

# SARAH (part II): Working in Mathematica

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- In this session we want to obtain **analytical information about a given model** using SARAH and Mathematica, i.e.
  - Tree-level masses and tadpoles
  - Tree-level vertices
  - RGEs

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- It is shown how the information is exported into  $\text{\LaTeX}$

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  - RGEs
- It is shown how the information is exported into  $\text{\LaTeX}$

**We use as showcase the SMSSM which we have 'implemented' in the previous session.**

# Loading and running SARAH

SARAH can easily be loaded in Mathematica

---

```
<<"SARAH-4.12.2/SARAH.m";
```

---

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```
<<"SARAH-4.12.2/SARAH.m";
```

---

```
In[2]:= <<"SARAH-4.12.2/SARAH.m"
```

**SARAH 4.12.2**

by Florian Staub, 2017

contributions by M. D. Goodsell, K. Nickel

References:

Comput.Phys.Commun.181 (2010) 1077-1086. (arXiv:0909.2863[hep-ph])

Comput.Phys.Commun.182 (2011) 808-833. (arXiv:1002.0840[hep-ph])

Comput.Phys.Commun.184 (2013) 1792-1809. (arXiv:1207.0906[hep-ph])

Comput.Phys.Commun.185 (2014) 1773-1790. (arXiv:1309.7223[hep-ph])

Download and Documentation:

<http://sarah.hepforge.org>

Start evaluation of a model with:

`Start["Name of Model"]`

e.g. `Start["MSSM"]` or `Start["NMSSM","CKM"]`

To get a list with all installed models, use `ShowModels`

# Loading and running SARAH

A model is initialised afterwards via `Start ["MODEL"]`

---

```
Mathematica
```

---

```
Start ["SMSSM"] ;
```

---

# Loading and running SARAH

A model is initialised afterwards via `Start ["MODEL"]`

---

Mathematica

---

```
Start ["SMSSM"] ;
```

---

```
In[3]:= Start["SMSSM"]
Preparing arrays
... checking Directory: /home/fnstaub/Desktop/HEP-Tools/SARAH-4.12.2/Models/
Model files loaded
  Model      : SMSSM
  Author(s)  : F.Staub
  Date       : 2012-09-01
.....
Loading Susyno functions for the handling of Lie Groups
Based on Susyno v.2.0 by Renato Fonseca (1106.5016)
webpage: web.ist.ufl.pt/renato.fonseca/susyno.html
.....
Initialization
Checking model files: All files okay

...

All Done. SMSSM is ready:
(Model initialized in 25.052s)

Are you unfamiliar with SARAH? Use SARAH FirstSteps
```



During the **initialisation**, the following happens **automatically**:

- The model is checked for anomalies, charge conservation, etc.
- All gauge interactions are derived
- The Lagrangian for component fields is derived from superpotential
- The soft-breaking terms are added
- All field rotations are performed
- The gauge fixing terms are derived, ghost interactions are calculated
- The mass matrices and tadpole equations are derived

During the **initialisation**, the following happens **automatically**:

- The model is **checked for anomalies, charge conservation**, etc.
- All **gauge interactions** are derived
- The **Lagrangian for component fields** is derived from superpotential
- The **soft-breaking terms** are added
- All **field rotations** are performed
- The **gauge fixing terms** are derived, **ghost interactions** are calculated
- The **mass matrices and tadpole equations** are derived

It takes less than a minute before the model is ready and we can start playing

# Particles

In order to see all particles of the current model for a given set of eigenstates, use

---

```
Particles[ STATES ] Mathematica
```

---

**STATES** is the name of the eigenstates.

For each fields the given information is

---

```
{Field, First Generation, Last Generation, Indices }
```

---

In order to see all particles of the current model for a given set of eigenstates, use

---

Mathematica

---

```
Particles[ STATES ]
```

---

**STATES** is the name of the eigenstates.

For each fields the given information is

---

```
{Field, First Generation, Last Generation, Indices }
```

---

## Example:

```
In[4]= Particles[EWSB]
```

```
Out[4]= {{aB, 1, 1, A, {}}, {aWB, 1, 3, A, {{generation, 3}}, {VG, 1, 1, V, {{color, 8}, {lorentz, 4}},  
{gG, 1, 1, G, {{color, 8}}, {fG, 1, 1, F, {{color, 8}}, {aG, 1, 1, A, {{color, 8}},  
{AdL, 1, 3, A, {{generation, 3}, {color, 3}}, {AuL, 1, 3, A, {{generation, 3}, {color, 3}},  
{AeL, 1, 3, A, {{generation, 3}}, {FvL, 1, 3, F, {{generation, 3}}, {AvL, 1, 3, A, {{generation, 3}},  
{AHd0, 1, 1, A, {}}, {AHdm, 1, 1, A, {}}, {AHu0, 1, 1, A, {}}, {AHup, 1, 1, A, {}},  
{AdR, 1, 3, A, {{generation, 3}, {color, 3}}, {AuR, 1, 3, A, {{generation, 3}, {color, 3}},  
{AeR, 1, 3, A, {{generation, 3}}, {AsR, 1, 1, A, {}}, {vd, 1, 1, VEV, {}}, {vu, 1, 1, VEV, {}},  
{vS, 1, 1, VEV, {}}, {VP, 1, 1, V, {{lorentz, 4}}, {VZ, 1, 1, V, {{lorentz, 4}},  
{gP, 1, 1, G, {}}, {gZ, 1, 1, G, {}}, {Vwm, 1, 1, V, {{lorentz, 4}}, {gWm, 1, 1, G, {}},  
{gWmC, 1, 1, G, {}}, {Sd, 1, 6, S, {{generation, 6}, {color, 3}}, {Sv, 1, 3, S, {{generation, 3}},
```

```
...
```

# Parameters

In order to see all parameter of the current model, use

---

```
parameters
```

---

For each parameter the given information is

---

```
{Parameter, Indices, Index Ranges }
```

---

In order to see all parameter of the current model, use

---

`parameters`

---

For each parameter the given information is

{Parameter, Indices, Index Ranges }

---

## Example:

```
In[5]:= parameters
```

```
Out[5]:= {{g1, {}, {}}, {g2, {}, {}}, {g3, {}, {}}, {L1, {}, {}}, {L[L1], {}, {}}, {μ, {}, {}}, {B[μ], {}, {}}, {MS, {}, {}},  
{B[MS], {}, {}}, {Yd, {generation, generation}, {3, 3}}, {T[Yd], {generation, generation}, {3, 3}},  
{Ye, {generation, generation}, {3, 3}}, {T[Ye], {generation, generation}, {3, 3}}, {λ, {}, {}}, {T[λ], {}, {}},  
{κ, {}, {}}, {T[κ], {}, {}}, {Yu, {generation, generation}, {3, 3}}, {T[Yu], {generation, generation}, {3, 3}},  
{mq2, {generation, generation}, {3, 3}}, {ml2, {generation, generation}, {3, 3}},  
{mHd2, {}, {}}, {mHu2, {}, {}}, {md2, {generation, generation}, {3, 3}},  
{mu2, {generation, generation}, {3, 3}}, {me2, {generation, generation}, {3, 3}},  
{ms2, {}, {}}, {MassB, {}, {}}, {MassWB, {}, {}}, {MassG, {}, {}}, {vd, {}, {1}}, {vu, {}, {1}},  
{vS, {}, {1}}, {ZZ, {generation, generation}, {2, 2}}, {ZW, {generation, generation}, {2, 2}},  
{Zfw, {generation, generation}, {3, 3}}, {PhaseGlu, {}, {}}, {ZD, {generation, generation}, {6, 6}},  
{ZV, {generation, generation}, {3, 3}}, {ZU, {generation, generation}, {6, 6}},  
{ZE, {generation, generation}, {6, 6}}, {ZH, {generation, generation}, {3, 3}},  
{ZA, {generation, generation}, {3, 3}}, {ZP, {generation, generation}, {2, 2}},  
...
```

# Mass Matrices and Tadpoles

# Mass Matrix

The tree-level mass matrix is given by

---

```
Mathematica
```

---

```
MassMatrix[ FIELD ];
```

---

with `FIELD` is the name of the mass eigenstates.



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---

Mathematica

---

```
MassMatrix[ FIELD ];
```

---

with **FIELD** is the name of the mass eigenstates.

## Examples:

The mass matrix of the CP-even Higgs states is given by:

```
In[34]:= MassMatrix[hh] // FullSimplify
```

```
Out[34]:= {{ 1/8 (8 mHd2 + (g1^2 + g2^2) (3 vd^2 - vu^2) + 4 ((vS^2 + vu^2) λ + √2 vS μ) conj[λ] + 4 (√2 vS λ + 2 μ) conj[μ]),  
1/4 (- (g1^2 + g2^2) vd vu - 2 B[μ] - 2 λ conj[L1] - 2 L1 conj[λ] + 4 vd vu λ conj[λ] - vS^2 (λ conj[κ] + κ conj[λ]) -  
2 conj[B[μ]] - √2 vS (λ conj[MS] + MS conj[λ] + conj[T[λ]] + T[λ])), vd (vS λ + μ/√2) conj[λ] + vd λ conj[μ] / √2 -  
1/4 vu (√2 λ conj[MS] + 2 vS (λ conj[κ] + κ conj[λ]) + √2 (MS conj[λ] + conj[T[λ]] + T[λ]))},  
... }
```

The tree-level mass matrix is given by

```
_____ Mathematica _____  
MassMatrix[ FIELD ] ;  
_____
```

with **FIELD** is the name of the mass eigenstates.

## Examples:

The down-quark mass matrix is given by:

```
In[7]:= MassMatrix[Fd] // MatrixForm  
Out[7]//MatrixForm=
```

$$\begin{pmatrix} \frac{vd \ Yd[1,1]}{\sqrt{2}} & \frac{vd \ Yd[2,1]}{\sqrt{2}} & \frac{vd \ Yd[3,1]}{\sqrt{2}} \\ \frac{vd \ Yd[1,2]}{\sqrt{2}} & \frac{vd \ Yd[2,2]}{\sqrt{2}} & \frac{vd \ Yd[3,2]}{\sqrt{2}} \\ \frac{vd \ Yd[1,3]}{\sqrt{2}} & \frac{vd \ Yd[2,3]}{\sqrt{2}} & \frac{vd \ Yd[3,3]}{\sqrt{2}} \end{pmatrix}$$

For fermions names of Weyl or Dirac spinors can be used, i.e.

**MassMatrix[FDL]** or **MassMatrix[FDR]** give the same output.

The tree-level mass matrix is given by

---

Mathematica

---

```
MassMatrix[ FIELD ];
```

---

with **FIELD** is the name of the mass eigenstates.

## Examples:

The neutralino and chargino mass matrices are given by:

```
In[9]= MatrixForm /@ MassMatrix /@ {Cha, Chi}
```

$$\text{Out[9]} = \left\{ \begin{array}{c} \left( \begin{array}{cc} \text{MassWB} & \frac{g2 \text{vu}}{\sqrt{2}} \\ \frac{g2 \text{vd}}{\sqrt{2}} & \frac{\text{vS} \lambda}{\sqrt{2}} + \mu \end{array} \right), \left( \begin{array}{ccccc} \text{MassB} & 0 & -\frac{g1 \text{vd}}{2} & \frac{g1 \text{vu}}{2} & 0 \\ 0 & \text{MassWB} & \frac{g2 \text{vd}}{2} & -\frac{g2 \text{vu}}{2} & 0 \\ -\frac{g1 \text{vd}}{2} & \frac{g2 \text{vd}}{2} & 0 & -\frac{\text{vS} \lambda}{\sqrt{2}} - \mu & -\frac{\text{vu} \lambda}{\sqrt{2}} \\ \frac{g1 \text{vu}}{2} & -\frac{g2 \text{vu}}{2} & -\frac{\text{vS} \lambda}{\sqrt{2}} - \mu & 0 & -\frac{\text{vd} \lambda}{\sqrt{2}} \\ 0 & 0 & -\frac{\text{vu} \lambda}{\sqrt{2}} & -\frac{\text{vd} \lambda}{\sqrt{2}} & \text{MS} + \sqrt{2} \text{vS} \kappa \end{array} \right) \end{array} \right\}$$

The tree-level mass matrix is given by

---

```
Mathematica
```

---

```
MassMatrix[ FIELD ];
```

---

with **FIELD** is the name of the mass eigenstates.

## Examples:

Mass matrices involving **Goldstone bosons** come with **RXi** what denotes the **gauge dependent part** from the gauge fixing Lagrangian.

- **RXi [ \_ ]**  $\rightarrow 0$  corresponds to **Landau gauge**
- **RXi [ \_ ]**  $\rightarrow 1$  corresponds to **Feynman gauge**

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MassMatrix[ FIELD ];
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## Examples:

Mass matrices involving **Goldstone bosons** come with **RXi** what denotes the **gauge dependent part** from the gauge fixing Lagrangian.

- **RXi [ \_ ]** -> 0 corresponds to **Landau gauge**
- **RXi [ \_ ]** -> 1 corresponds to **Feynman gauge**

For instance, the **CP-odd Higgs** mass matrix reads

```
In[13]:= MassMatrix [ Ah ]
```

$$\text{Out[13]} = \left\{ \left\{ m_{Hd2} + \frac{g_1^2 v d^2}{8} + \frac{g_2^2 v d^2}{8} - \frac{g_1^2 v u^2}{8} - \frac{g_2^2 v u^2}{8} + \frac{1}{2} v s^2 \lambda \text{conj}[\lambda] + \right. \right.$$
$$\left. \frac{1}{2} v u^2 \lambda \text{conj}[\lambda] + \frac{v s \mu \text{conj}[\lambda]}{\sqrt{2}} + \frac{v s \lambda \text{conj}[\mu]}{\sqrt{2}} + \mu \text{conj}[\mu] + \frac{1}{4} g_2^2 v d^2 \text{Cos}[\text{ThetaW}] \text{Rxi}[Z] + \right.$$

...

The tree-level mass matrix is given by

```
_____ Mathematica _____  
MassMatrix[ FIELD ] ;  
_____
```

with **FIELD** is the name of the mass eigenstates.

## Examples:

For vector bosons the name of one of the mass eigenstate can be used, i.e. `MassMatrix[VP]` and `MassMatrix[VZ]` give the same result:

```
In[10]:= MassMatrix [VP] // MatrixForm
```

```
Out[10]//MatrixForm=
```

$$\begin{pmatrix} \frac{g_1^2 v d^2}{4} + \frac{g_1^2 v u^2}{4} & -\frac{1}{4} g_1 g_2 v d^2 - \frac{1}{4} g_1 g_2 v u^2 \\ -\frac{1}{4} g_1 g_2 v d^2 - \frac{1}{4} g_1 g_2 v u^2 & \frac{g_2^2 v d^2}{4} + \frac{g_2^2 v u^2}{4} \end{pmatrix}$$

The tree-level mass matrix is given by

---

```
MassMatrix[ FIELD ];
```

---

with **FIELD** is the name of the mass eigenstates.

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`MassMatrix[VP]` and `MassMatrix[VZ]` give the same result:

```
In[10]:= MassMatrix [VP] // MatrixForm
```

```
Out[10]/MatrixForm=
```

$$\begin{pmatrix} \frac{g_1^2 v d^2}{4} + \frac{g_1^2 v u^2}{4} & -\frac{1}{4} g_1 g_2 v d^2 - \frac{1}{4} g_1 g_2 v u^2 \\ -\frac{1}{4} g_1 g_2 v d^2 - \frac{1}{4} g_1 g_2 v u^2 & \frac{g_2^2 v d^2}{4} + \frac{g_2^2 v u^2}{4} \end{pmatrix}$$

one can check immediately that one state is massless

```
In[35]:= MassMatrix [VP];
```

```
Eigenvalues [%]
```

```
Out[36]= {0, (g1^2 + g2^2) (vd^2/4 + vu^2/4)}
```

The tree-level mass matrix is given by

```
Mathematica  
-----  
MassMatrix[ FIELD ];  
-----
```

with **FIELD** is the name of the mass eigenstates.

## Examples:

The **sfermion mass matrices** are general  $6 \times 6$  matrices which are a bit lengthy, i.e. the down squark matrix reads

```
In[12]:= MassMatrix[Sd]  
Out[12]= {{ -1/24 g1^2 vd^2 - g2^2 vd^2/8 + g1^2 vu^2/24 + g2^2 vu^2/8 + mq2[1,1] + 1/2 vd^2 sum[j1,1,3, conj[Yd[j1,1]] Yd[j1,1]],  
mq2[1,2] + 1/2 vd^2 sum[j1,1,3, conj[Yd[j1,1]] Yd[j1,2]],  
mq2[1,3] + 1/2 vd^2 sum[j1,1,3, conj[Yd[j1,1]] Yd[j1,3]],  
-1/2 vS vu lambda conj[Yd[1,1]] - vu mu conj[Yd[1,1]]/sqrt(2) + vd conj[T[Yd][1,1]]/sqrt(2),  
... }
```



# Tadpole Equations I

The tadpole equations corresponding to a scalar or VEV is returned by

---

```
TadpoleEquation[ X ];
```

---

with  $X$  is the name of VEV (or the corresponding field).

# Tadpole Equations I

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```
TadpoleEquation[ X ];
```

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## Examples:

The three tadpole equations in the SMSSM read:

```
In[38]:= TadpoleEquation[ vd ]
```

$$\begin{aligned} \text{Out[38]} = & \text{mHd2} \text{ vd} + \frac{g1^2 \text{ vd}^3}{8} + \frac{g2^2 \text{ vd}^3}{8} - \frac{1}{8} g1^2 \text{ vd} \text{ vu}^2 - \frac{1}{8} g2^2 \text{ vd} \text{ vu}^2 - \frac{1}{2} \text{vu} \text{B}[\mu] - \\ & \frac{1}{2} \text{vu} \lambda \text{ conj}[\text{L1}] - \frac{\text{vS} \text{vu} \lambda \text{ conj}[\text{MS}]}{2 \sqrt{2}} - \frac{1}{4} \text{vS}^2 \text{vu} \lambda \text{ conj}[\kappa] - \frac{1}{2} \text{L1} \text{vu} \text{ conj}[\lambda] - \frac{\text{MS} \text{vS} \text{vu} \text{ conj}[\lambda]}{2 \sqrt{2}} - \\ & \frac{1}{4} \text{vS}^2 \text{vu} \kappa \text{ conj}[\lambda] + \frac{1}{2} \text{vd} \text{vS}^2 \lambda \text{ conj}[\lambda] + \frac{1}{2} \text{vd} \text{vu}^2 \lambda \text{ conj}[\lambda] + \frac{\text{vd} \text{vS} \mu \text{ conj}[\lambda]}{\sqrt{2}} + \\ & \frac{\text{vd} \text{vS} \lambda \text{ conj}[\mu]}{\sqrt{2}} + \text{vd} \mu \text{ conj}[\mu] - \frac{1}{2} \text{vu} \text{ conj}[\text{B}[\mu]] - \frac{\text{vS} \text{vu} \text{ conj}[\text{T}[\lambda]]}{2 \sqrt{2}} - \frac{\text{vS} \text{vu} \text{T}[\lambda]}{2 \sqrt{2}} = 0 \end{aligned}$$

# Tadpole Equations I

The tadpole equations corresponding to a scalar or VEV is returned by

```
TadpoleEquation[ X ];
```

with  $X$  is the name of VEV (or the corresponding field).

## Examples:

The three tadpole equations in the SMSSM read:

```
In[15]:= TadpoleEquation[vu]
```

```
Out[15]= mHu2 vu - 1/8 g1^2 vd^2 vu - 1/8 g2^2 vd^2 vu + g1^2 vu^3/8 + g2^2 vu^3/8 - 1/2 vd B[mu] -
```

$$\frac{1}{2} \text{vd } \lambda \text{ conj}[L1] - \frac{\text{vd } vS \lambda \text{ conj}[MS]}{2\sqrt{2}} - \frac{1}{4} \text{vd } vS^2 \lambda \text{ conj}[\kappa] - \frac{1}{2} L1 \text{vd conj}[\lambda] - \frac{MS \text{vd } vS \text{ conj}[\lambda]}{2\sqrt{2}} -$$
$$\frac{1}{4} \text{vd } vS^2 \kappa \text{ conj}[\lambda] + \frac{1}{2} \text{vd}^2 \text{vu } \lambda \text{ conj}[\lambda] + \frac{1}{2} vS^2 \text{vu } \lambda \text{ conj}[\lambda] + \frac{vS \text{vu } \mu \text{ conj}[\lambda]}{\sqrt{2}} +$$

```
vd vu lambda conj[mu] / sqrt(2) + vu mu conj[mu] - 1/2 vd conj[B[mu]] - vd vS conj[T[lambda]] / (2*sqrt(2)) - vd vS T[lambda] / (2*sqrt(2)) == 0
```

# Tadpole Equations I

The tadpole equations corresponding to a scalar or VEV is returned by

```
_____ Mathematica _____  
TadpoleEquation[ X ];
```

with  $X$  is the name of VEV (or the corresponding field).

## Examples:

The three tadpole equations in the SMSSM read:

```
In[16]:= TadpoleEquation[vS]
```

$$\begin{aligned} \text{Out[16]} = & m_s^2 v_S + \frac{1}{2} v_S B[MS] + \frac{MS \text{ conj}[L1]}{\sqrt{2}} + v_S \kappa \text{ conj}[L1] + \frac{L1 \text{ conj}[MS]}{\sqrt{2}} + MS v_S \text{ conj}[MS] + \\ & \frac{3 v_S^2 \kappa \text{ conj}[MS]}{2 \sqrt{2}} - \frac{vd vu \lambda \text{ conj}[MS]}{2 \sqrt{2}} + L1 v_S \text{ conj}[\kappa] + \frac{3 MS v_S^2 \text{ conj}[\kappa]}{2 \sqrt{2}} + v_S^3 \kappa \text{ conj}[\kappa] - \\ & \frac{1}{2} vd v_S vu \lambda \text{ conj}[\kappa] - \frac{MS vd vu \text{ conj}[\lambda]}{2 \sqrt{2}} - \frac{1}{2} vd v_S vu \kappa \text{ conj}[\lambda] + \frac{1}{2} vd^2 v_S \lambda \text{ conj}[\lambda] + \\ & \frac{1}{2} v_S vu^2 \lambda \text{ conj}[\lambda] + \frac{vd^2 \mu \text{ conj}[\lambda]}{2 \sqrt{2}} + \frac{vu^2 \mu \text{ conj}[\lambda]}{2 \sqrt{2}} + \frac{vd^2 \lambda \text{ conj}[\mu]}{2 \sqrt{2}} + \frac{vu^2 \lambda \text{ conj}[\mu]}{2 \sqrt{2}} + \frac{1}{2} v_S \text{ conj}[B[MS]] + \\ & \frac{\text{conj}[L[L1]]}{\sqrt{2}} + \frac{v_S^2 \text{ conj}[T[\kappa]]}{2 \sqrt{2}} - \frac{vd vu \text{ conj}[T[\lambda]]}{2 \sqrt{2}} + \frac{L[L1]}{\sqrt{2}} + \frac{v_S^2 T[\kappa]}{2 \sqrt{2}} - \frac{vd vu T[\lambda]}{2 \sqrt{2}} = 0 \end{aligned}$$

# Tadpole Equations I

The tadpole equations corresponding to a scalar or VEV is returned by

---

```
TadpoleEquation[ X ];
```

---

with  $X$  is the name of VEV (or the corresponding field).

## Examples:

The three tadpole equations in the SMSSM read:

The same results are obtained from

- `TadpoleEquation[phid]`
- `TadpoleEquation[phiu]`
- `TadpoleEquation[phiS]`

# Tadpole Equations II

A list with all tadpole equations for a given set of eigenstates is stored in

---

```
TadpoleEquations[ STATES ];
```

---

with **STATES** is the name of eigenstates.

# Tadpole Equations II

A list with all tadpole equations for a given set of eigenstates is stored in

---

TadpoleEquations [ STATES ] ;

---

with STATES is the name of eigenstates.

**Example:** The three tadpole equations after EWSB are stored in:

In[17]:= TadpoleEquations [EWSB]

Out[17]:= 
$$\left\{ \begin{aligned} & \text{mHd2} \text{ vd} + \frac{g_1^2 \text{ vd}^3}{8} + \frac{g_2^2 \text{ vd}^3}{8} - \frac{1}{8} g_1^2 \text{ vd vu}^2 - \frac{1}{8} g_2^2 \text{ vd vu}^2 - \frac{1}{2} \text{vu B}[\mu] - \\ & \quad \dots \\ & \frac{\text{vd vS} \lambda \text{ conj}[\mu]}{\sqrt{2}} + \text{vd} \mu \text{ conj}[\mu] - \frac{1}{2} \text{vu conj}[\text{B}[\mu]] - \frac{\text{vS vu conj}[\text{T}[\lambda]]}{2\sqrt{2}} - \frac{\text{vS vu T}[\lambda]}{2\sqrt{2}}, \\ & \text{mHu2 vu} - \frac{1}{8} g_1^2 \text{ vd}^2 \text{vu} - \frac{1}{8} g_2^2 \text{ vd}^2 \text{vu} + \frac{g_1^2 \text{vu}^3}{8} + \frac{g_2^2 \text{vu}^3}{8} - \frac{1}{2} \text{vd B}[\mu] - \frac{1}{2} \text{vd} \lambda \text{ conj}[\text{L1}] - \\ & \quad \dots \\ & \frac{\text{vS vu} \lambda \text{ conj}[\mu]}{\sqrt{2}} + \text{vu} \mu \text{ conj}[\mu] - \frac{1}{2} \text{vd conj}[\text{B}[\mu]] - \frac{\text{vd vS conj}[\text{T}[\lambda]]}{2\sqrt{2}} - \frac{\text{vd vS T}[\lambda]}{2\sqrt{2}}, \\ & \text{ms2 vS} + \frac{1}{2} \text{vS B}[\text{MS}] + \frac{\text{MS conj}[\text{L1}]}{\sqrt{2}} + \text{vS} \kappa \text{ conj}[\text{L1}] + \frac{\text{L1 conj}[\text{MS}]}{\sqrt{2}} + \text{MS vS conj}[\text{MS}] + \frac{3 \text{vS}^2 \kappa \text{ conj}[\text{MS}]}{2\sqrt{2}} - \\ & \quad \dots \\ & \frac{\text{conj}[\text{L}[\text{L1}]]}{\sqrt{2}} + \frac{\text{vS}^2 \text{ conj}[\text{T}[\kappa]]}{2\sqrt{2}} - \frac{\text{vd vu conj}[\text{T}[\lambda]]}{2\sqrt{2}} + \frac{\text{L}[\text{L1}]}{\sqrt{2}} + \frac{\text{vS}^2 \text{T}[\kappa]}{2\sqrt{2}} - \frac{\text{vd vu T}[\lambda]}{2\sqrt{2}} \end{aligned} \right\}$$

# Application: Pseudo-Scalar Mass

One can **combine** the information of the **mass matrix with the tadpole equations** to check the **masses of the Goldstones**:



# Application: Pseudo-Scalar Mass

One can combine the information of the mass matrix with the tadpole equations to check the masses of the Goldstones:

---

```
solTadpoles = Mathematica Solve[TadpoleEquations[EWSB] == 0,  
Simplify[Eigenvalues[MassMatrix[Ah] /. solTadpoles]]]
```

---

- 1 We use the Solve command of Mathematica to get the solutions of the tadpole equations for  $m_{H_d}^2$ ,  $m_{H_u}^2$ ,  $m_S^2$
- 2 We insert the solutions in the pseudo-scalar mass matrix
- 3 We use Eigenvalues to get CP odd masses

# Application: Pseudo-Scalar Mass

One can combine the information of the mass matrix with the tadpole equations to check the masses of the Goldstones:

```
In[39]:= solTadpoles = Solve[TadpoleEquations[EWSB] == 0, {mHd2, mHu2, ms2}][[1]];
Simplify[Eigenvalues[MassMatrix[Ah] /. solTadpoles]]
```

$$\text{Out[40]} = \left\{ \frac{1}{4} (v_d^2 + v_u^2) \text{Rxi}[Z] (g_2 \cos[\text{ThetaW}] + g_1 \sin[\text{ThetaW}])^2, \frac{1}{8 v_d v_u v_u} \right. \\
\left. \begin{aligned} & (-4 v_d v_u v_u B[MS] + 2 v_d^2 v_s B[\mu] + 2 v_s v_u^2 B[\mu] - 2 \sqrt{2} MS v_d v_u \text{conj}[L1] - 8 v_d v_s v_u \kappa \text{conj}[L1] + \\ & 2 v_d^2 v_s \lambda \text{conj}[L1] + 2 v_s v_u^2 \lambda \text{conj}[L1] - 2 \sqrt{2} L1 v_d v_u \text{conj}[MS] - \sqrt{2} v_d v_s^2 v_u \kappa \text{conj}[MS] + \\ & \sqrt{2} v_d^2 v_s^2 \lambda \text{conj}[MS] + \sqrt{2} v_d^2 v_u^2 \lambda \text{conj}[MS] + \sqrt{2} v_s^2 v_u^2 \lambda \text{conj}[MS] - 8 L1 v_d v_s v_u \text{conj}[\kappa] - \\ & \sqrt{2} MS v_d v_s^2 v_u \text{conj}[\kappa] + v_d^2 v_s^3 \lambda \text{conj}[\kappa] + 4 v_d^2 v_s v_u^2 \lambda \text{conj}[\kappa] + v_s^3 v_u^2 \lambda \text{conj}[\kappa] + \\ & 2 L1 v_d^2 v_s \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_s^2 \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_u^2 \text{conj}[\lambda] + 2 L1 v_s v_u^2 \text{conj}[\lambda] + \\ & \sqrt{2} MS v_s^2 v_u^2 \text{conj}[\lambda] + v_d^2 v_s^3 \kappa \text{conj}[\lambda] + 4 v_d^2 v_s v_u^2 \kappa \text{conj}[\lambda] + v_s^3 v_u^2 \kappa \text{conj}[\lambda] - \\ & \sqrt{2} v_d^3 v_u \mu \text{conj}[\lambda] - \sqrt{2} v_d v_u^3 \mu \text{conj}[\lambda] - \sqrt{2} v_d^3 v_u \lambda \text{conj}[\mu] - \sqrt{2} v_d v_u^3 \lambda \text{conj}[\mu] - \\ & 4 v_d v_s v_u \text{conj}[B[MS]] + 2 v_d^2 v_s \text{conj}[B[\mu]] + 2 v_s v_u^2 \text{conj}[B[\mu]] - 2 \sqrt{2} v_d v_u \text{conj}[L[L1]] - \\ & 3 \sqrt{2} v_d v_s^2 v_u \text{conj}[T[\kappa]] + \sqrt{2} v_d^2 v_s^2 \text{conj}[T[\lambda]] + \sqrt{2} v_d^2 v_u^2 \text{conj}[T[\lambda]] + \sqrt{2} v_s^2 v_u^2 \text{conj}[T[\lambda]] - \\ & 2 \sqrt{2} v_d v_u L[L1] - 3 \sqrt{2} v_d v_s^2 v_u T[\kappa] + \sqrt{2} v_d^2 v_s^2 T[\lambda] + \sqrt{2} v_d^2 v_u^2 T[\lambda] + \sqrt{2} v_s^2 v_u^2 T[\lambda] - \\ & \sqrt{((-4 v_d v_s v_u B[MS] + 2 v_s (v_d^2 + v_u^2) B[\mu] - 2 \sqrt{2} MS v_d v_u \text{conj}[L1] - 8 v_d v_s v_u \kappa \text{conj}[L1] + 2 v_d^2 v_s \lambda \text{conj}[ \\ & L1] + 2 v_s v_u^2 \lambda \text{conj}[L1] - 2 \sqrt{2} L1 v_d v_u \text{conj}[MS] - \sqrt{2} v_d v_s^2 v_u \kappa \text{conj}[MS] + \sqrt{2} v_d^2 v_s^2 \lambda \text{conj}[MS] + \\ & \sqrt{2} v_d^2 v_u^2 \lambda \text{conj}[MS] + \sqrt{2} v_s^2 v_u^2 \lambda \text{conj}[MS] - 8 L1 v_d v_s v_u \text{conj}[\kappa] - \sqrt{2} MS v_d v_s^2 v_u \text{conj}[\kappa] + \\ & v_d^2 v_s^3 \lambda \text{conj}[\kappa] + 4 v_d^2 v_s v_u^2 \lambda \text{conj}[\kappa] + v_s^3 v_u^2 \lambda \text{conj}[\kappa] + 2 L1 v_d^2 v_s \text{conj}[\lambda] + \\ & \sqrt{2} MS v_d^2 v_s^2 \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_u^2 \text{conj}[\lambda] + 2 L1 v_s v_u^2 \text{conj}[\lambda] + \sqrt{2} MS v_s^2 v_u^2 \text{conj}[\lambda] + \\ & v_d^2 v_s^3 \kappa \text{conj}[\lambda] + 4 v_d^2 v_s v_u^2 \kappa \text{conj}[\lambda] + v_s^3 v_u^2 \kappa \text{conj}[\lambda] - \sqrt{2} v_d^3 v_u \mu \text{conj}[\lambda] - \sqrt{2} v_d v_u^3 \mu \text{conj}[\lambda] - \\ & \sqrt{2} v_d^3 v_u \lambda \text{conj}[\mu] - \sqrt{2} v_d v_u^3 \lambda \text{conj}[\mu] - 4 v_d v_s v_u \text{conj}[B[MS]] + 2 v_d^2 v_s \text{conj}[B[\mu]] + 2 v_s v_u^2 \text{conj}[B[\mu]] - 2 \sqrt{2} v_d v_u \text{conj}[L[L1]] - \\ & 3 \sqrt{2} v_d v_s^2 v_u \text{conj}[T[\kappa]] + \sqrt{2} v_d^2 v_s^2 \text{conj}[T[\lambda]] + \sqrt{2} v_d^2 v_u^2 \text{conj}[T[\lambda]] + \sqrt{2} v_s^2 v_u^2 \text{conj}[T[\lambda]] - \\ & 2 \sqrt{2} v_d v_u L[L1] - 3 \sqrt{2} v_d v_s^2 v_u T[\kappa] + \sqrt{2} v_d^2 v_s^2 T[\lambda] + \sqrt{2} v_d^2 v_u^2 T[\lambda] + \sqrt{2} v_s^2 v_u^2 T[\lambda] - \end{aligned} \right.$$

**Masses and Tadpoles**

# Application: Pseudo-Scalar Mass

One can combine the information of the mass matrix with the tadpole equations to check the masses of the Goldstones:

```
In[39]:= solTadpoles = Solve[TadpoleEquations[EWSB] == 0, {mHd2, mHu2, ms2}][[1]];
Simplify[Eigenvalues[MassMatrix[Ah] /. solTadpoles]]
```

$$\text{Out[39]} = \left\{ \frac{1}{4} (v_d^2 + v_u^2) \text{Rxi}[Z] (g_2 \cos[\text{ThetaW}] + g_1 \sin[\text{ThetaW}])^2, \frac{1}{8 v_d v_s v_u} \right.$$

$$\left. \begin{aligned} & (-4 v_d v_s v_u B[MS] - 2 v_d^2 v_s B[\mu] - 2 v_s v_u^2 B[\mu] - 2 \sqrt{2} MS v_d v_u \text{conj}[L1] - 8 v_d v_s v_u \kappa \text{conj}[L1] + \right. \\ & 2 v_d^2 v_s \lambda \text{conj}[L1] + 2 v_s v_u^2 \lambda \text{conj}[L1] - 2 \sqrt{2} L1 v_d v_u \text{conj}[MS] - \sqrt{2} v_d v_s^2 v_u \kappa \text{conj}[MS] + \\ & \sqrt{2} v_d^2 v_s^2 \lambda \text{conj}[MS] + \sqrt{2} v_d^2 v_u^2 \lambda \text{conj}[MS] + \sqrt{2} v_s^2 v_u^2 \lambda \text{conj}[MS] - 8 L1 v_d v_s v_u \text{conj}[\kappa] - \\ & \sqrt{2} MS v_d v_s^2 v_u \text{conj}[\kappa] + v_d^2 v_s^3 \lambda \text{conj}[\kappa] + 4 v_d^2 v_s v_u^2 \lambda \text{conj}[\kappa] + v_s^3 v_u^2 \lambda \text{conj}[\kappa] + \\ & 2 L1 v_d^2 v_s \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_s^2 \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_u^2 \text{conj}[\lambda] + 2 L1 v_s v_u^2 \text{conj}[\lambda] + \\ & \sqrt{2} MS v_s^2 v_u^2 \text{conj}[\lambda] + v_d^2 v_s^3 \kappa \text{conj}[\lambda] + 4 v_d^2 v_s v_u^2 \kappa \text{conj}[\lambda] + v_s^3 v_u^2 \kappa \text{conj}[\lambda] - \\ & \sqrt{2} v_d^3 v_u \mu \text{conj}[\lambda] - \sqrt{2} v_d v_u^3 \mu \text{conj}[\lambda] - \sqrt{2} v_d^3 v_u \lambda \text{conj}[\mu] - \sqrt{2} v_d v_u^3 \lambda \text{conj}[\mu] - \\ & 4 v_d v_s v_u \text{conj}[B[MS]] + 2 v_d^2 v_s \text{conj}[B[\mu]] + 2 v_s v_u^2 \text{conj}[B[\mu]] - 2 \sqrt{2} v_d v_u \text{conj}[L[L1]] - \\ & 3 \sqrt{2} v_d v_s^2 v_u \text{conj}[T[\kappa]] + \sqrt{2} v_d^2 v_s^2 \text{conj}[T[\lambda]] + \sqrt{2} v_d^2 v_u^2 \text{conj}[T[\lambda]] + \sqrt{2} v_s^2 v_u^2 \text{conj}[T[\lambda]] - \\ & 2 \sqrt{2} v_d v_u L[L1] - 3 \sqrt{2} v_d v_s^2 v_u T[\kappa] + \sqrt{2} v_d^2 v_s^2 T[\lambda] + \sqrt{2} v_d^2 v_u^2 T[\lambda] + \sqrt{2} v_s^2 v_u^2 T[\lambda] - \end{aligned} \right.$$

→ The Goldstone mass squared is  $R_\xi M_Z^2$

$$\left. \begin{aligned} & \sqrt{2} v_d^4 v_u^4 \lambda \text{conj}[MS] + \sqrt{2} v_s^4 v_u^4 \lambda \text{conj}[MS] - 8 L1 v_d v_s v_u \text{conj}[\kappa] - \sqrt{2} MS v_d v_s^4 v_u \text{conj}[\kappa] + \\ & v_d^2 v_s^3 \lambda \text{conj}[\kappa] + 4 v_d^2 v_s v_u^2 \lambda \text{conj}[\kappa] + v_s^3 v_u^2 \lambda \text{conj}[\kappa] + 2 L1 v_d^2 v_s \text{conj}[\lambda] + \\ & \sqrt{2} MS v_d^2 v_s^2 \text{conj}[\lambda] + \sqrt{2} MS v_d^2 v_u^2 \text{conj}[\lambda] + 2 L1 v_s v_u^2 \text{conj}[\lambda] + \sqrt{2} MS v_s^2 v_u^2 \text{conj}[\lambda] + \\ & v_d^2 v_s^3 \kappa \text{conj}[\lambda] + 4 v_d^2 v_s v_u^2 \kappa \text{conj}[\lambda] + v_s^3 v_u^2 \kappa \text{conj}[\lambda] + \sqrt{2} v_d^3 v_u \mu \text{conj}[\lambda] - \end{aligned} \right.$$

Masses and Tadpoles

# Vertices

# Vertex command

A vertex for a given set of particles is calculated via

---

Mathematica

---

```
Vertex[{Field 1, Field 2, ...}, Options]
```

---

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A vertex for a given set of particles is calculated via

---

*Mathematica*

---

```
Vertex[{Field 1, Field 2, ...}, Options]
```

---

Possible options are

- `Eigenstates` -> `STATES`: for which set of states the vertices shall be calculated
- `UseDependencies` -> `True/False`: shall dependencies defined in `parameters.m` be applied

A vertex for a given set of particles is calculated via

---

Mathematica

---

Vertex[{Field 1, Field 2, ...}, Options]

---

## Examples:

The down-quark Higgs vertex is given by

```
In[20]:= Vertex[{bar[Fd], Fd, hh}]
```

```
Out[20]= {{bar[Fd[{gt1, ct1}], Fd[{gt2, ct2}], hh[{gt3}], {-1/sqrt(2) i conj[ZH[gt3, 1]] Delta[ct1, ct2]
sum[j2, 1, 3, conj[ZDL[gt2, j2]] sum[j1, 1, 3, conj[ZDR[gt1, j1]] Yd[j1, j2]], PL],
{-i conj[ZH[gt3, 1]] Delta[ct1, ct2] sum[j2, 1, 3, sum[j1, 1, 3, conj[Yd[j1, j2]] ZDR[gt2, j1]] ZDL[gt1, j2]]
PR}}
```

- The **indices** to external states **are added**
- For each vertex the **Lorentz dependent and independent parts are separated** (PL, PR are projection operators)
- $\text{Delta}[i, j]$  is the Kronecker Delta

# Vertex command

A vertex for a given set of particles is calculated via

---

Mathematica

---

Vertex[{Field 1, Field 2, ...}, Options]

---

## Examples:

The down-quark Higgs vertex **for third generation quarks and without flavour violation** is given by

```
In[21]:= Vertex[{bar[Fd], Fd, hh]} /. {gt1 -> 3, gt2 -> 3} /. ZDL -> Delta /. ZDR -> Delta
Out[21]:= {{bar[Fd[{3, ct1}], Fd[{3, ct2}], hh[{gt3}], {-frac[i conj[ZH[gt3, 1]] Delta[ct1, ct2] Yd[3, 3]]{sqrt(2)}, PL]},
{-frac[i conj[Yd[3, 3]] conj[ZH[gt3, 1]] Delta[ct1, ct2]]{sqrt(2)}, PR}}
```



# Vertex command

A vertex for a given set of particles is calculated via

---

Mathematica

---

Vertex[{Field 1, Field 2, ...}, Options]

---

## Examples:

The squark-quark gluon vertex is calculated via

```
In[28]:= Vertex[{bar[Fd], Glu, Sd}]
Out[28]:= {{bar[Fd[{gt1, ct1}], Glu[{ct2}], Sd[{gt3, ct3}],
  {
    
$$\frac{i g_3 \text{PhaseGlu Lam}[ct2, ct1, ct3] \text{sum}[j1, 1, 3, \text{conj}[ZD[gt3, 3 + j1]] \text{conj}[ZDR[gt1, j1]]]}{\sqrt{2}}, \text{PL}$$

    ,
    
$$- \frac{i g_3 \text{conj}[PhaseGlu] Lam}[ct2, ct1, ct3] \text{sum}[j1, 1, 3, \text{conj}[ZD[gt3, j1]] ZDL[gt1, j1]]}{\sqrt{2}}, \text{PR}}}}$$

```

- Lam are the Gell-Mann matrices (fSU3 would be the structure constants)

# Vertex command

A vertex for a given set of particles is calculated via

---

Mathematica

---

```
Vertex[{Field 1, Field 2, ...}, Options]
```

---

## Examples:

The quark-gluon vertex is calculated via

```
In[22]:= Vertex[{bar[Fd], Fd, VG}]
Out[22]:= {{bar[Fd[{gt1, ct1}], Fd[{gt2, ct2}], VG[{ct3, lt3}],
  {-1/2 i g3 Delta[gt1, gt2] Lam[ct3, ct1, ct2], LorentzProduct[gamma[lt3], PL]},
  {-1/2 i g3 Delta[gt1, gt2] Lam[ct3, ct1, ct2], LorentzProduct[gamma[lt3], PR]}}
```

- `gamma[x]` is  $\gamma_x$
- `LorentzProduct` defines a non-commutative product

- SARAH usually expresses the vertices in fundamental quantities.
- Relations can be defined in `parameters.m` and used via `UseDependences->True`

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- Relations can be defined in `parameters.m` and used via `UseDependencies->True`

## Example:

The standard form of the quark-photon vertex is

```
In[23]= Vertex[{bar[Fd], Fd, VP}]
Out[23]= {bar[Fd[{gt1, ct1}], Fd[{gt2, ct2}], VP[{lt3}],
  {-1/6 i Delta[ct1, ct2] Delta[gt1, gt2] (g1 Cos[ThetaW] - 3 g2 Sin[ThetaW]), LorentzProduct[gamma[lt3], PL]},
  {1/3 i g1 Cos[ThetaW] Delta[ct1, ct2] Delta[gt1, gt2], LorentzProduct[gamma[lt3], PR]}}
```

- SARAH usually expresses the vertices in fundamental quantities.
- Relations can be defined in `parameters.m` and used via `UseDependencies->True`

## Example:

One can use  $e, \Theta_W$  instead of  $g_1, g_2$  via

```
In[24]:= Vertex[{bar[Fd], Fd, VP}, UseDependencies -> True]
Out[24]= {{bar[Fd[{gt1, ct1}], Fd[{gt2, ct2}], VP[{lt3}],
  {1/3 i e Delta[ct1, ct2] Delta[gt1, gt2], LorentzProduct[gamma[lt3], PL]},
  {1/3 i e Delta[ct1, ct2] Delta[gt1, gt2], LorentzProduct[gamma[lt3], PR]}}
```

- All dependencies are stored in the list `subDependencies`
- This list can be applied also after calculating vertices

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- This list can be applied also after calculating vertices

## Example:

For the SMSSM the defined dependencies are

```
In[25]= subDependencies
Out[25]= {g1 → e Sec[ThetaW], g2 → e Csc[ThetaW],
  ZP[index1_Integer, index2_Integer] => {{-Cos[beta], Sin[beta]}, {Sin[beta], Cos[beta]}}[[index1, index2]}
```

- All dependencies are stored in the list `subDependencies`
- This list can be applied also after calculating vertices

## Example:

The standard form of the four charged Higgs vertex is

```
In[41]:= Vertex[{Hpm, conj[Hpm], Hpm, conj[Hpm]}]
Out[41]:= {Hpm[gt1], conj[Hpm[gt2]], Hpm[gt3], conj[Hpm[gt4]]},
{1/4 i (ZP[gt1, 2] ((g1^2 + g2^2 - 4 λ conj[λ]) ZP[gt2, 1] ZP[gt3, 1] ZP[gt4, 2] +
ZP[gt2, 2] ((g1^2 + g2^2 - 4 λ conj[λ]) ZP[gt3, 1] ZP[gt4, 1] - 2 (g1^2 + g2^2) ZP[gt3, 2] ZP[gt4, 2])) +
ZP[gt1, 1] ((g1^2 + g2^2 - 4 λ conj[λ]) ZP[gt2, 2] ZP[gt3, 2] ZP[gt4, 1] +
ZP[gt2, 1] (-2 (g1^2 + g2^2) ZP[gt3, 1] ZP[gt4, 1] + (g1^2 + g2^2 - 4 λ conj[λ]) ZP[gt3, 2] ZP[gt4, 2]))), 1}}
```



- All dependencies are stored in the list `subDependencies`
- This list can be applied also after calculating vertices

## Example:

One can fix the generation indices and replace afterwards the rotation matrix by the defined angle

```
In[42]:= Vertex[{Hpm, conj[Hpm], Hpm, conj[Hpm]} /. {gt1 -> 1, gt2 -> 1, gt3 -> 1, gt4 -> 1} /. subDependencies //  
FullSimplify  
Out[42]:= {{Hpm[{1}], conj[Hpm[{1}]], Hpm[{1}], conj[Hpm[{1}]]},  
{-1/4 i Csc[ThetaW]^2 Sec[ThetaW]^2 (e^2 (1 + Cos[4 beta]) + lambda conj[lambda] Sin[2 ThetaW]^2 Sin[2 beta]^2), 1}}
```

# All vertices at once

It is possible to calculate all vertices at once using

---

```
MakeVertexList [ STATES, Options ]
```

---

with

- STATES: the eigenstates for which all vertices shall be calculated
- One can define a subset of generic classes which should be considered, e.g. GenericClasses->{FFS, FFV}

# All vertices at once

It is possible to calculate all vertices at once using

---

```
MakeVertexList [ STATES, Mathematica Options ]
```

---

with

- **STATES**: the eigenstates for which all vertices shall be calculated
- One can define a subset of generic classes which should be considered, e.g. `GenericClasses->{FFS, FFV}`

## Output:

- The results are lists `SA'VertexList[TYPE]` for each generic class: `{FFS, FFV, SSS, SSV, SVV, SSVV, VVV, VVVV, GGS, GGV}`
- The results are stored in the output directory of the model
- The content of these lists are the information which gets exported into `FeynArts`, `CalcHep` or `UFO` files

# Example:

```
In[29]= MakeVertexList[EWSB]
Generate Directories
Building Particle List
Calculate all vertices
  Three Scalar - Interactions
    Found 26 potential vertices. Calculating 26/26. (All done in 17.468s; 17 are non-vanishing)
  Four Scalar - Interactions
    Found 65 potential vertices. Calculating 65/65. (All done in 30.724s; 44 are non-vanishing)
  Two Scalar - One Vector Boson - Interactions
    Found 66 potential vertices. Calculating 66/66. (All done in 14.8s; 20 are non-vanishing)
  One Scalar - Two Vector Boson - Interactions
    Found 36 potential vertices. Calculating 36/36. (All done in 8.856s; 6 are non-vanishing)
  Two Scalar - Two Vector Boson - Interactions
    Found 178 potential vertices. Calculating 178/178. (All done in 46.32s; 46 are non-vanishing)
  Three Vector Boson - Interactions
    Found 21 potential vertices. Calculating 21/21. (All done in 3.716s; 3 are non-vanishing)
  Two Fermion - One Scalar - Interactions
    Found 84 potential vertices. Calculating 84/84. (All done in 31.076s; 36 are non-vanishing)
  Two Fermion - One Vector Boson - Interactions
    Found 59 potential vertices. Calculating 59/59. (All done in 24.212s; 19 are non-vanishing)
  Four Vector Boson - Interactions
    Found 36 potential vertices. Calculating 36/36. (All done in 12.14s; 5 are non-vanishing)
  Two Ghost - One Vector Boson - Interactions
    Found 65 potential vertices. Calculating 65/65. (All done in 4.78s; 13 are non-vanishing)
  Two Ghost - One Scalar - Interactions
    Found 48 potential vertices. Calculating 48/48. (All done in 3.588s; 12 are non-vanishing)
  Two Scalar - One Auxiliary - Interactions
    Found 102 potential vertices. Calculating 102/102. (All done in 14.492s; 100 are non-vanishing)

Simplify Vertices
Writing vertices to files
All vertices calculated. (Time needed: 212.776s)
The vertices are saved in /home/frnstaub/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/EWSB/Vertices/
```

## Example:

the FFS vertices are stored in `SA`VertexList[FFS]`:

```
In[43]:= SA`VertexList[FFS]
```

```
Out[43]= {{{bar[Cha[{gt1}], Cha[{gt2}], Ah[{gt3}]},  
  {  
    - $\frac{1}{\sqrt{2}}$  (g2 conj[UM[gt2, 1]] conj[UP[gt1, 2]] conj[ZA[gt3, 2]] + conj[UM[gt2, 2]]  
      (g2 conj[UP[gt1, 1]] conj[ZA[gt3, 1]] -  $\lambda$  conj[UP[gt1, 2]] conj[ZA[gt3, 3]]), PL},  
    {  
       $\frac{1}{\sqrt{2}}$  (g2 conj[ZA[gt3, 1]] UM[gt1, 2] UP[gt2, 1] + (g2 conj[ZA[gt3, 2]] UM[gt1, 1] -  
        conj[ $\lambda$ ] conj[ZA[gt3, 3]] UM[gt1, 2]) UP[gt2, 2]), PR}}, {Chi[{gt1}], Chi[{gt2}], Ah[{gt3}]},  
    {  
       $\frac{1}{2}$  (-conj[ZA[gt3, 2]] (conj[ZN[gt1, 4]] (g1 conj[ZN[gt2, 1]] - g2 conj[ZN[gt2, 2]]) +  
         $\sqrt{2}$   $\lambda$  conj[ZN[gt1, 5]] conj[ZN[gt2, 3]] + g1 conj[ZN[gt1, 1]] conj[ZN[gt2, 4]] -  
        g2 conj[ZN[gt1, 2]] conj[ZN[gt2, 4]] +  $\sqrt{2}$   $\lambda$  conj[ZN[gt1, 3]] conj[ZN[gt2, 5]]) -  
      conj[ZA[gt3, 1]] (conj[ZN[gt1, 3]] (-g1 conj[ZN[gt2, 1]] + g2 conj[ZN[gt2, 2]]) -  
        g1 conj[ZN[gt1, 1]] conj[ZN[gt2, 3]] + g2 conj[ZN[gt1, 2]] conj[ZN[gt2, 3]] +  
         $\sqrt{2}$   $\lambda$  conj[ZN[gt1, 5]] conj[ZN[gt2, 4]] +  $\sqrt{2}$   $\lambda$  conj[ZN[gt1, 4]] conj[ZN[gt2, 5]]) -  
      ...  
    }  
  }  
}
```

