Cosmic-ray propagation in self-generated turbulence

Carmelo Evoli

Gran Sasso Science Institute - L'Aquila

Three-elephants in the gamma-ray sky @Garmisch-Partenkirchen



in collaboration with R. Aloisio, P. Blasi and G. Morlino (GSSI)

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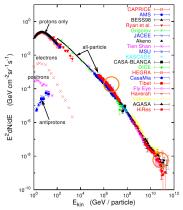
Non-linear CR propagation

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The cosmic-ray spectrum

Credit: Hillas+06

Energies and rates of the cosmic-ray particles



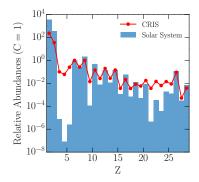
- Almost a perfect power-law over 12 energy decades.
- Evidence of the "knee" and "ankle" features.
- Observed at energy higher than terrestrial laboratories.
- Direct measurements versus air-cascade reconstructions.
- Composition:
 - $\sim 98\%$ are nuclei
 - $\sim 87\%~{\rm protons}$
 - $\sim 12\%~{\rm He}$
 - $\sim 1\%$ heavier nuclei
 - $\sim 2 \%$ are electrons

 $\sim 0.1\%$ are anti-matter particles (positrons and antiprotons)

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The grammage pillar



From this plot it follows the more robust evidence of diffusion so far:

$$c au_{
m esc} = rac{X(E)}{ar{n}_{
m ISM} \mu} \sim 10^3 \, {
m kpc} \, \gg \, {
m Galaxy} \, {
m size}$$

and it suggests that SN explosions can sustain the galactic CR population:

$$L_{\rm CR} = \frac{\epsilon_{\rm CR} V_{\rm MW}}{\tau_{\rm esc}} \sim 0.1 \div 0.5 L_{\rm SN}$$

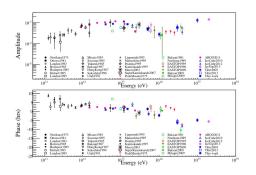
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The anisotropy puzzle

Di Sciascio & luppa, 2014



- dipole amplitude increases up to ~10 TeV and then it decreases
- phase of dipole steadily migrates and suddenly flips

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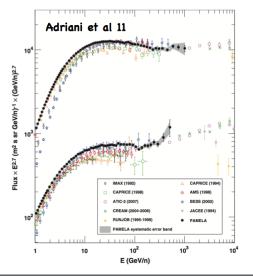
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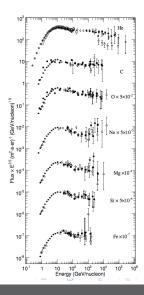
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$$A \sim \frac{v_A}{c} \sim 10^{-4} \frac{v_A}{30 \,\mathrm{km/s}}$$

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Elemental spectra



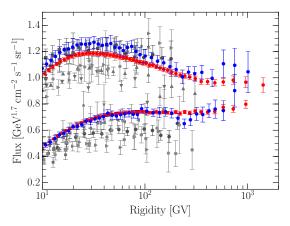


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PAMELA and AMS-02: The precision era

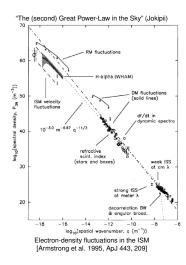
Adriani et al., Science, 2011; Aguilar et al., PRL, 2015



PAMELA and AMS-02 measurements of proton and helium (x10) spectrum.

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The interstellar turbulence



- Turbulence is stirred by Supernovae at a typical scale $L \sim 10 100 \text{ pc}$
- Fluctuations of velocity and magnetic field are Alfvénic
- ► They have a Kolmogorov $k^{-5/3}$ spectrum (density is a passive tracer so it has the same spectrum: $\delta n_e \sim \delta B$):

$$W(k)dk \equiv \frac{\langle \delta B \rangle^2(k)}{B_0^2} = \frac{2}{3} \frac{\eta_B}{k_0} \left(\frac{k}{k_0}\right)^{-5/3}$$

• where $k_0 = L^{-1}$ and the *level of turbulence* is

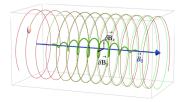
$$\eta_B = \int_{k_0}^\infty dk \, W(k) \sim 0.1 \div 0.01$$

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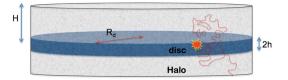
Charged particle in a turbulent field



- The turbulent field amplitude is a small fluctuation with respect to the regular component
- ▶ Resonant interaction wave-particle: $k_{\rm res}^{-1} \sim r_L(p)$
- It follows:

$$D_{\rm xx}(p) = \frac{vr_L}{3} \frac{1}{k_{\rm res}W(k_{\rm res})} \sim \frac{3 \times 10^{27}/\eta_B \,{\rm cm}^2/{\rm s}}{3 \times 10^{28} \,{\rm cm}^2/{\rm s}} \left(\frac{p}{{\rm GeV/c}}\right)^{1/3}$$

