

Telescoping Deconstruction and Collinear Drop

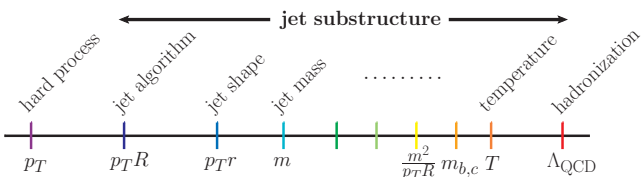
Yang-Ting Chien

LHC-TI Fellow, MIT Center for Theoretical Physics

July 19, BOOST 2018, Paris

In collaboration with A. Emerman, S.-C. Hsu, S. Meehan, Z. Montague (arXiv:1711.11041)
R. Elayavalli (arXiv:1803.03589), and I. Stewart (to appear soon)

The goals of jet substructure studies

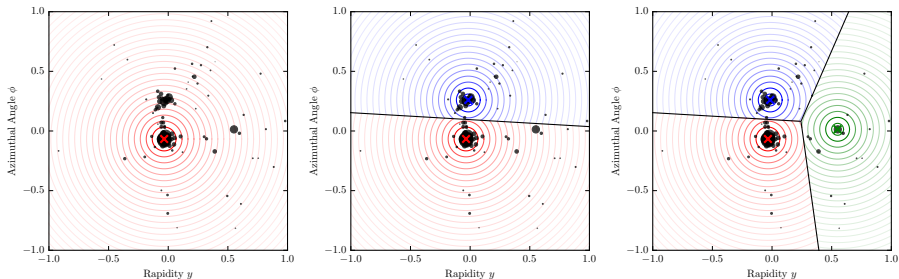


Relate precise jet modifications to medium properties

- ▶ Heavy ion jet physics: jet physics in $X + Y$ collisions, $X, Y \in \{e, p, A\}$
- ▶ The ultimate arena for jet and QCD/SM studies
- ▶ Precision is the key

Goals	TD	CD
Jet grooming	yes	included
Jet-initiating particle tagging	yes	yes
Test the accuracy of the description of soft radiation in MC	yes, →	yes
Test methods of predicting hadronization corrections	yes, →	yes
Illuminate the inner working of Quark-Gluon Plasma	hopefully	hopefully
Goal in the world cup	no	no

Telescoping Deconstruction: a complete subjet fragmentation basis



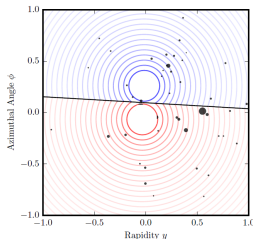
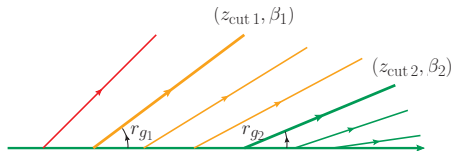
- ▶ A fixed-order N subjet expansion with subjet kinematics (p_T, y, ϕ, m)
 - ▶ identify dominant energy flow directions using N soft recoil-free axes
 - ▶ reconstruct subjets around the axes with multiple subjet radii R_T
 - ▶ TD observables represent *subjet topology* and *subjet substructure*
- ▶ Closely related to perturbative expansion and parton shower picture

arXiv:1803.03589, thanks to P. Komiske and E. Metodiev for early stage collaboration

Collinear Drop: veto energetic, collinear particles

- ▶ Understanding soft QCD is the goal
- ▶ Monte Carlo accuracy limited by soft radiation and hadronization modeling
- ▶ Heavy ion medium scale is low
- ▶ Want to *directly* probe soft physics by disentangling hard components of jets
- ▶ Implementation examples: (i)TD, (ii)soft drop and (iii)flattened angularity

(i) TD

(ii) Two soft-drop settings $z_{\text{cut}1} < z_{\text{cut}2}$, $\beta_1 \geq \beta_2$ (iii) Flattened angularity

$$\tau_\omega = \sum_{i \in \text{jet}} z_i \omega(\theta_i)$$

Suppress collinear and wide-angle radiation

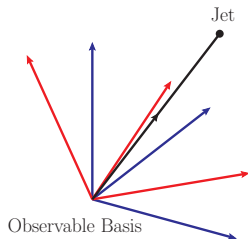
$$\omega(\theta) \rightarrow 0, \quad \theta \rightarrow 0, R$$

