

# Temperature dependence of $\eta/s$ : uncertainties from the equation of state

Jussi Auvinen

Institute of Physics Belgrade

in collaboration with

Kari J. Eskola, Pasi Huovinen, Harri Niemi, Risto Paatelainen  
and Peter Petreczky

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# Equations of state

Hadron resonance gas at low temperatures, fit to lattice at high  $T$

- **s95p**:  $M < 2$  GeV HG particles from 2005 PDG summary tables  
P. Huovinen and P. Petreczky, NPA 837, 26 (2010)  
Lattice data used for fitting: Bazavov et al., PRD 80, 014504 (2009); Cheng et al., PRD 77, 014511 (2008)
- **s83z**: Hadron gas particles from 2016 PDG summary tables  
Lattice data: Borsanyi et al., PLB 730, 99 (2014)
- **s87r**:  $M < 2$  GeV HG particles from 2005 PDG summary tables  
Lattice data: Bazavov et al., PRD 90, 094503 (2014); Bazavov et al., PRD 97, 014510 (2018)
- **s88s**: HG particles from 2016 PDG summary tables  
Lattice data: Bazavov et al., PRD 90, 094503 (2014); Bazavov et al., PRD 97, 014510 (2018)

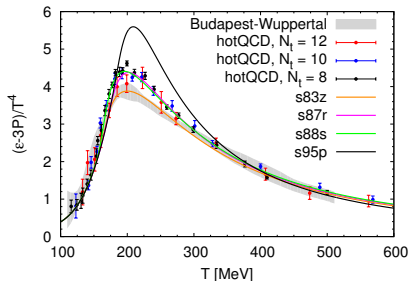


Figure: P. Huovinen

## EKRT+hydrodynamics model

- Initial energy density from the EKRT minijet saturation model

Paatelainen et al., PRC 87, 044904 (2013); PLB 731, 126 (2014)

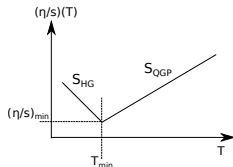
$$e(\vec{r}_T, \tau_s(\vec{r}_T)) = \frac{K_{sat}}{\pi} [p_{sat}(\vec{r}_T, K_{sat})]^4; \tau_s(\vec{r}_T) = 1/p_{sat}(\vec{r}_T, K_{sat})$$

- 2+1D viscous hydrodynamics with linear temperature dependence on shear viscosity coefficient  $\eta/s$

Niemi et al., PRC 93, 024907 (2016)

$$\eta/s(T) = S_{HG}(T_{min} - T) + (\eta/s)_{min}, \quad T < T_{min}$$

$$\eta/s(T) = S_{QGP}(T - T_{min}) + (\eta/s)_{min}, \quad T > T_{min}$$



Kinetic decoupling temperature  $T_{dec} = 120$  MeV

# Bayesian analysis

Model parameters (input):  $\vec{x} = (x_1, \dots, x_n)$

$(K_{sat}, T_{min}, (\eta/s)_{min}, S_{HG}, S_{QGP})$



Model output  $\vec{y} = (y_1, \dots, y_m) \Leftrightarrow$  Experimental values  $\vec{y}^{\text{exp}}$

$(N_{\text{ch}}(\sqrt{s_{NN}}, \text{centrality}), v_2(\sqrt{s_{NN}}, \text{centrality}))$

Bayes' theorem:

Posterior probability  $\propto$  Likelihood  $\cdot$  Prior knowledge

## Bayesian analysis

Posterior probability  $\propto$  Likelihood  $\cdot$  Prior knowledge

- Prior knowledge: Range of input parameter values to investigate (uniform hypercube)
- Likelihood:  $\mathcal{L}(\vec{x}) \propto \exp\left(-\frac{1}{2}(\vec{y}(\vec{x}) - \vec{y}^{\text{exp}})\Sigma^{-1}(\vec{y}(\vec{x}) - \vec{y}^{\text{exp}})^T\right)$ , where  $\Sigma$  is the covariance matrix

Use Markov chain Monte Carlo to sample the posterior probability

"emcee" ensemble sampler by D. Foreman-Mackey: <http://dfm.io/emcee/>

- Random walk in input parameter space, guided by the likelihood, constrained by the prior
- Running the full hydrodynamics model during random walk to compute likelihood infeasible timewise  
 $\Rightarrow$  use Gaussian process (GP) emulator to estimate model output for likelihood computations [http://scikit-learn.org/stable/modules/gaussian\\_process.html](http://scikit-learn.org/stable/modules/gaussian_process.html)

