ttbb 4F with Powheg+OpenLoops

Tomáš Ježo

University of Zürich

In collaboration with:

J. Lindert, N. Moretti and S. Pozzorini

Meeting of the LHC Higgs WG1 ttH/tH subgroup July 5th



Universität Zürich^{UZH}

Outline and goal



- We perform NLO+PS matching of $pp \rightarrow t\bar{t}b\bar{b}$ with the POWHEG method
 - Using 4 flavour number scheme
 - ▶ *b*-quark mass effects included
 - \blacktriangleright makes the full $t\bar{t}b$ -jet phase space available
 - ► Employing the brand new POWHEG BOX RES

[T.J., Nason 2015]

- allows the choice of the most suitable FKS mapping
- Here I present a progress update
 - ► Comaprison against Sherpa+OpenLoops

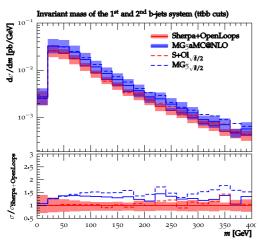
[Cascioli et al. 2013]

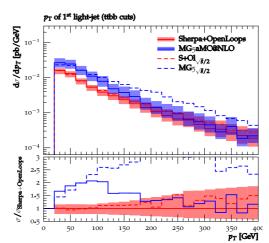
► A study of hdamp dependence

Motivations



Dependence on resummation scale μ_Q





Nominal MG5_aMC and Sherpa+OpenLoops predictions in YR4

• MG5_aMC supports only* $\mu_Q = f(\xi)\sqrt{\hat{s}} \Rightarrow$ smearing function restricted to $0.1 < f(\xi) < 0.25$ to mimic recommended $\mu_Q = H_T/2$ implemented in Sherpa

μ_Q variations enhance the discrepancy

- $\mu_Q = \sqrt{\hat{s}}/2$ in Sherpa to mimic MG5_aMC default choice $0.1 < f(\xi) < 1$
- strong μ_Q -sensitivity of MG5_aMC \Rightarrow much more pronounced deviations
- * Ongoing studies with new MG5 version supporting $\mu_Q=H_T/2$. See talk by C. Neu.

Stefano's talk





• POWHEG BOX RES

[T.J., Nason 2015]

- ▶ Implements "resonance aware" NLO+PS matching providing the means of
 - choosing a FKS mapping appropriate to the matrix element
 - separating the calculation into regions in which the above is possible
- ► Features a new interface with OpenLoops at a very advanced level of automation [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- Processes matched to PS using POWHEG BOX RES:
 - ▶ $pp \rightarrow \mu^+ \nu_{\mu} j_b j$

 - ▶ HV and HV+jet (QCD+EW)
- Improvements towards 4F $t\bar{t}b\bar{b}$ @ NLO+PS in POWHEG BOX RES[†]:
 - ▶ hdamp for massive final state emitters
 - ▶ Partial unweighting, Rivet interface, ...



• POWHEG BOX RES

- [T.J., Nason 2015]
- ▶ Implements "resonance aware" NLO+PS matching providing the means of
 - choosing a resonance virtuality preserving FKS mapping
 - separating the calculation into regions dominated by a single resonance history
- ► Features a new interface with OpenLoops at a very advanced level of automation [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- Processes matched to PS using POWHEG BOX RES:
 - ▶ $pp \rightarrow \mu^+ \nu_\mu j_b j$ @ NLO QCD (t-channel single top)
 - ▶ $pp \rightarrow \ell^+ \nu_{\ell} l^- \bar{\nu}_{t} b \bar{b}$ @ NLO QCD (tt & Wt)
 - ▶ HV and HV+jet @ NLO QCD+EW
- Improvements towards 4F $t\bar{t}b\bar{b}$ @ NLO+PS in POWHEG BOX RES[†]:
 - ▶ hdamp for massive final state emitters
 - ▶ Partial unweighting, Rivet interface, ...



• POWHEG BOX RES

[T.J., Nason 2015]

- ▶ Implements "resonance aware" NLO+PS matching providing the means of
 - choosing a FKS mapping appropriate to the matrix element
 - separating the calculation into regions in which the above is possible
- ► Features a new interface with OpenLoops at a very advanced level of automation [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- Processes matched to PS using POWHEG BOX RES:
 - ▶ $pp \rightarrow \mu^+ \nu_\mu j_b j$ @ NLO QCD (t-channel single top)
 - ▶ $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_\ell b \bar{b}$ @ NLO QCD (tt & Wt)
 - ► HV and HV+jet @ NLO QCD+EW
- Improvements towards 4F $t\bar{t}b\bar{b}$ @ NLO+PS in POWHEG BOX RES[†]:
 - ▶ hdamp for massive final state emitters
 - ▶ Partial unweighting, Rivet interface, ...



