



Tile Dual Readout and Beyond

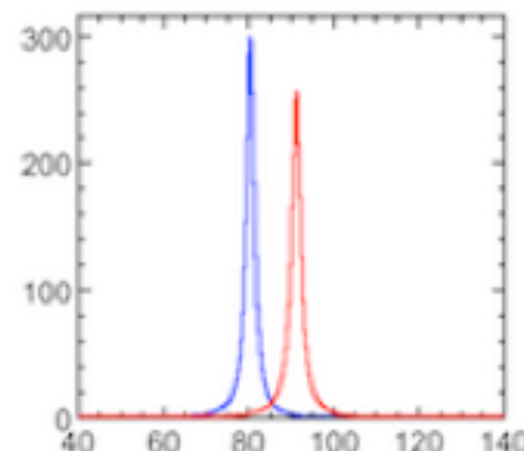
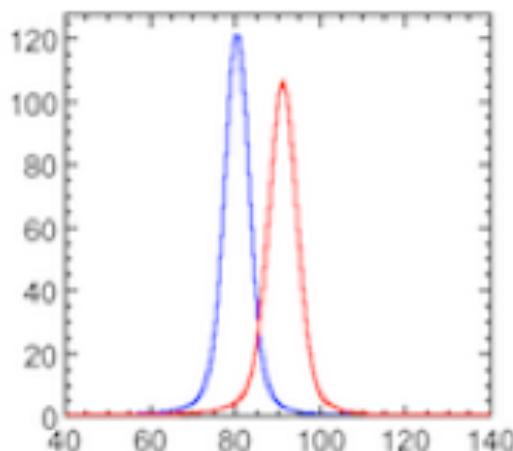
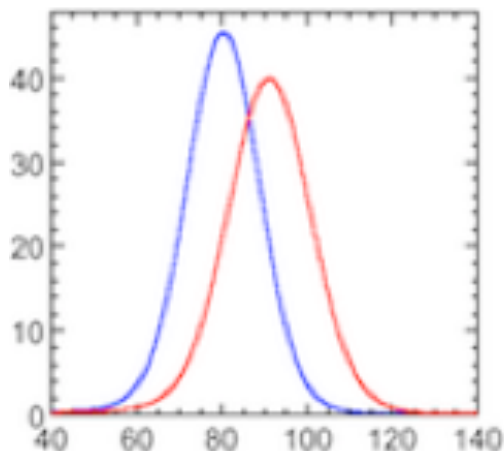
David R Winn, Fairfield/CMS

Yasar Onel, Iowa/CMS

(*Achtung! Disclaimer: Not a CMS Presentation!*)

One Physics Example: Precision jets

- - **Jet-Jet masses:** Goal for future experiments: SM $Z \rightarrow \text{jetjet}$; $W \rightarrow \text{jetjet}$
- Ratio $W, Z \rightarrow jj$ to $W, Z \rightarrow \text{leptons}$ $\sim 6-7$
- Reconstruct AND Separate(+SM, E_{Tmiss} , jet tags, V-V scattering, BSM, W' , Z' ...)
- Separation of W from Z: $\sigma_{E_{\text{jet}}}/E_{\text{jet}} \sim 3\%$ necessary at 100 GeV, with typical single particle energies ~ 10 GeV [*ASIDE: during collision crossing times which may be as small as ~ 10 's of ns, pileup events $\sim 200/\text{crossing}$ and raddam exceeding 50-100MRad.] A 3%-4% jet energy resolution from 50-500 GeV gives $2.6-2.3\sigma$ W/Z separation. J-J mass resolution is very important In searches for heavy W'/Z' , vector boson scattering, triple VVV....*



- W/Z- \rightarrow jet-jet separation: **Left** - calorimeter $\sigma_E/E = 60\%/ \sqrt{E}$; **Middle** $\sigma_E/E = 22\%/ \sqrt{E}$ (3% @ 50 GeV) $\sim 2.6\sigma$ separation; **Right** -perfect resolution: $\sim 4.5\sigma$ separation.

Dual Readout: Cerenkov Compensation

- 1st Quantitative MC Study:
 - GEANT MC in 1988 [*Compensating Hadron Calorimeters with Cerenkov Light*, D.R. Winn, W.Worstell, *IEEE Trans. Nuc. Sci.*, V1 NS-36, 334(1989)]
 - Idea: Use differences in response to e-m fluctuations between
 - Cerenkov Medium (transparent – LAr, H₂O, SiO₂) vs
 - Ionization Medium (scintillator, LAr ion collection,..)
 to reduce hadron shower fluctuations and make $e/h \rightarrow 1$
- DREAM Collaboration/Richard Wigmans et al.
 - Excellent progress in real tests!
 - Thorough Analysis of Dual Parallel Fiber Calorimeters!
- MC: 18%/√E seems possible.
- DREAM: 30%/√E: Parallel Scintillating+Q Fibers

Parallel Fiber Dual Deficits - 1

- **1. Constant Term – unavoidable issue** – scintillator light attenuation $\sim 2+$ m of fiber.
- **2. Pointing/Projective Geometry problematic** in a practical parallel fiber calorimeter over a substantial solid angle. The mechanics + fiber packing of fully projective (θ, ϕ) very difficult for (pitch, yaw) more than $\sim 5^\circ$. Streaming down fiber holes lowered the resolution in DREAM, even at a 2° pitch. Packing extra fibers from the back or conical fibers: \rightarrow Constant term, \rightarrow Calibration Issues
- **3. Scintillator Fiber & Photodetector Raddam:** At present, there are no good examples of scintillator fibers which have proven sufficient raddam resistance or speed to be useful for hadron calorimetry at most future colliders.
- **4. Fiber Bundle & Photodetector Punchthrough:** Huge fiber bundles, $>33\%$ of the back of the fiber dual calorimeter area, are directly behind the calorimeter. Large punchthrough backgrounds are generated by these fibers, photodetectors ($\sim 1/800$ incident π/K quasi-elastic scatter through a 10 Lint calorimeter).



- **4. E-M and Hadronic Components of Incident Jets:** Parallel fibers ~no ability to detect + separate incident *direct e-m component inside of a jet, since there is no longitudinal segmentation.*
- **5. High Resolution EM Front End.** The parallel fiber dual readout jet calo ~no ability to make a *compensated* high-Z high sampling EM front end.
- **6. Calibration:** Parallel fiber geometry difficult to calibrate, as radiation damage & attenuation varies w/ length. (Contrast w/ longitudinally segmented calorimeters)
- **7. Timing & Pileup:** Longitudinal fibers store the information of jet/em showers: the signal is over the time for the light to traverse the fibers. The light generated *at the back of the calorimeter arrives at the photodetector first.* Fiber calorimeters measure the falling edge of the shower, a less precise measurement
- **8. Longitudinal Segmentation:** Fiber dual readout is incompatible with true longitudinal segmentation even with waveform electronics, and cannot be easily rebuilt for front raddam or implement 4,5 above.

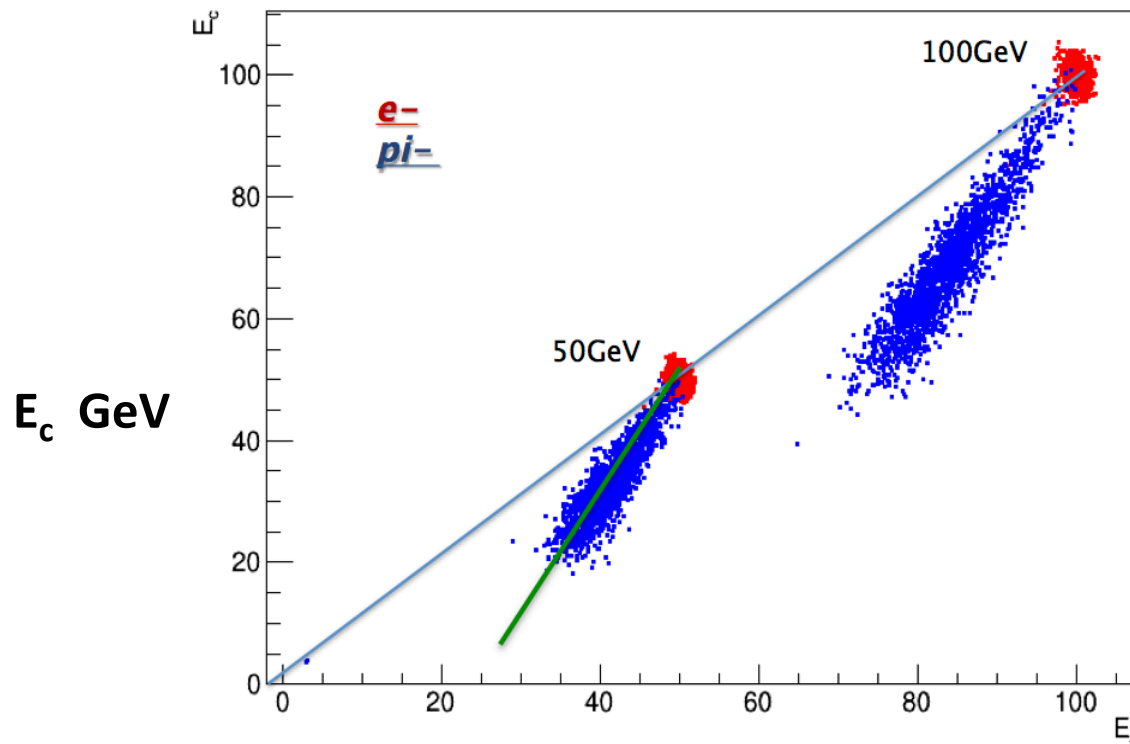
- **19. Cerenkov (Fiber) Index of Refraction:** High Radiation Resistant Cerenkov fibers limited to quartz, with $n=1.46 - h/e_c \sim 0.25-0.20$ – limiting resolution. Low n tiles: even lower h/e_c ratio (aerogels, Teflon AF, Siloxanes, Fluoride glasses,...)
- **10. Particle Flow/Energy Flow Calorimetry:** improving jet $\delta\theta/\theta$, core ID of jets, isolation/ID of leptons/photons in jets and pileup, and neutral particle (K^0 , n) ID, all especially under pileup. *Tile* dual readout is fully compatible, and can be easily added to particle flow calorimeters
- **11. Other Sensors/Dual, Triple:** Parallel fibers especially cannot readily use other hadronic or n -enhanced sensors, and e-m energy sensors. Examples: ionization detectors (solid – Si, Diamond, GaAs; liquid- LArgon; gasses – micromegas..); $\beta \rightarrow 1$ sensitive detectors such TRD, or ultra-low-index materials(aerogels, MgF_2 , water, perfluoros, silicones,..); secondary emission sensors with higher response to slow particles $\beta \rightarrow 0$ and minimal response to minimum ionizing energy (new large MCP); inorganic non-hydrogenous scintillators (LYSO, $PbWO_4$ et al.), and 6Li , ^{10}B or 3He containing materials.
- **12. Cost:** the cost of tiles is significantly less per mass or volume of sensitive material than that of fibers, and the cost of a fabricated tile absorber matrix is considerably less than the parallel fiber Swiss cheese.



MC Study: Tile Dual Readout



- GEANT4 MC on a simple tile calorimeter: 0.5 cm thick each of quartz, plastic scintillator, and Cu absorber tiles.
- Two energies (50, 100 GeV) each of 1000 electrons (red dots) and of ~800-1000 pions (blue dots)] were sent into the 50x50 cm calorimeter, 12 Lint deep.
- Nphotons 325-650nm generated in the Cerenkov and in the scintillator tiles were counted. 0.5% at random were assigned as converted to p.e.
- Scintillator photons ~120x Cerenkov photons; photostatistics not limiting factors.
- Means of histograms of the electron shower p.e. in quartz and in scintillator were used to convert/normalize the number of collected p.e. in Cerenkov light and in Scintillator light to normaized energies E_{Cerenkov} and $E_{\text{Scintillator}}$, and then plotted as a scatter plot of E_c vs E_s for each electron.
- Pions of 50, 100 GeV were then simulated, converted to E_c vs E_s



- Scatter plots:

E_s GeV

Electrons: E_c vs E_s (red) lie along line shown schematically as $E_c = E_s$.

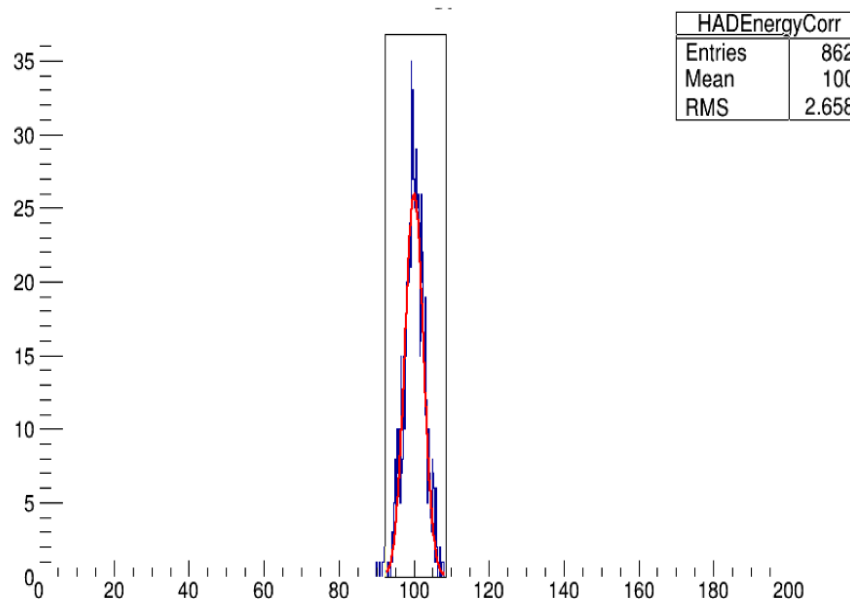
Pions: E_c vs E_s scatter-plotted (blue) lie mainly below the $E_c = E_s$ electron line with correlation between E_c vs E_s fitted as a line (green, shown only 50 GeV points).

- As the shower fluctuates more to hadrons, E_c falls faster than E_s .

- A Simple analysis: Linear fit to hadron scatter points (Green line), with slope R , corrects the energy: Project the scatter points as a histogram perpendicular to the linear correlation, the energy distribution becomes Gaussian & narrower.

Dual Correction

- Pion Energy E (*first order*): $E = E_s + [\text{a correction term proportional to the difference } (E_s - E_c)]$
 $E = E_s + \alpha(E_s - E_c)$ with α given by slope R as $R = (1 + \alpha)/\alpha$ or $\alpha = 1/(1 - R)$.
 - The angle between the line $E_c = E_s$ and fitted π scatter plot line: $\theta = \arctan(R) - \pi/4$.
- $(E_s - E_c)$ grows as shower fluctuates into nuclear/hadronic energies.
- As slope R gets steeper, the correction term $\alpha(E_s - E_c)$ becomes more important. When Cerenkov E_c is the same as scintillation E_s (e's or π 's exchange to π^0 's), then $(E_s - E_c) \sim 0$, $E = E_s$.



(mean, rms) = (100, 2.66) GeV energy resolution that could enable $W \rightarrow \text{jet-jet}$; $Z \rightarrow \text{jet-jet}$

- **Higher order terms**- $\alpha_2(E_s - E_c)^2 + \alpha_3(E_s - E_c)^3 + \dots$ and energy dependent α_n – there is a continuous mapping(vector field) of the points in E_c vs E_s space to the line $E_c = E_s = E$.



Tile Dual Summary/Discussion

– Rules of Thumb:

- (0) An intrinsic limit of normal hadron calorimetry: $\sigma_E/E > 11\text{-}13\%/ \sqrt{E}$, given by the ratio of detectable neutron energy to the fluctuations in lost nuclear binding energy.
- (1) Contrast between h_i/e_i (i =ionization) and h_c/e_c (C =Cerenkov) for hadrons h and e-m energy: the ratio of ratios $[h_i/e_i]/[h_c/e_c] \geq 4$ in order to reach incident hadron energy resolutions below $30\%/ \sqrt{E}$, with $18\%/ \sqrt{E}$ being a reasonable target to achieve using plastic scintillator and low index materials;
- (2) h_i/e_i : as large as possible \rightarrow hydrogenous or n-sensitized ionization detection media.
- (3) e-m energy resolution in Cerenkov light $< 70\%/ \sqrt{E}$ to achieve $< 20\%/ \sqrt{E}$;
- (4) Resolution scales $\sim \sqrt{f_{\text{sample}}/f_{\text{frequency}}}$.
- (5) Compensation can be achieved by enhancing neutron(hydrogenous or n-absorbing) or ion fragment sensitivity and/or by suppression electromagnetic component by tuning the absorber thickness relative to sampling media (f_{sample} typically $\sim 1/10$ but at a loss of potential ultimate resolution).

- **Adding sensor tiles** relatively insensitive to MIPs, OR more sensitive to $\gamma\beta \rightarrow 0$ increases the contrast between e-m and hadronic energy (enhancing the low energy hadronic signal) – one such sensor is Secondary Emission; its signal scales as dE/dx , with a MIP SE signal $\sim 100x$ less than that of the energy of the peak signal (peak signal for protons occurs at $\sim 200\text{KeV}$ - $n+p \rightarrow p+n$ knock-on protons).
- **Homogeneous non-hydrogenous dense inorganic scintillators** (LYSO, PbWO_4 , CeF_3)
 - $h_i/e_i \sim 0.4$ and $h_c/e_c \sim 0.25$, or $[h_i/e_i]/[h_c/e_c] \sim 1.6$:
 -> *Homogeneous calorimeters cannot achieve dual readout compensation better than $\sim 50\text{-}60\%/ \sqrt{E}$ on hadrons, even with perfect separation between scintillator & Cerenkov light in the homogeneous detector. [Note: LAr/Ch4]*

Theoretical $\sim 15\text{-}18\%/ \sqrt{E}$ on jets: *scintillator sensors with $h_i/e_i \sim 0.6\text{-}0.8$ (likely hydrogenous & n-sensitive), and Cerenkov sensors with $h_c/e_c \leq 0.2$ are needed. To achieve $h_c/e_c < 0.2$, lower n (index of refraction) Cerenkov radiators are required (i.e. $\beta_{\text{thresh}} \rightarrow 1$), but require enough photons to achieve an e-m resolution $< 70\%/ \sqrt{E}(\text{GeV})$ or $N_{pe} > 2 \text{ pe/GeV}$.*

Beyond/Extended Dual Readout -1

Extend E-M Response by higher sensitivity to $\beta > 1$
Results in high contrast ratio with $h_c/e_c > 0.15$ (i.e. $e_c/h_c > 6-7$)
(Lessens low energy Hadron, n, and nuclear fragment Sensitivity)

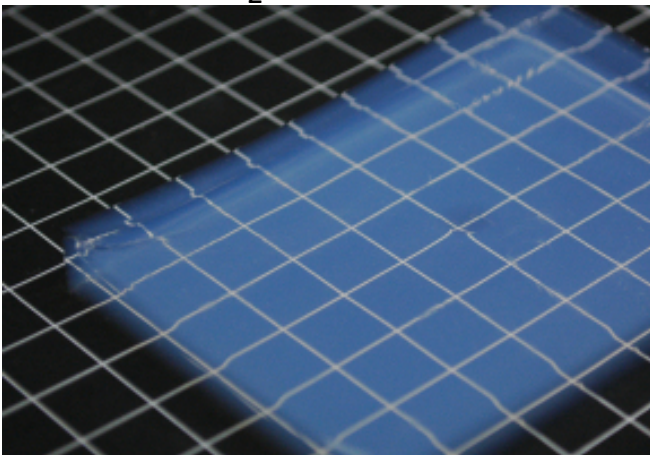
(A) TRD

- Straw tubes,..... Low mass issue for calorimetry)

(B) Low index Tiles

($1.1 < n < 1.35$) tiles :

- silica aerogels ($n=1.05-1.3$)
- TeflonAF ($n=1.29$, 12 Mrad) (amorphous form; water-clear)
- polysiloxanes ($n=1.35$, 100 Mrad)
- MgF_2 (1.37);

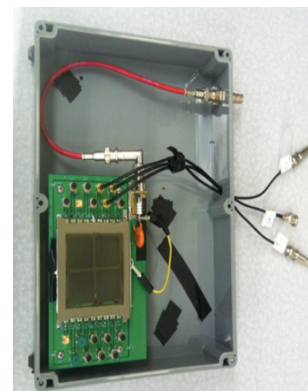
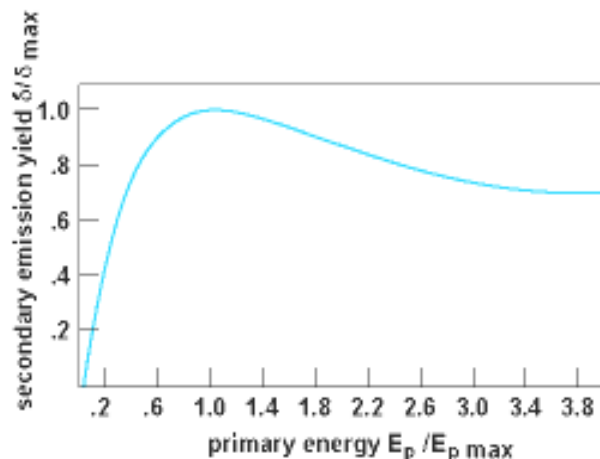
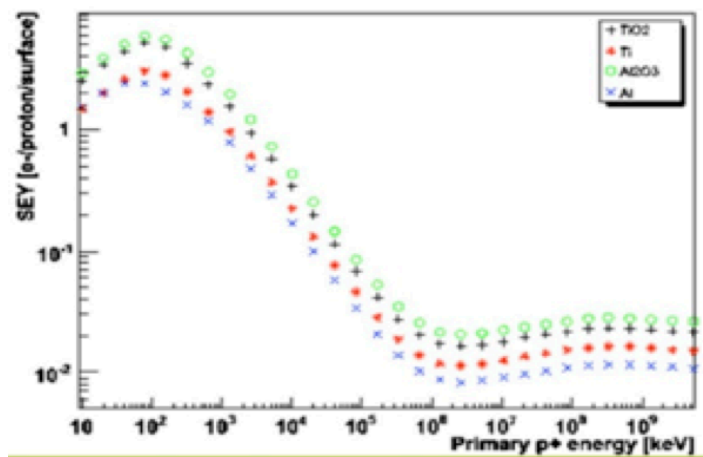


Beyond/Extended Dual Readout -2

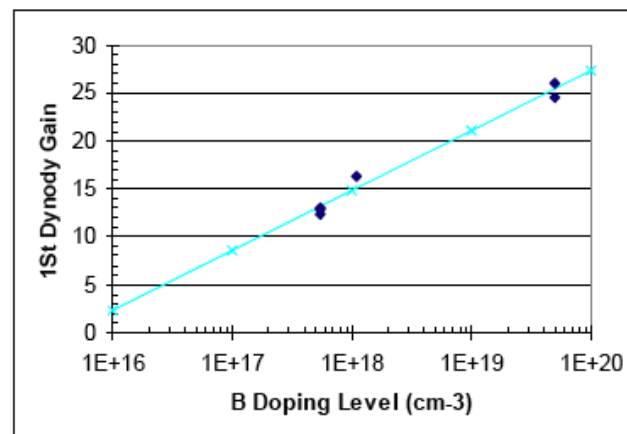
- **Secondary Emission (SE).** Secondary Emission(SE) tiles are more sensitive to $\gamma\beta \rightarrow 0$ particles than to MIPs - scales as dE/dx . MIP SE signal ~ 100 - $200\times$ less than at peak $\gamma\beta$ SE signal – **the opposite of Cerenkov light**. SE tiles for correction for heavy fragments, lost neutron energies, slow hadrons. (triple/quadruple readout.) 1-2 MeV alpha particle: secondary yield \sim at max.

LHC SE
Monitors:
SE Yield vs
 E_p (KeV)

10^{20} p/cm² !



12x12 MCP
FNAL Test Beam



B-doped
Nanoxtal
Diamond
1 μ m thick
On dynode

- **Triple Readout and beyond:** 3 tiles to improve dual readout: non-hydrogenous scintillator, hydrogenous/neutron-sensitive scintillator, 2 indices of Cerenkov tile(s), SE tiles.... Compare less-sensitive neutron scintillators [non-hydrogenous scintillators; inorganic and perfluorocompounds] to more neutron-sensitive H or n-absorbing/converting scintillator tiles.
- **Combined Dual/Triple Readout with Particle Flow:** Add Cerenkov tiles, TRD tiles, SE tiles, or others to existing Particle/Energy flow calorimeter prototypes.
- **Neutron-enhanced detecting scintillator tiles** thin film coatings - ^{10}B , ^6Li , hydrogenous materials [^6LiH] – thin clear film, buffered w/ alumina films; interesting: $\text{Li}^6\text{B}^{10}\text{H}_4$ which would be transparent if deposited as thin films between clear buffers. ^{10}B SE yield dynodes.
- **Liquids:** very large homogeneous detectors: LB, cosmic neutrinos or prot-rot.
 - 1) water “tiles”(n=1.29-1.31 TeflonAF light pipe) + LS tiles – no absorber
 - 2) LArgon drifted ions + Cerenkov light detection. The index n good e/h contrast; scintillation light at 128nm will not penetrate PMT windows.