

VH production with $H \rightarrow b\bar{b}$ decay in full NNLO QCD

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**LHC Higgs Cross Section Working Group
VH subgroup – June 29th 2017**

In collaboration with: M. Grazzini, G. Somogyi & F. Tramontano

Motivations

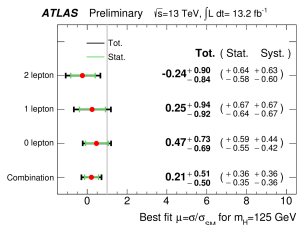
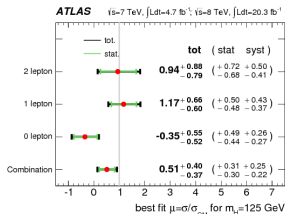
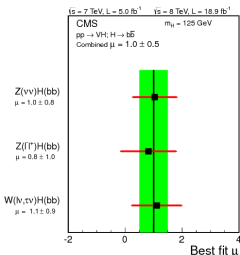
Associated vector boson Higgs (VH) production (with $H \rightarrow b\bar{b}$ and $V \rightarrow l_1 l_2$ decay)

- Important channel at the LHC (access to $Hb\bar{b}$ coupling) but challenging.

LHC experiments are close to the SM $H \rightarrow b\bar{b}$ sensitivity.

To improve these results, precise theoretical predictions are needed

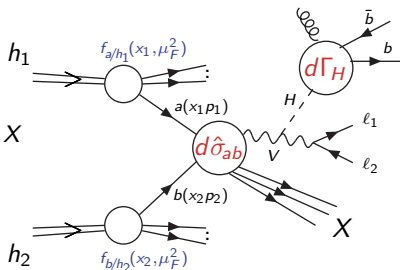
⇒ computation of higher-order QCD corrections.



Associated VH production with $H \rightarrow b\bar{b}$ decay

$$h_1(p_1) + h_2(p_2) \rightarrow V + H + X \rightarrow \ell_1 \ell_2 + b\bar{b} + X$$

where $V = Z^0, W^\pm$ and $\ell_1 \ell_2 = \ell^+ \ell^- , \nu \ell$



QCD factorization formula

$$d\sigma = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}_{ab}(x_1 p_1, x_2 p_2; \mu_F^2) + \mathcal{O}\left(\frac{\Lambda_{QCD}}{Q}\right)^p$$

- By using the zero width approximation ($\Gamma_H \ll m_H$)

$$d\sigma_{VH \rightarrow Vb\bar{b}} = d\sigma_{VH} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}}{\Gamma_H} = d\sigma_{VH} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}}{\Gamma_{H \rightarrow b\bar{b}}} \times \text{Br}(H \rightarrow b\bar{b}),$$

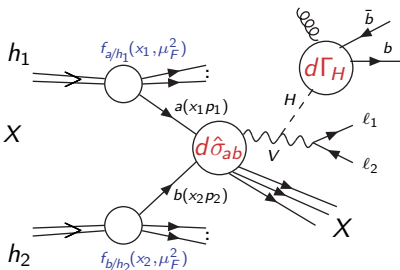
- Perturbative expansion gives

$$d\sigma_{VH \rightarrow Vb\bar{b}}^{LO+l_0} = d\sigma_{VH}^{(0)} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)}} \times \text{Br}(H \rightarrow b\bar{b}),$$

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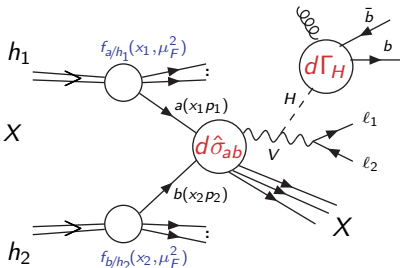
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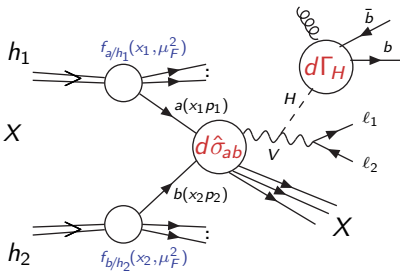
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$$d\sigma_{VH \rightarrow Vb\bar{b}}^{NLO+lo} = \left(d\sigma_{VH}^{(0)} + d\sigma_{VH}^{(1)} \right) \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)}} \times \text{Br}(H \rightarrow b\bar{b}),$$