• POWHEG BOX RES

- [T.J., Nason 2015]
- ▶ Implements "resonance aware" NLO+PS matching providing the means of
 - choosing a FKS mapping appropriate to the matrix element
 - separating the calculation into regions in which the above is possible
- ▶ Features a new interface with OpenLoops at a very advanced level of automation [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- Processes matched to PS using POWHEG BOX RES:
 - ▶ $pp \rightarrow \mu^+ \nu_{\mu} j_b j$

[T.J., Nason 2015]

- ▶ $pp \rightarrow \ell^+ v_\ell l^- \bar{v}_\iota b \bar{b}$ [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- ▶ HV and HV+jet (QCD+EW) [Granata, Lindert, Oleari, Pozzorini 2017]
- Improvements towards 4F $t\bar{t}bb$ @ NLO+PS in POWHEG BOX RES[†]:
 - ▶ hdamp for massive final state emitters
 - ▶ Partial unweighting, Rivet interface, ...



• POWHEG BOX RES

- [T.J., Nason 2015]
- ▶ Implements "resonance aware" NLO+PS matching providing the means of
 - choosing a FKS mapping appropriate to the matrix element
 - separating the calculation into regions in which the above is possible
- ▶ Features a new interface with OpenLoops at a very advanced level of automation [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- Processes matched to PS using POWHEG BOX RES:

- [T.J., Nason 2015]
- ▶ $pp \rightarrow \ell^+ v_\ell l^- \bar{v}_\ell b \bar{b}$ [T.J., Lindert, Nason, Oleari, Pozzorini 2016]
- ► HV and HV+jet (QCD+EW) [Granata, Lindert, Oleari, Pozzorini 2017]
- Improvements towards 4F $t\bar{t}b\bar{b}$ @ NLO+PS in POWHEG BOX RES[†]:
 - ▶ hdamp for massive final state emitters
 - ▶ Partial unweighting, Rivet interface, ...

† RES features not used in the indended way

Setup



- LHC 13 TeV
- 4F scheme
- NNPDF30_nlo_as_0118_nf_4
- α_S from 4F PDFs
- $m_t = 172.5 \text{ GeV}$, $m_b = 4.75 \text{ GeV}$

•
$$\mu_r = \mu_f = \sqrt{0.5\sqrt{E_T^b E_T^{\bar{b}}}\sqrt{E_T^t E_T^{\bar{t}}}}$$

- Top decay off, MPI off, Hadronization off, QED shower off
- Sherpa
 - $\mu_Q = H_T/2$
 - Default shower settings
- POWHEG BOX RES
 - ▶ hdamp = $1.5m_t$
 - ▶ Pythia 8.2 with PowhegHooks, A14 tune

Setup



- LHC 13 TeV
- 4F scheme
- NNPDF30_nlo_as_0118_nf_4
- α_S from 4F PDFs
- $m_t = 172.5 \text{ GeV}$, $m_b = 4.75 \text{ GeV}$

•
$$\mu_r = \mu_f = \sqrt{0.5\sqrt{E_T^b E_T^{\bar{b}}}}\sqrt{E_T^t E_T^{\bar{t}}}$$

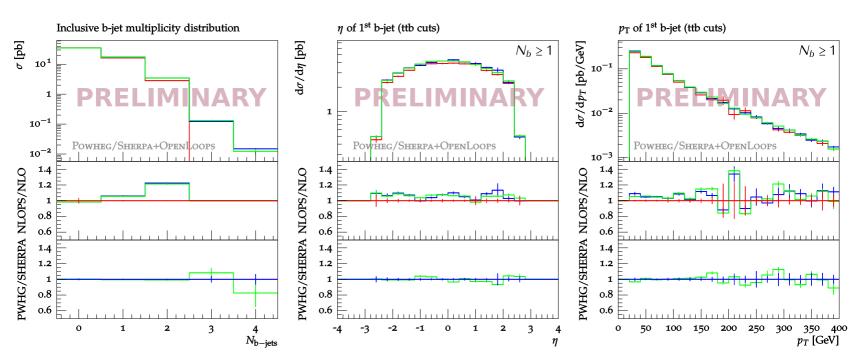
- Top decay off, MPI off, Hadronization off, QED shower off
- Sherpa
 - $\mu_Q = H_T/2$
 - Default shower settings
- POWHEG BOX RES
 - ▶ hdamp = $1.5m_t$
 - ▶ Pythia 8.2 with PowhegHooks, A14 tune

improves perturbative

convergence



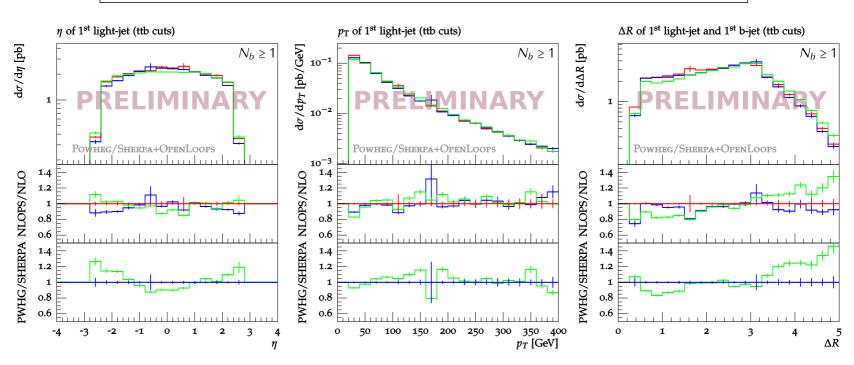




- Remarkable agreement for NLO accurate ttb observables
 - ► Agreement well under 5%; expected scale uncertainity ~20%
- Good agreement also in LOPS accurate bins 3 and 4 of $\sigma(N_{b-{\rm jets}})$



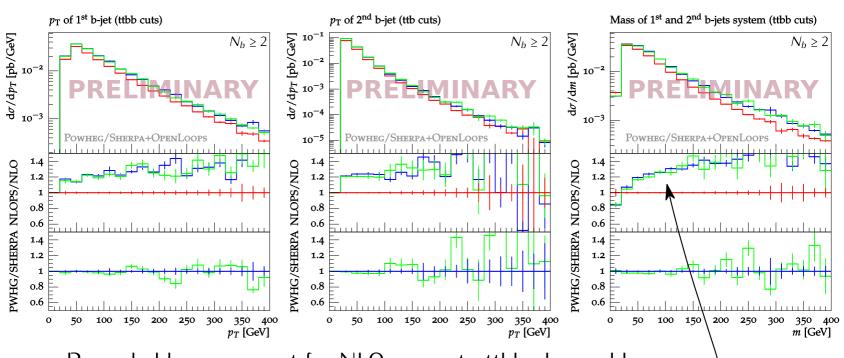
→ PWHG NLO → PWHG NLOPS → SHERPA NLOPS



- Good agreement for LOPS accurate ttbj observables
 - ► Agreement to ~20%; expected scale uncertainity ~50%



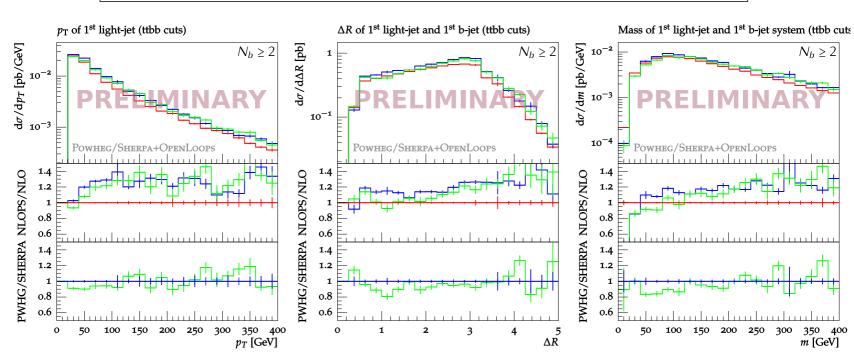




- Remarkable agreement for NLO accurate ttbb observables
 - ► Agreement well under 5%; expected scale uncertainity ~20%
- POWHEG BOX RES confirms the "double splitting" enhancement







- Good agreement for LOPS accurate ttbbj observables
 - ► Agreement to ~20%; expected scale uncertainity ~50%



$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$

- where $R_{\alpha} = R_{\alpha}^{s} + R_{\alpha}^{r}$
- Separation of the real contribution introduced to deal with "Born zeroes"
 - if (r0.gt.5*abs(rc+rs-rcs)) then ... R_{α}^{r}
 - \blacktriangleright else ... R_{α}^{s}
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^{s} = R_{\alpha}F(k_{T}^{2})$$
, $R_{\alpha}^{r} = R_{\alpha}\left[1 - F(k_{T}^{2})\right]$, $F(k_{T}^{2}) = \frac{h^{2}}{k_{T}^{2} + h^{2}}$

- ullet In top-pair production chosing hdamp $\sim m_t$ improves the description of the data
 - \blacktriangleright ATLAS tunes hdamp = 1.5 m_t , CMS sets to the same value



$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \underbrace{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}_{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$
• where $R_{\alpha} = R_{\alpha}^s + R_{\alpha}^r$ finite

- Separation of the real contribution introduced to deal with "Born zeroes"
 - if (r0.gt.5*abs(rc+rs-rcs)) then ... R_{α}^{r}
 - else ... R_{α}^{s}
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^{s} = R_{\alpha}F(k_{T}^{2})$$
, $R_{\alpha}^{r} = R_{\alpha}\left[1 - F(k_{T}^{2})\right]$, $F(k_{T}^{2}) = \frac{h^{2}}{k_{T}^{2} + h^{2}}$

- ullet In top-pair production chosing ${\tt hdamp} \sim m_t$ improves the description of the data
 - \blacktriangleright ATLAS tunes hdamp = 1.5 m_t , CMS sets to the same value



$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$

- where $R_{\alpha} = R_{\alpha}^{s} + R_{\alpha}^{r}$
- Separation of the real contribution introduced to deal with "Born zeroes"
 - if (r0.gt.5*abs(rc+rs-rcs)) then ... R_{α}^{r}
 - \blacktriangleright else ... R_{α}^{s}
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^{s} = R_{\alpha}F(k_{T}^{2})$$
, $R_{\alpha}^{r} = R_{\alpha}\left[1 - F(k_{T}^{2})\right]$, $F(k_{T}^{2}) = \frac{h^{2}}{k_{T}^{2} + h^{2}}$

- In top-pair production chosing $hdamp \sim m_t$ improves the description of the data
 - \blacktriangleright ATLAS tunes hdamp = 1.5 m_t , CMS sets to the same value



hdamp

• POWHEG radiation formula:

$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$

- where $R_{\alpha} = R_{\alpha}^{s} + R_{\alpha}^{r}$
- Separation of the real contribution introduced to deal with "Born zeroes"
 - if (r0.gt.5*abs(rc+rs-rcs)) then ... R_{α}^{r}
 - \triangleright else ... R_{α}^{s}
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^{s} = R_{\alpha}F(k_{T}^{2})$$
, $R_{\alpha}^{r} = R_{\alpha}\left[1 - F(k_{T}^{2})\right]$, $F(k_{T}^{2}) = \frac{h^{2}}{k_{T}^{2} + h^{2}}$

maybe be thought of as an analogue to μ_Q in MC@NLO

- In top-pair production chosing hdamp $\sim m_t$ improves the description of the
 - \triangleright ATLAS tunes hdamp = 1.5 m_t , CMS sets to the same value



$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$

- where $R_{\alpha} = R_{\alpha}^{s} + R_{\alpha}^{r}$
- Separation of the real contribution introduced to deal with "Born zeroes"
 - if (r0.gt.5*abs(rc+rs-rcs)) then ... R_{α}^{r}
 - \blacktriangleright else ... R_{α}^{s}
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^{s} = R_{\alpha}F(k_{T}^{2})$$
, $R_{\alpha}^{r} = R_{\alpha}\left[1 - F(k_{T}^{2})\right]$, $F(k_{T}^{2}) = \frac{h^{2}}{k_{T}^{2} + h^{2}}$

- ullet In top-pair production chosing hdamp $\sim m_t$ improves the description of the data
 - \blacktriangleright ATLAS tunes hdamp = 1.5 m_t , CMS sets to the same value



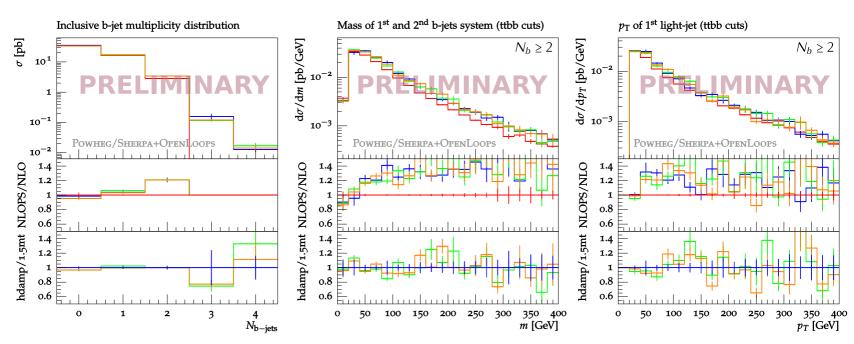
$$\mathrm{d}\,\sigma = \overline{B}(\Phi_B)\,\mathrm{d}\,\Phi_B \left[\Delta(q_\mathrm{cut}) + \sum_\alpha \Delta(k_T^\alpha) \frac{R_\alpha^s(\Phi_\alpha(\Phi_B,\Phi_\mathrm{rad}))}{B(\Phi_B)}\,\mathrm{d}\,\Phi_\mathrm{rad} \right] + \left(R_\alpha^r\mathrm{contr.}\right)$$
 where
$$R_\alpha^s = R_\alpha F(k_T^2) \;, \qquad F(k_T^2) = \frac{h^2}{k_T^2 + h^2}$$

- In $t\bar{t}b\bar{b}$:
 - ▶ Default behaviour of hdamp needs modifying:
 - ▶ Default "hdamp applied only to IS" manifests convergence issues
 - \triangleright We apply hdamp also to massive FS, with hdamp_{IS} and hdamp_{FS} independent
 - Further investigation underway
 - New POWHEG BOX RES features could be exploited for better understanding of the hdamp dependence

hdamp dependence



 \longrightarrow NLOPS hdamp = 1.5 m_t \longrightarrow NLOPS hdamp = 5 m_t \longrightarrow NLOPS hdamp = 20 m_t



- hdamp applied also to final state massive emitters
- Both NLO and LOPS accurate observables very stable with respect to hdamp
 - Variations of at most 10% are observed

Summary



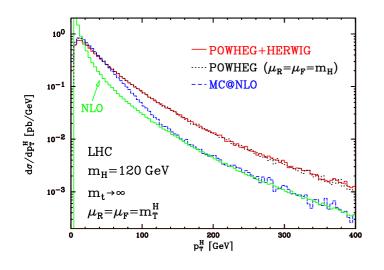
- Use POWHEG BOX RES framework to match 4F $t\bar{t}b\bar{b}$ at NLO QCD with PS
 - ► Matrix elements obtained from OpenLoops using the new OpenLoops+POWHEG BOX RES interface
 - ► The default hdamp behaviour extended in order to account for final state *b*-jets
 - ▶ Showered with Pythia8.2
- Compare our first results against Sherpa+OpenLoops
 - ▶ We observe a very good agreement
 - ▶ POWHEG BOX RES+OpenLoops confirms the "double splitting" enhancement
 - Comparison with YR benchmarks soon
- Study hdamp dependence
 - ▶ Predictions show very good stability with respect to hdamp variations
 - Further investigation under way

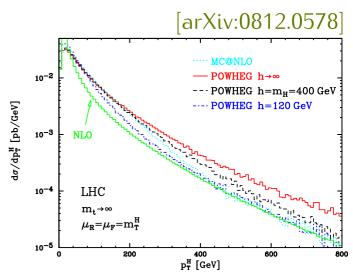


• POWHEG radiation formula:

$$\mathrm{d}\,\sigma = \overline{B}(\Phi_B)\,\mathrm{d}\,\Phi_B \left[\Delta(q_\mathrm{cut}) + \sum_{\alpha} \Delta(k_T^\alpha) \frac{R_\alpha^s(\Phi_\alpha(\Phi_B, \Phi_\mathrm{rad}))}{B(\Phi_B)} \,\mathrm{d}\,\Phi_\mathrm{rad} \right] + \left(R_\alpha^r \mathrm{contr.} \right)$$
 where
$$R_\alpha^s = R_\alpha F(k_T^2) \;, \qquad F(k_T^2) = \frac{h^2}{k_T^2 + h^2}$$

• In Higgs production:







• POWHEG radiation formula:

$$d \sigma = \overline{B}(\Phi_B) d \Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d \Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{contr.})$$
where $R_{\alpha}^s = R_{\alpha}F(k_T^2)$, $F(k_T^2) = \frac{h^2}{k_T^2 + h^2}$

In Higgs production:

