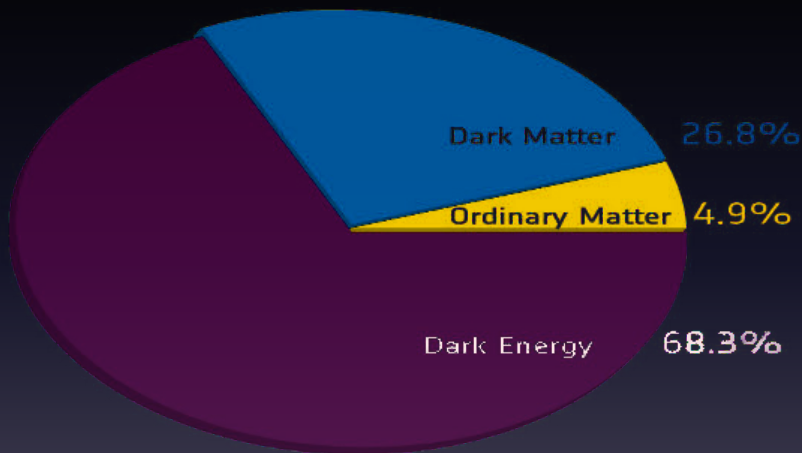


A photograph of a city skyline at night, featuring several tall skyscrapers with illuminated windows. A bright, jagged lightning bolt strikes one of the buildings, creating a dramatic focal point against the dark sky.

Laboratory fifth
force experiments
probe dark energy

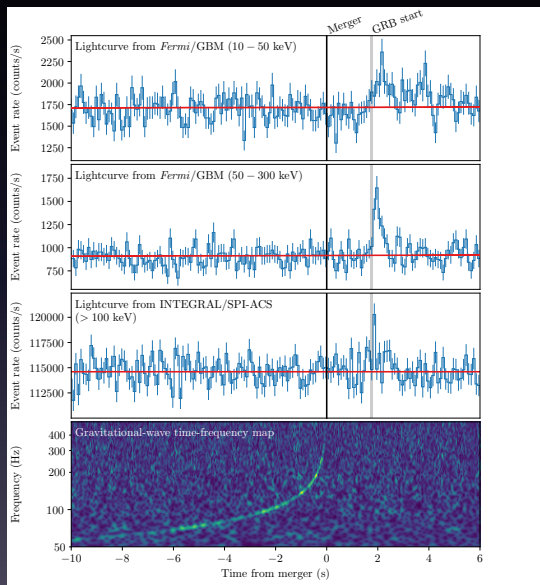
Amol Upadhye
UNSW Sydney
Nov. 30, 2017

The greatest mystery in cosmology



2/3 of the effective energy density is a cosmological constant Λ , a “dark energy,” or some modification to gravity.

What does LIGO tell us?



LIGO, Virgo, Fermi, and INTEGRAL: $\Delta t = 1.74$ sec

1 Mpc/c $\approx 10^{14}$ sec, so:

$$\alpha_{GW} = \frac{c_{GW}^2 - c^2}{c^2} \lesssim \frac{\text{few seconds}}{40 \text{ Mpc}/c} \sim 10^{-15}$$

Tuning?

- Δt from emission, escape?
- α_{GW} small only for $z < 0.01$?

What does LIGO tell us?

Most general scalar-tensor gravity with 2nd-order e.o.m.:

$$S = \int d^4x \sqrt{-g} [\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4 + \mathcal{L}_5 + \mathcal{L}_M]$$

$$\mathcal{L}_2 = G_2 \quad \mathcal{L}_3 = G_3 \square \phi$$

$$\mathcal{L}_4 = G_{4,X} \left[(\square \phi)^2 - \nabla_\mu \nabla_\nu \phi \nabla^\mu \nabla^\nu \phi \right] + G_4 R$$

$$\mathcal{L}_5 = -\frac{1}{6} G_{5,X} \left[(\square \phi)^3 - 3 \nabla^\mu \nabla^\nu \phi \nabla_\mu \nabla_\nu \phi \square \phi + 2 \nabla^\nu \nabla_\mu \phi \nabla^\alpha \nabla_\nu \nabla^\mu \nabla_\alpha \phi \right] \\ + G_5 G_{\mu\nu} \nabla^\mu \nabla^\nu \phi$$

where the G_i are functions of ϕ and $X = -\frac{1}{2} \nabla_\mu \phi \nabla^\mu \phi$. Also, $G_{i,\phi}$ and $G_{i,X}$ are $\partial G_i / \partial \phi$ and $\partial G_i / \partial X$.

