

# EFT in VBF and VBS

Michael Rauch | HL-/HE-LHC Workshop, 31 Oct 2017

INSTITUTE FOR THEORETICAL PHYSICS



- Vector-Boson Fusion (VBF) and Vector-Boson Scattering (VBS)  
important processes for HL-/HE-LHC

[→ talks by Jan Kieseler, Claire Lee]

- searches for new physics important task
- useful tool for heavy new physics: **effective field theory (EFT)**

assumption: new physics is heavy

can integrate out heavy, non-SM degrees of freedom

higher-dimensional operators appearing in Lagrangian

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_{d>4} \sum_i \frac{f_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}$$

- operators  $\mathcal{O}$  contain SM fields only
- respect SM gauge symmetries
- suppressed by  $1/\Lambda^{d-4}$  ( $\Lambda$ : scale of new physics)  
→ keep only leading order(s) (lowest dimension  $d = 6$ )
- building blocks:
  - Higgs field  $\Phi$
  - (covariant) derivative  $\partial^\mu, D^\mu$
  - field strength tensors  $G^{\mu\nu}, W^{\mu\nu}, B^{\mu\nu}$
  - fermion fields  $\psi$

- linear realization of the EFT
- D6: 59 operators when assuming
  - baryon/lepton-number conservation
  - flavour universality

[Buchmüller, Wyler; Hagiwara et al; Grzadkowski et al; ...]

## List of Operators (only gauge and Higgs couplings)

$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu} \Phi$$

$$\mathcal{O}_{WWW} = \text{Tr} \left[ \widehat{W}^\mu{}_\nu \widehat{W}^\nu{}_\rho \widehat{W}^\rho{}_\mu \right]$$

$$\mathcal{O}_{\widetilde{W}} = (D_\mu \Phi)^\dagger \widetilde{W}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_{\widetilde{W}W} = \Phi^\dagger \widetilde{W}_{\mu\nu} \widehat{W}^{\mu\nu} \Phi$$

$$\mathcal{O}_{\widetilde{W}WW} = \text{Tr} \left[ \widetilde{W}^\mu{}_\nu \widehat{W}^\nu{}_\rho \widehat{W}^\rho{}_\mu \right]$$

$$\mathcal{O}_B = (D_\mu \Phi)^\dagger \widehat{B}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_{BB} = \Phi^\dagger \widehat{B}_{\mu\nu} \widehat{B}^{\mu\nu} \Phi$$

$$\mathcal{O}_{\phi,2} = \partial_\mu (\Phi^\dagger \Phi) \partial^\mu (\Phi^\dagger \Phi)$$

$$\mathcal{O}_{\widetilde{B}} = (D_\mu \Phi)^\dagger \widetilde{B}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_{\widetilde{B}B} = \Phi^\dagger \widetilde{B}_{\mu\nu} \widehat{B}^{\mu\nu} \Phi$$

One constraint on CP-odd operators

$$\mathcal{O}_{\widetilde{W}} + \frac{1}{2} \mathcal{O}_{\widetilde{W}W} = \mathcal{O}_{\widetilde{B}} + \frac{1}{2} \mathcal{O}_{\widetilde{B}B}$$

Additional CP-even operator

$$\mathcal{O}_{\phi W} = \text{Tr} [W^{\mu\nu} W_{\mu\nu}] \Phi^\dagger \Phi \equiv 2\mathcal{O}_{WW}$$

