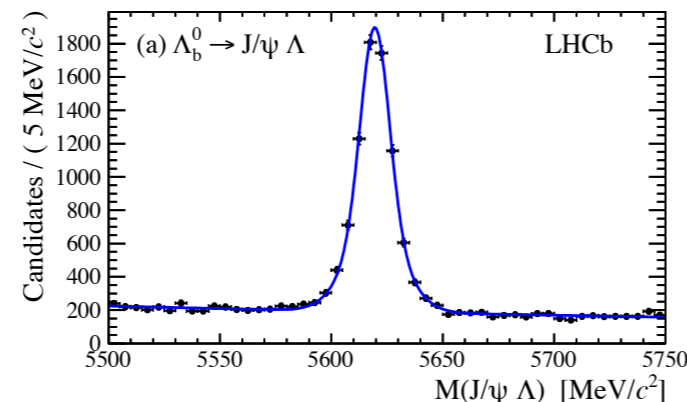
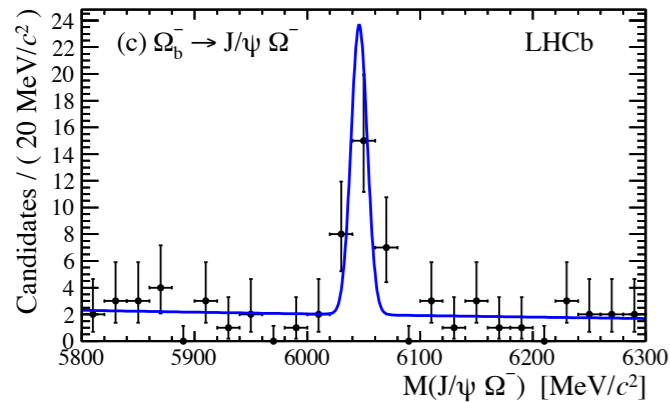
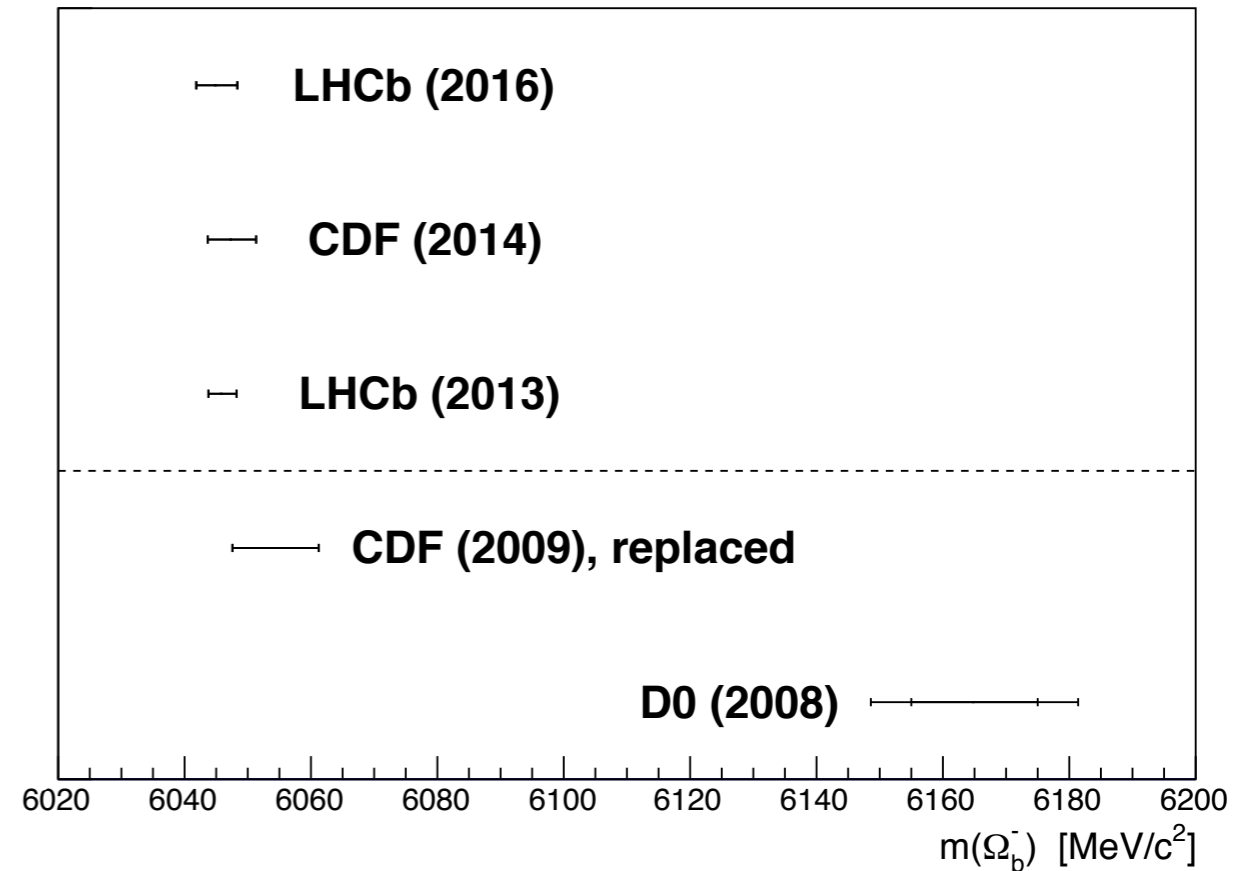
- Lastly, Ω_b^- . Here a few measurements contribute:



Signal mode



Control mode

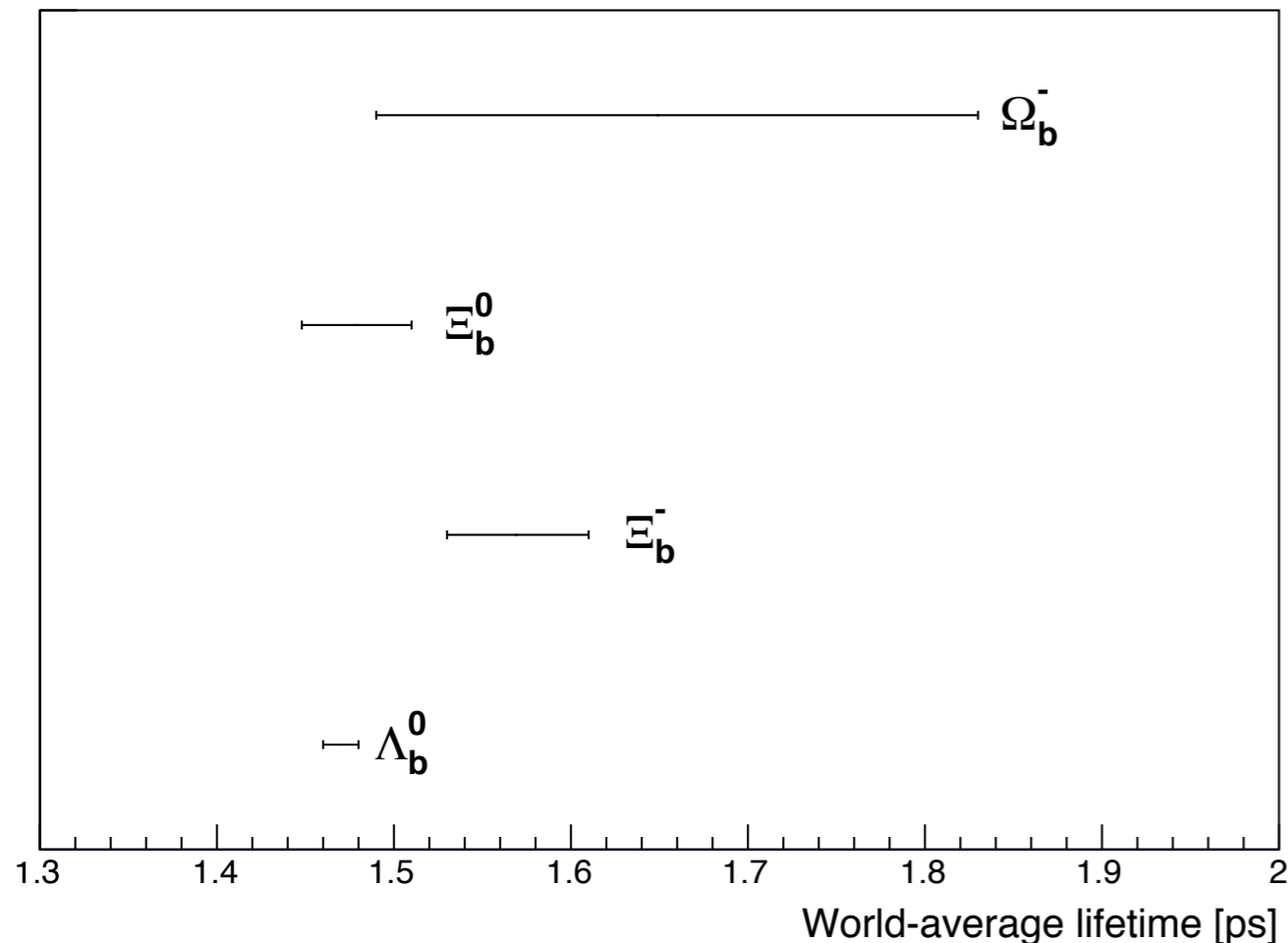


[LHCb-PAPER-2016-008](#) ; [arXiv:1604.01412](#)
[LHCb-PAPER-2012-048](#) ; [arXiv:1302.1072](#)

CDF: [arXiv:1403.8126](#) ; PRD 89, 072014 (2014)
D0: [arXiv:0808.4142](#) ; PRL 101:232002 (2008)

b-baryon ground state lifetimes

- Technique of measuring lifetime ratios relative to control mode used extensively
 - Λ_b^0 lifetime dominated by Λ_b^0/B^0 ratio ($\Lambda_b^0 \rightarrow J/\psi p K^- / \bar{B}^0 \rightarrow J/\psi \pi^+ K^-$)
 - Ξ_b^- and Ξ_b^0 dominated by Ξ_b/Λ_b^0 ratios (same as for masses)
 - Ω_b^- lifetime measured via ratio w.r.t. Ξ_b^- , also absolute via $\Omega_b^- \rightarrow J/\psi \Omega^-$
- Choose similar decay modes for numerator & denominator
- Many systematics cancel, especially if lifetimes similar



[LHCb-PAPER-2014-003](#) ; [arXiv:1402.6242](#)
[LHCb-PAPER-2014-048](#) ; [arXiv:1409.8568](#)
[LHCb-PAPER-2014-021](#) ; [arXiv:1405.7223](#)
[LHCb-PAPER-2016-008](#) ; [arXiv:1604.01412](#)
[LHCb-PAPER-2014-010](#) ; [arXiv:1405.1543](#)