

HL/HE questions in strangeness+tau

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taus: Lusiani Pich Passemar

Rare n Strange 2017: strange physics at LHCb

GD, Lewis Tunstall, Diego Martinez Santos, Veronika Chobanova,
Xabier Cid Vidal, Francesco Dettori, Marc-Olivier Bettler

Collaboration with Teppei Kitahara arXiv:
1707.06999 PRL

Collaboration with Teppei Kitahara, Isabel Fernández Suárez, Miriam Lucio Martínez,
Diego Martínez Santos, Veronika Georgieva Chobanova

Collaboration with Crivellin, A., Kitahara, T and Nierste, U.
e-Print: arXiv:1703.05786

Collaboration with Cappiello, L. E. Greynat, D. EPJC

Collaboration with Abhishek Iyer

hyperons: Jorge is the master

Outline

- $K^0 \rightarrow \mu\mu$ Golden mode
- $K_S \rightarrow \pi\mu\mu$
- $K_S \rightarrow \pi\pi ee$
- $K_S \rightarrow \mu\mu\mu\mu$
- Questions on hyperons and taus

Rare decay modes of the K mesons in gauge theories

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(Received 4 March 1974)

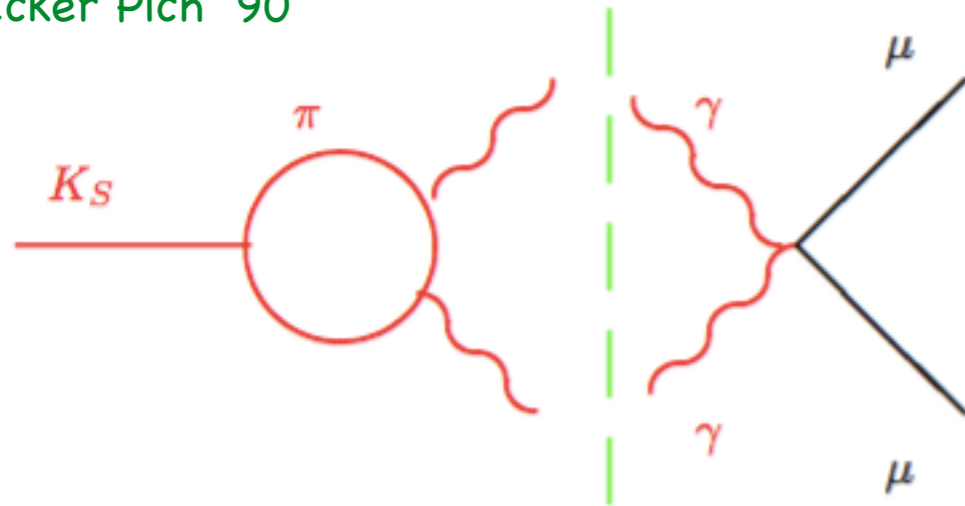
Rare decay modes of the kaons such as $K \rightarrow \mu\bar{\mu}$, $K \rightarrow \pi\nu\bar{\nu}$, $K \rightarrow \gamma\gamma$, $K \rightarrow \pi\gamma\gamma$, and $K \rightarrow \pi e\bar{e}$ are of theoretical interest since here we are observing higher-order weak and electromagnetic interactions. Recent advances in unified gauge theories of weak and electromagnetic interactions allow in principle unambiguous and finite predictions for these processes. The above processes, which are "induced" $|\Delta S|=1$ transitions, are a good testing ground for the cancellation mechanism first invented by Glashow, Iliopoulos, and Maiani (GIM) in order to banish $|\Delta S|=1$ neutral currents. The experimental suppression of $K_L \rightarrow \mu\bar{\mu}$ and nonsuppression of $K_L \rightarrow \gamma\gamma$ must find a natural explanation in the GIM mechanism which makes use of extra quark(s). The procedure we follow is the following: We deduce the effective interaction Lagrangian for $\lambda + \pi \rightarrow l + \bar{l}$ and $\lambda + \bar{\pi} \rightarrow \gamma + \gamma$ in the free-quark model; then the appropriate matrix elements of these operators between hadronic states are evaluated with the aid of the principles of conserved vector current and partially conserved axial-vector current. We focus our attention on the Weinberg-Salam model. In this model, $K \rightarrow \mu\bar{\mu}$ is suppressed due to a fortuitous cancellation. To explain the small $K_L - K_S$ mass difference and nonsuppression of $K_L \rightarrow \gamma\gamma$, it is found necessary to assume $m_{\phi}/m_{\phi'} \ll 1$, where m_{ϕ} is the mass of the proton quark and $m_{\phi'}$ the mass of the charmed quark, and $m_{\phi'} < 5$ GeV. We present a phenomenological argument which indicates that the average mass of charmed pseudoscalar states lies below 10 GeV. The effective interactions so constructed are then used to estimate the rates of other processes. Some of the results are the following: $K_S \rightarrow \gamma\gamma$ is suppressed; $K_S \rightarrow \pi\gamma\gamma$ proceeds at a normal rate, but $K_L \rightarrow \pi\gamma\gamma$ is suppressed; $K_L \rightarrow \pi\nu\bar{\nu}$ is very much forbidden, and $K^+ \rightarrow \pi^+\nu\bar{\nu}$ occurs with the branching ratio of $\sim 10^{-10}$; $K^+ \rightarrow \pi^+e\bar{e}$ has the branching ratio of $\sim 10^{-6}$, which is comparable to the presently available experimental upper bound. The predictions of other models are briefly discussed. Relevant renormalization

| VALUE (10^{-9}) | CL% | DOCUMENT ID | TECN |
|---|-----|-------------------|------------|
| < 9 | 90 | ¹ AAIJ | 2013G LHCb |
| ••• We do not use the following data for averages, fits, limits, etc. ••• | | | |
| $< 0.032 \times 10^4$ | 90 | GJESDAL | 1973 ASPK |
| $< 0.7 \times 10^4$ | 90 | HYAMS | 1969B OSPK |

¹ AAIJ 2013G uses 1.0 fb^{-1} of pp collisions at $\sqrt{s} = 7$ TeV. They obtained $B(K_S^0 \rightarrow \mu^+\mu^-) < 11 \times 10^{-9}$ at 95% C.L.

$K_S \rightarrow \mu\mu$

Ecker Pich '90



No CP conserving Short Distance due to Furry Theorem

Gaillard Lee

LD 5×10^{-12} 30% TH err

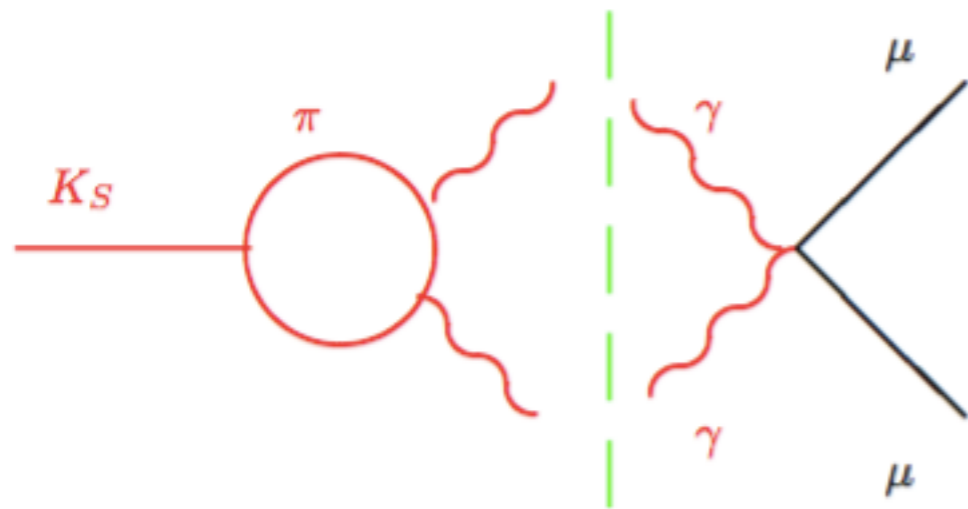
Short Distance

| | |
|----|--|
| SM | $10^{-5} \Im(V_{ts}^* V_{td}) ^2 \sim 10^{-13}$ |
| NP | few 10^{-11} allowed |

LHCb

$< 8 \times 10^{-10}$ 90%CL

$K_S \rightarrow \mu\mu$: how to improve the THEORY error



Dispersive treatment of $K_S \rightarrow \gamma\gamma$ and $K_S \rightarrow \gamma l^+ l^-$

Gilberto Colangelo, Ramon Stucki, and Lewis C. Tunstall

LD 5×10^{-12} 20% TH err

$$K_S \rightarrow \gamma\mu\mu$$

$$K_S \rightarrow \mu\mu\mu\mu$$

$$K_S \rightarrow e e \mu\mu$$

$$K_S \rightarrow \gamma\gamma$$

$K_L \rightarrow \mu\mu$

$$\cdot \Gamma(K_L^0 \rightarrow \mu^+\mu^-) / \Gamma(K_L^0 \rightarrow \pi^+\pi^-)$$

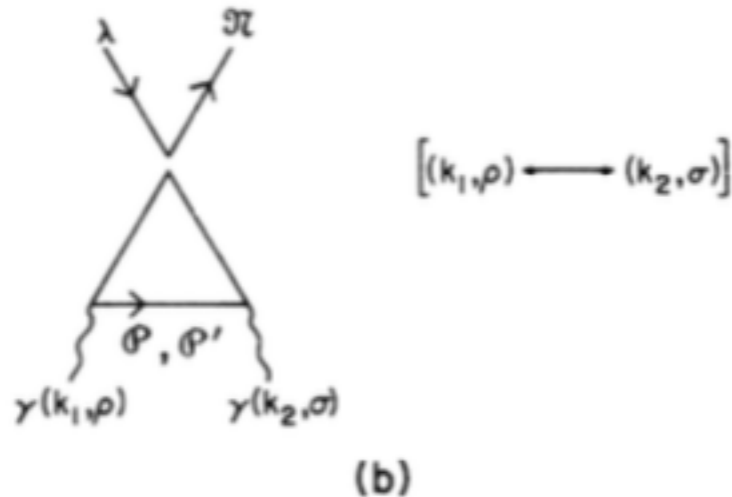
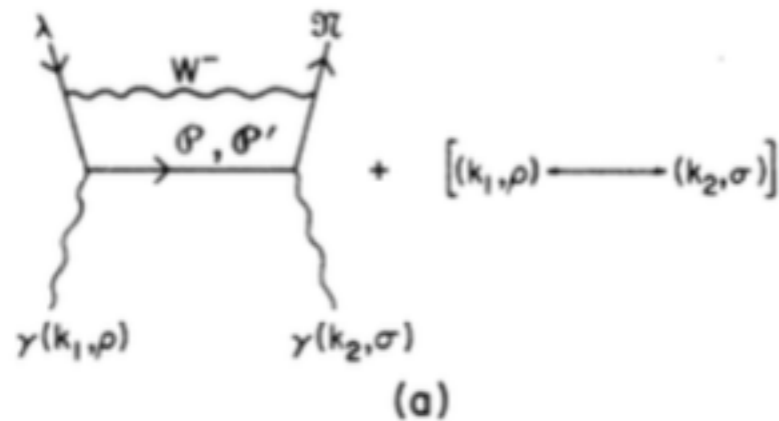


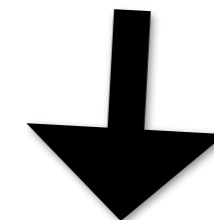
FIG. 7. Leading contributions to $\lambda + \pi^- \rightarrow \gamma + \gamma$. To leading order in M_W^{-2} , the diagrams in (a) reduce to those of (b).

Gaillard Lee

| VALUE (10^{-6}) | EVTS | DOCUMENT ID | TECN | CO |
|---|--------------------|--------------------|-------|------|
| 3.48 ± 0.05 | OUR AVERAGE | | | |
| 3.474 ± 0.057 | 6210 | AMBROSE | 2000 | B871 |
| 3.87 ± 0.30 | 179 | ¹ AKAGI | 1995 | SPEC |
| 3.38 ± 0.17 | 707 | HEINSON | 1995 | B791 |
| ... We do not use the following data for averages, fits, limits, etc. ... | | | | |
| $3.9 \pm 0.3 \pm 0.1$ | 178 | ² AKAGI | 1991B | SPEC |

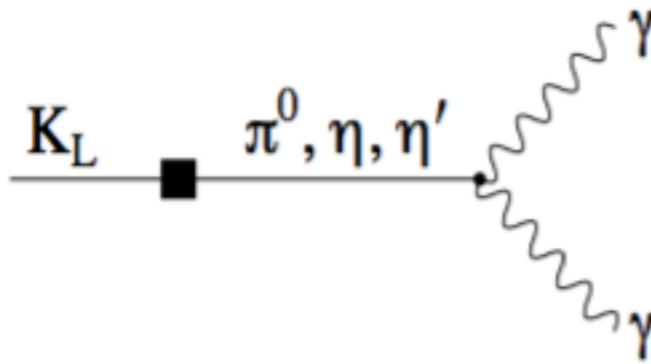
$$\mathcal{B}(K_L \rightarrow \mu^+\mu^-)_{\text{exp}} = (6.84 \pm 0.11) \times 10^{-9}$$

