

Thermal Dark Matter and the Higgs

SB, Carena, Shah, Wagner; 1712.09873

Sebastian Baum

Oskar Klein Centre for Cosmoparticle Physics
Stockholm University



Stockholm
University



UCLA Dark Matter 2018

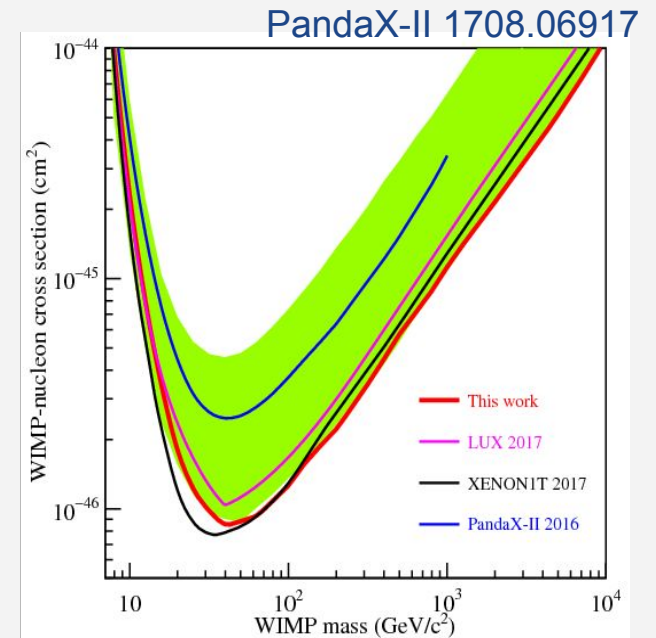
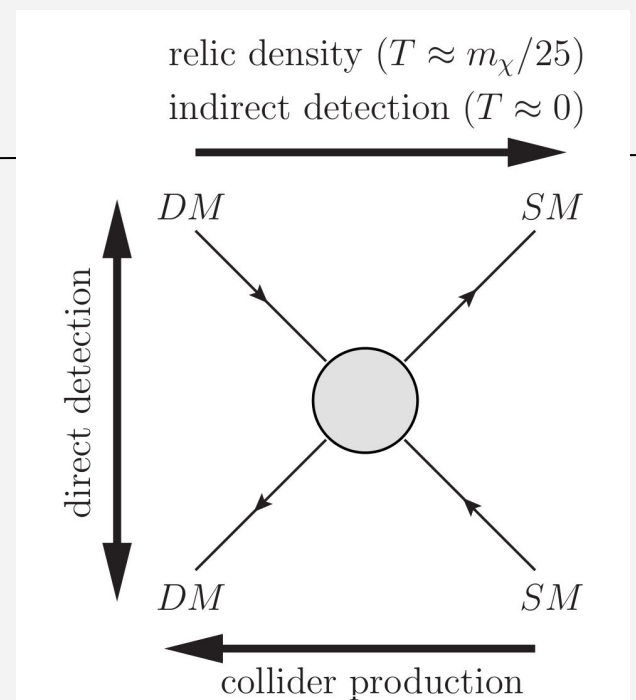
WIMP Dark Matter

Vanilla thermal production:

$$\langle\sigma v\rangle \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

but direct detection limits are very strong

Production and direct detection cross-sections must be decoupled (unless $m_{\text{DM}} \gg 100 \text{ GeV}$)



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- resonant annihilation

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Both options require extended BSM sectors to supply either

- co-annihilation partners
- Multiple mediators allowing for destructive interference

Model-building

- Dark Matter candidate: Majorana fermion χ
- Type II 2 Higgs Doublet Model (2HDM) (H_u, H_d)
- Invoke a Z_3 symmetry under which all scalars/fermions transform as $\psi \rightarrow e^{2\pi i} \psi$ to forbid Majorana mass for χ
- Weak scale mass for χ can be obtained via the vacuum expectation value (vev) of an extra SM-singlet (complex) scalar S
 - Generates weak scale mass for χ if vev $\langle S \rangle$ from same mechanism as EWSB
 - S couples to SM only via mixing with neutral components of H_u, H_d after EWSB

Model-building

- Write down all operators to dimension $d \leq 6$ arising from integrating out a heavy $SU(2)$ -doublet Dirac fermion

(ignoring charged gauge boson interactions)

$$\begin{aligned} \mathcal{L} = & -\delta \frac{\chi\chi}{\mu} (H_u \cdot H_d) \left(1 - \frac{\lambda \hat{S}}{\mu} \right) - \kappa S \chi\chi \left(1 + \xi \frac{H_d^\dagger H_d + H_u^\dagger H_u}{|\mu|^2} \right) + \text{h.c.} \\ & + \frac{\alpha}{|\mu|^2} \left\{ \chi^\dagger H_u^\dagger \bar{\sigma}^\mu \left[i\partial_\mu - \frac{g_1}{s_W} (T_3 - Q s_W^2) Z_\mu \right] (\chi H_u) \right. \\ & \left. + \chi^\dagger H_d^\dagger \bar{\sigma}^\mu \left[i\partial_\mu - \frac{g_1}{s_W} (T_3 - Q s_W^2) Z_\mu \right] (\chi H_d) \right\}, \end{aligned}$$

- DM mass $m_\chi = 2\kappa \langle S \rangle$
- Heavy $SU(2)$ -doublet fermion mass: $\mu = \lambda \langle S \rangle$

Model phenomenology. Higgs sector

- Rotate scalars to Higgs basis

$$H^{\text{SM}} = \sqrt{2}\text{Re}(\sin\beta H_u^0 + \cos\beta H_d^0), \quad \leftarrow \text{SM-like}$$

$$G^0 = \sqrt{2}\text{Im}(\sin\beta H_u^0 - \cos\beta H_d^0), \quad \leftarrow \text{Neutral Goldstone (Z)}$$

$$H^{\text{NSM}} = \sqrt{2}\text{Re}(\cos\beta H_u^0 - \sin\beta H_d^0), \quad \leftarrow \text{MSSM-like scalar}$$

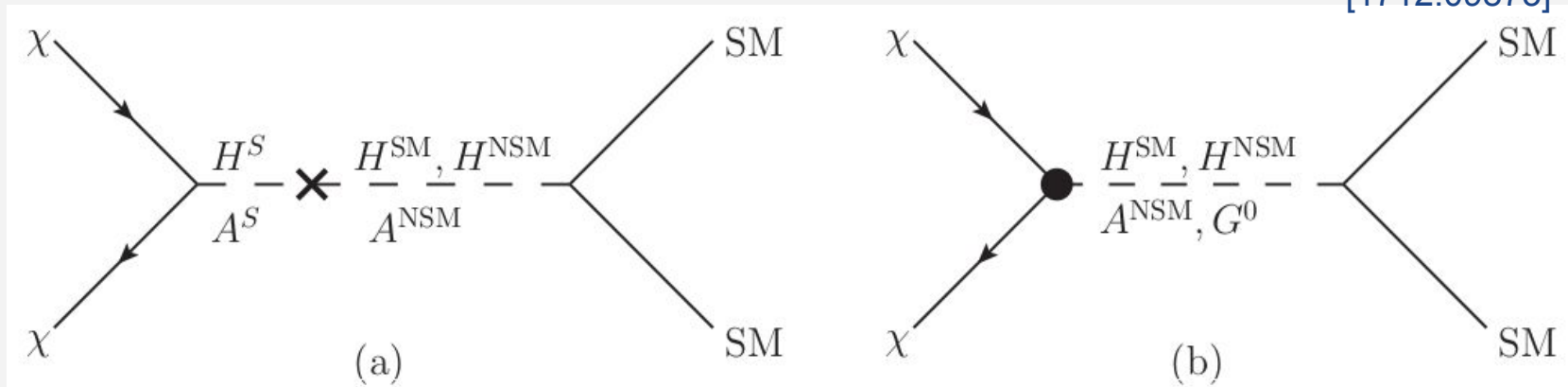
$$A^{\text{NSM}} = \sqrt{2}\text{Im}(\cos\beta H_u^0 + \sin\beta H_d^0), \quad \leftarrow \text{MSSM-like pseudoscalar}$$

+ scalar (H^S) and pseudoscalar (A^S) from singlet

- 125 GeV Higgs pheno requires approximate alignment: $h_{125} \sim H^{\text{SM}}$

Model phenomenology. DM

[1712.09873]



mixing

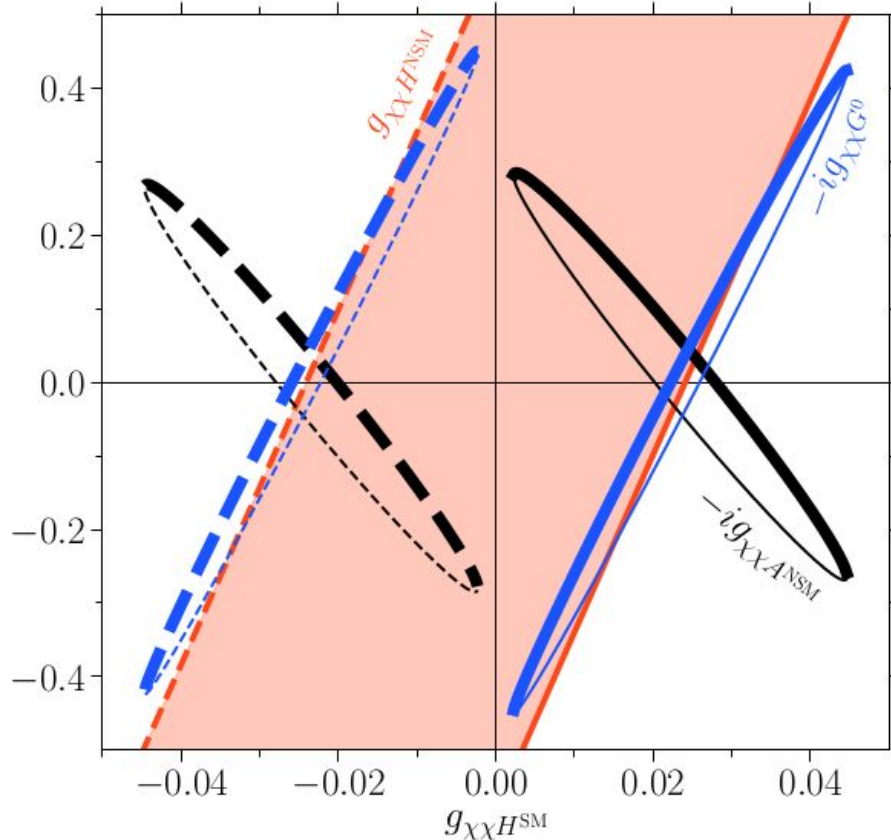
$d \geq 5$ operators

- Thermal production via
 - $\chi\chi \rightarrow q\bar{q}$ (important contribution from longitudinal Z/neutral Goldstone)
 - Usually small contributions from $\chi\chi \rightarrow h_i h_j / h_i a_j / WW / ZZ$
- SI Direct Detection mediated predominantly by CP-even Higgs bosons
 - EFT operators allow for blind spot where $g_{\chi\chi h} \rightarrow 0$
 - Presence of 3 CP-even Higgses allows for destructive interference

Model phenomenology. DM

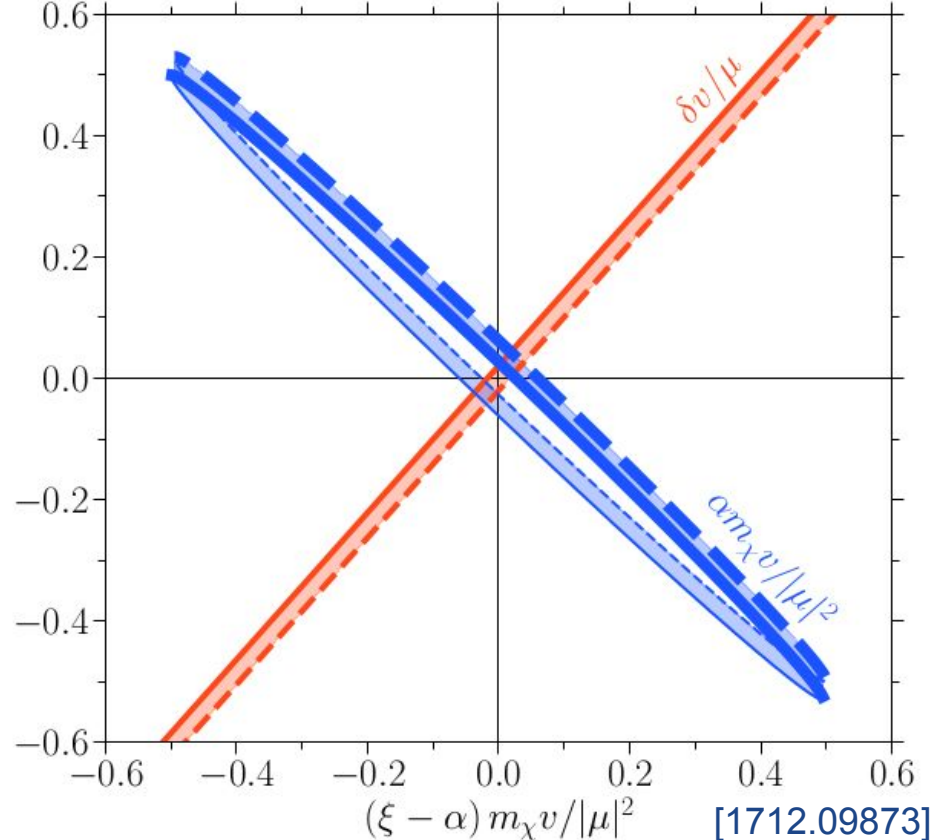
$$m_\chi = 300 \text{ GeV}; m_{H^{\text{NSM}}} = m_{A^{\text{NSM}}} = 500 \text{ GeV}; \tan \beta = 2;$$

$$\sigma_p^{\text{SI}} \leq 3.0 \times 10^{-10} \text{ pb}; \langle \sigma v \rangle_{x_F} = 2.3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$



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Example of allowed couplings / EFT parameters

(standard thermal production & current Direct Detection limits)

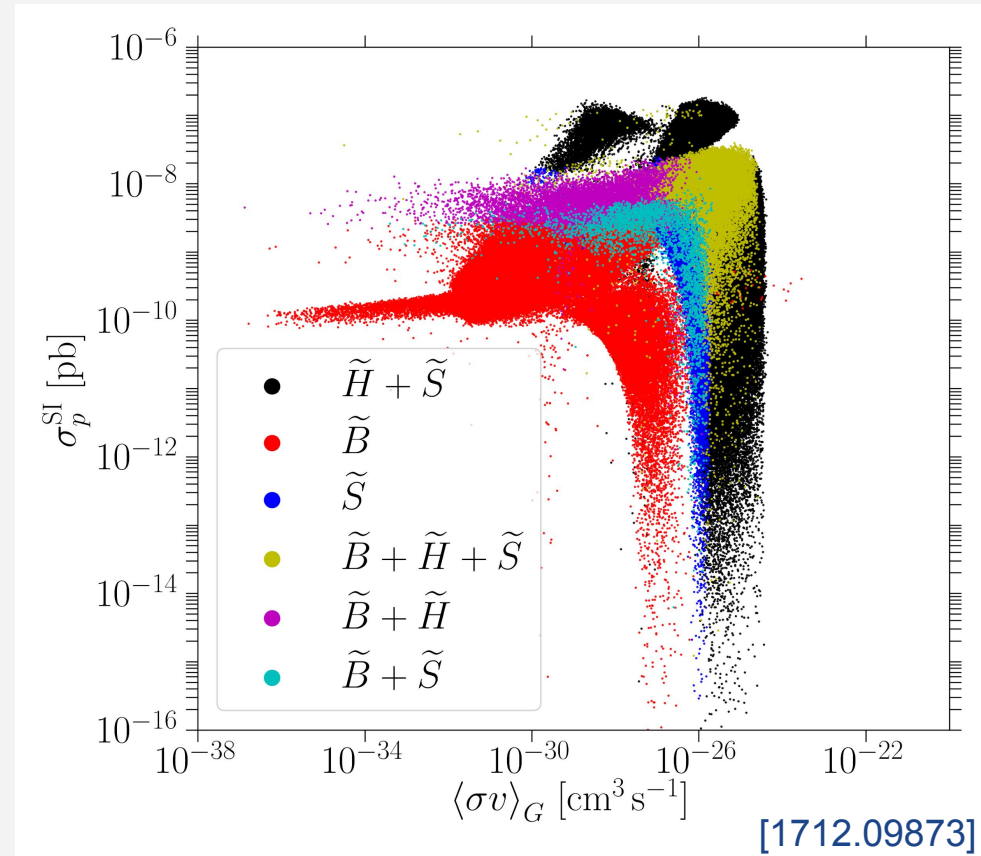
Does this really work? (aka UV completion)

(Z_3 -invariant) Next to minimal Supersymmetric Standard Model (NMSSM)

- MSSM particle content + one extra SM-singlet chiral superfield
- Same Higgs sector as our model (2HDM+S)
- Fermionic component of singlet ('singlino') plays role of our DM candidate
- Mature numerical tools exist to study DM & collider phenomenology
- Besides the singlino, the NMSSM contains a second SM-singlet fermion, the bino
 - Smaller couplings to Higgs sector than singlino
 - EFT model can be extended in simple fashion to include bino

NMSSM: Singlino DM. Production

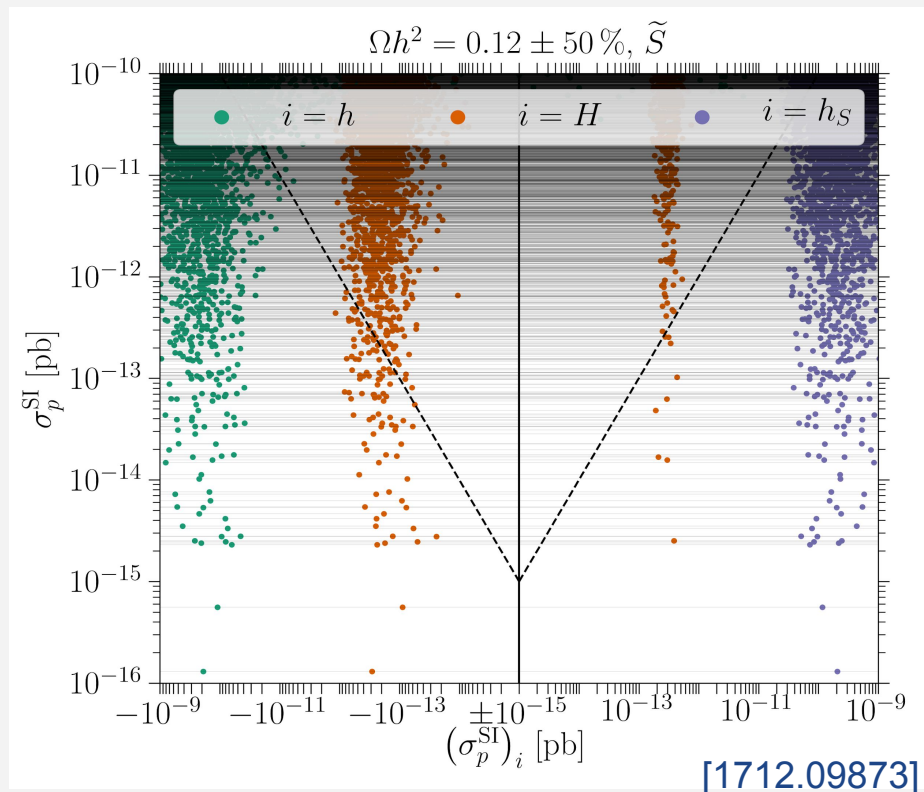
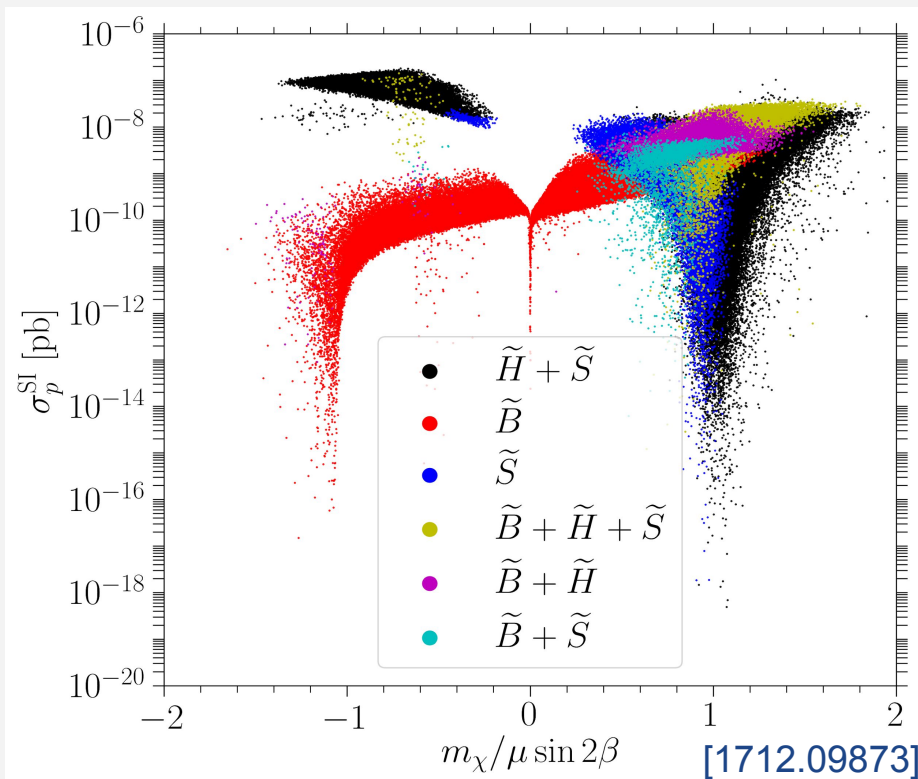
- predominantly vanilla thermal production dominated by annihilations into top quarks
- Longitudinal Z (neutral Goldstone mode) plays prominent role



NMSSM: Singlino DM. Direct Detection

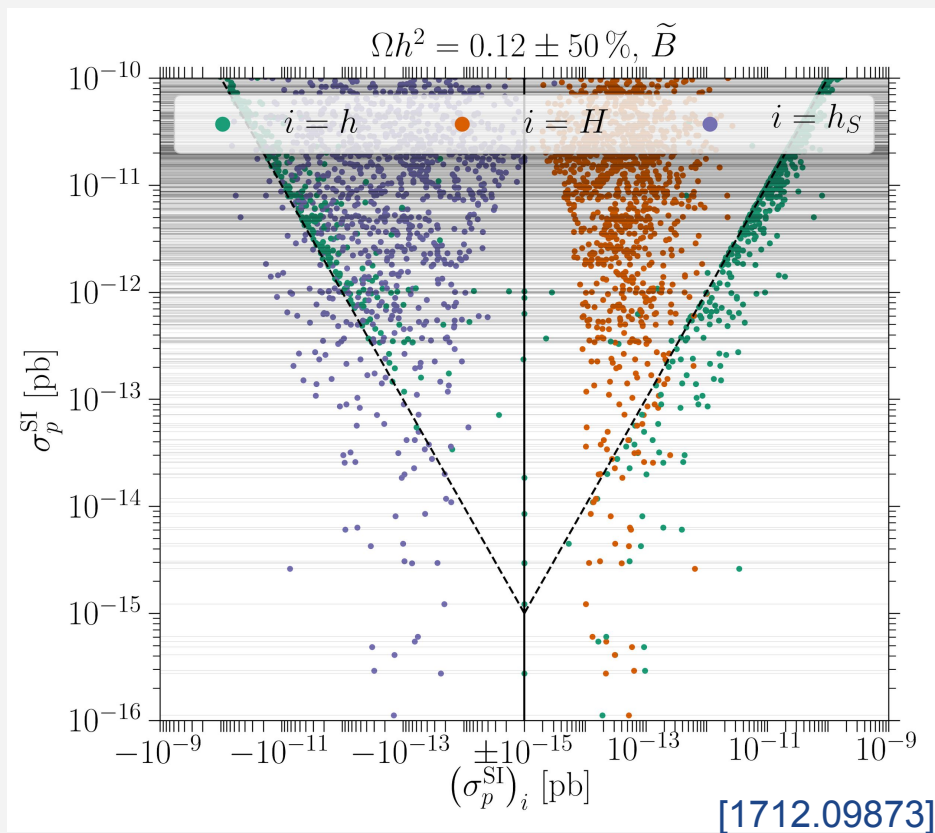
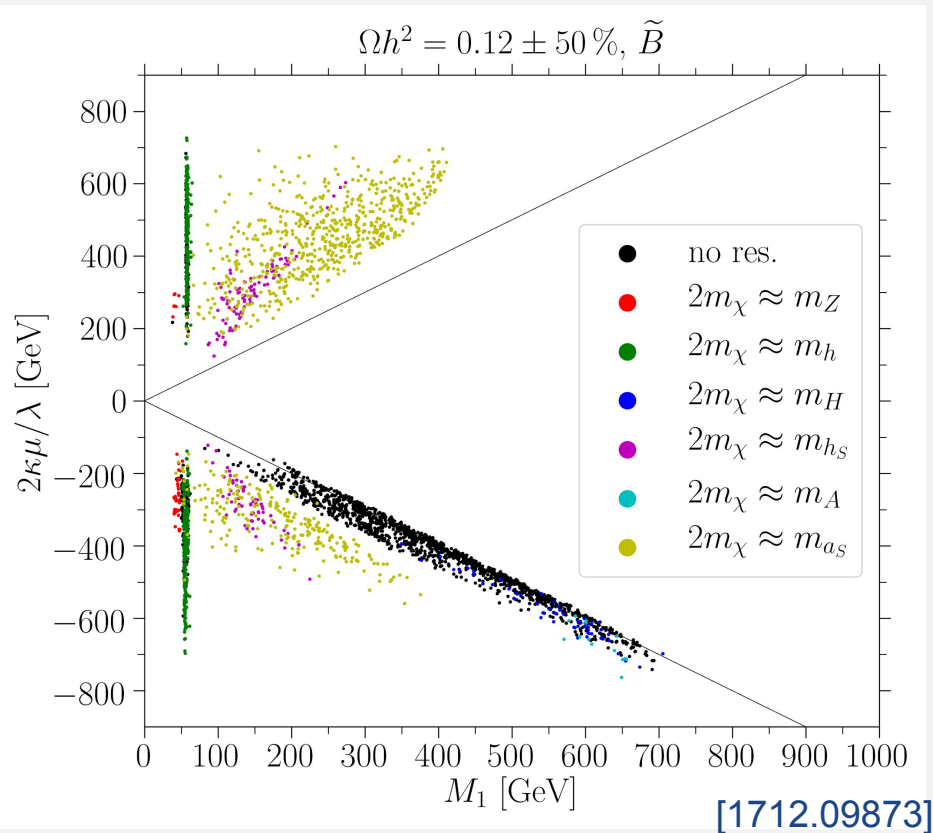
SI Direct Detection cross section suppressed by:

- Proximity to h_{125} blind spot (Cheung+ 1406.6372, Badziak+ 1512.02472)
- Destructive interference between CP-even Higgs mass eigenstates

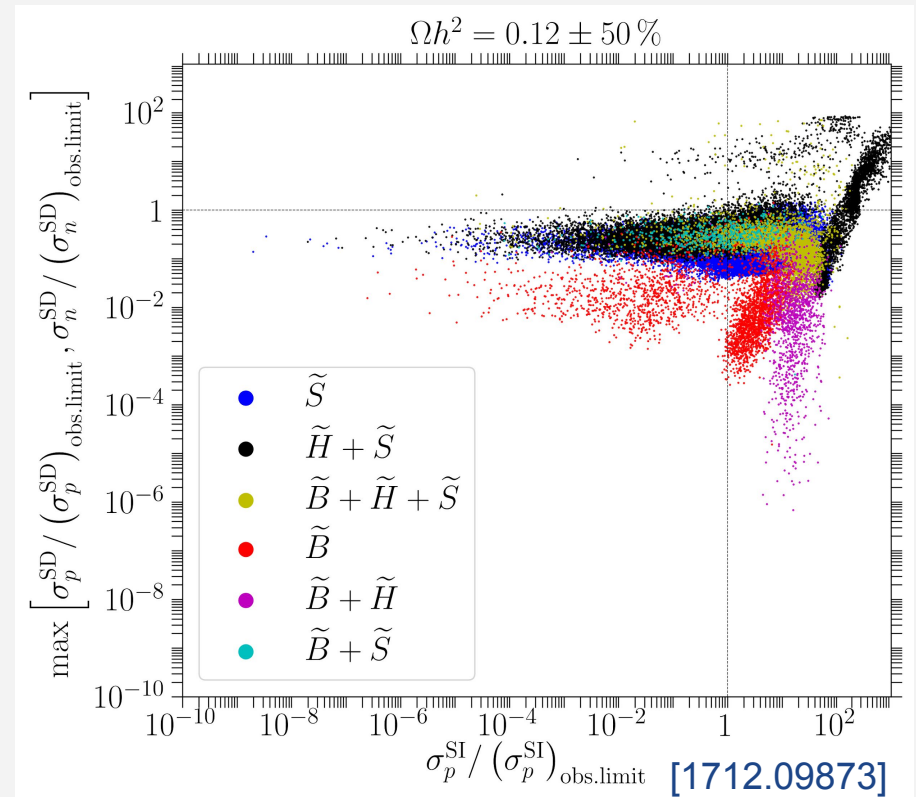
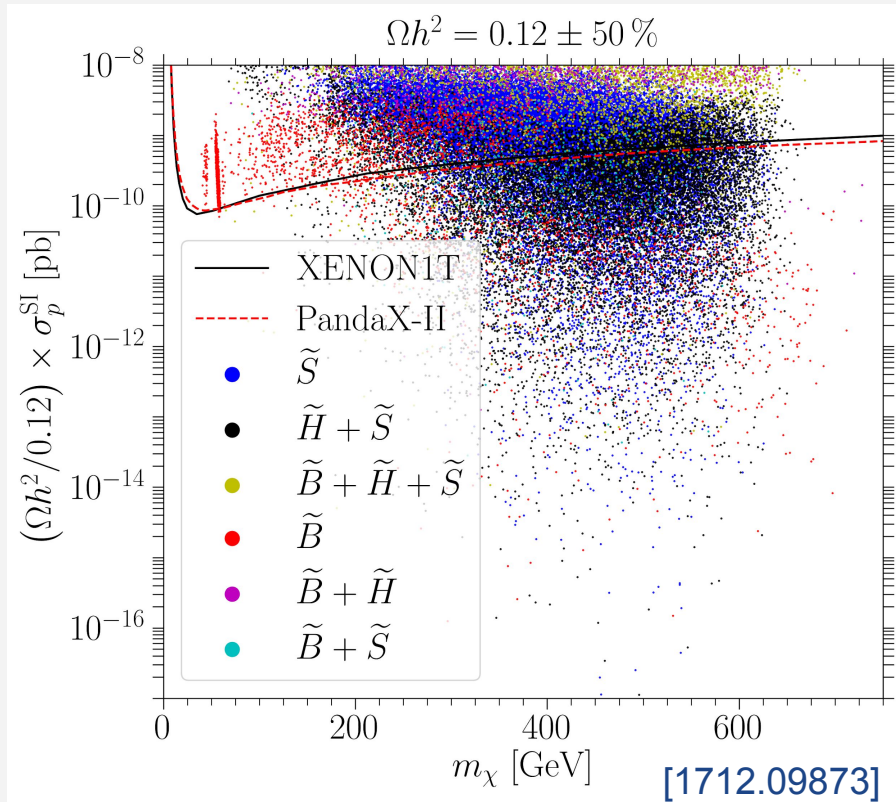


NMSSM: (new 'well-tempered') Bino DM

- Thermal production via co-annihilation with singlino
- SI Direct Detection cross section suppressed by h_{125} blind spot proximity



NMSSM: DM. Direct Detection Potential



- As yet, SI limits are much more relevant
- But, increasing sensitivity of SD experiments by ~ 2 orders of magnitude would probe most of the remaining parameter space!

Conclusions

- Stringent Direct Detection limits require decoupling of SI detection cross section and annihilation cross section
 - Either, one enhances the production cross-section by co-annihilation or resonant annihilation,
 - Or, one suppresses the SI detection cross section by destructive interference
- Qualitative features studied in Majorana DM + (2HDM+S) model
- Tested validity of results and extended numerical study in NMSSM
 - Singlino region: Vanilla thermal production (via the longitudinal Z), h_{125} blind spot + destructive interference for direct detection
 - New 'well-tempered' Bino region: co-annihilation with singlino for thermal production, blind-spot for direct detection
 - Both can be probed by SD experiments in near future!