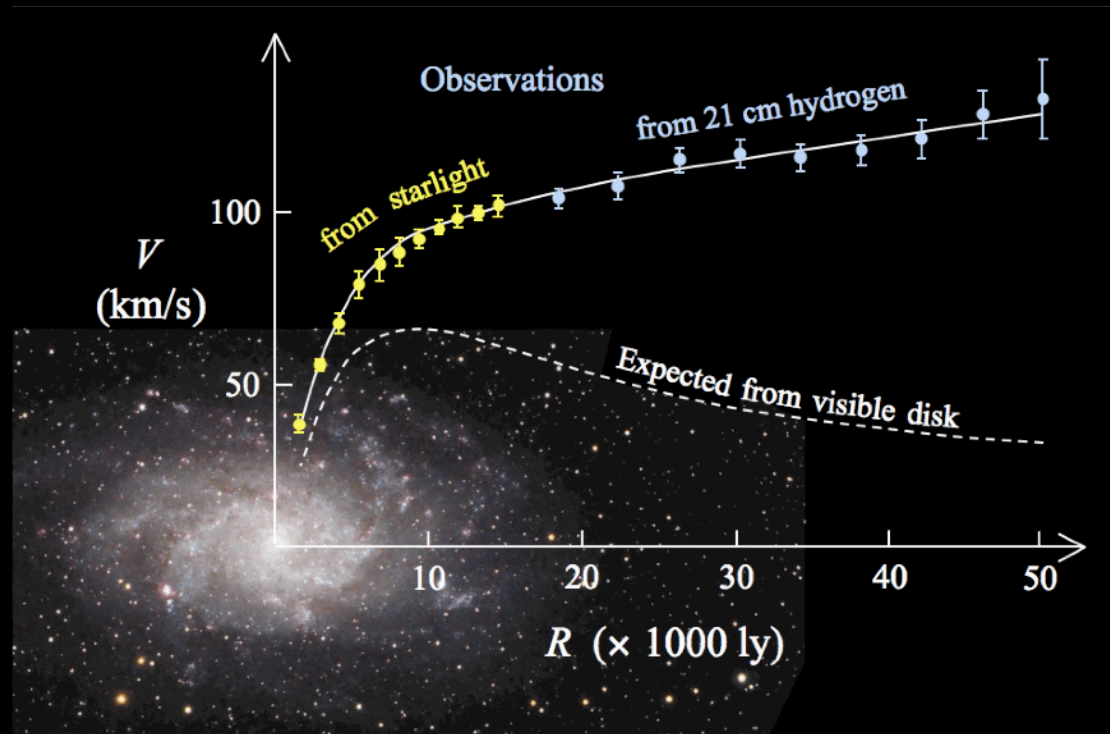


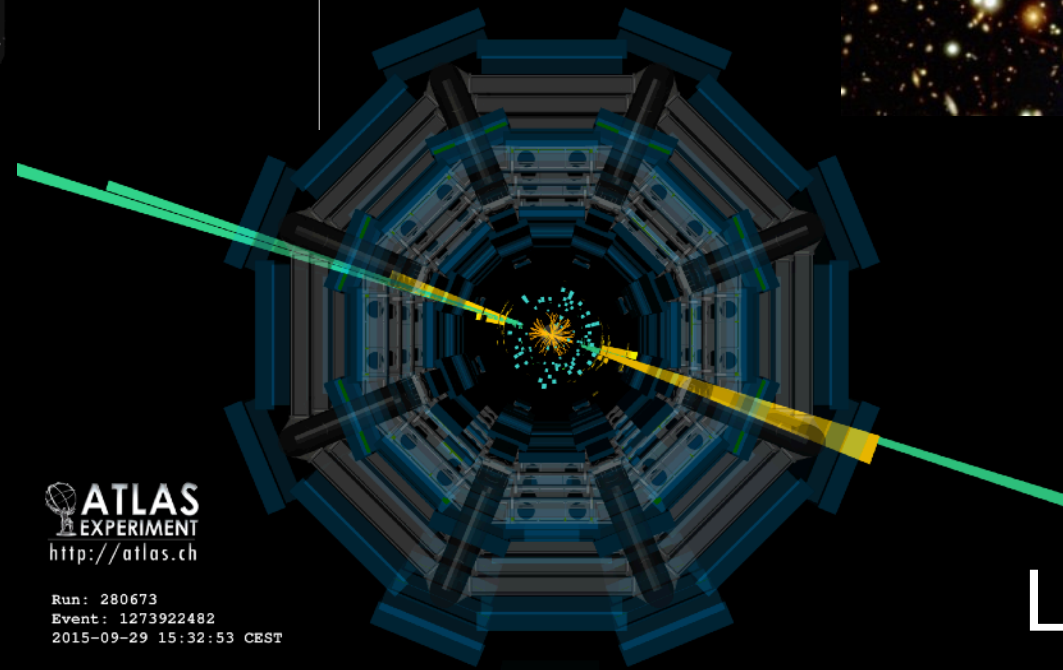
Dark Matter searches with the ATLAS Detector



DIS Kobe
18.04.18

ATLAS
EXPERIMENT
<http://atlas.ch>

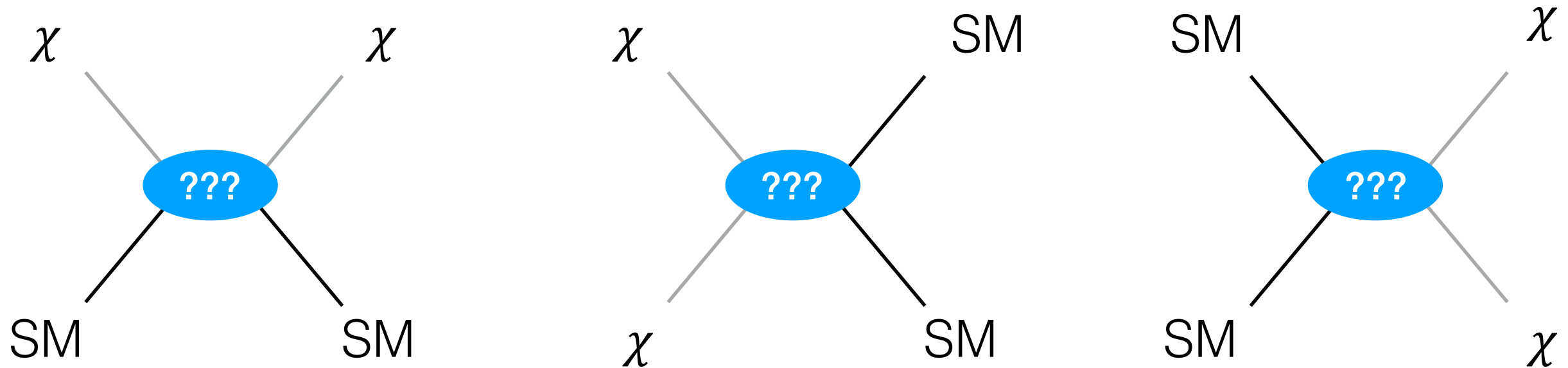
Run: 280673
Event: 1273922482
2015-09-29 15:32:53 CEST



Will Kalderon,
Lund University (SE)

on behalf of the ATLAS Collaboration

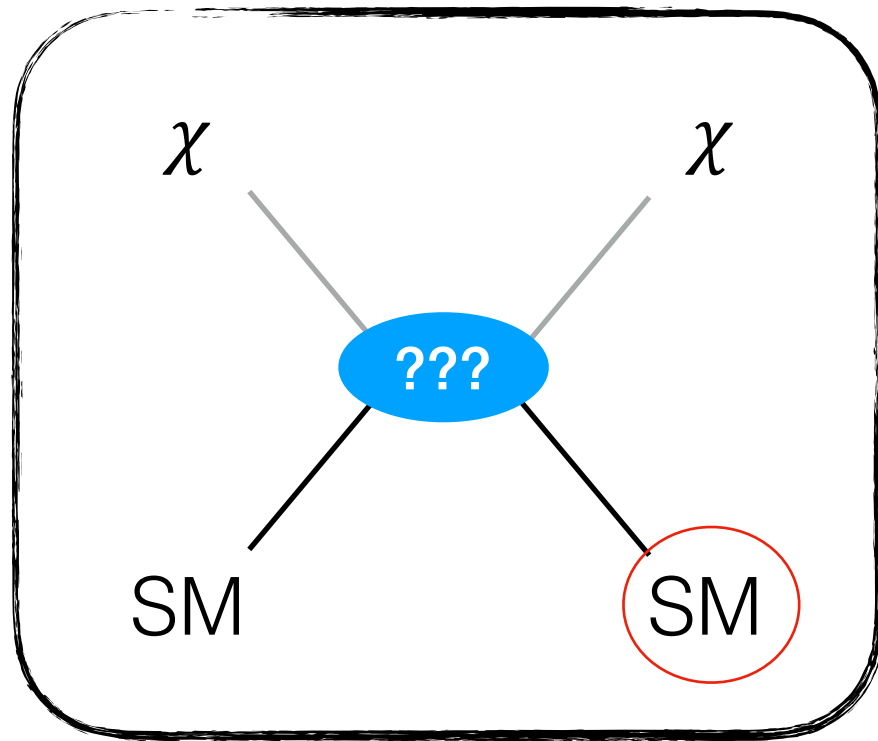
DM signatures



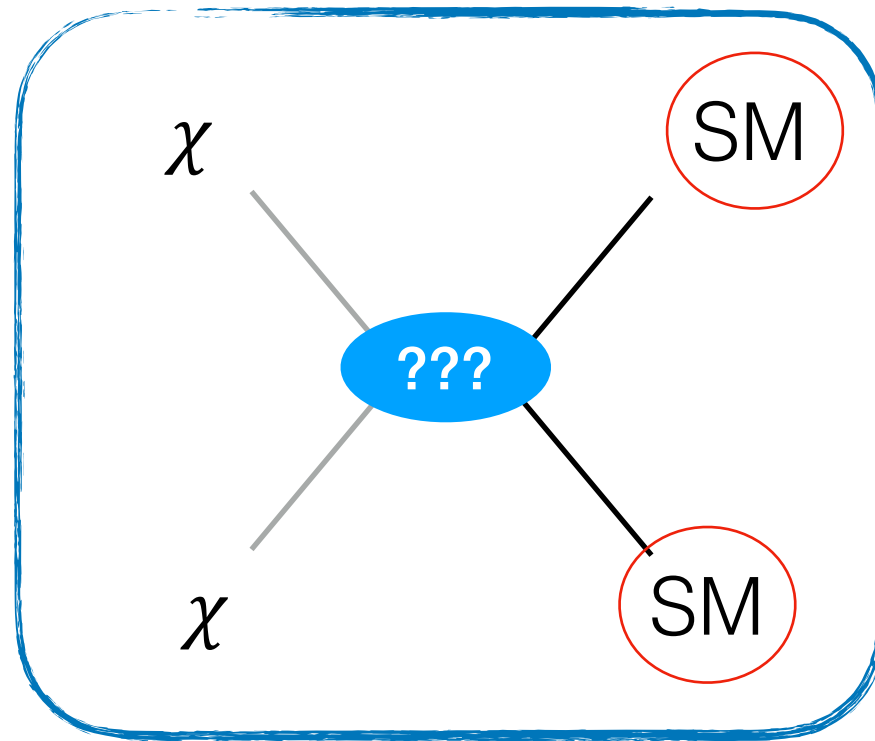
??? =

Gravitational	✓
Electromagnetic	✗
Strong	✗
Weak-strength	?

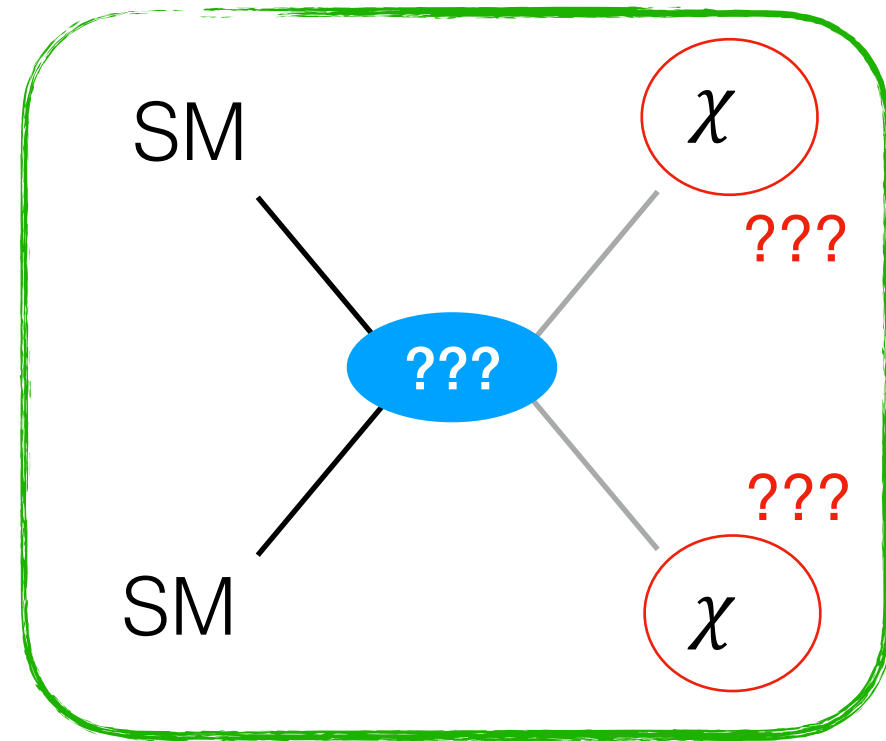
How to search for them



Direct detection



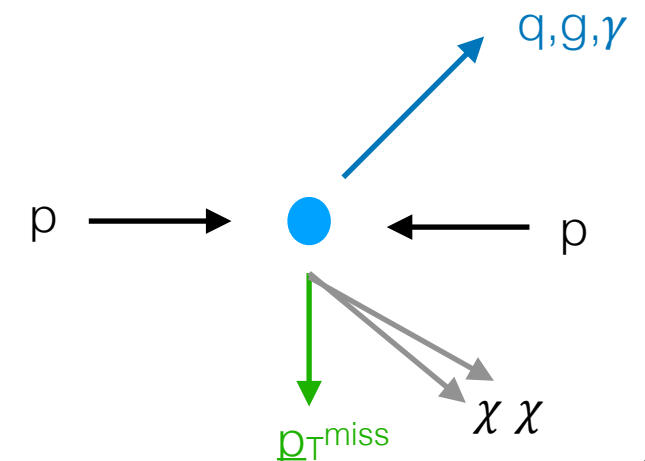
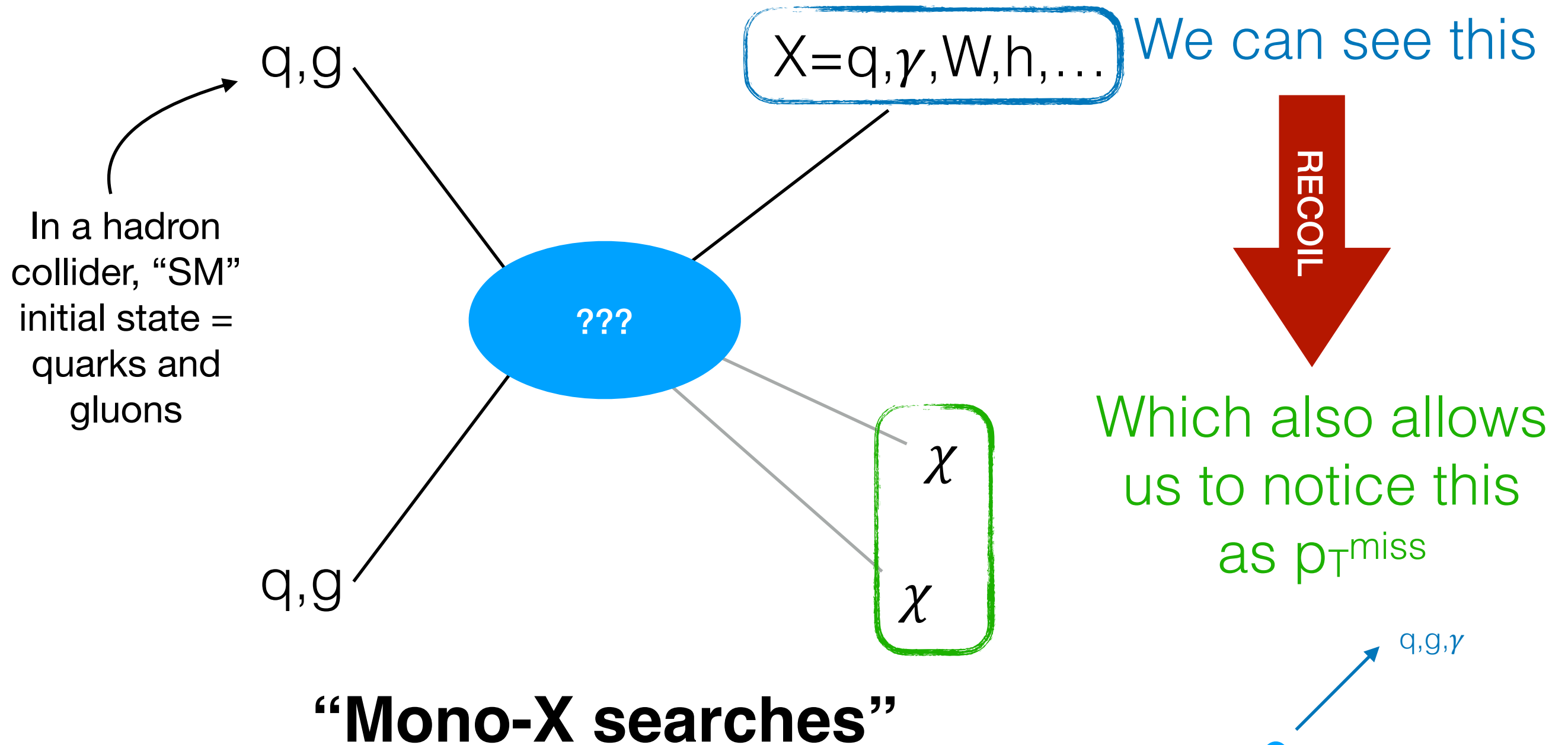
Indirect detection



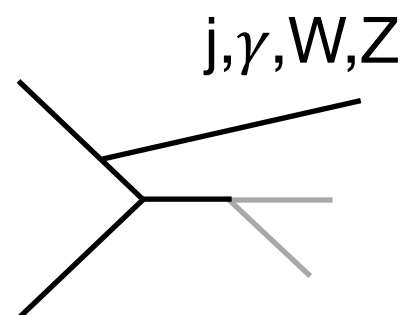
Collider production

Collider: how do we search for nothing?