$K_L \rightarrow \gamma\gamma$ |_{exp} **known**



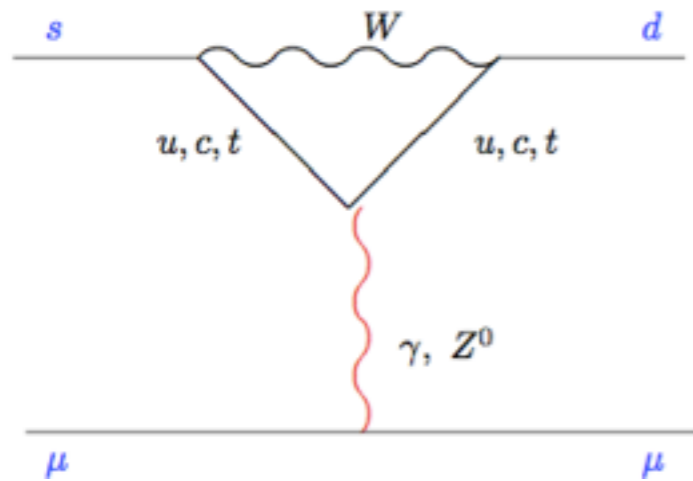
Dispersive calculation: **Re A**, Im A

We do not know the sign of $A(K_L \rightarrow \gamma\gamma)$

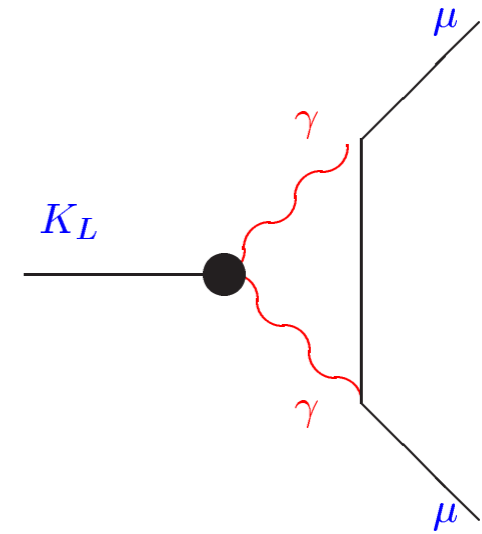


$$\begin{aligned}
 A(K_L \rightarrow 2\gamma_{\perp})_{O(p^4)} &= A(K_L \rightarrow \pi^0 \rightarrow 2\gamma_{\perp}) + A(K_L \rightarrow \eta_8 \rightarrow 2\gamma_{\perp}) \\
 &= A(K_L \rightarrow \pi^0)A(\pi^0 \rightarrow 2\gamma_{\perp}) \left[\frac{1}{M_K^2 - M_{\pi}^2} + \frac{1}{3} \cdot \frac{1}{M_K^2 - M_8^2} \right] \simeq 0
 \end{aligned}$$

$K_L \rightarrow \mu\mu$



\ll



$$\frac{\Gamma(K_L \rightarrow \mu\bar{\mu})}{\Gamma(K_L \rightarrow \gamma\gamma)} \sim |ReA|^2 + |ImA|^2$$

Absorptive calculation
model independent

27.14

Subtracting from expt. the Absorptive contribution

$$0.98 \pm 0.55 = |ReA|^2 = (\chi_{\gamma\gamma}(M_\rho) + \chi_{\text{short}} - 5.12)^2$$

$$|\chi_{\text{short}}^{\text{SM}}| = 1.96(1.11 - 0.92\bar{\rho})$$

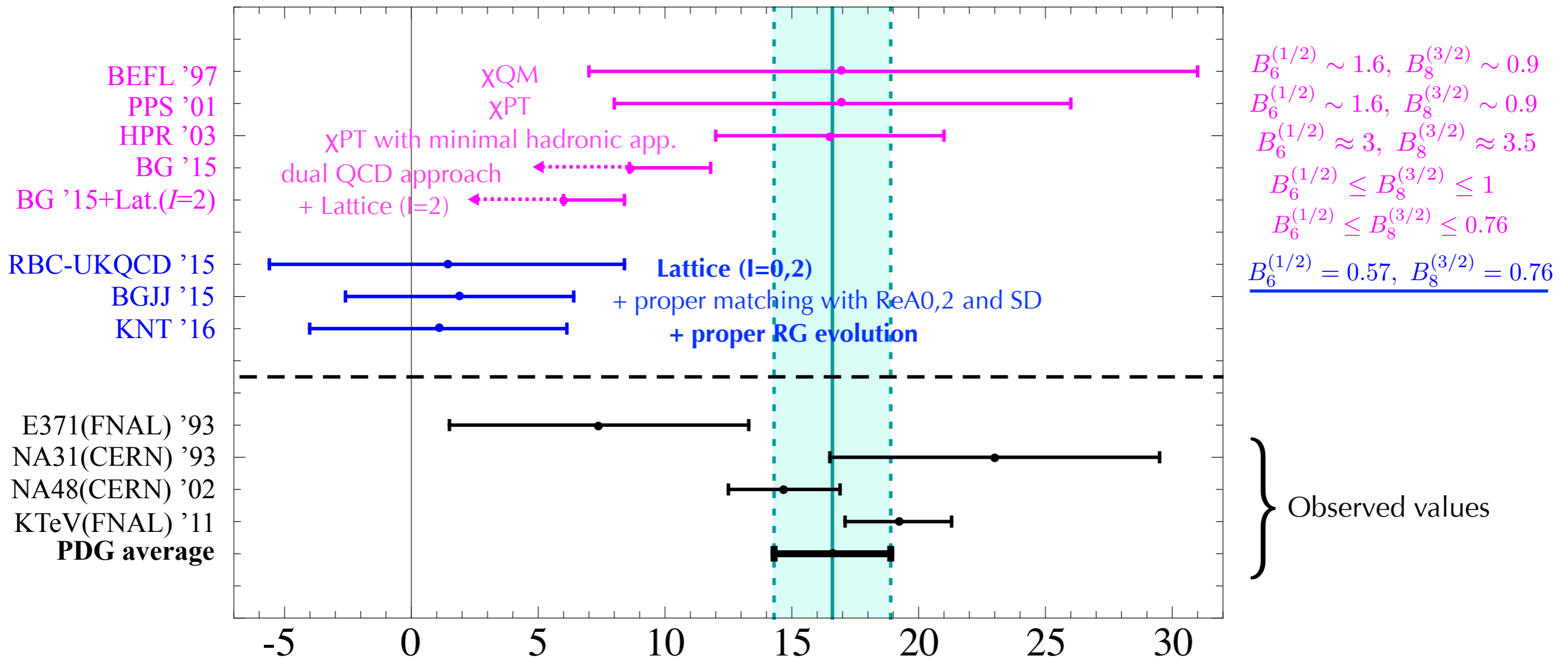
$$\varepsilon'/\varepsilon = (1.4 \pm 7.0) \cdot 10^{-4}$$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right) = 31.0 \pm 6.6$$

$$\left(\varepsilon'/\varepsilon \right)_{\text{exp}} = (16.6 \pm 2.3) \cdot 10^{-4}$$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{\text{exp}} = 22.4$$

Current situation of $\epsilon'_K / \epsilon_K \propto \text{Im}A_0 - \left(\frac{\text{Re}A_0}{\text{Re}A_2}\right) \text{Im}A_2$



$B_6^{(1/2)} \sim 1.6, B_8^{(3/2)} \sim 0.9$
 $B_6^{(1/2)} \sim 1.6, B_8^{(3/2)} \sim 0.9$
 $B_6^{(1/2)} \approx 3, B_8^{(3/2)} \approx 3.5$
 $B_6^{(1/2)} \leq B_8^{(3/2)} \leq 1$
 $B_6^{(1/2)} \leq B_8^{(3/2)} \leq 0.76$
 $B_6^{(1/2)} = 0.57, B_8^{(3/2)} = 0.76$

} Observed values

large N limit (convention)

$$B_6^{(1/2)} = B_8^{(3/2)} = 1$$

dual QCD prediction

$$B_6^{(1/2)} \leq B_8^{(3/2)} < 1, B_8^{(3/2)} = 0.8$$

$$\text{Re } \epsilon'_K / \epsilon_K \times 10^4$$

| | Exp. | χ PT | dual QCD | Lattice (I=0,2) |
|--|------------------|-----------|----------------|-----------------|
| $\left(\frac{\text{Re}A_0}{\text{Re}A_2}\right)$ | 22.45 ± 0.05 | ~ 14 | 16.0 ± 1.5 | 31.0 ± 11.1 |

Models solving ϵ'/ϵ anomaly

- Several new physics models have been studied to explain ϵ'/ϵ anomaly

| | |
|----------------------------------|--|
| MSSM -- chargino Z penguin | <i>[M. Endo, S. Mishima, D. Ueda and KY, PLB762(2016)493]</i> |
| -- gluino Z penguin | <i>[M. Tanimoto and KY, PTEP(2016)no.12,123B02]</i> |
| -- gluino box | <i>[T.Kitahara, U.Nierste and P.Trempfer, PRL117(2016)no.9, 091802 A.Crivellin, G.D'Ambrosio, T.Kitahara and U.Nierste, 1703.05786]</i> |
| Vector-like quarks | <i>[C.Bobeth, A.J.Buras, A.Celis and M.Jung, JHEP1704(2017)079]</i> |
| Little Higgs Model with T-parity | <i>[M.Blanke, A.J.Buras and S.Recksiegel, EPJ.C76 (2016)no.4,182]</i> |
| 331 model | <i>[A.J.Buras and F.De Fazio, JHEP1603(2016)010 & JHEP1608 (2016) 115]</i> |
| Right handed current | <i>[V.Cirigliano, W.Dekens, J.de Vries and E.Mereghetti, PLB 767 (2017) 1 S.Alioli, V.Cirigliano, W.Dekens, J.de Vries and E.Mereghetti, JHEP1705 (2017)086]</i> |

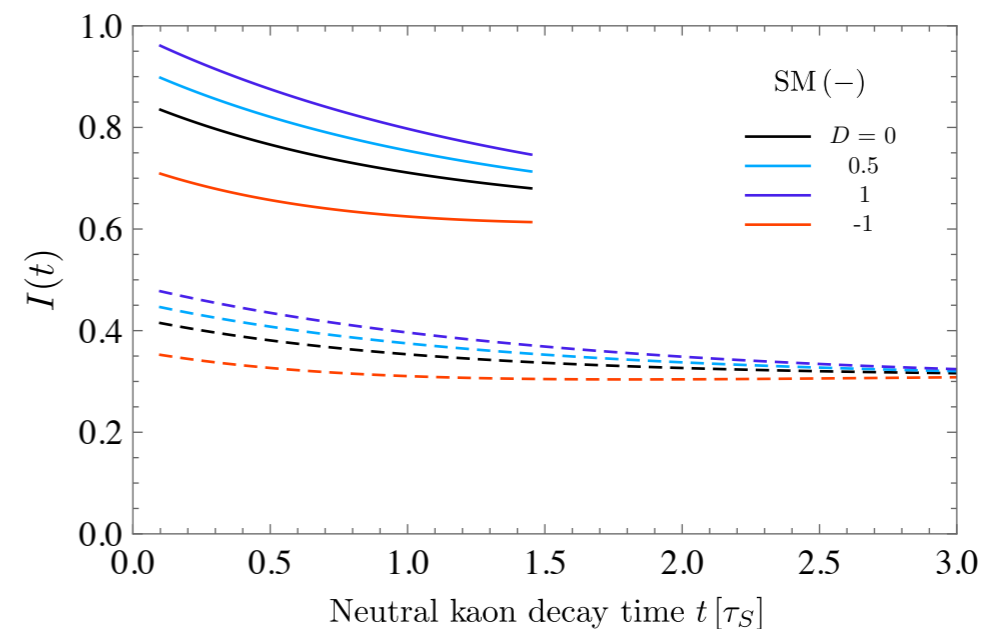
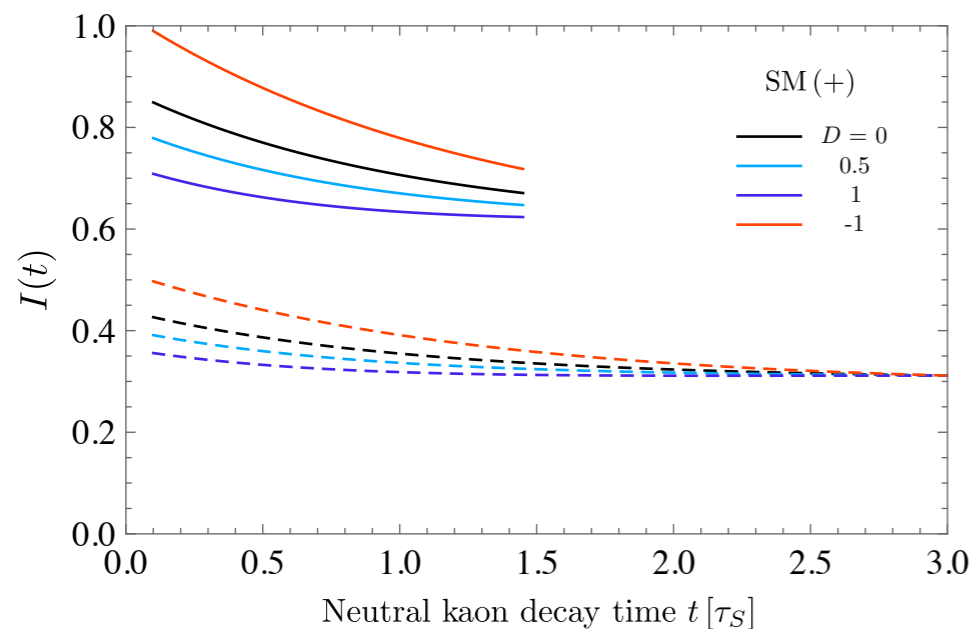
- Different implications (correlations & predictions) for other observables appear depending on models \Rightarrow Possibility of model discriminations

Can we study $K^0(t)$?