With I. Stewart, to appear soon

Use quark and gluon jets to probe heavy ion collisions



- ▶ Classify quark jets and gluon jets in pp and AA
- ▶ Different multivariate techniques suit different jet representations
 - ▶ list of physics-motivated observables
 - ▶ unbiased and raw input
 - ▶ complete basis and expansion
- ▶ Modern computation power and deep learning tools help benchmark jet feature identification

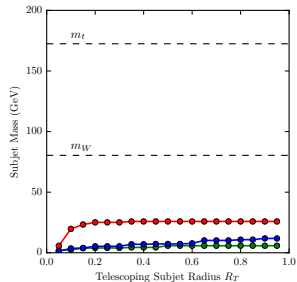
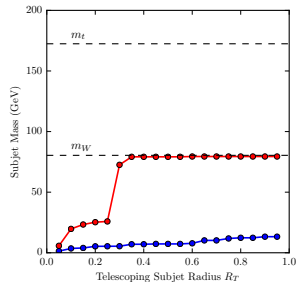
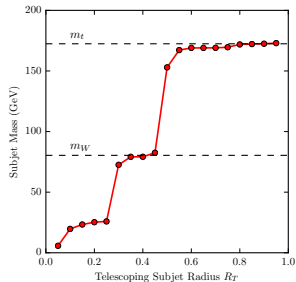
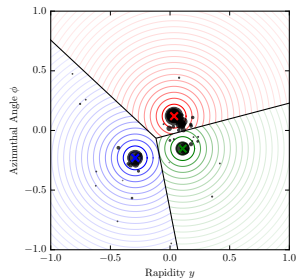
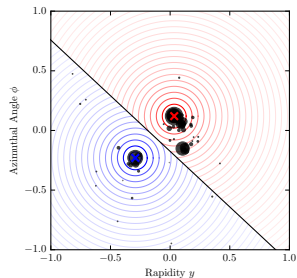
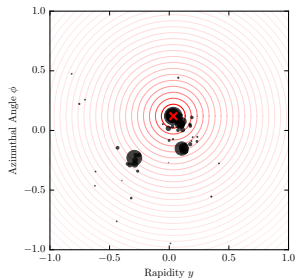


Outline

- ▶ Telescoping deconstruction
 - ▶ boosted W , top, quark v.s. gluon tagging
 - ▶ collinear QCD splitting
 - ▶ heavy ion jet modification
- ▶ Collinear drop
 - ▶ designed for soft physics studies
 - ▶ color singlet jet isolation
- ▶ Conclusions

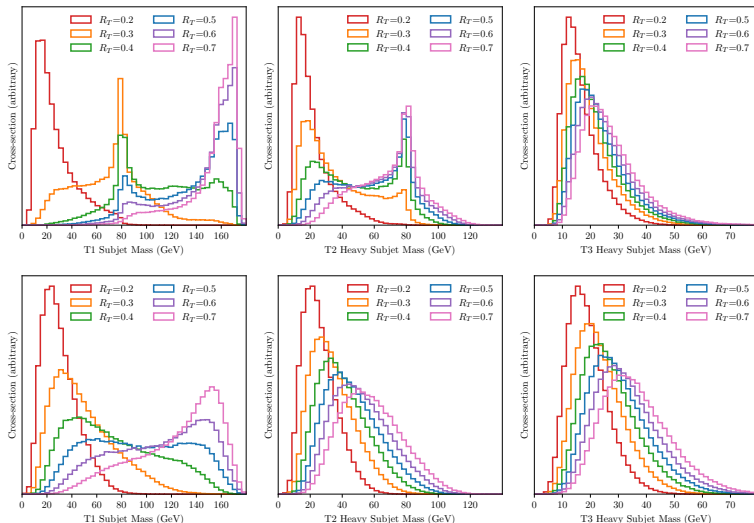
Subject masses of a top jet, resonance finder

arXiv:1803.03589

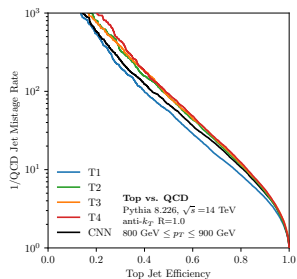
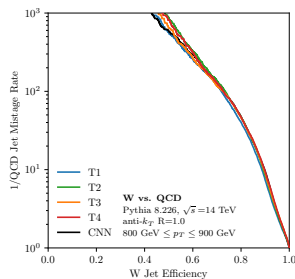
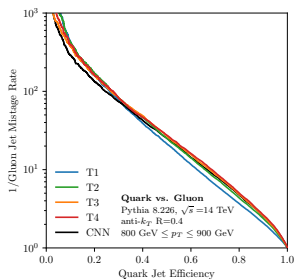


