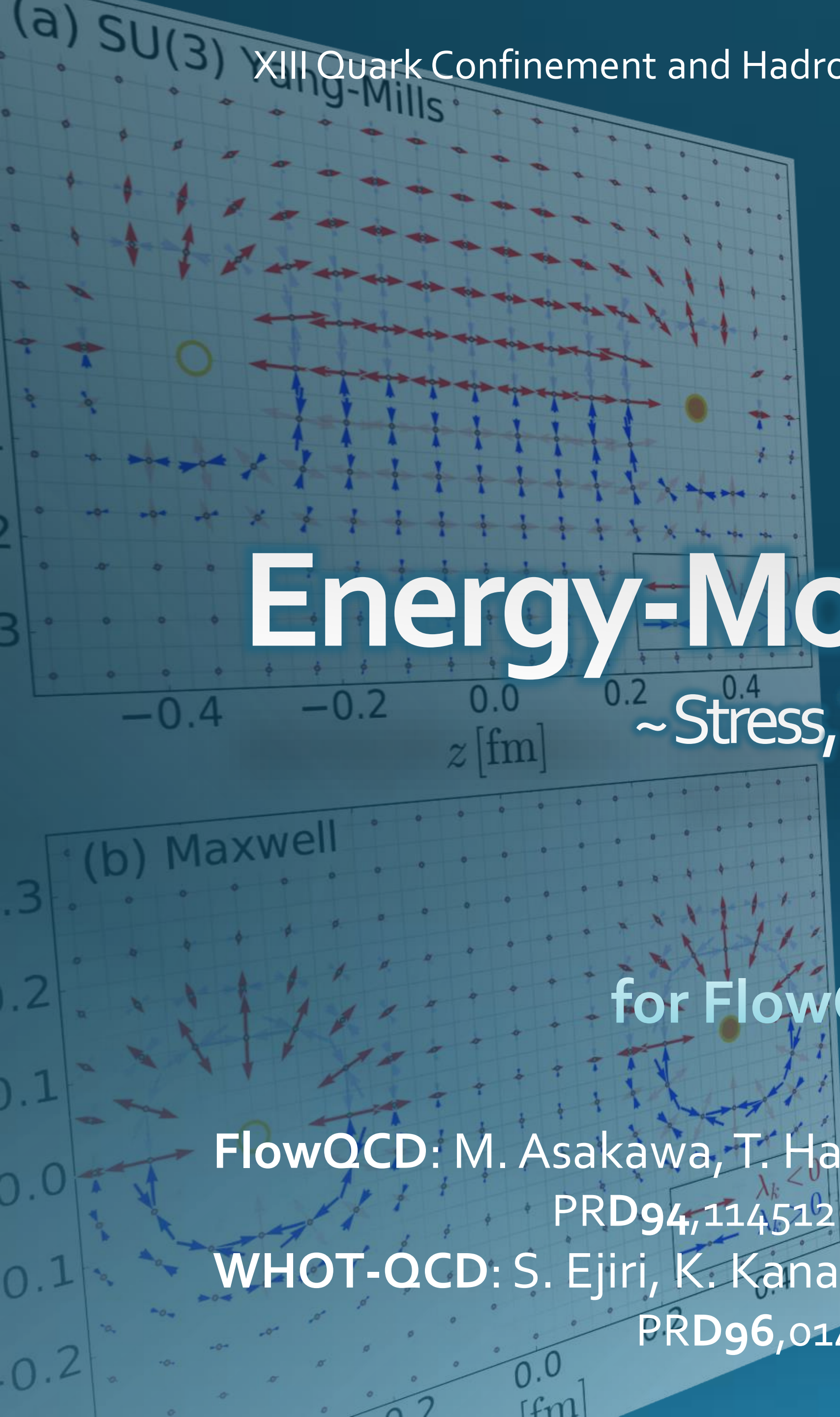
XIII Quark Confinement and Hadron Spectroscopy, 4 Aug., 2018, Maynooth, Ireland



Exploring Non-Abelian Gauge Theory with Energy-Momentum Tensor

~ Stress, Thermodynamics and Correlations ~

Masakiyo Kitazawa

for FlowQCD / WHOT-QCD Collaborations

FlowQCD: M. Asakawa, T. Hatsuda, **T. Iritani**, H. Suzuki, **R. Yanagihara**

PR**D94**,114512(2016); PR**D96**,111502(2017); arXiv:1803.05656

WHOT-QCD: S. Ejiri, K. Kanaya, H. Suzuki, **Y. Taniguchi**, T. Umeda, ...

PR**D96**,014509(2017); arXiv:1710.10015; arXiv:1711.02262

Energy-Momentum Tensor

One of the most fundamental quantities in physics

$$T_{\mu\nu}$$

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = \kappa T_{\mu\nu}$$

$$\partial_{\mu} T_{\mu\nu} = 0$$

Energy-Momentum Tensor

energy momentum

$$T_{\mu\nu} = \begin{bmatrix} \boxed{T_{00}} & T_{01} & T_{02} & T_{03} \\ T_{10} & T_{11} & T_{12} & T_{13} \\ T_{20} & T_{21} & T_{22} & T_{23} \\ T_{30} & T_{31} & T_{32} & T_{33} \end{bmatrix}$$

pressure

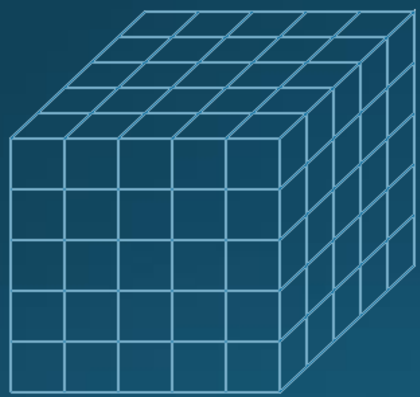
stress

All components are important physical observables!

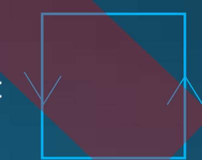
$T_{\mu\nu}$: nontrivial observable on the lattice

① Definition of the operator is nontrivial

because of the explicit breaking of Lorentz symmetry



ex: $T_{\mu\nu} = F_{\mu\rho}F_{\nu\rho} - \frac{1}{4}\delta_{\mu\nu}FF$

$F_{\mu\nu} =$ 

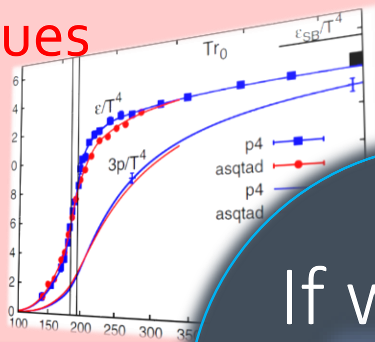
② Its measurement is extremely noisy

due to high dimensionality and etc.

Thermodynamics

direct measurement of expectation values

$\langle T_{00} \rangle, \langle T_{ii} \rangle$



Fluctuations and Correlations

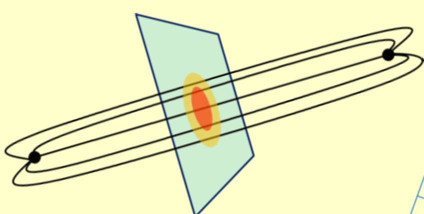
viscosity, specific heat, …

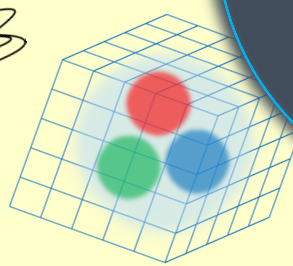
$$c_V \sim \langle \delta T_{00}^2 \rangle$$

$$\eta = \langle T_{12} ; T_{12} \rangle$$

If we have

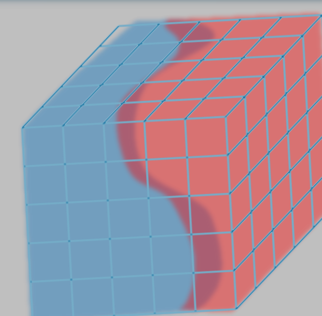
$$T_{\mu\nu}$$





- flux tube / hadrons
- stress distribution

Hadron Structure

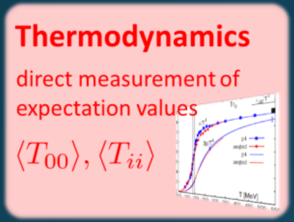


- vacuum configuration
- mixed state on 1st transition

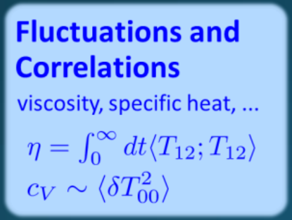
Vacuum Structure

Contents

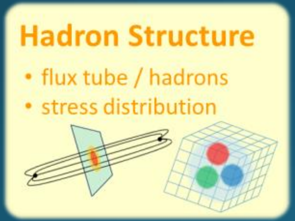
Constructing EMT on the lattice



Thermodynamics



EMT Correlation Function



Stress distribution in $\bar{q}q$ system

EMT on the Lattice: Conventional

Lattice EMT Operator

Caracciolo+, 1990

$$T_{\mu\nu} = Z_6 T_{\mu\nu}^{[6]} + Z_3 T_{\mu\nu}^{[3]} + Z_1 \left(T_{\mu\nu}^{[1]} - \langle T_{\mu\nu}^{[1]} \rangle \right)$$

$$T_{\mu\nu}^{[6]} = (1 - \delta_{\mu\nu}) F_{\mu\rho}^a F_{\nu\rho}^a, \quad T_{\mu\nu}^{[3]} = \delta_{\mu\nu} \left(F_{\mu\rho}^a F_{\nu\rho}^a - \frac{1}{4} F_{\rho\sigma}^a F_{\rho\sigma}^a \right), \quad T_{\mu\nu}^{[1]} = \delta_{\mu\nu} F_{\rho\sigma}^a F_{\rho\sigma}^a$$

- Fit to thermodynamics: Z_3, Z_1
- Shifted-boundary method: Z_6, Z_3 — Giusti, Meyer, 2011; 2013; Giusti, Pepe, 2014~; Borsanyi+, 2018

Multi-level algorithm

- effective in reducing statistical error of correlator — Meyer, 2007; Borsanyi, 2018; Astrakhantsev+, 2018

Yang-Mills Gradient Flow

$$\frac{\partial}{\partial t} A_\mu(t,x) = -\frac{\partial S_{\mathrm{YM}}}{\partial A_\mu}$$

Luscher 2010
Narayanan, Neuberger, 2006
Luscher, Weiss, 2011

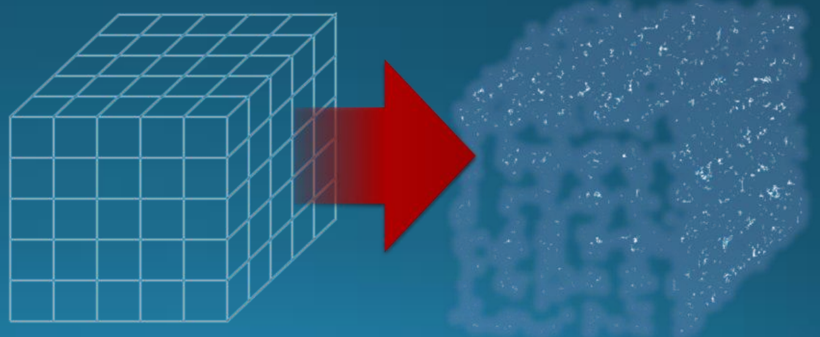
t: "flow time"
dim:[length²]

$$A_\mu(0,x) = A_\mu(x)$$

leading

$$\boxed{\partial_t A_\mu} = D_\nu G_{\mu\nu} = \boxed{\partial_\nu \partial_\nu A_\mu} + \cdots$$

- diffusion equation in 4-dim space
- diffusion distance $d \sim \sqrt{8t}$
- "continuous" cooling/smearing
- No UV divergence at t>0



Yang-Mills Gradient Flow

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leading

$$\boxed{\partial_t A_\mu} = D_\nu G_{\mu\nu} = \boxed{\partial_\nu \partial_\nu A_\mu} + \cdots$$

Applications

scale setting / topological charge / running coupling
noise reduction / defining operators / ...

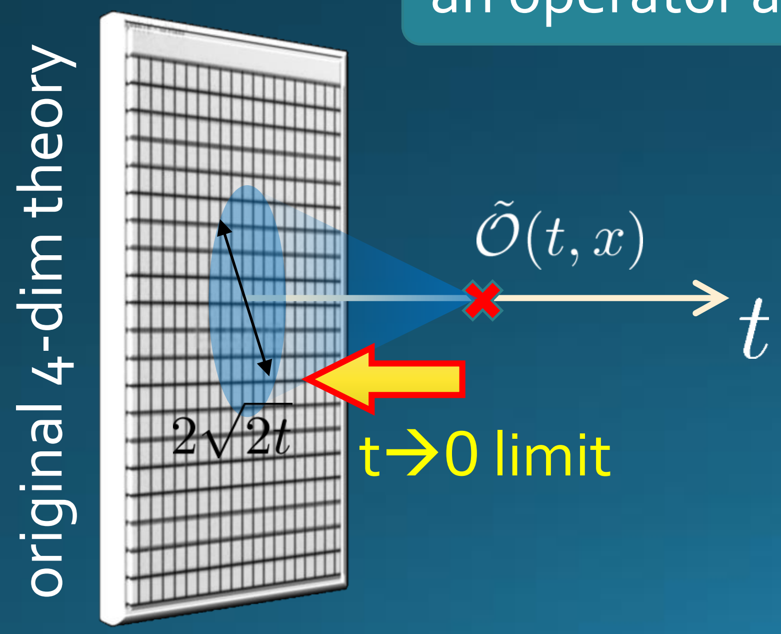
Small Flow-Time Expansion

Luescher, Weisz, 2011
Suzuki, 2013

$$\tilde{\mathcal{O}}(t,x) \xrightarrow[t\to 0]{} \sum_i c_i(t)\mathcal{O}_i^R(x)$$

an operator at t>0

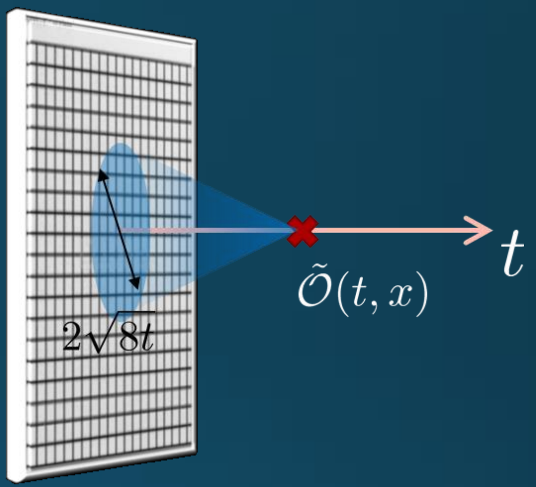
remormalized operators of original theory



Constructing EMT 1

Suzuki, 2013

$$\tilde{\mathcal{O}}(t,x) \xrightarrow[t\to 0]{} \sum_i c_i(t)\mathcal{O}_i^R(x)$$



- Gauge-invariant dimension 4 operators

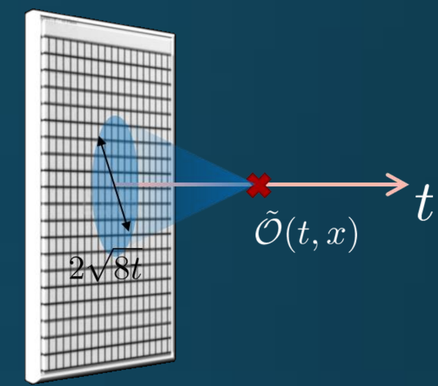
$$\begin{cases} U_{\mu\nu}(t,x) = G_{\mu\rho}(t,x)G_{\nu\rho}(t,x) - \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t,x)G_{\mu\nu}(t,x) \\ E(t,x) = \frac{1}{4}\delta_{\mu\nu}G_{\mu\nu}(t,x)G_{\mu\nu}(t,x) \end{cases}$$

Constructing EMT 2

Suzuki, 2013

$$U_{\mu\nu}(t,x) = \alpha_U(t)\left[T^R_{\mu\nu}(x) - \frac{1}{4}\delta_{\mu\nu}T^R_{\rho\rho}(x)\right] + \mathcal{O}(t)$$

$$E(t,x) = \langle E(t,x)\rangle + \alpha_E(t)T^R_{\rho\rho}(x) + \mathcal{O}(t)$$



Suzuki coeffs. $\left\{\begin{aligned} \alpha_U(t) &= g^2\left[1 + 2b_0 s_1 g^2 + O(g^4)\right] \\ \alpha_E(t) &= \frac{1}{2b_0}\left[1 + 2b_0 s_2 g^2 + O(g^4)\right] \end{aligned}\right.$