GD, Kitahara
1707.06999 PRL

$$pp \rightarrow K^0 K^- X$$

$$pp \rightarrow K^{*+} X \rightarrow K^0 \pi^+ X$$



$$|\bar{K}^0(t)\rangle = \frac{1}{\sqrt{2}(1 \pm \bar{\epsilon})} \left[e^{-iH_S t} (|K_1\rangle + \bar{\epsilon}|K_2\rangle) \right. \\ \left. \pm e^{-iH_L t} (|K_2\rangle + \bar{\epsilon}|K_1\rangle) \right]$$

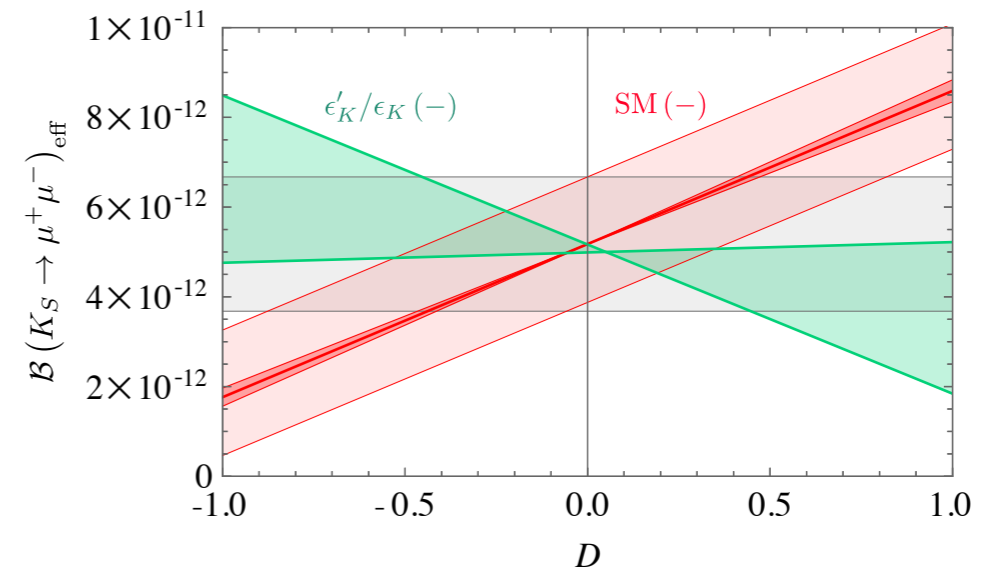
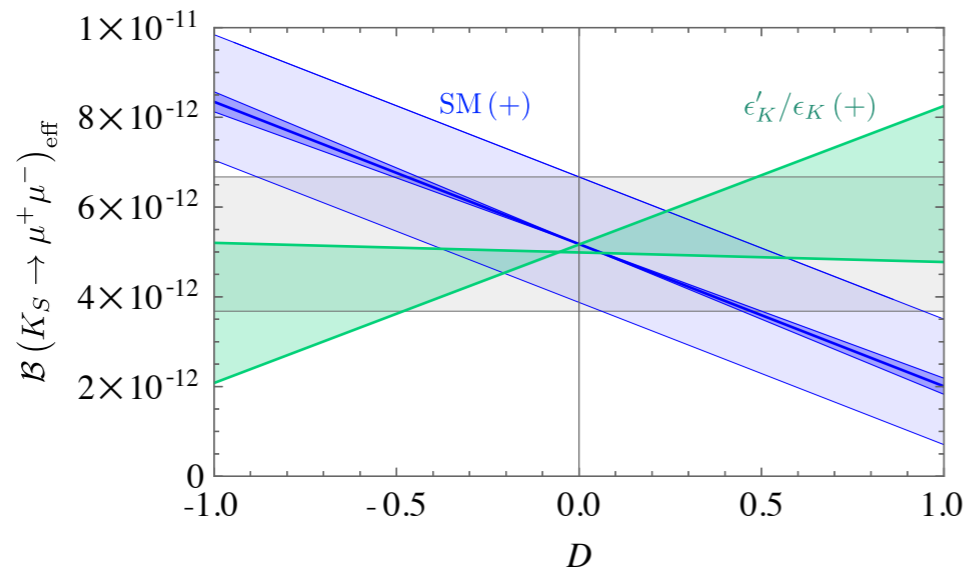
$$D = \frac{K^0 - \bar{K}^0}{K^0 + \bar{K}^0}$$

- **Short distance interfering** with Large CP conserving LD contribution !
- We may be able to study the time evolution of K^0 by tracking the associated particles (K^-)

$$\sum_{\text{spin}} \mathcal{A}(K_1 \rightarrow \mu^+ \mu^-)^* \mathcal{A}(K_2 \rightarrow \mu^+ \mu^-)$$

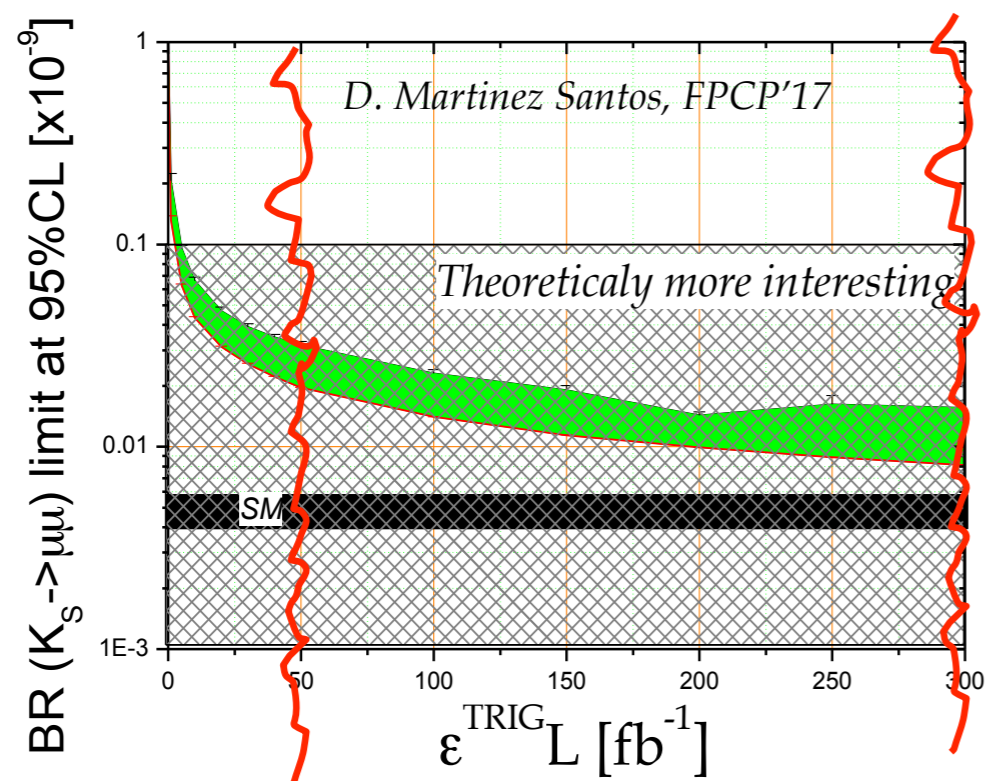
$$\sim \text{Im}[\lambda_t] y'_{7A} \left\{ A_{L\gamma\gamma}^\mu - 2\pi \sin^2 \theta_W (\text{Re}[\lambda_t] y'_{7A} + \text{Re}[\lambda_c] y_c) \right\}$$

Short distance window



$$\begin{aligned} & \mathcal{B}(K_S \rightarrow \mu^+ \mu^-)_{\text{eff}} \\ &= \tau_S \left[\int_{t_{\min}}^{t_{\max}} dt \left(\Gamma(K_1) e^{-\Gamma_S t} + \frac{D}{8\pi M_K} \sqrt{1 - \frac{4m_\mu^2}{M_K^2}} \sum_{\text{spin}} \text{Re} \left[e^{-i\Delta M_K t} \mathcal{A}(K_1)^* \mathcal{A}(K_2) \right] e^{-\frac{\Gamma_S + \Gamma_L}{2} t} \right) \varepsilon(t) \right] \\ & \times \left(\int_{t_{\min}}^{t_{\max}} dt e^{-\Gamma_S t} \varepsilon(t) \right)^{-1}, \end{aligned}$$

$K_S \rightarrow \mu \mu$ prospects



LHCb-upgrade

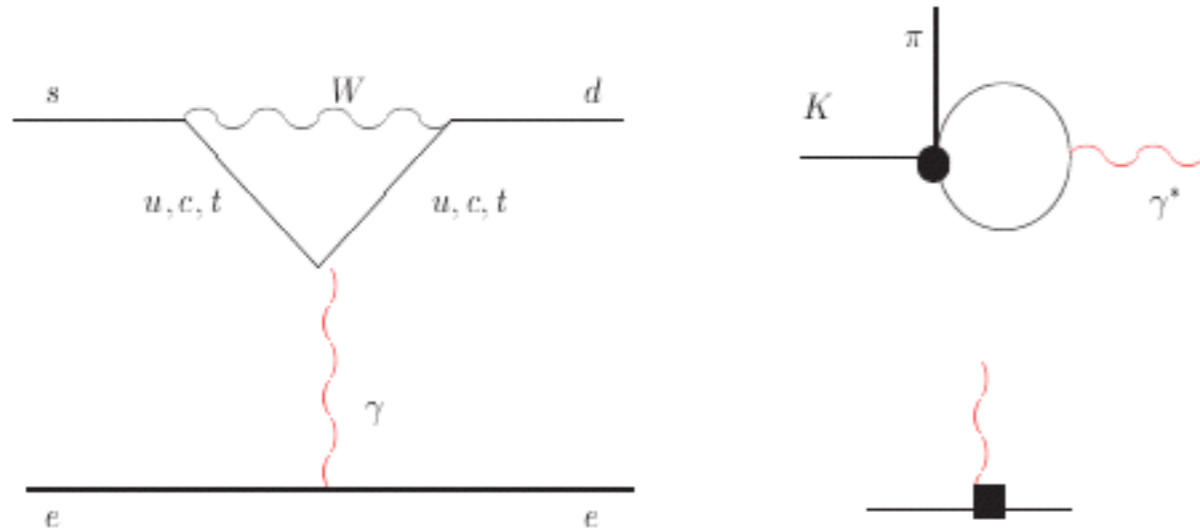
Phase-II-upgrade? <https://arxiv.org/abs/1707.06999>

- Extrapolating from Run-I result
- The most interesting region can be achieved by LHCb upgrade with trigger improvements
- Starting to investigate tagged decays, which would allow to access NP in the K_S - K_L interference
[D'Ambrosio&Kitahara <https://arxiv.org/abs/1707.06999>]

$$K^\pm(K_S) \rightarrow \pi^\pm(\pi^0)\ell^+\ell^-$$

- short distance \ll long distance

LD described by form factor W



$$W^i = G_F m_K^2 (a_i + b_i z) + W_{\pi\pi}^i(z)$$

$$i = \pm, S$$

$$a_i, b_i \sim O(1), \quad z = \frac{q^2}{m_K^2}$$

- Observables $\Gamma(K^+ \rightarrow \pi^+ e^+ e^-)$, $\Gamma(K^+ \rightarrow \pi^+ \mu^+ \mu^-)$, slopes

- $a_i \sim O(p^4)$ $a_+ \sim N_{14} - N_{15}$, $a_S \sim 2N_{14} + N_{15}$

Ecker, Pich, de Rafael

- $b_i \sim O(p^6)$

G.D., Ecker, Isidori, Portoles

- a_+, b_+ in general not related to a_S, b_S

averaging flavour

$$a_+^{\text{exp.}} = -0.578 \pm 0.016$$

$$b_+^{\text{exp.}} = -0.779 \pm 0.066$$

$K_S \rightarrow \pi^0 l^+ l^-$ at **NA48/1 Collaboration at CERN**

- $K_S \rightarrow \pi^0 e^+ e^-$ **7** evts observed (with 0.15 expected bkg evts)

$$B(K_S \rightarrow \pi^0 e^+ e^-)_{m_{ee} > 165 \text{ MeV}} = (3.0_{-1.2}^{+1.5} \pm 0.2) \times 10^{-9}$$

$$|a_S| = 1.08_{-0.21}^{+0.26}$$

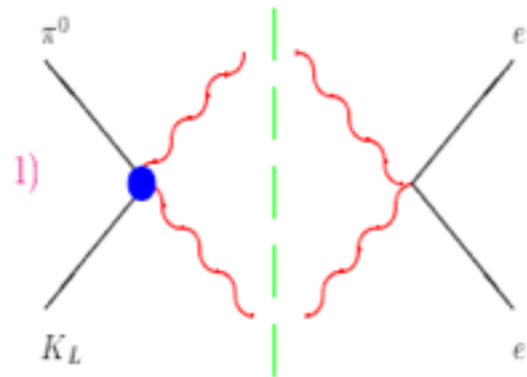
- $K_S \rightarrow \pi^0 \mu^+ \mu^-$ **6** events observed

$$B(K_S \rightarrow \pi^0 \mu^+ \mu^-) = (2.9_{-1.2}^{+1.5} \pm 0.2) \times 10^{-9}$$

$$|a_S|_{\mu\mu} = 1.54_{-0.32}^{+0.40} \pm 0.06$$

$K_L \rightarrow \pi^0 e^+ e^-$: summary

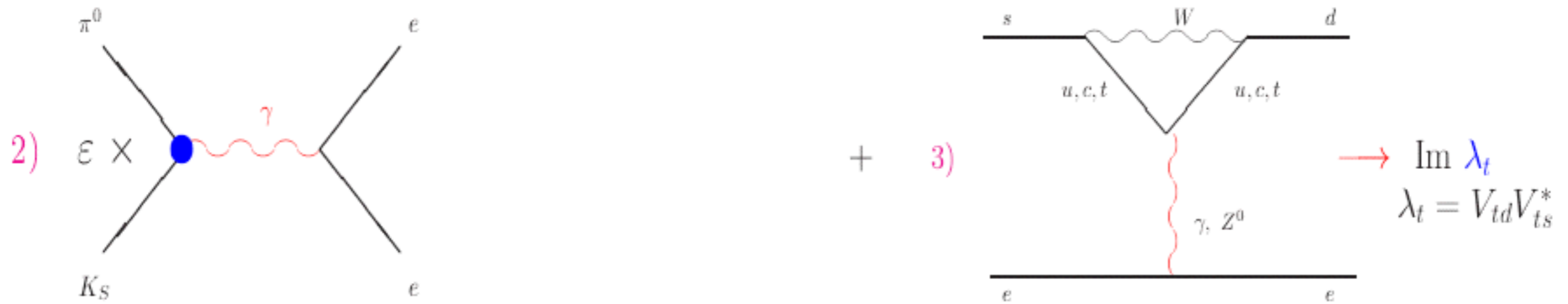
$$\text{Br}(K_L \rightarrow \pi^0 e^+ e^-) \leq 2.8 \cdot 10^{-10} \text{ at 90\% CL} \quad \text{KTeV}$$



