



# 2HDM Pseudoscalar + DM Benchmarks

(intended summary of DM WG  
2HDM Discussions over last ~2 Weeks)

Jose Miguel No  
King's College London

LHC DM WG MEETING 20/06/17

# 2HDM + P<sub>(pseudoscalar)</sub>

## Parameter Executive Summary

$$\begin{aligned}
 V_{2\text{HDM}} = & \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 & + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 & + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[ (H_1^\dagger H_2)^2 + \text{h.c.} \right] \\
 V_P = & \frac{1}{2} m_P^2 P^2 + \kappa (i P H_1^\dagger H_2 + \text{h.c.}) \\
 & + \lambda_{P1} P^2 |H_1|^2 + \lambda_{P2} P^2 |H_2|^2
 \end{aligned}$$

- 12 Parameters +  $m_x + y_x$
- 2 fixed by EWSB + 1 "fixed" by Higgs properties + 1 "fixed" by EWPO

### "Free" Parameters

$m_a$   
 $m_A$   
 $m_H$

$\sin \theta$

$\tan \beta$

$\lambda_3, \lambda_{P1}, \lambda_{P2}$

enter

$g_{aAh}, g_{Haa}$

$g_{haa}$

Do not affect  
Kinematics

$m_x$   
 $y_x$

# 2HDM + $P_{(\text{pseudoscalar})}$

## (Further) "Relevant" Constraints

- Boundedness from Below of Scalar Potential  
(+ Absolute Stability of EW Vacuum)
- Unitarity of  $2 \rightarrow 2$  scattering processes

I will discuss  
these in detail

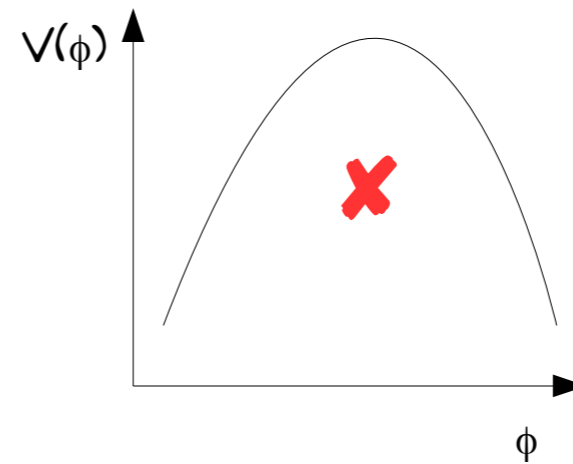
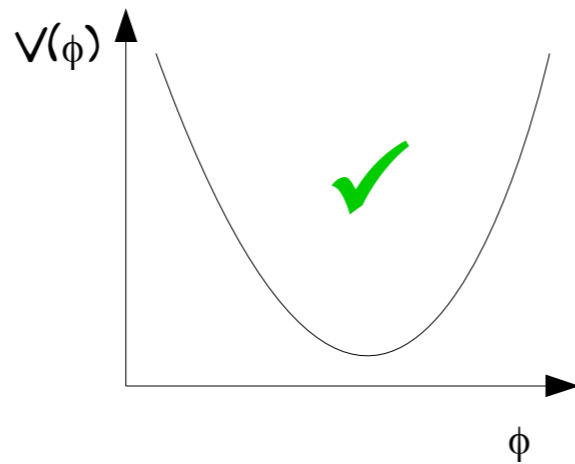
- Flavour (e.g.  $\bar{B} \rightarrow X_s \gamma$  ,  $B_s \rightarrow \mu^+ \mu^-$ )

...



# Boundedness from Below (BFB)

If ~~BFB~~ in  $V(\phi)$ , field(s) run away to  $\mp\infty$  (with  $V \rightarrow -\infty$ )

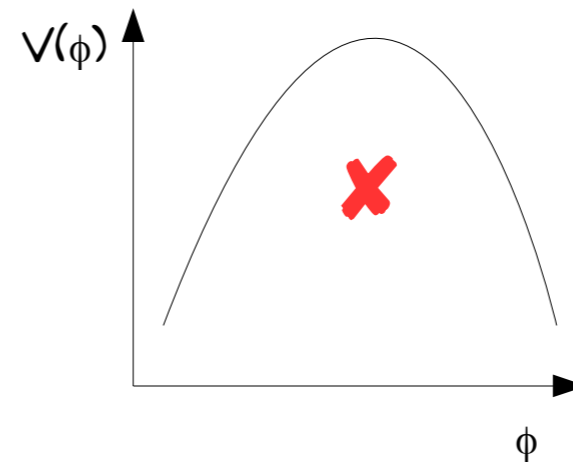
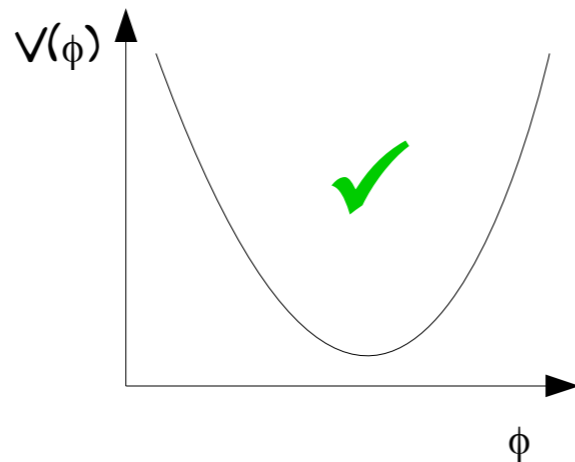


$$\begin{aligned} V_{2\text{HDM}} &= \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\ &+ \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\ &+ \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[ (H_1^\dagger H_2)^2 + \text{h.c.} \right] \end{aligned}$$

$$\begin{aligned} V_P &= \frac{1}{2} m_P^2 P^2 + \kappa (i P H_1^\dagger H_2 + \text{h.c.}) \\ &+ \lambda_{P1} P^2 |H_1|^2 + \lambda_{P2} P^2 |H_2|^2 + \lambda P^4 \end{aligned}$$

# Boundedness from Below (BFB)

If ~~BFB~~ in  $V(\phi)$ , field(s) run away to  $\mp\infty$  (with  $V \rightarrow -\infty$ )



$$\begin{aligned}
 V_{2\text{HDM}} = & \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 & + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 & + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} [(H_1^\dagger H_2)^2 + \text{h.c.}]
 \end{aligned}$$

$$\begin{aligned}
 V_P = & \frac{1}{2} m_{H^\pm}^2 |H^\pm|^2 + \kappa (i P H_1^\dagger H_2 + \text{h.c.}) \\
 & + \lambda_{H^\pm} |H^\pm|^4 + \lambda_{P^2} P^2 |H_2|^2 + \lambda P^4
 \end{aligned}$$

*Here I will not discuss these*

How to Check?

Look at  $V(\phi)$  at large field values  $\phi \gg v$

hep-ph/0207010

$$\text{BFB: } \lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > -\sqrt{\lambda_1 \lambda_2}, \quad \lambda_3 + \lambda_4 - |\lambda_5| > -\sqrt{\lambda_1 \lambda_2}$$

# Boundedness from Below (BFB)

$$\text{BFB: } \lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > -\sqrt{\lambda_1 \lambda_2}, \quad \lambda_3 + \lambda_4 - |\lambda_5| > -\sqrt{\lambda_1 \lambda_2}$$