$g = g(1/\sqrt{8t})$

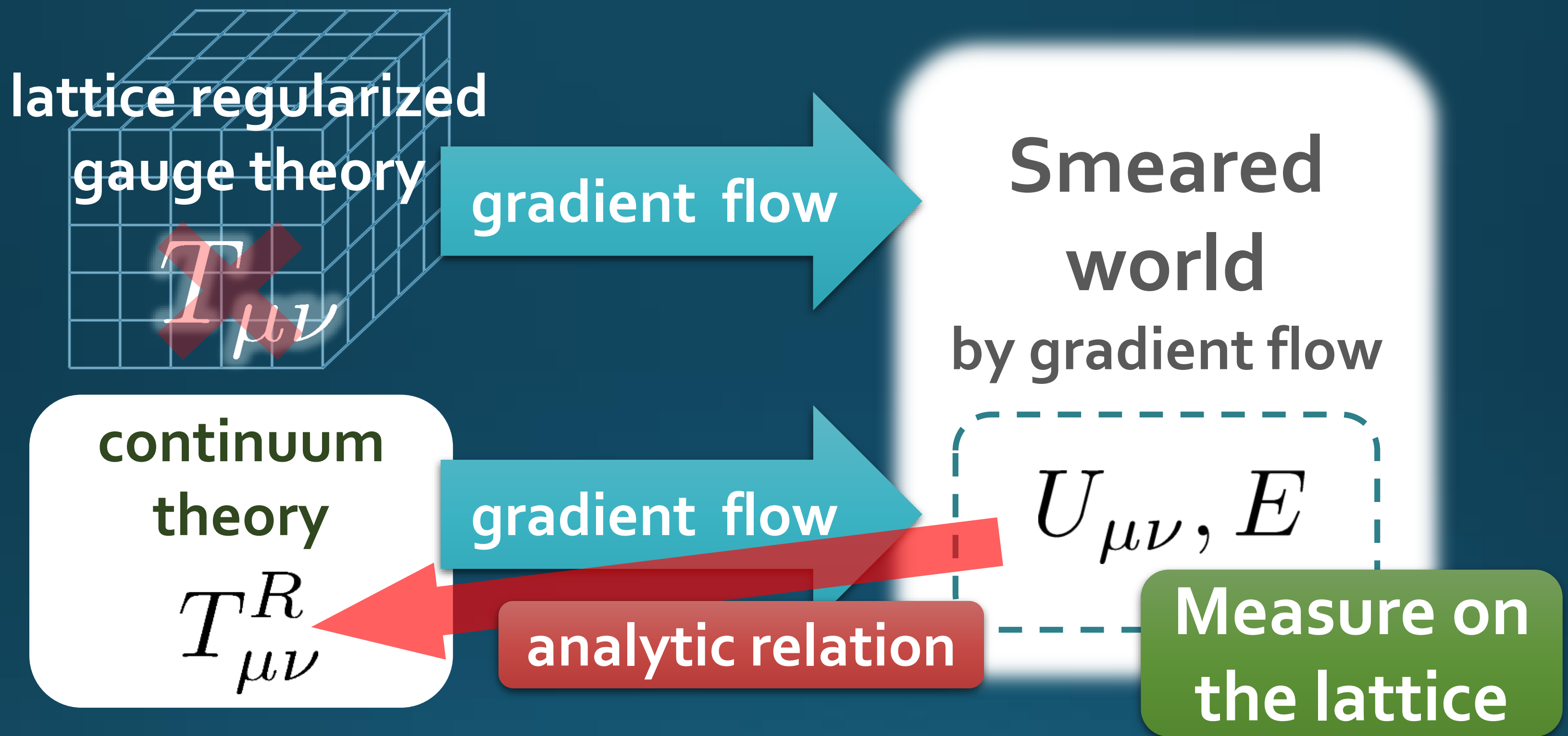
$s_1 = 0.03296\ldots$

$s_2 = 0.19783\ldots$

Remormalized EMT

$$T^R_{\mu\nu}(x) = \lim_{t\to 0}\left[\frac{1}{\alpha_U(t)}U_{\mu\nu}(t,x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)}E(t,x)_{\mathrm{subt.}}\right]$$

Gradient Flow Method



Take Extrapolation (t,a)→(0,0)

$$\langle T_{\mu\nu}(t)\rangle_{\rm latt} = \langle T_{\mu\nu}(t)\rangle_{\rm phys} + \boxed{C_{\mu\nu} t} + \boxed{D_{\mu\nu}\frac{a^2}{t}} + \cdots$$

O(t) terms in SFTE   lattice discretization

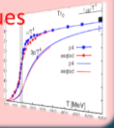
Contents

Constructing EMT on the lattice

Thermodynamics

direct measurement of expectation values

$\langle T_{00}\rangle, \langle T_{ii}\rangle$



Thermodynamics

Fluctuations and Correlations

viscosity, specific heat, ...

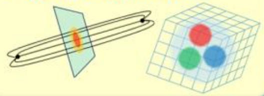
$\eta = \int_0^\infty dt \langle T_{12}; T_{12}\rangle$

$c_V \sim \langle \delta T_{00}^2\rangle$

EMT Correlation Function

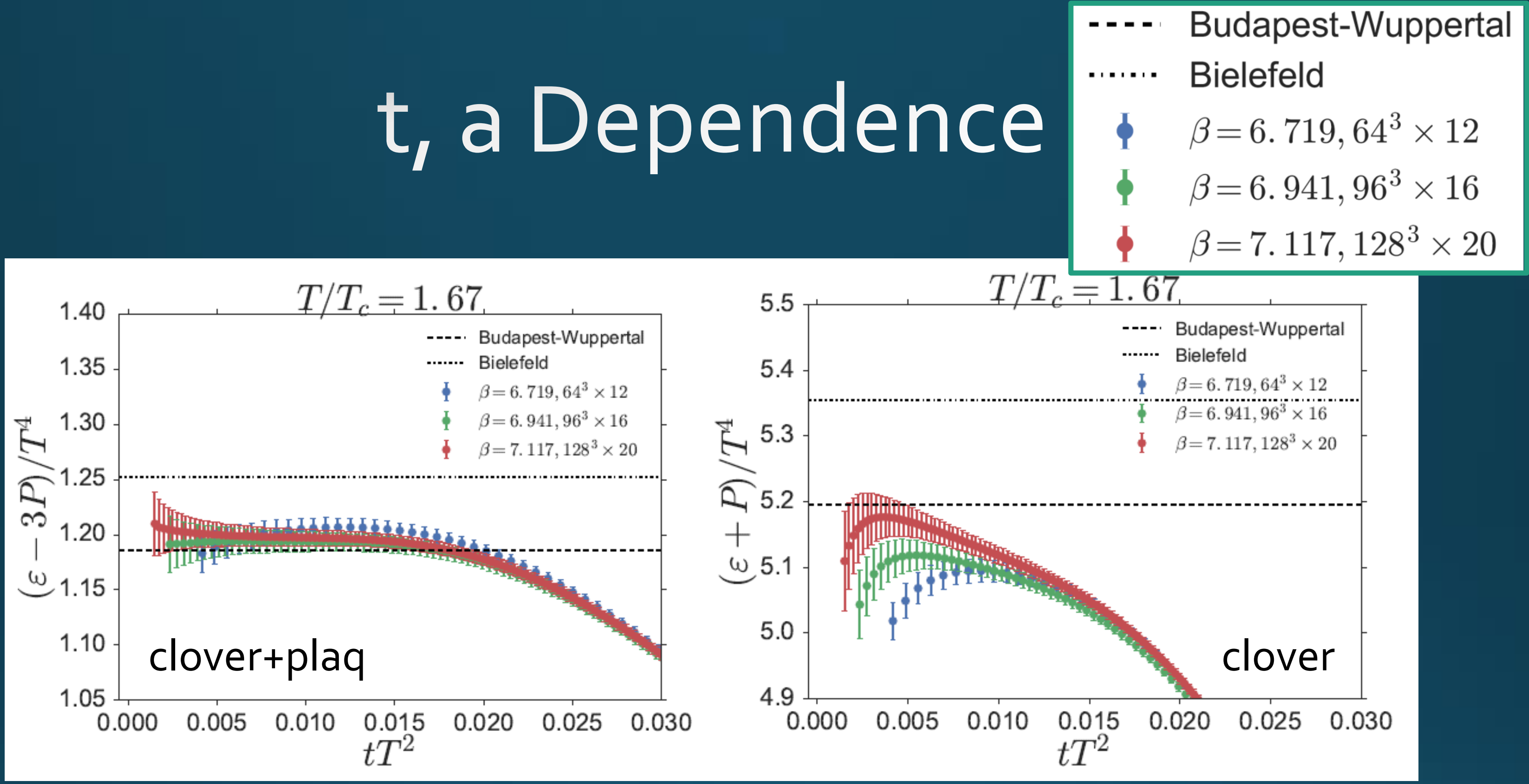
Hadron Structure

- flux tube / hadrons
- stress distribution



Stress distribution in $\bar{q}q$ system

t, a Dependence



$$\begin{cases} \sqrt{8t} < a & \text{: strong discretization effect} \\ \sqrt{8t} > 1/(2T) & \text{: over smeared} \end{cases}$$

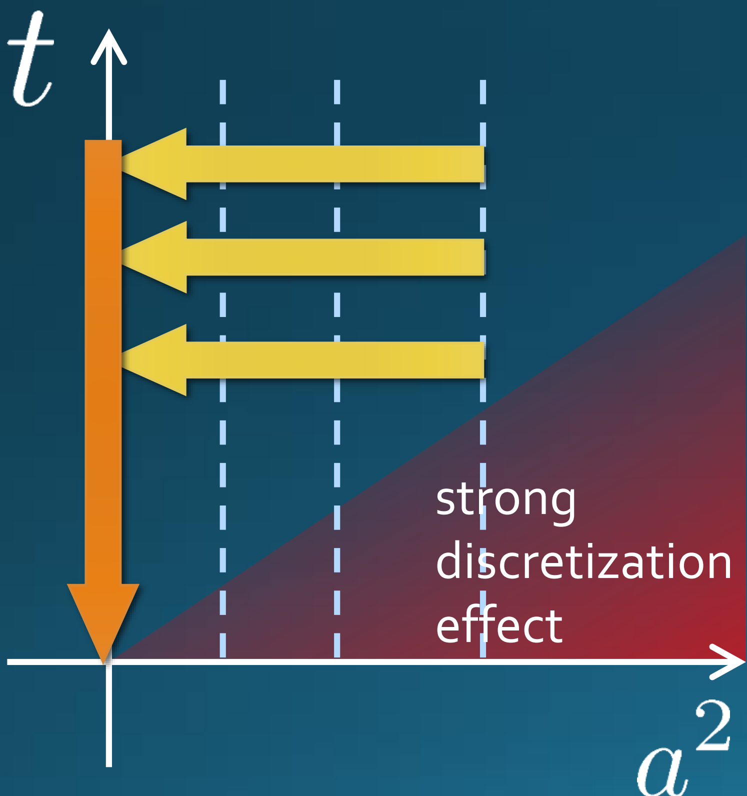
$a < \sqrt{8t} < 1/(2T)$: Linear t dependence

Double Extrapolation

t→0, a→0

$$\langle T_{\mu\nu}(t)\rangle_{\rm latt} = \langle T_{\mu\nu}(t)\rangle_{\rm phys} + \boxed{C_{\mu\nu} t} + \boxed{D_{\mu\nu}(t)\frac{a^2}{t}}$$

O(t) terms in SFTE lattice discretization



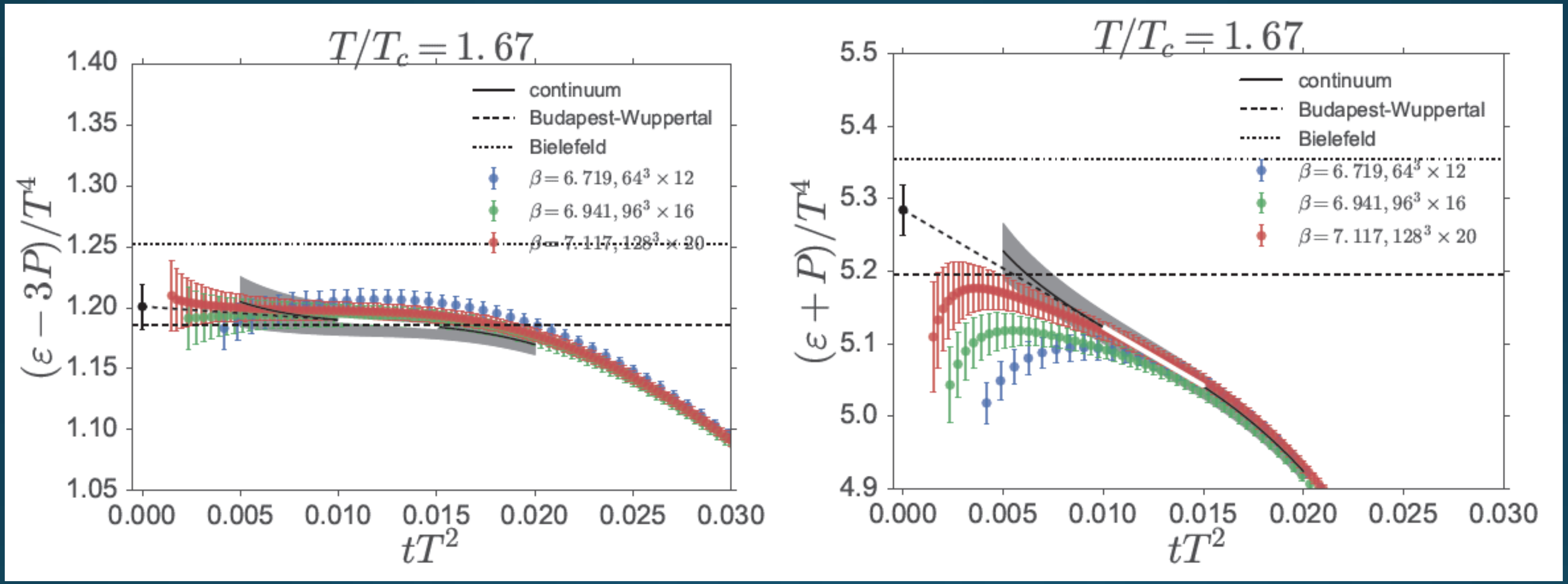
Continuum extrapolation

$$\langle T_{\mu\nu}(t)\rangle_{\rm cont} = \langle T_{\mu\nu}(t)\rangle_{\rm lat} + C(t)a^2$$

Small t extrapolation

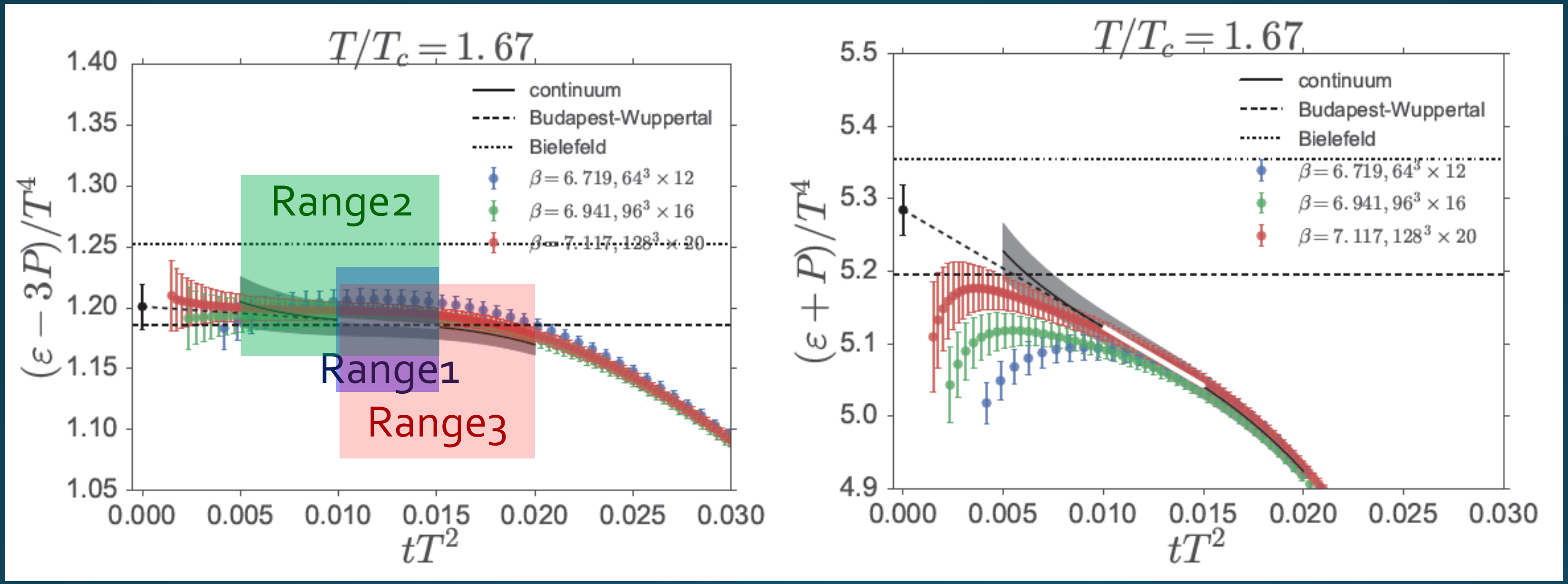
$$\langle T_{\mu\nu}\rangle = \langle T_{\mu\nu}(t)\rangle + C't$$

Double Extrapolation



Black line: continuum extrapolated

Double Extrapolation

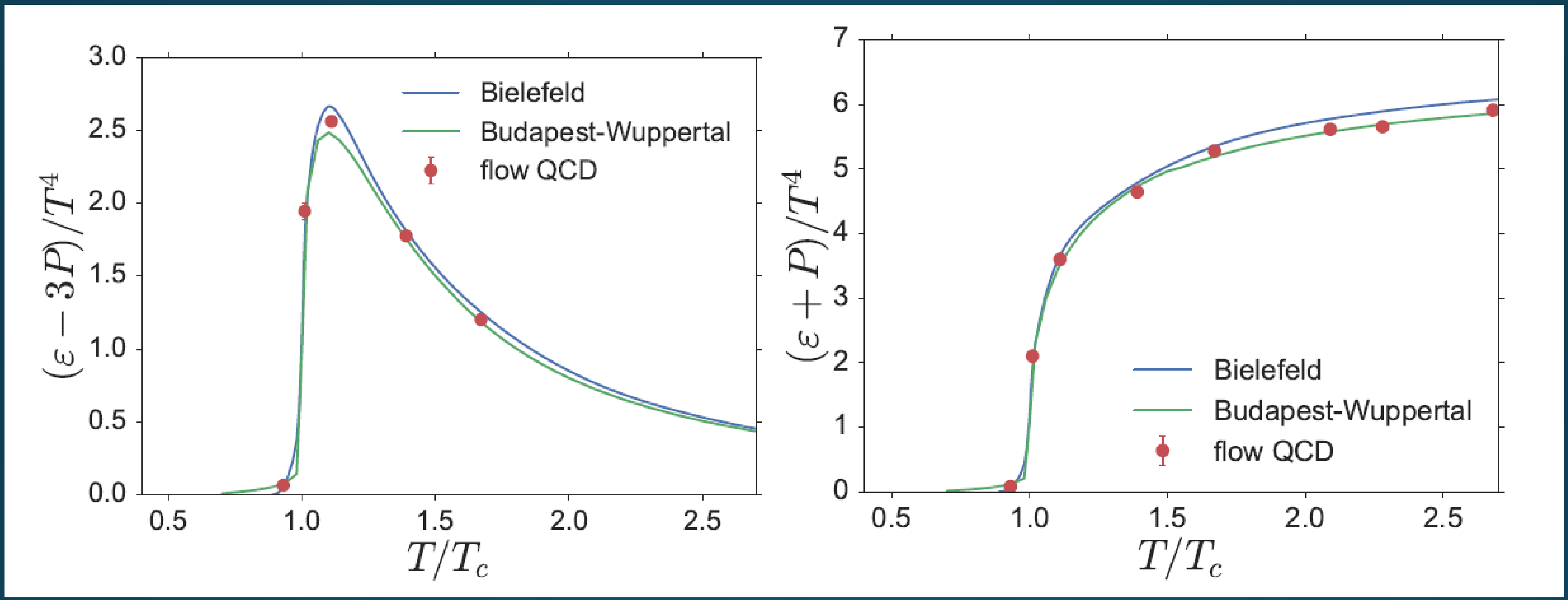


Black line: continuum extrapolated

- Fitting ranges:
  - range-1: $0.01 < tT^2 < 0.015$
  - range-2: $0.005 < tT^2 < 0.015$
  - range-3: $0.01 < tT^2 < 0.02$

Systematic error from the choice of fitting range ≈ statistical error

Temperature Dependence



Error includes

- statistical error
- choice of t range for t→0 limit
- uncertainty in aΛ_{MS}

total error <1.5% for T>1.1T_c

- Excellent agreement with integral method
- High accuracy only with ~2000 confs.