The transport equation



$$\frac{\partial N_i}{\partial t} - \vec{\nabla} \cdot \left(D_{\rm xx} \cdot \vec{\nabla} N_i - \vec{u} N_i \right) = \underbrace{Q_{\rm SN}}_{\rm gain/sinks} = Q_{\rm gain/sinks}$$

- ▶ Spatial diffusion: $\vec{\nabla} \cdot \vec{J} = Q$
- Advection by Galactic winds/outflows: $\vec{u} = u_z^w$
- Source term proportional to Galactic SN profile
- Energy losses: ionization, Bremsstrahlung, IC, Synchrotron, ...
- Production of light nuclei due to the inelastic scattering or the decay of heavier species

Predictions of the standard picture

For a primary CR species (e.g., H, C, O) at **high energy** we can ignore energy gain/losses, and the transport equation can be simplified as:

$$\frac{\partial N}{\partial t} = Q_0(p)\delta(z) + \frac{\partial}{\partial z} \left[D \frac{\partial N}{\partial z} \right]$$

For $z \neq 0$ one has:

$$D\frac{\partial N}{\partial z} = \text{constant} \rightarrow N(z) = N_0 \left(1 - \frac{z}{H}\right)$$

where we used the definition of a *halo*: $N(z = \pm H) = 0$. The typical solution gives:

$$N_0(p) = \frac{Q_0(p)}{2A_{\rm d}} \frac{H}{D(p)} \sim p^{-\gamma - \delta}$$

For a secondary (e.g., Li, Be, B) the source term is proportional to the primary density:

$$Q_B \sim \bar{n}_{\rm ISM} c \sigma_{C \to B} N_C \to \frac{N_B}{N_C} \sim \frac{H}{D_0} p^{-\delta}$$

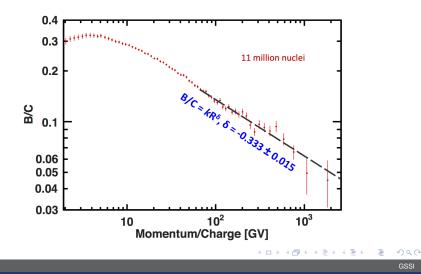
where we use $\bar{n}_{ISM} = n_{disk}h/H$.

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Comparison with B/C as measured by AMS-02

Aguilar et al., PRL, 2016



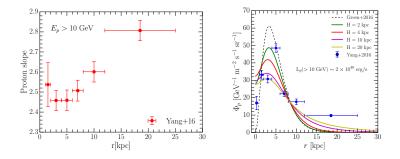
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Predictions of the standard picture

- By solving the transport equation we obtain a featureless (at least up to the knee) propagated spectrum for each primary species, at the odds with observations.
- This result remains true even in more sophisticated approach as GALPROP or USINE or DRAGON or PICARD (in order of appearance)
- What is missing in our physical picture?

Limits of standard approach: not-local observations

Yang, Aharonian, Evoli, PRD, 2016



The one-zone model for transport implemented in most of the numerical approaches fails in reproducing not-local observation as γ -ray diffuse and anisotropy.

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The halo size H

- Assuming N(z = H) = 0 reflects the requirement of lack of diffusion (infinite diffusion coefficient)
- ▶ May be because $B \rightarrow 0$, or because turbulence vanishes (in both cases *D* cannot be spatially constant!)
- Vanishing turbulence may reflect the lack of sources
- ▶ Can be *H* dependent on *p*?
- ▶ What is the physical meaning of *H*?

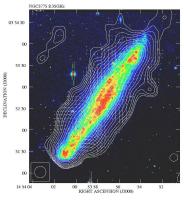
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The radio halo in external galaxies

Credit: MPIfR Bonn



Total radio emission and B-vectors of edge-on galaxy NGC891, observed at 3.6 cm wavelength with the Effelsberg telescope



Total radio intensity and B-vectors of edge-on galaxy NGC 5775, combined from observations at 3.6 cm wavelength with

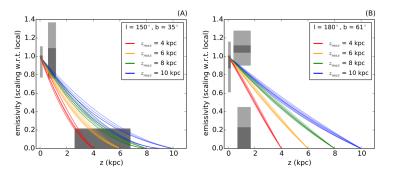
the VLA and Effelsberg telescopes

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The γ -halo in our Galaxy

Tibaldo et al., 2015, ApJ



- Using high-velocity clouds one can measure the emissivity per atom as a function of z (proportional to N)
- Indication of a halo with $H \sim$ few kpc

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Non-linear cosmic ray transport

Skilling71, Wentzel74

- CR energy density is ~ 1 eV/cm⁻³ in equipartition with: starlight, turbulent gas motions and magnetic fields.
- In these conditions, low energy can self-generate the turbulence for their scattering (notice that self-generated waves are k ~ r_L)
- Waves are amplified by CRs through streaming instability:

$$\Gamma_{\rm CR} = \frac{16\pi^2}{3} \frac{v_A}{kW(k)B_0^2} \left[p^4 v(p) \frac{\partial N}{\partial z} \right]$$

and are damped by wave-wave interactions that lead the development of a turbulent cascade (NLLD):

$$\Gamma_{\rm NLLD} = (2c_k)^{-3/2} k v_A (kW)^{1/2}$$

What is the typical scale/energy up to which self-generated turbulence is dominant?