$$m_{H^\pm} = m_H$$

$$\frac{m_h^2}{v^2} (1 - t_\beta^{-2}) + \lambda_3 t_\beta^{-2} > 0$$

$$\frac{m_h^2}{v^2} (1 - t_\beta^2) + \lambda_3 t_\beta^2 > 0$$

$$\Rightarrow \lambda_3 \geq \frac{m_h^2}{v^2}$$



# Boundedness from Below (BFB)

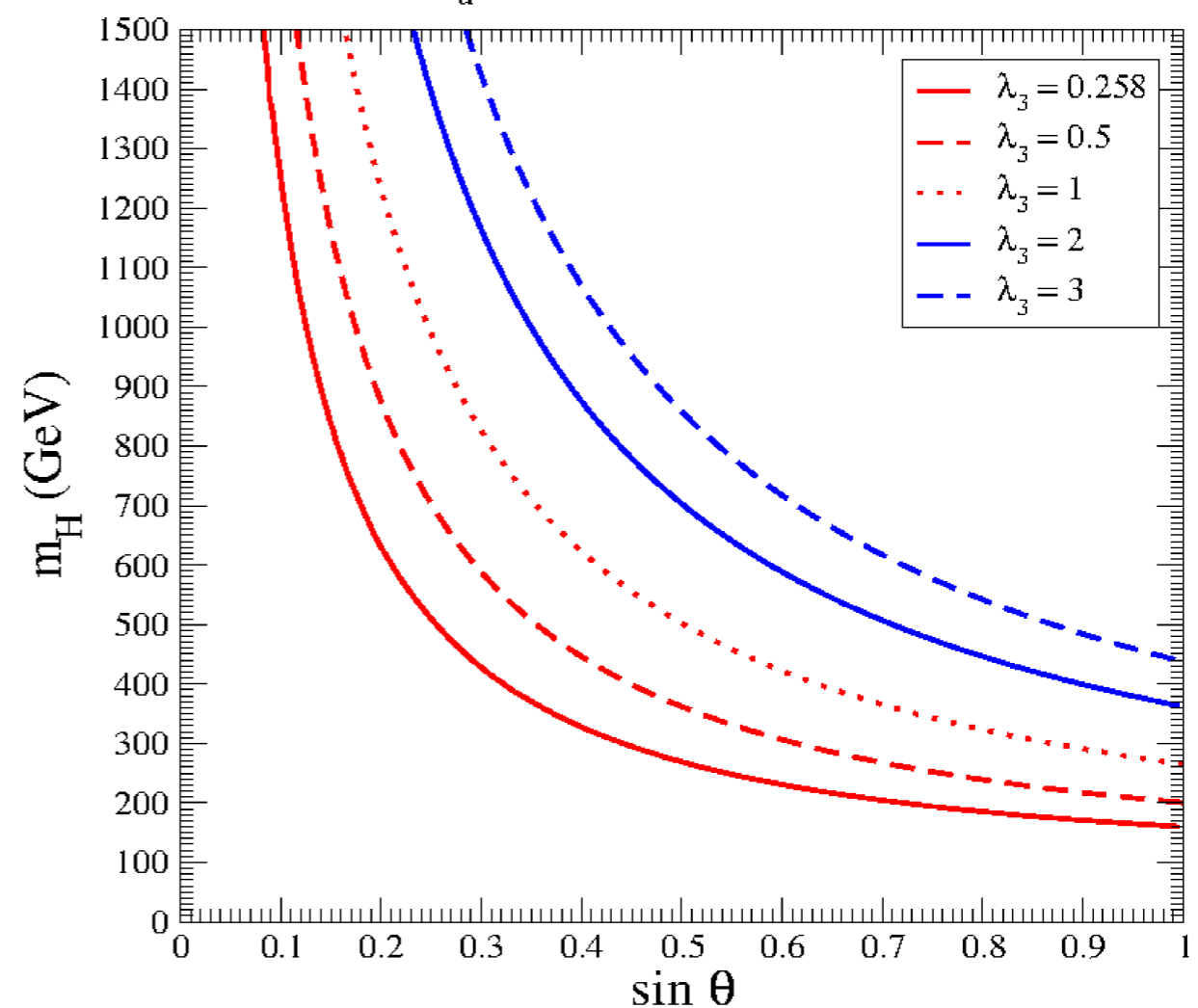
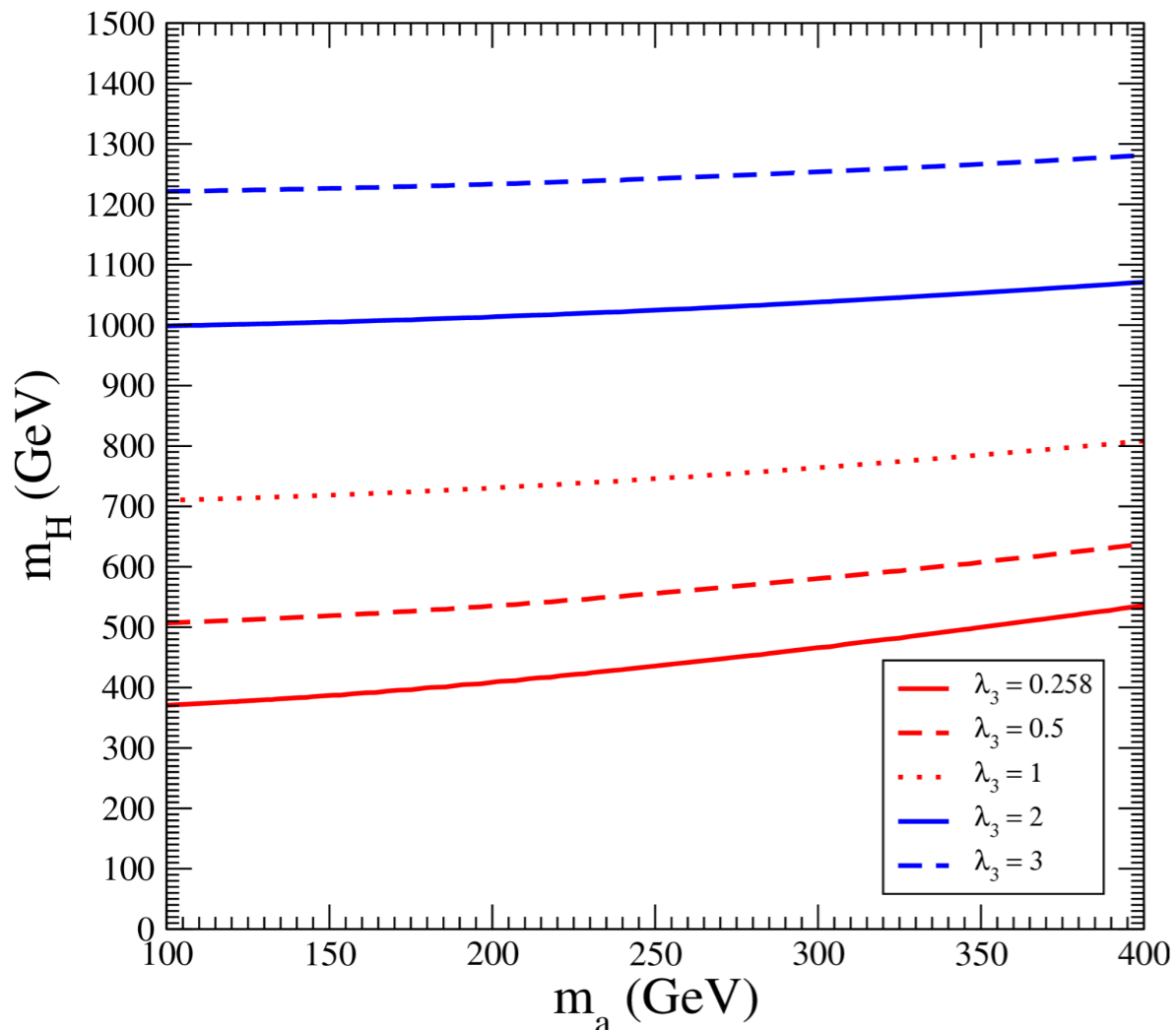
**BFB:**  $\lambda_1 > 0, \quad \lambda_2 > 0, \quad \lambda_3 > -\sqrt{\lambda_1\lambda_2}, \quad \lambda_3 + \lambda_4 - |\lambda_5| > -\sqrt{\lambda_1\lambda_2}$