## List of Operators (only gauge and Higgs couplings)

$$\begin{aligned}
 \mathcal{O}_W &= (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi) & \mathcal{O}_B &= (D_\mu \Phi)^\dagger \widehat{B}^{\mu\nu} (D_\nu \Phi) \\
 \mathcal{O}_{WW} &= \Phi^\dagger \widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu} \Phi & \mathcal{O}_{BB} &= \Phi^\dagger \widehat{B}_{\mu\nu} \widehat{B}^{\mu\nu} \Phi \\
 \mathcal{O}_{WWW} &= \text{Tr} \left[ \widehat{W}^\mu{}_\nu \widehat{W}^\nu{}_\rho \widehat{W}^\rho{}_\mu \right] & \mathcal{O}_{\phi,2} &= \partial_\mu (\Phi^\dagger \Phi) \partial^\mu (\Phi^\dagger \Phi) \\
 \mathcal{O}_{\widetilde{W}} &= (D_\mu \Phi)^\dagger \widetilde{W}^{\mu\nu} (D_\nu \Phi) & \mathcal{O}_{\widetilde{B}} &= (D_\mu \Phi)^\dagger \widetilde{B}^{\mu\nu} (D_\nu \Phi) \\
 \mathcal{O}_{\widetilde{W}W} &= \Phi^\dagger \widetilde{W}_{\mu\nu} \widehat{W}^{\mu\nu} \Phi & \mathcal{O}_{\widetilde{B}B} &= \Phi^\dagger \widetilde{B}_{\mu\nu} \widehat{B}^{\mu\nu} \Phi \\
 \mathcal{O}_{\widetilde{W}WW} &= \text{Tr} \left[ \widetilde{W}^\mu{}_\nu \widehat{W}^\nu{}_\rho \widehat{W}^\rho{}_\mu \right]
 \end{aligned}$$

Modification of corresponding triple-gauge-coupling vertices:

	$\mathcal{O}_{WWW}$	$\mathcal{O}_W$	$\mathcal{O}_B$	$\mathcal{O}_{WW}$	$\mathcal{O}_{BB}$	$\mathcal{O}_{\phi,2}$	$\mathcal{O}_{\widetilde{W}WW}$	$\mathcal{O}_{\widetilde{W}}$	$\mathcal{O}_{\widetilde{B}}$	$\mathcal{O}_{\widetilde{W}W}$	$\mathcal{O}_{\widetilde{B}B}$
<i>WWZ</i>	X	X	X				X	X	X		
<i>WW<math>\gamma</math></i>	X	X	X				X	X	X		
<i>HWW</i>		X		X		X		X		X	
<i>HZZ</i>		X	X	X	X	X		X	X	X	X
<i>HZ<math>\gamma</math></i>		X	X	X	X	(X)		X	X	X	X
<i>H<math>\gamma\gamma</math></i>				X	X	(X)				X	X
<i>WWW</i>	X	X					X				
<i>WWZ</i>	X	X					X				
<i>WWZ<math>\gamma</math></i>	X	X					X				
<i>WW<math>\gamma\gamma</math></i>	X						X				

## Dimension-8

Bosonic dimension-8 operators

[Eboli, Gonzalez-Garcia]

(D6 could be loop-induced  $\rightarrow$  D8 effects can become sizable [Arzt, Einhorn, Wudka])

$$\mathcal{O}_{S,0} = \left[ (D_\mu \Phi)^\dagger D_\nu \Phi \right] \times \left[ (D^\mu \Phi)^\dagger D^\nu \Phi \right]$$

$$\mathcal{O}_{S,1} = \left[ (D_\mu \Phi)^\dagger D^\mu \Phi \right] \times \left[ (D_\nu \Phi)^\dagger D^\nu \Phi \right]$$

$$\mathcal{O}_{S,2} = \left[ (D_\mu \Phi)^\dagger D_\nu \Phi \right] \times \left[ (D^\nu \Phi)^\dagger D^\mu \Phi \right]$$

$$\mathcal{O}_{M,0} = \text{Tr} \left[ \widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu} \right] \times \left[ (D_\beta \Phi)^\dagger D^\beta \Phi \right]$$

$$\mathcal{O}_{M,1} = \text{Tr} \left[ \widehat{W}_{\mu\nu} \widehat{W}^{\nu\beta} \right] \times \left[ (D_\beta \Phi)^\dagger D^\mu \Phi \right]$$

$$\mathcal{O}_{M,2} = \left[ \widehat{B}_{\mu\nu} \widehat{B}^{\mu\nu} \right] \times \left[ (D_\beta \Phi)^\dagger D^\beta \Phi \right]$$

$$\mathcal{O}_{M,3} = \left[ \widehat{B}_{\mu\nu} \widehat{B}^{\nu\beta} \right] \times \left[ (D_\beta \Phi)^\dagger D^\mu \Phi \right]$$

$$\mathcal{O}_{M,4} = \left[ (D_\mu \Phi)^\dagger \widehat{W}_{\beta\nu} D^\mu \Phi \right] \times \widehat{B}^{\beta\nu}$$