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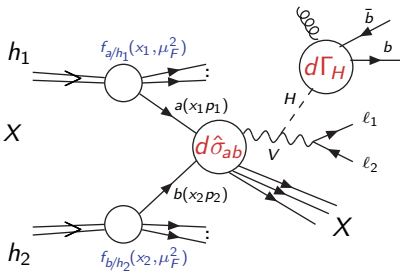
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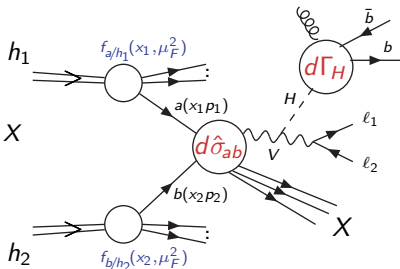
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VH production and decay in full NNLO QCD

G.F., Somogyi, Tramontano arXiv:1705.0304

Fully differential NNLO calculation for VH production including $H \rightarrow b\bar{b}$ at NNLO and $V \rightarrow l_1 l_2$ decays with spin correlations.

- NNLO calculation for $h_1 h_2 \rightarrow VH + X$ production calculated in [G.F., Grazzini, Tramontano('11, '15)] within the q_T -subtraction formalism [Catani, Grazzini('07)] requires:
 - Up to $d\sigma_{NLO}^{VH+jets}$.
 - $\mathcal{H}^{VH(1)}$ and $\mathcal{H}^{VH(2)}$ [Catani, Cieri, de Florian, G.F., Grazzini('09, '12)]: contains the finite-part of the one- and two-loops amplitude $c\bar{c} \rightarrow VH$.
 - Up to $d\sigma_{NLO}^{CT}$: depends by the (universal) q_T -resummation coefficients [Bozzi, Catani, de Florian, Grazzini('09, '12)].
- $H \rightarrow b\bar{b}$ decay at NNLO calculated by [DeL Duca, Duhr, Somogyi, Tramontano, Trocsanyi('15)] with CoLoRFulNNLO method [DeL Duca, Somogyi, Trocsanyi('07)].
- Fully inclusive QCD effects in the H decay taken into account by normalizing the $Hb\bar{b}$ branching fraction to the LHCHSWG-YR result.

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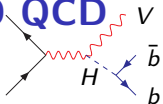
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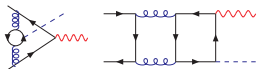
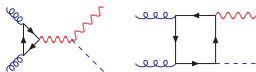
VH production and decay in full NNLO QCD



Our fully differential calculation implemented in the parton level code **HVNNLO**.

For VH prod. we have consistently included:

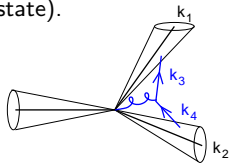
- NNLO DY-like QCD corrections
(bulk of NNLO correction for WH)
[VanNeerven et al.('91)]
- $gg \rightarrow HZ$ top-loop $\sim g^2 \lambda_t^2 \alpha_S^2$
(non DY-like) corrections [Kniehl('90)]
(important at the LHC due to large gg luminosity).
- NNLO top-mediated contributions
 $\sim g^3 \lambda_t \alpha_S^2$ to VH
[Brein,Harlander,Wiesemann,Zirke('11)]
(we included only the terms calculated with the full m_t dependence)



b-quark jets identification

We are interested in the identification of the b -quark jet which originate from the Higgs boson (b -quark treated in massless approximation).

- We consistently include b -quark emissions from initial and final state partons (at NNLO up to four b -quarks in the final state).
- Standard jet alg. not infrared and collinear safe definition of flavoured jets: splitting of a gluon in a soft or collinear (massless) $b\bar{b}$ pair affect the flavour of a jet.



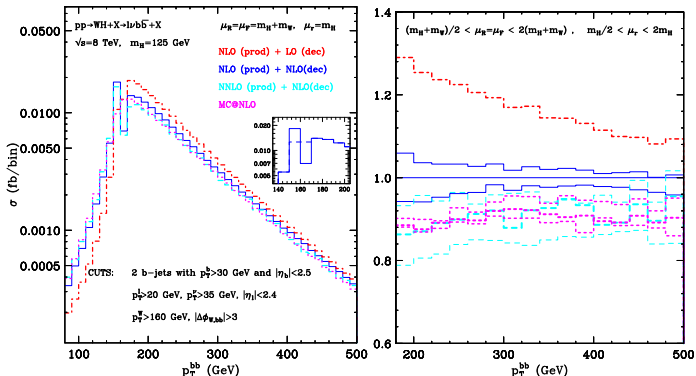
- **Collinear unsafety** removed by defining “b-jet” if contains $N(b) - N(\bar{b}) \neq 0$.
- **Infrared unsafety** removed by using the “flavour- k_T ” algorithm [Banfi, Salam, Zanderighi('06)]

$$d_{ij}^{(F)} = (\Delta\eta_{ij}^2 + \Delta\phi_{ij}^2) \times \begin{cases} \max(k_{ti}^2, k_{tj}^2), & \text{softer of } i, j \text{ is flavoured} \\ \min(k_{ti}^2, k_{tj}^2), & \text{softer of } i, j \text{ is flavourless} \end{cases}$$

(numerical difference with respect to standard alg. small in our case).

Numerical results at the LHC

WH production

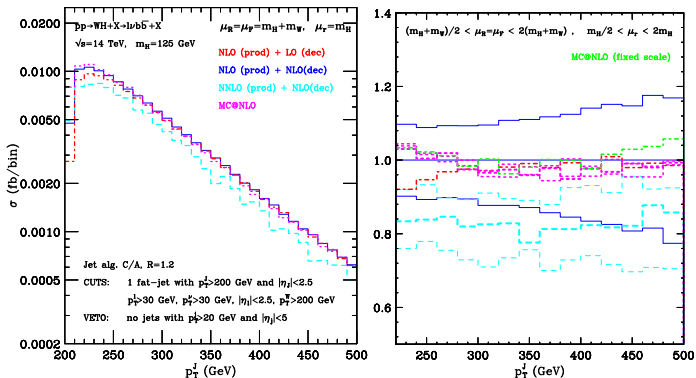


LHC at 8 TeV analysis

Left panel: p_T spectrum of the b -jets pair.

Right panel: Spectra normalized to the full NLO results (perturbative scale - μ_R, μ_F, μ_{Rdec} uncertainty bands are shown).

WH production

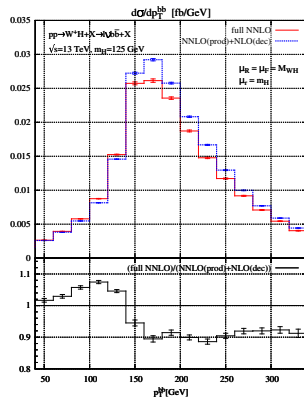
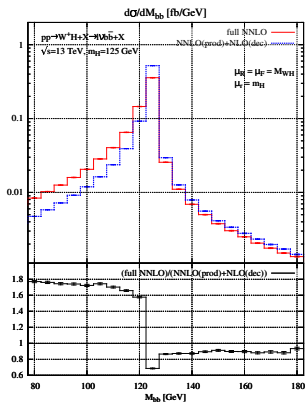


LHC at 14 TeV “fat-jet” analysis. Left panel: p_T spectrum of the fat jet. Right panel: Spectra normalized to the full NLO results (Perturbative scale - μ_R, μ_F, μ_{Rdec} uncertainty bands are shown).

NEW: WH production at full NNLO

| σ (fb) | NNLO(prod)+NLO(dec) | full NNLO |
|---------------|---------------------|-----------------|
| W^+H | 4.23 ± 0.02 | 3.96 ± 0.02 |

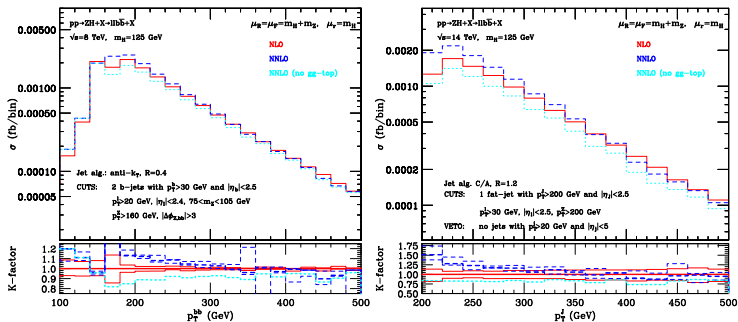
$$K_{fact} \sim -6.5\%$$



LHC13 analysis: $p_T^l > 15$ GeV, $|\eta_l| < 2.5$, $p_T^W > 150$ GeV, 2 b -jets $p_T^b > 25$ GeV, $|\eta_b| < 2.5$, flavour- k_T $R = 0.5$.

Left panel: M_{bb} spectrum of the b -jets pair. Right panel: p_T^{bb} spectrum of the b -jets pair. Lower panels: spectra normalized to the NNLO+nlo results.