## Example:

Mathematica commands can be used to filter & select subgroups of vertices, e.g. all FFS vertices with a pseudo-scalar are returned by

```
In[44]:= Select[SA`VertexList[FFS], FreeQ[#, Ah] == False &]
```

```
Out[44]= {{bar[Cha[{{gt1}}], Cha[{{gt2}}], Ah[{{gt3}}]},
```

$$\left\{ -\frac{1}{\sqrt{2}} (g_2 \text{conj}[\text{UM}[\text{gt}_2, 1]] \text{conj}[\text{UP}[\text{gt}_1, 2]] \text{conj}[\text{ZA}[\text{gt}_3, 2]] + \text{conj}[\text{UM}[\text{gt}_2, 2]] \right. \\ \left. (g_2 \text{conj}[\text{UP}[\text{gt}_1, 1]] \text{conj}[\text{ZA}[\text{gt}_3, 1]] - \lambda \text{conj}[\text{UP}[\text{gt}_1, 2]] \text{conj}[\text{ZA}[\text{gt}_3, 3]]), \text{PL} \right\},$$

$$\left\{ \frac{1}{\sqrt{2}} (g_2 \text{conj}[\text{ZA}[\text{gt}_3, 1]] \text{UM}[\text{gt}_1, 2] \text{UP}[\text{gt}_2, 1] + (g_2 \text{conj}[\text{ZA}[\text{gt}_3, 2]] \text{UM}[\text{gt}_1, 1] - \right. \\ \left. \text{conj}[\lambda] \text{conj}[\text{ZA}[\text{gt}_3, 3]] \text{UM}[\text{gt}_1, 2]) \text{UP}[\text{gt}_2, 2]), \text{PR} \right\}, \left\{ \text{Chi}[\{\{\text{gt}_1\}\}], \text{Chi}[\{\{\text{gt}_2\}\}], \text{Ah}[\{\{\text{gt}_3\}\}], \right.$$

$$\left. \frac{1}{2} \left( -\text{conj}[\text{ZA}[\text{gt}_3, 2]] \left( \text{conj}[\text{ZN}[\text{gt}_1, 4]] (g_1 \text{conj}[\text{ZN}[\text{gt}_2, 1]] - g_2 \text{conj}[\text{ZN}[\text{gt}_2, 2]]) + \right. \right. \\ \left. \left. \sqrt{2} \lambda \text{conj}[\text{ZN}[\text{gt}_1, 5]] \text{conj}[\text{ZN}[\text{gt}_2, 3]] + g_1 \text{conj}[\text{ZN}[\text{gt}_1, 1]] \text{conj}[\text{ZN}[\text{gt}_2, 4]] - \right. \right. \\ \left. \left. g_2 \text{conj}[\text{ZN}[\text{gt}_1, 2]] \text{conj}[\text{ZN}[\text{gt}_2, 4]] + \sqrt{2} \lambda \text{conj}[\text{ZN}[\text{gt}_1, 3]] \text{conj}[\text{ZN}[\text{gt}_2, 5]] \right) - \right. \\ \left. \text{conj}[\text{ZA}[\text{gt}_3, 1]] \left( \text{conj}[\text{ZN}[\text{gt}_1, 3]] (-g_1 \text{conj}[\text{ZN}[\text{gt}_2, 1]] + g_2 \text{conj}[\text{ZN}[\text{gt}_2, 2]]) - \right. \right. \\ \left. \left. g_1 \text{conj}[\text{ZN}[\text{gt}_1, 1]] \text{conj}[\text{ZN}[\text{gt}_2, 3]] + g_2 \text{conj}[\text{ZN}[\text{gt}_1, 2]] \text{conj}[\text{ZN}[\text{gt}_2, 3]] + \right. \right. \\ \left. \left. \sqrt{2} \lambda \text{conj}[\text{ZN}[\text{gt}_1, 5]] \text{conj}[\text{ZN}[\text{gt}_2, 4]] + \sqrt{2} \lambda \text{conj}[\text{ZN}[\text{gt}_1, 4]] \text{conj}[\text{ZN}[\text{gt}_2, 5]] \right) - \right. \\ \left. \sqrt{2} \text{conj}[\text{ZA}[\text{gt}_3, 3]] (\lambda \text{conj}[\text{ZN}[\text{gt}_1, 4]] \text{conj}[\text{ZN}[\text{gt}_2, 3]] + \lambda \text{conj}[\text{ZN}[\text{gt}_1, 3]] \text{conj}[\text{ZN}[\text{gt}_2, 4]] - \right. \\ \left. \left. 2 \kappa \text{conj}[\text{ZN}[\text{gt}_1, 5]] \text{conj}[\text{ZN}[\text{gt}_2, 5]]) \right), \text{PL} \right\},$$

$$\left\{ \frac{1}{2} \left( -\text{conj}[\text{ZA}[\text{gt}_3, 1]] \left( \text{ZN}[\text{gt}_1, 3] (g_1 \text{ZN}[\text{gt}_2, 1] - g_2 \text{ZN}[\text{gt}_2, 2]) + g_1 \text{ZN}[\text{gt}_1, 1] \text{ZN}[\text{gt}_2, 3] - g_2 \text{ZN}[\text{gt}_1, 2] \right. \right. \right.$$

...

# Renormalisation Group Equations

The one- and two-loop RGEs for a model are calculated via

---

```
CalcRGES[ Options]
```

---

The possible options are

- TwoLoop -> True/False
- ReadLists -> True/False (reading previous results)
- VariableGenerations -> FIELDS (consider number of generations as free variable)
- NoMatrixMultiplication -> True/False (use explicit sums instead of matrix multiplication)
- IgnoreAt2Loop -> PARAMETERS (ignore some parameters at two-loop)
- WriteFunctionsToRun->True/False (write a file to evaluate the RGEs numerically in Mathematica)



In[45]:= **CalcRGEs []**

Calculate supersymmetric RGEs

Making Lists of Particles and Couplings

Calculating anomalous Dimensions

Calculate Beta Functions for trilinear Superpotential parameters

Calculating 5/5. (All done in 1.004s)

Calculate Beta Functions for bilinear Superpotential parameters

Calculating 2/2. (All done in 0.208s)

...

Calculate Beta Functions for VEVs

Calculating 3/3. (All done in 0.396s)

Writing Mathematica code to evaluate RGEs

Finished with the calculation of the RGEs. Time needed: 43.664s

The results are saved in `/home/fnstaub/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/RGEs/`

# Results

The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

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---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

The names of the arrays **for SUSY models** are

- **Gij**: Anomalous dimensions of all chiral superfields
- **BetaWijkl**: Quartic superpotential parameters
- **BetaYijk**: Trilinear superpotential parameters
- **BetaMuij**: Bilinear superpotential parameters
- **BetaLi**: Linear superpotential parameters
- **BetaQijkl**: Quartic soft-breaking parameters
- **BetaTijk**: Trilinear soft-breaking parameters
- **BetaBij**: Bilinear soft-breaking parameters
- **BetaSLi**: Linear soft-breaking parameters
- **Beta $m_{2ij}$** : Scalar squared masses
- **BetaMi**: Majorana Gaugino masses
- **BetaGauge**: Gauge couplings
- **BetaVEVs**: VEVs
- **BetaDGi**: Dirac gaugino mass terms

The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

The names of the arrays **for non-SUSY models** are

- **G<sub>ij</sub>**: Anomalous dimensions of all fermions and scalars
- **BetaGauge**: Gauge couplings
- **BetaLijkl**: Quartic scalar couplings
- **BetaYijk**: Interactions between two fermions and one scalar
- **BetaTijk**: Cubic scalar interactions
- **BetaMuij**: Bilinear fermion term
- **BetaBij**: Bilinear scalar term
- **BetaVEVs**: Vacuum expectation values

The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

## Example:

The  $\beta$ -functions of the gauge couplings are

In[46]:= **BetaGauge**

```
Out[46]:= {{g1,  $\frac{33}{5} g1^3$ ,  $\frac{1}{25} g1^3 (199 g1^2 + 135 g2^2 + 440 g3^2 -$   
30  $\lambda \text{ conj}[\lambda] - 70 \text{ trace}[\text{Yd}, \text{Adj}[\text{Yd}]] - 90 \text{ trace}[\text{Ye}, \text{Adj}[\text{Ye}]] - 130 \text{ trace}[\text{Yu}, \text{Adj}[\text{Yu}]]$ }},  
{g2,  $g2^3$ ,  $\frac{1}{5} g2^3 (9 g1^2 + 125 g2^2 + 120 g3^2 - 10 \lambda \text{ conj}[\lambda] - 30 \text{ trace}[\text{Yd}, \text{Adj}[\text{Yd}]] -$   
10  $\text{ trace}[\text{Ye}, \text{Adj}[\text{Ye}]] - 30 \text{ trace}[\text{Yu}, \text{Adj}[\text{Yu}]]$ }},  
{g3,  $-3 g3^3$ ,  $\frac{1}{5} g3^3 (11 g1^2 + 45 g2^2 + 70 g3^2 - 20 \text{ trace}[\text{Yd}, \text{Adj}[\text{Yd}]] - 20 \text{ trace}[\text{Yu}, \text{Adj}[\text{Yu}]]$ )}}
```

- Coefficients  $1/16\pi^2$ ,  $1/(16\pi^2)^2$  are dropped

The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt.}

---

## Example:

The  $\beta$ -functions of the gauge couplings are

In[46]:= **BetaGauge**

```
Out[46]:= {{g1,  $\frac{33 g_1^3}{5}$ ,  $\frac{1}{25} g_1^3 (199 g_1^2 + 135 g_2^2 + 440 g_3^2 - 30 \lambda \text{conj}[\lambda] - 70 \text{trace}[\text{Yd}, \text{Adj}[\text{Yd}]] - 90 \text{trace}[\text{Ye}, \text{Adj}[\text{Ye}]] - 130 \text{trace}[\text{Yu}, \text{Adj}[\text{Yu}]]]$ },  
{g2,  $\frac{g_2^3}{5}$ ,  $\frac{1}{5} g_2^3 (9 g_1^2 + 125 g_2^2 + 120 g_3^2 - 10 \lambda \text{conj}[\lambda] - 30 \text{trace}[\text{Yd}, \text{Adj}[\text{Yd}]] - 10 \text{trace}[\text{Ye}, \text{Adj}[\text{Ye}]] - 30 \text{trace}[\text{Yu}, \text{Adj}[\text{Yu}]]]$ },  
{g3,  $-3 g_3^3$ ,  $\frac{1}{5} g_3^3 (11 g_1^2 + 45 g_2^2 + 70 g_3^2 - 20 \text{trace}[\text{Yd}, \text{Adj}[\text{Yd}]] - 20 \text{trace}[\text{Yu}, \text{Adj}[\text{Yu}]]]$ }}
```

- Coefficients  $1/16\pi^2$ ,  $1/(16\pi^2)^2$  are dropped
- 1-loop, 2-loop

The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

## Example:

The  $\beta$ -functions of  $T_k$  are

```
In[49]:= BetaTijk[[4]]
Out[49]:= {T[k], 6 (3 κ conj[k] T[k] + conj[λ] (λ T[k] + 2 κ T[λ])),
  - 6/5 (100 κ^2 conj[k]^2 T[k] + 10 λ conj[λ]^2 (λ T[k] + 4 κ T[λ]) +
  conj[λ] (λ T[k] (-3 g1^2 - 15 g2^2 + 60 κ conj[k] + 15 trace[Yd, Adj[Yd]] + 5 trace[Ye, Adj[Ye]] +
  15 trace[Yu, Adj[Yu]]) + 2 κ (T[λ] (-3 g1^2 - 15 g2^2 + 20 κ conj[k] + 15 trace[Yd, Adj[Yd]] +
  5 trace[Ye, Adj[Ye]] + 15 trace[Yu, Adj[Yu]]) + λ (3 g1^2 MassB + 15 g2^2 MassWB +
  15 trace[Adj[Yd], T[Yd]] + 5 trace[Adj[Ye], T[Ye]] + 15 trace[Adj[Yu], T[Yu]]))}}
```

# Results

The results are stored in three dimensional arrays containing

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

## Example:

For soft masses often traces  $\text{Tr}X[Y]$  appear, e.g. for  $m_{H_u}^2$

```
In[48]:= Betam2ij[4]
```

$$\begin{aligned} \text{Out[48]} = & \left\{ m_{H_u}^2, -\frac{6}{5} g_1^2 \text{MassB} \text{conj}[\text{MassB}] - 6 g_2^2 \text{MassWB} \text{conj}[\text{MassWB}] + 2 m_{H_d}^2 \lambda \text{conj}[\lambda] + \right. \\ & 2 m_{H_u}^2 \lambda \text{conj}[\lambda] + 2 m_{s_2}^2 \lambda \text{conj}[\lambda] + 2 \text{conj}[\text{T}[\lambda]] \text{T}[\lambda] + \sqrt{\frac{3}{5}} g_1 \text{Tr1}[1] + 6 m_{H_u}^2 \text{trace}[\text{Yu}, \text{Adj}[\text{Yu}]] + \\ & 6 \text{trace}[\text{conj}[\text{T}[\text{Yu}]], \text{Tp}[\text{T}[\text{Yu}]]] + 6 \text{trace}[\text{mq2}, \text{Adj}[\text{Yu}], \text{Yu}] + 6 \text{trace}[\text{mu2}, \text{Yu}, \text{Adj}[\text{Yu}]], \\ & \frac{1}{25} \left( g_1^2 \text{conj}[\text{MassB}] \left( 621 g_1^2 \text{MassB} + 90 g_2^2 \text{MassB} + 45 g_2^2 \text{MassWB} + 80 \text{MassB} \text{trace}[\text{Yu}, \text{Adj}[\text{Yu}]] - \right. \right. \\ & 40 \text{trace}[\text{Adj}[\text{Yu}], \text{T}[\text{Yu}]] \left. \right) + 5 \left( 3 g_2^2 \left( 55 g_2^2 \text{MassWB} + 3 g_1^2 (\text{MassB} + 2 \text{MassWB}) \right) \text{conj}[\text{MassWB}] - \right. \\ & 2 \left( 30 (m_{H_d}^2 + m_{H_u}^2 + m_{s_2}^2) \lambda^2 \text{conj}[\lambda]^2 + 10 \text{conj}[\kappa] \right. \\ & \left. \left. \left( (m_{H_d}^2 + m_{H_u}^2 + 4 m_{s_2}^2) \kappa \lambda \text{conj}[\lambda] + \text{conj}[\text{T}[\lambda]] (\lambda \text{T}[\kappa] + \kappa \text{T}[\lambda]) \right) - 15 g_2^4 \text{Tr2}[2] - 3 g_1^2 \text{Tr2U1}[1, 1] - \right. \right. \\ & \left. \left. 2 \sqrt{15} g_1 \text{Tr3}[1] + 15 \text{conj}[\text{T}[\lambda]] \text{T}[\lambda] \text{trace}[\text{Yd}, \text{Adj}[\text{Yd}]] + 5 \text{conj}[\text{T}[\lambda]] \text{T}[\lambda] \text{trace}[\text{Ye}, \text{Adj}[\text{Ye}]] - \right. \right. \\ & \left. \left. \dots \right. \right. \end{aligned}$$



The results are stored in three dimensional arrays containing

---

{Parameter, 1-loop Beta-Fkt., 2-loop Beta-Fkt}

---

## Example:

The expressions for the traces are given in TraceAbbr

```
In[78]= TraceAbbr
Out[78]= {{{Tr1[1],  $\sqrt{\frac{3}{5}}$  g1 (-mHd2 + mHu2 + trace[md2] + trace[me2] - trace[ml2] + trace[mq2] - 2 trace[mu2])}},
  {{Tr2U1[1, 1],  $\frac{1}{10}$  g1^2 (3 mHd2 + 3 mHu2 + 2 trace[md2] + 6 trace[me2] + 3 trace[ml2] + trace[mq2] + 8 trace[mu2])}},
  {{Tr3[1],  $\frac{1}{20\sqrt{15}}$  g1 (-9 g1^2 mHd2 - 45 g2^2 mHd2 + 9 g1^2 mHu2 + 45 g2^2 mHu2 + 30 (mHd2 - mHu2)  $\lambda$  conj[ $\lambda$ ] +
    4 (g1^2 + 20 g3^2) trace[md2] + 36 g1^2 trace[me2] - 9 g1^2 trace[ml2] - 45 g2^2 trace[ml2] +
    g1^2 trace[mq2] + 45 g2^2 trace[mq2] + 80 g3^2 trace[mq2] - 32 g1^2 trace[mu2] - 160 g3^2 trace[mu2] +
    90 mHd2 trace[Yd, Adj[Yd]] + 30 mHd2 trace[Ye, Adj[Ye]] - 90 mHu2 trace[Yu, Adj[Yu]] -
    60 trace[Yd, Adj[Yd], conj[md2]] - 30 trace[Yd, conj[mq2], Adj[Yd]] - 60 trace[Ye, Adj[Ye], conj[me2]] +
    30 trace[Ye, conj[ml2], Adj[Ye]] + 120 trace[Yu, Adj[Yu], conj[mu2]] - 30 trace[Yu, conj[mq2], Adj[Yu])}},
  {{Tr2[2],  $\frac{1}{2}$  (mHd2 + mHu2 + trace[ml2] + 3 trace[mq2])}, {Tr2[3],  $\frac{1}{2}$  (trace[md2] + 2 trace[mq2] + trace[mu2])}}}}
```

- SARAH writes the RGEs also in a **format which can be used directly with Mathematica**
- This format is saved in the file `RunRGEs.m`
- Also a **function `RunRGEs to run the RGEs`** is provided in this file.  
The syntax is

---

```
RunRGEs[ input, scale1, Mathematica scale2, Options]
```

---

- SARAH writes the RGEs also in a **format which can be used directly with Mathematica**
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The syntax is

---

```
RunRGEs[ input, scale1, scale2, Options]
```

---

- **non-vanishing boundary conditions at the scale where the running starts**

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- This format is saved in the file `RunRGEs.m`
- Also a **function `RunRGEs to run the RGEs`** is provided in this file.  
The syntax is

---

```
RunRGEs[ input, scale1, Mathematica scale2, Options ]
```

---

- **non-vanishing boundary conditions at the scale where the running starts**
- **log of the scale where the running starts**

- SARAH writes the RGEs also in a **format which can be used directly with Mathematica**
- This format is saved in the file `RunRGEs.m`
- Also a **function `RunRGEs to run the RGEs`** is provided in this file.  
The syntax is

---

```
RunRGEs [ input, scale1, Mathematica scale2, Options ]
```

---

- non-vanishing boundary conditions at the scale where the running starts
- log of the scale where the running starts
- log of the scale where the running stops

- SARAH writes the RGEs also in a **format which can be used directly with Mathematica**
- This format is saved in the file `RunRGEs.m`
- Also a **function `RunRGEs to run the RGEs`** is provided in this file.  
The syntax is

---

```
RunRGEs [ input, scale1, scale2, Options ]
```

---

- non-vanishing boundary conditions at the scale where the running starts
- log of the scale where the running starts
- log of the scale where the running stops
- option to turn off two-loop running (`TwoLoop->False`)

# Examples:

- Loading RunRGEs .m
- Running the gauge couplings as well as  $Y_t$  &  $\lambda$  from  $10^3$  to  $10^{17}$  GeV

```
In[52]:= << "SARAH-4.12.2/Output/SMSSM/RGEs/RunRGEs.m"
```

```
In[54]:= runRGEs = RunRGEs[{g1 -> 0.47, g2 -> 0.65, g3 -> 1.1, Yu[3, 3] -> 0.9, T[lam] -> 0, lam -> 0.6}, 3, 17][[1]]
```

```
Out[54]:= {L[L1] -> InterpolatingFunction[+ [Domain: {{3., 17.}}  
Output: scalar],
```

```
B[mu] -> InterpolatingFunction[+ [Domain: {{3., 17.}}  
Output: scalar],
```

→ **the results are stored as interpolating function**

These functions are used as

---

Mathematica

---

PARAMTER[SCALE] /. InterpolatingFunction

---

to get the value of a parameter at any scale



# Examples:

One can use the function to get values of the **running gauge couplings at the GUT scale**:

```
In[55]:= {g1[16], g2[16], g3[16], Yu[3, 3][16], λ[16]} /. runRGEs  
Out[55]:= {0.709092, 0.720759, 0.719522, 0.644797, 0.956999}
```

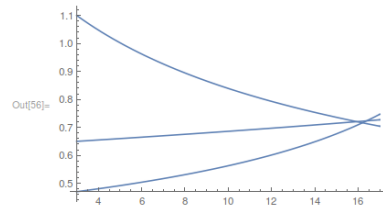
# Examples:

One can use the function to get values of the **running gauge couplings at the GUT scale**:

```
In[55]:= {g1[16], g2[16], g3[16], Yu[3, 3][16],  $\lambda$ [16]} /. runRGEs  
Out[55]:= {0.709092, 0.720759, 0.719522, 0.644797, 0.956999}
```

One can also make plots to show the running:

```
In[56]:= Plot[{g1[t], g2[t], g3[t]} /. runRGEs, {t, 3, 17}]
```



SUSY boundary conditions at the GUT scale can be set:

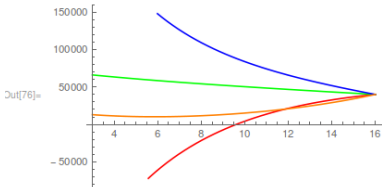
```
In[66]:= runRGESUSY =  
  RunRGES[({g1 → (g1[16] /. runRGES), g2 → (g2[16] /. runRGES), g3 → (g3[16] /. runRGES), T[λ] → 0,  
    λ → (λ[16] /. runRGES), Yu[3, 3] → (Yu[3, 3][16] /. runRGES), mq2[a__] → Delta[a] m0 ^ 2,  
    mu2[a__] → Delta[a] m0 ^ 2, md2[a__] → Delta[a] m0 ^ 2, me2[a__] → Delta[a] m0 ^ 2,  
    ml2[a__] → Delta[a] m0 ^ 2, ms2 → m0 ^ 2, mHd2 → m0 ^ 2, mHu2 → m0 ^ 2, MassB → M12, MassWB → M12,  
    MassG → M12} /. {m0 → 200, M12 → 300}, 16, 3][[1]];
```

SUSY boundary conditions at the GUT scale can be set:

```
In[66]:= runRGEsSUSY =  
  RunRGEs[{g1 → (g1[16] /. runRGEs), g2 → (g2[16] /. runRGEs), g3 → (g3[16] /. runRGEs), T[λ] → 0,  
    λ → (λ[16] /. runRGEs), Yu[3, 3] → (Yu[3, 3][16] /. runRGEs), mq2[a__] → Delta[a] m0^2,  
    mu2[a__] → Delta[a] m0^2, md2[a__] → Delta[a] m0^2, me2[a__] → Delta[a] m0^2,  
    ml2[a__] → Delta[a] m0^2, ms2 → m0^2, mHd2 → m0^2, mHu2 → m0^2, MassB → M12, MassWB → M12,  
    MassG → M12} /. {m0 → 200, M12 → 300}, 16, 3][[1]];
```

... and the running of the masses can be plotted:

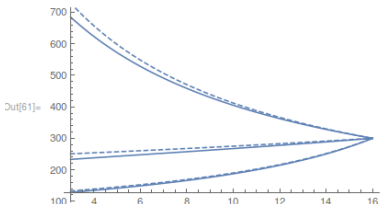
```
In[76]:= Plot[{mHu2[t] /. runRGEsSUSY, mHd2[t] /. runRGEsSUSY, mu2[3, 3][t] /. runRGEsSUSY, ms2[t] /. runRGEsSUSY},  
  {t, 16, 3}, PlotStyle → {Red, Green, Blue, Orange}]
```



Also one- and two-loop running can be compared, e.g. for the gauginos:

```
In[60]:= runRGEsSUSY1L =  
  RunRGEs[{g1 → (g1[16] /. runRGEs), g2 → (g2[16] /. runRGEs), g3 → (g3[16] /. runRGEs),  
    Yu[3, 3] → (Yu[3, 3][16] /. runRGEs), mq2[a__] → Delta[a] m0 ^ 2, mu2[a__] → Delta[a] m0 ^ 2,  
    md2[a__] → Delta[a] m0 ^ 2, me2[a__] → Delta[a] m0 ^ 2, ml2[a__] → Delta[a] m0 ^ 2, mHd2 → m0 ^ 2,  
    mHu2 → m0 ^ 2, MassB → M12, MassWB → M12, MassG → M12} /. {m0 → 200, M12 → 300}, 16, 3, TwoLoop → False][[  
    1]];
```

```
In[61]:= Show[Plot[{MassB[t], MassWB[t], MassG[t]} /. runRGEsSUSY, {t, 16, 3}],  
  Plot[{MassB[t], MassWB[t], MassG[t]} /. runRGEsSUSY1L, {t, 16, 3}, PlotStyle → Dashed]]
```



# L<sup>A</sup>T<sub>E</sub>X Output

All information which we have obtained so far can be exported into L<sup>A</sup>T<sub>E</sub>X using

---

MakeTeX[ [Options](#) ] *Mathematica*

---

The possible options are

- FeynmanDiagrams -> True/False (feynman diagrams for all vertices?)
- ShortForm -> True/False (write vertices in a more compact form)
- WriteSARAH -> True/False (write information about the model implementation in SARAH)

All information which we have obtained so far can be exported into L<sup>A</sup>T<sub>E</sub>X using

---

```
MakeTeX[ Options]
```

---

The possible options are

- `FeynmanDiagrams` -> True/False (feynman diagrams for all vertices?)
- `ShortForm` -> True/False (write vertices in a more compact form)
- `WriteSARAH` -> True/False (write information about the model implementation in SARAH)

About Feynman diagrams:

- To draw Feynman diagrams, the package `FeynMF` must be installed
- A batch script is provided to compile all diagrams automatically



## Generate the $\text{\LaTeX}$ files

In[77]: **MakeTeX** []

### Generate LaTeX files

Writing Superfields and Superpotential to TeX-File

Writing Particle Content to TeX-File

Write VEVs to TeX-File

Write Flavor Decomposition to TeX-File

Writing Mass Matrices to TeX-File

Writing Tadpole Equations to TeX-File

Writing RGEs to TeX-File

**TeXOutput :**

Loop corrections not calculated so far. Skipping this parts.

Use CalcLoopCorrections[States] to calculate RGEs and start MakeTeX again to include them in the output.

Write Clebsch-Gordan Coefficients

Writing Vertices to TeX-File

Done. Output is in /home/fnstaub/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/EWSB/TeX/

Use Script MakePDF.sh (Linux) or MakePDF.bat (Windows) to generate pdf file.

# Example

make the script executable and run it

```
fnstaud@fnstaud-ThinkPad-T450s ~/Desktop/HEP-Tools $ cd SARAH-4.12.2/Output/SMSSM/EWSB/TeX/
fnstaud@fnstaud-ThinkPad-T450s ~/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/EWSB/TeX $ chmod 775 MakePDF.sh
fnstaud@fnstaud-ThinkPad-T450s ~/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/EWSB/TeX $ ./MakePDF.sh

. . .

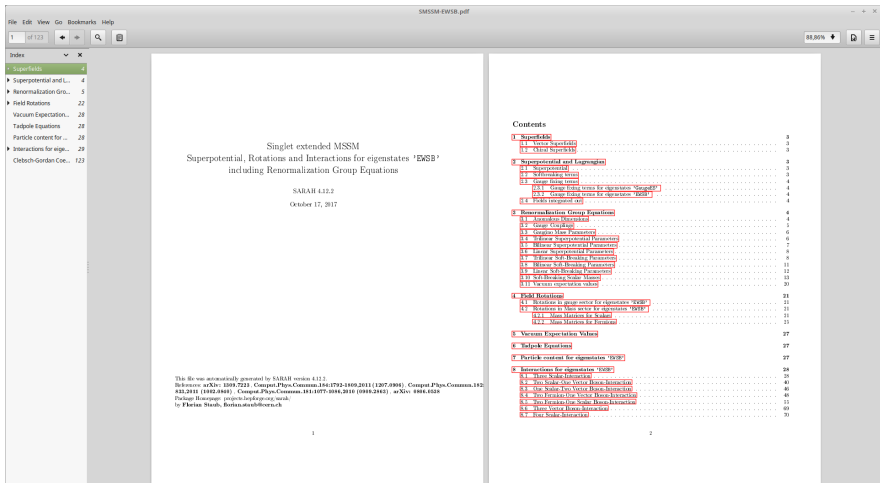
Output written on SMSSM-EWSB.pdf (123 pages, 525889 bytes).
Transcript written on SMSSM-EWSB.log.

PDF for Model finished
Thanks for using SARAH

fnstaud@fnstaud-ThinkPad-T450s ~/Desktop/HEP-Tools/SARAH-4.12.2/Output/SMSSM/EWSB/TeX $
```

# Example

Take a look ...



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SMSSM-EWSB.pdf

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Singlet extended MSSM  
Superpotential, Rotations and Interactions for eigenstates 'EWSB'  
including Renormalization Group Equations

SARAH 4.3.2.2  
October 17, 2017

This file was automatically generated by SARAH version 4.3.2.2  
Bibliography arXiv: 1209.7223, Comput.Phys.Commun.144:1792-1809,2011(1207.0994), Comput.Phys.Commun.182:  
422,2011(1102.1885), Comput.Phys.Commun.181:1077-1084,2010(0908.2062), arXiv: 0804.4526  
Desktop Homepage: <http://www.fys.kit.edu/~staud>  
by Florian Staub, [florian.staub@kit.edu](mailto:florian.staub@kit.edu)

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# Example

## Fields and superpotential ...

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### 1 Superfields

#### 1.1 Vector Superfields

SF	Spin	Spn. 4	Spn. 1	SU(3)	SU(2)	U(1)	Source
$\hat{F}$	$\frac{1}{2}$	0	0	1	1	0	Supercharge
$\hat{W}$	$\frac{1}{2}$	0	0	3	1	0	color

#### 1.2 Chiral Superfields

SF	Spin	Spn. 4	Generators	$U(1) \oplus SU(2) \oplus SU(3)$
$\hat{\psi}$	$\frac{1}{2}$	0		(1, 2, 3)
$\hat{f}$	$\frac{1}{2}$	0		(-1/2, 2, 1)
$\hat{H}_d$	$\frac{1}{2}$	0		(-1/2, 2, 1)
$\hat{H}_u$	$\frac{1}{2}$	0		(1/2, 2, 1)
$\hat{\phi}$	0	0		(1, 1, 3)
$\hat{u}$	0	0		(-1/3, 1, 3)
$\hat{d}$	0	0		(1/3, 1, 3)
$\hat{e}$	0	0		(0, 1, 1)
$\hat{\nu}$	0	0		(0, 1, 1)

### 2 Superpotential and Lagrangian

#### 2.1 Superpotential

$$W = \Delta_1 \hat{\psi} + \mu \hat{H}_d \hat{H}_u + \frac{1}{2} M_{12} \hat{\psi}^2 - Y_d \hat{\phi} \hat{d} \hat{H}_d - Y_e \hat{\phi} \hat{e} \hat{H}_d + \lambda \hat{H}_d \hat{H}_u \hat{\psi} + \frac{1}{3} \hat{\psi}^3 + Y_\nu \hat{\phi} \hat{\nu} \hat{H}_d \quad (1)$$

#### 2.2 Softbreaking terms

$$-\mathcal{L}_{SB} = + \frac{1}{2} m_{12}^2 \hat{\psi}^2 - M_1^2 \hat{F}^2 + M_2^2 \hat{W}^2 + M_3^2 \hat{\phi}^2 + M_4^2 \hat{\psi}^2 + M_5^2 \hat{f}^2 + M_6^2 \hat{H}_d^2 + M_7^2 \hat{H}_u^2 + M_8^2 \hat{\phi}^2 + M_9^2 \hat{u}^2 + M_{10}^2 \hat{d}^2 + M_{11}^2 \hat{e}^2 + M_{12}^2 \hat{\nu}^2 + M_{13}^2 \hat{H}_d \hat{H}_u + M_{14}^2 \hat{H}_d \hat{\psi} + M_{15}^2 \hat{H}_d \hat{\phi} + M_{16}^2 \hat{H}_d \hat{u} + M_{17}^2 \hat{H}_d \hat{d} + M_{18}^2 \hat{H}_d \hat{e} + M_{19}^2 \hat{H}_d \hat{\nu} + M_{20}^2 \hat{H}_u \hat{\psi} + M_{21}^2 \hat{H}_u \hat{\phi} + M_{22}^2 \hat{H}_u \hat{u} + M_{23}^2 \hat{H}_u \hat{d} + M_{24}^2 \hat{H}_u \hat{e} + M_{25}^2 \hat{H}_u \hat{\nu} + M_{26}^2 \hat{\phi} \hat{u} + M_{27}^2 \hat{\phi} \hat{d} + M_{28}^2 \hat{\phi} \hat{e} + M_{29}^2 \hat{\phi} \hat{\nu} + M_{30}^2 \hat{u} \hat{d} + M_{31}^2 \hat{u} \hat{e} + M_{32}^2 \hat{u} \hat{\nu} + M_{33}^2 \hat{d} \hat{e} + M_{34}^2 \hat{d} \hat{\nu} + M_{35}^2 \hat{e} \hat{\nu} \quad (2)$$

L<sup>A</sup>T<sub>E</sub>X Output

Florian Staub – SARAH (part II): Working in Mathematica (Tools Bootcamp, 23.10.17)

25/26

MSMMA-EWSB.pdf

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**3.3 Gaugino Mass Parameters**

$$M_1^0 = \frac{60}{5} M_1 \quad (28)$$

$$M_2^0 = \frac{2}{3} M_1 (10g_2^2 M_1 + 120g_2^2 M_1 + 440g_2^2 M_1 + 880g_2^2 M_1 + 120g_2^2 M_2 - 31V(M_1, \lambda - T)) - 70M_1 T (V_1 V_2^2) \quad (29)$$

$$M_3^0 = -90M_1 T (V_1 V_2^2) - 134M_1 T (V_1 V_2^2) + 70T (V_1^2 V_2) + 90T (V_1^2 V_2) + 130T (V_1^2 V_2) \quad (30)$$

$$M_4^0 = 2g_2^2 M_1 \quad (31)$$

$$M_5^0 = \frac{2}{3} M_1 (8g_2^2 M_1 + 120g_2^2 M_1 + 9g_2^2 M_1 + 250g_2^2 M_1 + 120g_2^2 M_2 - 10V(M_1, \lambda - T)) - 30M_1 T (V_1 V_2^2) - 10M_1 T (V_1 V_2^2) - 30M_1 T (V_1 V_2^2) + 30T (V_1^2 V_2) + 30T (V_1^2 V_2) \quad (32)$$

$$M_6^0 = -9g_2^2 M_1 \quad (33)$$

$$M_7^0 = \frac{2}{3} M_1 (11g_2^2 M_1 + 11g_2^2 M_1 + 40g_2^2 M_1 + 140g_2^2 M_1 + 45g_2^2 M_2 - 20M_1 T (V_1 V_2^2) - 20M_1 T (V_1 V_2^2) + 20T (V_1^2 V_2) + 20T (V_1^2 V_2)) \quad (34)$$

**3.4 Trilinear Superpotenzial Parameters**

$$A_1^0 = 2V_1 V_2 + V_1 (-3g_2^2 + 2T (V_1 V_2^2)) - \frac{16}{3} M_1^2 - \frac{7}{3} M_1^2 + 10T (V_1 V_2^2) + V_1 V_2 V_2 \quad (35)$$

$$A_2^0 = \frac{2}{3} M_1^2 (2V_1 V_2^2 + V_1 (-10V_1 V_2^2 V_2 - 0V_1 V_2 V_2 V_2 - 2V_1 V_2 V_2 V_2) - 2V_1 V_2 V_2 V_2 + V_1 V_2 V_2 V_2 + V_1 V_2 V_2 V_2 (-3V_1^2 - 3T (V_1 V_2^2) + 6g_2^2 - 9T (V_1 V_2^2) + \frac{2}{3} M_1^2) - 30M_1 T (V_1 V_2^2) \quad (36)$$

$$+ V_1 (2M_1^2 g_2^2 + 6g_2^2 + \frac{13}{3} M_1^2 + \frac{8}{3} M_1^2 + 8g_2^2 - \frac{16}{3} M_1^2 - 21M_1^2 V - 3V_1^2 V^2 - \frac{2}{3} M_1^2 T (V_1 V_2^2) + 16g_2^2 T (V_1 V_2^2) + \frac{2}{3} M_1^2 T (V_1 V_2^2) - 3M_1^2 T (V_1 V_2^2) - 9T (V_1 V_2^2 V_2 V_2^2) - 9T (V_1 V_2 V_2 V_2^2)) \quad (37)$$

$$A_3^0 = 3V_1 V_2^2 + V_1 (-3g_2^2 + 2T (V_1 V_2^2)) - \frac{16}{3} M_1^2 + 10T (V_1 V_2^2) \quad (38)$$

$$A_4^0 = -9V_1 V_2 V_2^2 + V_1 V_2^2 (-3V_1^2 - 3T (V_1 V_2^2) + 6g_2^2 - 9T (V_1 V_2^2)) \quad (39)$$

$$+ V_1 (\frac{27}{3} M_1^2 + \frac{8}{3} M_1^2 + \frac{13}{3} M_1^2 - 21M_1^2 V - 3V_1^2 V^2 - \frac{2}{3} M_1^2 T (V_1 V_2^2) + 16g_2^2 T (V_1 V_2^2) + \frac{2}{3} M_1^2 T (V_1 V_2^2) - 3M_1^2 T (V_1 V_2^2) - 9T (V_1 V_2 V_2 V_2^2) - 9T (V_1 V_2 V_2 V_2^2)) \quad (40)$$

**3.5 Bilinear Superpotenzial Parameters**

$$B_1^0 = 2M_1^2 V - 3g_2^2 + 30T (V_1 V_2^2) + 30T (V_1 V_2^2) - \frac{2}{3} M_1^2 + 2T (V_1 V_2^2) \quad (41)$$

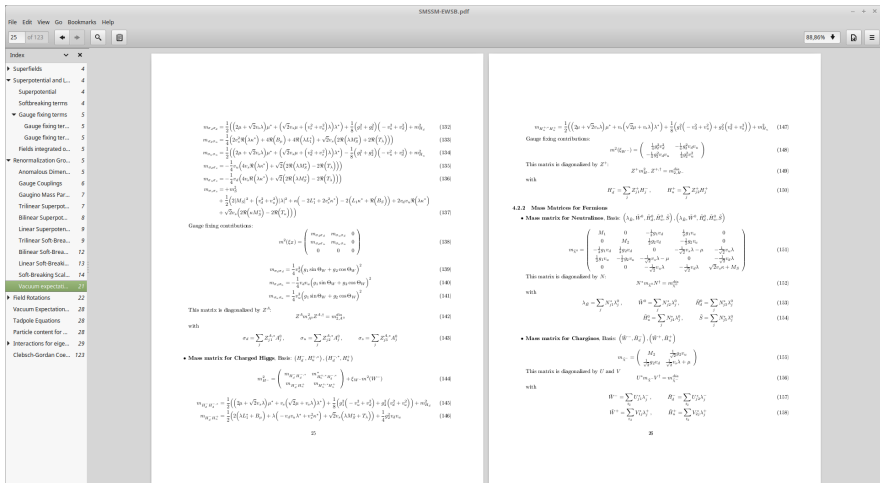
$$B_2^0 = \frac{1}{20} (20M_1^2 V + 90g_2^2 V + 375g_2^2 - 200M_1^2 V - 200M_1^2 V^2 - 20M_1^2 V^2) - 80g_2^2 T (V_1 V_2^2) - 80g_2^2 T (V_1 V_2^2) - 30M_1^2 T (V_1 V_2^2) - 10T (V_1 V_2^2) - 30T (V_1 V_2^2) - 80g_2^2 T (V_1 V_2^2) - 80g_2^2 T (V_1 V_2^2) - 30M_1^2 T (V_1 V_2^2) - 30M_1^2 T (V_1 V_2^2) - 30M_1^2 T (V_1 V_2^2) - 30M_1^2 T (V_1 V_2^2) \quad (42)$$

$$B_3^0 = \frac{2}{3} M_1^2 (2V_1 V_2^2 + V_1 (-10V_1 V_2^2 V_2 - 0V_1 V_2 V_2 V_2 - 2V_1 V_2 V_2 V_2) - 2V_1 V_2 V_2 V_2 + V_1 V_2 V_2 V_2 + V_1 V_2 V_2 V_2 (-3V_1^2 - 3T (V_1 V_2^2) + 6g_2^2 - 9T (V_1 V_2^2) + \frac{2}{3} M_1^2) - 30M_1 T (V_1 V_2^2) - 30M_1 T (V_1 V_2^2) - 30M_1 T (V_1 V_2^2) + 30T (V_1^2 V_2) + 30T (V_1^2 V_2)) \quad (43)$$

$$B_4^0 = 6M_1^2 (V^2 + V^2) \quad (44)$$

# Example

## Mass matrices ...



The image shows a PDF viewer window titled "SMSSM-EWSB.pdf" with a zoom level of 88.80%. The left sidebar contains a table of contents with "Nonlinear Equations" highlighted in green. The main content area displays two pages of mathematical derivations.

**Page 1 (Left):**

- Equation (132): 
$$m_{\nu_{\mu\nu}} = \frac{1}{2}((3\alpha + \sqrt{3}\alpha_3)\nu^2 + (\sqrt{3}\alpha_3 + (\tilde{d} + \tilde{e})\lambda)\nu^2) + \frac{1}{2}(\tilde{d} + \tilde{e})(-\tilde{e} + \tilde{d}) + m_{\tilde{d}\tilde{e}}$$
- Equation (133): 
$$m_{\nu_{\tau\nu}} = \frac{1}{2}(2\alpha_3(\lambda\nu^2) + 4\theta(\lambda\nu) + 4\theta(\lambda\tilde{e}) + \sqrt{3}\alpha_3(2\theta(\lambda\tilde{d}) + 2\theta(\tilde{e}\lambda)))$$
- Equation (134): 
$$m_{\nu_{\tau\mu}} = \frac{1}{2}((3\alpha + \sqrt{3}\alpha_3)\nu^2 + (\sqrt{3}\alpha_3 + (\tilde{d} + \tilde{e})\lambda)\nu^2) - \frac{1}{2}(\tilde{d} + \tilde{e})(-\tilde{e} + \tilde{d}) + m_{\tilde{d}\tilde{e}}$$
- Equation (135): 
$$m_{\nu_{\tau\nu}} = -\frac{1}{2}\nu_1(\nu_1\theta(\lambda\nu^2) + \sqrt{3}(2\theta(\lambda\tilde{d}) - 2\theta(\tilde{e}\lambda)))$$
- Equation (136): 
$$m_{\nu_{\tau\mu}} = -\frac{1}{2}\nu_1(\nu_1\theta(\lambda\nu^2) + \sqrt{3}(2\theta(\lambda\tilde{d}) - 2\theta(\tilde{e}\lambda)))$$
- Equation (137): 
$$= \frac{1}{2}(3(\lambda\tilde{d}\tilde{e}^2 + (\tilde{d} + \tilde{e})\lambda^2) + \nu_1^2(-3\tilde{d} + 2\tilde{e}\lambda^2) - 2(\tilde{e}\lambda^2 + \theta(\lambda\nu_1)) + 2\nu_1\theta(\lambda\nu^2) + \sqrt{3}\alpha_3(2\theta(\lambda\tilde{d}) - 2\theta(\tilde{e}\lambda)))$$
- Grouping contribution: 
$$m^{\nu}(x) = \begin{pmatrix} m_{\nu_{\mu\nu}} & m_{\nu_{\tau\nu}} & 0 \\ m_{\nu_{\tau\nu}} & m_{\nu_{\tau\mu}} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$
 (138)
- Equation (139): 
$$m_{\nu_{\mu\nu}} = \frac{1}{2}\nu_1^2(\nu_1 \cos \theta_{\nu\mu} + \nu_2 \sin \theta_{\nu\mu})^2$$
- Equation (140): 
$$m_{\nu_{\tau\nu}} = -\frac{1}{2}\nu_1^2(\nu_1 \sin \theta_{\nu\mu} + \nu_2 \cos \theta_{\nu\mu})^2$$
- Equation (141): 
$$m_{\nu_{\tau\mu}} = -\frac{1}{2}\nu_1^2(\nu_1 \sin \theta_{\nu\mu} + \nu_2 \cos \theta_{\nu\mu})^2$$
- Text: "This matrix is diagonalized by  $Z^{\nu}$ " (142)
- Equation (143): 
$$Z^{\nu} \nu_{\nu}^{\nu} Z^{\nu\dagger} = \nu_{\nu}^{\nu}$$
- Text: "with" (144)
- Equation (145): 
$$\nu_{\mu} = \sum_{\nu} Z_{\mu\nu}^{\nu} \nu_{\nu}^{\nu}, \quad \nu_{\tau} = \sum_{\nu} Z_{\tau\nu}^{\nu} \nu_{\nu}^{\nu}, \quad \nu_{\nu} = \sum_{\nu} Z_{\nu\nu}^{\nu} \nu_{\nu}^{\nu}$$
- Section: "• Mass matrix for Charged Higgs Bos. ( $H_1^{\pm}, H_2^{\pm}$ ), ( $H_3^{\pm}, H_4^{\pm}$ )" (146)
- Equation (147): 
$$m_{H_{1,2}^{\pm}} = \begin{pmatrix} m_{H_1^{\pm}} & m_{H_1^{\pm} H_2^{\pm}} \\ m_{H_1^{\pm} H_2^{\pm}} & m_{H_2^{\pm}} \end{pmatrix} + (2\theta - \nu_1^2 \theta^2 \mathbb{1}^{\pm})$$
 (146)
- Equation (148): 
$$m_{H_{1,2}^{\pm}} = \frac{1}{2}((3\alpha + \sqrt{3}\alpha_3)\nu^2 + \nu_1(\sqrt{3}\alpha_3 + \lambda)\nu^2) + \frac{1}{2}(\tilde{d}(-\tilde{e} + \tilde{d}) + \tilde{e}(\tilde{d} + \tilde{e})) + m_{\tilde{d}\tilde{e}}$$
 (148)
- Equation (149): 
$$m_{H_{3,4}^{\pm}} = \frac{1}{2}(\tilde{d}(\lambda\tilde{e}_1 + \tilde{e}_1\lambda) + \lambda(-\nu_1\nu_2\nu^2 + 2\tilde{e}\nu^2) + \sqrt{3}\alpha_3(\lambda M_2 + \tilde{e}_1^2)) + \frac{1}{2}\tilde{d}^2\nu_1$$
 (149)

**Page 2 (Right):**

- Equation (147): 
$$m_{H_{3,4}^{\pm}} = \frac{1}{2}((3\alpha + \sqrt{3}\alpha_3)\nu^2 + \nu_1(\sqrt{3}\alpha_3 + \lambda)\nu^2) + \frac{1}{2}(\tilde{d}(-\tilde{e} + \tilde{d}) + \tilde{e}(\tilde{d} + \tilde{e})) + m_{\tilde{d}\tilde{e}}$$
 (147)
- Text: "Grouping contribution:" (148)
- Equation (149): 
$$m^{\nu}(x) = \begin{pmatrix} m_{H_1^{\pm}} & m_{H_1^{\pm} H_2^{\pm}} \\ m_{H_1^{\pm} H_2^{\pm}} & m_{H_2^{\pm}} \end{pmatrix} + \begin{pmatrix} -\frac{1}{2}\tilde{d}^2\nu_1 & 0 \\ 0 & -\frac{1}{2}\tilde{d}^2\nu_1 \end{pmatrix}$$
 (149)
- Text: "This matrix is diagonalized by  $Z^{\nu}$ " (150)
- Equation (151): 
$$Z^{\nu} \nu_{H_1^{\pm}}^{\nu} Z^{\nu\dagger} = \nu_{H_1^{\pm}}^{\nu}$$
 (151)
- Text: "with" (152)
- Equation (153): 
$$\lambda_{\pm} = \sum_{\nu} Z_{\nu\pm}^{\nu} \lambda_{\nu}^{\nu}, \quad \theta_{\pm}^{\nu} = \sum_{\nu} Z_{\nu\pm}^{\nu} \theta_{\nu}^{\nu}, \quad \tilde{d}_{\pm}^{\nu} = \sum_{\nu} Z_{\nu\pm}^{\nu} \tilde{d}_{\nu}^{\nu}$$
 (153)
- Equation (154): 
$$\tilde{e}_{\pm}^{\nu} = \sum_{\nu} Z_{\nu\pm}^{\nu} \tilde{e}_{\nu}^{\nu}, \quad \delta_{\pm}^{\nu} = \sum_{\nu} Z_{\nu\pm}^{\nu} \delta_{\nu}^{\nu}$$
 (154)
- Section: "• Mass matrix for Charged Bos. ( $H^{\pm}, A_1^{\pm}$ )" (155)
- Equation (156): 
$$m_{H^{\pm}} = \begin{pmatrix} M_1 & \frac{2\theta\nu_1\nu_2}{\sqrt{2}} \\ \frac{2\theta\nu_1\nu_2}{\sqrt{2}} & \frac{2\theta^2\nu_1\nu_2}{\sqrt{2}} + \nu^2 \end{pmatrix}$$
 (155)
- Text: "This matrix is diagonalized by U and V" (156)
- Equation (157): 
$$U^{\nu} \nu_{H^{\pm}}^{\nu} U^{\nu\dagger} = \nu_{H^{\pm}}^{\nu}, \quad B_1^{\pm} = \sum_{\nu} U_{\nu 1}^{\nu} \nu_{\nu}^{\nu}$$
 (157)
- Equation (158): 
$$V^{\nu} \nu_{A_1^{\pm}}^{\nu} V^{\nu\dagger} = \nu_{A_1^{\pm}}^{\nu}, \quad B_2^{\pm} = \sum_{\nu} V_{\nu 2}^{\nu} \nu_{\nu}^{\nu}$$
 (158)

# Example

## Vertices ...

SMSSM-EWSB.pdf



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
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$$\frac{1}{2} \epsilon_{123} (\sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} + \lambda (\kappa_1 \tilde{A}_1^2 + \kappa_2 \tilde{A}_2^2) \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i})$$
$$+ \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} \tilde{A}_{1i} \quad (181)$$

$$\frac{1}{2} (\epsilon_{123}^2 \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} + \lambda (\kappa_1 \tilde{A}_1^2 + \kappa_2 \tilde{A}_2^2) \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i})$$
$$+ \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} \tilde{A}_{1i} \quad (182)$$


$$\frac{1}{2} \epsilon_{123} (\sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} + \lambda (\kappa_1 \tilde{A}_1^2 + \kappa_2 \tilde{A}_2^2) \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i})$$
$$+ \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} - \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} \tilde{A}_{1i} \quad (183)$$

$$\frac{1}{2} (\epsilon_{123}^2 (\kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} + \lambda (\kappa_1 \tilde{A}_1^2 + \kappa_2 \tilde{A}_2^2) \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i})$$
$$+ 2 \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} + \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$+ 2 \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} + 2 \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- 2 \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i} - 2 \kappa_1 \kappa_2^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- 2 \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 \sum_{k=1}^6 Y_{i,j,k} \tilde{T}_{i,j,k} \tilde{A}_{1i} - 2 \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$- 2 \sqrt{2} \kappa_1^2 \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i}$$
$$+ \lambda (\kappa_1^2 \tilde{A}_1^2 + (\kappa_1 \tilde{A}_1^2 + \kappa_2 \tilde{A}_2^2) \sum_{i=1}^6 \sum_{j=1}^6 Y_{i,j} \kappa_{i,j}^2 \tilde{A}_{1i})$$
$$+ \kappa_1^2 (\lambda (-2 \kappa_1 + \kappa_2) \tilde{A}_1^2 - 2 (2 \kappa_1 + \sqrt{2} \kappa_2) \tilde{A}_1^2 + (2 \kappa_1 + \sqrt{2} \kappa_2) \tilde{A}_2^2)$$

- SARAH makes it easy to get analytical information about a given model
- RGE running can easily be performed within `Mathematica` at the two-loop level
- The  $\text{\LaTeX}$  output provides many information in a human readable form