$$\mathcal{O}_{M,5} = \left[ (D_\mu \Phi)^\dagger \widehat{W}_{\beta\nu} D^\nu \Phi \right] \times \widehat{B}^{\beta\mu}$$

$$\mathcal{O}_{M,7} = \left[ (D_\mu \Phi)^\dagger \widehat{W}_{\beta\nu} \widehat{W}^{\beta\mu} D^\nu \Phi \right]$$

$$\mathcal{O}_{T,0} = \text{Tr} \left[ \widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu} \right] \times \text{Tr} \left[ \widehat{W}_{\alpha\beta} \widehat{W}^{\alpha\beta} \right]$$

$$\mathcal{O}_{T,1} = \text{Tr} \left[ \widehat{W}_{\alpha\nu} \widehat{W}^{\mu\beta} \right] \times \text{Tr} \left[ \widehat{W}_{\mu\beta} \widehat{W}^{\alpha\nu} \right]$$

$$\mathcal{O}_{T,2} = \text{Tr} \left[ \widehat{W}_{\alpha\mu} \widehat{W}^{\mu\beta} \right] \times \text{Tr} \left[ \widehat{W}_{\beta\nu} \widehat{W}^{\nu\alpha} \right]$$

$$\mathcal{O}_{T,5} = \text{Tr} \left[ \widehat{W}_{\mu\nu} \widehat{W}^{\mu\nu} \right] \times \widehat{B}_{\alpha\beta} \widehat{B}^{\alpha\beta}$$

$$\mathcal{O}_{T,6} = \text{Tr} \left[ \widehat{W}_{\alpha\nu} \widehat{W}^{\mu\beta} \right] \times \widehat{B}_{\mu\beta} \widehat{B}^{\alpha\nu}$$

$$\mathcal{O}_{T,7} = \text{Tr} \left[ \widehat{W}_{\alpha\mu} \widehat{W}^{\mu\beta} \right] \times \widehat{B}_{\beta\nu} \widehat{B}^{\nu\alpha}$$

$$\mathcal{O}_{T,8} = \widehat{B}_{\mu\nu} \widehat{B}^{\mu\nu} \widehat{B}_{\alpha\beta} \widehat{B}^{\alpha\beta}$$

$$\mathcal{O}_{T,9} = \widehat{B}_{\alpha\mu} \widehat{B}^{\mu\beta} \widehat{B}_{\beta\nu} \widehat{B}^{\nu\alpha}$$

$\rightarrow$  each operators contains  
at least four bosons

$\Rightarrow$  leading contribution to  
quartic gauge coupling

# Non-linear EFT

Also possible to use non-linear EFT (electroweak chiral Lagrangian)  
based on power counting in terms of canonical dimension

[Appelquist; Longhitano; ...; Alboteanu, Kilian, Reuter; ...]

$$\mathcal{L}_4 = \alpha_4 (\text{Tr} [V_\mu V_\nu])^2 ,$$

$$\mathcal{L}_5 = \alpha_5 (\text{Tr} [V_\mu V^\mu])^2 ,$$

$$\mathcal{L}_{S,0} = F_{S,0} \text{Tr} [(D_\mu \hat{H})^\dagger D_\nu \hat{H}] \times \text{Tr} [(D^\mu \hat{H})^\dagger D^\nu \hat{H}] ,$$

$$\mathcal{L}_{S,1} = F_{S,1} \text{Tr} [(D_\mu \hat{H})^\dagger D^\mu \hat{H}] \times \text{Tr} [(D_\nu \hat{H})^\dagger D^\nu \hat{H}] .$$

with

$$V_\mu = \Sigma (D_\mu \Sigma)^\dagger = - (D_\mu \Sigma) \Sigma^\dagger ,$$

$$D_\mu \Sigma = \partial_\mu \Sigma + ig \frac{\sigma^a}{2} W_\mu^a \Sigma - ig' \Sigma B_\mu \frac{\sigma^3}{2} ,$$

$$\Sigma = \exp \left( -\frac{i}{v} \sigma^a w^a \right) \underset{\text{unitary gauge}}{=} 1$$