## Analysis procedure

Produce training data for emulator conditioning

- run simulations with  $\mathcal{O}(100)$  different parameter combinations
- rule of thumb:  $\sim 20$  points per parameter

One GP  $\leftrightarrow$  One observable

Reduce the dimension of the output space  
with principal component analysis (PCA)

<http://scikit-learn.org/stable/modules/decomposition.html>

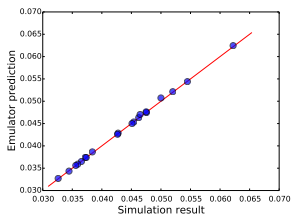


Condition the emulators on training data; check  
prediction quality



Calibrate on experimental data by running MCMC

Walker density in parameter space  $\rightarrow$  Posterior  
probability distribution as number of steps  $\rightarrow \infty$

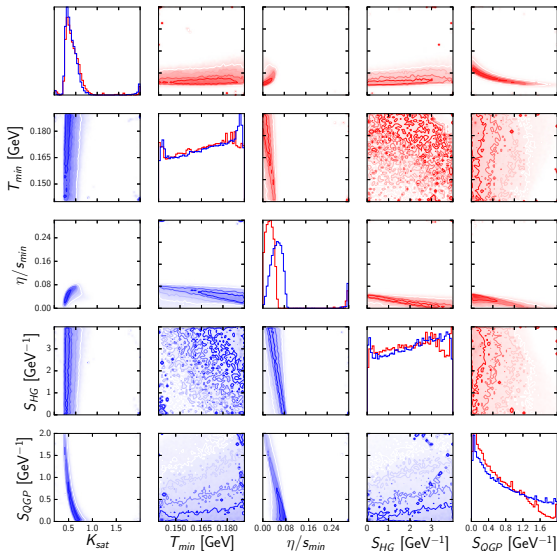


Emulator verification for  $v_2$  at

$\sqrt{s_{NN}} = 200$  GeV at (10-20)%

centrality, using s95p EoS

# Posterior probability distributions



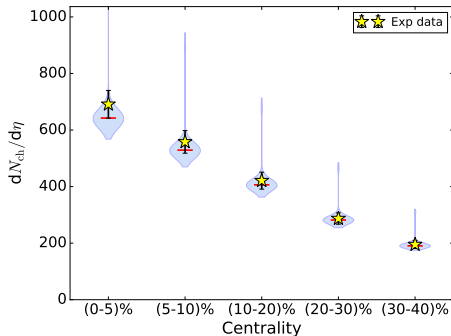
s83z

s95p

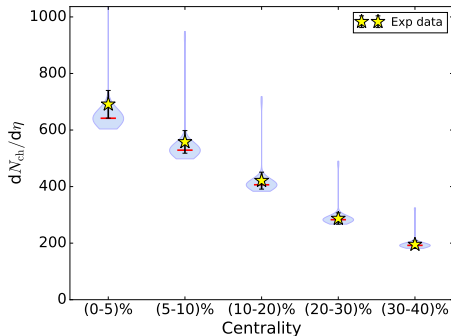
# Observables at $\sqrt{s_{NN}} = 200$ GeV

Emulator estimates for samples from the posterior distribution  
 Charged particle yield vs. centrality

s83z



s95p



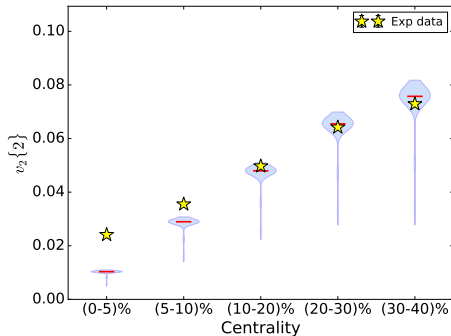
Experimental data: STAR, PRC 79, 034909 (2009)



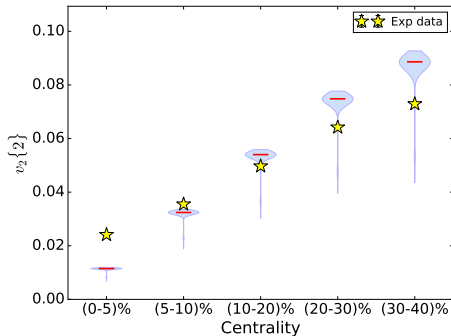
# Observables at $\sqrt{s_{NN}} = 200$ GeV