CP conserving NA48

$$\text{Br}(K_L \rightarrow \pi^0 e^+ e^-) < 3 \cdot 10^{-12}$$

$$V-A \otimes V-A \Rightarrow \langle \pi^0 e^+ e^- | (\bar{s}d)_{V-A} (\bar{e}e)_{V-A} | K_L \rangle \text{ violates CP}$$



$$\uparrow \text{B}(K_S \rightarrow \pi^0 e^+ e^-) = 4.6 a_S^2 \times 10^{-9}$$

Possible large interference: $a_S < -0.5$ or $a_S > 1$; short distance probe even for a_S large

$$|2) + 3)|^2 = \left[15.3 a_S^2 - 6.8 \frac{\text{Im} \lambda_t}{10^{-4}} a_S + 2.8 \left(\frac{\text{Im} \lambda_t}{10^{-4}} \right)^2 \right] \cdot 10^{-12}$$

$$[17.7 \pm \quad 9.5 + \quad 4.7] \cdot 10^{-12}$$

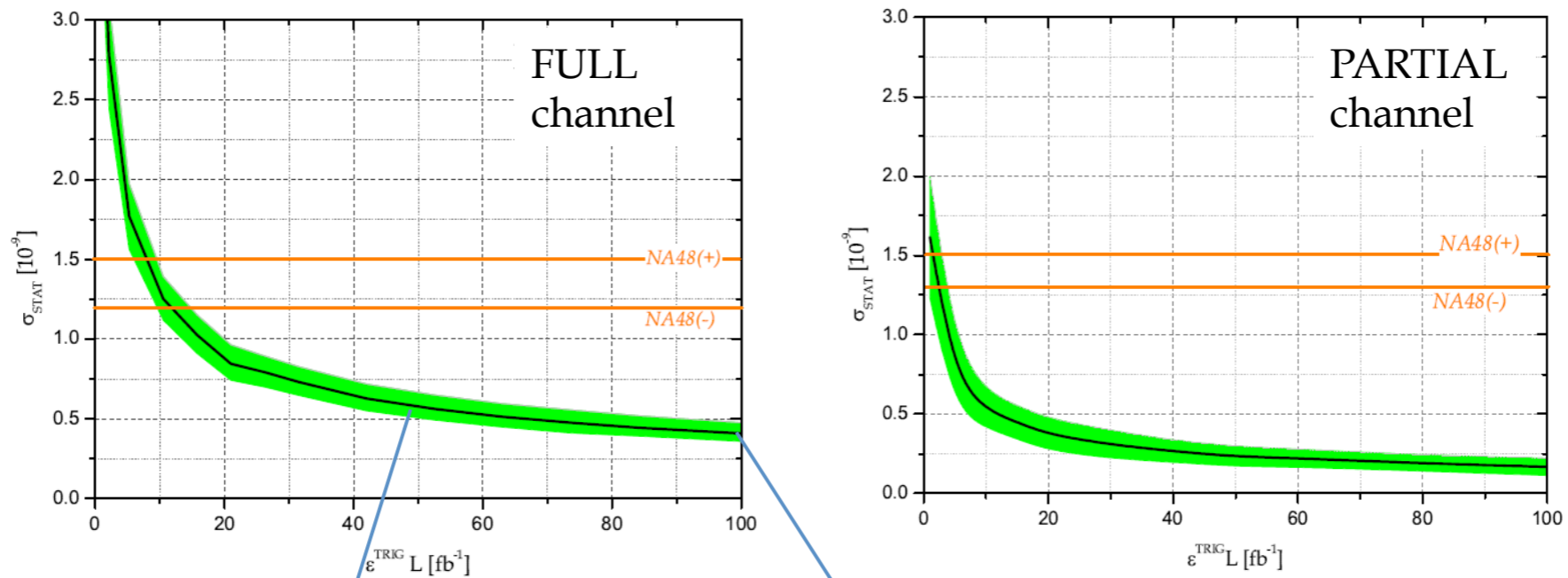
$K_S \rightarrow \pi^0 \mu \mu$ sensitivity study

CERN-LHCb-PUB-2016-017

$K_S \rightarrow \pi^0 \mu \mu$ sensitivity study

V. Chobanova et al,
CERN-LHCb-PUB-2016-017

Normalize yields to $K_S \rightarrow \pi \pi \rightarrow$ sensitivity to BR



LHCb Upgrade, 100 % eff

LHCb HL-LHC
100 fb⁻¹ 100 % eff
or 300 fb⁻¹ 33 % eff

$K_S \rightarrow \pi^0 \mu\mu$ short distance?

$$A_{FB}^{K_S^{\pi^0 \mu\mu}}(z) = \frac{\Gamma(\cos \theta_{K\mu} > 0) - \Gamma(\cos \theta_{K\mu} < 0)}{\Gamma(\cos \theta_{K\mu} > 0) + \Gamma(\cos \theta_{K\mu} < 0)}$$

- sensitive to short distance physics
- several background to control $(\epsilon_K \quad |a_S|^2)$

$$K_S \rightarrow \pi^+ \pi^- \mu \mu \sim 10^{-14}$$

$$K_S \rightarrow \pi^+ \pi^- e e \sim 10^{-5}$$

• $\Gamma(K_S^0 \rightarrow \pi^+ \pi^- e^+ e^-) / \Gamma_{\text{total}}$

| VALUE (10^{-5}) | EVTS | DOCUMENT ID | TECN | COMMENT |
|---------------------|--------------------|---------------------|------------|----------------|
| 4.79 ± 0.15 | OUR AVERAGE | | | |
| 4.83 ±0.11 ±0.14 | 23k | ¹ BATLEY | 2011 NA48 | 2002 data |
| 4.69 ±0.30 | 676 | ² LAI | 2003C NA48 | 1998+1999 data |

• • • We do not use the following data for averages, fits, limits, etc. • • •

$$K(p_K) \rightarrow \pi(p_1)\pi(p_2)\gamma(q)$$

- Lorentz + gauge invariance \Rightarrow Electric (E) and Magnetic (M) amplitude

$$A(K \rightarrow \pi\pi\gamma) = F^{\mu\nu} [E \partial_\mu K \partial_\nu \pi + M \varepsilon_{\mu\nu\rho\sigma} \partial^\rho K \partial^\sigma \pi]$$

- Unpolarized photons

$$\frac{d^2\Gamma}{dz_1 dz_2} \sim |E|^2 + |M|^2$$

$$|E^2| = |E_{IB}|^2 + 2\text{Re}(E_{IB}^* E_D) + |E_D|^2$$

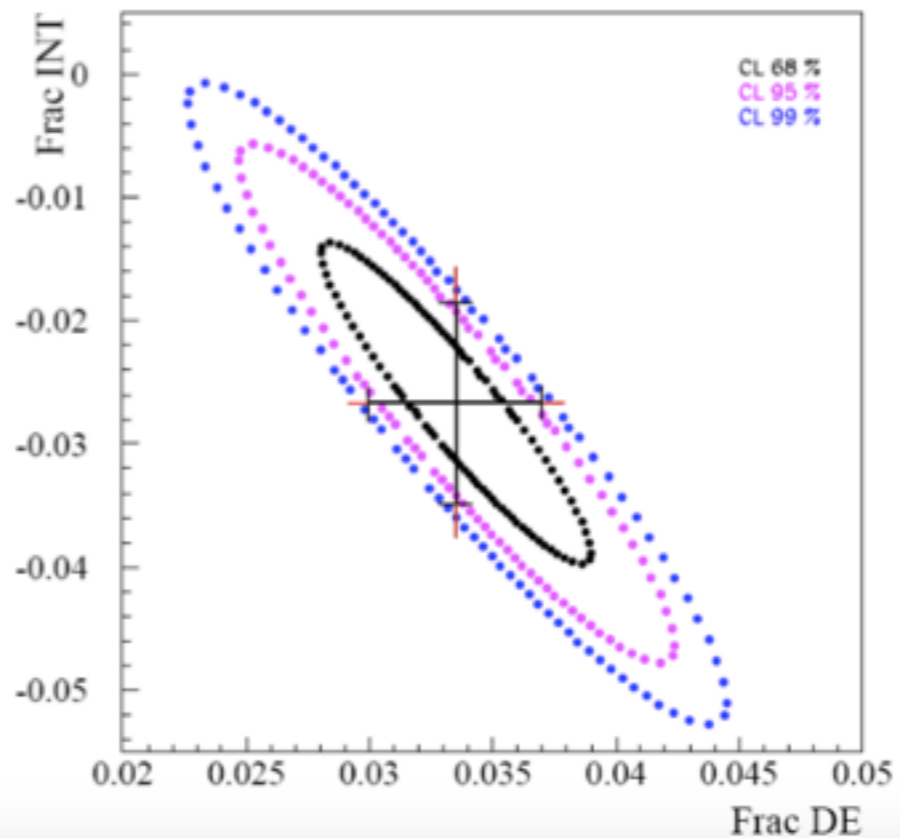
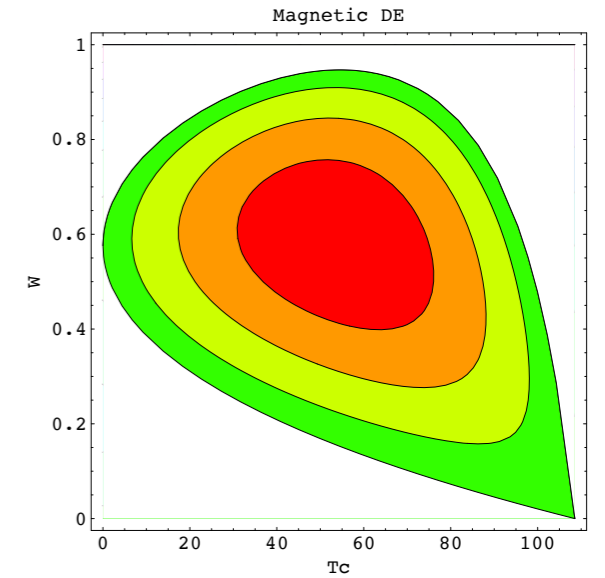
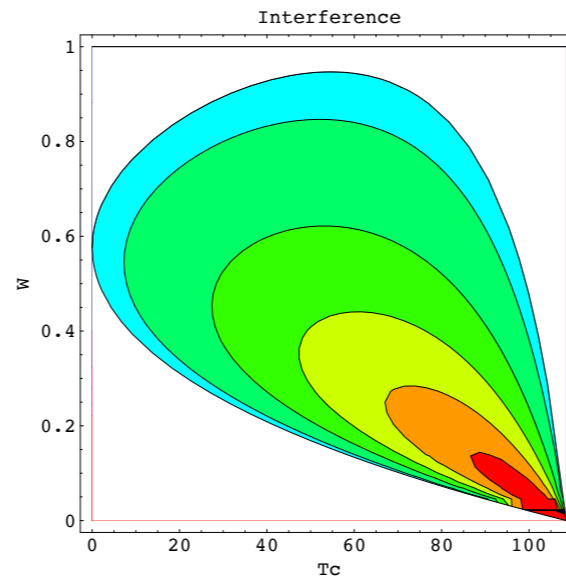
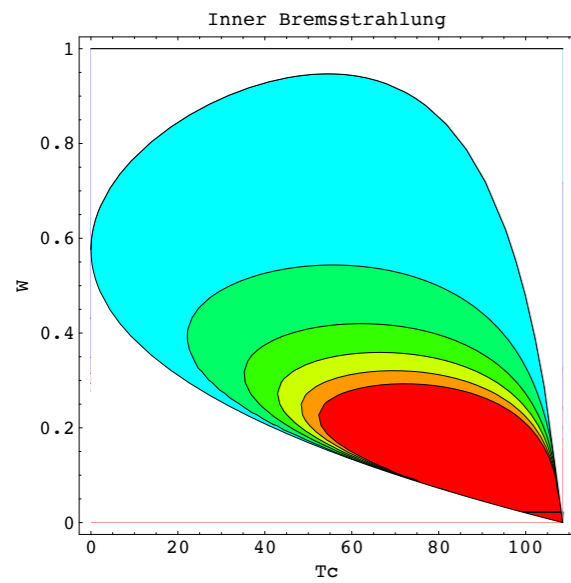
↓

$$\text{Low Theorem} \Rightarrow E_{IB} \sim \frac{1}{E_\gamma^*} + c \qquad E_D, M \text{ chiral}$$