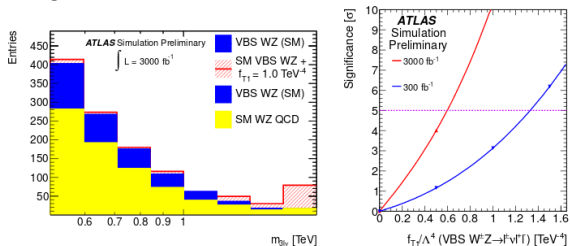
$$\hat{H} = \frac{1}{2} \begin{pmatrix} v + H - iw^3 & -i(w^1 - iw^2) \\ -i(w^1 + iw^2) & v + H + iw^3 \end{pmatrix} \underset{\text{unitary gauge}}{=} \frac{v + H}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

# Experimental Projections

Some projection studies available for HL-LHC

[↔ talk by Claire Lee]

ATLAS:



Parameter	dimension	channel	$\Lambda_{UV}$ [TeV]	300 fb <sup>-1</sup>		3000 fb <sup>-1</sup>	
				5 $\sigma$	95% CL	5 $\sigma$	95% CL
$c_{\phi W}/\Lambda^2$	6	ZZ	1.9	34 TeV <sup>-2</sup>	20 TeV <sup>-2</sup>	16 TeV <sup>-2</sup>	9.3 TeV <sup>-2</sup>
$f_{S0}/\Lambda^4$	8	W <sup>±</sup> W <sup>±</sup>	2.0	10 TeV <sup>-4</sup>	6.8 TeV <sup>-4</sup>	4.5 TeV <sup>-4</sup>	0.8 TeV <sup>-4</sup>
$f_{T1}/\Lambda^4$	8	WZ	3.7	1.3 TeV <sup>-4</sup>	0.7 TeV <sup>-4</sup>	0.6 TeV <sup>-4</sup>	0.3 TeV <sup>-4</sup>
$f_{T8}/\Lambda^4$	8	Z $\gamma\gamma$	12	0.9 TeV <sup>-4</sup>	0.5 TeV <sup>-4</sup>	0.4 TeV <sup>-4</sup>	0.2 TeV <sup>-4</sup>
$f_{T9}/\Lambda^4$	8	Z $\gamma\gamma$	13	2.0 TeV <sup>-4</sup>	0.9 TeV <sup>-4</sup>	0.7 TeV <sup>-4</sup>	0.3 TeV <sup>-4</sup>

CMS (WZjj):

Significance	3 $\sigma$	5 $\sigma$
SM EWK scattering discovery	75 fb <sup>-1</sup>	185 fb <sup>-1</sup>
$f_{T1}/\Lambda^4$ at 300 fb <sup>-1</sup>	0.8 TeV <sup>-4</sup>	1.0 TeV <sup>-4</sup>
$f_{T1}/\Lambda^4$ at 3000 fb <sup>-1</sup>	0.45 TeV <sup>-4</sup>	0.55 TeV <sup>-4</sup>

→ test energy scales up to

- $f = 4\pi$ :  $\Lambda = 2.5$  TeV
- $f = 1$ :  $\Lambda = 1.4$  TeV
- $f = \frac{1}{16\pi^2}$ :  $\Lambda = 0.4$  TeV

Important gauge **cancellations** between different diagram types

- longitudinal  $W$  scattering through quartic gauge boson vertex



**high energy limit:** centre-of-mass energy  $\sqrt{s} \rightarrow \infty$

$$\mathcal{M}_{\text{quartic vertex}} \propto s^2 \rightarrow \text{cross section diverges} \quad \sigma \propto s^4/s = s^3 \rightarrow \infty$$

Gauge boson with momentum  $\vec{p} = (0, 0, p)^T$ :

longitudinal polarization vector:

$$\epsilon_L^\mu = \frac{1}{M} (p, 0, 0, E)^T \stackrel{E \gg M}{\approx} \frac{p^\mu}{M}$$



Important gauge **cancellations** between different diagram types

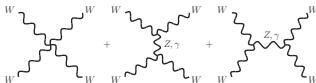
- longitudinal  $W$  scattering through quartic gauge boson vertex



**high energy limit:** centre-of-mass energy  $\sqrt{s} \rightarrow \infty$

$\mathcal{M}_{\text{quartic vertex}} \propto s^2 \rightarrow$  **cross section diverges**  $\sigma \propto s^4/s = s^3 \rightarrow \infty$

