

# Structure of pion and kaon from lattice QCD

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In collaboration with:

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## Progress on our lattice calculation of pion and kaon PDFs

[work done in collaboration with: Colin Egerer, Joe Karpie, Kostas Orginos, Jianwei Qiu, David Richards, Raza Sufian]

Calculation of pion electromagnetic form factor at high momentum transfer

[work done in collaboration with: Raul Briceno, Robert Edwards, David Richards]

- "Light-like" separated hadonic tensor-K-F Liu et al Phys. Rev. Lett. 72 1790 (1994), Phys. Rev. D62 (2000) 074501 A Chambers et.al (2017) 1703.01153
- Position-space correlators: V. M. Braun, D. Müller, EPJ (2008)
- Quasi PDF using LaMET:

X. Ji, Phys.Rev.Lett. 110, (2013), C Alexandrou et.al. (2018) 1803.02685 J.-W. Chen et.al. (2018) 1803.04393

• Pseudo PDF :

A. Radyushkin Phys.Lett. B767 (2017), K. Orginos et. al. (2017) 1706.05373

• New concept based on QCD factorisation— "Good lattice cross-sections" Y.-Q. Ma J.-W. Qiu (2014) 1404.6860, Y.-Q. Ma, J.-W. Qiu (2017) 1709.03018

# PDFs from "Good" Lattice Cross-sections

"Good" lattice cross sections [Ma & Qiu, Phys.Rev.Lett. 120 (2018) no.2, 022003]

If we can form single hadron matrix elements of renormalizable nonlocal operators such that -

(1) it is calculable in lattice- QCD with an Euclidean time,

$$\sigma_n(\omega,\xi^2,P^2) = \langle P|T\{\mathcal{O}_n(\xi)\}|P\rangle$$
$$\mathcal{O}_{j_1j_2}(\xi) \equiv \xi^{d_{j_1}+d_{j_2}-2} Z_{j_1} Z_{j_2} j_1(\xi) j_2(0)$$

(2) it has a well-defined continuum limit (UV finite) at the lattice spacing zero limit,

(3) it has the same and factorizable logarithmic collinear (CO) divergences as that of PDFs



the single hadron matrix element can be factorized into PDFs to all orders IR-safe hard coefficients with controllable power corrections, provided  $\xi^2$  is sufficiently small,

$$\sigma_n(\omega, \xi^2, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) \\ \times K_n^a(x\omega, \xi^2, x^2 P^2, \mu^2) + O(\xi^2 \Lambda_{\text{QCD}}^2)$$

loffe time:  $\omega = P.\xi$ 

B. L. loffe, Phys. Lett. 30B, 123 (1969)

$$f_{\bar{a}}(x,\mu^2) = -f_a(-x,\mu^2)$$

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 $\mathcal{O}_{j_1 j_2}(\xi) \equiv \xi^{d_{j_1} + d_{j_2} - 2} Z_{j_1} Z_{j_2} j_1(\xi) j_2(0)$ 

 $\mathcal{O}_S(\xi) = \xi^4 Z_S^2[\overline{\psi}_q \psi_q](\xi) [\overline{\psi}_q \psi_q](0) ,$ 

 $\mathcal{O}_V(\xi) = \xi^2 Z_V^2 [\overline{\psi}_q \not \xi \psi_q](\xi) [\overline{\psi}_q \not \xi \psi_q](0) ,$ 

#### What currents we can use?

 $\sigma_n(\omega,\xi^2,P^2) = \langle P|T\{\mathcal{O}_n(\xi)\}|P\rangle$ 

Including flavor changing and gluonic currents with any single hadronic state

# **Motivation for pion and Kaon PDF calculation**

- Large-x behaviour of pion valence distribution an unresolved problem- (1-x)<sup>2</sup> or (1-x)<sup>1</sup>?
- C12-15-006 experiment at JLab to explore large-x behavior
- Limited Old data from Drell-Yan for Kaon PDF
- JLab approved pion-TDIS will also measure Kaon structure functions for the first time ever at small x
- Long-searched-for Sullivan process for accessing the kaon structure function



EIC will add large (x,Q<sup>2</sup>) landscape for both pion and kaon! Recent DSE calculations of pion and Kaon PDFs and ratios



First-principle lattice QCD calculation of needed – First calculation of Kaon PDF on the lattice

- Need calculation of four point functions, with 'red' being different heavier quarks [similar technique by Bali et. al. 1807.03073]
- Isoclover configurations 32<sup>3</sup> x 96, lattice spacing 0.127 fm, pion mass around 405 MeV
- Current renormalisations known from previous calculation, arXiv:1611.07452
- Two-exponential fits for ratios of 4pt and 2pt functions





# **Preliminary lattice results for pion**

- Analysis shown here on isoClover with 450 Configurations
- Momentum smearing used for higher momentum

[Gunnar S. Bali, et al (PRD 2016)]

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V-A matrix element

T is the source sink separation



## **Preliminary lattice results for pion**



V-V matrix element

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## "Global" fit for lattice data - technique

$$\sigma_n(\nu, z^2, p^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(x\nu, z^2, x^2p^2, \mu^2) + O(z^2\Lambda_{QCD}^2)$$
Many such lattice
"cross-sections"
To be calculated
$$f_{\bar{a}}(x, \mu^2) = -f_a(-x, \mu^2)$$
to obtain valence quark and
anti-quark distribution
$$Perturbative Kernel known,$$
Oscillatory [being calculated at LO and NLO]
$$e Need many P.\xi \text{ data to have a good global fit} for f(x)$$

$$A \mod I fit form [JAM Collaboration, Image description]$$

 $f_{abcd}(x) = N_{abcd} x^{a} (1-x)^{b} (1 + c \sqrt{x} + d x)$ 

arXiv:1804.01965

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## **Preliminary lattice results for pion**