tests

$$K^+ \rightarrow \pi^+ \pi^0 \gamma$$

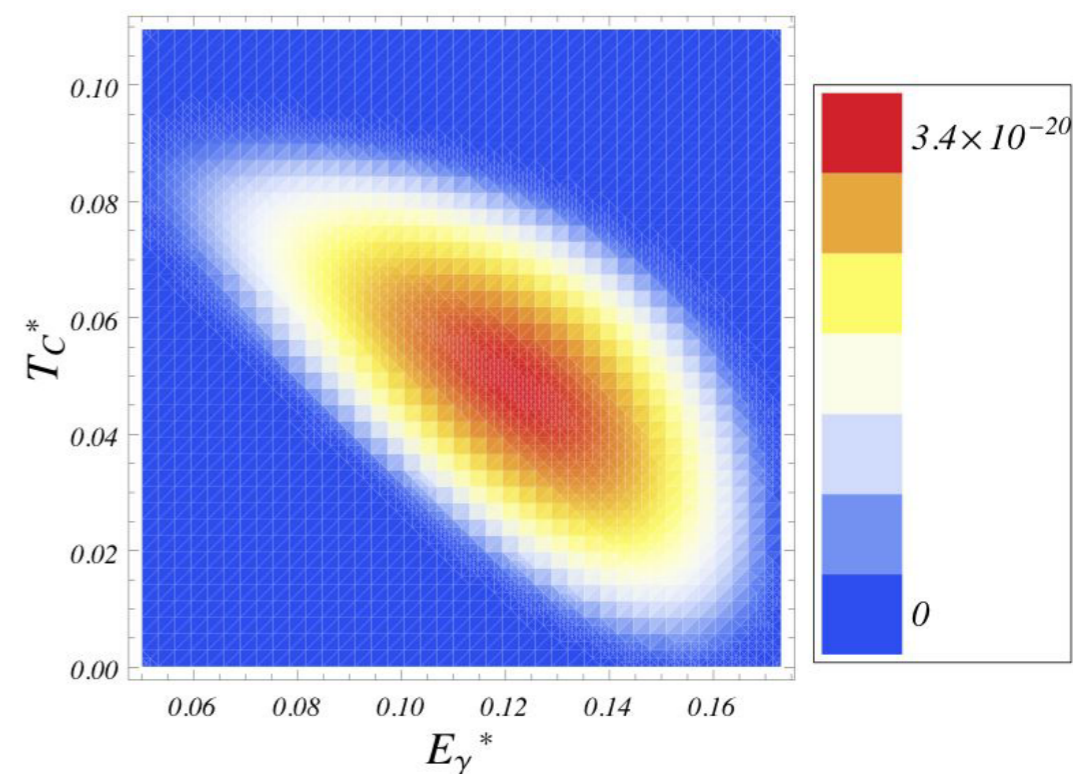
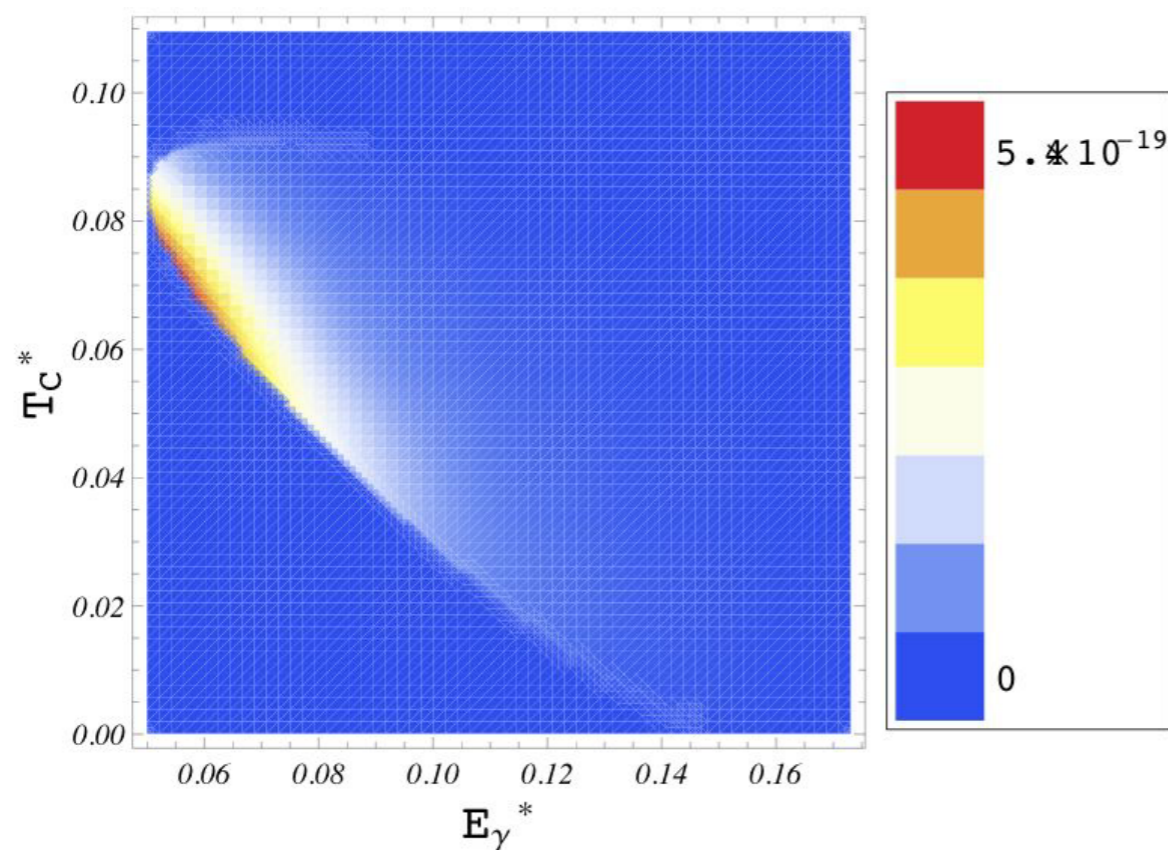
Dalitz plot NA48/2



$$K^+ \rightarrow \pi^+ \pi^0 e e$$

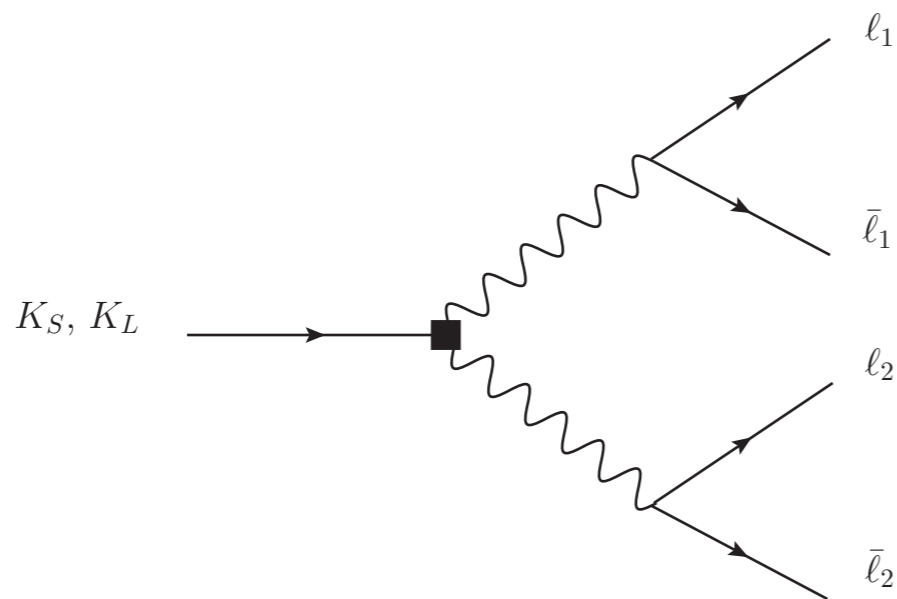
Starting from CP conserving IB, DE

| q_c (MeV) | B [10^{-8}] | B/M | B/E | B/BE | B/BM |
|-------------|-----------------|-----|------|------|------|
| $2m_l$ | 418.27 | 71 | 4405 | 128 | 208 |
| 55 | 5.62 | 12 | 118 | 38 | 44 |
| 100 | 0.67 | 8 | 30 | 71 | 36 |
| 180 | 0.003 | 12 | 5 | -19 | 44 |



Interesting channels

| | | | |
|--------------------------------|---|-------|--------------------------|
| $K_S \rightarrow \mu\mu\mu\mu$ | — | SM LD | $\sim 2 \times 10^{-14}$ |
| $K_S \rightarrow e e \mu \mu$ | — | | $\sim 10^{-11}$ |
| $K_S \rightarrow e e e e$ | — | | $\sim 10^{-10}$ |



GD, Greynat, Vulvert

Hyperons

- Hyperon semileptonic decays/ V_{us} determination
- Search for resonance (Hypercp) $\Sigma \rightarrow p\mu\mu$

$\Xi^- \rightarrow \Lambda \mu^- \bar{\nu}_\mu$

$\Gamma(\Xi^- \rightarrow \Lambda \mu^- \bar{\nu}_\mu) / \Gamma(\Xi^- \rightarrow \Lambda \pi^-)$

| VALUE (10^{-3}) | CL% | EVTS | DOCUMENT ID | TECN | COMMENT |
|--|-----|------|---------------|------|-----------------------|
| $0.35^{+0.35}_{-0.22}$ | | | | | OUR FIT |
| 0.35 ± 0.35 | | 1 | YEH 1974 | HBC | Effective denom.=2859 |
| *** We do not use the following data for averages, fits, limits, etc *** | | | | | |
| <2.3 | 90 | 0 | THOMPSON 1980 | ASPK | Effective denom.=1017 |
| <1.3 | | | DAUBER 1969 | HBC | |
| <12 | | | BERGE 1966 | HBC | |

References

THOMPSON 1980 PR D21 25 Studies in the BNL 21 GeV/c Negative Hyperon Beam.

YEH 1974 PR D10 3545 Observation of Rare Decay Modes of the Ξ Hyperons

DAUBER 1969 PR 179 1262 Production and Decay of Cascade Hyperons

BERGE 1966 PR 147 945 Some Properties of Ξ^- and Ξ^0 Hyperons Produced in 1.7 GeV/c

see fit info

$\Lambda \rightarrow p \mu^- \bar{\nu}_\mu$

$\Gamma(\Lambda \rightarrow p \mu^- \bar{\nu}_\mu) / \Gamma(\Lambda \rightarrow N \pi)$

| VALUE (10^{-4}) | EVTS | DOCUMENT ID | TECN | COMMENT |
|---------------------|------|---------------|------|--------------------|
| 1.57 ± 0.35 | | | | OUR FIT |
| 1.57 ± 0.35 | | | | OUR AVERAGE |
| 1.4 ± 0.5 | 14 | BAGGETT 1972B | HBC | $K^- p$ at rest |
| 2.4 ± 0.8 | 9 | CANTER 1971B | HBC | $K^- p$ at rest |
| 1.3 ± 0.7 | 3 | LIND 1964 | RVUE | |
| 1.5 ± 1.2 | 2 | RONNE 1964 | FBC | |

References

BAGGETT 1972B ZPHY 252 362 Measurement of the $\Lambda \rightarrow p \mu^- \bar{\nu}_\mu$ Branching Ratio

CANTER 1971B PRL 27 59 Branching Ratio for $\Lambda \rightarrow p \mu^- \bar{\nu}_\mu$

LIND 1964 PR 135 B1483 Experimental Investigation of $V - A$ in Leptonic Λ Decay

RONNE 1964 PL 11 357 Branching Ratio for the Muonic Decay Mode of the Λ Hyperon

see fit info

- Most data in the μ -channel is very old (60's and 70's)! $\delta\text{Br}/\text{Br} \sim 10\% - 100\%$

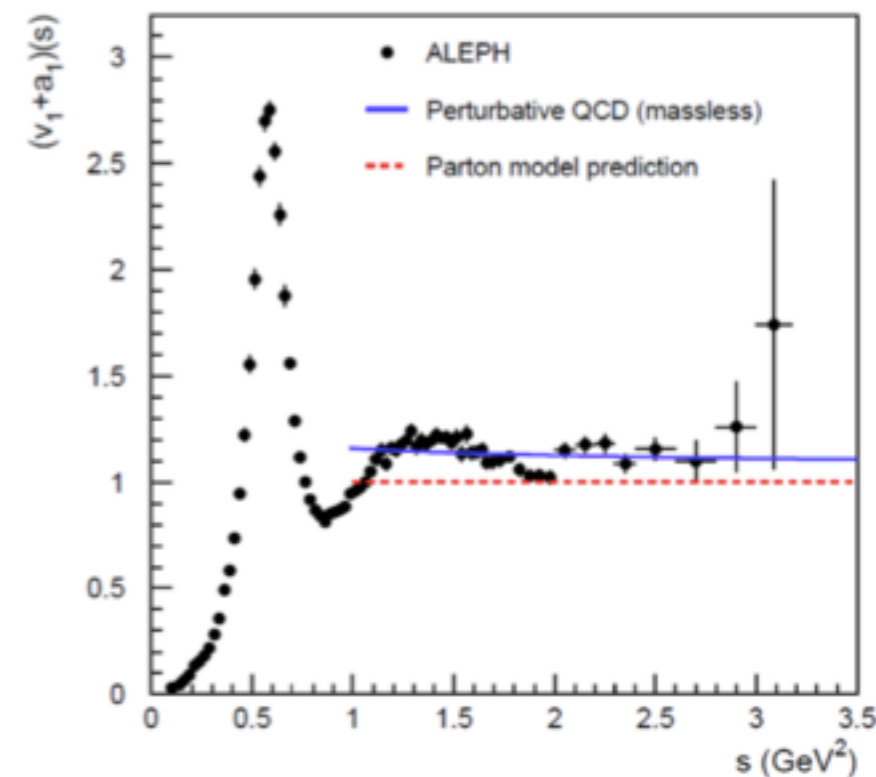
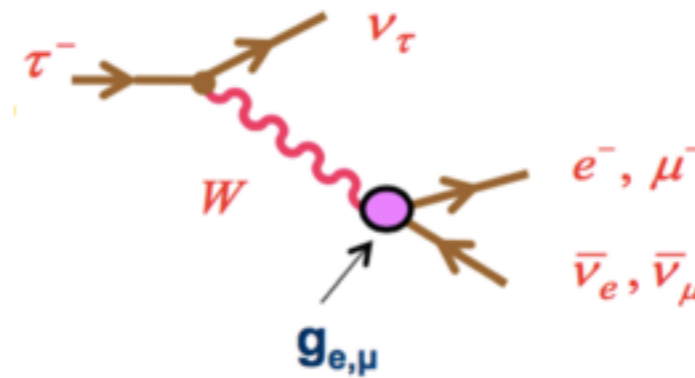
| | $\Lambda \rightarrow p$ | $\Sigma^- \rightarrow n$ | $\Xi^0 \rightarrow \Sigma^+$ | $\Xi^- \rightarrow \Lambda$ |
|--------|-------------------------|--------------------------|------------------------------|-----------------------------|
| Expt. | 0.189(41) | 0.442(39) | 0.0092(14) | 0.6(5) |
| SM-NLO | 0.153(8) | 0.444(22) | 0.0084(4) | 0.275(14) |

