

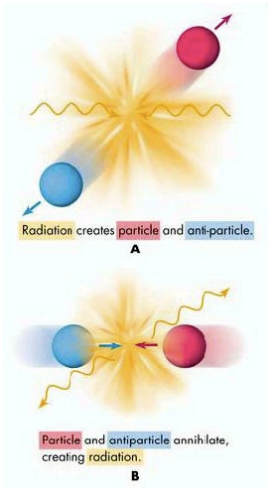
# HOW TO MASTER LIGHT MESON DYNAMICS:

from hadron spectroscopy to CP phases

June 15, 2018 | Christoph Hanhart | IKP/IAS Forschungszentrum Jülich



# MATTER EXCESS



In the early universe:

matter:antimatter as  
1 000 000 001:1 000 000 000

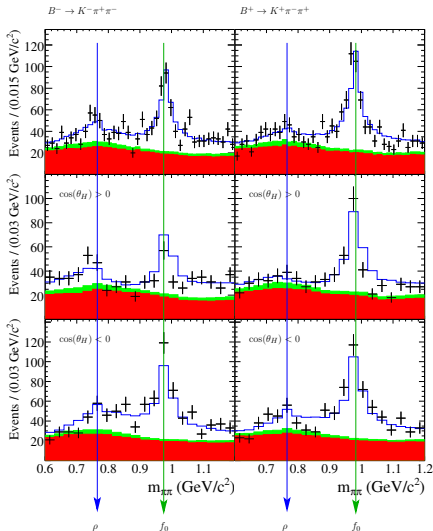
**Problem:** The Standard Model falls short by several orders of magnitude!

→ What is missing?

Look for the needle in the haystack



# SENSITIVE OBSERVABLES



Here:  $B^\pm \rightarrow K^\pm \pi^+ \pi^- \pi^\pm$

BABAR (2008); Belle (2006)

- $3.7\sigma$  in  $\rho K$
- $3.5\sigma$  in  $f_2(1270)K$
- $1.8\sigma$  in  $f_0(980)K$

Consistent with SM

Important next steps:

- Improve analysis and thus sensitivity
- Study  $D$  decays  
→ SM prediction tiny



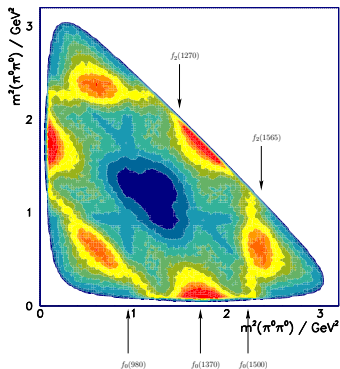
# WHY DALITZ PLOTS?

CP violation shows up as a phase

⇒ most visible in environments **with strong phase variation**

Example Dalitz plot:  $\bar{p}p \rightarrow 3\pi^0$  from Chrystal Barrel

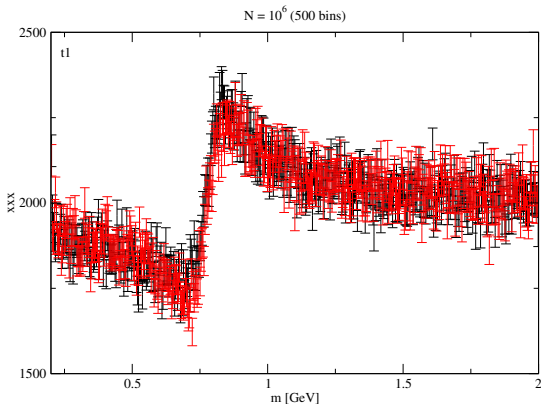
from Chrystal Barrel ( $\bar{p}p$  at rest); download: <http://www-meg.phys.cmu.edu/cb/pi03.html>



# ILLUSTRATION

Employing hadronic input drastically increases sensitivity!

$$N/\bar{N} = \alpha + \beta \operatorname{Re} \left\{ \exp(\pm i\delta_{CP}) \frac{1}{m^2 - M_{\text{res}}^2 + iM_{\text{res}}\Gamma_{\text{res}}} \right\}$$

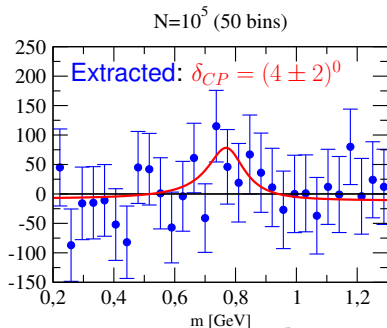
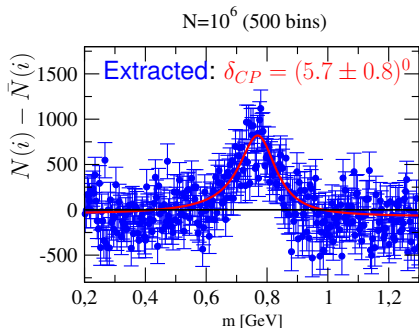


# ILLUSTRATION

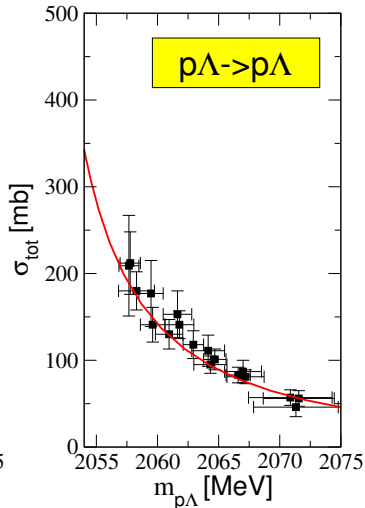
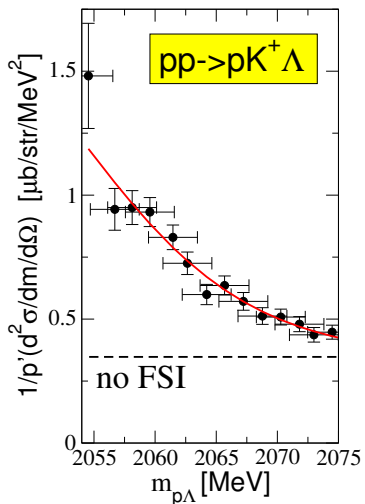
with this

$$N - \bar{N} = \delta_{CP} \times \left( \frac{2\beta M_{\text{res}} \Gamma_{\text{res}}}{(m^2 - M_{\text{res}}^2)^2 + (M_{\text{res}} \Gamma_{\text{res}})^2} \right) + \mathcal{O}(\delta_{CP}^2)$$

Input:  $\delta_{CP} = 5^0$ ;  $M_{\text{res}} = 0.77 \text{ GeV}$ ;  $\Gamma_{\text{res}} = 0.15 \text{ GeV}$



# PRODUCTION VS. SCATTERING



R. Siebert et al. (1994); G. Alexander et al. (1968)

In general: not identical, but related



# MODELING HADRON PHYSICS

Standard treatment: **sum of Breit-Wigners**

Propagator:  $iG_k(s) = \overline{\text{---}}_k = i/(s - M_k^2 + iM_k\Gamma_k)$

Scattering:  $\sum_k \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} = \sum_k ig_k^2 G_k(s)$

Production:  $\sum_k \otimes \text{---} \bullet \text{---} \otimes + \otimes = (\sum_k ig_k G_k(s)\alpha_k) + i\beta$

**Problems:**

- Wrong threshold behavior (cured by  $\Gamma = \Gamma(s)$ )
- Violates unitarity  $\rightarrow$  **wrong phase motion**
- Parameters reaction dependent  
**only pole positions and residues universal!**





# IMPOSING UNITARITY

- Properly constructed models fit to large sets of data

a la Bonn-Gatchina, J'ulich-Bonn, SAID

- or recent parametrization of near threshold cross sections

C.H. et al., PRL 115(2015)202001 and Guo et al. PRD93(2016)074031

- or **Dispersion Theory**: Starting point: Im-part of form factor  $F_i$

$$\text{Im}(F_i) = \sum_k T_{ik}^* \sigma_k F_k \quad \rightarrow \text{Dispersion Integral(s)}$$

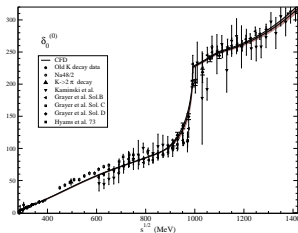
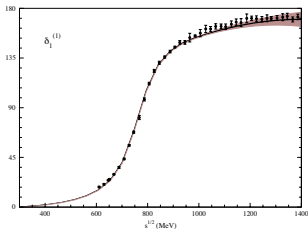
for **single channel**  $\rightarrow$  **Watson theorem** and **Omnès function**

$$\Omega(s) = \exp\left(\frac{s}{\pi} \int_{4m_\pi^2}^{\infty} ds' \frac{\delta(s')}{s'(s' - s - i\epsilon)}\right) \quad \text{and} \quad F(s) = P(s)\Omega(s)$$

- $\Omega(s)$  is **universal and fixed** in elastic regime
- $P(s)$  **reaction specific** and contains e.g.  
higher thresholds, inelastic resonances, left-hand cuts



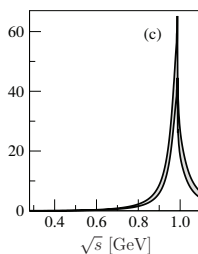
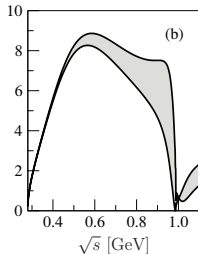
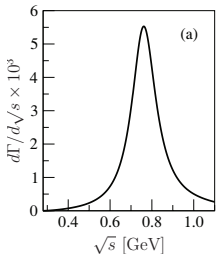
# STATUS $\Omega(s)$ : $\pi\pi$ S- AND P-WAVES



Give for  $\tau \rightarrow \mu\pi\pi$   $\Omega_1^+(s)(p+p')^\mu = \langle \pi\pi | \bar{q}\gamma^\mu q | 0 \rangle$

$\Gamma_\pi^n(s) = \langle \pi\pi | (\bar{u}u + \bar{d}d) / 2 | 0 \rangle$

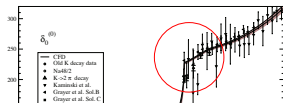
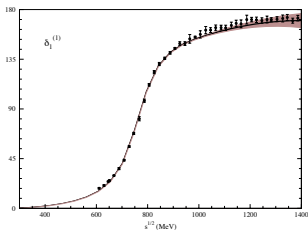
$\Gamma_\pi^S(s) = \langle \pi\pi | \bar{s}s | 0 \rangle$



$\delta$ : Garcia-Martin et al., PRD83(2011)074004; FF's: Daub et al., JHEP01(2013)179



# STATUS $\Omega(s)$ : $\pi\pi$ S- AND P-WAVES



Region of uncertainty - see also

Caprini et al., EPJC72(2012)1860

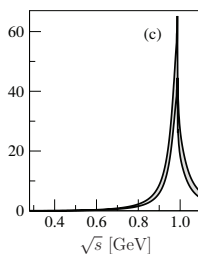
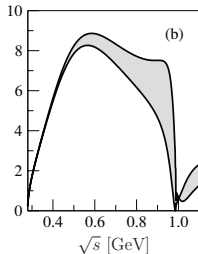
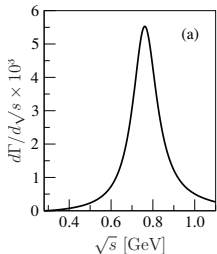
Büttiker et al., EPJC33(2004)409

Dai, Pennington, PRD90(2014)036004

Give for  $\tau \rightarrow \mu\pi\pi$   $\Omega_1^+(s)(\rho+\rho')^\mu = \langle \pi\pi | \bar{q}\gamma^\mu q | 0 \rangle$

$\Gamma_\pi^n(s) = \langle \pi\pi | (\bar{u}u + \bar{d}d) / 2 | 0 \rangle$

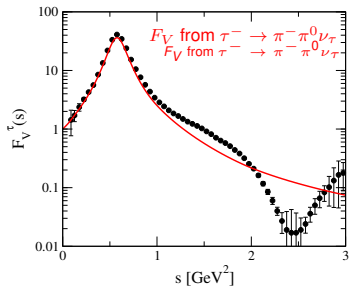
$\Gamma_\pi^S(s) = \langle \pi\pi | \bar{s}s | 0 \rangle$



$\delta$ : Garcia-Martin et al., PRD83(2011)074004; FF's: Daub et al., JHEP01(2013)179



# UNIVERSALITY OF FSI



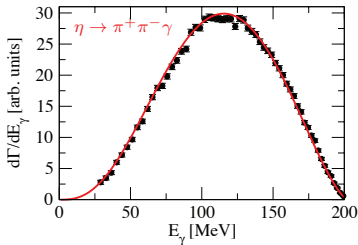
red lines: p-wave Omnes

× kinematic factors

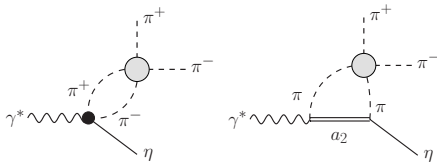
- bulk described properly

- there are deviations

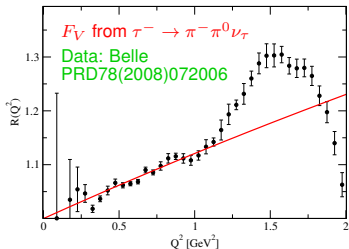
→  $P(s)$  not constant!



We need to add for  $\eta$ -decay:



# UNIVERSALITY OF FSI

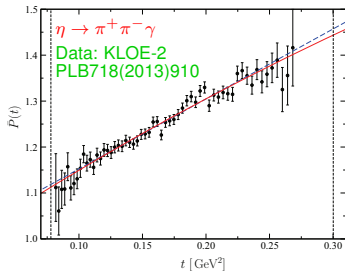


We write  $F_V(Q^2) = R(Q^2)\Omega(Q^2)$

We find

- $R(Q^2)$  linear for  $Q^2 < 1 \text{ GeV}^2$
- deviations by  $\rho'$  &  $\rho''$

C.H. et al., EPJC73(2013)2668



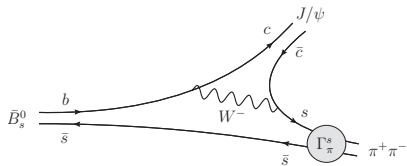
Inclusion of left-hand cut from  
 $a_2(1320)$  for  $\eta \rightarrow \pi\pi\gamma$

$$F_{a_2}(s_{\pi\pi}) = \Omega(s_{\pi\pi}) \left\{ A(1 + \alpha_\Omega[a_2] s_{\pi\pi}) + \frac{\cos \delta_1(s') \tilde{F}_{a_2}(s')}{|\Omega(s')|} + \frac{s_{\pi\pi}^2}{\pi} \int_{4m_\pi^2}^{\infty} \frac{ds'}{s'^2} \frac{\sin \delta_1(s') \tilde{F}_{a_2}(s')}{|\Omega(s')|(s' - s_{\pi\pi})} \right\},$$

Kubis and Plenter, EPJC75(2015)283



# APPLICATION: $\bar{B}_{s/d}^0 \rightarrow J/\psi \pi \pi$



- $\bar{B}_s^0$ : clean  $\bar{s}s$  source
- $\bar{B}_d^0$ : clean  $\bar{d}d$  source
- $\pi J/\psi$  interactions negligible  
→ no left-hand cuts

Ideal testing ground for formalism

e.g. for  $\bar{B}^0$ -decay:

- phenomenological analysis: inclusion of BW-functions for  $f_0(500)$ ,  $\rho(770)$ ,  $\omega(782) \rightarrow$  14 parameters
- dispersive approach for  $S$ - and  $P$ -waves  $\rightarrow$  3-4 parameters  
3 normalizations; eventually allowing for additional slope

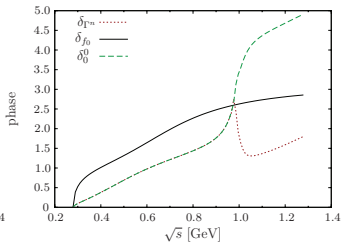
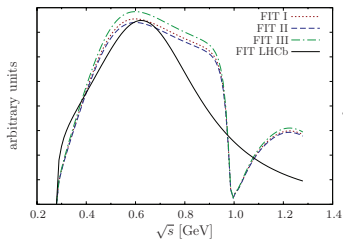
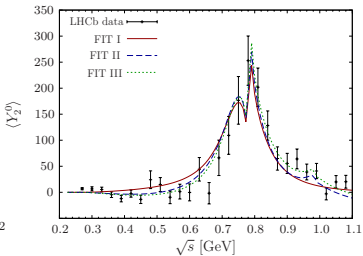
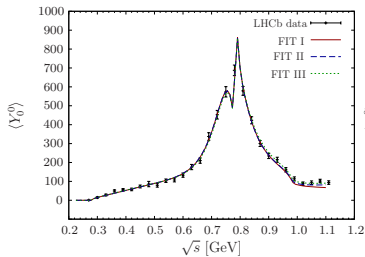
LHCb, PRD90(2014)012003

Daub, C.H., Kubis, JHEP1602(2016)009

→ Fits with even of higher quality and proper phase motion!



# RESULTS FOR $\bar{B}^0 \rightarrow J/\psi\pi\pi$



# RESULTS FOR $\bar{B}^0 \rightarrow J/\psi\pi\pi$

- The absence of the  $f_0(980)$  in the LHCb fit was interpreted as incompatibility with tetraquark structure (at  $8\sigma$ )

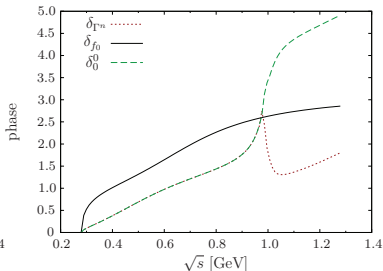
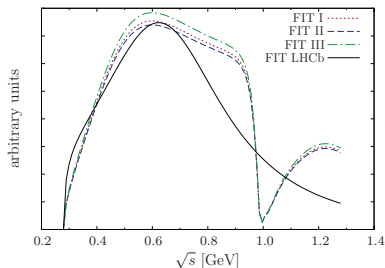
R. Aaij et al. (LHCb Collaboration) PRD90(2014)012003

- Re-analysis shows that this conclusion was premature

Daub, C.H., Kubis, JHEP1602(2016)009

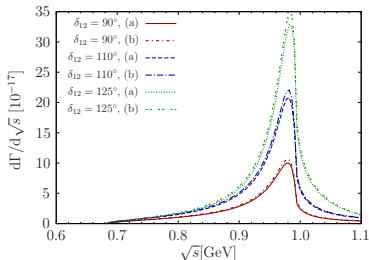
- Light scalars consistent with two-meson states

For a review on had.-molecules: F. K. Guo et al., Rev. Mod. Phys. 90(2018)015004





# OUTLOOK

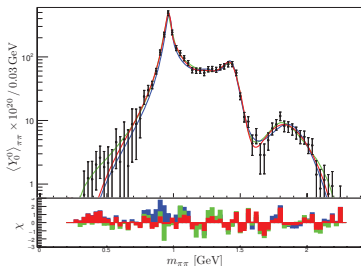


Data for  $\bar{B}_d^0 \rightarrow J/\psi\pi\eta$  will allow one to **fix final uncertainty in  $\eta\pi$  scattering phase shifts**

Albaladejo et al., JHEP1704(2017)010

employing

Albaladejo & Moussallam EPJC75(2015)488



Analysis of  $B_s \rightarrow J/\psi\pi\pi$  extended to higher energies: shown fits **by LHCb, 2 res., 3 res.**

Ropertz et al., in preparation

Data: LHCb, PRD89(2014)092006

adapting formalism of

C.H., PLB715(2012)170



# SUMMARY

- **States** are characterized by their **pole positions and residues**
- **Breit-Wigner** fits should in general be **avoided**
- Where ever possible **phase information** should be employed
  - either via **dispersion theory**

for other examples see, e.g.,

$\gamma\gamma \rightarrow \pi\pi$ : Hoferichter et al. 2011, Moussallam 2011, Mao et al. 2009, Pennington et al. 2008

$\omega/\phi \rightarrow \pi\pi\pi$ : Niecknig, Kubis, Schneider 2012, Danilkin et al. (2015)

- or via **models consistent with analyticity and unitarity**

coupled channel analyses for  $\pi N$ ,  $\pi\pi N$ ,  $\gamma N$  a la Bonn-Gatchina, Jülich-Bonn, SAID

The theoretical tools are becoming available to

- **extract the hadron spectrum** from data in a controlled way
- treat hadronic **final state interactions rigorously**
- **hunt for physics beyond the Standard Model**