Heavy subjet mass distribution

► Top (top row) v.s. QCD (bottom row)

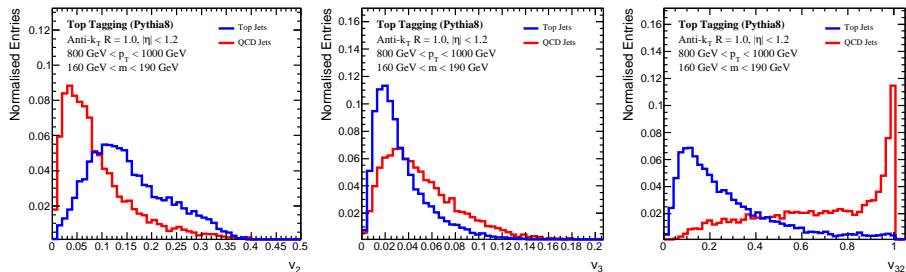


Good performances in multivariate analysis



- ▶ TN performances converge quickly to optimal performance
- ▶ TD faithfully represents jet information

Variability

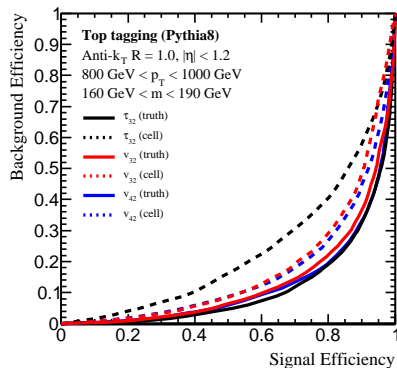
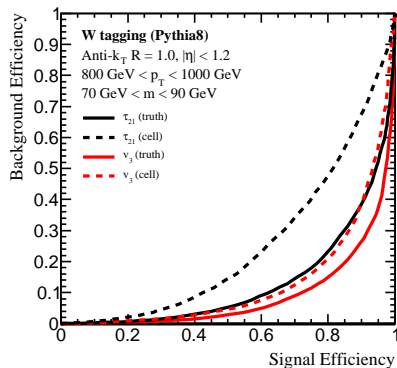


- ▶ The variation of the invariant masses $m_N(R_T)$ of the sum of N subjects can be quantified by the *variability*, motivated by the Qjet volatility (Ellis et al).

$$v_N = \frac{\sigma(\{m_N(R_T)\})}{\langle \{m_N(R_T)\} \rangle}$$

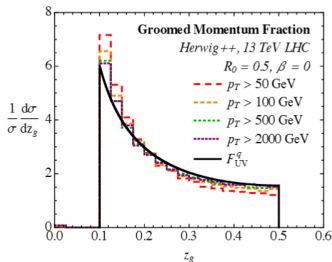
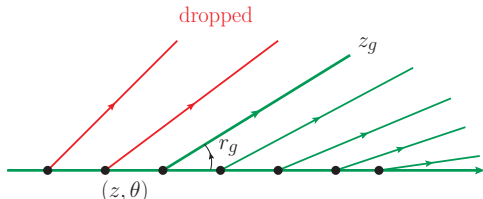
- ▶ In the case of top tagging, we also consider $v_{N2} = v_N/v_2$ for $N = 3, 4$ to optimize the variability performance, motivated by the N -subjettiness (Thaler et al) and energy correlation function (Larkoski et al) ratios.

Variability in W and top tagging



- ▶ ν_3 and ν_{42} have good performance and are robust against substructure smearing, which is meant to destroy the N -prong structure
- ▶ Manifestation of W jet isolation due to its color-singlet nature
 - ▶ A dominant feature over the two-prong structure
 - ▶ Net color charge: leading-order contribution in multipole expansion