- Good agreement between SM and data \Rightarrow Bounds on ϵ_S^{sl} and ϵ_T^{sl}

Possible tau issues

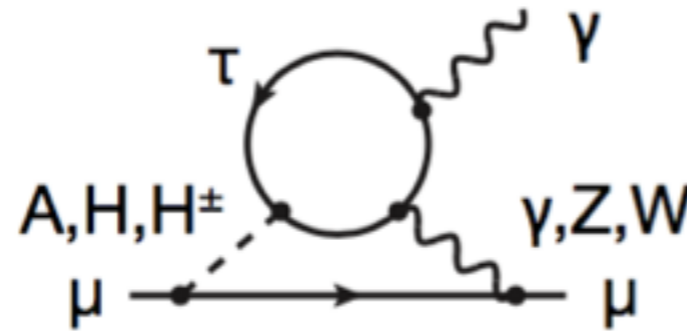
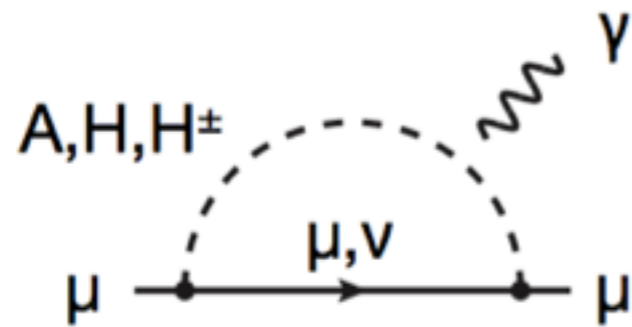
- Leptonic universality
- QCD tests
- V_{us}
- Lepton flavour universality

| Experiment | Number of τ pairs |
|------------|------------------------|
| LEP | $\sim 3 \times 10^5$ |
| CLEO | $\sim 1 \times 10^7$ |
| BaBar | $\sim 5 \times 10^8$ |
| Belle | $\sim 9 \times 10^8$ |
| Belle II | $\sim 10^{12}$ |



- The lepton universality tests give strong constraints on type-X (lepton-specific) 2HDMs → Model favoured to explain the g-2 discrepancy

M. Endo@b2tip'15



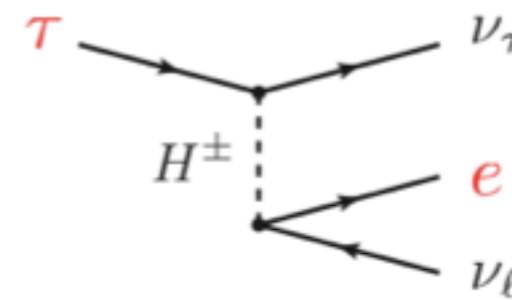
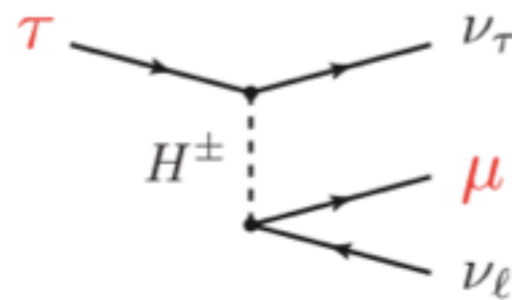
Barr-Zee contribution with A enhances muon g-2

→ Contribution to LU :

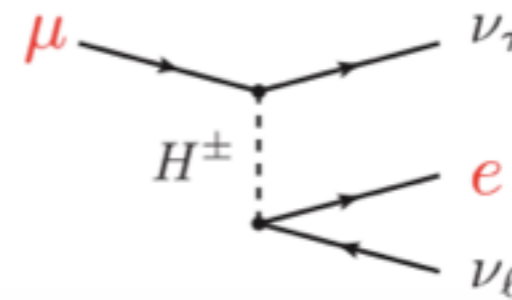
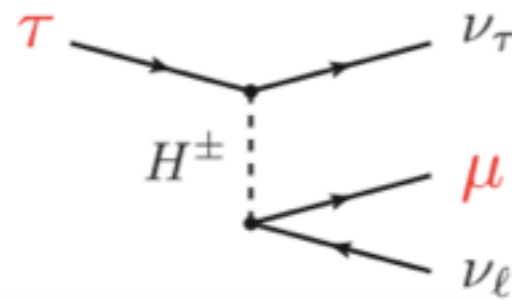
Large

Negligible

$$\frac{g_\mu}{g_e} = \frac{1 + \Delta_{\tau \rightarrow \mu}}{1 + \Delta_{\tau \rightarrow e}} = 1.0018 \pm 0.0014$$



$$\frac{g_\tau}{g_e} = \frac{1 + \Delta_{\tau \rightarrow \mu}}{1 + \Delta_{\mu \rightarrow e}} = 1.0029 \pm 0.0015$$



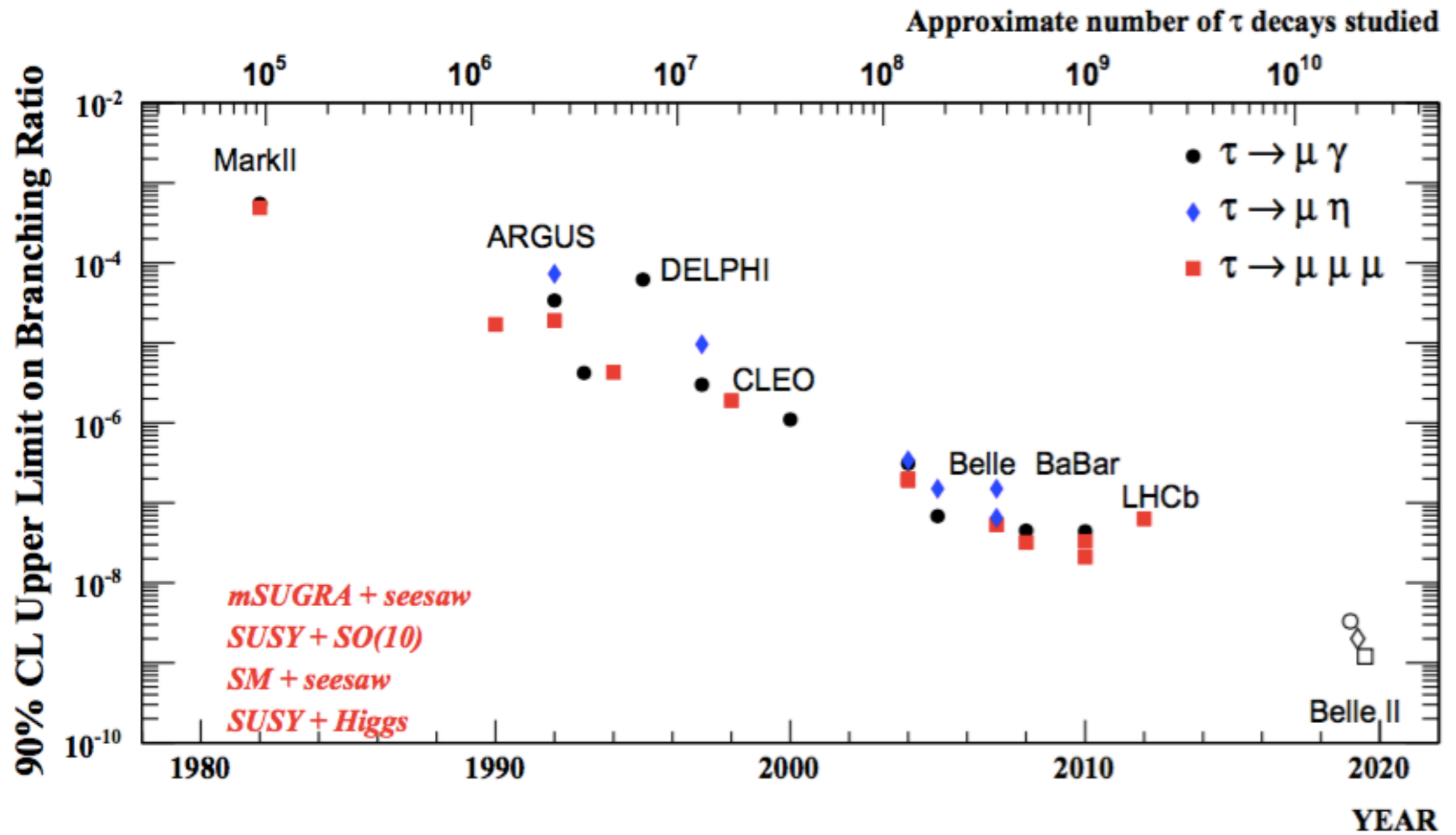
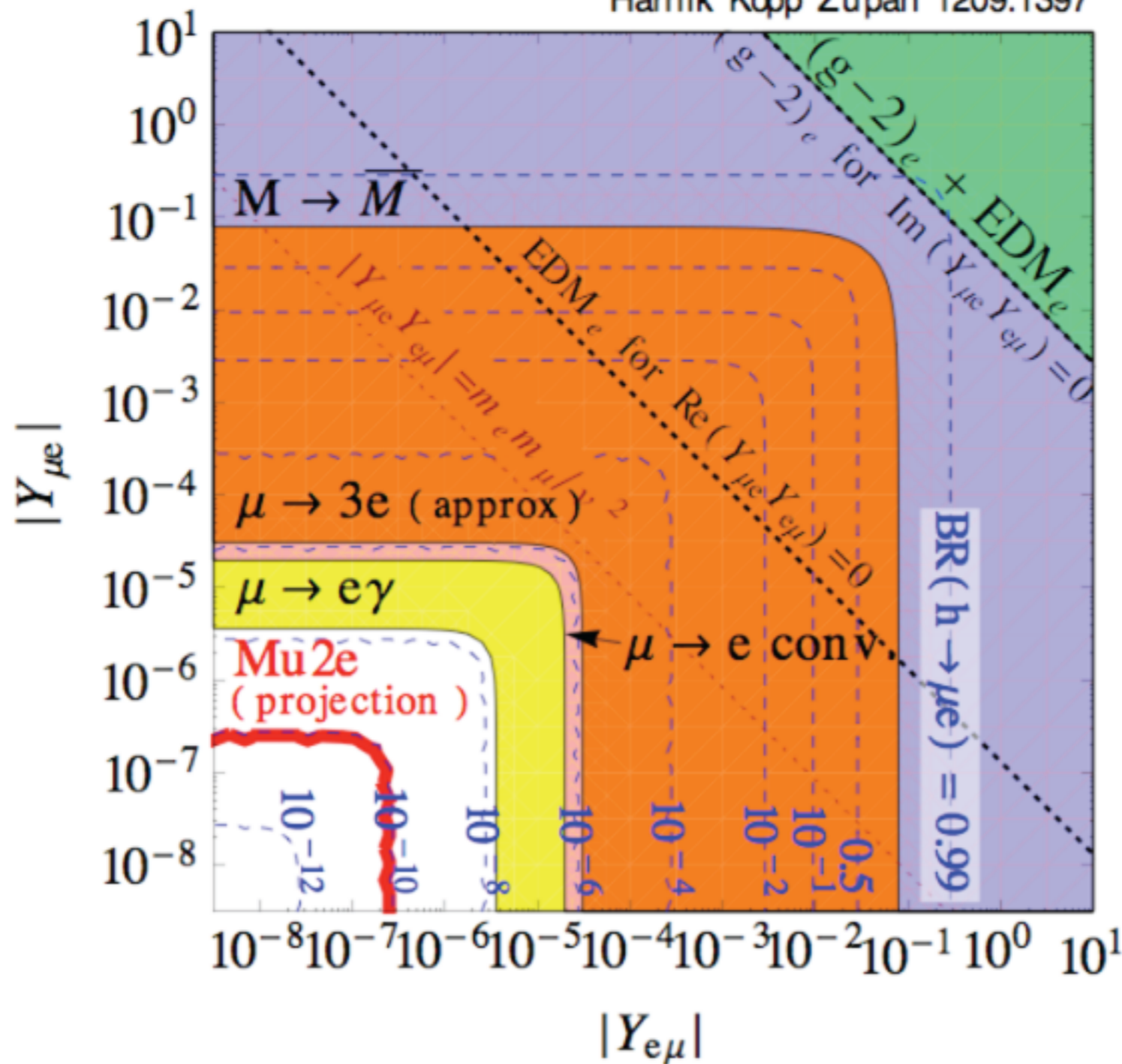


Figure 5: Evolution of 90% CL upper limit on LFV in τ decays vs. approximate number of τ decays studied.

Higgs couplings to μe

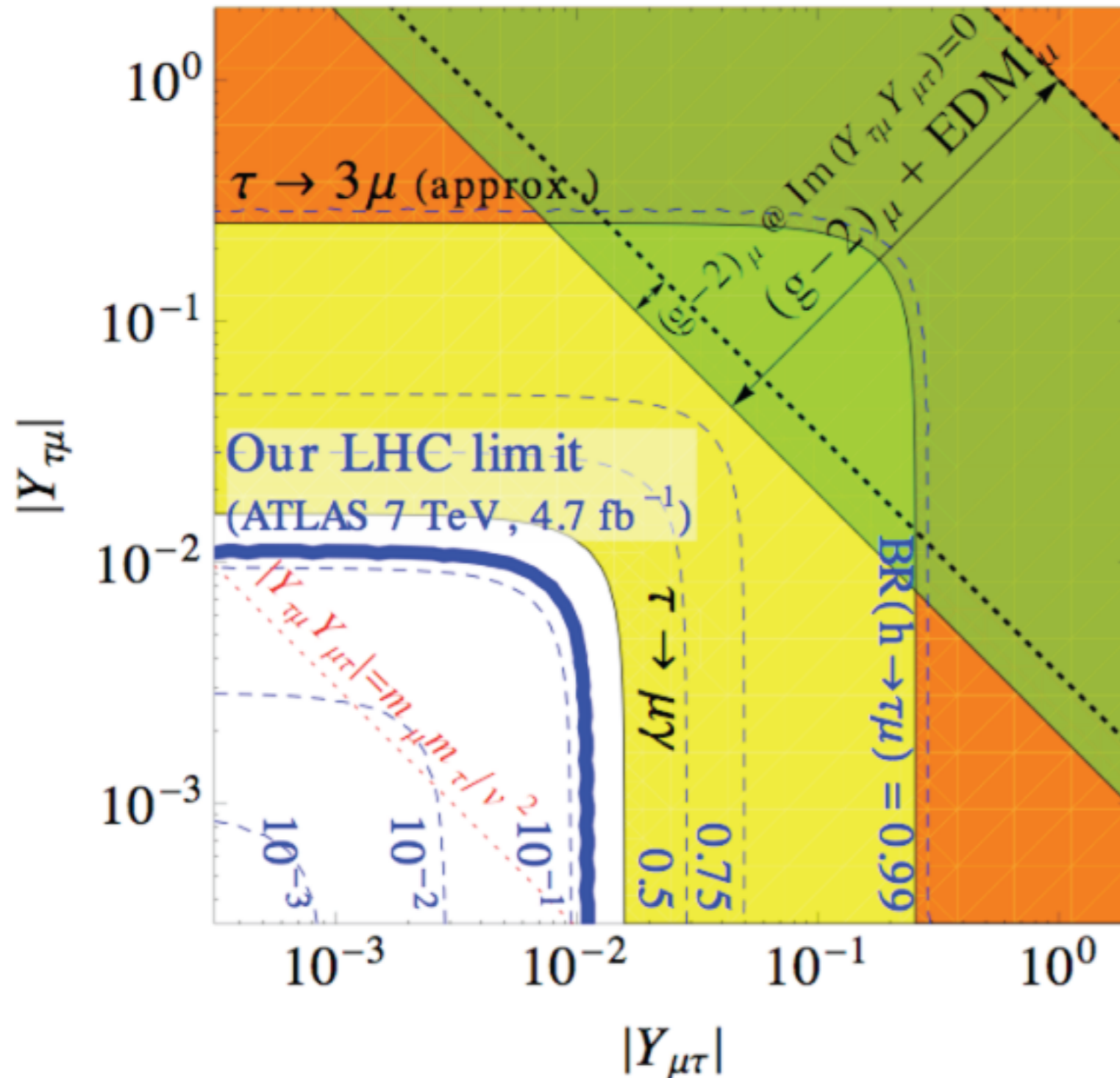
Harnik Kopp Zupan 1209.1397



Outside of LHC reach.