ZH production

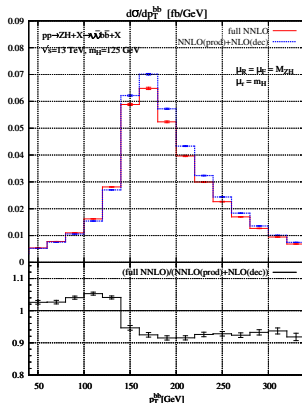
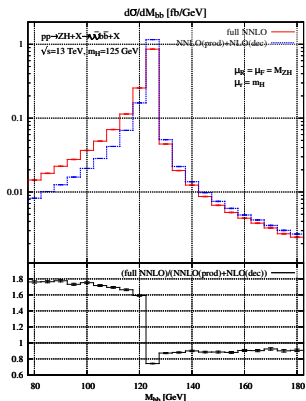


Left panel: **LHC8 analysis** p_T spectrum of the b -jets pair. Right panel: **LHC14 fat-jet analysis** p_T spectrum of the fat jet. Lower panels: spectra normalized to the full NLO results (perturbative scale - μ_R, μ_F uncertainty bands are shown).

NEW: ZH production at full NNLO

| σ (fb) | NNLO(prod)+NLO(dec) | full NNLO |
|---------------|---------------------|-----------------|
| ZH | 8.58 ± 0.02 | 8.10 ± 0.02 |

$K_{fact} \sim -5.5\%$



LHC13 analysis: $E_T^{miss} > 150$ GeV, 2 b -jets $p_T^b > 25$ GeV, $|\eta_b| < 2.5$, flavour- k_T $R = 0.5$.

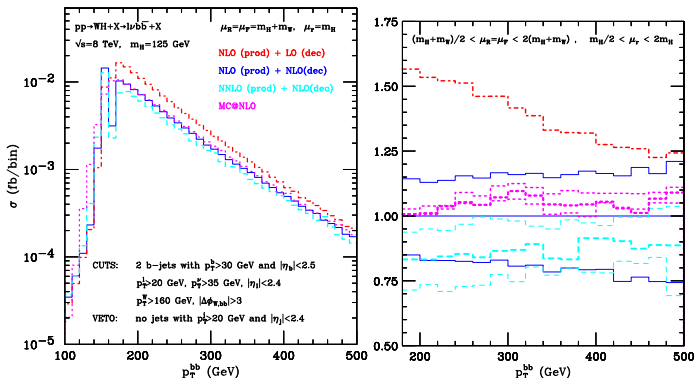
Left panel: M_{bb} spectrum of the b -jets pair. Right panel: p_T^{bb} spectrum of the b -jets pair. Lower panels: spectra normalized to the NNLO+nlo results.

Conclusions

- Associated vector boson Higgs (VH) production important (but challenging) channel at the LHC . ATLAS and CMS experiments are close to the SM sensitivity. Precise theoretical predictions are important.
- Calculation of **full NNLO QCD** corrections to VH production with $H \rightarrow \bar{b}b$ decay in hadron collision using the q_T -subtraction and **CoLoRFuINNLO** formalism, included in the **fully-exclusive** parton-level Monte Carlo code **HVNNLO**.
- Perturbative corrections are important and strongly depend on the experimental selection cuts. Illustrative phenomenological results with typical cuts applied in ATLAS and CMS analysis at the LHC.
- **Full NNLO** QCD corrections significantly reduces the accepted cross section ($\sim -6\%$) and substantially affect the distributions with respect to the partial NNLO result.

Back up slides

Associated WH production

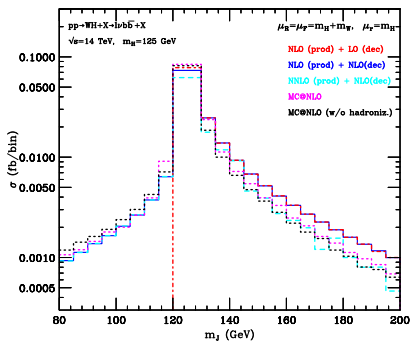


LHC8 with CUTS + VETO (on light jets)

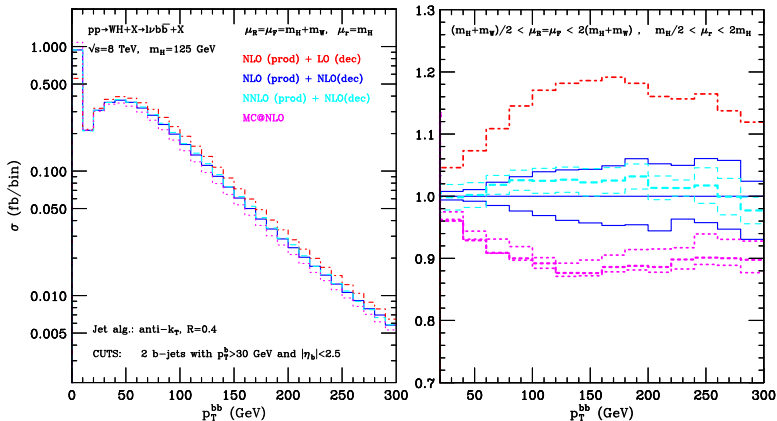
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Associated WH production



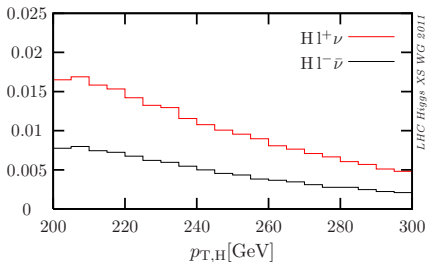
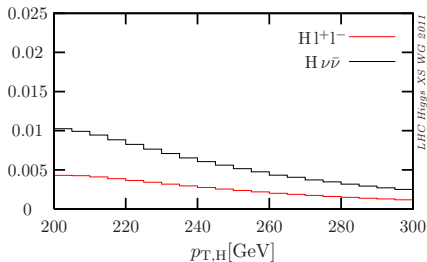
LHC14 fat-jet analysis. Invariant mass distribution of the **fat jet** computed at fixed-order QCD and at MC@NLO without hadronization (black dots) and with default MC@NLO (magenta dots).



LHC8 with NO CUTS (except the b -jet selection).

Left panel: p_T spectrum of the b -jets pair.

Right panel: Spectra normalized to the full NLO results (Perturbative scale - μ_R, μ_F, μ_{Rdec} uncertainty bands are shown).

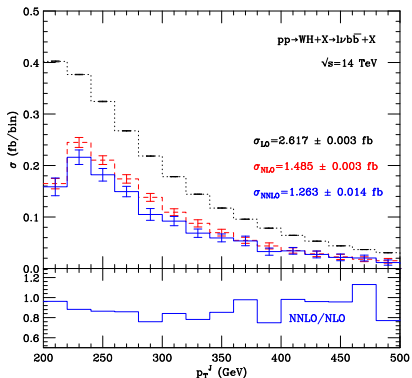
$d\sigma/dp_{T,H}[\text{fb/GeV}]$  $d\sigma/dp_{T,H}[\text{fb/GeV}]$ 

Yellow Report II:arXiv:1201.3084

Distributions in $p_{T,H}$ for $pp \rightarrow WH \rightarrow l\nu H$ (NNLO QCD + NLO EW) and for $pp \rightarrow ZH \rightarrow ll̄\nu\nū H$ (NLO QCD + NLO EW) at $\sqrt{s} = 7 \text{ TeV}$.

Boosted setup: $|\eta_l| < 2.5$, $p_{T,l} > 20 \text{ GeV}$, $p_{T,\nu} > 25 \text{ GeV}$, $p_{T,H} > 200 \text{ GeV}$, $p_{T,W/Z} > 190 \text{ GeV}$.

We produced similar results at $\sqrt{s} = 8 \text{ TeV}$.



$pp \rightarrow WH(\rightarrow l\nu b\bar{b})$

p_T spectra of the fat jet at the LHC@14TeV for $m_H = 120\text{GeV}$ at LO (dots), NLO (dashes) and NNLO (solid).

- Selection strategy of [Butterworth et al. ('08)]: search a large- p_T Higgs boson thorough a collimated $b\bar{b}$ pair decay.
Cuts:
Leptons: $p_T^l > 30\text{GeV}$, $|\eta^l| < 2.5$,
 $p_T^{\text{miss}} > 30\text{GeV}$, $p_T^W > 200\text{GeV}$.
Jets: Cambridge/Aachen algorithm with $R=1.2$.
Fat jet (contain the $b\bar{b}$) $p_T^J > 200\text{GeV}$,
 $|\eta^J| < 2.5$
Jet veto: No other jets with $p_T > 20\text{GeV}$ and $|\eta| < 5$.
- Large negative higher-order corrections: NLO (NNLO) effects -52%/-36% (-6%/-19%), depending on the scale choice (factor two around $\mu_F = \mu_R = m_W + m_H$).
- Jet veto strongly affect the higher order corrections \Rightarrow stability of fixed order calculation challenged.