$m_{H^\pm} = m_H$      $\lambda_3 \geq \frac{m_h^2}{v^2} = 0.258$

$m_A = m_H$     (Better for  $m_A > m_H$ , worse for  $m_A < m_H$ )

$\sin \theta = 0.35, \quad \tan \beta = 1$

$m_a = 100 \text{ GeV}, \quad \tan \beta = 1$



# Boundedness from Below (BFB)

$$m_{H^\pm} = m_H$$

$$m_A = m_H$$

$$\lambda_3 \geq 1 \quad (\text{e.g. } \lambda_3 \in [1, 3])$$

Increasing  $\lambda_3$  may decrease  $g_{aAh}$  (weakens mono-Higgs sensitivity)

$$g_{aAh} = \frac{c_\theta s_\theta}{m_H v} [m_h^2 + m_H^2 - m_a^2 - 2(\lambda_3 - \lambda_{P1} c_\beta^2 - \lambda_{P2} s_\beta^2)v^2]$$

**Solution** (Martin Bauer & Oleg Brandt discussion):

$$\lambda_3 = \lambda_{P1} = \lambda_{P2}$$

Impact on  $g_{haa}$  (very minor effect on  $h \rightarrow \text{Inv}$ , given  $m_a$  scan)

Choice of  $\lambda_3$  impacts  $g_{Haa}$  (effect on mono-Z)

$$g_{Haa} = \frac{1}{m_H v} [2t_{2\beta}^{-1} s_\theta^2 (m_h^2 - \lambda_3 v^2) + s_{2\beta} c_\theta^2 v^2 (\lambda_{P1} - \lambda_{P2})]$$



- Absolute Stability of EW Vacuum (MS) 1303.5098

Even if BFB, EW vacuum may not be absolute minimum

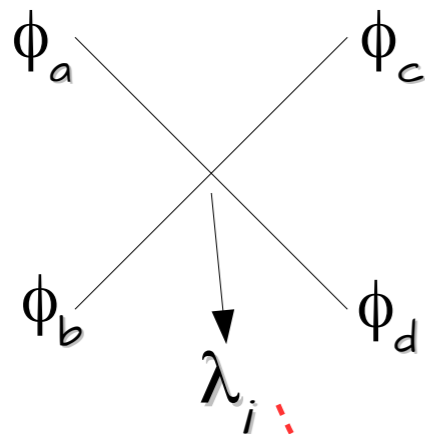
For alignment/decoupling limit, MS requires  $2m_{H^\pm}^2 - m_H^2 + m_h^2 > \lambda_3 v^2$

(for  $m_{H^\pm} = m_H$  &  $\lambda_3 = 3$ ,

$m_H > 400$  GeV OK)

- Unitarity hep-ph/0508020

2 → 2 Scatterings



Need to respect S-Matrix unitarity

Upper bound on 2HDM quartics

Upper bound on Mass Splittings

Generically less constraining than BFB

# Summary & Remarks

*“Implications from all indirect constraints - be it flavour, electroweak precision constraints or stability requirements - should be treated as preferred parameter space in a simplified model framework. It would contradict the idea of simplified models if these constraints were taken at face value.”*

Well - justified to relax an indirect constraint if impact of the analysis (e.g. modified kinematics) and/or new phenomenology arises.

Initial (set of) 2HDM+P benchmarks

$$m_{H^\pm} = m_A = m_H \quad , \quad \lambda_3 = \lambda_{P1} = \lambda_{P2} \in [1, 3] \quad ?$$

Scalar widths beyond 30% only if ~~BFB~~

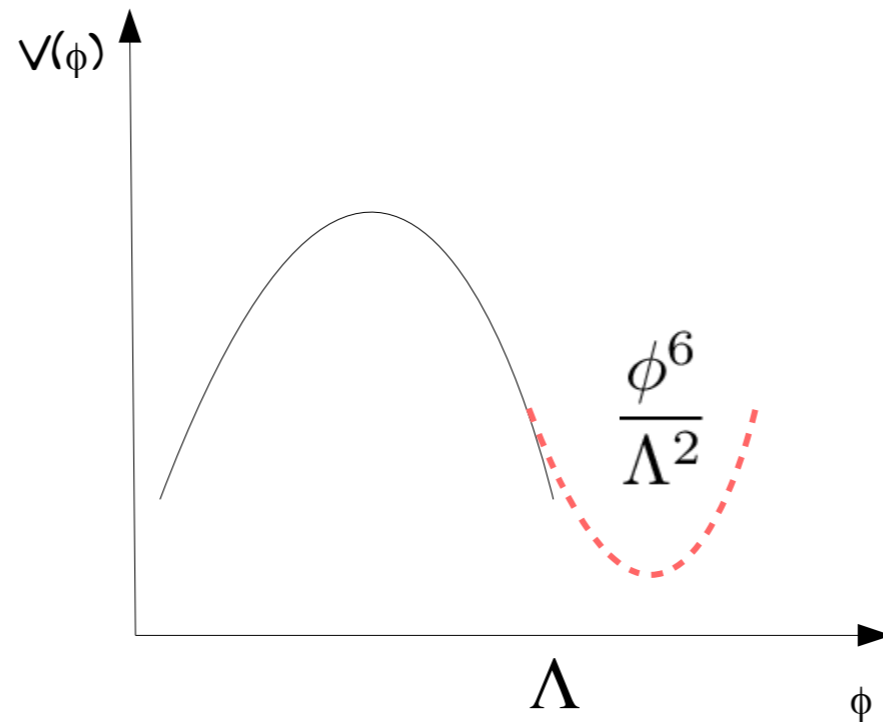
Maximum value of  $m_H$  in scan ? Minimum value of  $m_a$  ?

For later (to study in the meantime):

- Impact of mass splittings among heavy states
- Away from alignment/decoupling limit

# A (tiny) bit more on BFB...

Possible for BFB conditions to be sensitive to UV physics



This UV sensitivity of BFB generically leads to meta-stable EW vacuum