PROBING "natural" models.

Higgs couplings to $\tau\mu$



LHC $h \rightarrow \tau\mu$ gives dominant bound.
 (currently just a theorist's re-interpretation)

"natural models" are within reach.

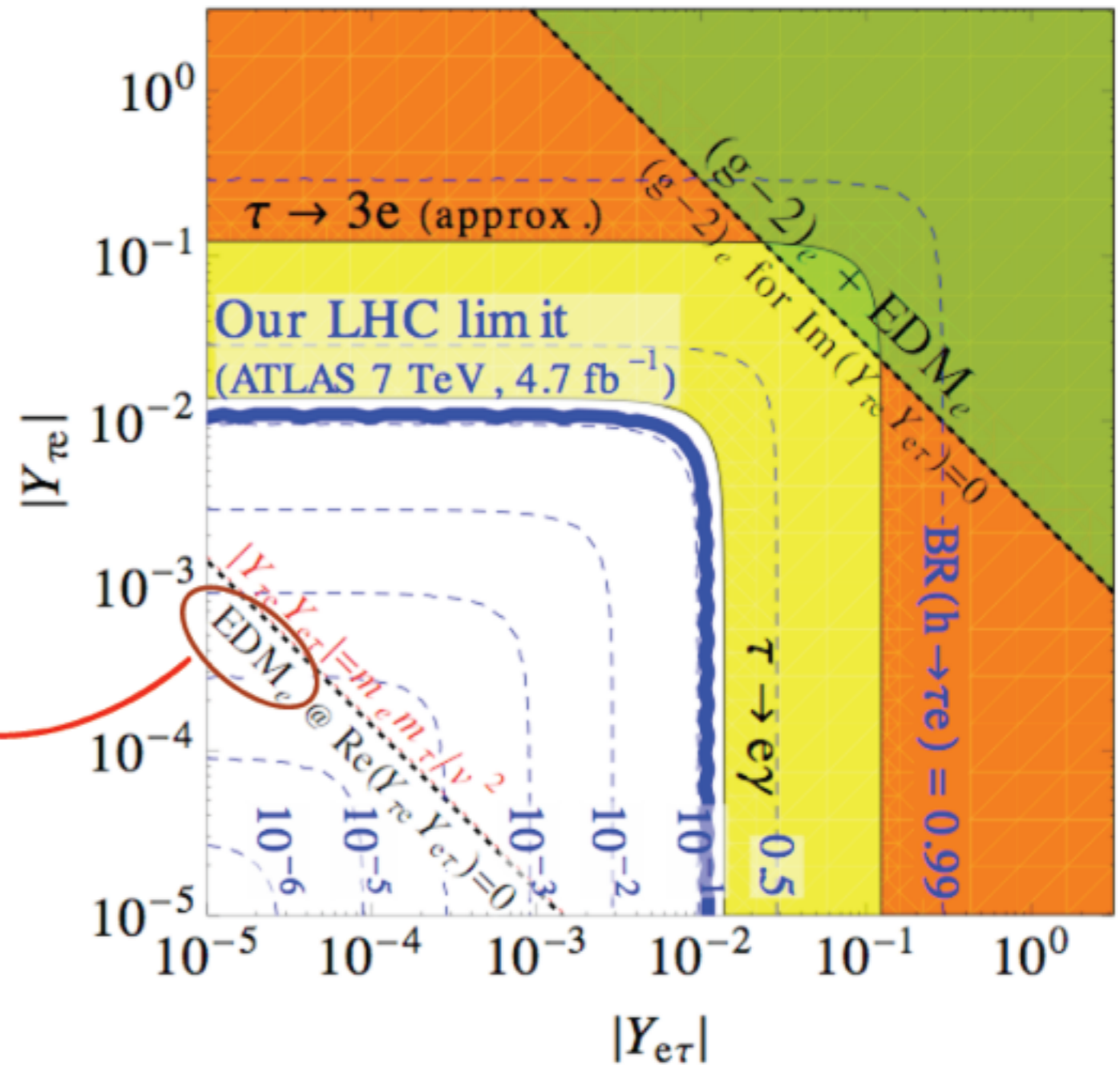
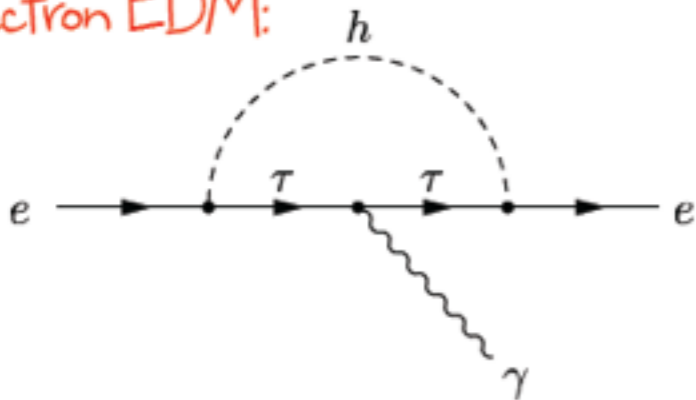
Higgs couplings to τe

* τe is similar to $\tau\mu$... but:

Harnik, Kopp Zupan 1209.1397

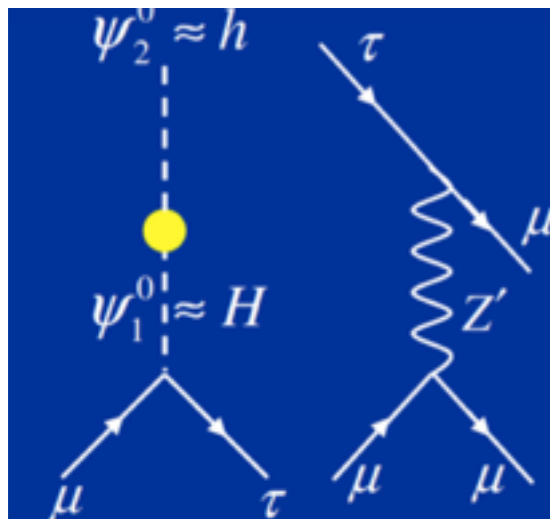
Electron EDM is interesting here!

electron EDM:



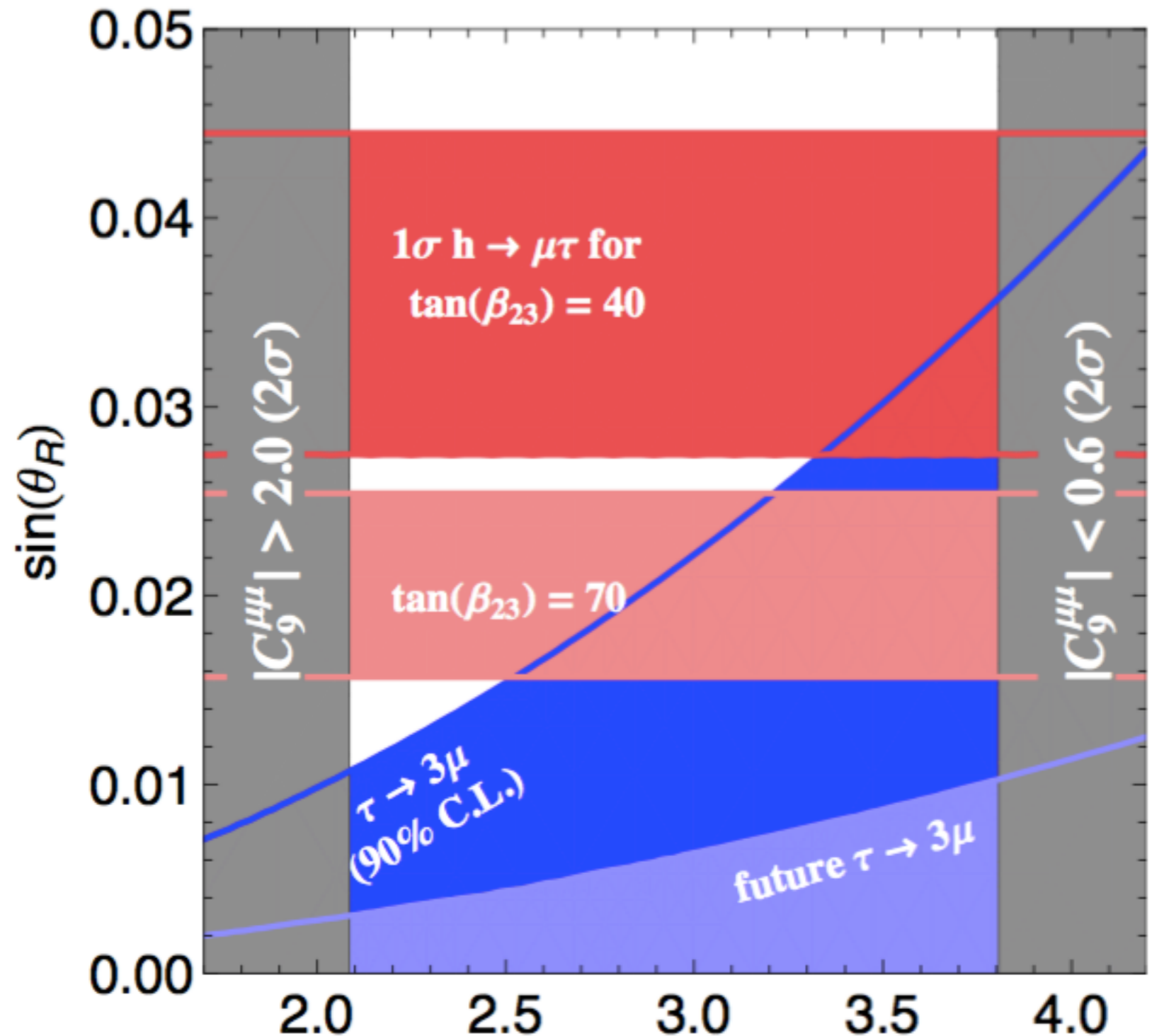
Interplay

Motivated by B-anomalies
Z'-model with 2 higgs in
Lmu-Ltau model



$$H \rightarrow \mu\tau \quad \tau \rightarrow 3\mu$$

$$\cos(\alpha_{23} - \beta_{23}) = 0.25, a = 1/3$$



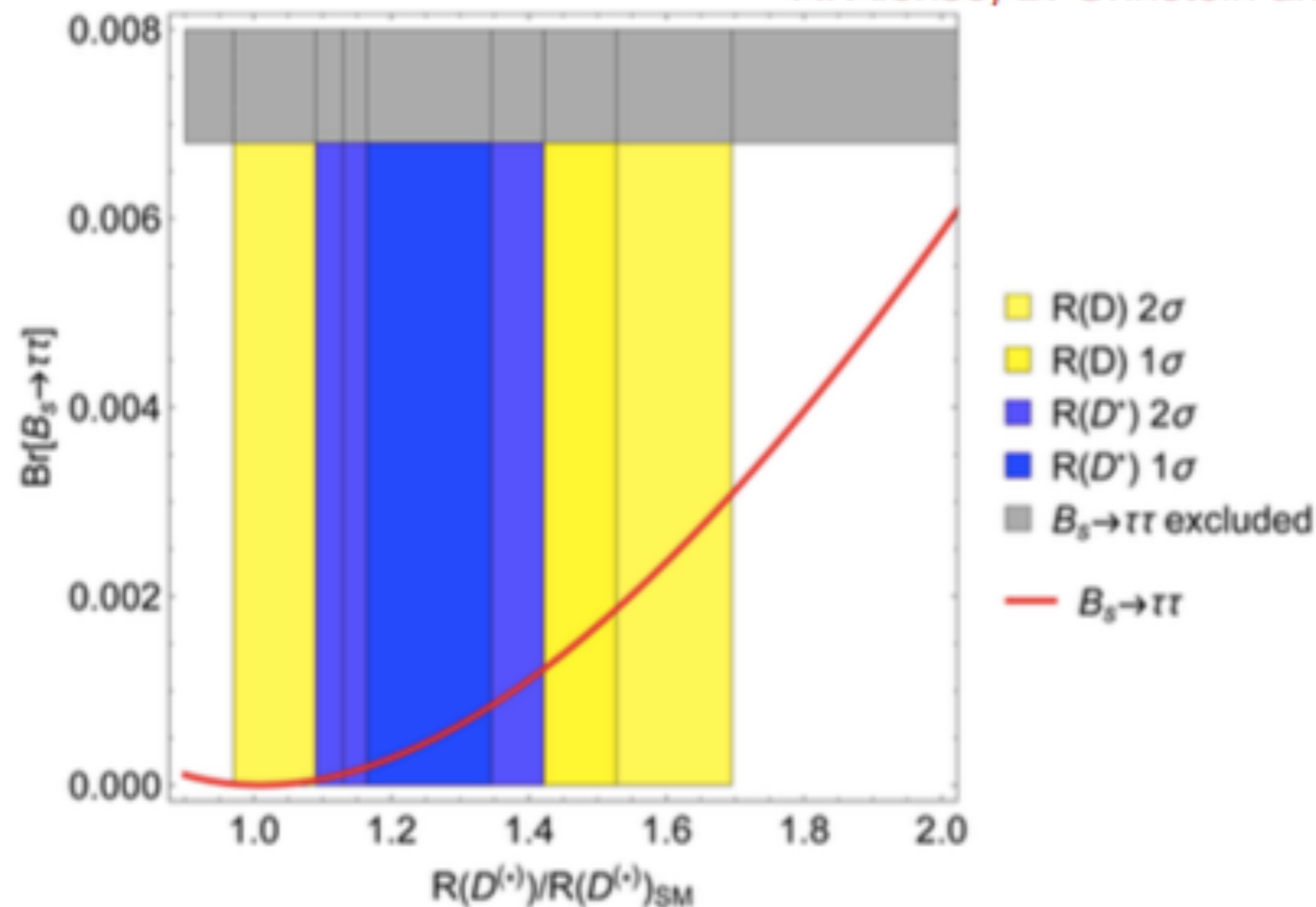
Conclusions

- My personal view: we are just at the beginning in exploring the potentiality: See $KS \rightarrow \mu\mu$, also the improvement for Ecal (I did not discuss much ee in the final state)
- New ideas in this sector came out recently