Velocity correction: $\alpha_{GW} \propto 2G_{4,X} - 2G_{5,\phi} - (\ddot{\phi} - \dot{\phi}H)G_{5,X}$

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$\alpha_{GW} = 0$, no tuning: What's left?

LIGO-preferred extensions to cosmological constant Λ :

Not-so-slow roll:

Conformal coupling $G_4(\phi)$
small to suppress $\partial G_N / \partial t$

- Quintessence: 
Wetterich 1988; Peebles, Ratra 1988
 $G_2 = X + V(\phi)$
 $w = P/\rho = \frac{X-V}{X+V} \neq 1$
observed cosmologically
- Cubic Galileon:
Nicolis, Rattazzi, Trincherini 2008
 $G_3 \propto X$, 5th force screened
through kinetic term

Conformal coupling:

Rolling $\Delta\phi/M_{\text{Pl}} \ll 1$ to suppress fifth
force from $G_4(\phi)$

- Chameleons, $f(R)$ gravity:
Khoury, Weltman 2004
5th force screened through
effective mass $m_{\text{eff}} = V''$
- Symmetrons:
Olive, Pospelov 2007; Hinterbichler, Khoury 2010
 G_2, G_4 invariant under $\phi \rightarrow -\phi$;
5th force screened through
symmetry restoration

Screening of fifth forces

“Hide” 5th forces in dense regions such as the Solar System.
Consider small perturbations about a background: $\phi = \bar{\phi} + \delta\phi$:

$$\left[Z(\bar{\phi}) \square - m(\bar{\phi})^2 \right] \delta\phi = \frac{1}{M_{\text{Pl}}} \beta(\bar{\phi}) \delta\rho$$

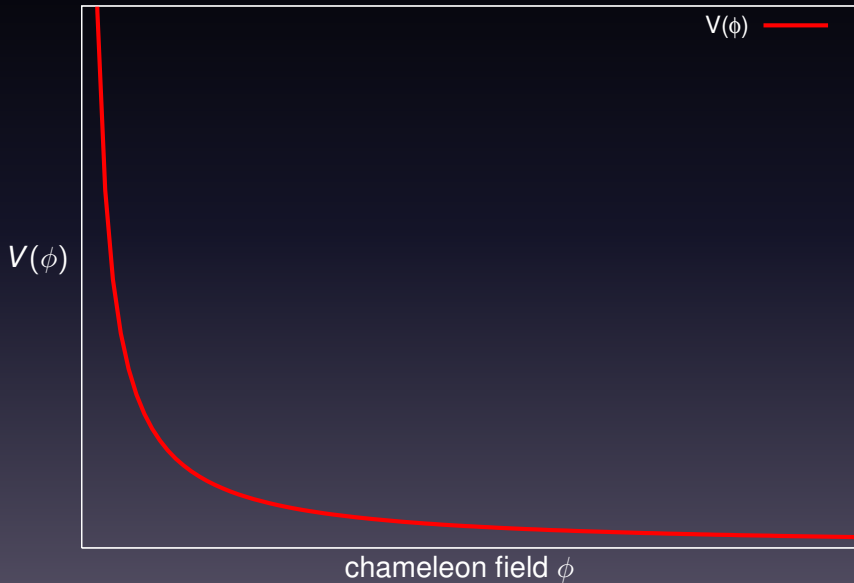
Chameleon screening: Increase m at high ρ .
e.g., $f(R)$ gravity. Reduces force range. Violates Equivalence Principle by making force density/geometry dependent.

Kinetic screening: Increase Z at high ρ .
e.g., Galileons. Deviations from Poisson's equation in high-density regions.

Dynamical coupling: Decrease β at high ρ .
e.g., symmetrons. Uncoupled phase at high densities.

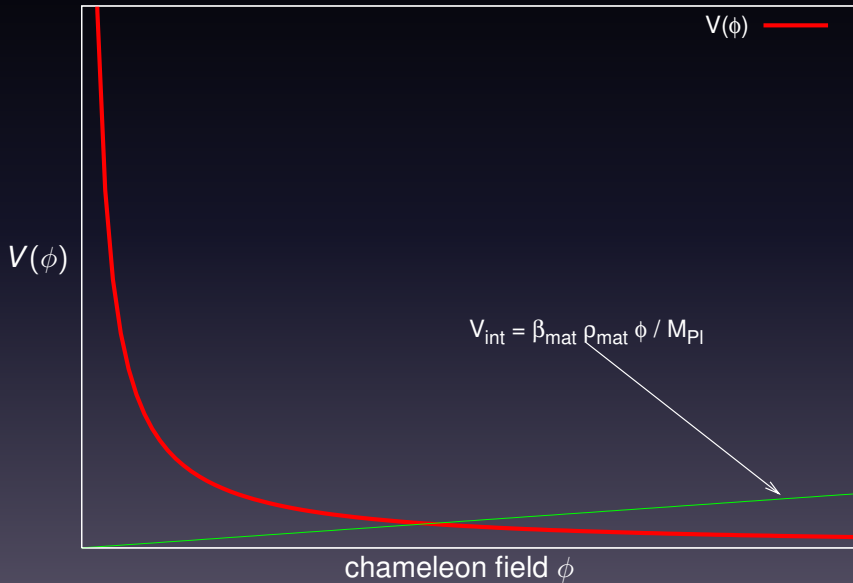
Chameleon mechanism

effective potential: $V_{\text{eff}}(\phi, \rho) = V(\phi) + \beta\rho\phi/M_{\text{Pl}}$



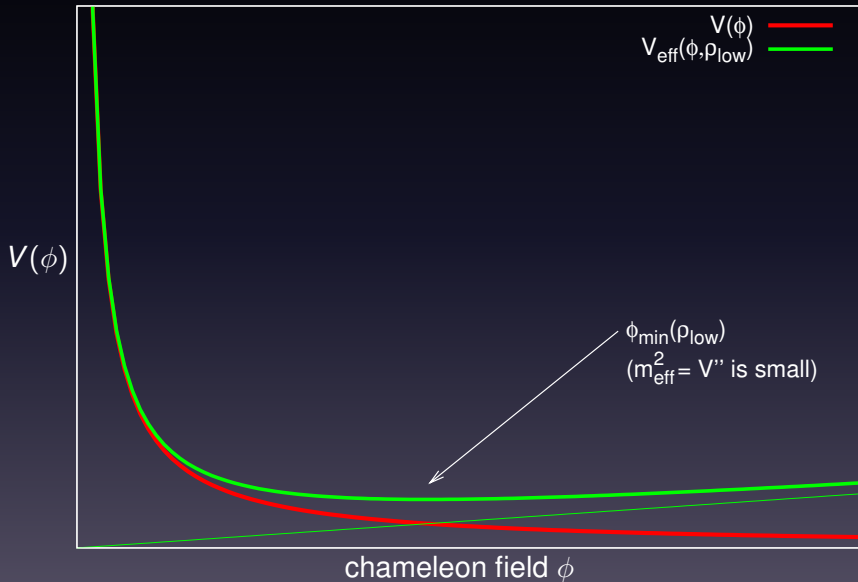
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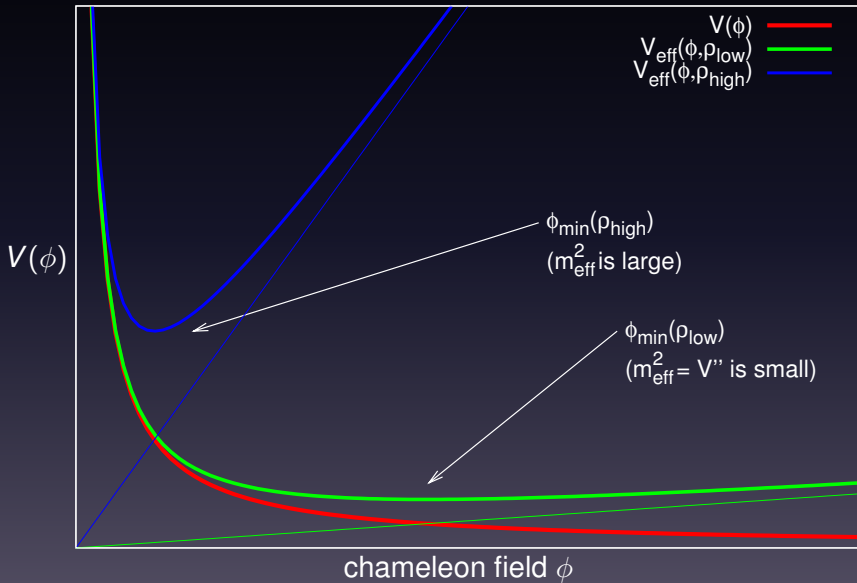
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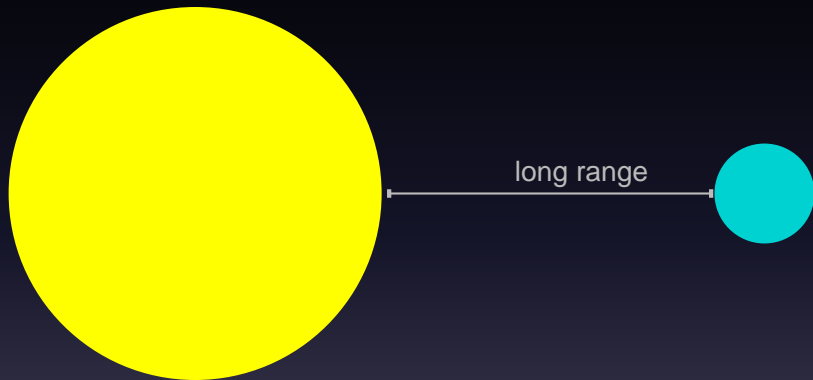


Chameleon mechanism

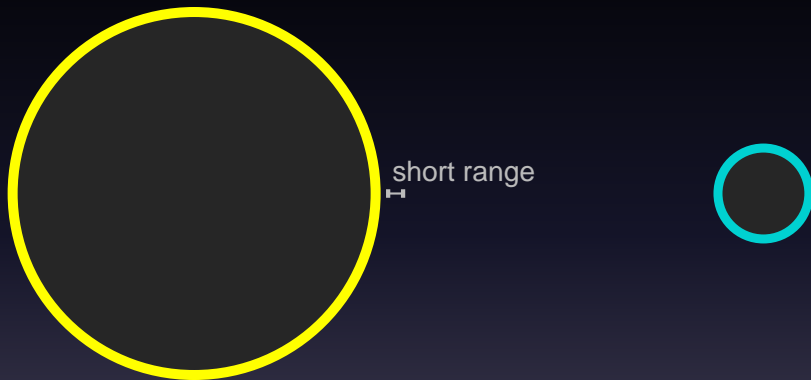
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Effects of a large mass



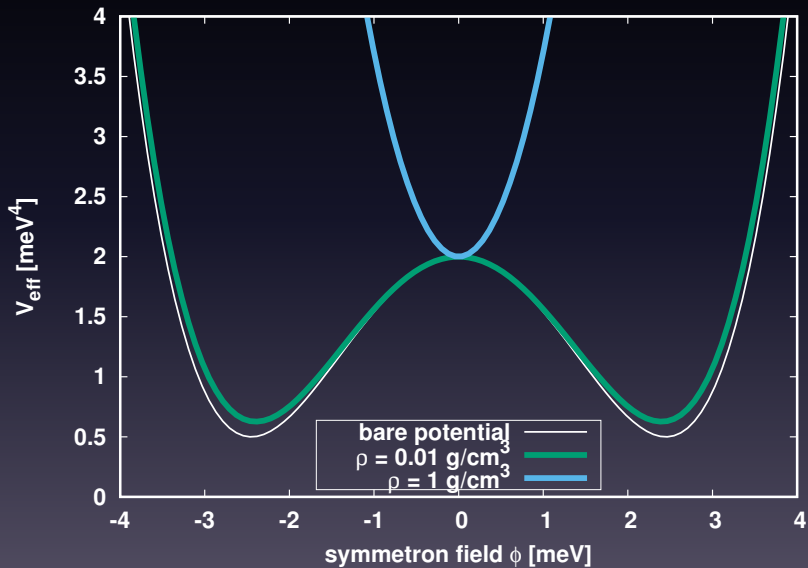
Effects of a large mass



- Compton wavelength (force range) $\propto 1/m$
- Thin-shell effect: ϕ “sees” outer shell of thickness $\sim 1/m$.

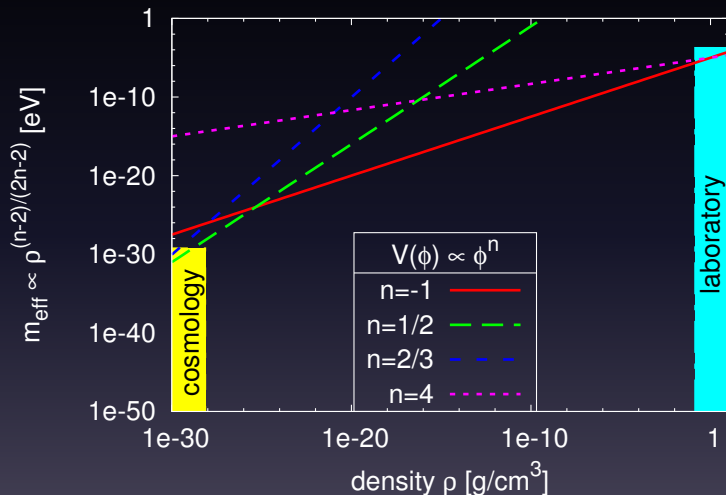
Symmetron screening

$$V_{\text{eff}} = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{\lambda}{4!} \phi^4$$



At which scale should we probe each model?

$$V(\phi) \propto \phi^n + \text{const.} \Rightarrow m_{\text{eff}} \propto \rho^{\frac{n-2}{2n-2}} \quad (\text{use lab for } n \lesssim -\frac{1}{2}, n > 2)$$



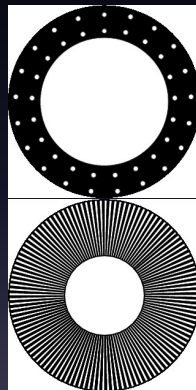
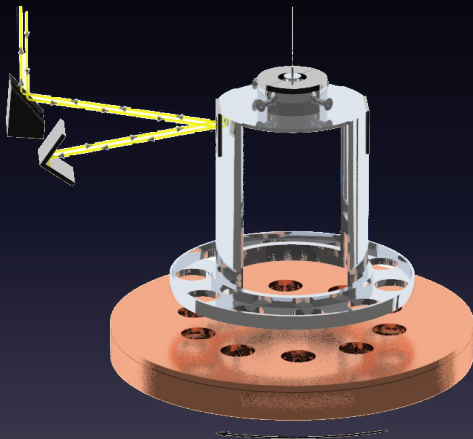
Finding hidden fifth forces in the lab



- 1 force with m_{eff} best detected using objects of size $1/m_{\text{eff}}$
- 2 barrier between source and test masses weakens fifth force
- 3 chamber walls and surroundings contribute to screening
 - chameleons: consider effects of Earth, solar system, galaxy
 - symmetrons: no forces unless chamber radius $\gg \pi/\mu$

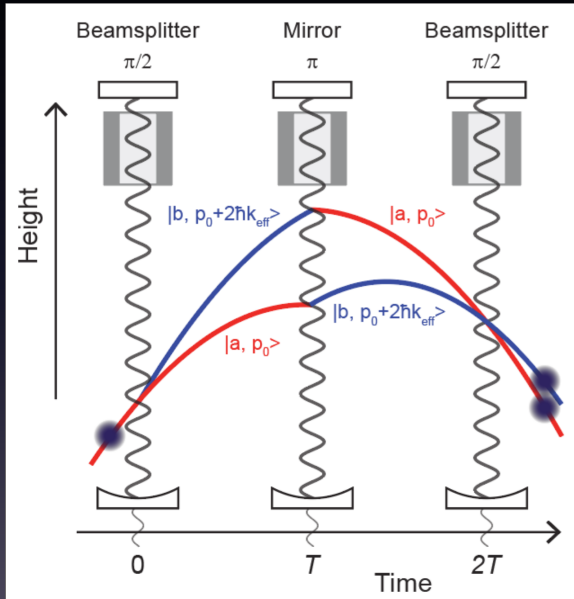
Fifth force tests using a torsion pendulum

Eöt-Wash Experiment

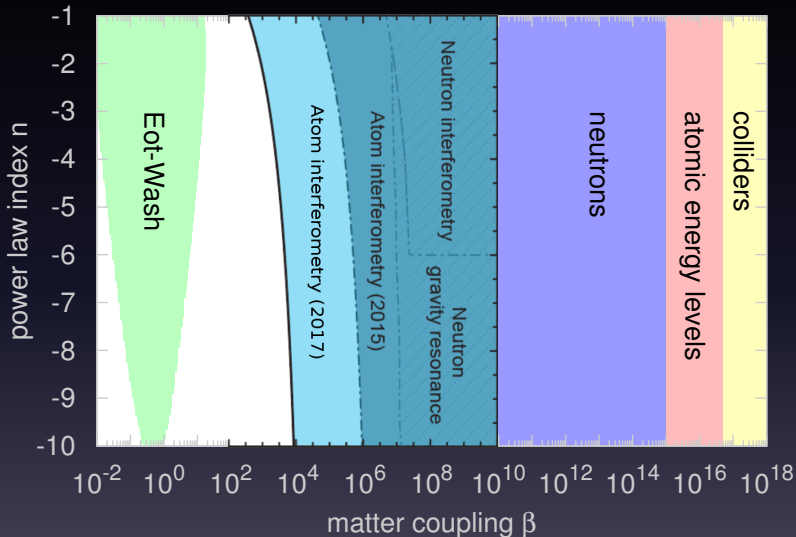


<http://www.npl.washington.edu/eotwash>

Fifth force tests using an atom interferometer



Constraints on $V(\phi) = M_\Lambda^{4-n} \phi^n$ ($M_\Lambda = 2.4$ meV)

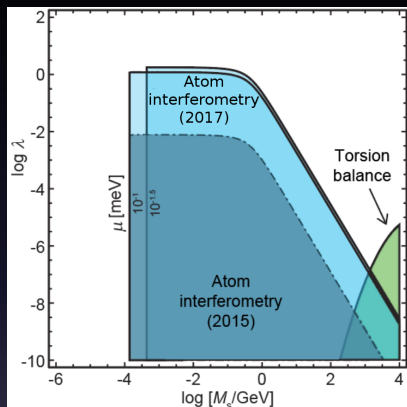
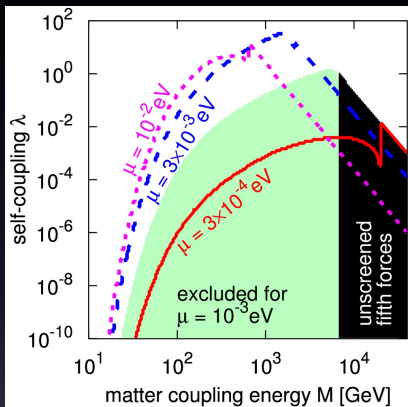


AU, PRD **86**:102003(2012)[*arXiv*:1209.0211]

Adelberger, Heckel, Hoedl, Hoyle, Kapner, AU, PRL **98**:131104(2007)[*arXiv*:hep-ph/0611223]

Jaffe, Haslinger, Xu, Hamilton, AU, Elder, Khoury, Müller, Nat. Phys. **13**:938(2017)[*arXiv*:1612.05171]

Constraints on symmetrons



Symmetron effective potential: $V_{\text{eff}} = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{\lambda}{4!} \phi^4$

Eöt-Wash probes $\lambda \sim 1$, $\mu \sim 10^{-3} \text{ eV}$ (dark energy),

$M \sim 1 \text{ TeV}$ (beyond the Standard Model)

Interferometers extend constraints to lower μ , M

AU, PRL 110:031301(2013)[arXiv:1210.7804]

Jaffe, Haslinger, Xu, Hamilton, AU, Elder, Khoury, Müller, Nat. Phys. 13:938(2017)[arXiv:1612.05171]

Conclusions

- 1 Astrophysics suggests two extensions to Λ : rolling models and conformally-coupled models.
- 2 Chameleons and symmetrons are conformally-coupled models whose fifth forces are screened at high densities.
- 3 Laboratory experiments are close to the dark energy scale $(8\pi G_N/\Lambda)^{1/4} \approx 80 \mu\text{m}$ and are the best ways to probe certain classes of chameleon and symmetron models.
- 4 Carefully-designed torsion pendulum experiments and atom interferometers have excluded orders of magnitude in parameter space.