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**Interference effects in heavy Higgs
searches**

**Workshop on the physics of HL-LHC,
and perspectives at HE-LHC**

Geneva – November 2017



Outline

- 1 Classification of interferences
- 2 High-energy LHC
- 3 High-luminosity LHC
- 4 Conclusions

Relevance of off-shell SM Higgs decays to heavy gauge bosons $H \rightarrow VV^{(*)}$:

- ▷ First discussion: Off-shell contributions in $H \rightarrow VV^{(*)}$

[1206.4803; Kauer Passarino]

[1305.2092, 1310.7011; Kauer]

- ▷ Access to the Higgs width Γ_H

[1307.4935; Caola Melnikov]

[1311.3589, 1312.1628, 1408.1723;
Campbell Ellis Williams]

Application: see Roberto's talk!

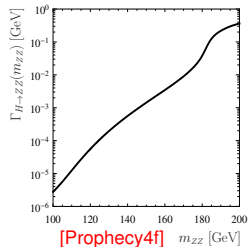
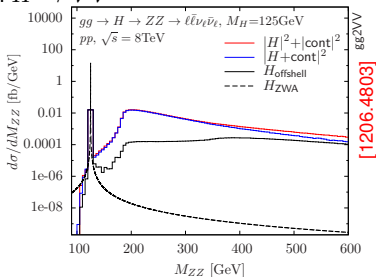
$$\sigma_{\text{OS}}^{VV} \propto \frac{(g_{ggH}^{\text{OS}} g_{HV}^{\text{OS}})^2}{\Gamma_H}$$

$$\frac{d\sigma^{VV}}{dm_{VV}} \propto (g_{ggH}^{\text{OFF}} g_{HV}^{\text{OFF}})^2$$

- ▷ A lot of effort went into the NLO description of the interferences!

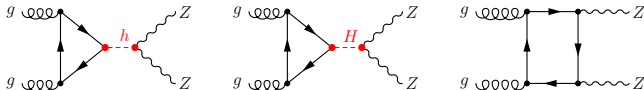
- ▷ Limitations of Higgs width constraint and opportunities:

[1310.1397, 1405.0285, 1405.1925, 1406.1757,
1406.6338, 1410.5440, 1410.5806, 1710.02149, etc.]



How does this discussion interfere with the searches for heavy Higgs bosons in VV channels (with $V \in \{W, Z\}$)?

Example in a Two-Higgs-Doublet Model (2HDM): [1512.07232 Greiner SL Weiglein]



→ The heavy Higgs H (mass m_H) can interfere with h and background B .

Classification of interferences depending on SM final state in $gg \rightarrow \Phi_i \rightarrow F$:
Heavy (CP-admixed) Higgs bosons ϕ_i , SM Higgs boson H , SM background B

$F =$	$b\bar{b}/\tau^+\tau^-$	$t\bar{t}$	$\gamma\gamma$	VV	HH	ZH
Int. between ϕ_i and ϕ_j	✓	✓	✓	✓	✓	✓
Int. between ϕ_i and B	(✓)	✓	✓	✓	✓	✓
Int. between ϕ_i and H			✓	✓	✓	✓

Similar tables can be produced for the other production mechanism, mainly VBF (NB ZH is listed as decay!)

Non-exhaustive list regarding interferences at the LHC (last 3 years):

▷ Final state $t\bar{t}/gg/\gamma\gamma$:

[1605.00542 Djouadi Ellis Quevillon]: $gg \rightarrow \phi \rightarrow t\bar{t}$ and $gg \rightarrow \phi \rightarrow \gamma\gamma$

[1608.07282 Carena Liu]: $gg \rightarrow \phi \rightarrow t\bar{t}$

[1606.04149 Hespel Maltoni Vryonidou]: $gg \rightarrow \phi \rightarrow t\bar{t}$ (2HDM, NLO)

[1707.06760 Franzosi Vryonidou Zhang]: $gg \rightarrow \phi \rightarrow t\bar{t}$ (NLO advanced)

[1606.03026 Martin]: $pp \rightarrow \phi \rightarrow gg$

[1511.05584 Bernreuther Galler Mellein Si Uwer]: $gg \rightarrow \phi \rightarrow t\bar{t}$

[1702.06063 Bernreuther Galler Mellein Si Uwer]: $gg \rightarrow \phi \rightarrow t\bar{t}$ (polarization, spin)

[1505.00291 Jung Song Yoon]: Generic discussion with complex phase (also $b\bar{b}$)

▷ Final state VV : (Consistent model due to unitarity needed!)

[1501.02139 Maina]: $gg \rightarrow \phi \rightarrow VV$ (SM+singlet)

[1502.04113 Kauer O'Brien]: $gg \rightarrow \phi \rightarrow VV$ (SM+singlet)

[1506.02257 Ballestrero Maina]: $VBF \rightarrow \phi \rightarrow VV$ (SM+singlet)

[1506.01694 Kauer O'Brien Vryonidou]: $gg \rightarrow \phi \rightarrow VV \rightarrow 4l$ (SM)

[1510.03450 Jung Song Yoon]: $gg \rightarrow \gamma\gamma/ZZ$ (2HDM)

[1512.07232 Greiner SL Weiglein]: $gg \rightarrow VV \rightarrow 4l$ (2HDM)

▷ Final state HH :

[1407.0281 Hespel Lopez-Val Vryonidou]: $gg \rightarrow \Phi \rightarrow HH$ (2HDM, NLO)

[1508.05397 Dawson Lewis]: $gg \rightarrow \Phi \rightarrow HH$ (SM+singlet, NLO)

▷ Interferences among heavy Higgs bosons:

[1411.4652 1705.05757 Fuchs Weiglein]: ϕ 's of the MSSM

Various works on Higgsstrahlung, not listed: $\gamma\gamma$ peak shift for light Higgs!

see the poster for m_t determination from $\gamma\gamma$ peak-dip structure by H. Yokoya

What happens at a high-energy LHC?

Consider the LHC running at $\sqrt{s} = 27$ rather than 13 TeV and exactly the same scenario of BSM physics:

The **relative interference** I among signal S and background B are mostly **identical**, since all contributions start with the same initial state (at a certain x of the PDFs). Thus, the overall cross section of S , B and I are different, but their relative fractions stay.

↔ **Interference** studies can be carried out at the partonic CMS energy $\sqrt{\hat{s}}$.

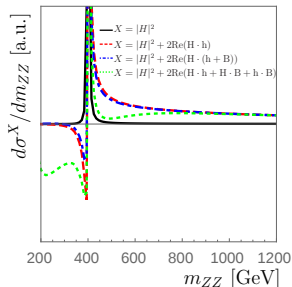
Effect at higher energies \sqrt{s} :

Probe different BSM scenarios with larger m_ϕ .

→ Higher masses m_ϕ usually come with larger

decay width Γ_ϕ and larger Γ_ϕ/m_ϕ .

→ Tendency to increase the relevance of **interference** effects (following NWA $\Gamma_\phi/m_\phi \ll 1$).



[1512.07232 Greiner SL Weiglein]

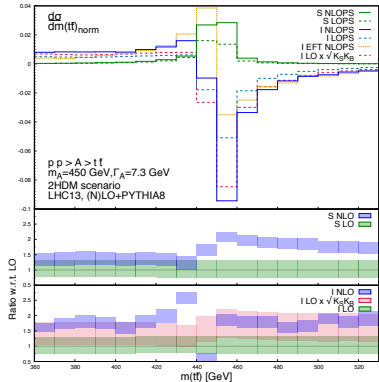
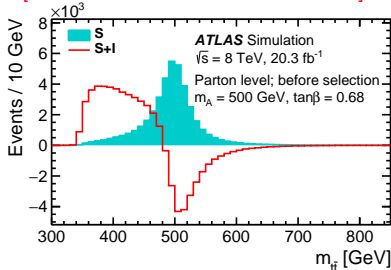
Of much larger relevance are high integrated luminosities, since they probe smaller signals S . Thus the relative fraction to the background S/B is smaller (B/S larger) and the **interference** I is larger. Rough estimate:

$$I \sim 2\text{Re}(\mathcal{A}_S \cdot \mathcal{A}_B)$$

$$\text{Fraction of interference over signal: } I/|\mathcal{A}_S|^2 \sim \mathcal{A}_B/\mathcal{A}_S$$

Already of relevance and taken into account for $t\bar{t}$:

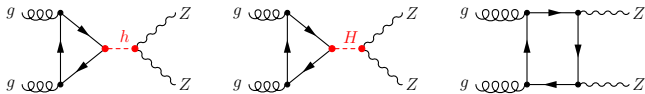
[ATLAS-CONF-2016-073 → 1707.06025]



[1707.06760 Franzosi Vryonidou Zhang]

At which level does the interference in VV play a role?