Thermodynamics on the Lattice

recent progress in SU(3) YM

- **Integral method**
  - Most conventional / established
  - Use themodynamic relations

    Boyd+ 1995; Borsanyi, 2012

$$p = \frac{T}{V} \ln Z$$

$$T \frac{\partial(p/T^4)}{\partial T} = \frac{\varepsilon - 3p}{T^4}$$

- Gradient-flow method
  - Take expectation values of EMT

    FlowQCD, 2014, 2016

$$\begin{cases} \varepsilon = \langle T_{00} \rangle \\ p = \langle T_{11} \rangle \end{cases}$$

- Moving-frame method

  Giusti, Pepe, 2014~
- Non-equilibrium method
  - Use Jarzynski's equality Caselle+, 2016;2018
- Differential method

  Shirogane+(WHOT-QCD), 2016~

SU(3) YM EoS: Comparison

$$\frac{e-3p}{T^4}$$

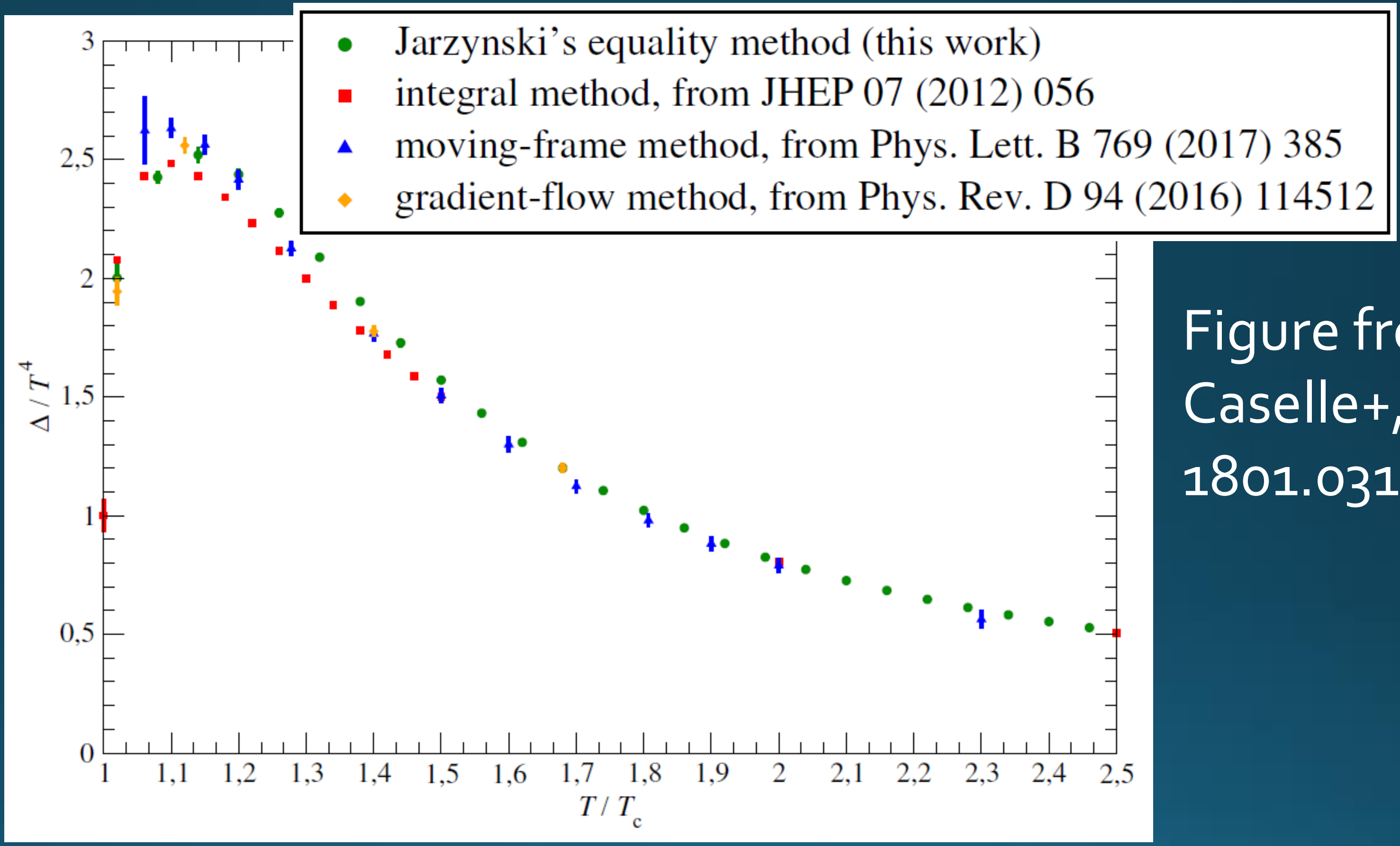


Figure from Caselle+, 1801.03110

- Measurement of thermodynamics with various methods.
- All results are in good agreement.
- But, non-negligible discrepancy at $T/T_c \approx 1\text{-}1.3$?

Gradient Flow for Fermions

$$\partial_t \psi(t,x) = D_\mu D_\mu \psi(t,x)$$

$$\partial_t \bar{\psi}(t,x) = \psi(t,x) \overleftarrow{D}_\mu \overleftarrow{D}_\mu$$

$$D_\mu = \partial_\mu + A_\mu(t,x)$$

Luscher, 2013
Makino, Suzuki, 2014
Taniguchi+ (WHOT) 2016; 2017

- Not "gradient" flow but a "diffusion" equation.

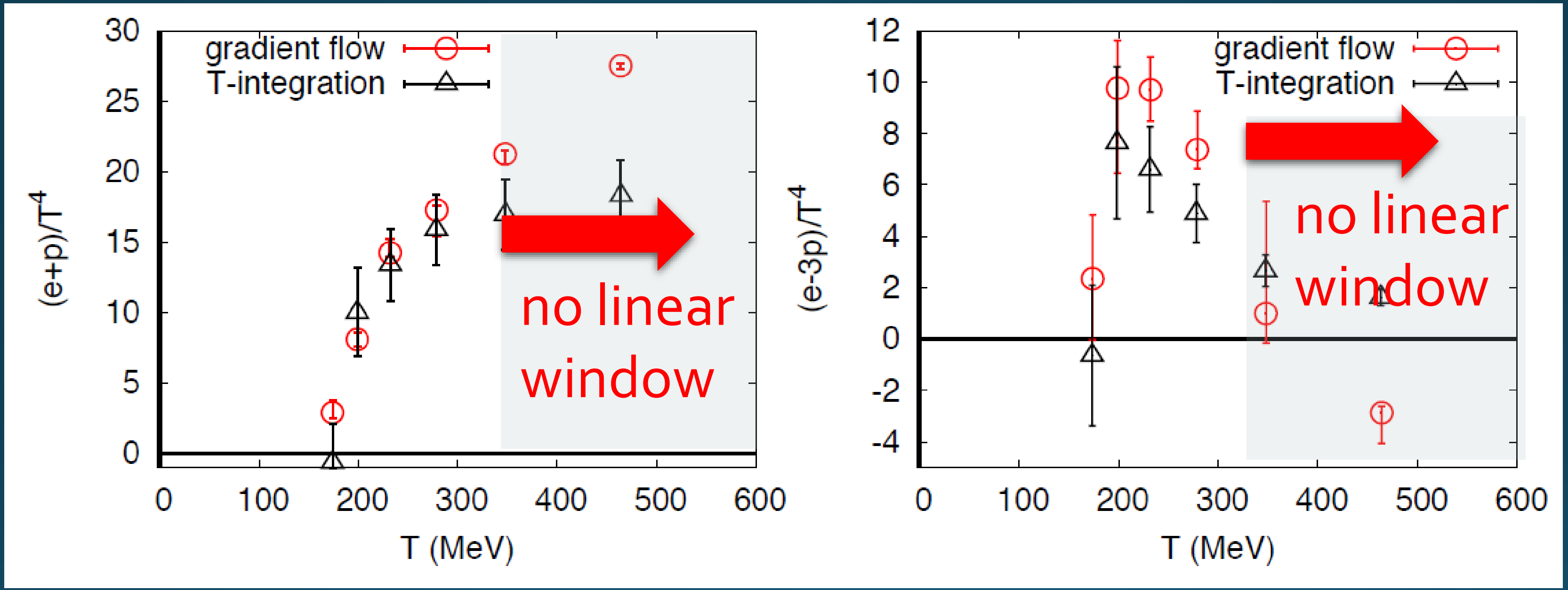
- Divergence in field renormalization of fermions.
- All observables are finite at t>0 once Z(t) is fixed.

$$\tilde{\psi}(t,x) = Z(t)\psi(t,x)$$

2+1 QCD EoS from Gradient Flow

Taniguchi+ (WHOT-QCD), PR**D96**, 014509 (2017)

$m_{PS}/m_V \approx 0.63$

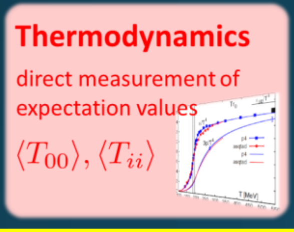


- Agreement with integral method except for N_t=4, 6
- No stable extrapolation for N_t=4, 6
- Statistical error is substantially suppressed!

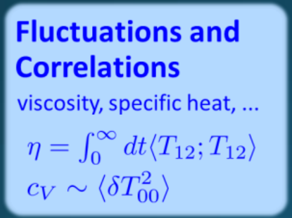
Physical mass: Kanaya+ (WHOT-QCD), 1710.10015

Contents

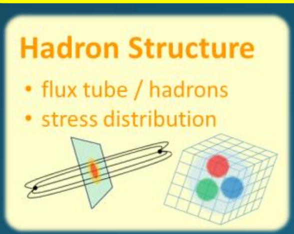
Constructing EMT on the lattice



Thermodynamics



EMT Correlation Function



Stress distribution in $\bar{q}q$ system

EMT Correlator: Motivation

- Transport Coefficient

Kubo formula → viscosity

$$\eta = \int_0^\infty dt \int_0^{1/T} d\tau \int d^3x \langle T_{12}(x, -i\tau) T_{12}(0, t) \rangle$$

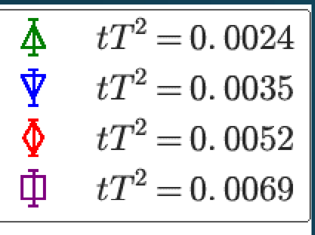
Karsch, Wyld, 1987
Nakamura, Sakai, 2005
Meyer; 2007, 2008
...
Borsanyi+, 2018
Astrakhantsev+, 2018

- Energy/Momentum Conservation

$\langle \bar{T}_{0\mu}(\tau) \bar{T}_{\rho\sigma}(0) \rangle$: τ-independent constant

- Fluctuation-Response Relations

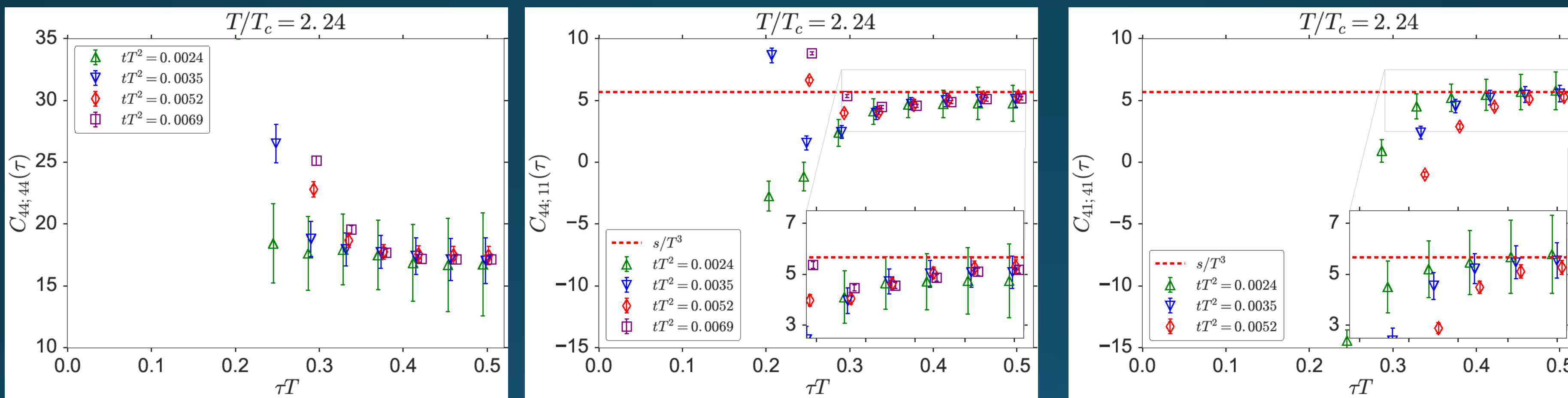
$$c_V = \frac{\langle \delta E^2 \rangle}{VT^2} \qquad E + p = \frac{\langle \bar{T}_{01}^2 \rangle}{VT} = \frac{\langle \bar{T}_{11} \bar{T}_{00} \rangle}{VT}$$



EMT Euclidean Correlator

FlowQCD, PR **D96**, 111502 (2017)

$\langle \bar{T}_{44}(\tau)\bar{T}_{44}(0)\rangle$ $\langle \bar{T}_{44}(\tau)\bar{T}_{11}(0)\rangle$ $\langle \bar{T}_{41}(\tau)\bar{T}_{41}(0)\rangle$

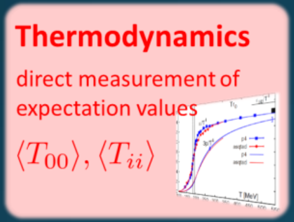


- τ-independent plateau in all channels ➔ conservation law
- Confirmation of fluctuation-response relations
- New method to measure c_V
  - Similar result for (41;41) channel: Borsanyi+, 2018
  - Perturbative analysis: Eller, Moore, 2018

Contents

$T_{\mu\nu}$

Constructing EMT on the lattice



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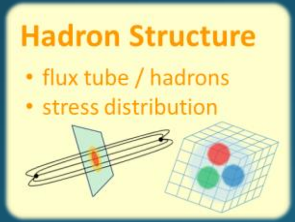
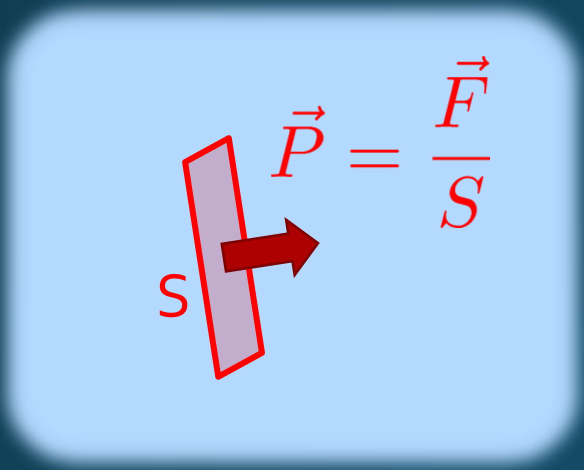