$R(D^{(*)})$ and $b \rightarrow s \tau \tau$ with Leptoquarks

- Large couplings to the 2nd generation needed in order to avoid collider bounds.
- Cancellation in $b \rightarrow s \nu \nu$ needed: $C^{(1)} = -C^{(3)}$

R. Alonso, B. Grinstein and J. Camalich, 1505.05164



$B_s \rightarrow \tau\tau$
very
strongly
enhanced

Kaon physics

Tests of CPV already among most stringent (ϵ_K, ϵ')

Near future improvements mostly due to theory (Lattice)

More progress foreseen in rare decays

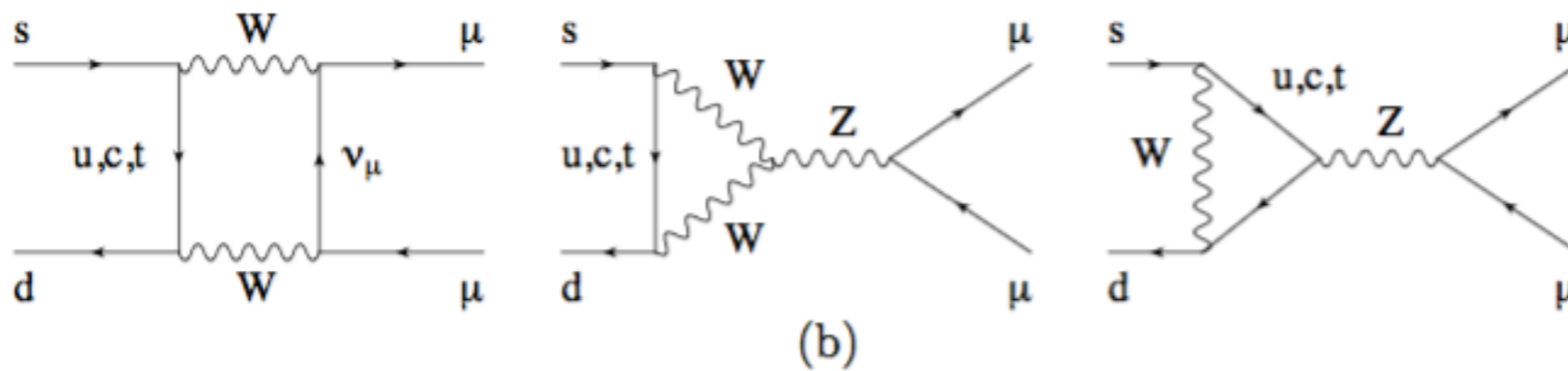
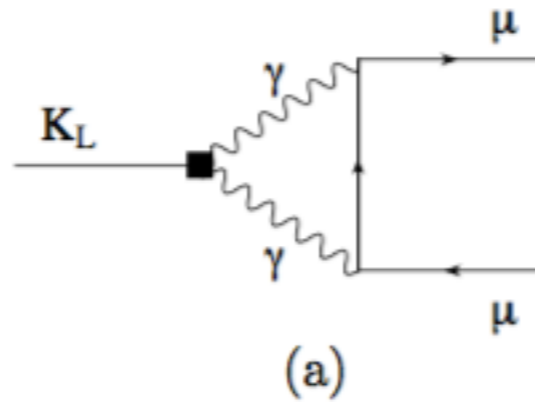
$$\Rightarrow K^+ \rightarrow \pi^+ \nu \bar{\nu}, K_L \rightarrow \pi^0 \nu \bar{\nu}$$

\Rightarrow rare K decays at HL-LHCb?

d'Ambrosio, PoS(FPCP2015)018

| | PDG | Prospects | |
|---|-------------------------------------|--------------------------------|-----------------|
| $K_S \rightarrow \mu\mu$ | $< 9 \times 10^{-9}$ at 90% CL (LD) | $(5.0 \pm 1.5) \cdot 10^{-12}$ | NP $< 10^{-11}$ |
| $K_L \rightarrow \mu\mu$ | $(6.84 \pm 0.11) \times 10^{-9}$ | difficult : SD \ll LD | |
| $K_S \rightarrow \mu\mu\mu\mu$ | — | SM LD $\sim 2 \times 10^{-14}$ | } NP? |
| $K_S \rightarrow ee\mu\mu$ | — | $\sim 10^{-11}$ | |
| $K_S \rightarrow eeee$ | — | $\sim 10^{-10}$ | |
| $K_S \rightarrow \pi^+ \pi^- \mu^+ \mu^-$ | — | SM LD $\sim 10^{-14}$ | |

SD depends from a crucial sign



$$\mathcal{B}(K_L \rightarrow \mu^+ \mu^-)_{\text{SM}} = \begin{cases} (6.85 \pm 0.80 \pm 0.06) \times 10^{-9} (+), \\ (8.11 \pm 1.49 \pm 0.13) \times 10^{-9} (-), \end{cases}$$

$$\begin{aligned} \mathcal{B}(K_S \rightarrow \mu^+ \mu^-)_{\text{SM}} &= (4.99 \text{ (LD)} + 0.19 \text{ (SD)}) \times 10^{-12} \\ &= (5.18 \pm 1.50 \pm 0.02) \times 10^{-12} \end{aligned}$$

We need FIGHT $DE/IB \sim 10^{-3}$

| | <i>IB</i> | <i>DE_{exp}</i> | |
|--------------------------------------|---|--------------------------------------|-------------------|
| $K_S \rightarrow \pi^+ \pi^- \gamma$ | 10^{-3} | $< 9 \cdot 10^{-5}$ | <i>E1</i> |
| $K^+ \rightarrow \pi^+ \pi^0 \gamma$ | 10^{-4} ($\Delta I = \frac{3}{2}$) | $(0.44 \pm 0.07) 10^{-5}$ PDG | <i>M1, E1</i> |
| $K_L \rightarrow \pi^+ \pi^- \gamma$ | 10^{-5} (CPV) | $(2.92 \pm 0.07) 10^{-5}$ KTeVnew | <i>M1,</i> VMD |

CPV is **only** from IB K_L (also measured in $K_L \rightarrow \pi^+ \pi^- e^+ e^-$)

BUT IB suppressed in K^+ and K_L .

$$K^+ \rightarrow \pi^+ \pi^0 \gamma$$

$$A(K \rightarrow \pi\pi\gamma) = F^{\mu\nu} [E \partial_\mu K \partial_\nu \pi + M \varepsilon_{\mu\nu\rho\sigma} \partial^\rho K \partial^\sigma \pi]$$

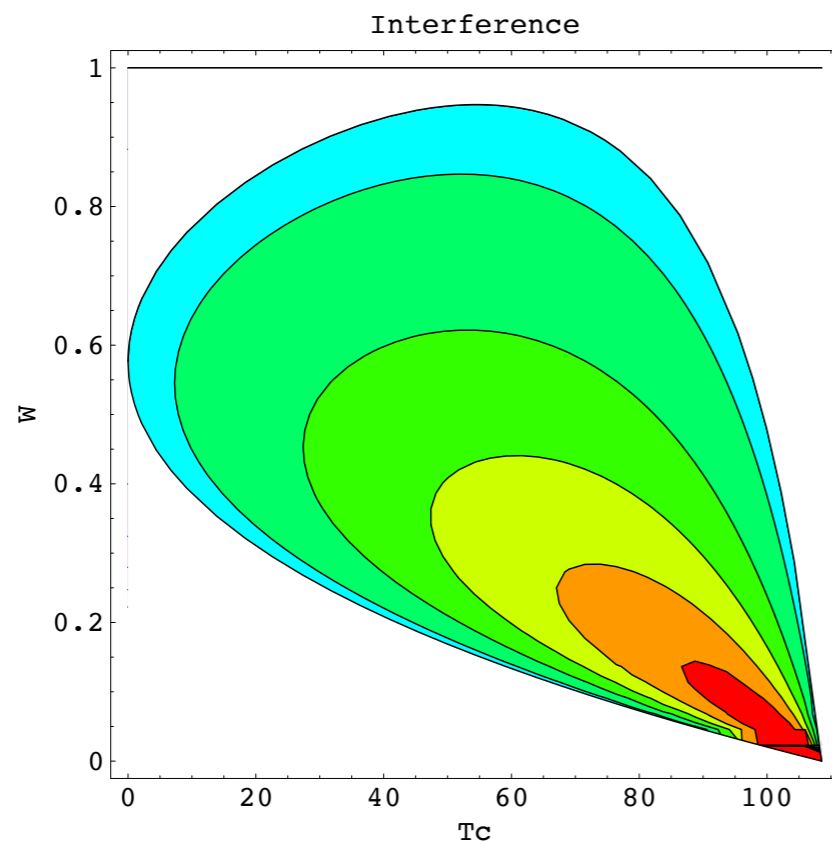
$E1$ and $M1$ are measured with Dalitz plot

$$\frac{\partial^2 \Gamma}{\partial T_c^* \partial W^2} = \frac{\partial^2 \Gamma_{IB}}{\partial T_c^* \partial W^2} \left[1 + \frac{m_{\pi^+}^2}{m_K^2} 2 \operatorname{Re} \left(\frac{E1}{eA} \right) W^2 + \frac{m_{\pi^+}^4}{m_K^2} \left(\left| \frac{E1}{eA} \right|^2 + \left| \frac{M1}{eA} \right|^2 \right) W^4 \right]$$

$$W^2 = (q \cdot p_K)(q \cdot p_+) / (m_\pi^2 m_K^2)$$

$$A = A(K^+ \rightarrow \pi^+ \pi^0)$$

NA48/2 CP violation



Dalitz plot analysis crucial

$$SM \leq \mathcal{O}(10^{-5})$$

Paver et al.

$$NP \leq \mathcal{O}(10^{-4})$$

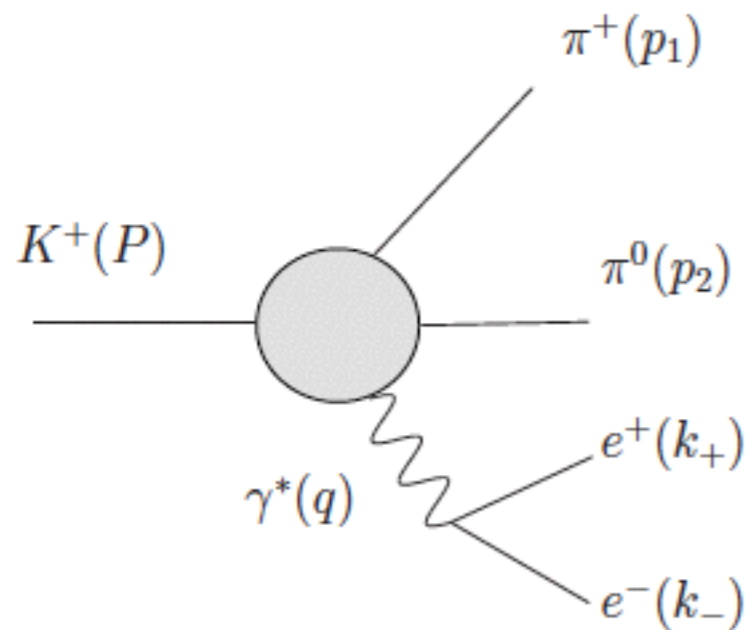
Colangelo et al.

$$NA48/2 < 1.5 \cdot 10^{-3} \text{ at } 90\% \text{ CL}$$

BUT NOT in the interesting interf. kin. region (statistics)

$$K_L \rightarrow \pi^+ \pi^- \gamma^* \rightarrow \pi^+ \pi^- e^+ e^-$$

Sehgal et al; Savage, Wise et al



- $\mathcal{M}_{LD} = \frac{e}{q^2} \bar{e} \gamma^\mu (1 - \gamma^5) e H_\mu$
- $H^\mu = F_1 p_1^\mu + F_2 p_2^\mu + F_3 \varepsilon^{\mu\nu\alpha\beta} p_{1\nu} p_{2\alpha} q_\beta$
- $F_{1,2} \sim E \quad F_3 \sim M$

- Interference $E \quad M$ novel compared to $K_L \rightarrow \pi^+ \pi^- \gamma$
- $E \quad M$ known from $K_L \rightarrow \pi^+ \pi^- \gamma$ (IB and DE)

$$K^+ \rightarrow \pi^+ \pi^0 \gamma^* \rightarrow \pi^+ \pi^0 e^+ e^-$$

Cappiello, Cata, G.D. and Gao,

- the asymm. , $\frac{\Re(E_B M^*)}{|E_B|^2 + |M|^2}$, not as lucky $E_B \gg M$:
- $B(K^+)_{IB} \sim 3.3 \times 10^{-6} \sim 50 B(K^+)_{M}$
- Short distance info without having simultaneously K^+ and K^- , asymm. in phase space, (P-violation) interesting! No ϵ -contamination
- interesting Dalitz plots (at fixed q^2) to disentangle M from E_B
- at $q^2 = 50\text{MeV}$ IB only 10 times larger than DE