Non-linear cosmic ray transport

Blasi, Amato & Serpico, PRL, 2012

Transition occurs at scale where external turbulence (e.g., from SNe) equals in energy density the self-generated turbulence

$$W_{\rm ext}(k_{\rm tr}) = W_{\rm CR}(k_{\rm tr})$$

where $W_{\rm CR}$ corresponds to $\Gamma_{\rm CR} = \Gamma_{\rm NLLD}$ Assumptions:

- Quasi-linear theory applies
- The external turbulence has a Kolmogorov spectrum
- Main source of damping is non-linear damping
- $\blacktriangleright\,$ Diffusion in external turbulence explains high-energy flux with SNR efficiency of $\epsilon\sim10\%$

$$E_{\rm tr} = 228 \, {\rm GeV} \, \left(\frac{R_{d,10}^2 H_3^{-1/3}}{\epsilon_{0.1} E_{51} \mathcal{R}_{30}} \right)^{3/2(\gamma_p - 4)} B_{0,\mu}^{(2\gamma_p - 5)/2(\gamma_p - 4)}$$

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The turbulence evolution equation

$$\frac{\partial W}{\partial t} = \frac{\partial}{\partial k} \left[D_{kk} \frac{\partial W}{\partial k} \right] + \frac{\partial}{\partial z} \left(v_A W \right) + \Gamma_{\rm CR} W + Q(k)$$

• Diffusion in *k*-space damping: $D_{kk} = c_k |v_A| k^{7/2} W^{1/2}$

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- ▶ Diffusion in *k*-space damping: $D_{kk} = c_k |v_A| k^{7/2} W^{1/2}$
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- Waves growt due to cosmic-ray streaming: $\Gamma_{\rm CR} \propto \partial N/\partial z$

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- External (e.g., SNe) source term $Q \sim \delta(z)\delta(k-k_0)$

The turbulence evolution equation

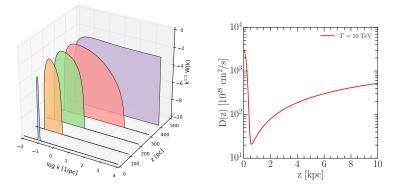
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- Advection of the Alfvén waves
- ▶ Waves growt due to cosmic-ray streaming: $\Gamma_{\rm CR} \propto \partial N / \partial z$
- External (e.g., SNe) source term $Q \sim \delta(z)\delta(k k_0)$
- ► In the absence of the instability, it returns a kolmogorov spectrum: $W(k) \sim k^{-5/3}$

Non-linear cosmic ray transport: the turbulent halo

Evoli+2017, in preparation



$$au_{
m cascade} = au_{
m adv} o rac{k_0^2}{D_{kk}} = rac{z_{
m peak}}{v_A} o z_{
m peak} \sim \mathcal{O}(
m kpc)$$

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Non-linear CR propagation

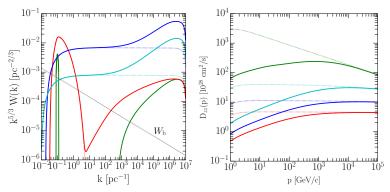
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Non-linear cosmic ray transport: a global picture

Evoli+2017, in preparation

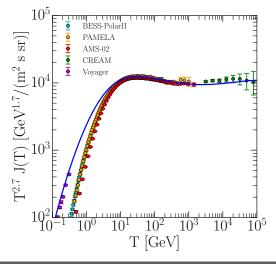


Turbulence spectrum without (dotted) and with (solid) CR self-generated waves at different distance from the galactic plane.

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Non-linear cosmic ray transport: a global picture

Evoli+2017, in preparation



- Pre-existing waves (Kolmogorov) dominates above the break
- Self-generated turbulence between 1-100 GeV
- Voyager data are reproduced with no additional breaks, but due to advection with self-generated waves
- ▶ No *H* is assumed here

The elephant in the glassware



In Fig: Cosmic rays with \sim GeV energy propagating in the interstellar turbulence. We assume that nothing is going to happen.

Conclusions

- Recent findings by PAMELA and AMS-02 (breaks in the spectra of primaries, B/C a la Kolmogorov, flat anti-protons, rising positron fraction) are challenging the standard scenario of CR propagation.
- Non-linearities might play an essential role for propagation (as they do for acceleration). They allow to reproduce local observables (primary spectra) without ad hoc breaks.
- ▶ We present a non-linear model in which SNRs inject: a) turbulence at a given scale with efficiency $\epsilon_{\rm w} \sim 10^{-4}$ and b) cosmic-rays with a single power-law and $\epsilon_{\rm CR} \sim 10^{-1}$. The turbulent halo and the change of slope at ~200 GV are obtained self-consistently.
- As a bonus, these models enable us a deeper understanding of the interplay between CR, magnetic turbulence and ISM in our Galaxy.