Stress distribution in $\bar{q}q$ system

Stress = Force per Unit Area

Stress = Force per Unit Area

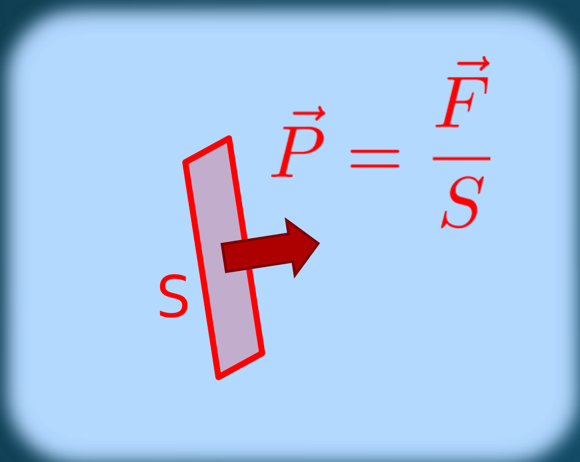
Pressure



$$\vec{P} = P\vec{n}$$

Stress = Force per Unit Area

Pressure

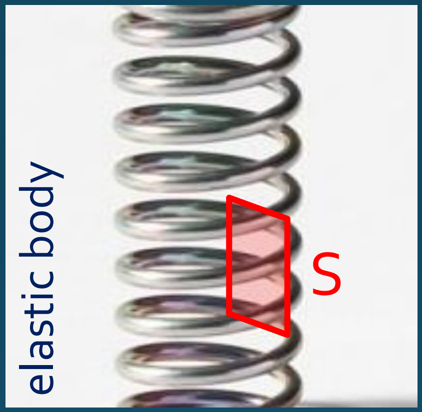


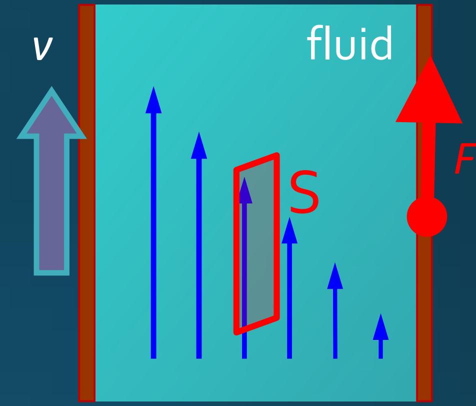
$$\vec{P} = P\vec{n}$$

In thermal medium

$$T_{ij} = P\delta_{ij}$$

Generally, F and n are not parallel





$$\frac{F_i}{S} = \sigma_{ij} n_j$$

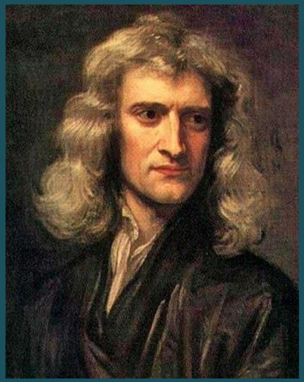
Stress Tensor

$$\sigma_{ij} = -T_{ij}$$

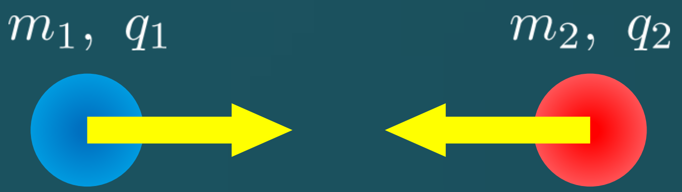
Landau
Lifshitz

Force

Action-at-a-distance



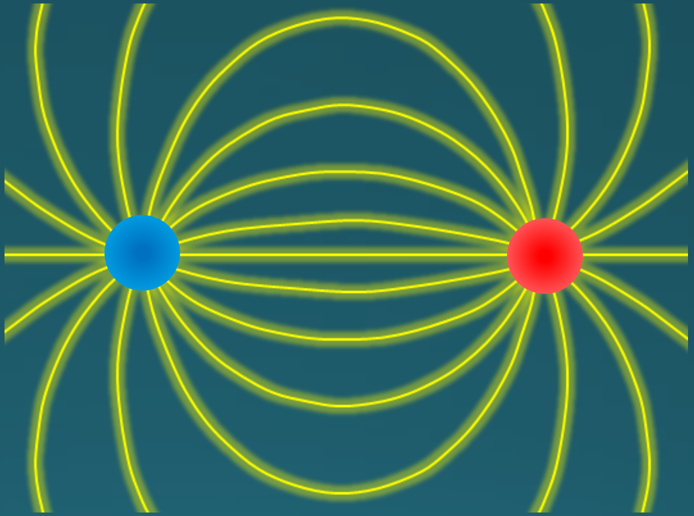
Newton
1687



$$F = -G\frac{m_1 m_2}{r^2} \qquad F = -\frac{1}{4\pi\epsilon_0}\frac{q_1 q_2}{r^2}$$

Local interaction

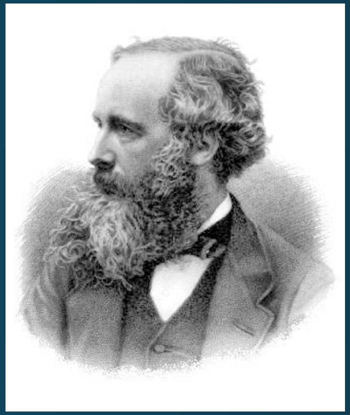
Faraday
1839



Maxwell Stress

(in Maxwell Theory)

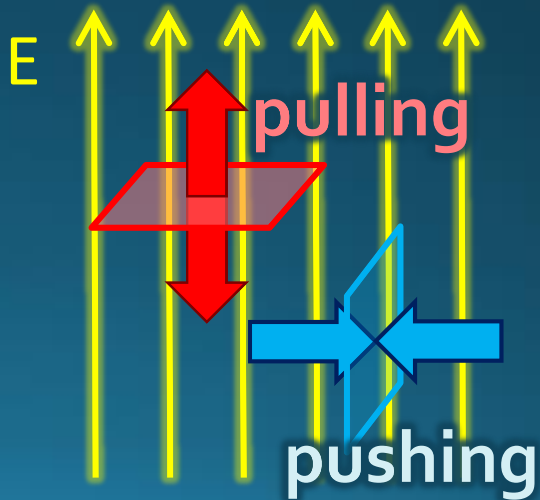
$$\sigma_{ij} = \varepsilon_0 E_i E_j + \frac{1}{\mu_0} B_i B_j - \frac{1}{2}\delta_{ij}\left(\varepsilon_0 E^2 + \frac{1}{\mu_0} B^2\right)$$



Maxwell

$$\vec{E} = (E, 0, 0)$$

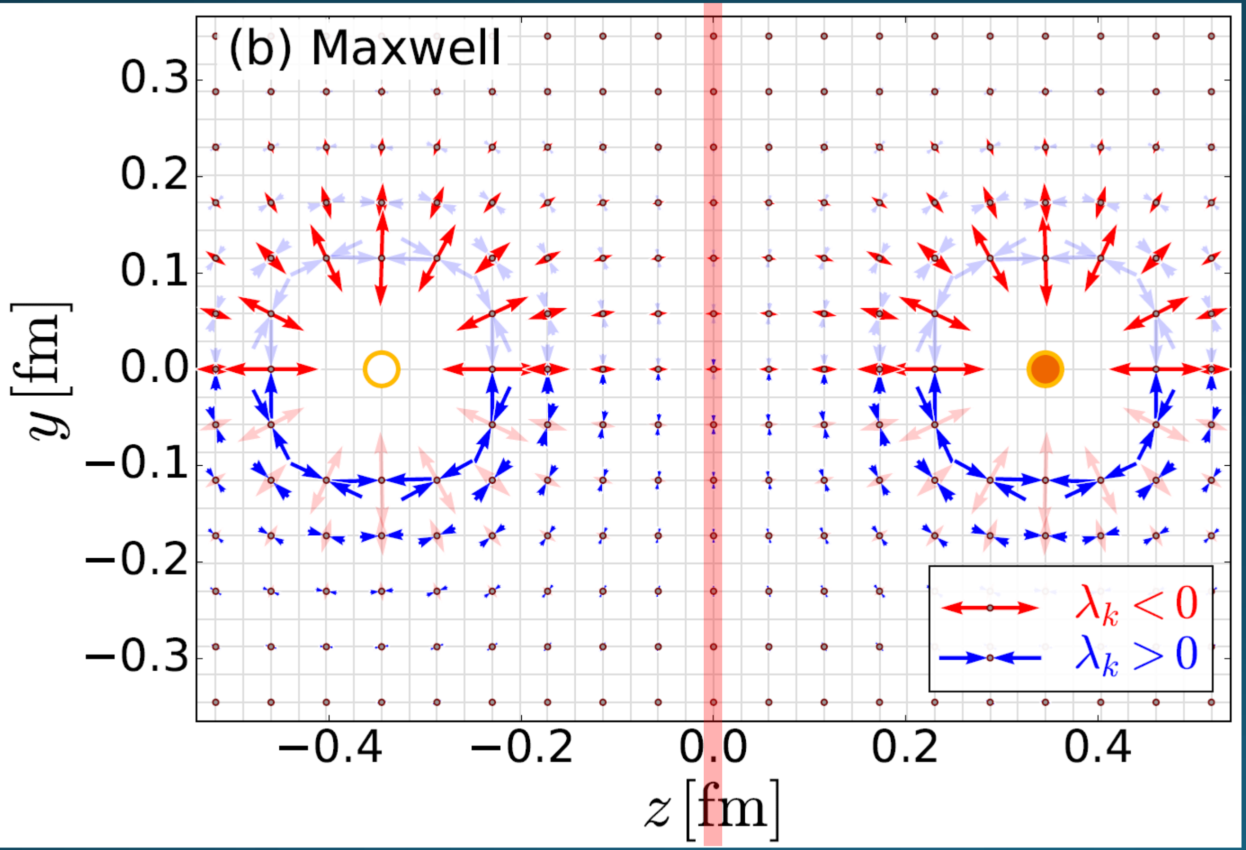
$$T_{ij} = \begin{pmatrix} -E^2 & 0 & 0 \\ 0 & E^2 & 0 \\ 0 & 0 & E^2 \end{pmatrix}$$



- Parallel to field: **Pulling**
- Vertical to field: **Pushing**

Maxwell Stress

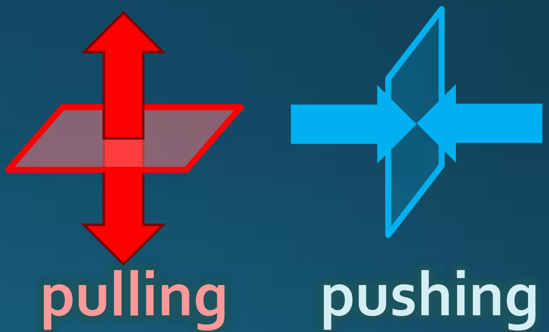
(in Maxwell Theory)



$$T_{ij} v_j^{(k)} = \lambda_k v_i^{(k)}$$

$$(k = 1, 2, 3)$$

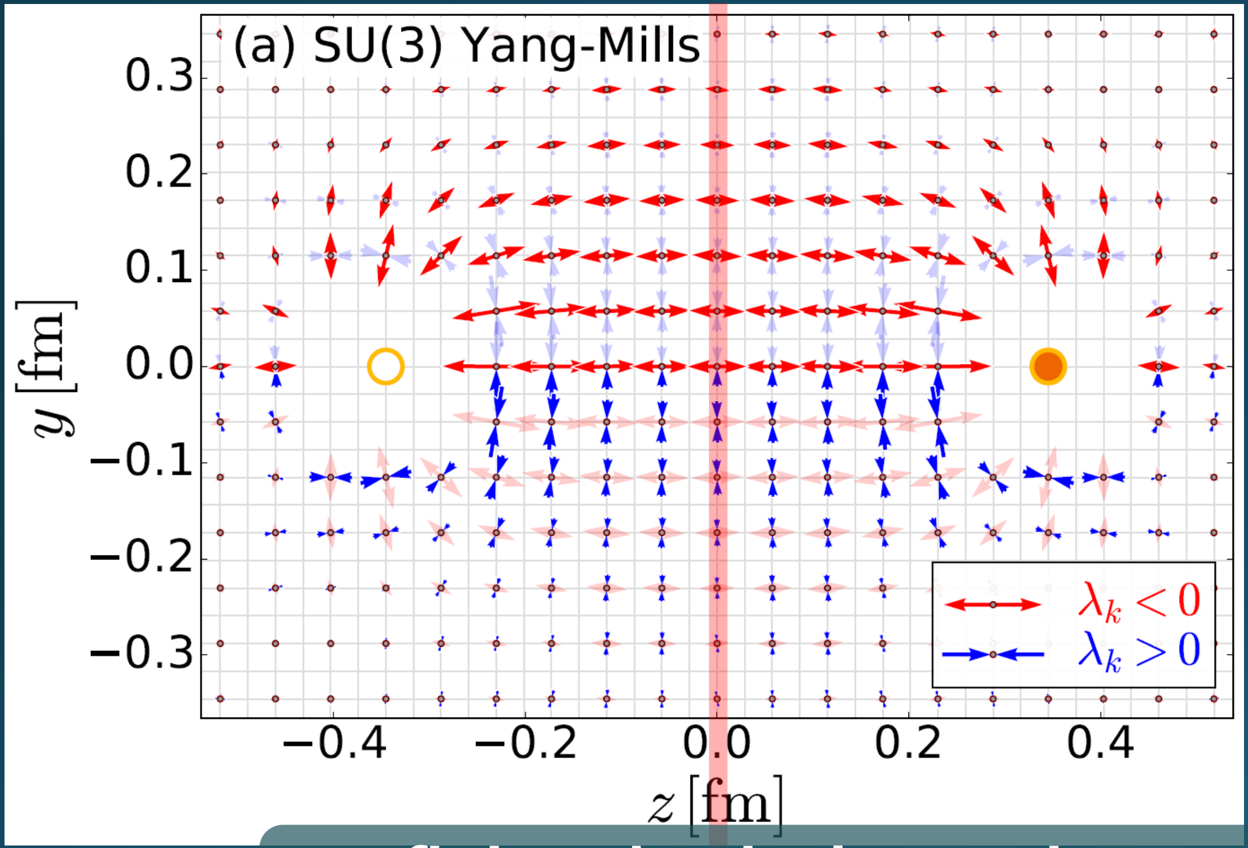
length: $\sqrt{|\lambda_k|}$



Definite physical meaning

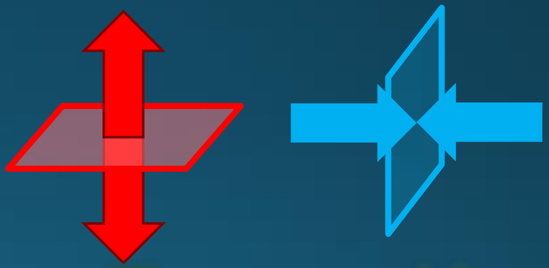
- ☐ Distortion of field, line of the force
- ☐ Propagation of the force as local interaction

Stress Tensor in $Q\bar{Q}$ System



Yanagihara+, 1803.05656

Lattice simulation
SU(3) Yang-Mills
a=0.029 fm
R=0.69 fm
t/a^2=2.0



Definite physical meaning

- ☐ Distortion of field, line of the force
- ☐ Propagation of the force as local interaction
- ☐ Manifestly gauge invariant

Lattice Setup

- ☐ SU(3) Yang-Mills (Quenched)
- ☐ Wilson gauge action
- ☐ Clover operator

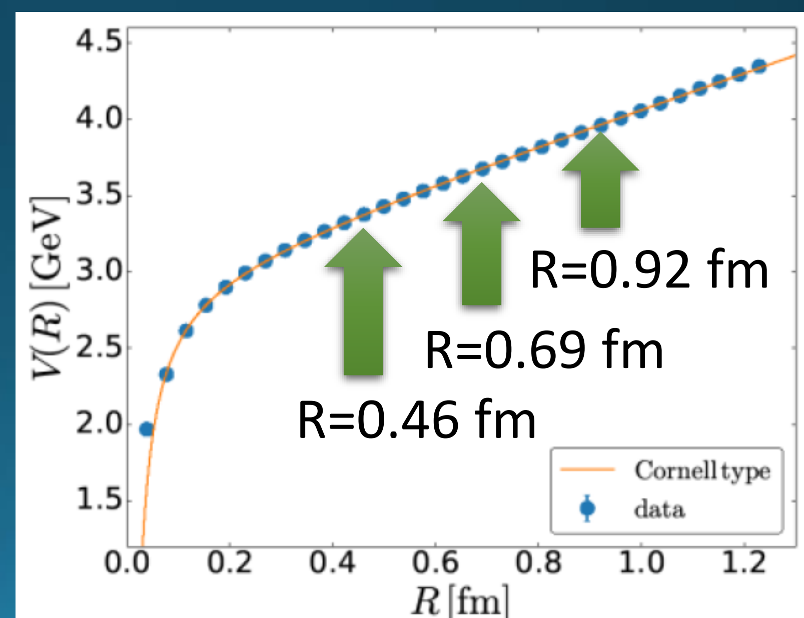
- ☐ APE smearing / multi-hit

- ☐ fine lattices (a=0.029-0.06 fm)
- ☐ continuum extrapolation

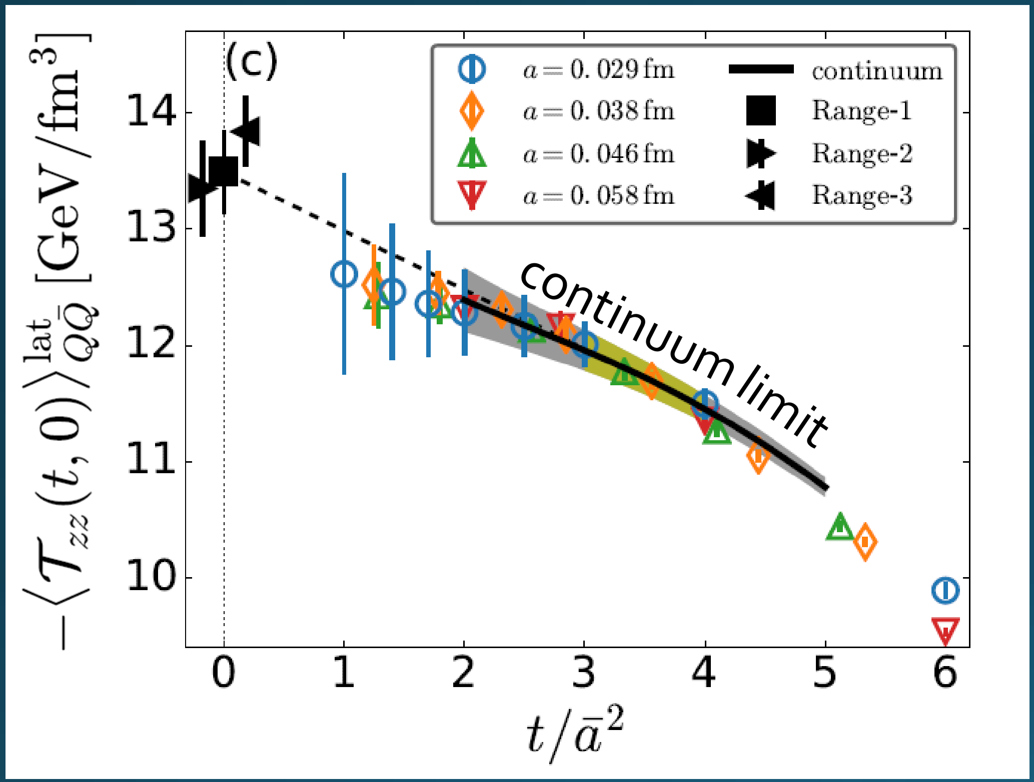
- ☐ Simulation: bluegene/Q@KEK

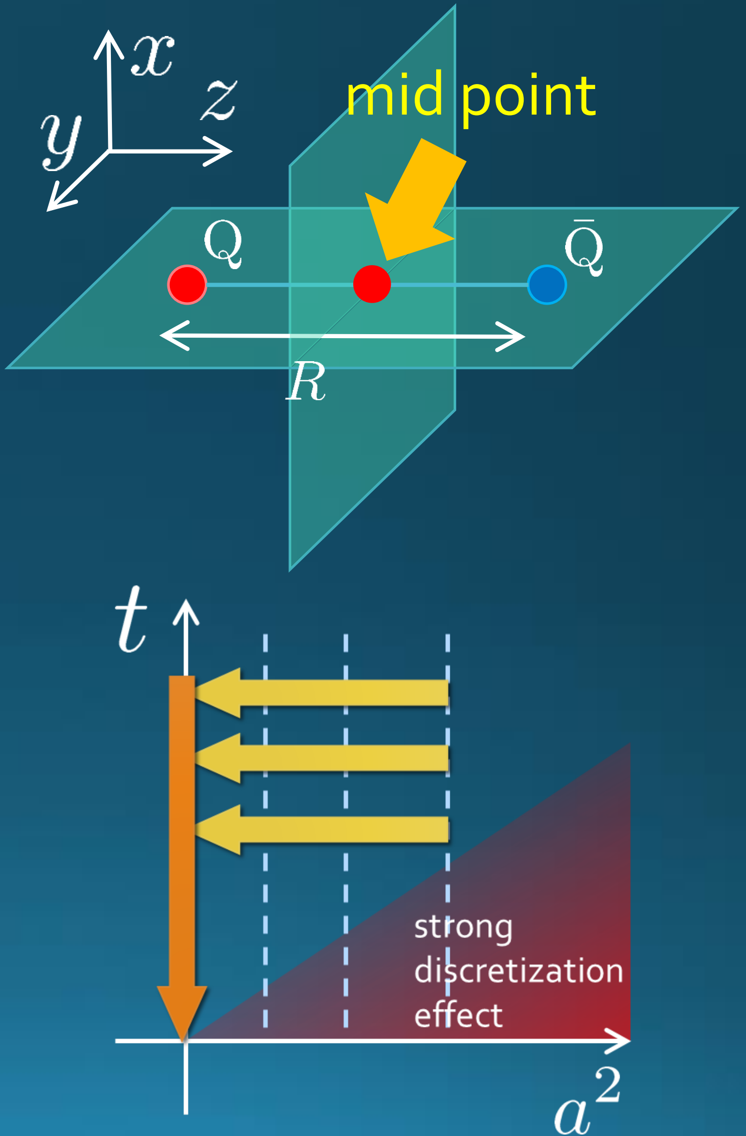
Yanagihara+, 1803.05656

β	a [fm]	$N_{\rm size}^4$	$N_{\rm conf}$	R/a		
6.304	0.058	48^4	140	8	12	16
6.465	0.046	48^4	440	10	–	20
6.513	0.043	48^4	600	–	16	–
6.600	0.038	48^4	1,500	12	18	24
6.819	0.029	64^4	1,000	16	24	32
			R [fm]	0.46	0.69	0.92



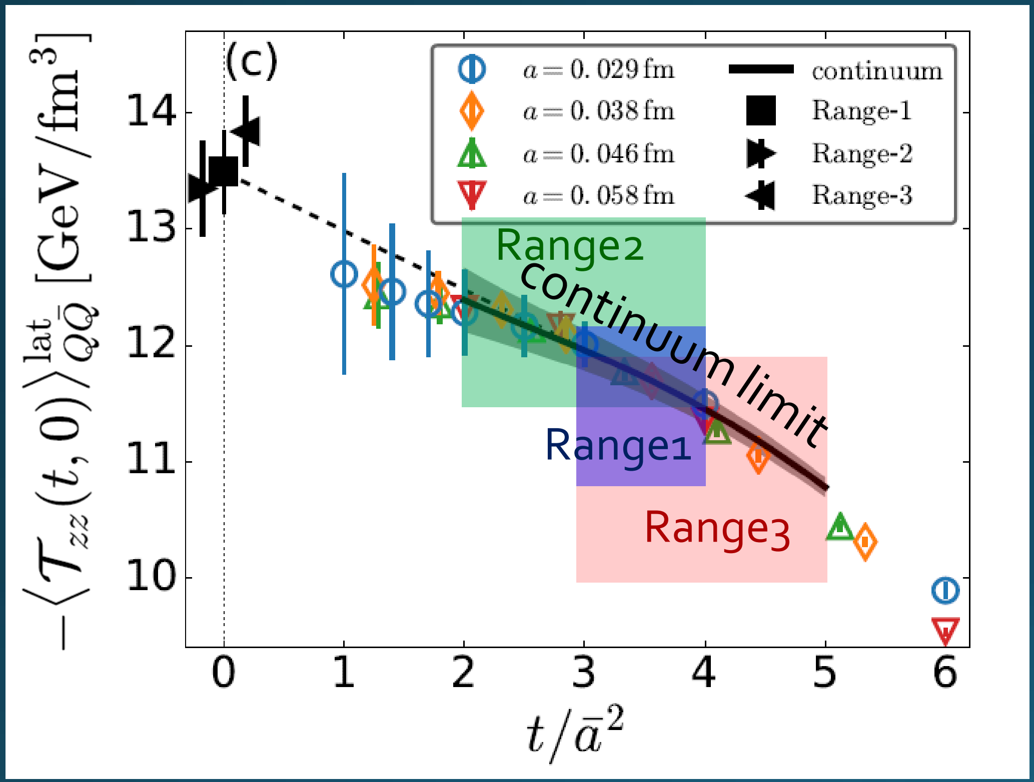
Continuum Extrapolation
at mid-point

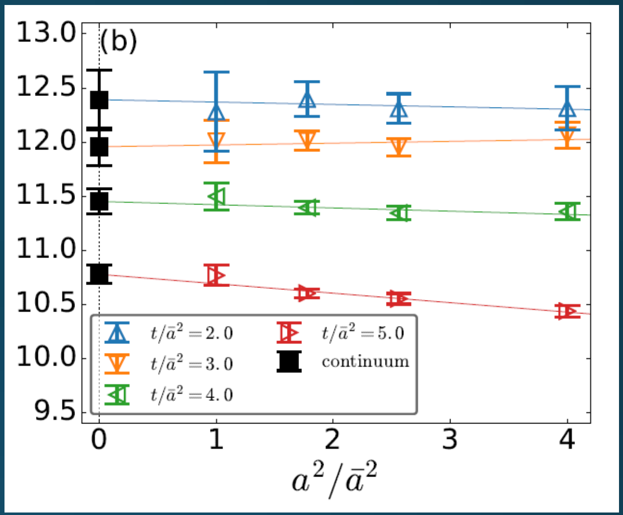




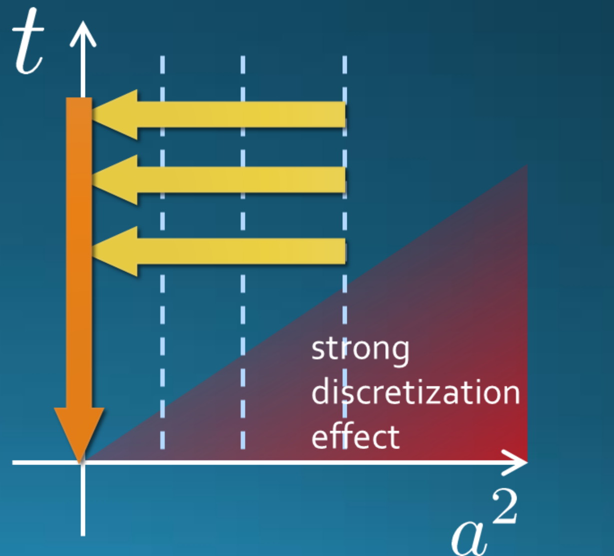
- ☐ a→0 extrapolation with fixed t

t→0 Extrapolation
at mid-point





- ☐ a→0 extrapolation with fixed t
- ☐ Then, t→0 with three ranges



Stress Distribution on Mid-Plane

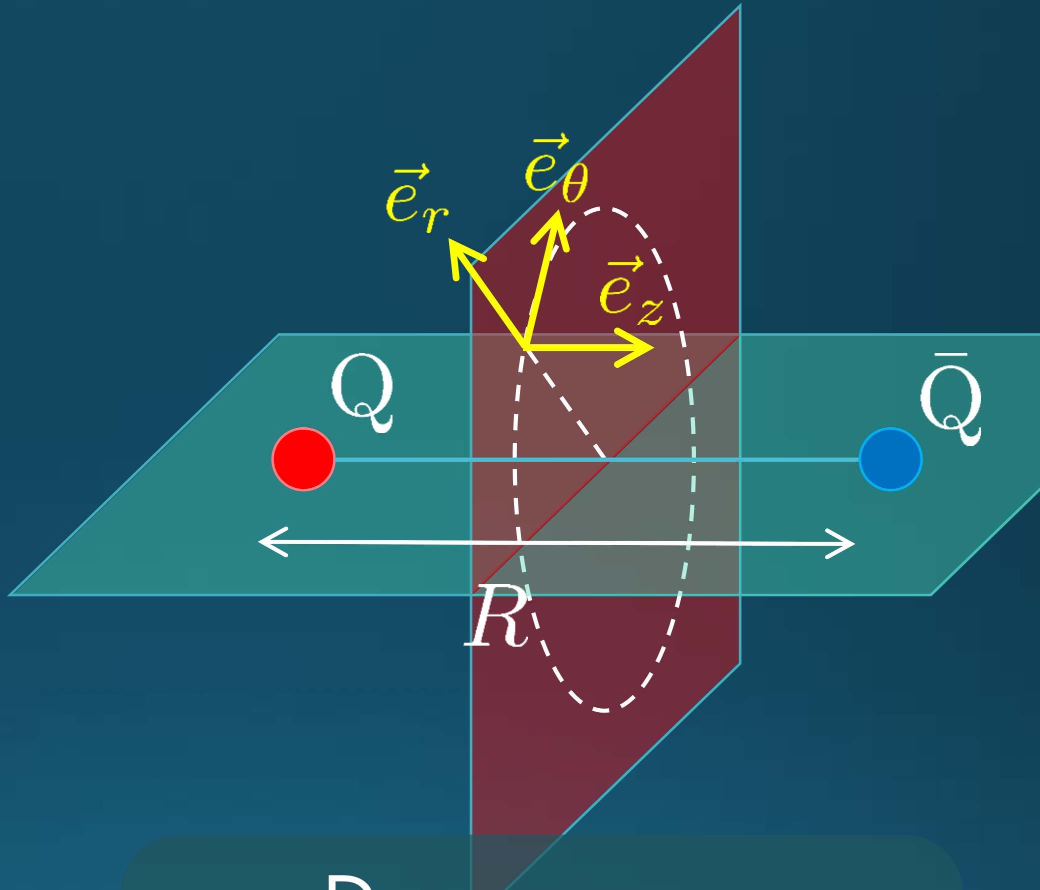
From rotational symm. & parity

EMT is diagonalized
in Cylindrical Coordinates

$$T_{cc'}(r) = \begin{pmatrix} T_{rr} & & & \\ & T_{\theta\theta} & & \\ & & T_{zz} & \\ & & & T_{44} \end{pmatrix}$$

$$T_{rr} = \vec{e}_r^{\,T} T \vec{e}_r$$

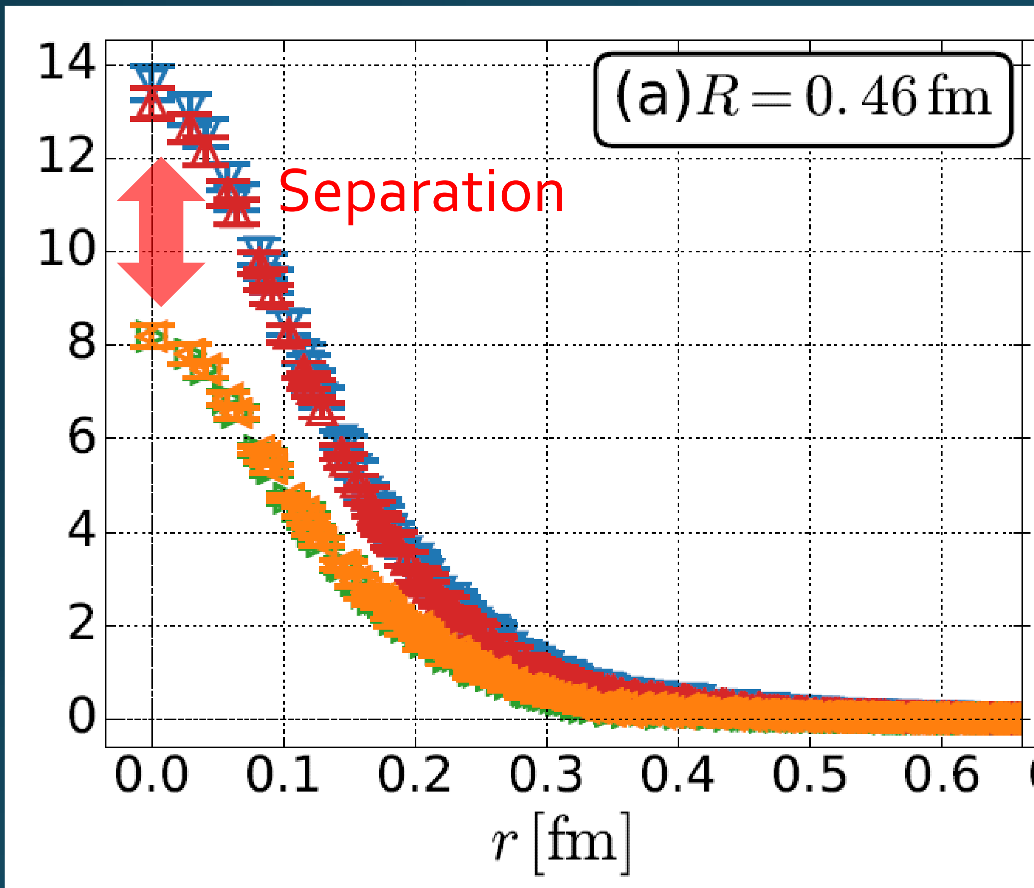
$$T_{\theta\theta} = \vec{e}_\theta^{\,T} T \vec{e}_\theta$$

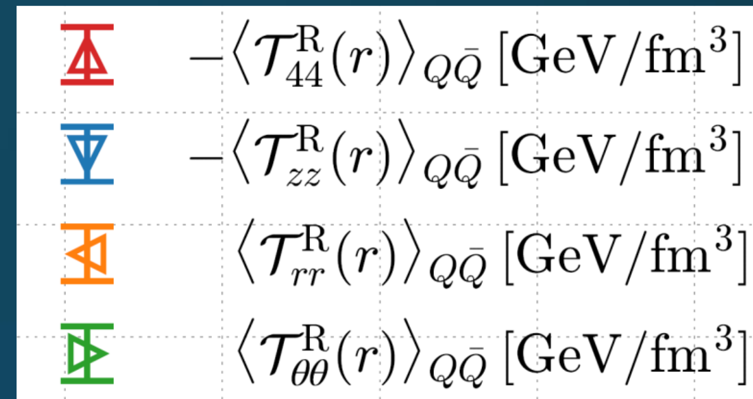


Degeneracy
in Maxwell theory

$$T_{rr} = T_{\theta\theta} = -T_{zz} = -T_{44}$$

Mid-Plane





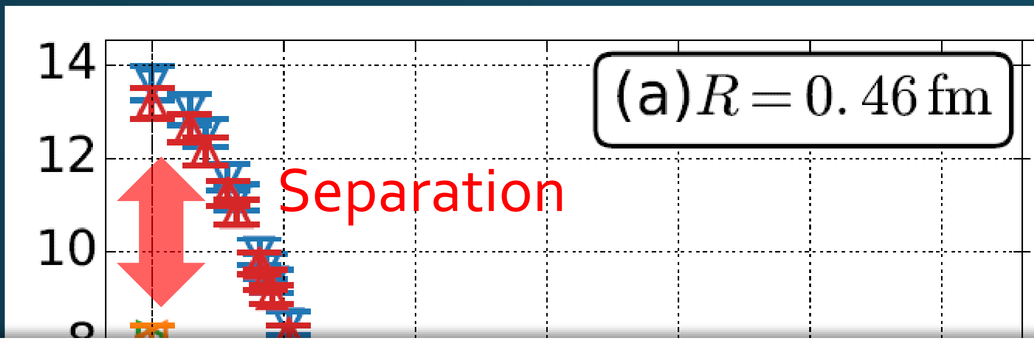
Continuum Extrapolated!

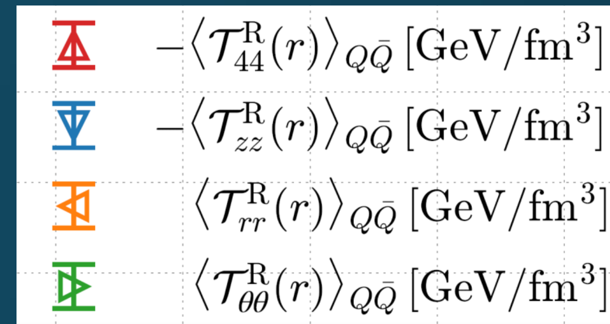
In Maxwell theory

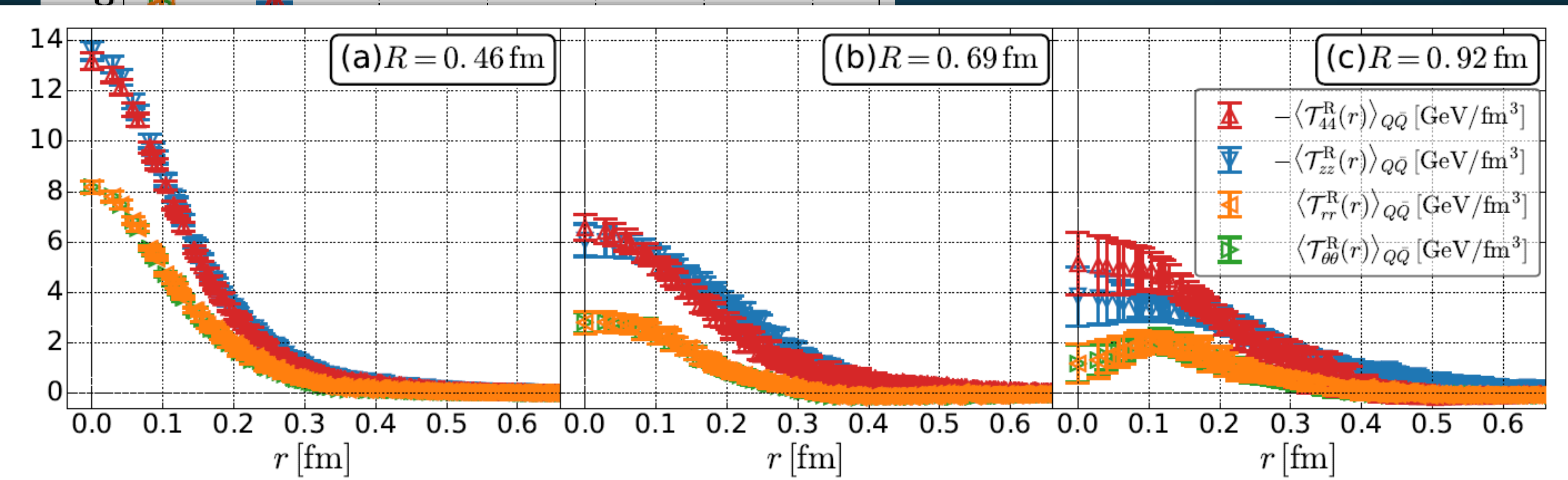
$$T_{rr} = T_{\theta\theta} = -T_{zz} = -T_{44}$$

- Degeneracy: $T_{44} \simeq T_{zz}, \quad T_{rr} \simeq T_{\theta\theta}$
- Separation: $T_{zz} \neq T_{rr}$
- Nonzero trace anomaly $\sum T_{cc} \neq 0$

Mid-Plane







- Degeneracy: $T_{44} \simeq T_{zz}, \quad T_{rr} \simeq T_{\theta\theta}$
- Separation: $T_{zz} \neq T_{rr}$
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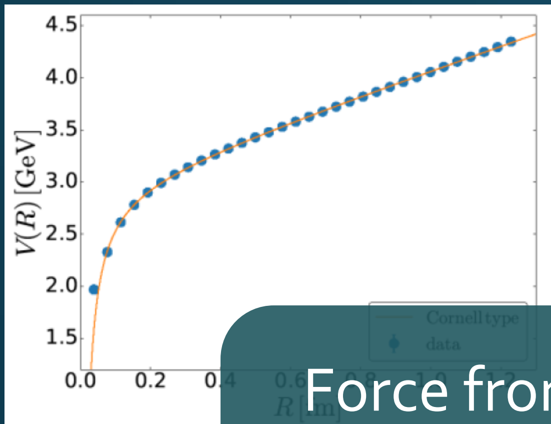
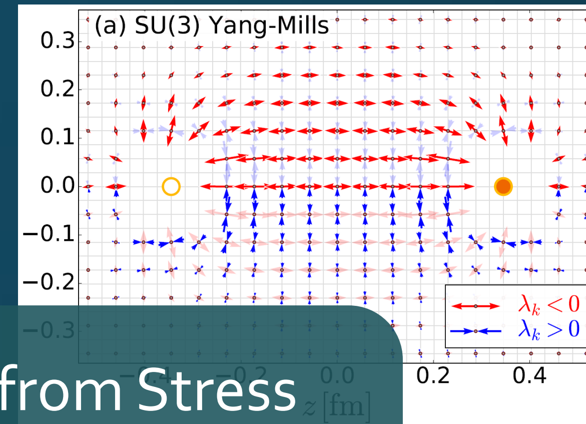
Force





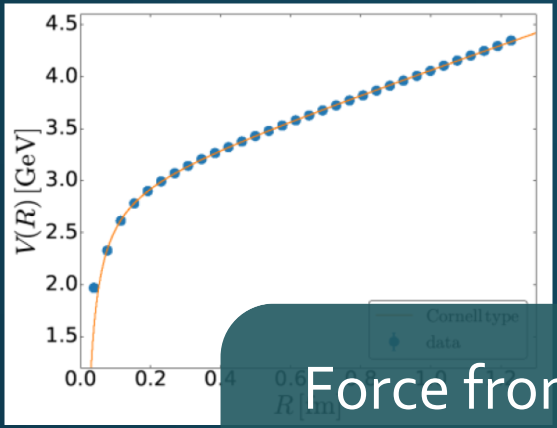
Force from Potential

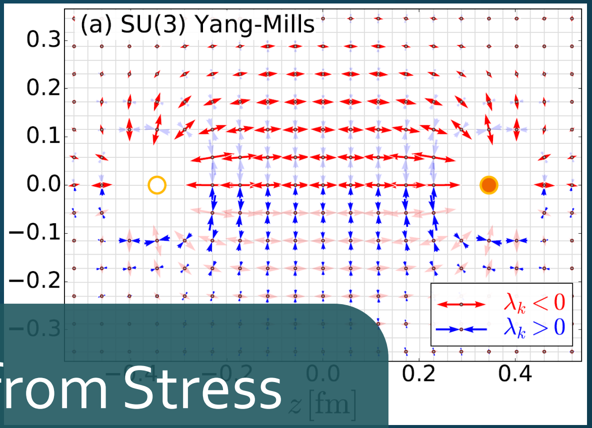
$$F_{\text{pot}} = -\frac{dV}{dR}$$

Force from Stress

$$F_{\text{stress}} = \int_{\text{mid.}} d^2x T_{zz}(x)$$

Force





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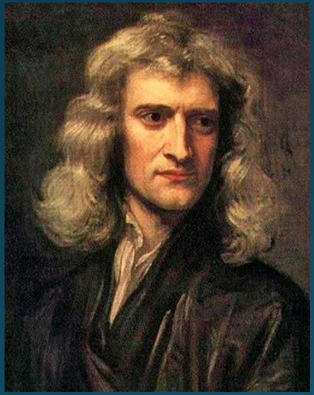
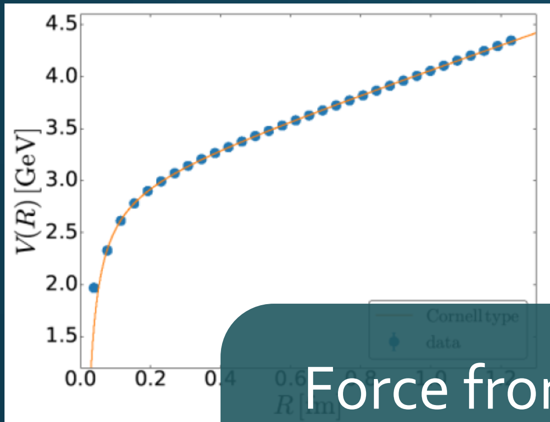
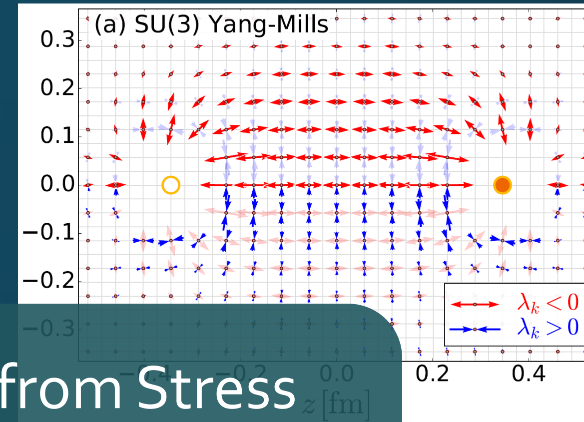


Newton
1687

Faraday
1839

Force





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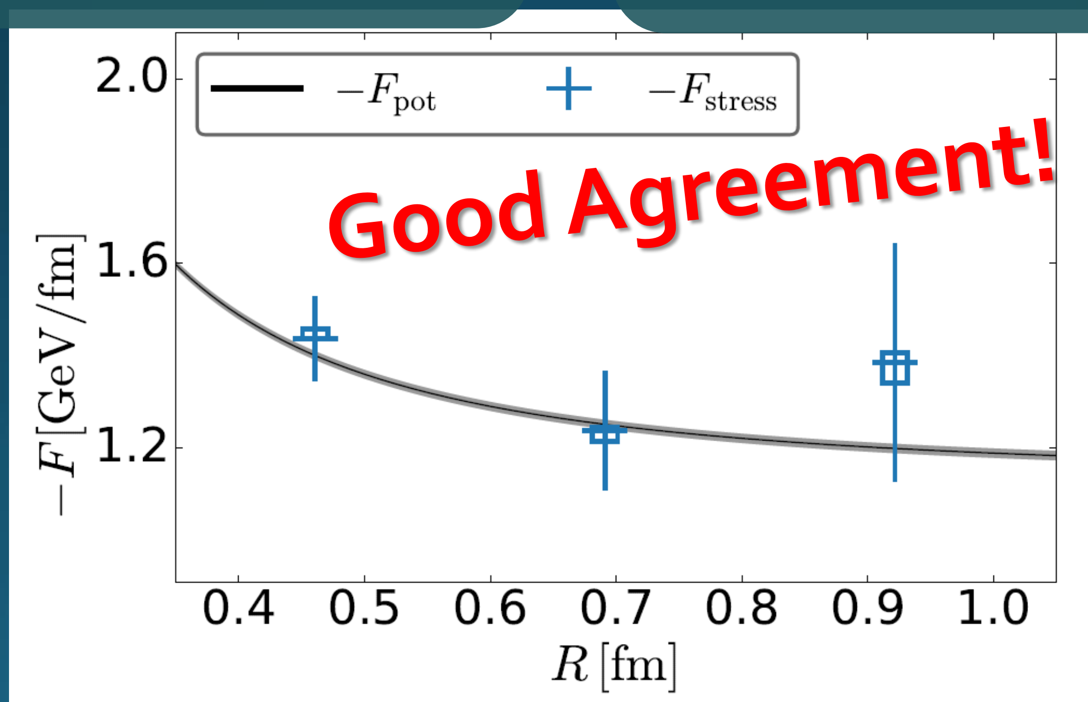
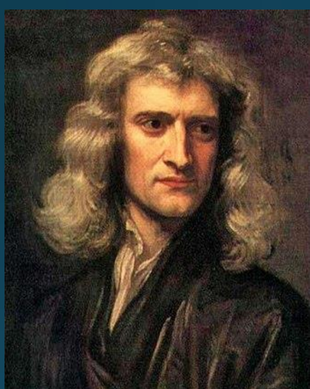




Newton
1687

Faraday
1839

Abelian-Higgs Model

Yanagihara, Iritani, MK, in prep.

Abelian-Higgs Model

$$\mathcal{L}_{\mathrm{AH}} = -\frac{1}{4}F_{\mu\nu}^2 + |(\partial_\mu + igA_\mu)\phi|^2 - \lambda(\phi^2 - v^2)^2$$

GL parameter: $\kappa = \sqrt{\lambda}/g$

- type-I : $\kappa < 1/\sqrt{2}$
- type-II : $\kappa > 1/\sqrt{2}$
- Bogomol'nyi bound : $\kappa = 1/\sqrt{2}$

Infinitely long tube

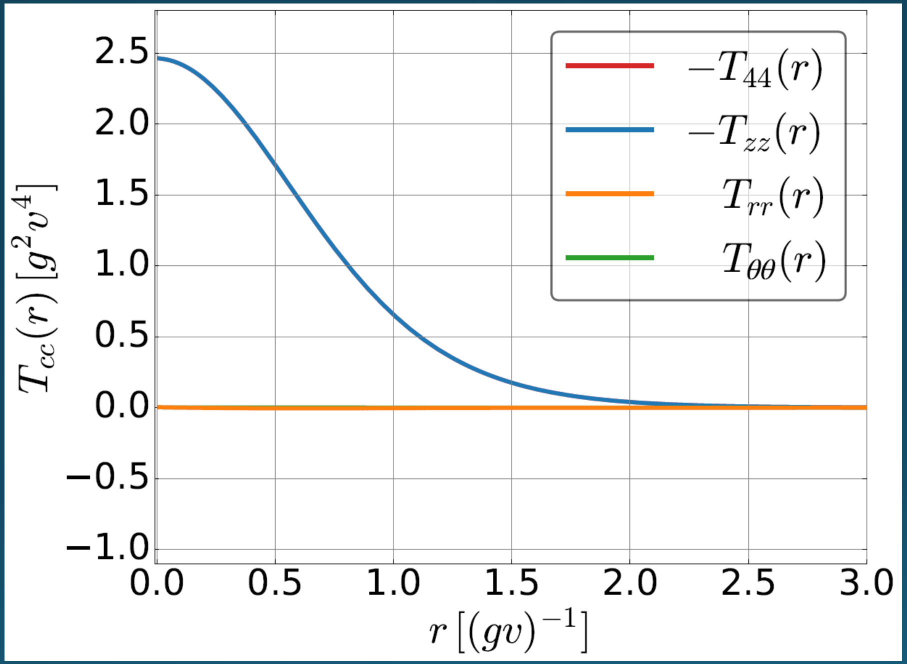
- degeneracy

  $T_{zz}(r) = T_{44}(r)$ Luscher, 1981
- momentum conservation

  $$\frac{d}{dr}(rT_{rr}) = T_{\theta\theta}$$

Stress Tensor in AH Model
infinitely-long flux tube

Bogomol'nyi bound : $\kappa = 1/\sqrt{2}$



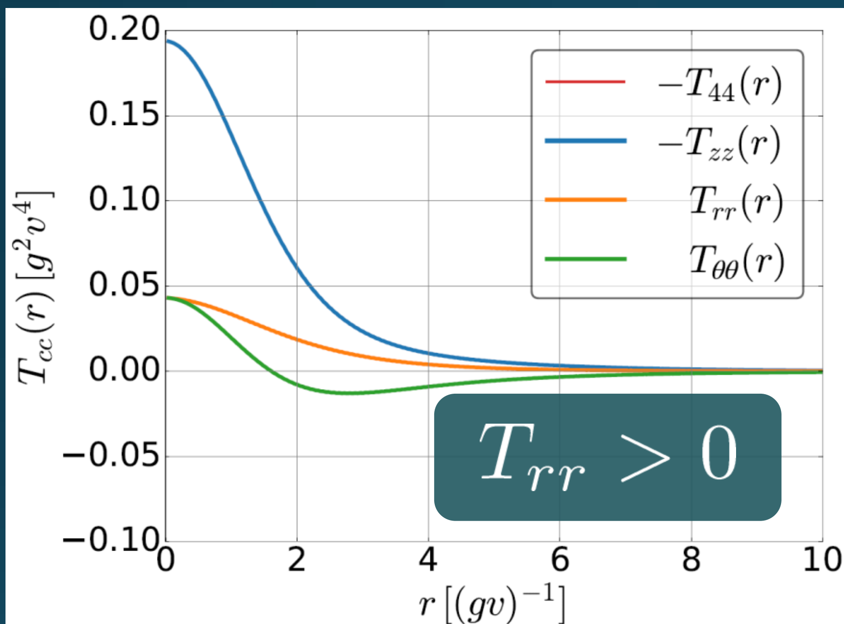
$$T_{rr} = T_{\theta\theta} = 0$$

de Vega, Schaposnik, PRD**14**, 1100 (1976).

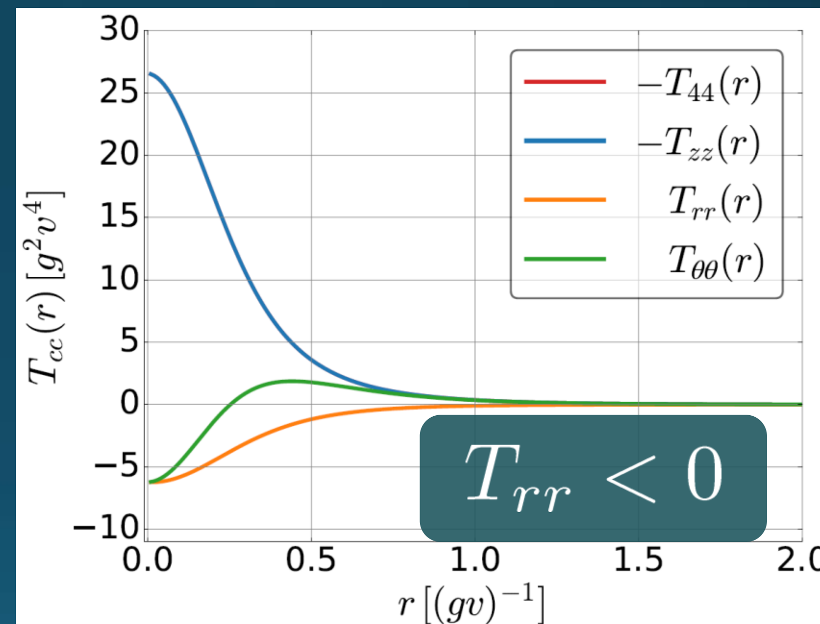
Stress Tensor in AH Model

infinitely-long flux tube

Type-I $\kappa = 0.1$



Type-II $\kappa = 3.0$



- ☐ No degeneracy bw T_{rr} & $T_{\theta\theta}$
- ☐ $T_{\theta\theta}$ changes sign

conservation law

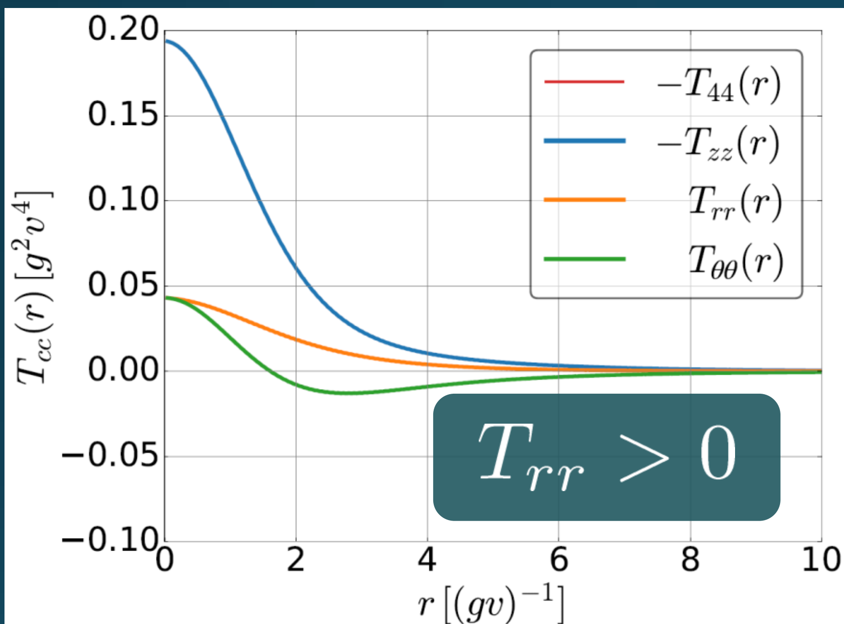
$$\frac{d}{dr}(rT_{rr}) = T_{\theta\theta}$$

Stress Tensor in AH Model

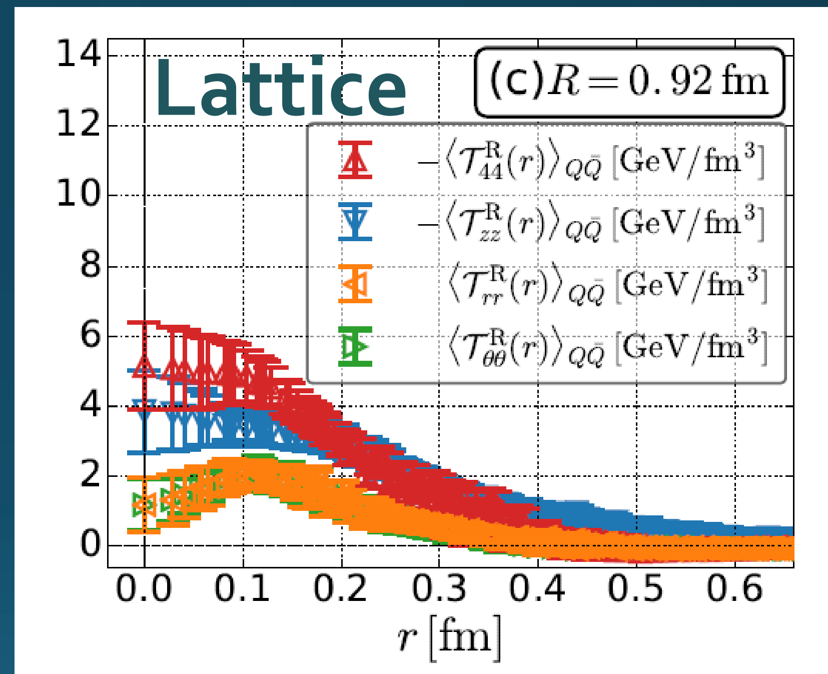
infinitely-long flux tube

Type-I

$\kappa = 0.1$



$T_{rr} > 0$



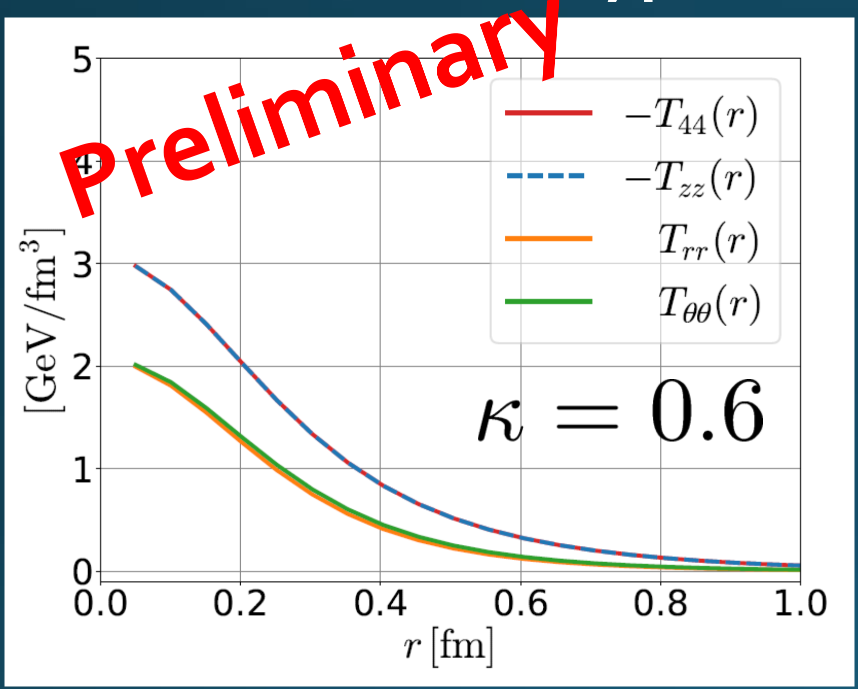
- ☐ No degeneracy bw T_{rr} & $T_{\theta\theta}$
- ☐ $T_{\theta\theta}$ changes sign

Inconsistent with lattice result

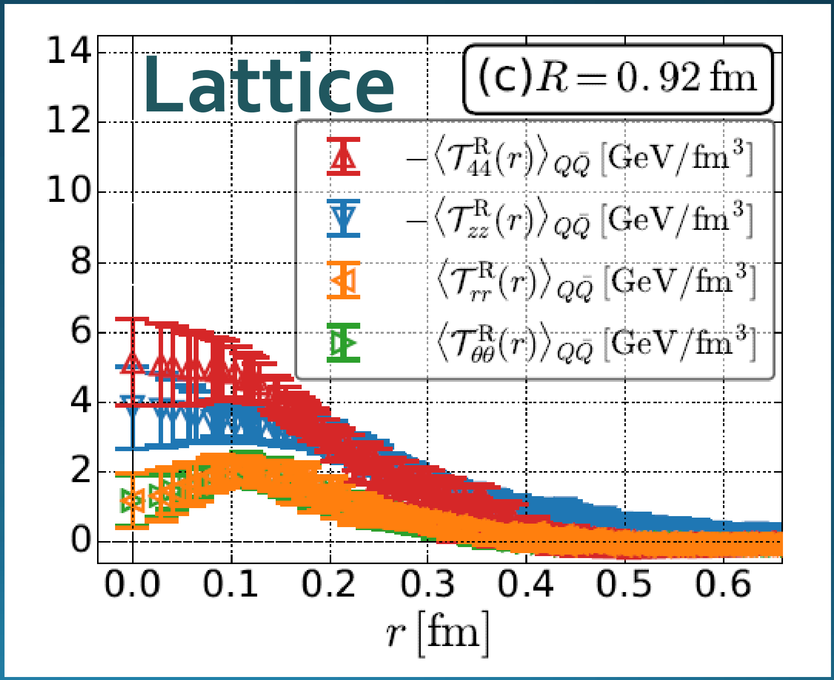
$$T_{rr} \simeq T_{\theta\theta}$$

Flux Tube with Finite Length

Finite R, weak Type-I



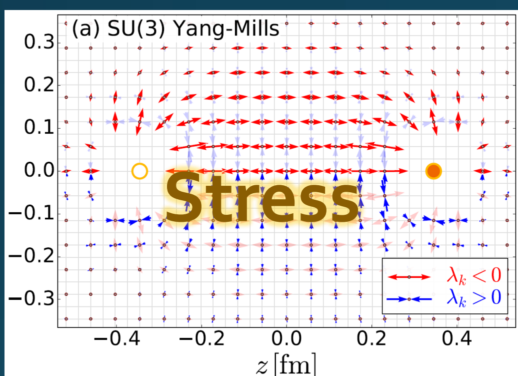
$R = 0.91\mathrm{fm}$
$\xi_A = 0.71\mathrm{fm}$
$\xi_\phi = 0.84\mathrm{fm}$

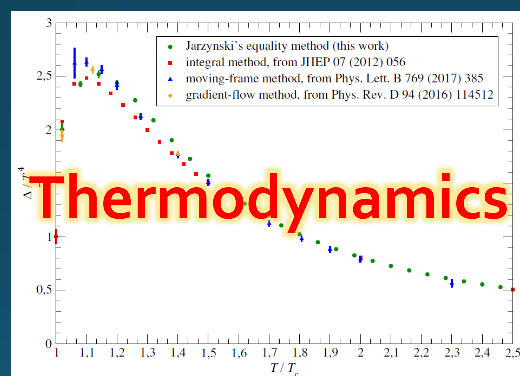


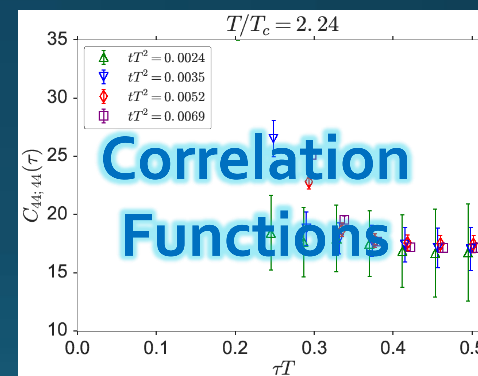
- Finite-length effect of the flux tube is crucial!

Summary

- The analysis of energy-momentum tensor on the lattice is now available, and various stuides are ongoing!
  - gradient flow method
  - determination of Z_6, Z_3, Z_1 / multilevel algorithm

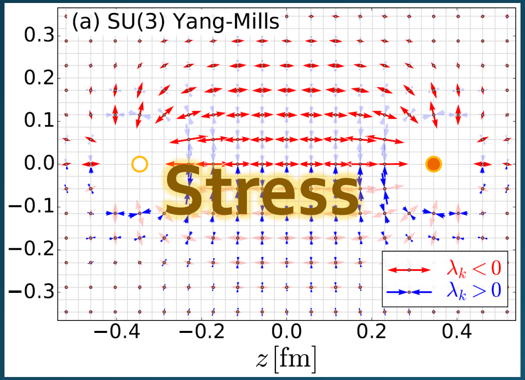


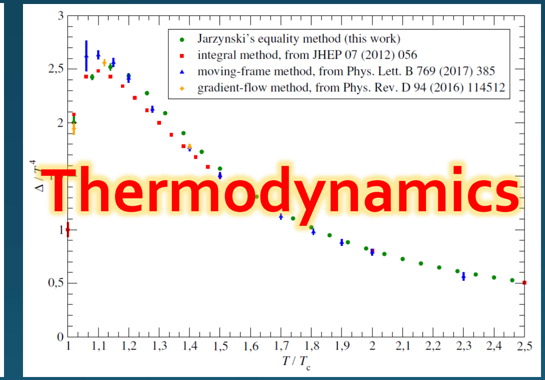


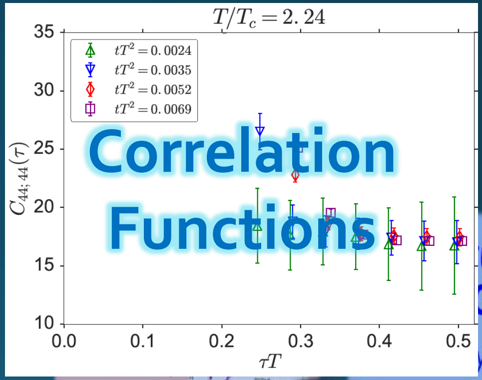


Summary

- The analysis of energy-momentum tensor on the lattice is now available, and various stuides are ongoing!
  - gradient flow method
  - determination of Z_6, Z_3, Z_1 / multilevel algorithm

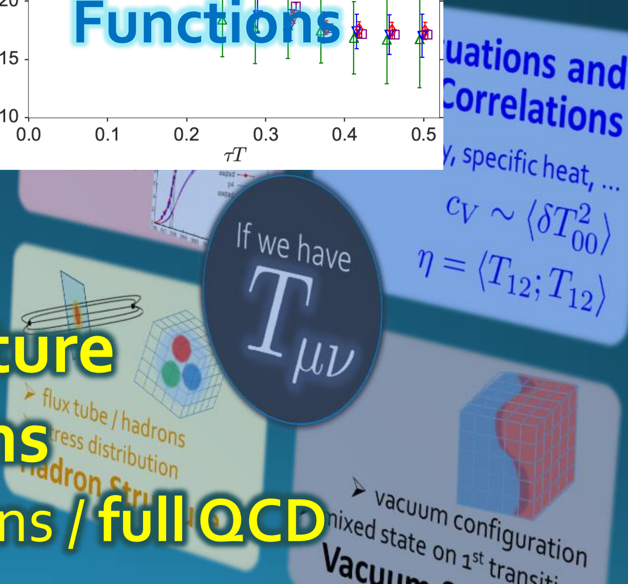






So many future studies

- Flux tube at **nonzero temperature**
- EMT distribution **inside hadrons**
- viscosity / other operators / instantons / **full QCD**



backup

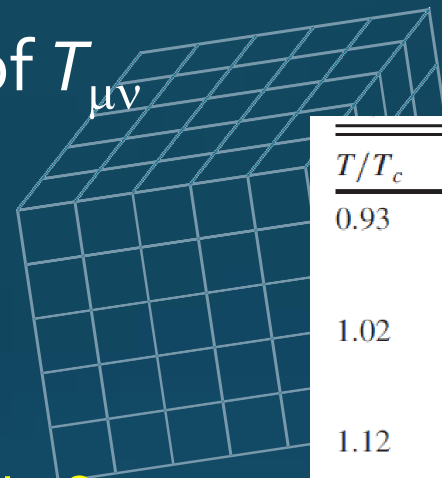
Numerical Simulation

FlowQCD,
PRD**94**, 114512 (2016)

- Expectation values of $T_{\mu\nu}$
- SU(3) YM theory
- Wilson gauge action
- Parameters:
  - N_t = 12, 16, 20-24
  - aspect ratio $5.3<N_s/N_t<8$
  - 1500~2000 configurations
- Scale from gradient flow

$$\rightarrow aT_c \text{ and } a\Lambda_{\mathrm{MS}}$$

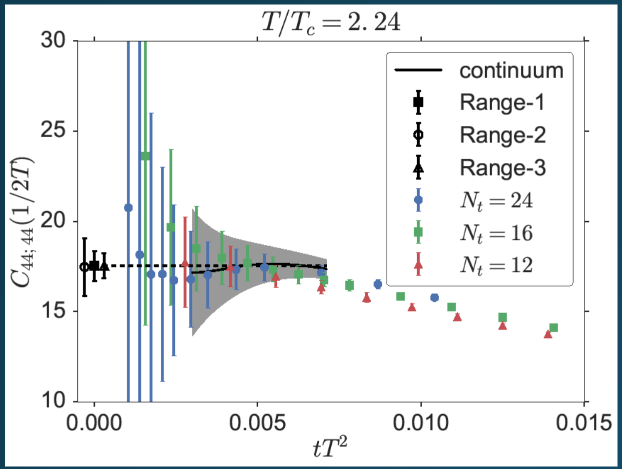
FlowQCD, 1503.06516

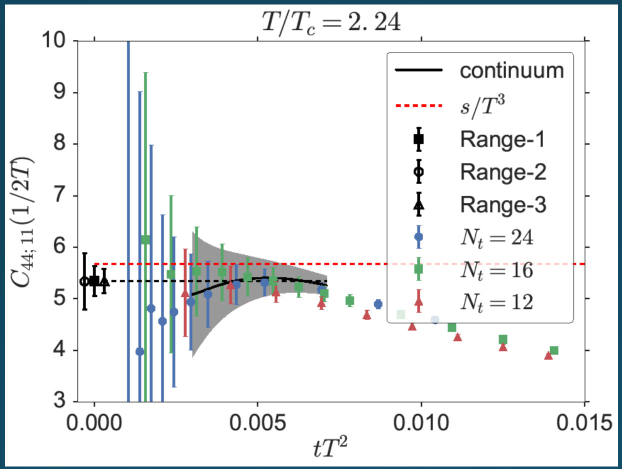


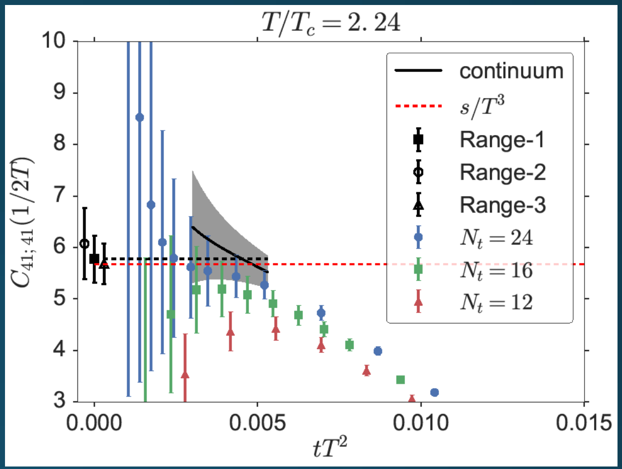
T/T_c	β	N_s	N_τ	Configurations
0.93	6.287	64	12	2125
	6.495	96	16	1645
	6.800	128	24	2040
1.02	6.349	64	12	2000
	6.559	96	16	1600
	6.800	128	22	2290
1.12	6.418	64	12	1875
	6.631	96	16	1580
	6.800	128	20	2000
1.40	6.582	64	12	2080
	6.800	128	16	900
	7.117	128	24	2000
1.68	6.719	64	12	2000
	6.941	96	16	1680
	7.117	128	20	2000
2.10	6.891	64	12	2250
	7.117	128	16	840
	7.296	128	20	2040
2.31	7.200	96	16	1490
	7.376	128	20	2020
	7.519	128	24	1970
2.69	7.086	64	12	2000
	7.317	96	16	1560
	7.500	128	20	2040

Mid-Point Correlator

$\langle T_{44}(\tau_{\text{mid}}) T_{44}(0) \rangle$ $\langle T_{44}(\tau_{\text{mid}}) T_{11}(0) \rangle$ $\langle T_{41}(\tau_{\text{mid}}) T_{41}(0) \rangle$







- (44;11), (41;41) channels : confirmation of FRR
- (44;44) channel: new measurement of c_V

$$c_V = \frac{\langle \delta E^2 \rangle}{VT^2}$$

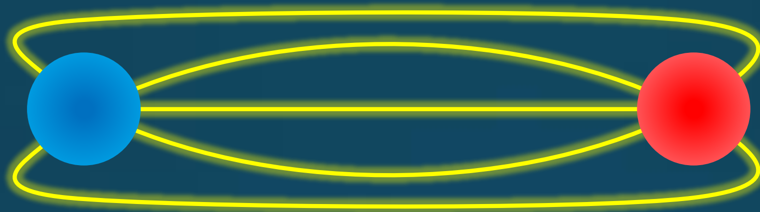
c_V/T^3

T/T_c	$C_{44;44}(\tau_m)$	Ref.[19]	Ref.[11]	ideal gas
1.68	$17.7(8)(^{+2.1}_{-0.4})$	$22.8(7)^*$	17.7	21.06
2.24	$17.5(0.8)(^{+0}_{-0.1})$	$17.9(7)^{**}$	18.2	21.06

2+1 QCD:
Taniguchi+ (WHOT-QCD), 1711.02262

Quark—Anti-quark system

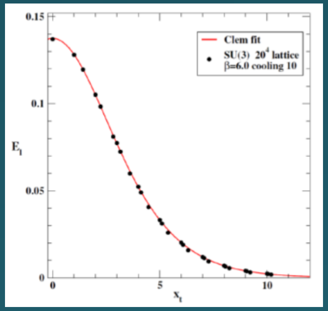
Formation of the flux tube → confinement



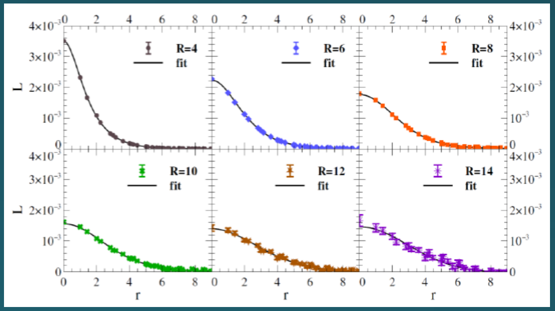
Previous Studies on Flux Tube

- ☐ Potential
- ☐ Action density
- ☐ Color-electric field

so many studies…



Cea+ (2012)

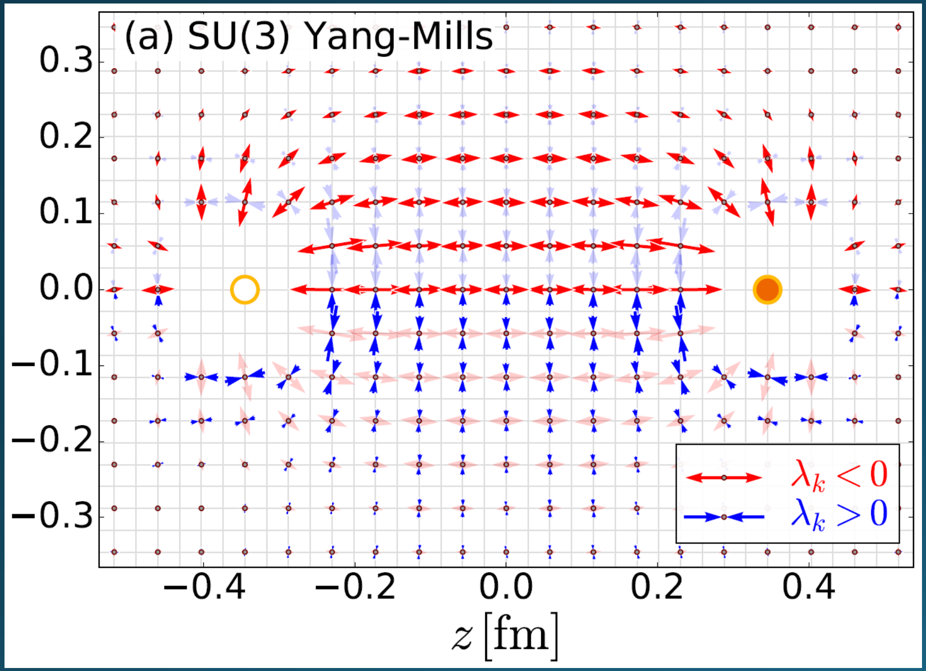


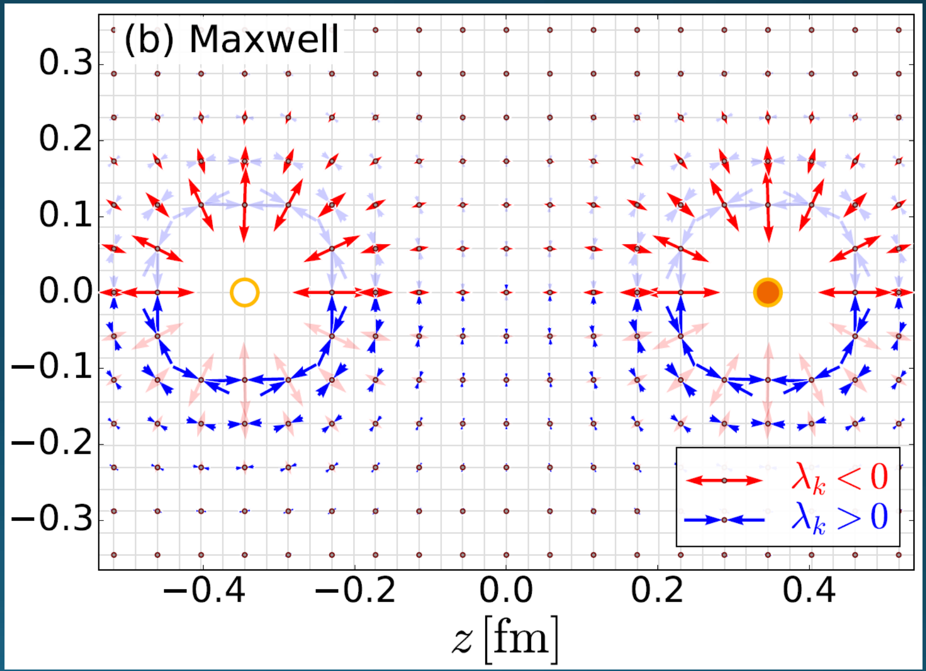
Cardoso+ (2013)

SU(3) YM vs Maxwell

SU(3) Yang-Mills
(quantum)

Maxwell
(classical)



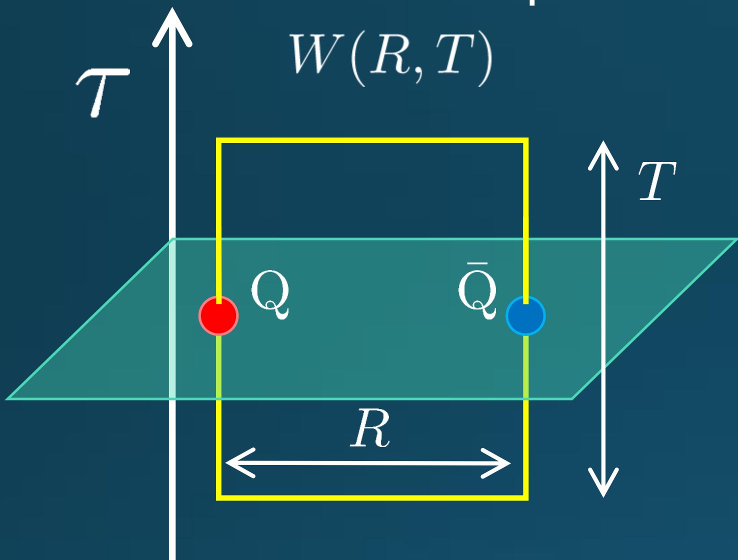


Propagation of the force is clearly different in YM and Maxwell theories!

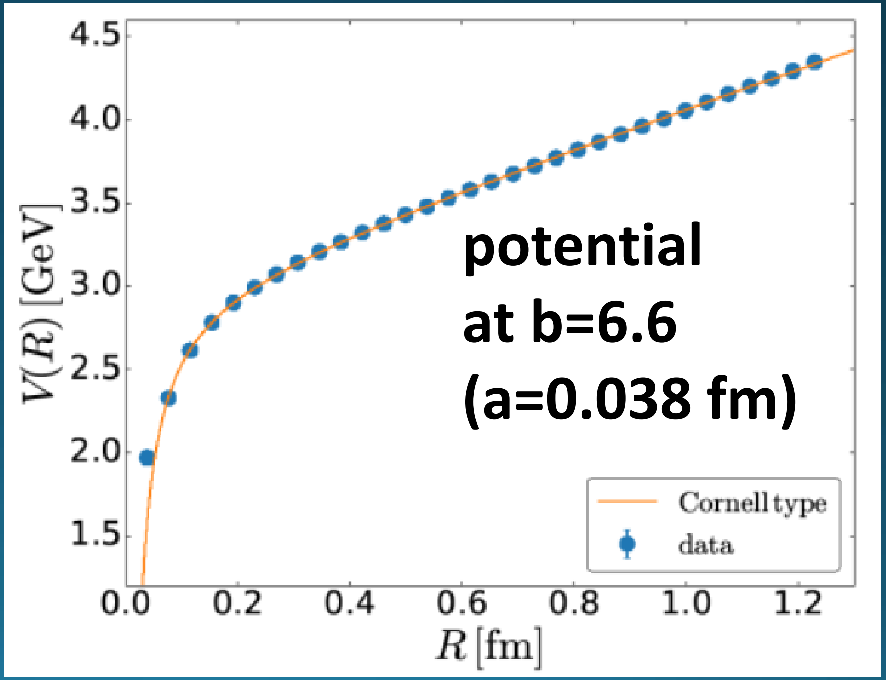
Preparing Static $Q\bar{Q}$

Wilson loop

$W(R,T)$



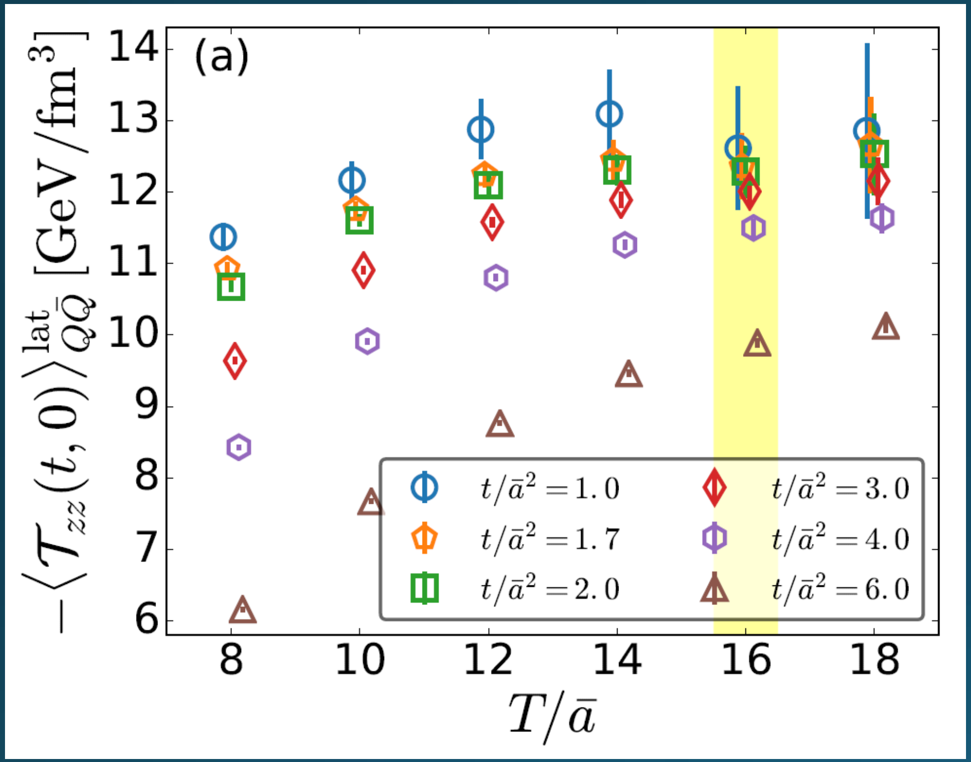
- ☐ APE smearing for spatial links
- ☐ Multi-hit for temporal links
- ☐ No gradient flow for W(R,T)

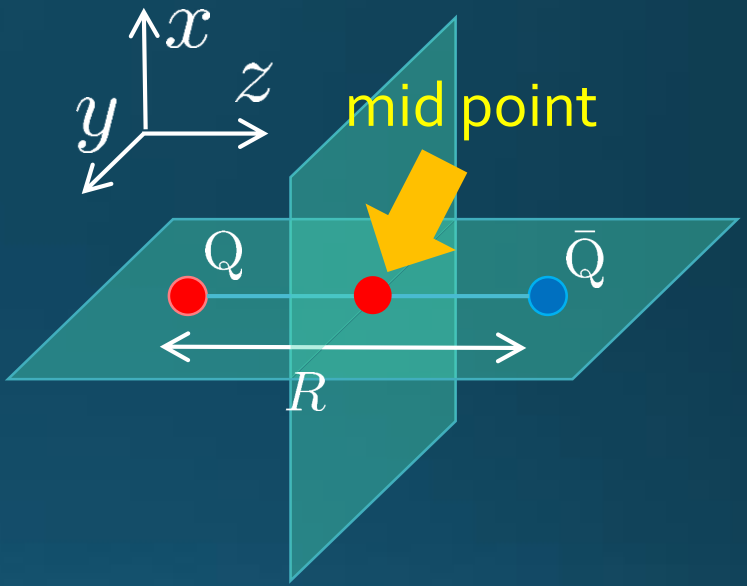


$$V(R) = -\lim_{T\to\infty} \log\langle W(R,T)\rangle$$

$$\langle O(x)\rangle_{Q\bar{Q}} = \lim_{T\to\infty} \frac{\langle \delta O(x)\delta W(R,T)\rangle}{\langle W(R,T)\rangle}$$

Ground State Saturation



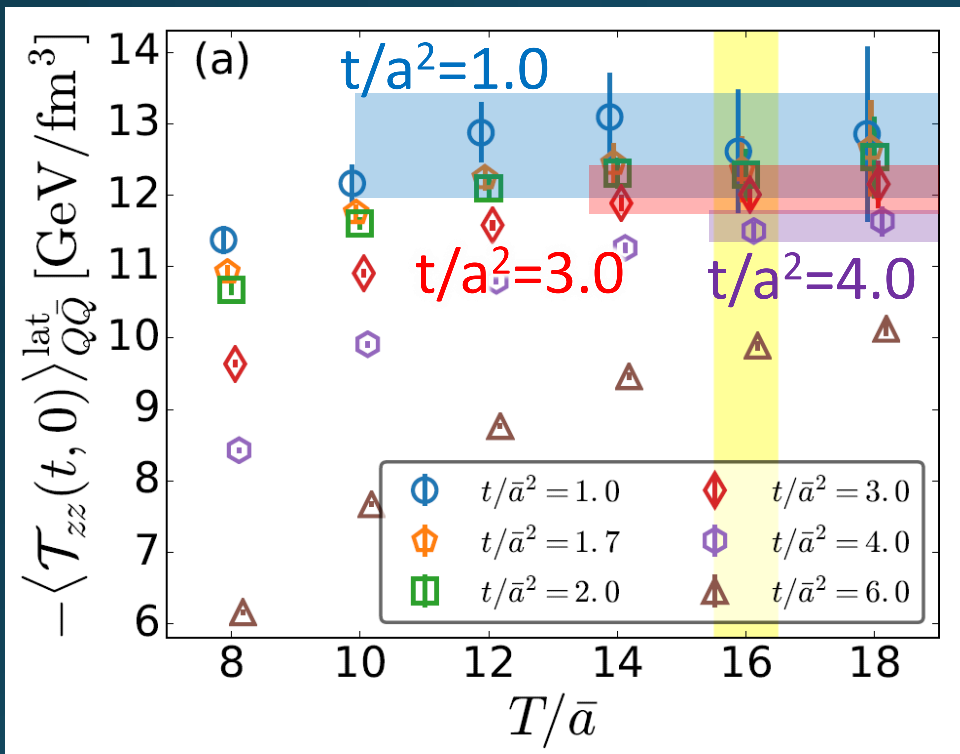


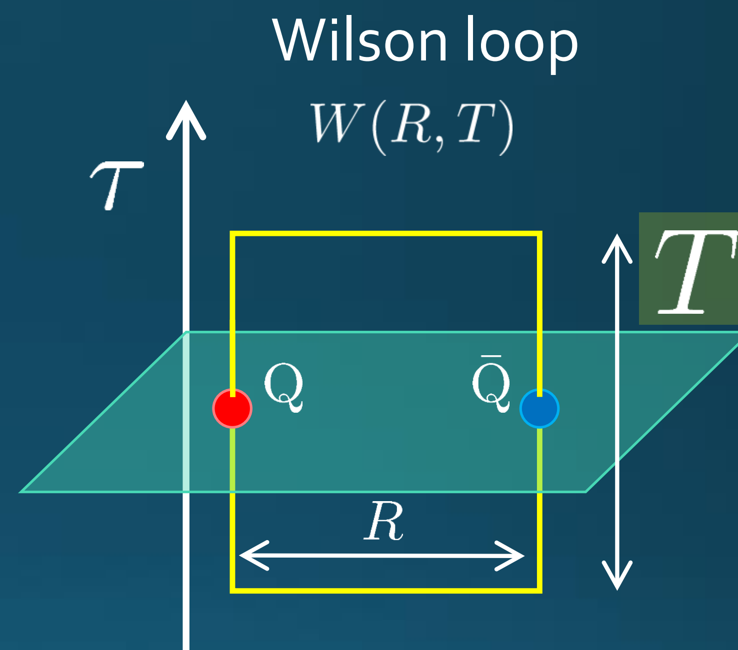
β=6.819 (a=0.029 fm), R=0.46 fm

Appearance of plateau for $t/a^2<4$, $T/a>15$

Grand state saturation under control

Ground State Saturation





β=6.819 (a=0.029 fm), R=0.46 fm

Appearance of plateau for $t/a^2<4$, $T/a>15$ → Grand state saturation under control

Abelian-Higgs Model

Abelian-Higgs Model

$$\mathcal{L}_{\mathrm{AH}} = -\frac{1}{4}F_{\mu\nu}^2 + |(\partial_\mu + igA_\mu)\phi|^2 - \lambda(\phi^2 - v^2)^2$$

GL parameter: $\kappa = \sqrt{\lambda}/g$

- type-I : $\kappa < 1/\sqrt{2}$
- type-II : $\kappa > 1/\sqrt{2}$
- Bogomol'nyi bound : $\kappa = 1/\sqrt{2}$

Infinitely long tube

- degeneracy

  $T_{zz}(r) = T_{44}(r)$ Luscher, 1981
- conservation law

$$\frac{d}{dr}(rT_{rr}) = T_{\theta\theta}$$