## Elliptic flow $v_2\{2\}$

### s83z



### s95p

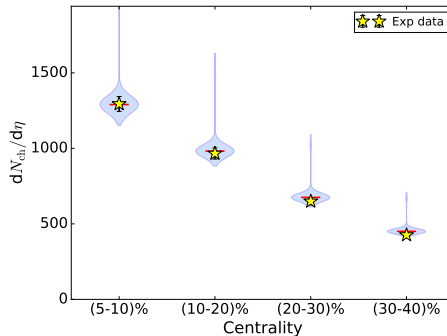


Experimental data: STAR, PRC 72, 014904 (2005)

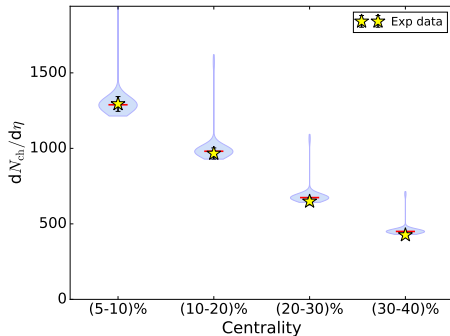
# Observables at $\sqrt{s_{NN}} = 2.76$ TeV

## Charged particle yield vs. centrality

s83z



s95p

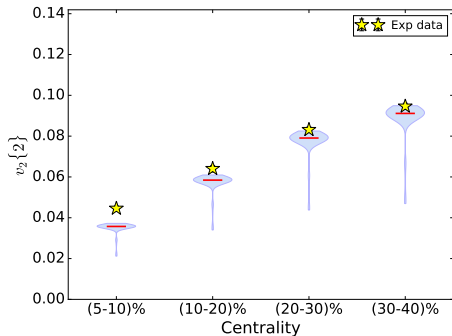


Experimental data: ALICE PRL 106, 032301 (2011)

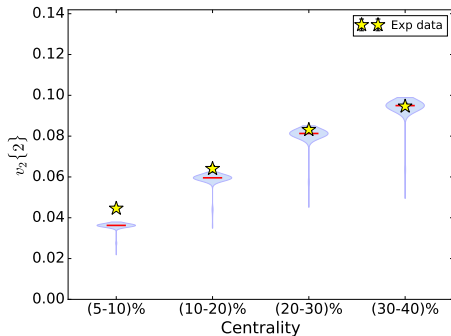
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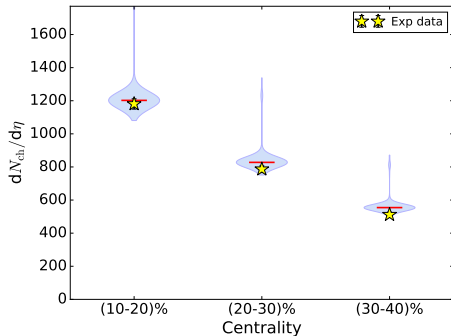


Experimental data: ALICE PRL 116, 132302 (2016)

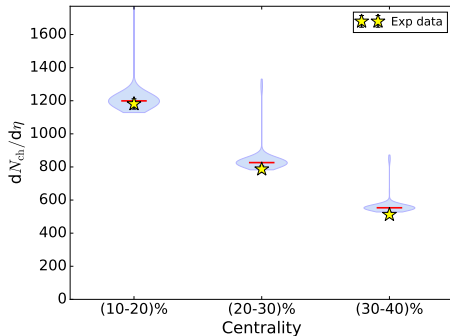
# Observables at $\sqrt{s_{NN}} = 5.02$ TeV

## Charged particle yield vs. centrality

s83z



s95p

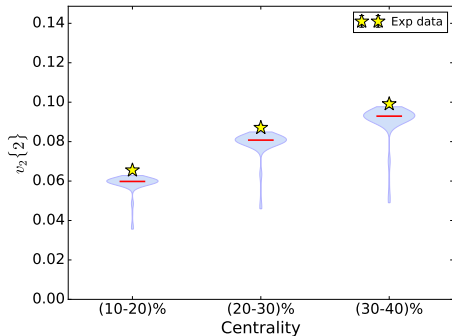


Experimental data: ALICE PRL 116, 222302 (2016)

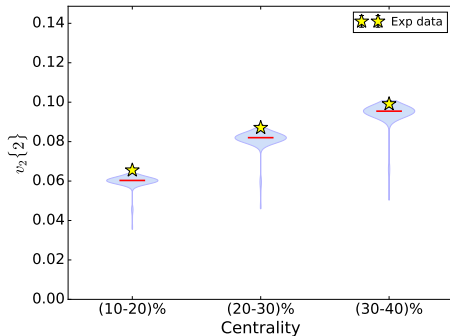
# Observables at $\sqrt{s_{NN}} = 5.02$ TeV

## Elliptic flow $v_2\{2\}$

### s83z

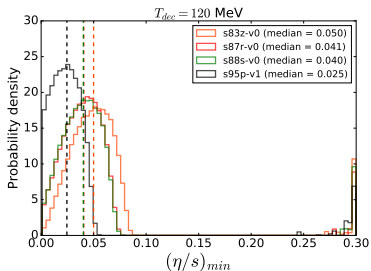
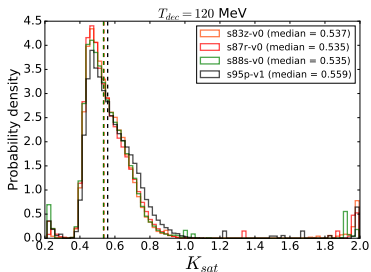


### s95p

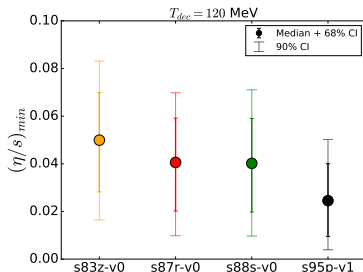


Experimental data: ALICE PRL 116, 132302 (2016)

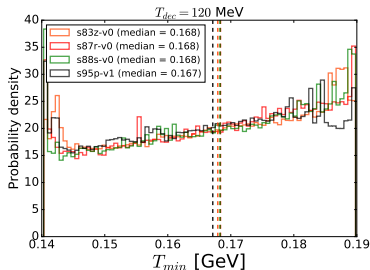
## Marginal posterior distribution comparisons



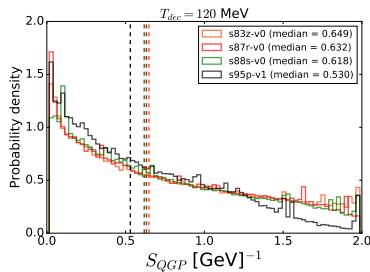
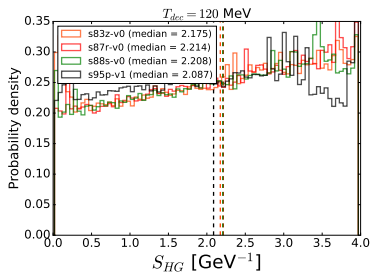
- $K_{sat}$  unaffected by EoS; consistent with the values used in previous studies ( $\approx 0.5 - 0.75$ )
- $(\eta/s)_{min}$  peak probability at higher values using newer equations of state compared to **s95p**
- Uncertainties remain large



## Marginal posterior distribution comparisons



- $T_{min}, S_{HG}$  not constrained by this analysis; considerable probability density on both edges  $\Rightarrow$  Prior range too small
- Smaller values of  $S_{QGP}$  favored; the skew is more pronounced for **s95p**



## Summary

- Bayesian analysis is a convenient method for simultaneously determining both the “best-fit” parameter values and their uncertainties.
- The choice of the equation of state has a notable effect on the extracted minimum value of  $\eta/s$ ; significant overlap remain in the posterior probability distributions, however
- Minijet saturation parameter  $K_{sat}$  unaffected by EoS
- Smaller slope values for  $eta/s(T)$  in QGP favored, weak dependence on EoS
- The place of the minimum (wrt temperature) and the linear slope of  $eta/s(T)$  in hadron gas remain unconstrained; prior ranges need to be extended
- Overall, the equations of state based on newer lattice data tend to require higher values of  $eta/s(T)$  for the best fit with experimental data



# Extra slides

## Gaussian process

Stochastic process: A parameterized collection of random variables  $\{y_t\}_{t \in \mathcal{T}}$  ( $\mathcal{T}$  possibly infinite).  
E.g. random walk over time.

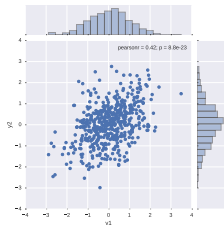
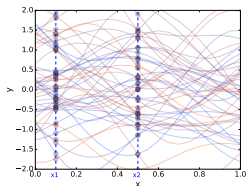
**Gaussian process:** A stochastic process, in which every finite set  $Y = \{y_i\}_{i=1}^N$  is a multivariate Gaussian random variable  $\mathcal{N}(\boldsymbol{\mu}, \boldsymbol{\Sigma})$ , where

$$\boldsymbol{\mu} = \boldsymbol{\mu}(X) = \{\mu(x_1), \dots, \mu(x_N)\}$$

is the mean and

$$\boldsymbol{\Sigma} = \sigma(X, X) = \begin{pmatrix} \sigma(\vec{x}_1, \vec{x}_1) & \cdots & \sigma(\vec{x}_1, \vec{x}_N) \\ \vdots & \ddots & \vdots \\ \sigma(\vec{x}_N, \vec{x}_1) & \cdots & \sigma(\vec{x}_N, \vec{x}_N) \end{pmatrix}$$

is the covariance matrix with **covariance function**  $\sigma(\vec{x}, \vec{x}')$ .



## Gaussian process

Choice: Squared-exponential covariance function with a noise term

$$\sigma(\vec{x}, \vec{x}') = \theta_0 \exp\left(-\sum_{i=1}^n \frac{(x_i - x'_i)^2}{2\theta_i^2}\right) + \theta_{\text{noise}}\delta_{\vec{x}\vec{x}'}$$

The *hyperparameters*  $\vec{\theta} = (\theta_0, \theta_1, \dots, \theta_n, \theta_{\text{noise}})$  are not known a priori and must be estimated from the given data

⇒ emulator **training**: Maximise the marginal likelihood (aka “evidence”)

$$\log P(Y|X, \vec{\theta}) = \underbrace{-\frac{1}{2} Y^T \Sigma^{-1}(X, \vec{\theta}) Y}_{\text{data fit}} \underbrace{-\frac{1}{2} \log |\Sigma(X, \vec{\theta})|}_{\text{complexity penalty}} \underbrace{-\frac{N}{2} \log(2\pi)}_{\text{normalization}}$$

## Gaussian process

To predict the model output  $y_0$  in an arbitrary point  $\vec{x}_0$ , we write the joint distribution

$$\begin{pmatrix} y_0 \\ Y \end{pmatrix} = \mathcal{N} \left( \begin{pmatrix} \mu(\vec{x}_0) \\ \mu(X) \end{pmatrix}, \begin{pmatrix} \Sigma_{0,0} & \Sigma_{0,X} \\ \Sigma_{X,0} & \Sigma_{X,X} \end{pmatrix} \right)$$

and calculate the conditional predictive mean

$$\bar{\mu}(\vec{x}_0) = \mu(\vec{x}_0) + \Sigma_{0,X} \Sigma_{X,X}^{-1} (Y - \mu(X)).$$

Typically we define the mean function  $\mu(\vec{x}) \equiv 0$  and the prediction is simply

$$\bar{\mu}(\vec{x}_0) = \Sigma_{0,X} \Sigma_{X,X}^{-1} Y$$

with the associated predictive (co)variance

$$\bar{\Sigma} = \Sigma_{0,0} - \Sigma_{0,X} \Sigma_{X,X}^{-1} \Sigma_{X,0}.$$

## Principal component analysis

$m$  observables  $\Rightarrow m$  Gaussian processes needed for model emulation

However,  $m$  can grow to  $\mathcal{O}(100)$  at top RHIC energies and at the LHC!  
Number of emulators can be reduced with **principal component analysis**

First principal component represents the direction of largest variance in output space, second PC the direction of second largest variance, etc.

- Fraction of variance explained by principal component  $p_q$ :

$$\text{Var}(p_q) = \frac{\lambda_q}{\sum_{i=1}^m \lambda_i}$$

- Select the number of principal components which together explain desired fraction of total variance; often only a few PCs are needed to explain 99% of the variance

## Likelihood function

The likelihood function used in MCMC:

$$\frac{1}{|2\pi(\Sigma_{\text{exp}} + \Sigma_{GP})|} \exp\left(-\frac{1}{2}(\vec{z}_{GP}^* - \vec{z}_{\text{exp}})(\Sigma_{\text{exp}} + \Sigma_{GP})^{-1}(\vec{z}_{GP}^* - \vec{z}_{\text{exp}})^T\right)$$

- $\vec{z}_{GP}^*$  is the emulator prediction at the input parameter point  $\vec{x}^*$
- $\vec{z}_{\text{exp}}$  is the experimental data transformed to principal component space
- $\Sigma_{GP}$  is the predictive variance (emulator uncertainty)
- $\Sigma_{\text{exp}}$  is the experimental error squared, transformed to PC space