- add triple gauge boson vertices



$\mathcal{M}_{\text{quartic+triple vertices}} \propto s \rightarrow$  **still divergent**

Important gauge **cancellations** between different diagram types

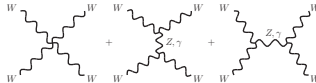
- longitudinal  $W$  scattering through quartic gauge boson vertex



**high energy limit:** centre-of-mass energy  $\sqrt{s} \rightarrow \infty$

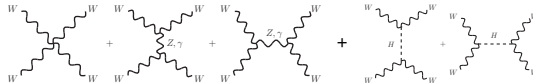
$\mathcal{M}_{\text{quartic vertex}} \propto s^2 \rightarrow$  **cross section diverges**  $\sigma \propto s^4/s = s^3 \rightarrow \infty$

- add triple gauge boson vertices



$\mathcal{M}_{\text{quartic+triple vertices}} \propto s \rightarrow$  **still divergent**

- additional Higgs diagrams



remove divergence exactly  $\mathcal{M} \propto \text{const.}$   $\sigma \propto 1/s \rightarrow 0$

Anomalous gauge couplings spoil cancellation  $\rightarrow$  stringent **tests**

Anomalous gauge couplings spoil cancellation

↔ effects can become large → **unitarity violation** → unphysical

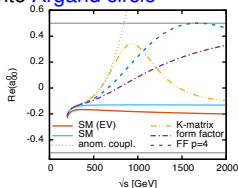
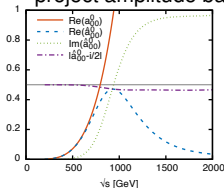
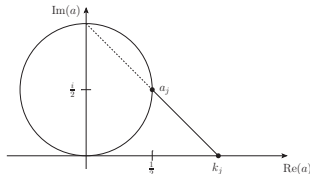
Several solutions:

- consider only unitarity-conserving phase-space regions  
loses some information → possibly reduced sensitivity  
cut on relevant region might not be directly accessible ( $m_{4\ell}$  vs. neutrinos)
- (dipole) **form factor** multiplying amplitudes

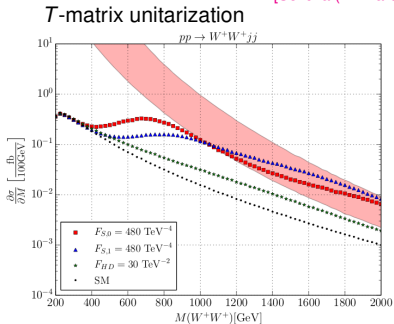
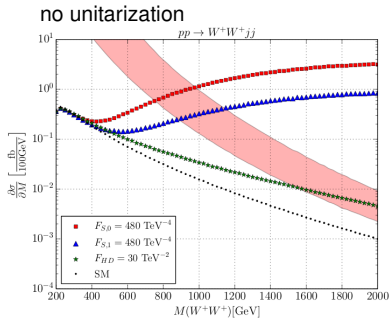
[...]

$$\mathcal{F}(s) = \frac{1}{\left(1 + \frac{s}{\Lambda_{\text{FF}}^2}\right)^n} \quad \Lambda_{\text{FF}}^2, n: \text{ free parameters}$$

- **K-/T-matrix unitarization** [Alboteanu, Kilian, Reuter, Sekulla]  
based on partial-wave analysis [Jacob, Wick]  
project amplitude back onto **Argand circle**



[Sekula (Whizard)]



red band: one/all partial-wave amplitudes saturated

# Impact of Current Limits

Investigate impact of D6 vs D8 operators on VBS

D6 input: Global Higgs and Gauge analysis of run-I data

[Butter, Eboli, Gonzalez-Fraile, Gonzalez-Garcia, Plehn, MR]

Take results and apply to vector-boson scattering

⇒ No contribution from  $\mathcal{O}_{GG}$  and fermionic operators

$f_x / \Lambda^2 [\text{TeV}^{-2}]$	LHC-Higgs + LHC-TGV + LEP-TGV	
	Best fit	95% CL interval
$f_{WW}$	-0.1	(-3.1, 3.7)
$f_{BB}$	0.9	(-3.3, 6.1)
$f_W$	1.7	(-0.98, 5.0)
$f_B$	1.7	(-11.8, 8.8)
$f_{WWW}$	-0.06	(-2.6, 2.6)
$f_{\phi,2}$	1.3	(-7.2, 7.5)

For simplicity: use pos. and neg. 95% CL bound with other parameters set to zero  
→ slightly larger effect than true 95% CL bound

Additionally:

effect from dimension-8 operator  $\mathcal{O}_{S,1}$

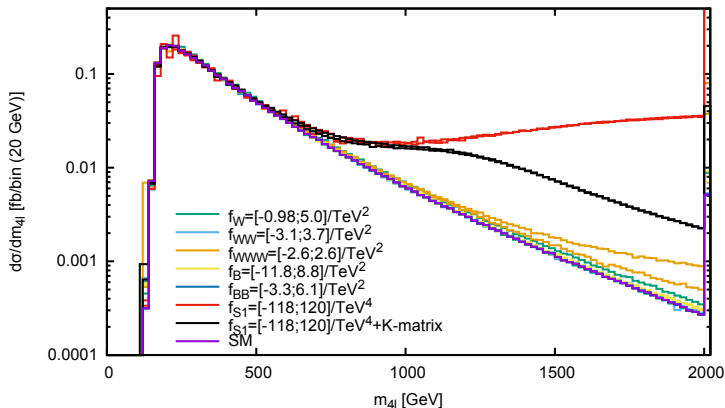
using CMS,  $W^\pm W^\pm jj$ ,  $\sqrt{S} = 8$  TeV, no unitarization

[arXiv:1410.6315]

$$f_{S,1}/\Lambda^4 \in (-118, 120)\text{TeV}^{-4} \quad (\text{for } f_{S,0}/\Lambda^4 = 0)$$

# Results

Process:  $pp \rightarrow W^+W^+jj \rightarrow \ell^+\nu\ell^+\nu jj$ ,  $\sqrt{S} = 13$  TeV, VBF cuts, NLO QCD

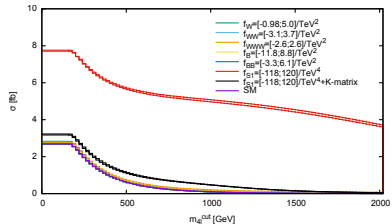
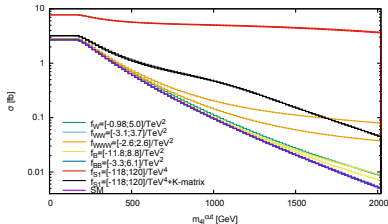


- last bin: overflow bin,  $m_{4\ell} > 2000$  GeV
- effect of D6 contributions in general small; largest one by  $\mathcal{O}_{WWW}$
- D8 operator clearly dominating

# Results

Process:  $pp \rightarrow W^+ W^+ jj \rightarrow \ell^+ \nu \ell^+ \nu jj$ ,  $\sqrt{S} = 13$  TeV, VBF cuts, NLO QCD

cross section when requiring  $m_{4\ell} > m_{4\ell}^{\text{cut}}$



- $\mathcal{O}_{WWW}$  contribution large only for very high  $m_{4\ell} \leftrightarrow$  low event counts

excess of 10 events for  $m_{4\ell} > 1$  TeV,  $\mathcal{L} = 100 \text{ fb}^{-1}$ , SM contrib. of 10 events  
other D6 operators below 1 event

$\leftrightarrow$  unitarity violating contributions (?)
- $\mathcal{O}_{S1}$  yielding large excess even without cuts on  $m_{4\ell}$

excess of almost 500 events for  $m_{4\ell} > 1$  TeV,  $\mathcal{L} = 100 \text{ fb}^{-1}$   
even after unitarization excess of 37 events

- EFT formulation useful tool to parametrize new-physics
- contribution to anomalous triple and quartic gauge couplings
- strong growth with energy
  - unitarity violation in experimentally probed region
  - ⇒ need some mitigation procedure (cut, form factor, T-matrix, . . .)
- effect of dimension-6 operators in VBF/VBS processes in general small  
13/14 TeV diboson data will further reduce the allowed contributions  
final accuracy vs effects in VBF/VBS
- effect of dimension-8 operators dominates
  - constraining power of experimental results
  - expectations for HE-LHC



# Backup

Hiccups of the original version:

vanish identically:

$$\mathcal{O}_{T,3} = \text{Tr} \left[ \widehat{W}_{\alpha\mu} \widehat{W}^{\mu\beta} \widehat{W}^{\nu\alpha} \right] \times \widehat{B}_{\beta\nu}$$

$$\mathcal{O}_{T,4} = \text{Tr} \left[ \widehat{W}_{\alpha\mu} \widehat{W}^{\alpha\mu} \widehat{W}^{\beta\nu} \right] \times \widehat{B}_{\beta\nu}$$

redundant:

$$\begin{aligned} \mathcal{O}_{M,6} &= \left[ (D_\mu \Phi)^\dagger \widehat{W}_{\beta\nu} \widehat{W}^{\beta\nu} D^\mu \Phi \right] \\ &= \frac{1}{2} \mathcal{O}_{M,0} \end{aligned}$$

missing:

$$\mathcal{O}_{S,2} = \left[ (D_\mu \Phi)^\dagger D_\nu \Phi \right] \times \left[ (D^\nu \Phi)^\dagger D^\mu \Phi \right]$$

Contribution to the different vertices:

	$\mathcal{O}_{S,0}$	$\mathcal{O}_{M,0}$	$\mathcal{O}_{M,2}$	$\mathcal{O}_{T,0}$	$\mathcal{O}_{T,5}$	
	$\mathcal{O}_{S,1}$	$\mathcal{O}_{M,1}$	$\mathcal{O}_{M,3}$	$\mathcal{O}_{T,1}$	$\mathcal{O}_{T,6}$	$\mathcal{O}_{T,8}$
	$\mathcal{O}_{S,2}$	$\mathcal{O}_{M,7}$	$\mathcal{O}_{M,4}$	$\mathcal{O}_{T,2}$	$\mathcal{O}_{T,7}$	$\mathcal{O}_{T,9}$
			$\mathcal{O}_{M,5}$			
WWWW	X	X		X		
WWZZ	X	X	X	X	X	
ZZZZ	X	X	X	X	X	X
WWZ $\gamma$		X	X	X	X	
WW $\gamma\gamma$		X	X	X	X	
ZZZ $\gamma$		X	X	X	X	X
ZZ $\gamma\gamma$		X	X	X	X	X
Z $\gamma\gamma\gamma$				X	X	X
$\gamma\gamma\gamma\gamma$				X	X	X

Relations between linear and non-linear EFT

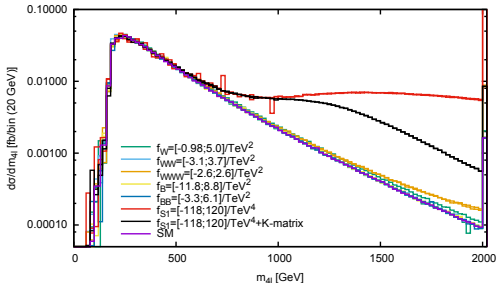
$$\alpha_4 = \frac{v^4}{16} F_{S,0} = \frac{v^4}{16} \frac{f_{S,0} + f_{S,2}}{\Lambda^4}, \quad f_{S,0} = f_{S,2}$$
$$\alpha_5 = \frac{v^4}{16} F_{S,1} = \frac{v^4}{16} \frac{f_{S,1}}{\Lambda^4}$$

Linear-EFT scenarios with  $f_{S,0} \neq f_{S,2}$  need additional operator

$$\mathcal{L}_6 = \alpha_6 \text{Tr}[V_\mu V_\nu] \text{Tr}[TV^\mu] \text{Tr}[TV^\nu]$$

with  $T = \Sigma \sigma_3 \Sigma^\dagger$

isospin-breaking



Process:

$$pp \rightarrow W^+ Z j j$$

$$\rightarrow \ell^+ \nu \ell^+ \ell^- j j,$$

$\sqrt{S} = 13 \text{ TeV}$ , VBF cuts,  
NLO QCD

exactly the same picture  
as in  $W^+ W^+ j j$  case