QCD subjet distribution in soft drop

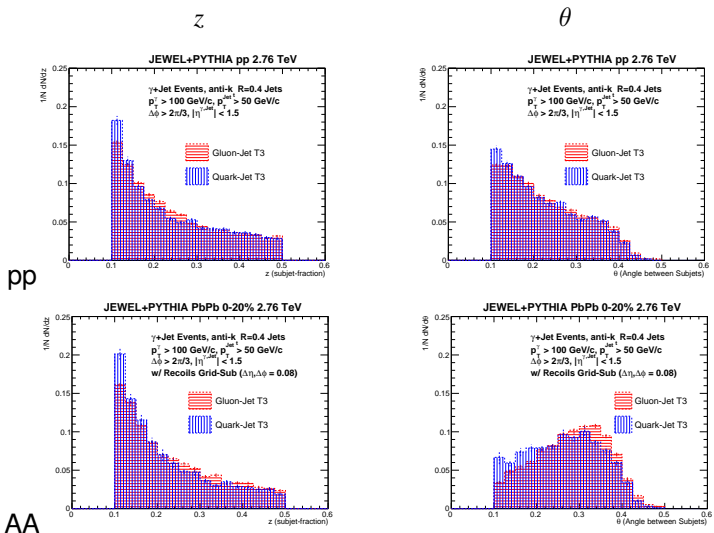


- ▶ Tree-based procedure to drop soft radiation (Larkoski et al)
 - ▶ Recluster a jet using C/A algorithm: angular ordered
 - ▶ For each branching, consider the p_T of each branch and the angle θ
 - ▶ Drop the soft branch if $z < z_{cut} \theta^\beta$, where $z = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$
- ▶ z_g : the momentum fraction of the soft branch. r_g : the angle between the branches

Telescoping subjet topology (z, θ)

arXiv:1803.03589

- Enhancement of soft, wide angle radiation in AA (simulated with Jewel (Zapp et al))

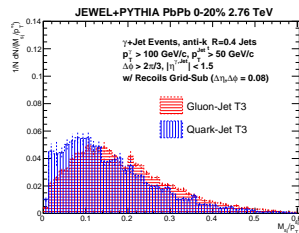
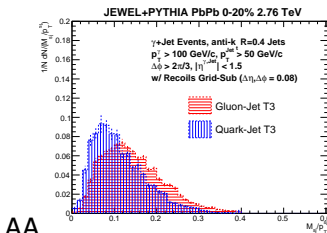
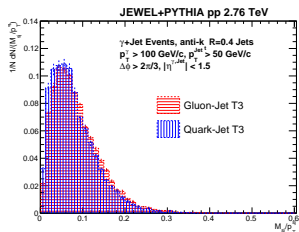
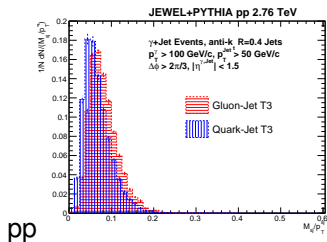


Telescoping subjet substructure m

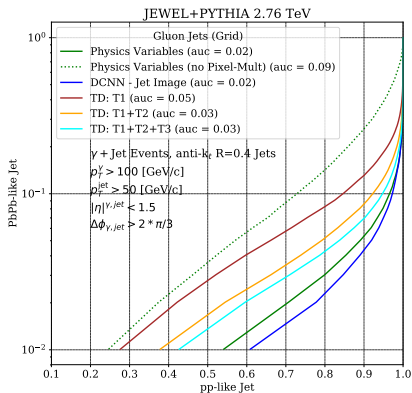
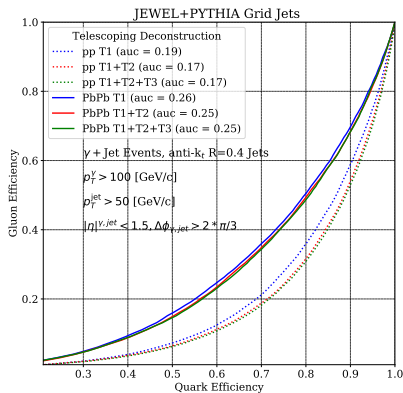
- Reveal subjet flavor dependence in first splitting $q \rightarrow qg$ and $g \rightarrow gg$ using $m^{\text{sub}}/p_T^{\text{sub}}$

Hard Subjet

Soft Subjet



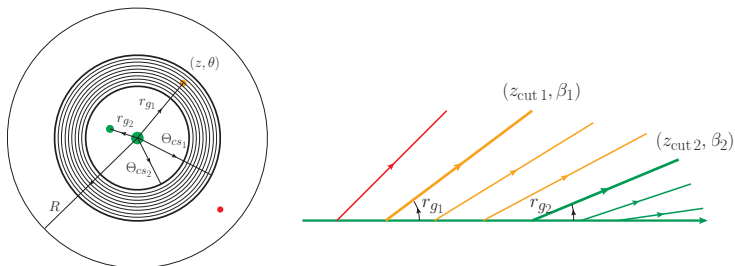
Quark gluon jet classification in pp and AA



- ▶ Quark/gluon discrimination performance drops in JEWEL-simulated AA collisions
- ▶ Information contained in subleading subjects is washed out
- ▶ Pixel multiplicity is the dominant feature distinguishing pp and AA jets

Collinear drop

One example of a CD observable can be constructed using a Soft Drop + an Anti Soft Drop



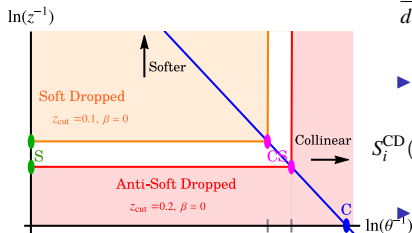
- ▶ Consider $\delta m^2 = m_{SD_1}^2 - m_{SD_2}^2$ to quantify the radiation distribution within the “ring”
 - ▶ TD version: $\delta m^2 = m^2(r_2) - m^2(r_1)$ where $r_2 > r_1$ (BOOST2015)
 - ▶ Difference of angularities, event shapes are not CD observables
- ▶ Contributions from collinear radiation to soft-drop jet mass cancel

$$m_{SD_i}^2 \sim p_c^2 + 2E_J n \cdot p_{cs_i}, \quad \delta m^2 \sim 2E_J n \cdot (p_{cs_1} - p_{cs_2})$$

- ▶ Phase space constraints on the kinematics of soft emissions,

$$z\theta^2 \approx \frac{\Delta m^2}{E_J^2}, \quad z_{cut 1} \left(\frac{\theta}{R_0} \right)^{\beta_1} \lesssim z \lesssim z_{cut 2} \left(\frac{\theta}{R_0} \right)^{\beta_2}$$

Factorization and resummation of δm^2 using SCET

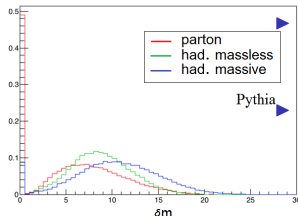
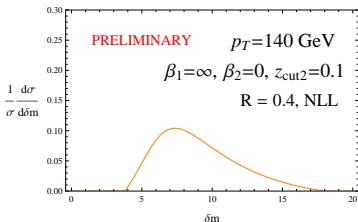


$$\frac{d\sigma}{d\delta m^2} = \sum_{i=q,g} N_i(\mu) J_{\text{un},i}^{\text{SD}}(z_{\text{cut}2}, \beta_2, \mu) S_i^{\text{CD}}(\delta m^2, z_{\text{cut}i}, \beta_i, \mu)$$

- ▶ If two soft-drop conditions are hierarchically separated, collinear-soft sector can be further factorized

$$S_i^{\text{CD}}(\delta m^2, \mu) = \int dk_i \overline{S_{C_2,i}}(k_2, \mu) S_{C_1,i}(k_1, \mu) \delta(\delta m^2 - 2E_J(k_1 + k_2))$$

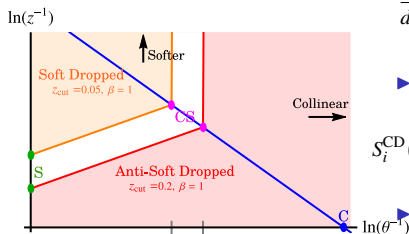
- ▶ Using SCET renormalization group techniques we can resum δm^2



▶ Massless hadrons are constructed using the p -scheme

▶ Significant hadronization correction is seen. Addressed in future work.

Factorization and resummation of δm^2 using SCET

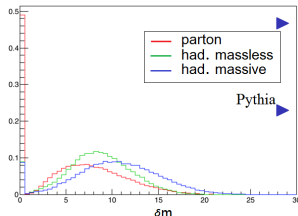
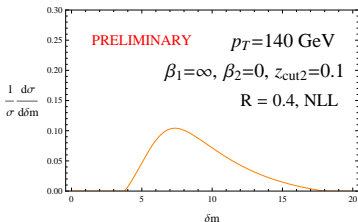


$$\frac{d\sigma}{d\delta m^2} = \sum_{i=q,g} N_i(\mu) J_{\text{un},i}^{\text{SD}}(z_{\text{cut}2}, \beta_2, \mu) S_i^{\text{CD}}(\delta m^2, z_{\text{cut}i}, \beta_i, \mu)$$

- ▶ If two soft-drop conditions are hierarchically separated, collinear-soft sector can be further factorized

$$S_i^{\text{CD}}(\delta m^2, \mu) = \int dk_i \overline{S_{C_2,i}}(k_2, \mu) S_{C_1,i}(k_1, \mu) \delta(\delta m^2 - 2E_J(k_1 + k_2))$$

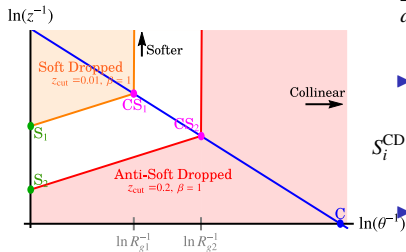
- ▶ Using SCET renormalization group techniques we can resum δm^2



Massless hadrons are constructed using the p -scheme

Significant hadronization correction is seen. Addressed in future work.

Factorization and resummation of δm^2 using SCET

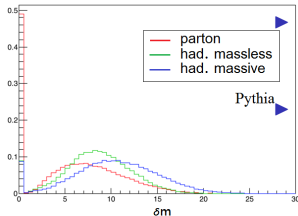
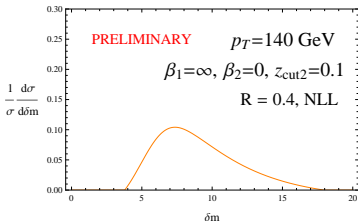


$$\frac{d\sigma}{d\delta m^2} = \sum_{i=q,g} N_i(\mu) J_{\text{un},i}^{\text{SD}}(z_{\text{cut}2}, \beta_2, \mu) S_i^{\text{CD}}(\delta m^2, z_{\text{cut}1}, \beta_1, \mu)$$

- ▶ If two soft-drop conditions are hierarchically separated, collinear-soft sector can be further factorized

$$S_i^{\text{CD}}(\delta m^2, \mu) = \int dk_i \overline{S_{C_2,i}}(k_2, \mu) S_{C_1,i}(k_1, \mu) \delta(\delta m^2 - 2E_J(k_1 + k_2))$$

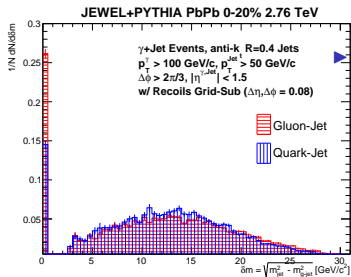
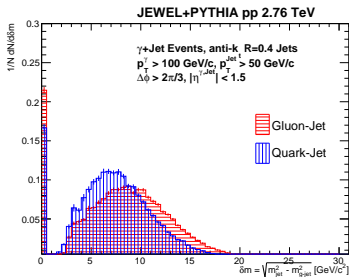
- ▶ Using SCET renormalization group techniques we can resum δm^2



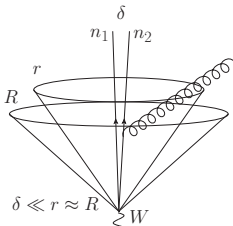
▶ Massless hadrons are constructed using the p -scheme

▶ Significant hadronization correction is seen. Addressed in future work.

Collinear-drop in heavy ion and W jet isolation



▶ Quark/gluon jet difference disappearing in JEWEL-simulated AA collisions



- ▶ In the boosted regime, $\delta \approx m_W/p_T \ll 1$
- ▶ At $\mathcal{O}(\alpha_s)$,

$$\frac{d\sigma}{d\delta m^2} = \delta^2 \frac{\alpha_s C_F}{\pi \delta m^2} \frac{R^2 - r^2}{R^2 r^2}, \text{ power suppressed by } \delta^2$$

- ▶ Key difference between W and QCD jets: net color charge (also seen in Jet Lund plane (F. Dreyer))

Conclusions

- ▶ Telescoping deconstruction is a systematic framework for jet studies
 - ▶ Can be easily extended to event-level studies
- ▶ Collinear-drop class of observables allows one to directly probe soft physics in Jets
 - ▶ studying hadronization in Jets and testing MC
 - ▶ studying the effect of the jet on the soft medium in heavy ion
 - ▶ discrimination studies
- ▶ Heavy ion jet modification can be comprehensively studied as a classification problem
- ▶ W jet isolation is identified as a dominant feature of boosted W jets