For details: talk by Raza Sufian at Lattice 2018

## (Very) preliminary lattice results for Kaon

## Calculation for Kaon PDF looks encouraging too



• Comparison between:

 $(J_1 = V, J_2 = A and vice versa)$  and different  $\xi = z$ 

• 
$$P_z = 1 ~(\sim 0.3 ~GeV)$$

- Real part of the 4pt correlator
- Strange quark is the active quark
- No signal after  $\xi = 6$



• Comparison between: ( $J_1 = V, J_2 = A$  and vice versa) and different  $\xi = z$ 

• 
$$P_z = 1 \ (\sim 0.3 \ \text{GeV})$$

- Real part of the 4pt correlator
- Light quark is the active quark
- No signal after  $\xi = 6$



- Need lot more P.ξ data points for a realistic f(x) fit
- And with higher P values
- Real part of the V-A 4pt correlator multiplied by vector current renormalisation and a factor ξ<sup>4</sup>
- Strange quark is the active quark

For details: talk by Bipasha Chakraborty at Lattice 2018

## Ongoing...

- We have demonstrated that this method works
- Higher momenta, more statistics, calculation at different lattice spacings, quark masses and volumes ongoing
- Methods to tackle ill posed inverse Fourier Transform being developed [ref. talk by Savvas Zafeiropoulos]
- In future, look into finite volume effects [ref. R. Briceno et. al. Phys. Rev. D 98, 014511 (2018)], although expect to be small
- Generalisation of other methods pseudo, quasi, T33
- We can also get polarised PDFs, moments

# Pion electromagnetic form factor at high $Q^2$ from lattice QCD

# Definition



#### Interplay between hard and soft scales





Need better understanding of the transition to the asymptotic region

## JLAB 12 GeV upgrade



G. Huber and D. Gaskell

 $F_{\pi}$  measurements at  $Q^2 \sim 6 - 10$  GeV<sup>2</sup>: E12-06-101 at JLAB Hall C, For EIC,  $Q^2 \sim 30$  GeV<sup>2</sup>

Can we get some insight from first principles lattice QCD calculations to the question - where does the transition to pQCD happen?



### **Two-point correlator construction**



$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^{\dagger}(0) | 0 \rangle$$

• Basis of operators

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \cdots \overleftrightarrow{D} \psi$$

- Optimized operator for state |n>  $\Omega_{\mathfrak{n}}^{\dagger}=\sum_{i}w_{i}^{(\mathfrak{n})}\mathcal{O}_{i}^{\dagger}$ 

in a variational sense by solving generalized eigenvalue problem-

$$C(t) v^{(\mathfrak{n})} = \lambda_{\mathfrak{n}}(t) C(t_0) v^{(\mathfrak{n})}$$

• Diagonalize the correlation matrix – eigenvalues

$$\lambda_n(t) = \exp[-En(t-t_0)]$$



#### Correlator Construction: smearing of quark fields - 'distillation' with

$$\Box_{\vec{x}\vec{y}}(t) = \sum_{n=1}^{N_D} \xi_{\vec{x}}^{(n)}(t) \; \xi_{\vec{y}}^{(n)\dagger}(t)$$

Extraction of low lying hadron states

#### Meson creation operator :

$$\mathcal{O}^{\dagger}(\vec{p}) = \bar{\psi}_{\vec{x}} \Box_{\vec{x}\vec{y}} e^{-i\vec{p}\cdot\vec{y}} \Gamma_{\vec{y}\vec{z}} \Box_{\vec{z}\vec{w}} \psi_{\vec{w}}$$



Parambulators by inverting the Dirac matrix

Operator construction with momentum projection



#### Form factor calculation

Need three-point correlator – calculated with Weighted operators  $\pi(\vec{p}_2)$  $\pi$  ( $\vec{p}_1$  $C_{\mathrm{f}\mu\mathrm{i}}(\Delta t, t) = \langle 0 | \mathcal{O}_{\mathrm{f}}(\Delta t) \, j_{\mu}(t) \, \mathcal{O}_{\mathrm{i}}^{\dagger}(0) | 0 \rangle$ Т  $\left(Z_{V}\right) < \pi^{+}(p_{2})|J^{\mu}_{\pi}(0)|\pi^{+}(p_{1}) > = e(p_{1} + p_{2})^{\mu}F_{\pi}(q^{2})$ Clover  $Z_V$  calculated using  $F_{\pi}(q^2 = 0) = 1$ discretised fermion action

### Pion electromagnetic form factor: up to $Q^2 = 1 \text{ GeV}^2$



Pion charge radius:

$$\langle r^2 \rangle \equiv -6 \frac{d}{dQ^2} F(Q^2) \big|_{Q^2 = 0}$$

Parametrising Q<sup>2</sup> dependence  $Q^2 < 0.3 \,\mathrm{GeV}^2$ 

$$F_{\pi}(Q^2) = F(0) \frac{1}{1 + Q^2/m^2}$$

 $m_{\pi} = 750 \text{ MeV} : 0.47(6) \text{ fm}$ 

 $m_{\pi} = 390 \text{ MeV} : 0.55 (10) \text{ fm}$ 

Lightest vector meson mass

Comparison of different lattice results:

