

UV conformal window for asymptotic safety

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- Motivation
- Interacting fixed points
- Asymptotically safe theory
- Fixed points and vacuum stability
- Competition of fluctuations
- Conclusions

- Asymptotic safety: interactive fixed point; residual interactions in the UV
- Critical phenomena: Wilson-Fisher fixed point (ϕ^4 theory in $d = 4 - \epsilon$ dimensions) (Wilson & Fisher, 1972)
- A quantum theory of gravity may be non-perturbatively renormalizable provided there is an interactive UV fixed point. (Weinberg, 1979)
- Necessary conditions and no-go theorems for asymptotic safety in gauge theories have been derived at weak coupling. (Bond & Litim, 2017a)
- There are exact proofs of existence in simple and semi-simple gauge groups, as well as SUSY. (Litim & Sanino, 2014)
(Bond & Litim, 2017b)
(Bond & Litim, 2018)
- Higher dimensional scalar operators are ok. Fixed points are available away from $4d$. (Buyukbese & Litim, 2017)
- Asymptotically safe model building now available. (Codello, Langæble, Litim, & Sannino, 2016)
- IR conformal windows in QCD-like theories have been extensively studied, and are known to extend past the domain of perturbation theory. (Bond, Hiller, Kowalska & Litim, 2017)
- What is the size of the conformal window of asymptotically safe theories? How many fields are required for asymptotic safety? (Banks & Zaks, 1982)
(Appelquist et al., 2008)
(Del Debbio., 2011)

Interacting fixed points

- Gauge renormalization group (RG) running at two loop in perturbation theory

$$\beta = \frac{d\alpha}{d\ln\mu} = -B\alpha^2 + C\alpha^3$$

- The RG flow vanishes at the fixed point

$$\beta^* = 0 \quad \alpha^* = \frac{B}{C}$$

- Physical coupling: B and C must have the same sign
- Weak coupling: B must be much smaller than C

Interacting fixed points

- Give-up asymptotic freedom: $B < 0$; $C < 0$
- $C > 0$ for any simple or semi simple gauge group, for any matter fields multiplicities and representation
- $C < 0$ can only be achieved with Yukawa interactions
- Necessary ingredients:
 - Gauge group (simple or semi-simple)
 - Fermions
 - Scalars
 - Yukawa interactions

(Bond & Litim, 2017a)

Asymptotically safe theory

$$L = L_{\text{YM}} + L_{\text{kin.}} + L_{\text{Yuk.}} + L_{\text{pot.}}$$

$$L_{\text{YM}} = -\frac{1}{2} \text{Tr} F^{\mu\nu} F_{\mu\nu}$$

$$L_{\text{kin.}} = \text{Tr} (\bar{Q} i \not{D} Q) + \text{Tr} (\partial_\mu H^\dagger \partial^\mu H)$$

$$L_{\text{Yuk.}} = -y \text{Tr} (\bar{Q}_L H Q_R) + \text{h.c.}$$

$$L_{\text{pot.}} = -u \text{Tr} (H^\dagger H H^\dagger H) - v (\text{Tr} H^\dagger H)^2$$

- $4d$ gauge theory
- NF fermions in the fundamental representation of $SU(\text{NC})$
- Meson-like complex scalar NF x NF matrix, uncharged
- Yukawa & scalar quartic interactions

Asymptotically safe theory

$$\alpha_g = \frac{g^2 N_C}{(4\pi)^2}; \quad \alpha_y = \frac{y^2 N_C}{(4\pi)^2}; \quad \alpha_u = \frac{u N_F}{(4\pi)^2}; \quad \alpha_v = \frac{v N_F^2}{(4\pi)^2}$$

- Take the Veneziano limit; field multiplicities are taken to infinity, keeping their ratio fixed
- One free parameter left, which can be taken to be perturbatively small
- The theory has a weakly coupled fixed point

$$N_C \rightarrow \infty; \quad N_F \rightarrow \infty$$

$$\epsilon = \frac{N_F}{N_C} - \frac{11}{2}$$

$$0 < \epsilon \ll 1$$

Asymptotically safe theory

- Weak coupling: series expansion in a small parameter is justified

$$\alpha_i^* = \lambda_{1i}\epsilon + \lambda_{2i}\epsilon^2 + \lambda_{3i}\epsilon^3 + O(\epsilon^n)$$

- The beta functions can then be expanded in epsilon. Exactly at the fixed point, all orders must vanish individually. The resulting equations determine the λ coefficients.
- In perturbation theory only the first few loop orders of the beta functions are known.
- Approximations have to be consistent.

Asymptotically safe theory

- The approximation is denoted by the number of loop orders retained:

$$\left. \begin{array}{l} n \text{ loop gauge} \\ m \text{ loop Yukawa} \\ l \text{ loop scalar} \end{array} \right\} \equiv (n, m, l)$$

- The approximation $(n + 1, n, n)$ completely determines the coefficient of order $O(\epsilon^n)$ in the series expansion of the fixed point.

$$\begin{array}{ccccccc} & & (3, 2, 2) & & (4, 3, 3) & & (n + 1, n, n) \\ & (2, 1, 1) & & & & & \\ & \searrow & \downarrow & \swarrow & \searrow & & \\ \alpha_i^* & = & \lambda_{1i}\epsilon & + & \lambda_{2i}\epsilon^2 & + & \lambda_{3i}\epsilon^3 & + & O(\epsilon^n) \end{array}$$

Fixed points and vacuum stability

- We have computed the running of the scalar self-interactions to two loop. This means our most advanced approximation is (322), an improvement over the original analysis (321).
- The fixed point can be determined analytically to order ϵ^2 . Numerically, we find:

$$\alpha_g^* = 0.4561\epsilon + 0.7808\epsilon^2 + O(\epsilon^3)$$

$$\alpha_y^* = 0.2105\epsilon + 0.5082\epsilon^2 + O(\epsilon^3)$$

$$\alpha_u^* = 0.1998\epsilon + 0.4403\epsilon^2 + O(\epsilon^3)$$

$$\alpha_v^* = -0.1373\epsilon - 0.6318\epsilon^2 + O(\epsilon^3)$$

Fixed points and vacuum stability

- In order to have a stable vacuum state, we require for the scalar potential to be bounded from below. In the present setting, this means:

$$\alpha_u^* > 0, \quad \alpha_u^* + \alpha_v^* > 0$$

(Litim, Mojaza & Sannino, 2016)

- Comparing with the previous approximations, we detect a change of sign in the subleading term. This implies that the vacuum becomes unstable at a finite value of epsilon.

$$\alpha_u^* + \alpha_v^*|_{(211)} = 0.0625\epsilon + O(\epsilon^3)$$

$$\alpha_u^* + \alpha_v^*|_{(321)} = 0.0625\epsilon + 0.1535\epsilon^2 + O(\epsilon^3)$$

$$\alpha_u^* + \alpha_v^*|_{(322)} = 0.0625\epsilon - 0.1915\epsilon^2 + O(\epsilon^3)$$

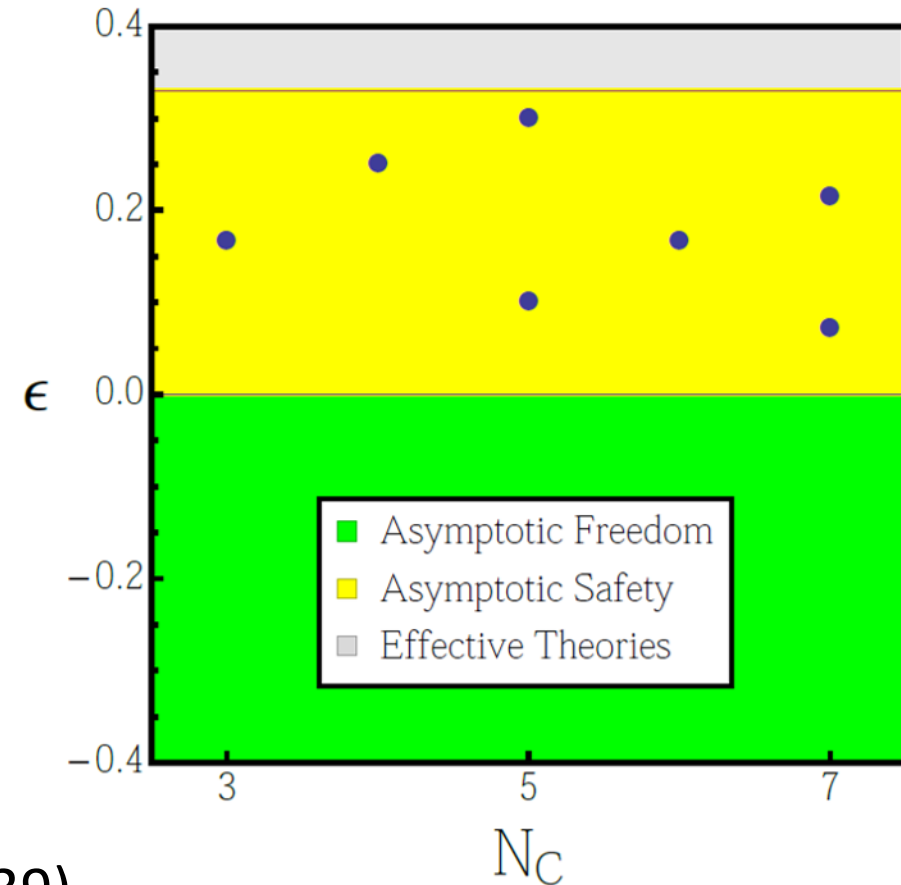
Fixed points and vacuum stability

- Vacuum instability imposes a physical constraint.

$$0 < \epsilon < 0.326$$

- This effect is driven by the two loop running of the scalar self-interactions.
- The plot shows the UV conformal window. Blue dots correspond to smallest field multiplicities:

$$(N_C, N_F) = (3, 17), (4, 23), (5, 28), (5, 29), (6, 34), (7, 39), \dots$$



Competition of fluctuations

- The size of the conformal window can be understood as a competition of fluctuations
- Consider the following qualitative indicator:

$$\beta_g|_{(322)} = 10.24\epsilon^5, \quad \beta_y|_{(322)} = -1.71\epsilon^4, \quad \beta_u|_{(322)} = 1.70\epsilon^4, \quad \beta_v|_{(322)} = 7.24\epsilon^4$$

- This is obtained by inserting the fixed point at order ϵ^2 back into the beta functions
- Positive shifts to the flow equations destabilize the fixed point, while negative ones stabilize it
- The running of the gauge and scalar interactions tend to close the conformal window, while the Yukawa interaction tends to open it

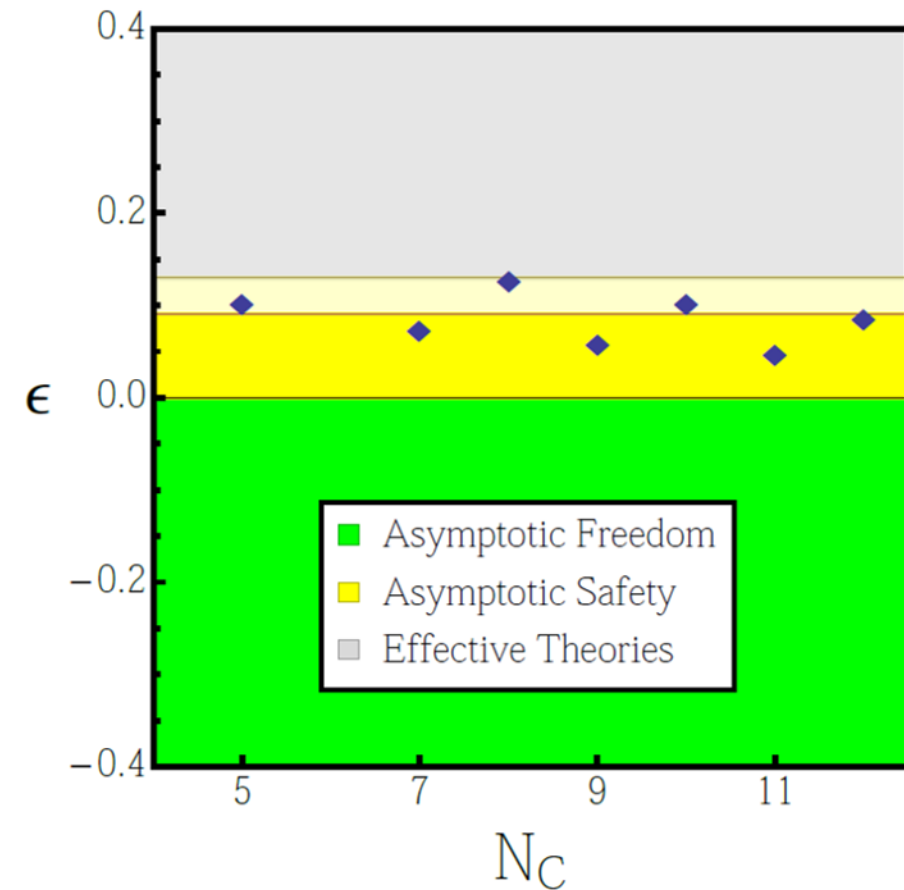
Competition of fluctuations

- We now consider the (incomplete) subleading terms in the beta functions
- The plot shows the conformal window of approximation (321) bounded by a fixed point merger (light yellow), and (322), bounded by vacuum instability (dark yellow).
- Vacuum instability still poses the tightest constraint at a numerically smaller value

$$\epsilon_{max} \approx 0.09 \dots 0.13$$

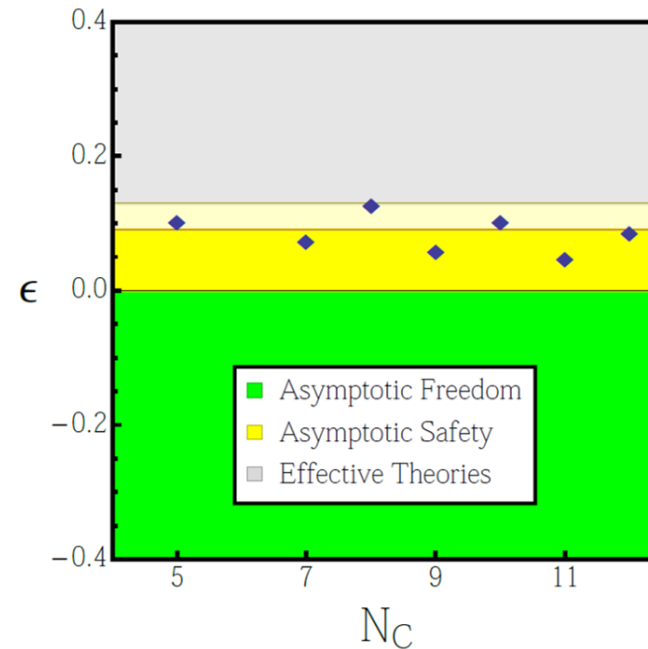
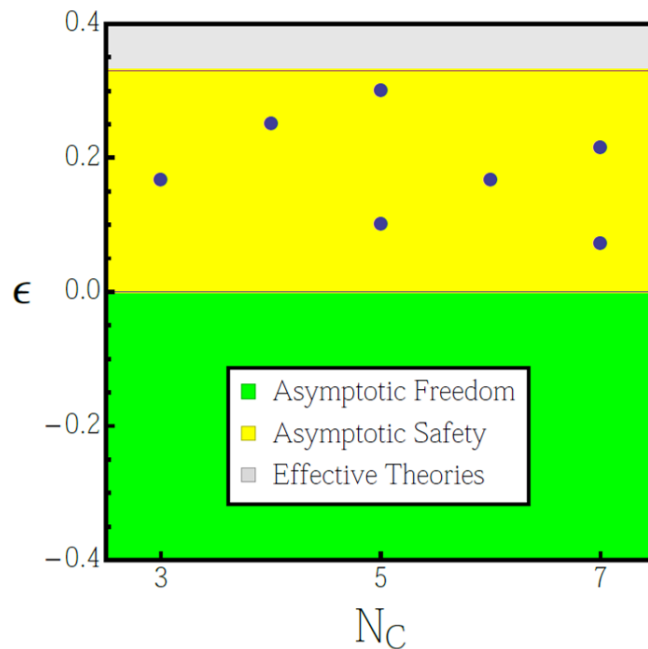
- Blue dots are:

$$(N_C, N_F) = (7, 39), (9, 50), (11, 61), (12, 67), \dots$$



Conclusions

- Vacuum instability imposes the tightest constraint in the theory. This is driven by the two loop running of the scalar self-interactions.
- Two pictures: Which one is right? Both are approximations that we use to understand the trend at higher loop orders.



- The UV conformal window remains roughly within the region of perturbation theory.
- It would be interesting to obtain the next approximation (433), then we can start to say something about the convergence of the boundary of the conformal window.

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