Consider again the 2HDM: [1512.07232 Greiner SL Weiglein]



→ The heavy Higgs H (mass m_H) can interfere with h and background B .

▷ Higgs mixing angle α from (H_1^0, H_2^0) to (h, H) and $\tan \beta = v_2/v_1$ defines

$$g_V^h = \sin(\beta - \alpha) = s_{\beta - \alpha} \quad \text{and} \quad g_V^H = \cos(\beta - \alpha) = c_{\beta - \alpha} \quad .$$

→ The combination $(g_V^h)^2 + (g_V^H)^2 = 1$ guarantees unitarization!

→ If $g_V^h \rightarrow 1$, then $g_V^H \approx 0$. Expect weak signal with large interferences!

Experimental analyses in heavy Higgs boson searches $H \rightarrow WW/ZZ$:

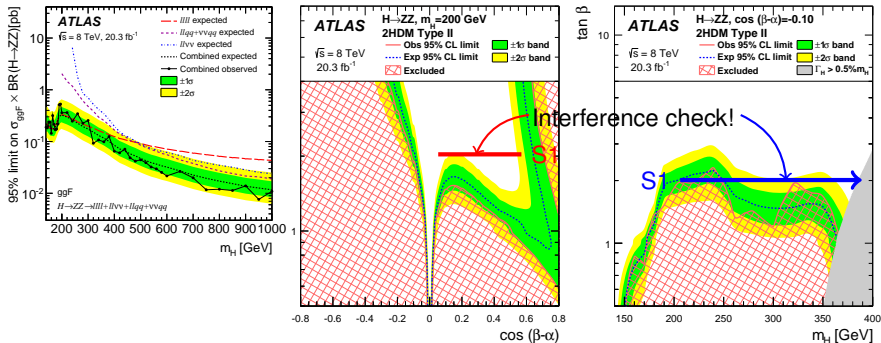
▷ ATLAS ($H \rightarrow ZZ$ in singlet+2HDM): No interference. [1507.05930]

▷ CMS ($H \rightarrow WW/ZZ$ in singlet): Rescaled SM $H \cdot B$ interference. [1504.00936]

Is the assumption of no interference justified?

- ▷ In Run I (and partially Run II) it was indeed ok.
- ▷ At HL-LHC experiments are getting sensitive to interferences in H searches.

Check of the 8 TeV ATLAS analysis (result of various channels): [1507.05930]



Implementation of $gg \rightarrow ZZ$ in a 2HDM in `vh@nnlo` [1210.5347, 1307.8122]

Implementation of $gg \rightarrow e^+e^-\mu^+\mu^-/e^+e^-\nu_l\nu_l$ in a 2HDM

with `GoSam` [1111.2034, 1404.7096]

with link to `2HDMC` [0902.0851] for the calculation of the Higgs width Γ_H

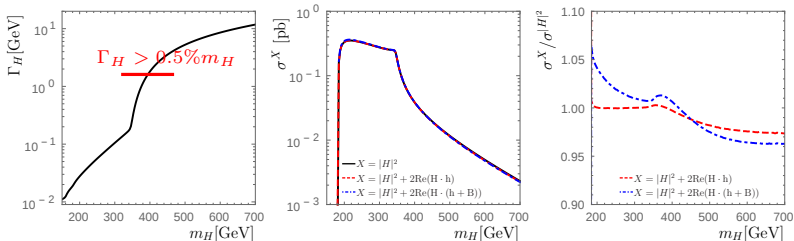
Scenario S1: 2HDM type II with $\tan \beta = 2$,

$\sin(\beta - \alpha) = -0.995 \leftrightarrow \cos(\beta - \alpha) = -0.10$, $m_H = 200 \text{ GeV}$, $\Gamma_H = 0.0277 \text{ GeV}$

Relevant parameters for variation: m_H , $\cos(\beta - \alpha)$ and $\tan \beta$

Let's start with the previously mentioned m_H variation (for $\sqrt{s} = 8 \text{ TeV}$):

$$\sigma^X = \int_{m_H - 15 \text{ GeV}}^{m_H + 15 \text{ GeV}} dm_{ZZ} \frac{d\sigma^X(gg \rightarrow ZZ)}{dm_{ZZ}}$$

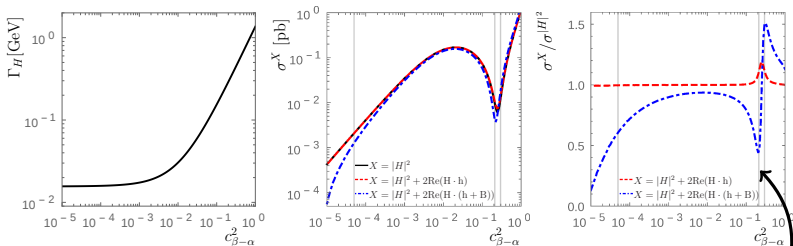


→ Relevance of the interferences remains below 10%.

It's not so much the mass m_H and width Γ_H which determine the relevance of the interferences.

We continue with a **variation of $\cos(\beta - \alpha)$** (for $\sqrt{s} = 8 \text{ TeV}$):

▷ Decay $H \rightarrow ZZ$ is determined by $g_V^H = \cos(\beta - \alpha) = c_{\beta-\alpha}$.

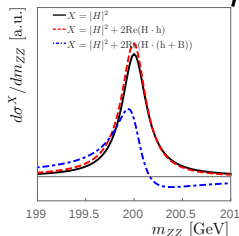


▷ Production $gg \rightarrow H$ involves g_t^H and g_b^H with

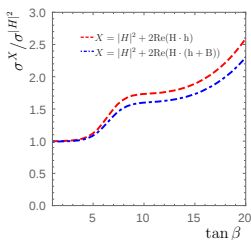
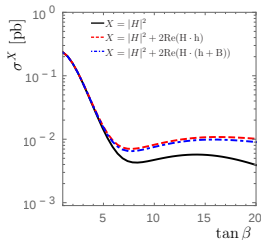
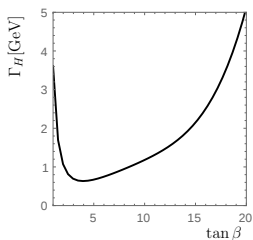
$$g_t^H = \frac{\sin \alpha}{\sin \beta} = -s_{\beta-\alpha} \frac{1}{\tan \beta} + c_{\beta-\alpha} \approx 0$$

$$\text{at } c_{\beta-\alpha}^2 = 0.2, \tan \beta = 2.$$

→ Where couplings are small (in production or decay) interferences are large.



We are left with a variation of $\tan \beta$ (for $\sqrt{s} = 13$ TeV):
 Scenario S2: 2HDM type II with $\tan \beta = 1$,
 $\sin(\beta - \alpha) = 0.990$, $m_H = 400$ GeV, $\Gamma_H = 3.605$ GeV



→ Pushing the bottom Yukawa g_b^H through $\tan \beta$ enhances interferences.

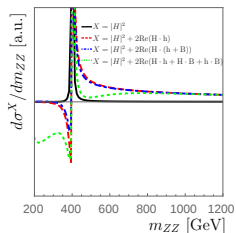
Not to forget at high invariant masses:

All interferences are of relevance:

$h \cdot B$ and $H \cdot B \leftrightarrow$ negative.

$h \cdot H$ can be large and have either sign!

Another effect: Sum of interferences can mimic a **broad peak-like structure** beyond m_H .

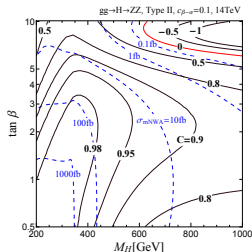
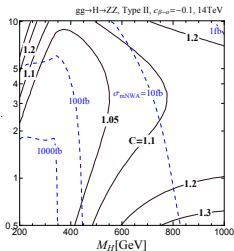
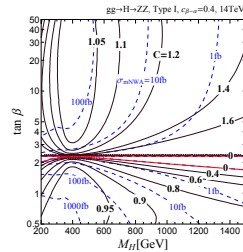
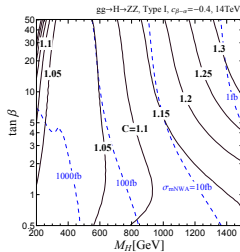
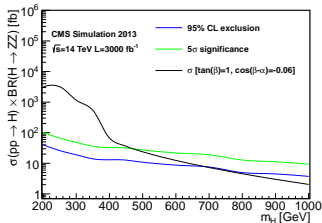


Similar conclusion for the 2HDM derived in: [1510.03450 Jung Song Yoon]

$$C = \frac{\sigma_{\text{mNWA}}}{\sigma_{\text{NWA}}}$$

mNWA = modified NWA
with interferences

[CMS-PAS-FTR-13-024]

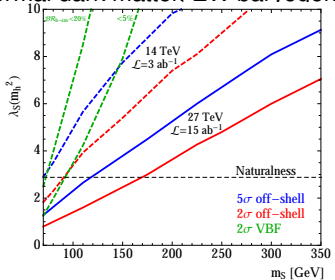
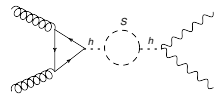


Coming back to the interference of the light Higgs H with the background:
Sensitivity to Higgs portal scenarios: [1710.02149 Gonçalves Han Mukhopadhyay]
Singlet S coupling only via the Higgs sector

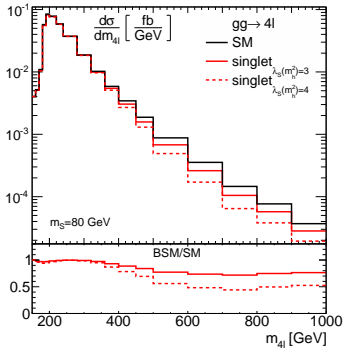
$$\mathcal{L} \supset \partial_\mu S \partial^\mu S^* - \mu^2 |S|^2 - \lambda_S |S|^2 |H|^2 \quad \rightarrow \text{see Matthew McCullough's talk}$$

based on prospects in:
[ATL-PHYS-PUB-2015-024]

can address: naturalness,
Thermal dark matter, EW barvoogenesis



\rightarrow typical in DM portals:
see Anne-Marie Magnan's talk



The relevance of **interference contributions** in the search for heavy Higgs bosons is process dependent. In the last years various publications looked at interferences at the LHC in different BSM models.

Two main points:

- ▷ Going to **higher energies** will not alter the relative interferences in a fixed BSM scenario. Though higher energies probe larger masses and thus Γ_ϕ/m_ϕ gets larger.
- ▷ Going to **higher luminosities** probes smaller signal cross sections and this has a profound effect on the relevance of interferences. Many processes will have to take into account interferences sooner (e.g. $t\bar{t}$) or later (VV).

Outlook:

- ▷ There is a LesHouches project ongoing to classify interferences in VV and HH (We are a bit behind schedule though...). We could add a higher mass scenario with a larger width.
- ▷ In the context of this group work out a summary of the interferences as a function of the integrated luminosity of a HL-LHC, if possible.

Thank you for your attention!

Theoretical issues:

▷ Precision for gluon fusion ($gg \rightarrow VV$):

S: meanwhile known at N³LO QCD (since 2016) (applicable as fct. of m_{VV})

B: meanwhile known at NLO QCD (since 2014+2015)!

[1404.4853; Gehrmann von Manteuffel Tancredi Weihs: Master integrals]

[1503.08759; Caola Henn Melnikov Smirnov Smirnov: Helicity amplitudes]

[1503.08835; von Manteuffel Tancredi: Helicity amplitudes]

towards pheno predictions in the heavy top-limit:

[1503.01274; Melnikov Dowling]

[1509.06734, 1511.08617, 1605.04610; Caola (Dowling) Melnikov Röntsch Tancredi (ZZ/WW)]

[1605.01380; Campbell Ellis Czakon Kirchner: interference + top-quark mass effects]

[1609.09719; Alioli Caola Luisoni Röntsch (ZZ matched to parton shower)]

Related work on $q\bar{q} \rightarrow VV$: [1408.6409, 1503.04812]

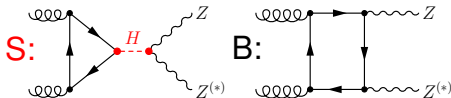
[1710.06294; Heinrich Jahn Jones Kerner Pires (ZZ NNLO)]

Previously K-factor

$$R_H^B = \frac{K(gg \rightarrow ZZ)}{K(gg \rightarrow H^{(*)} \rightarrow ZZ)}$$

in ATLAS **bound on Γ_H !**

[1503.01060]



▷ Other issues:

Dominant background $q\bar{q} \rightarrow VV$ known at NNLO QCD [MATRIX Grazzini et al.]

Interference of $WW \rightarrow 2l2\nu$ and $ZZ \rightarrow 2l2\nu$ known (e.g. $gg2VV$ [Kauer et al.]