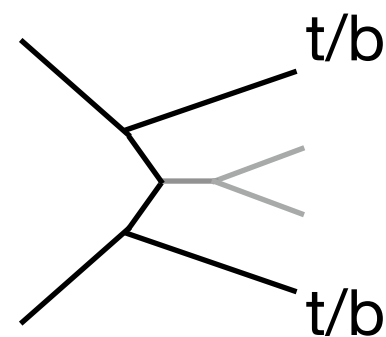
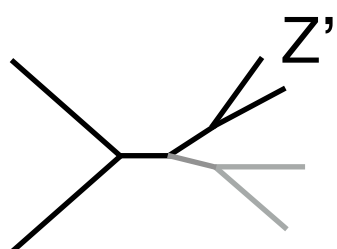
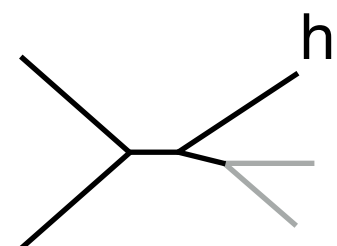
Option 1: require something to happen!



ATLAS mono-X / associated production



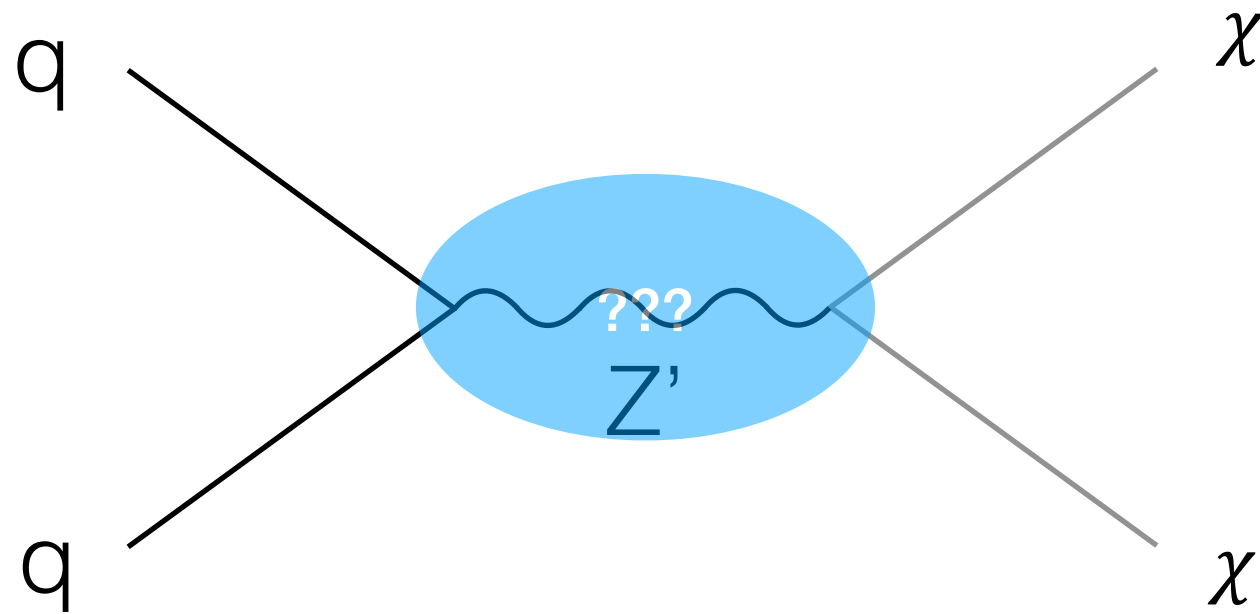
mono-X	Dataset	Reference
jet	36.1 fb ⁻¹	JHEP 01 (2018) 126
γ	36.1 fb ⁻¹	Eur. Phys. J. C 77 (2017) 393
$Z (\rightarrow \ell\ell)$	36.1 fb ⁻¹	PLB 776 (2017) 318
$W/Z (\rightarrow qq)$	3.2 fb ⁻¹	PLB 763 (2016) 251
$h (\rightarrow bb)$	36.1 fb ⁻¹	PRL 119 (2017) 181804
$h (\rightarrow \gamma\gamma)$	36.1 fb ⁻¹	PRD 96, 112004 (2017)
$Z' (\rightarrow qq)$	36.1 fb ⁻¹	ATLAS-CONF-2018-005



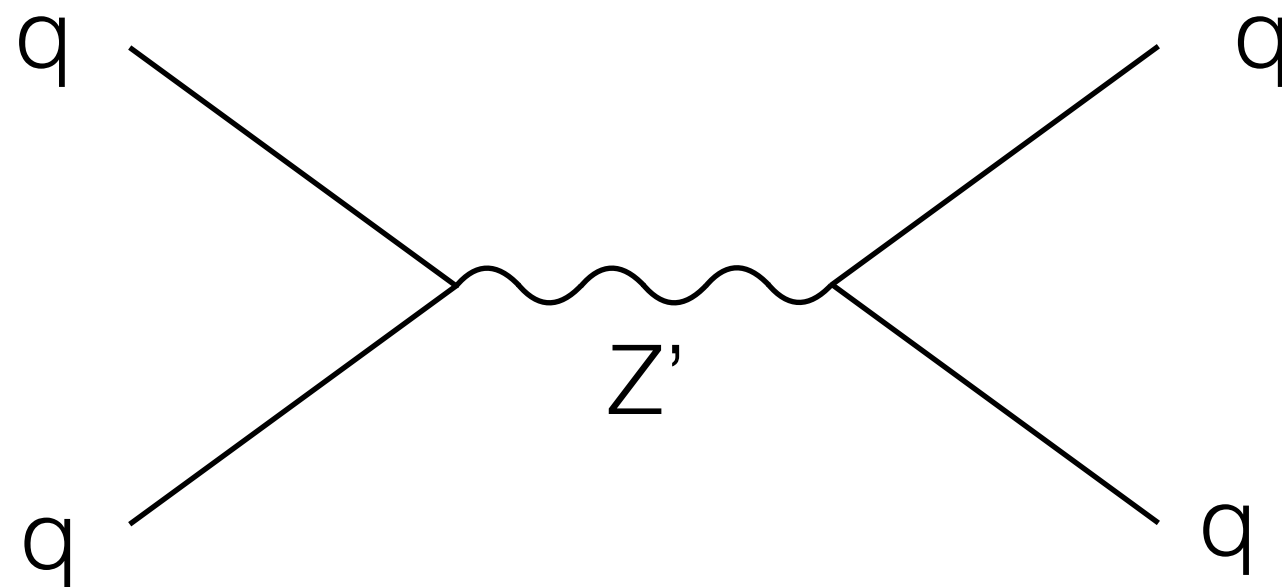
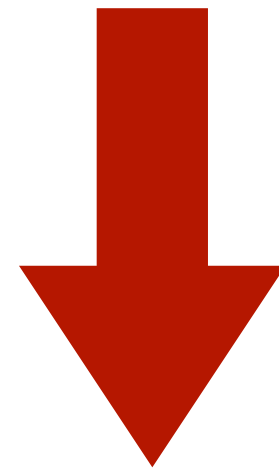
Associated production	Dataset	Reference
tt/bb/b+MET	36.1 fb ⁻¹	Eur. Phys. J. C 77 (2017) 393

scalar mediator,
3rd-gen couplings

Option 2: dark matter? What dark matter?



If there is a mediator
that couples to
quarks and DM...



... then we can
forget about the
DM and look for
the mediator

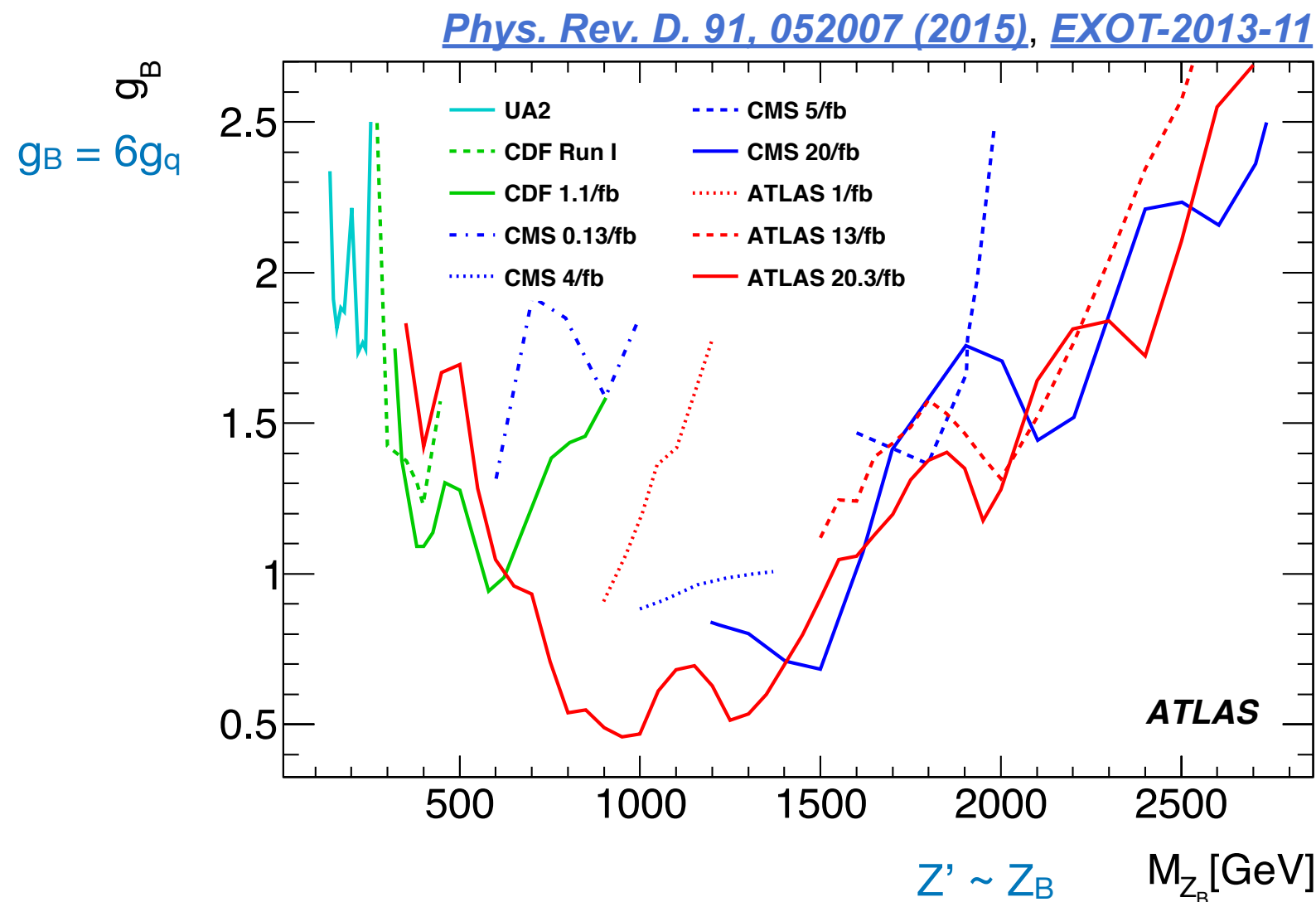
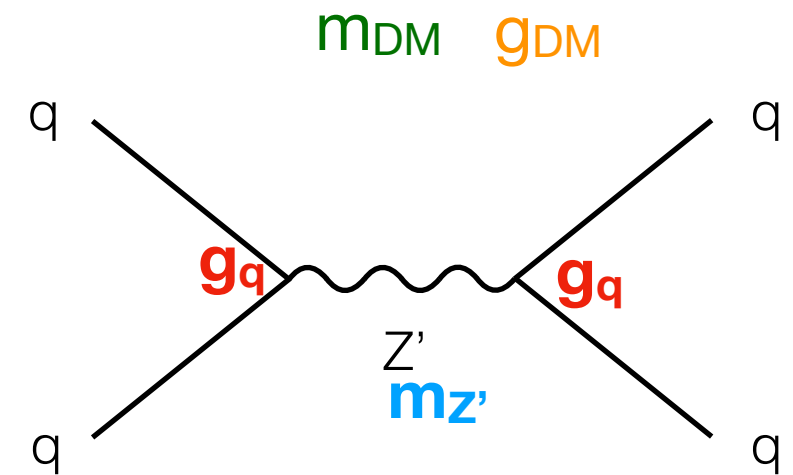
“Dijet* resonance searches”

*One can also imagine the Z' coupling to leptons \rightarrow dilepton resonances, lower BR to dijet

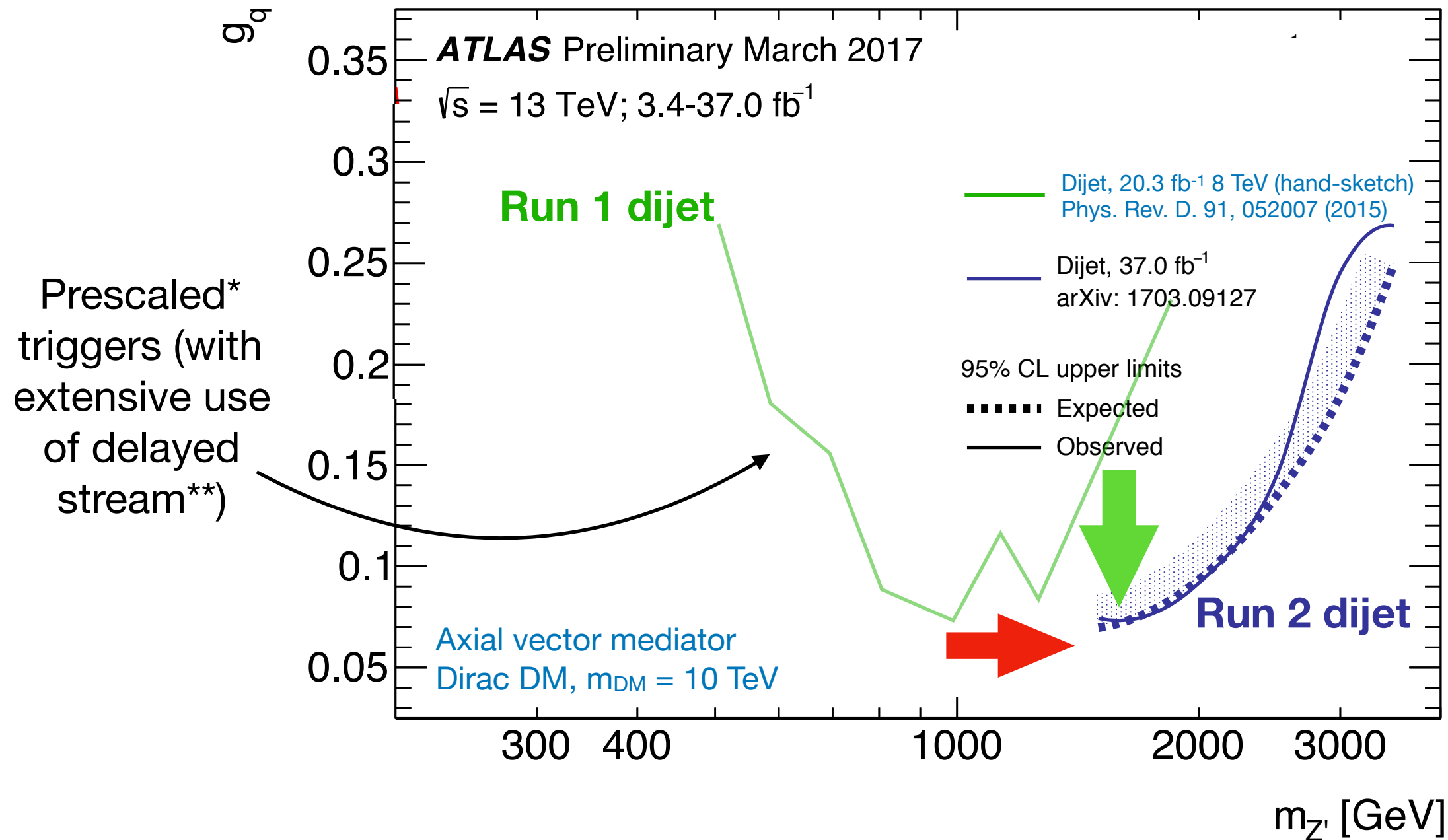
Dijet limits on Z' , at end of run 1

Model has **four parameters**:

1. Mediator coupling to quarks g_q (usually assumed universal, but dijets ignore $Z' \rightarrow t\bar{t}$)
2. Mediator mass $m_{Z'}$
3. Dark Matter mass m_{DM} - set well above $0.5 m_{Z'}$ (eg 10 TeV) \rightarrow kinematically inaccessible
4. Mediator coupling to Dark Matter, g_{DM} - not very relevant given 3, often set to 1



Dijet limits, run 1 vs run 2



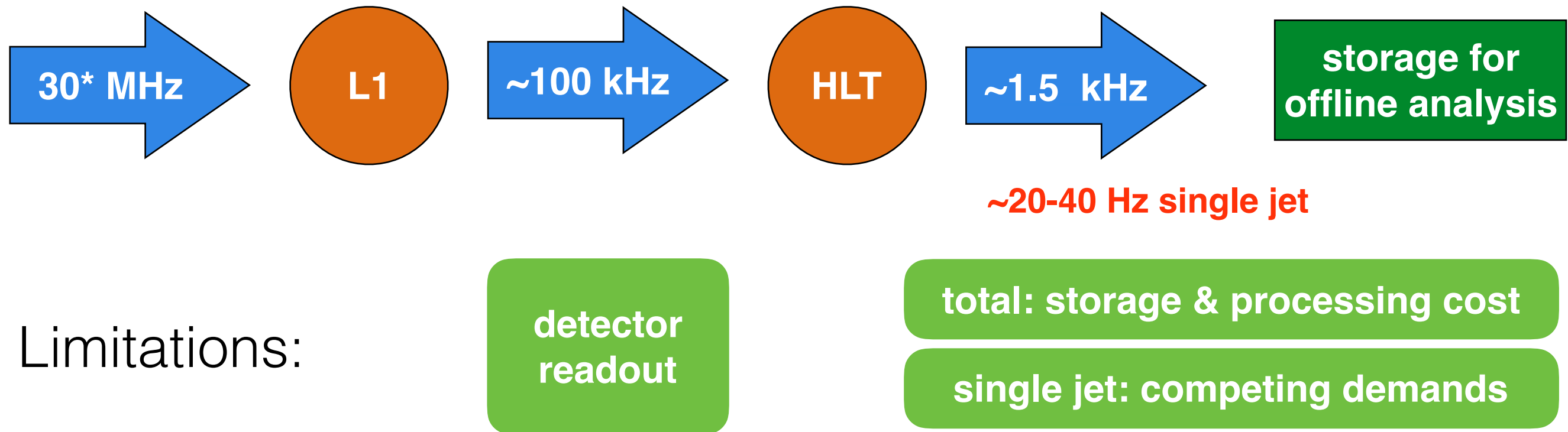
**Lower coupling g_q for given mass $m_{Z'}$:
 more data ($\sigma \sim g_q^2 \Rightarrow \text{limit}(g_q) \sim \text{data}^{1/4}$), better mass resolution**

Higher bottom mass edge to exclusion: trigger limitations

* Prescaled: only a fraction of events accepted by a trigger are recorded

** Delayed stream: events accepted by some triggers are written to a separate stream that is not reconstructed until computing resources become available over a shutdown

What limits the ATLAS trigger?

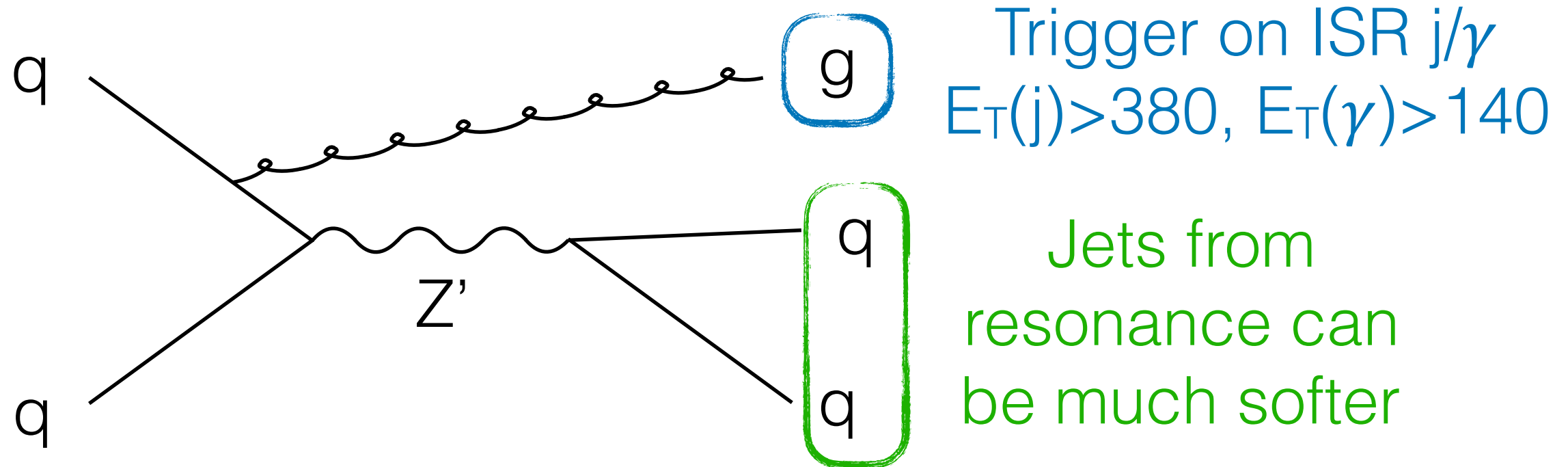


Higher instantaneous luminosity -> higher rate of high- p_T jet production

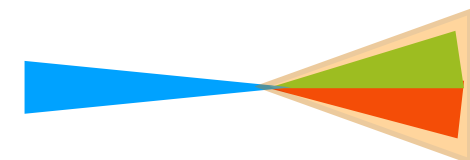
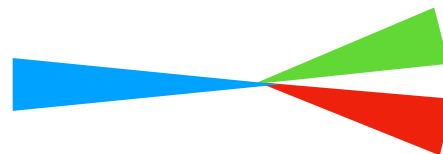
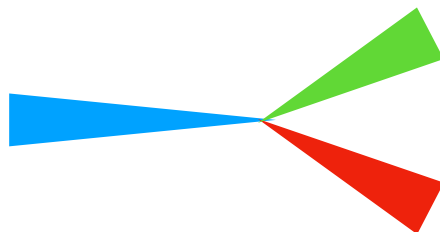
=> with rising instantaneous luminosity, must raise jet p_T threshold for recording events

- Empirical observation: at high p_T (> 100 GeV or so), rate $\sim p_T^{-5}$
- 2016: record events containing jets with $E_T > 380$ GeV -> efficient by $p_T > 440$ GeV in analysis

Overcoming trigger 1: ISR

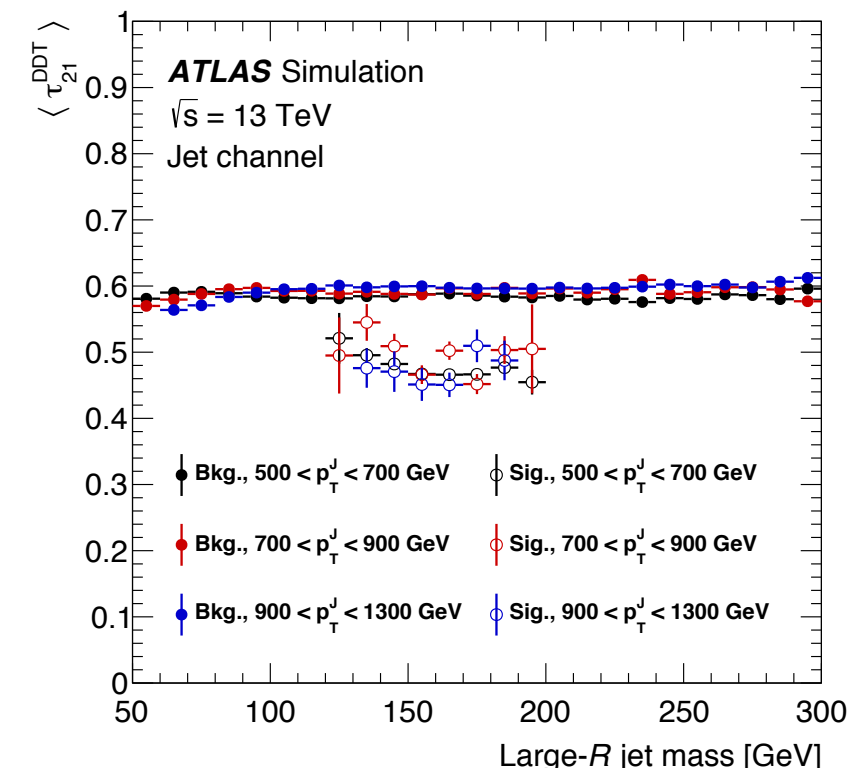


- ATLAS has preliminary results ([ATLAS-CONF-2016-070](#)) using photon and jet using initial state radiation to trigger on => resonance can be much lower p_T (lead resonance jet $p_T > 25$ GeV, vs 440 GeV)
- At Z' masses below ~ 200 GeV, resonance jets merge -> large-R jet

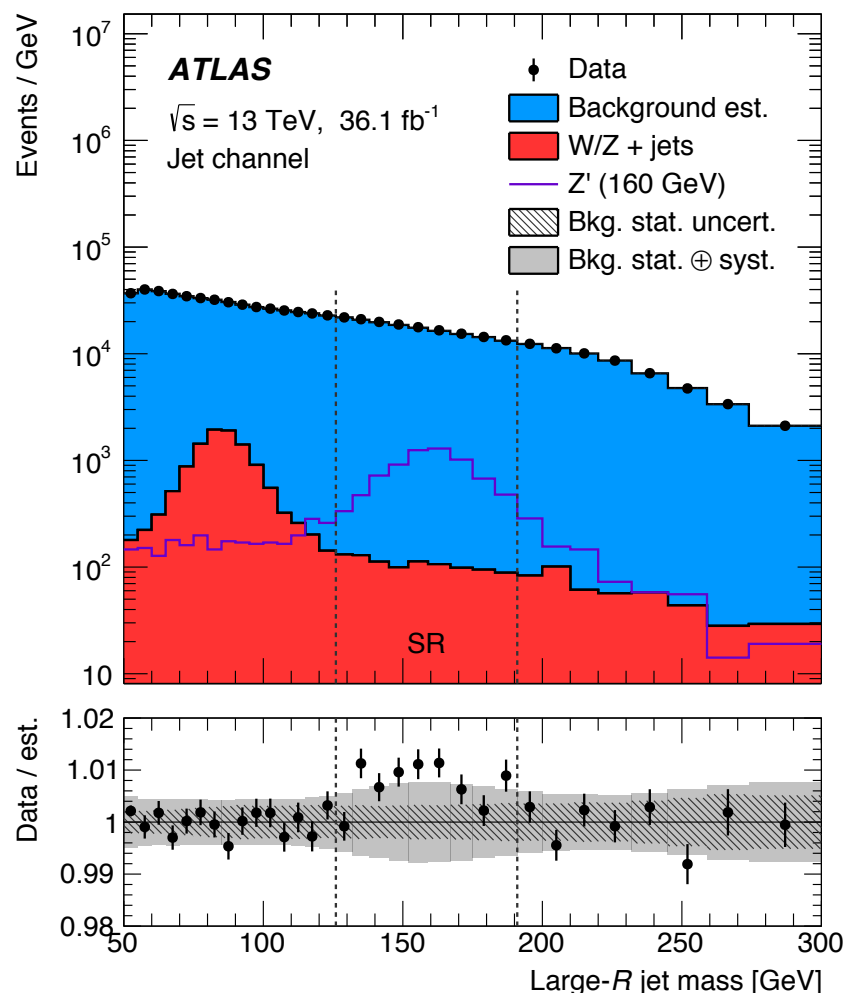


Quick overview: Large-R + ISR

arxiv: 1801.08768, [EXOT-2017-01](#)

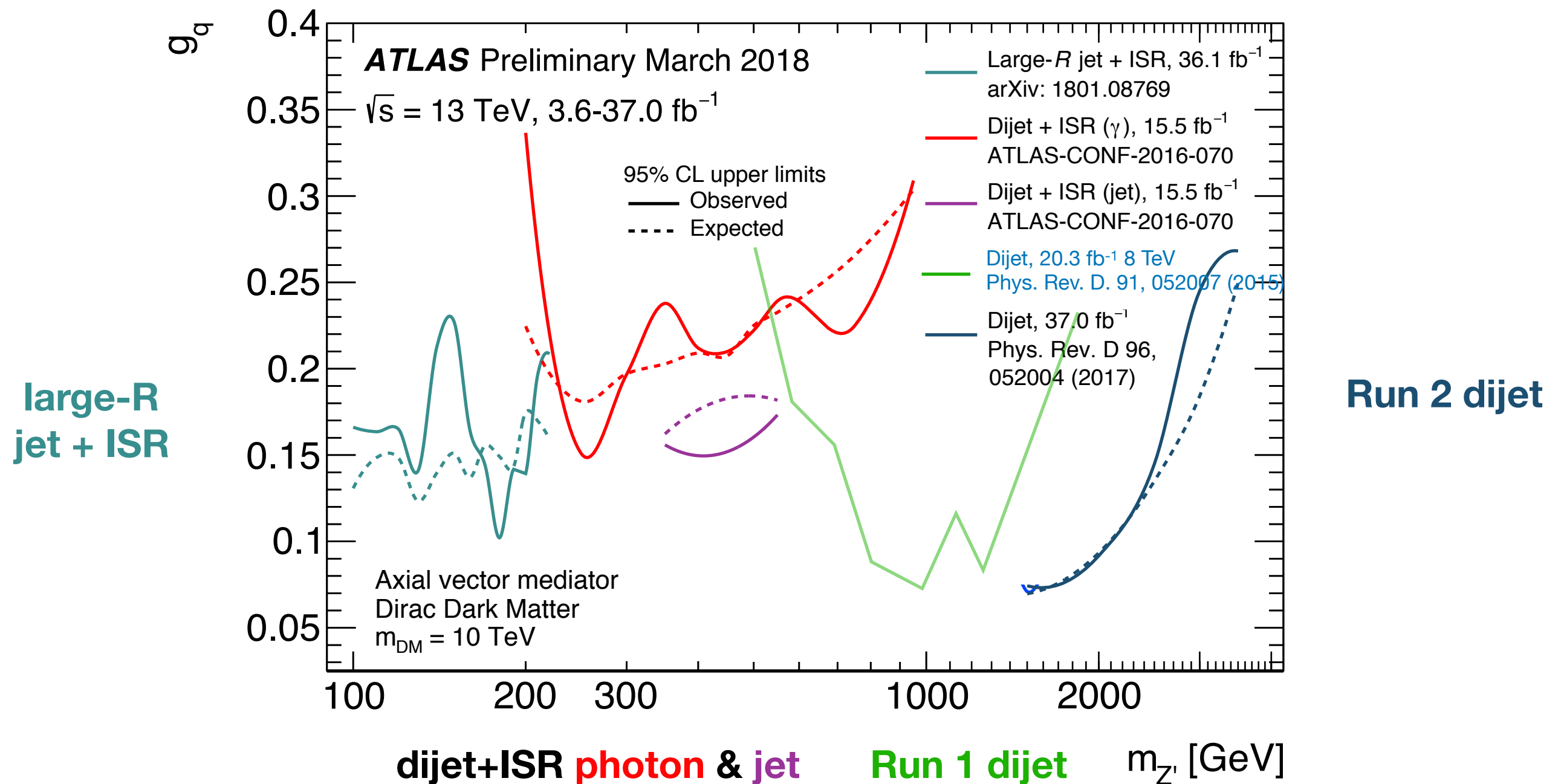


- Use substructure τ_{21} to distinguish 2-subjet signal from single-subjet QCD background
- Use “designed decorrelated tagger” method to decorrelate from jet mass



- Main background QCD
- Data-driven method for background estimation based on inverted τ_{21}^{DDT}
- Method validated on W/Z peak
- Separate signal region for each mass point

Dijet (merged & resolved) + ISR limits

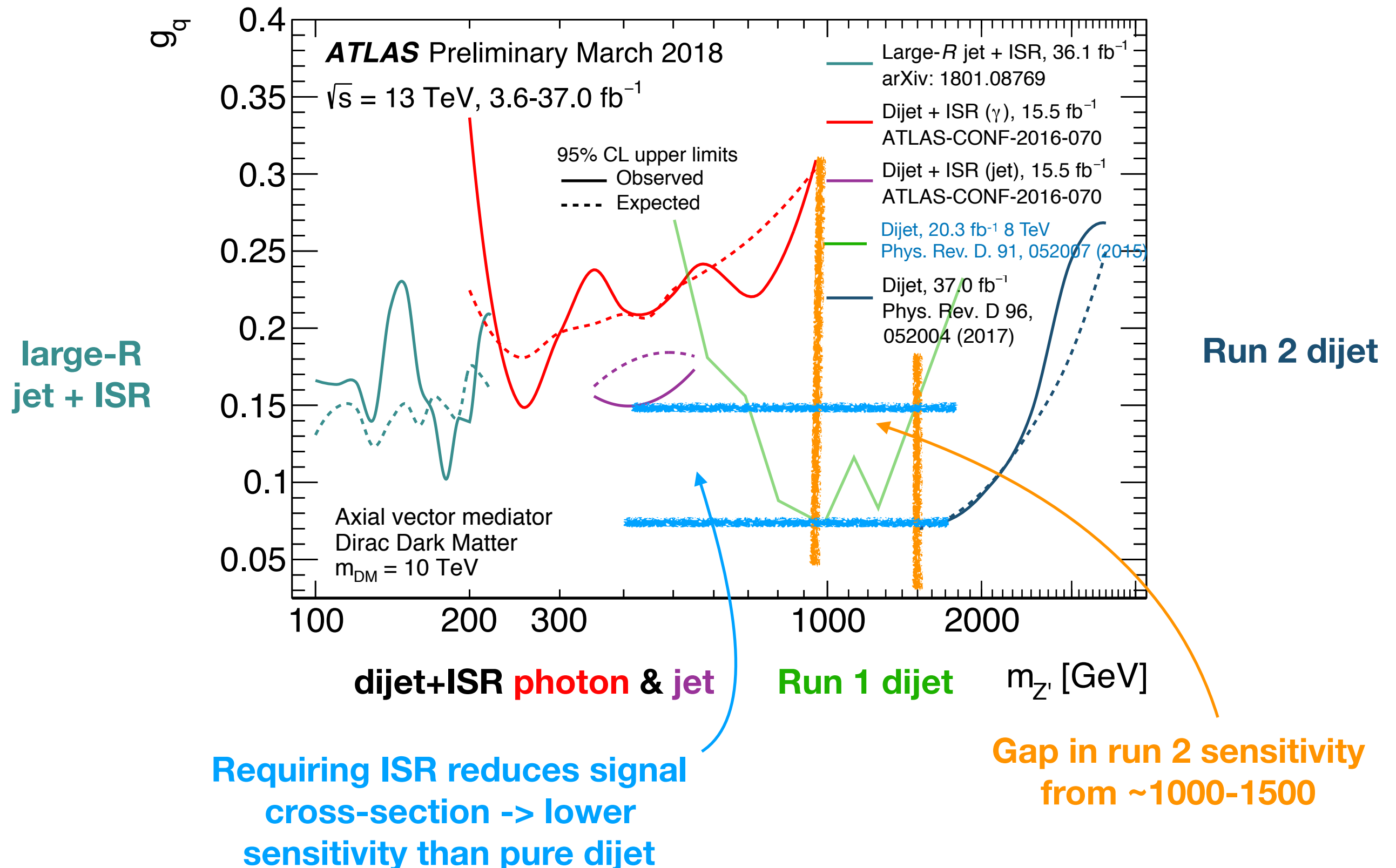


ISR -> sensitivity down to 200 GeV

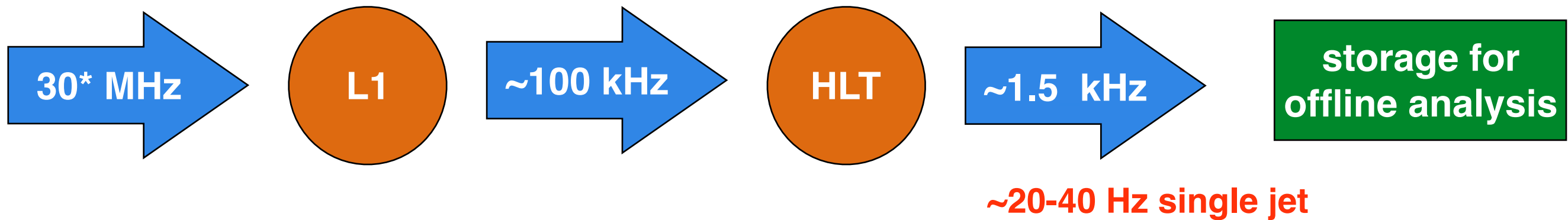
200 GeV = crossover between merged and resolved

Large- R jet -> takes this down to 100 GeV

Dijet (merged & resolved) + ISR limits



2: Revisit trigger limitations



Limitations:

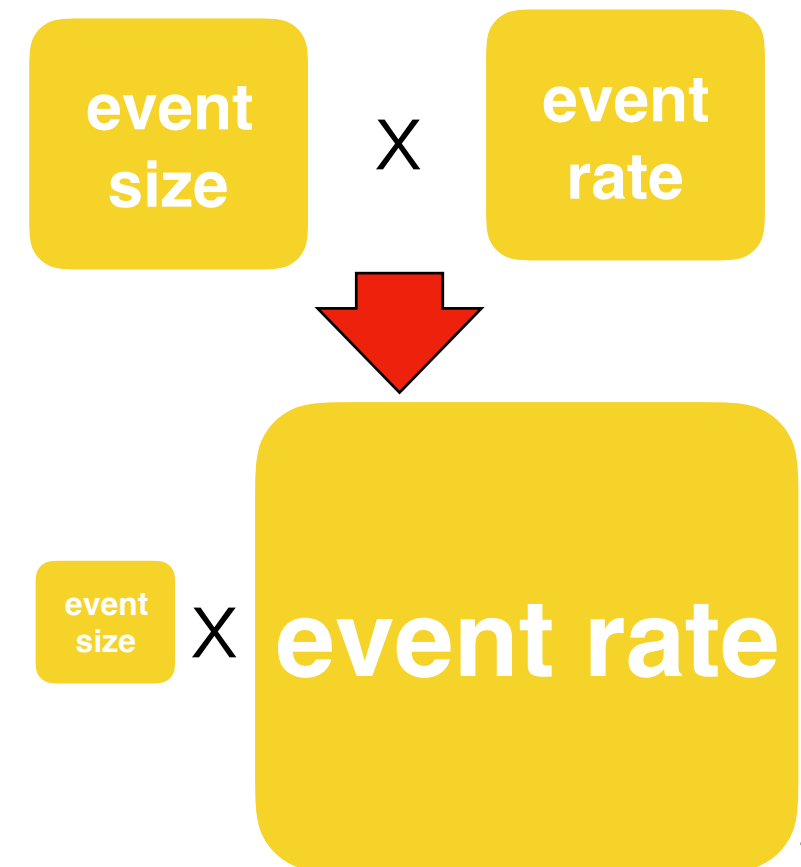
detector
readout

total: storage & processing cost

single jet: competing demands

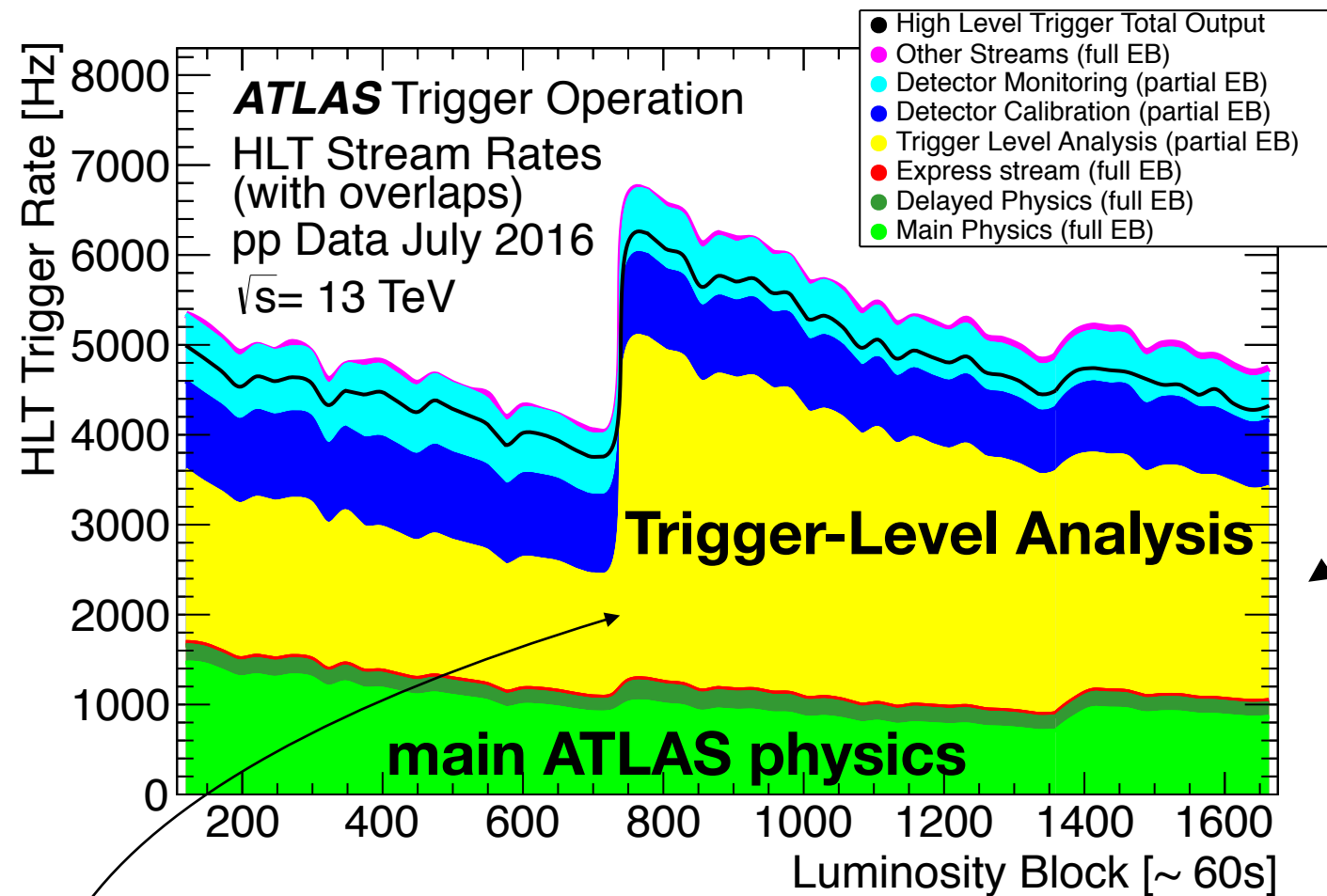
Storage and processing drives 1.5 kHz limit for ATLAS

- dijet resonance search only uses jets - no leptons, no p_T^{miss} , etc.
- we already build and calibrate jets in the trigger... just save these
- **record minimal events at high rate**



* 25ns bunch spacing gives 40 MHz, but the ring is not full

Evade trigger bandwidth limits



Instead of 20-40 Hz for a dijet resonance search, we now have 1-3 kHz!

jump: sometimes recorded more TLA data once the luminosity had fallen

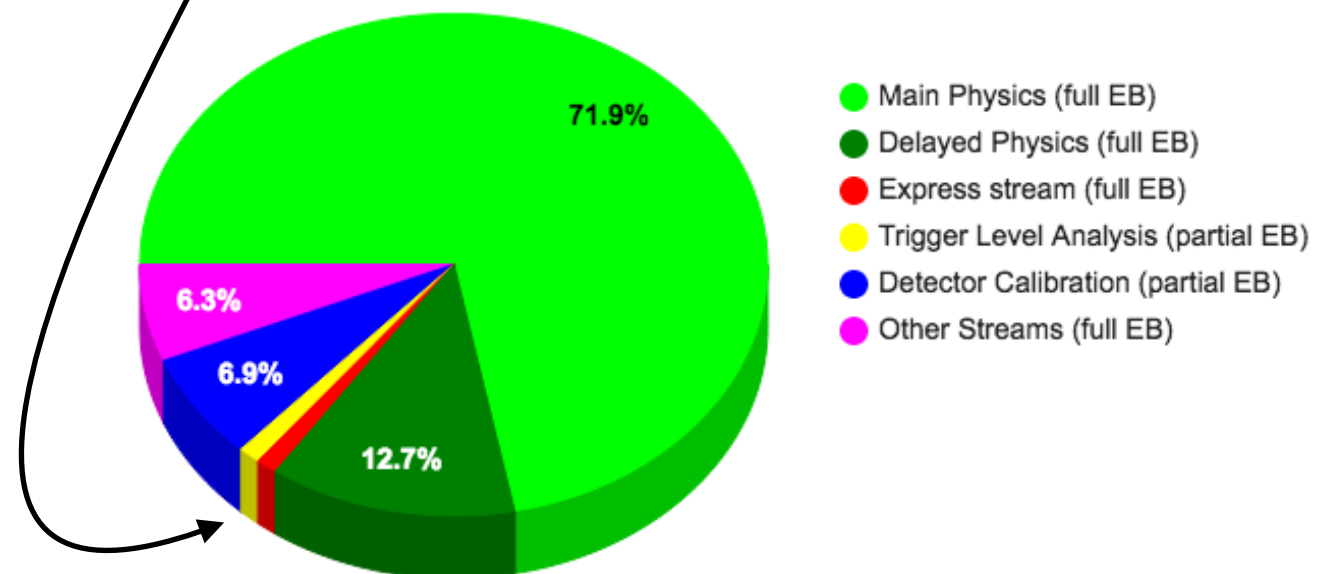
LHCb: “Turbo stream” [1]

CMS: “Data Scouting” [2]

ATLAS: “Trigger Level Analysis”
(arXiv: 1804.03496, April 11th!)

Huge TLA rate but tiny bandwidth since ~0.5% of full event size

ATLAS Trigger Operation
pp Data July 2016, $\sqrt{s} = 13$ TeV

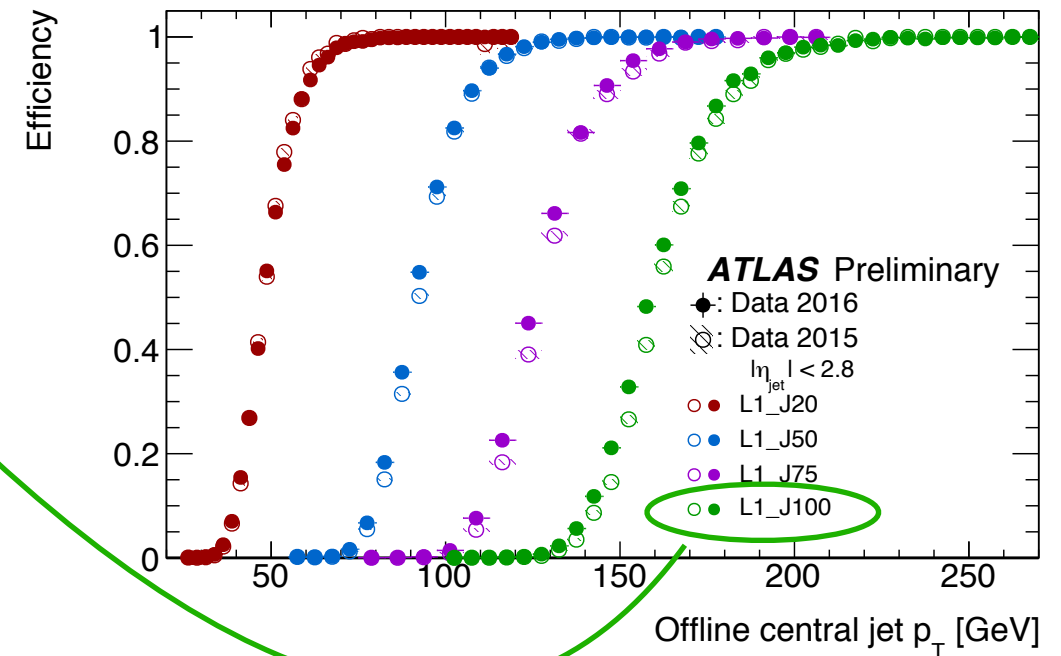


[1] LHCb Collaboration, *Tesla : an application for real-time data analysis in High Energy Physics*, *Comput. Phys. Commun.* **208** (2016) 35, arXiv: 1604.05596 [physics.ins-det].

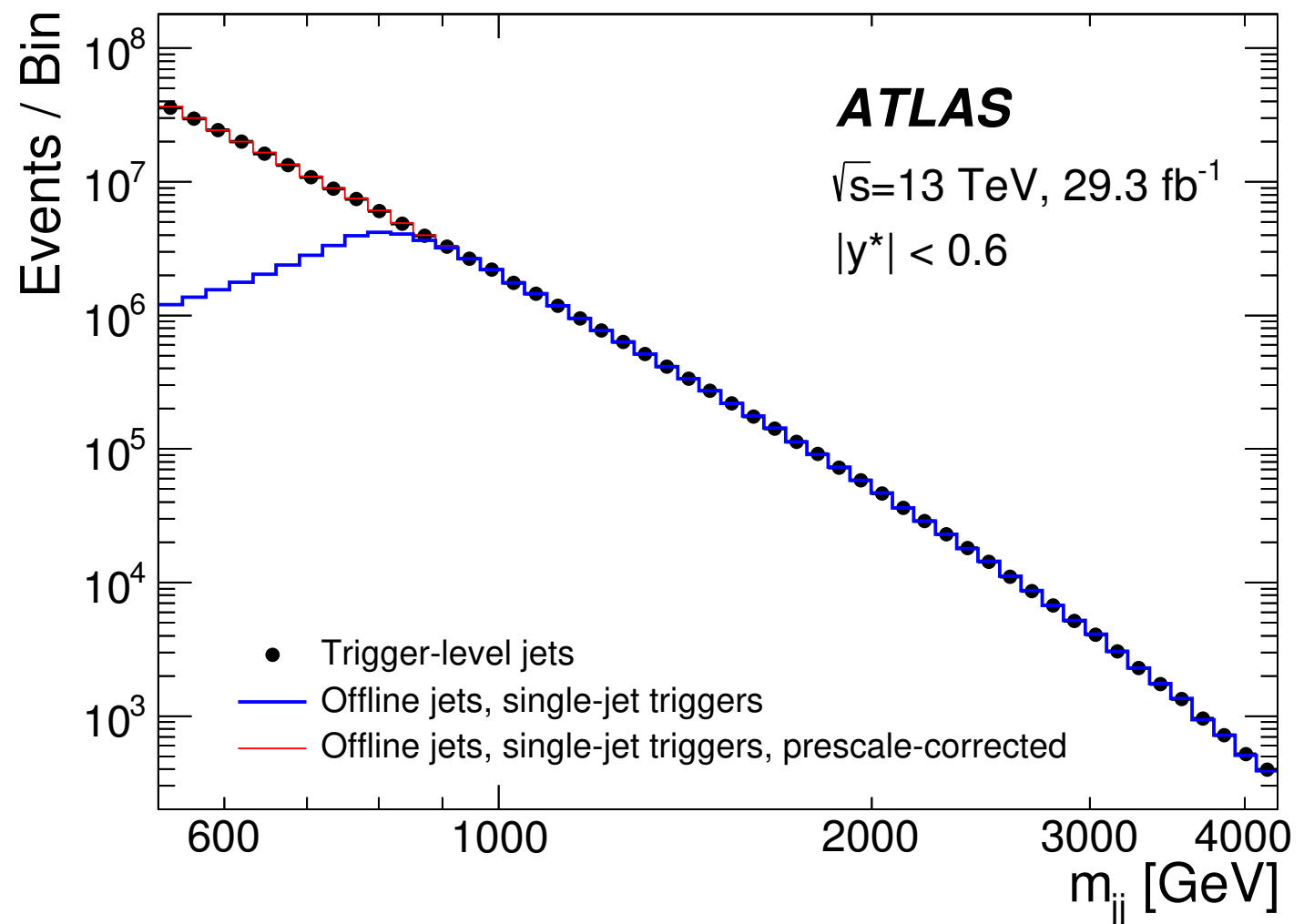
[2] CMS Collaboration, *Search for dijet resonances in proton–proton collisions at $\sqrt{s} = 13$ TeV and constraints on dark matter and other models*, *Phys. Lett. B* **769** (2017) 520, arXiv: 1611.03568 [hep-ex].

The payoff

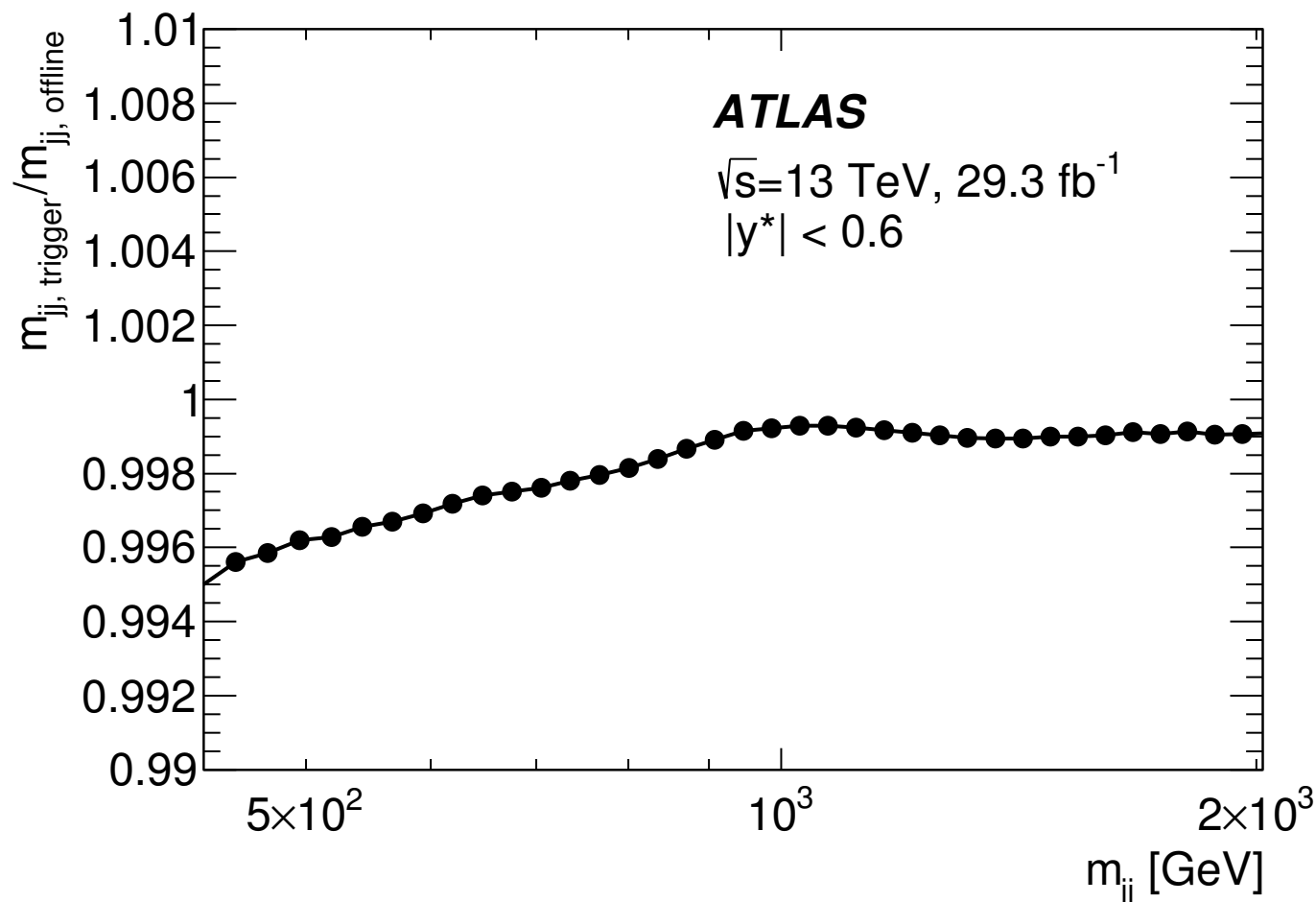
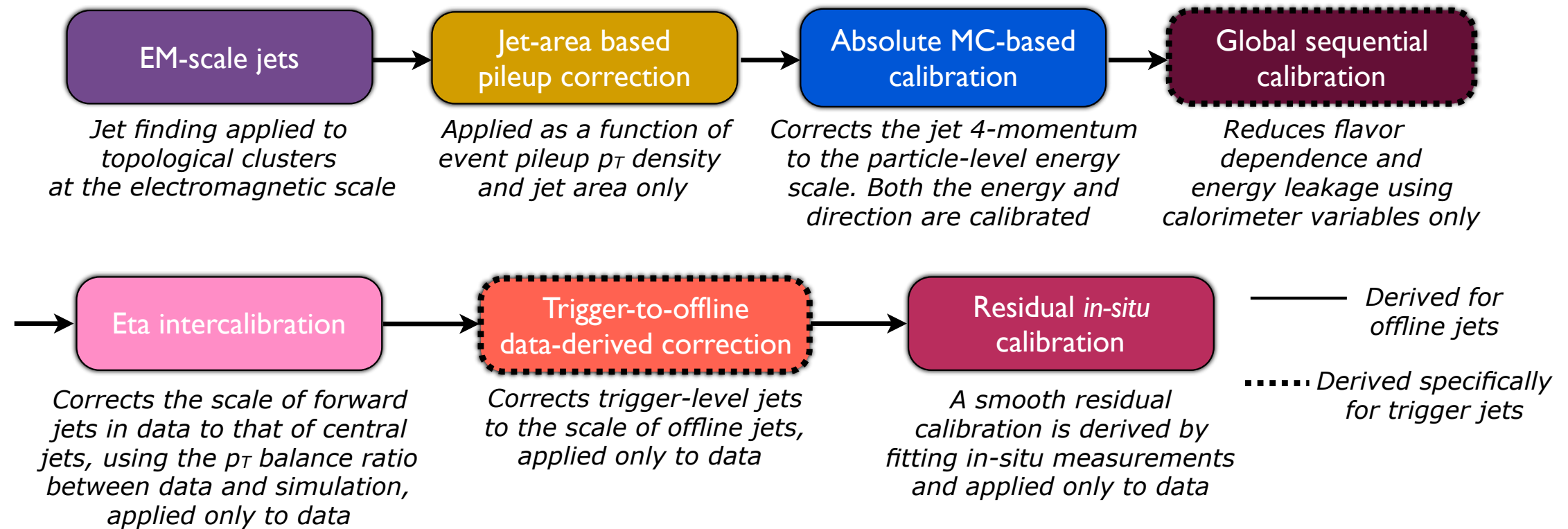
	“standard” dijet	TLA
lead jet $p_T >$	440	220
sublead jet $p_T >$	60	85
$m_{jj} >$	1100	520



**4×10^7 events in first bin
in 29.3 fb^{-1} of 2016 data**



TLA calibration

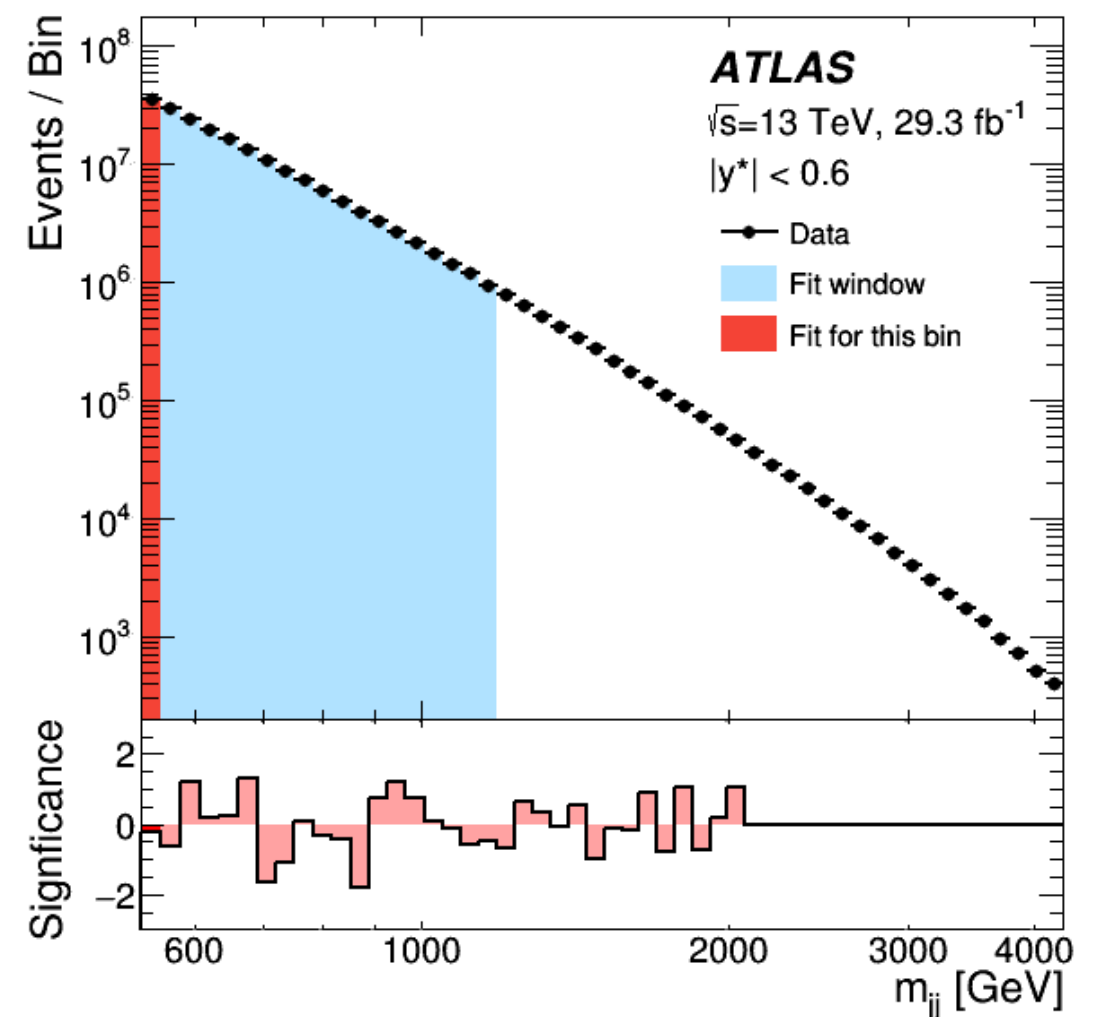


- Write out sufficient information to be able to redo calibration offline
- Some parts rederived since TLA data lacks eg track information
- End result: excellent agreement between offline and recalibrate trigger m_{jj}

Background estimation

- Fit to functional form
 - Choose one with best χ^2
- Very large number of events -> very little scope for QCD to deviate from functional form
- In 2015, could not fit whole m_{jj} range, hence truncated fit at 1250 GeV
- Solution, also used by high-mass dijet 37 fb⁻¹ result: fit sub-ranges

Functional form
$f(x) = p_1(1-x)^{p_2}x^{p_3}$
$f(x) = p_1(1-x)^{p_2}x^{p_3+p_4 \ln x}$
$f(x) = p_1(1-x)^{p_2}x^{p_3+p_4 \ln x} * p_5 \ln x^2$
$f(x) = \frac{p_1}{x^{p_2}} e^{-p_3 x - p_4 x^2}$

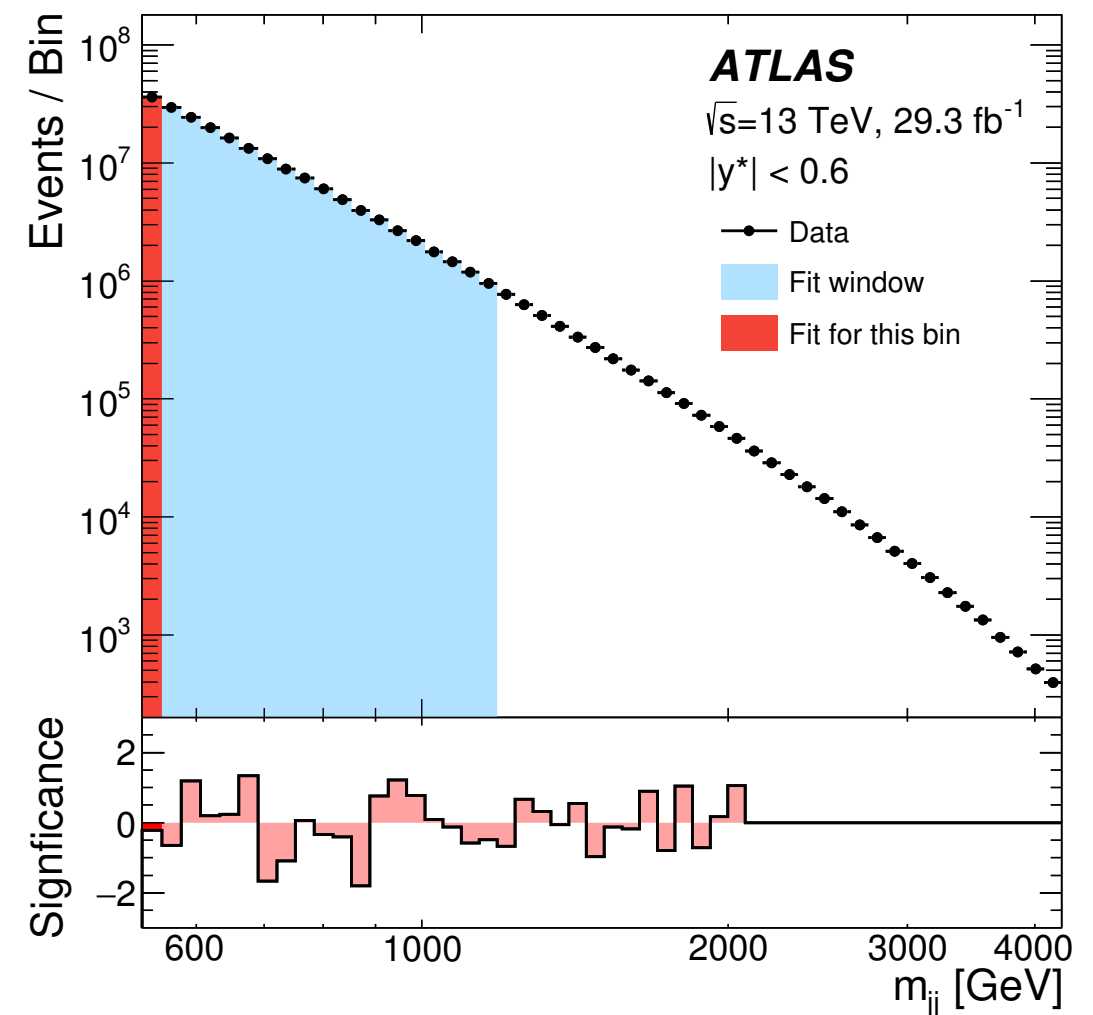


[animation here](#)

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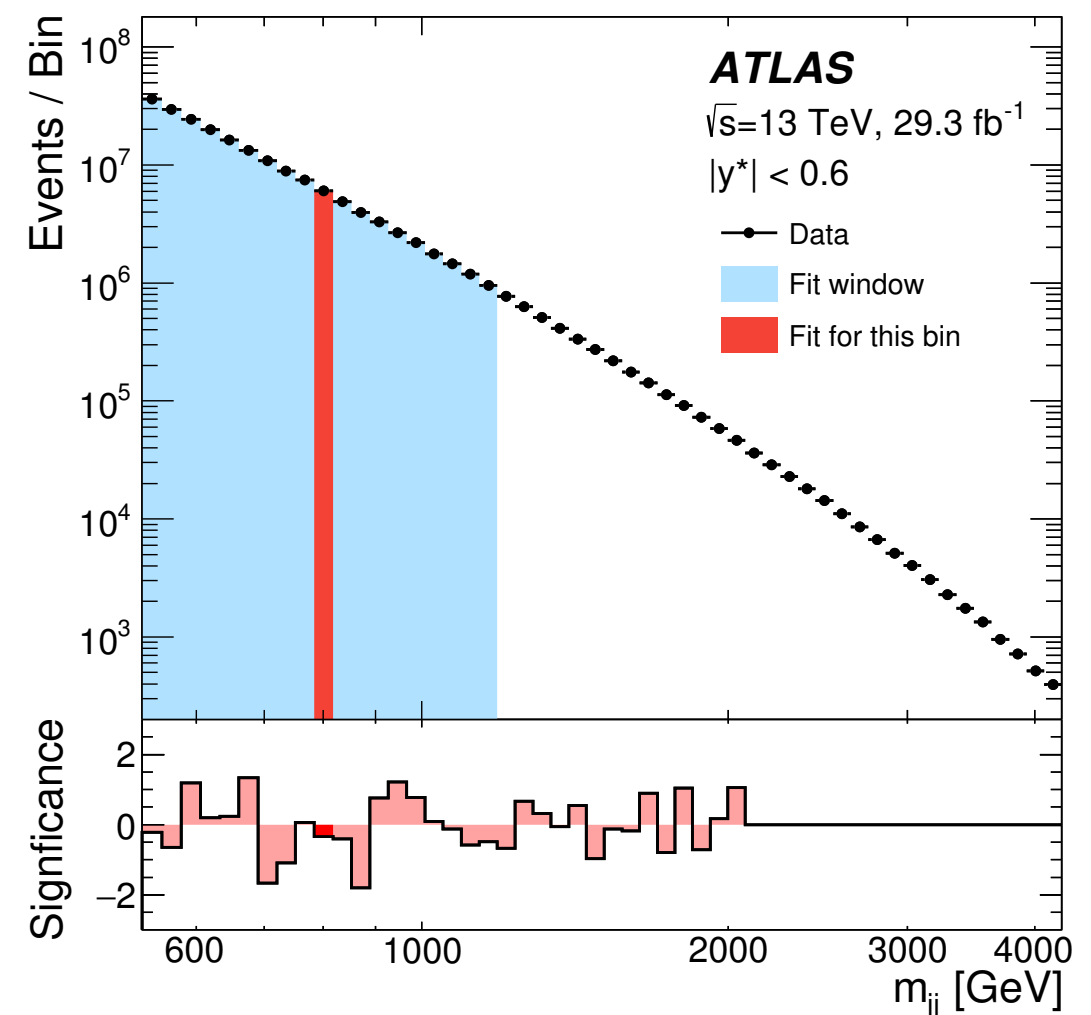


animation [here](#)

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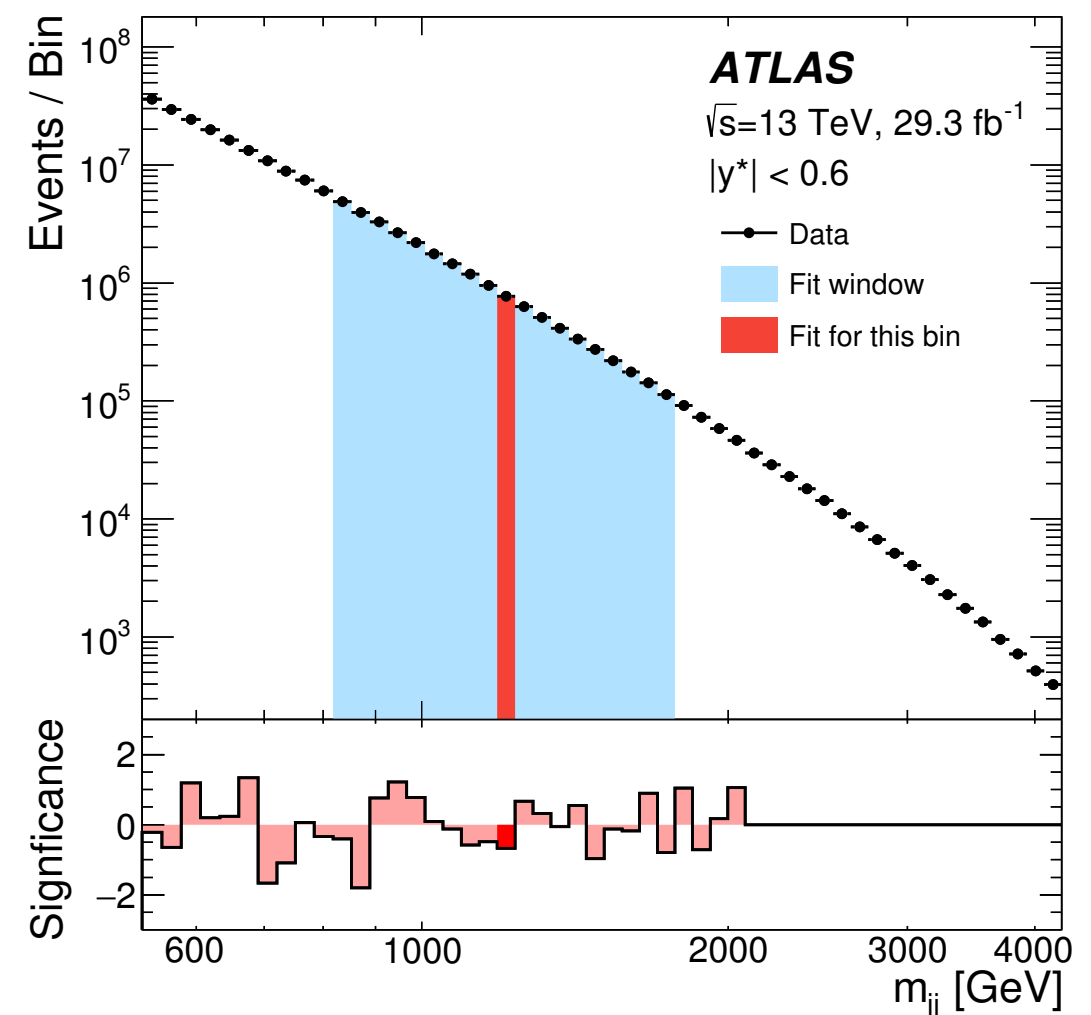


animation [here](#)

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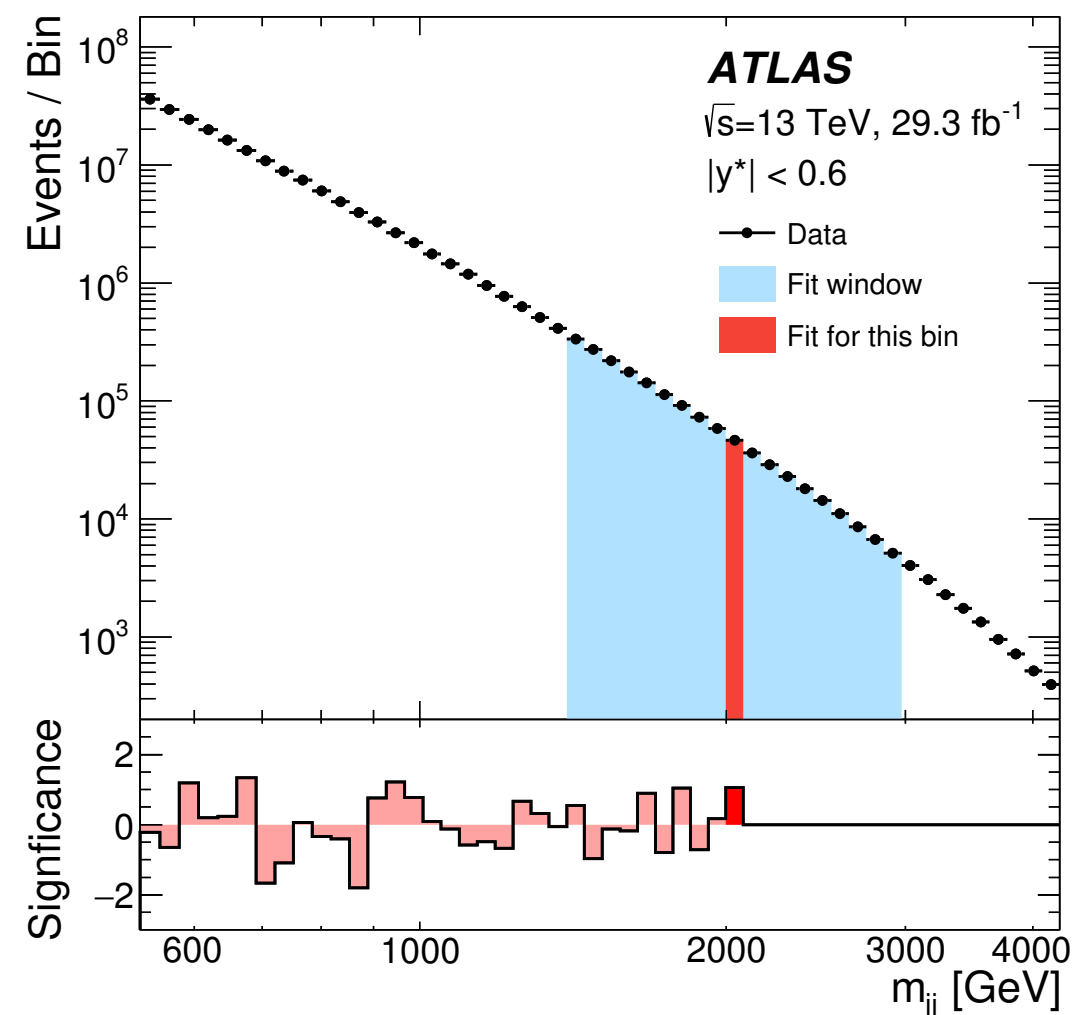


animation [here](#)

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- Solution, also used by high-mass dijet 37 fb⁻¹ result: fit sub-ranges
 - $|y^*| < 0.3$: 27 bins, $|y^*| < 0.6$: 19

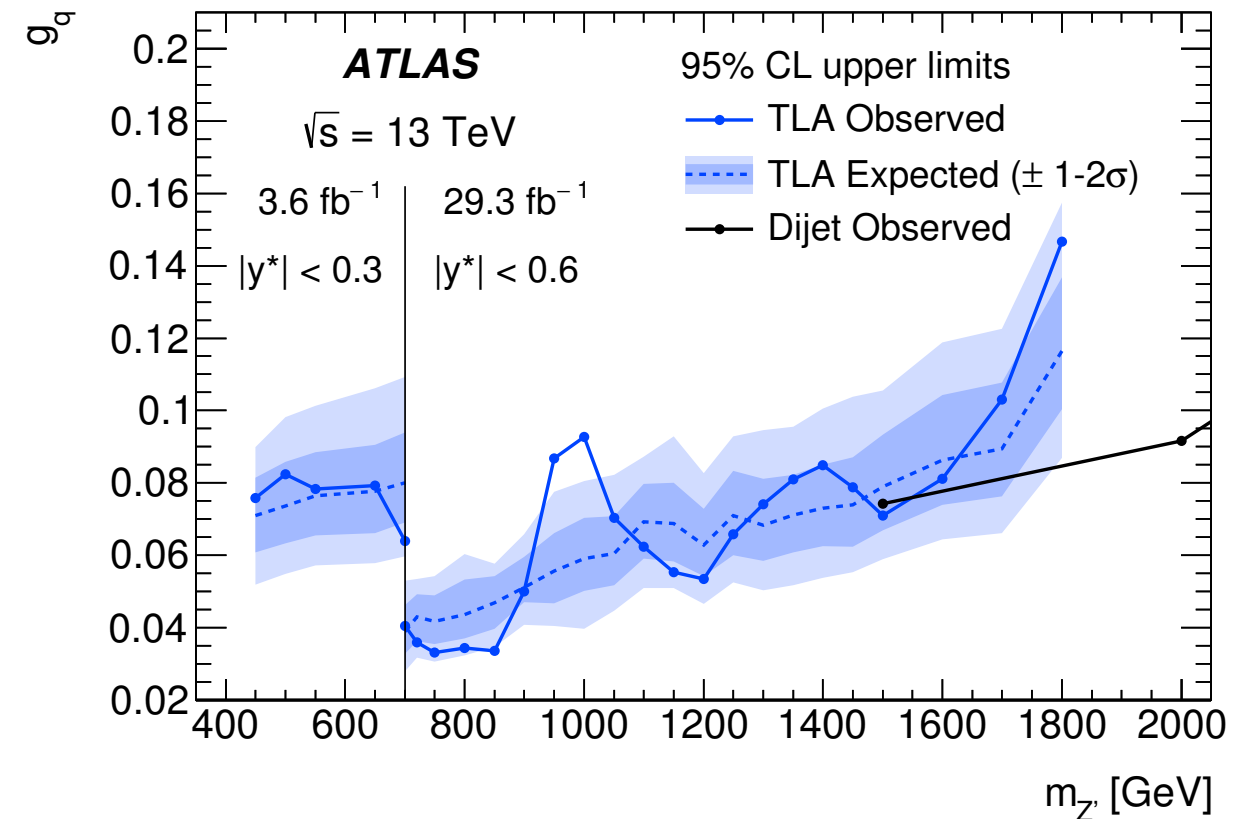
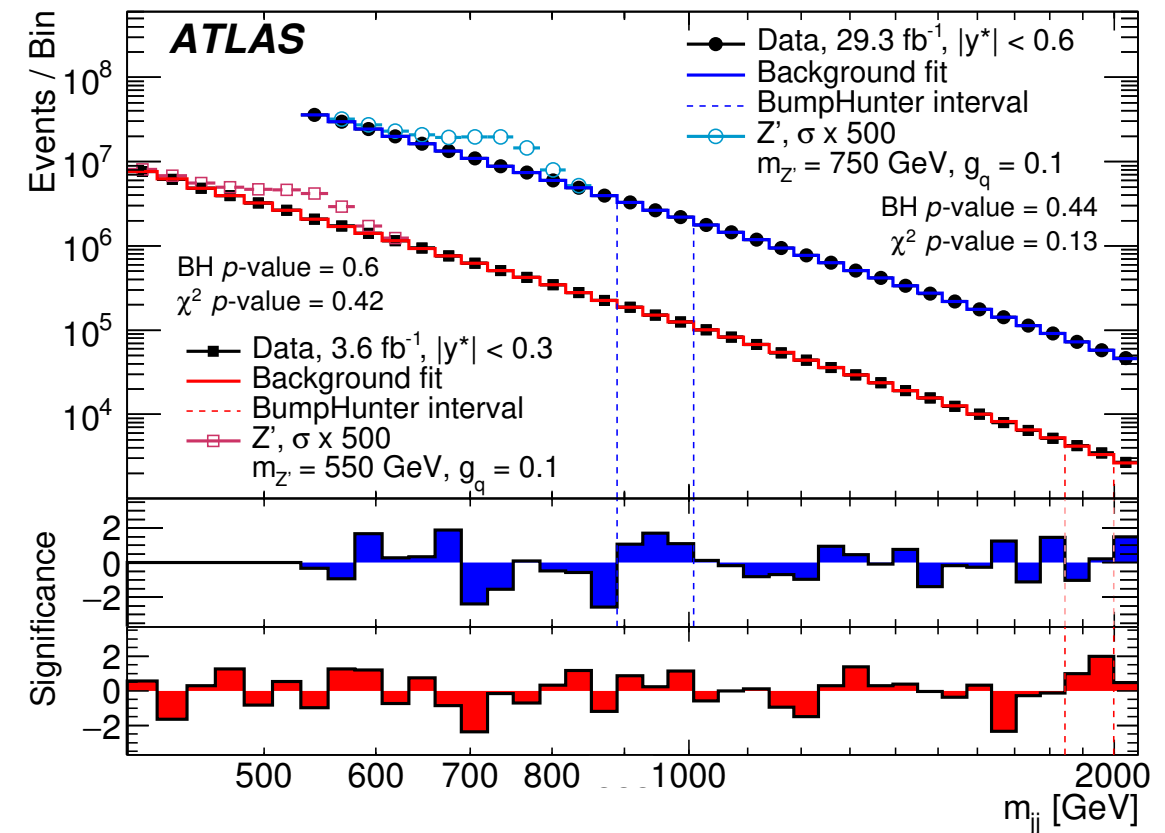
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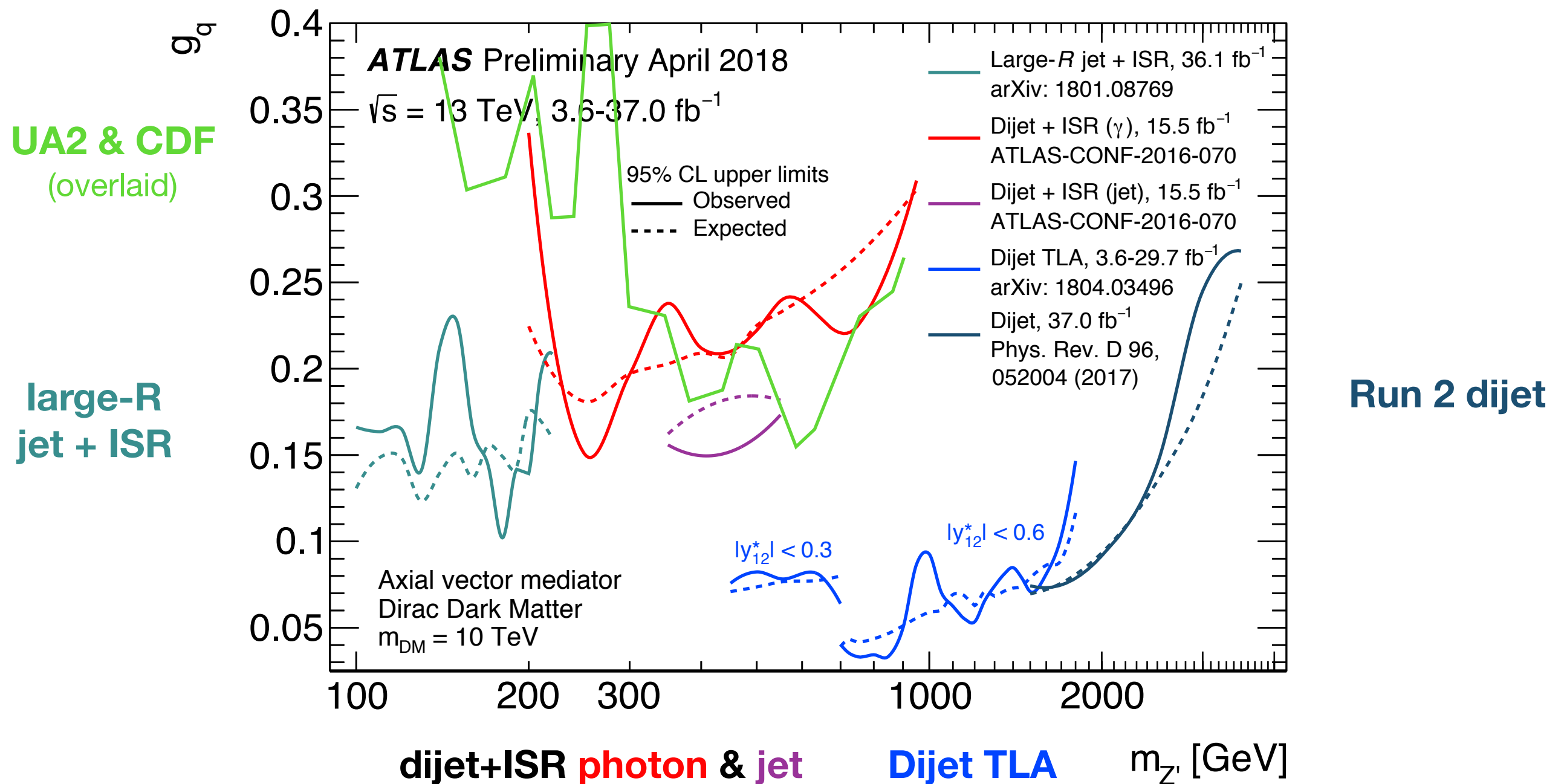
animation [here](#)

Results

- “BumpHunter” with background-only fit: no significant excesses found
- Signal + Background fit: set limits (areas of flexibility give observed - expected differences)
- Similar sensitivity to conventional dijet resonance search at 1.5 TeV
- Can go much lower in $m_{Z'}$
 - 450-700 GeV using dedicated signal region with L1_J75 for some of 2016



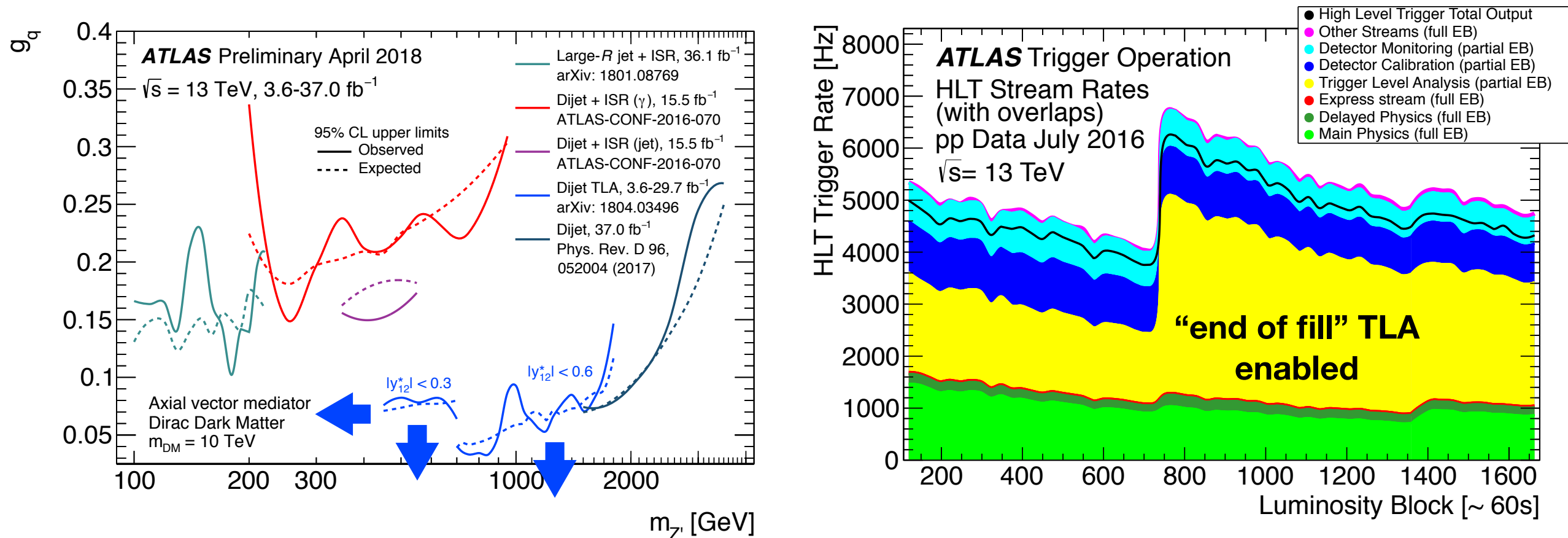
Limits, March 2018



Trigger-level analysis greatly improves sensitivity

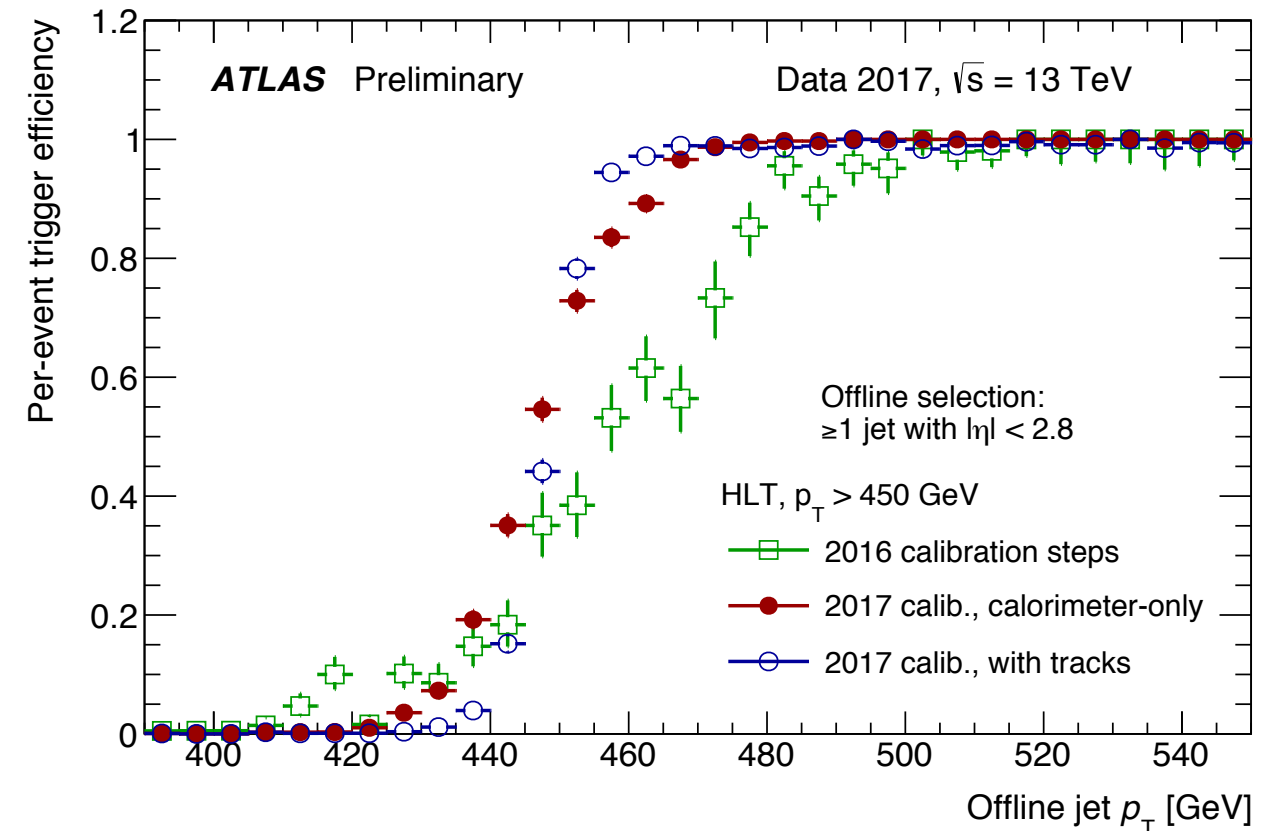
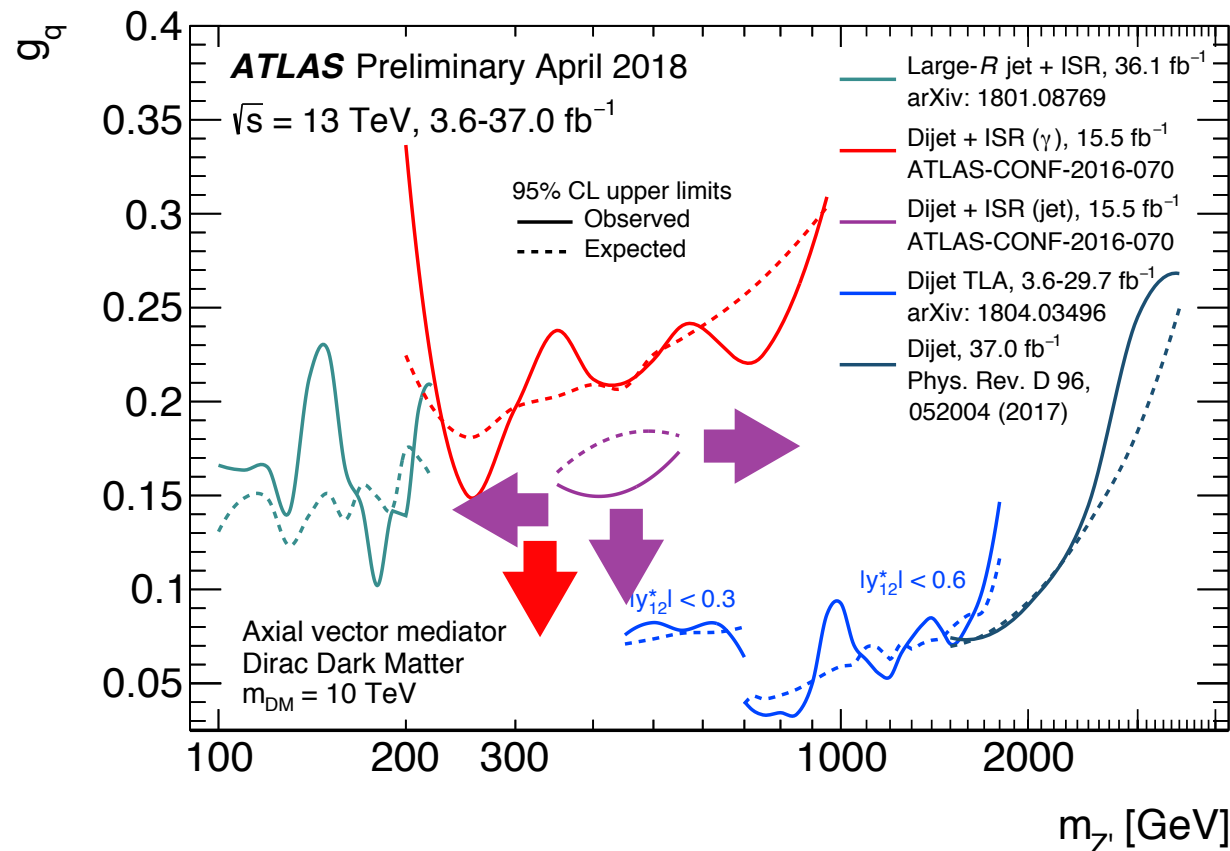
New results mean that we surpass pre-LHC constraints everywhere

Prospects, TLA



- 2017/8: improve calibration of trigger jets, take advantage of unused L1 rate towards end of fill to run new triggers allowing lower masses to be probed (J50 vs J75/J100)
- Run 3: improve reconstruction of L1 objects with new hardware => can probe lower mass for given rate
- Run 3: FTK -> full tracking at HLT -> pileup rejection possible -> can go well below 85 GeV

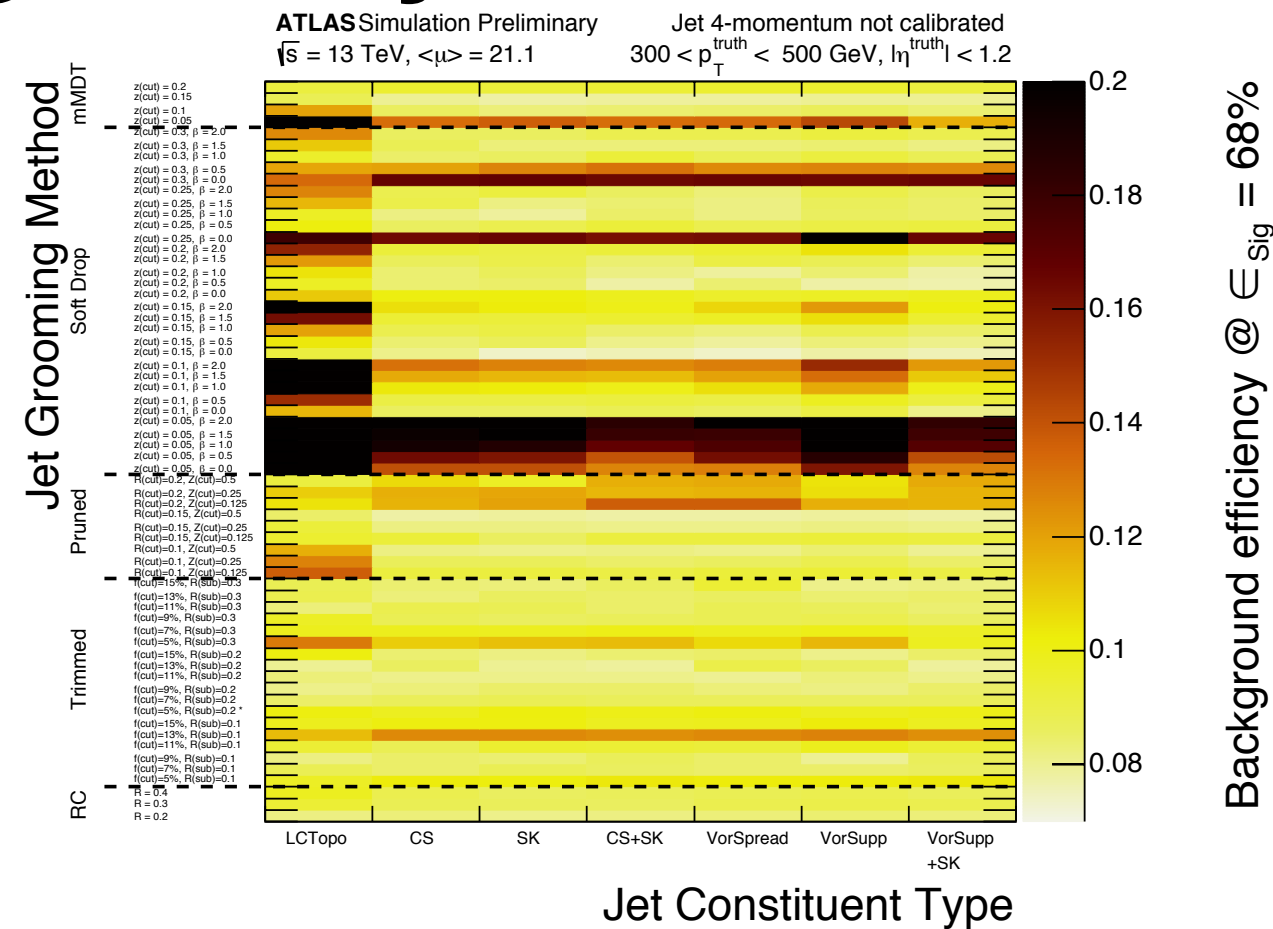
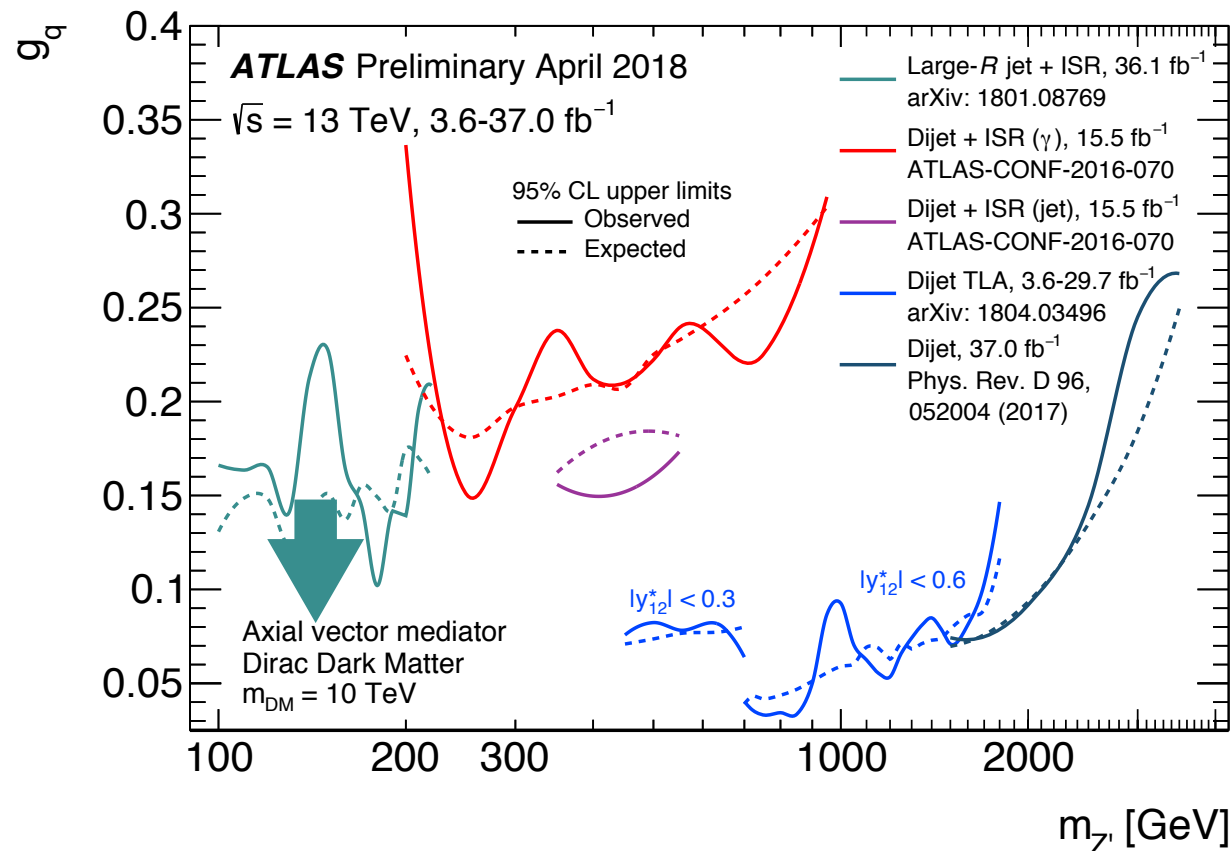
Prospects, resolved dijet + ISR



plateau of j380 in 2016 ~ j420 in 2017

- g_q limit scales as $\text{data}^{1/4} \Rightarrow 15.5 \text{ to } 120 \text{ fb}^{-1} = \text{factor } 1.7$
- Higher instantaneous luminosity \rightarrow higher trigger thresholds, mitigated by improved jet trigger performance
- Combinatorics in jet channel can improve mass reach and sensitivity
- Potential for TLA technique in run 3 with FTK

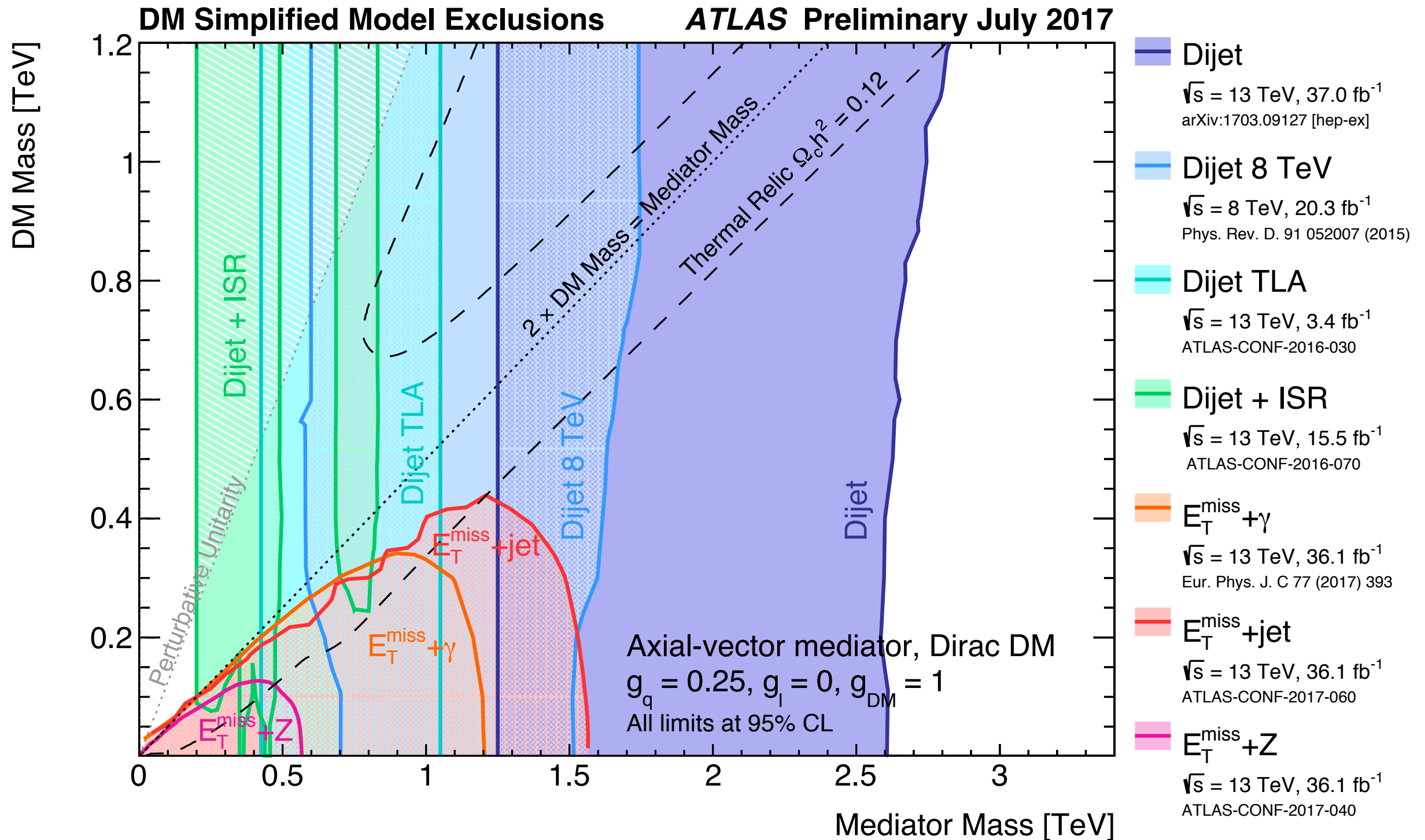
Prospects, merged dijet + ISR



- g_q limit scales as $\text{data}^{1/4} \Rightarrow 37 \text{ to } 120 \text{ fb}^{-1} = \text{factor } 1.3$
- New trigger strategies for large-R, including substructure information in the trigger (2017 has mass, run 3 will have more) \rightarrow much more data
- Optimised grooming methods [ATL-PHYS-PUB-2017-020](#) \rightarrow better S/B
- Also improvements in jet substructure resolution thanks to track information in jet reconstruction inputs [ATL-PHYS-PUB-2017-015](#)

More details on substructure in [Jason Veatch's WG4 talk yesterday](#)

Complementarity between DM searches

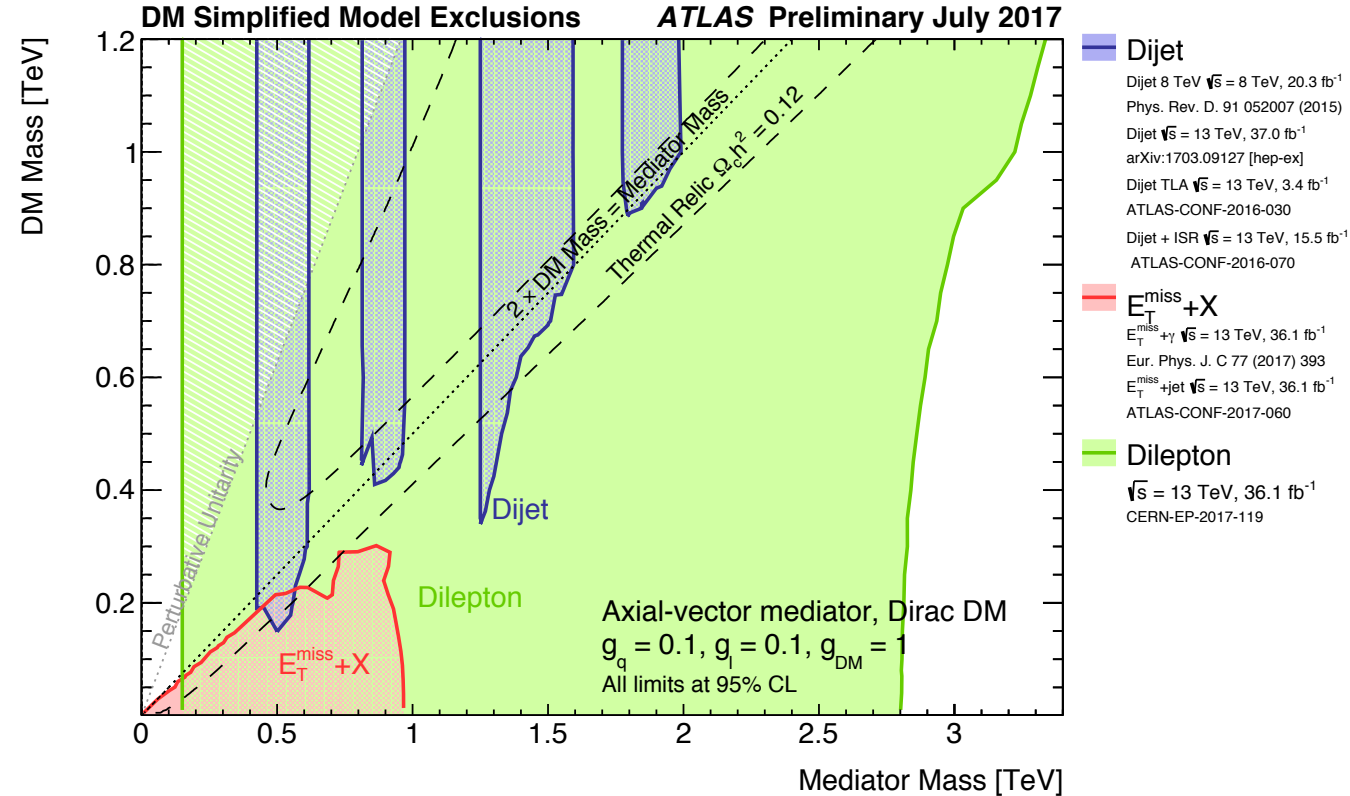
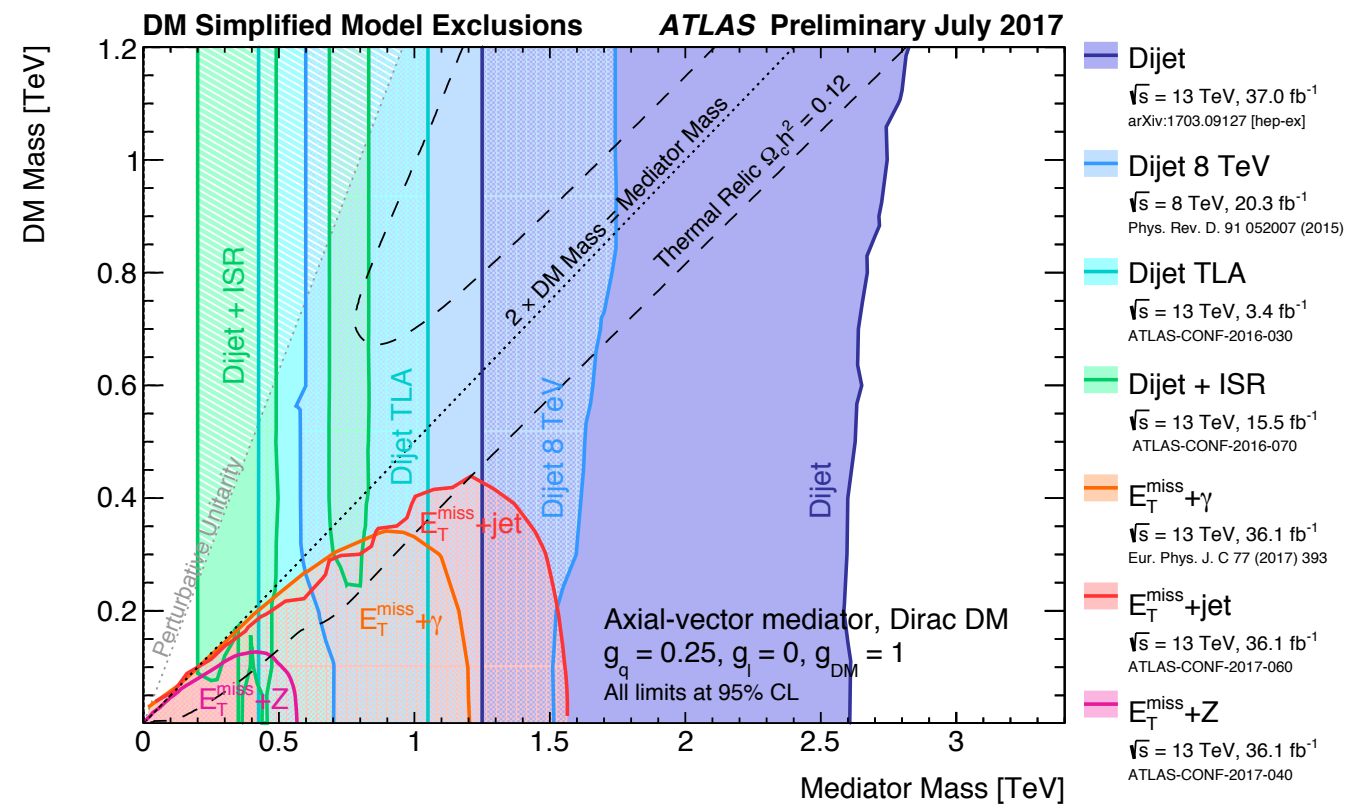


**mono-X and resonance searches
complement each other**

Caveats:

- plot is ~ 1 year old, doesn't include latest TLA, large-R+ISR or mono-X results

Complementarity between DM searches

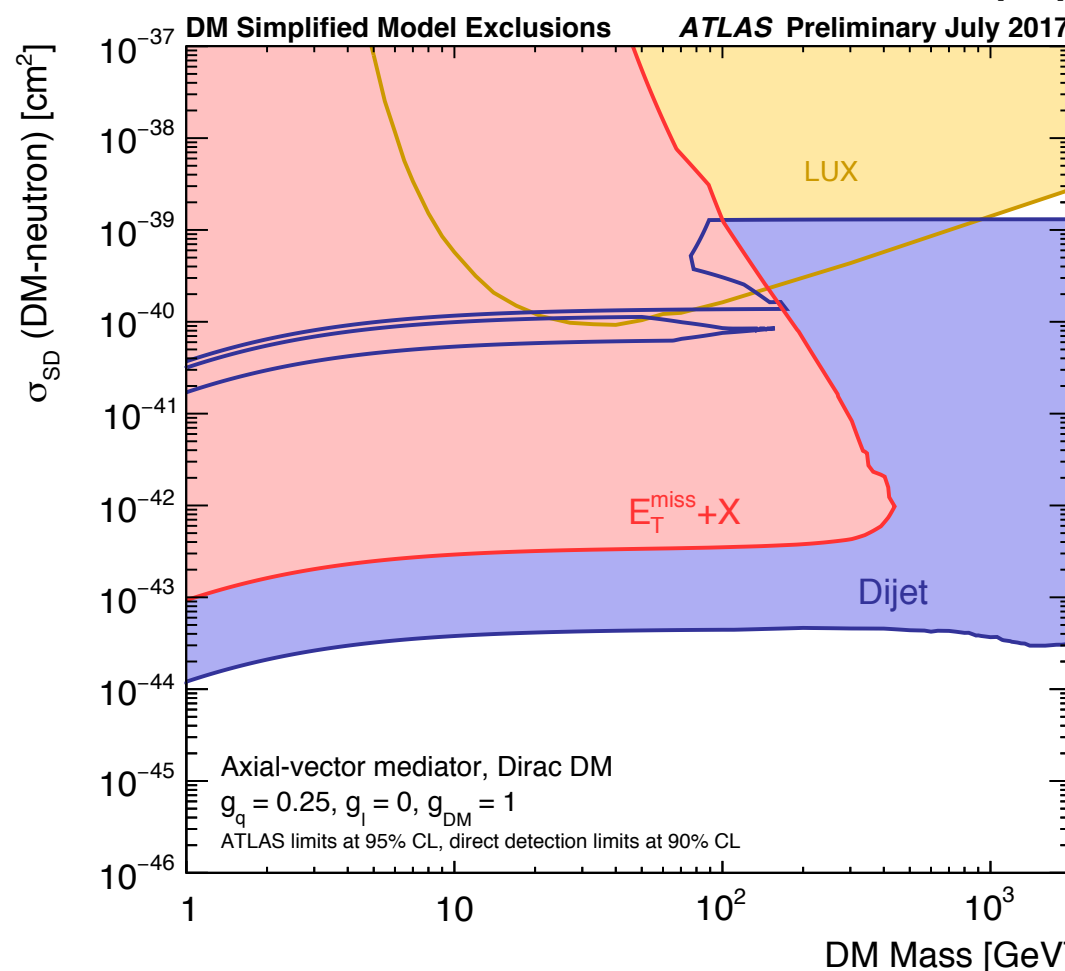
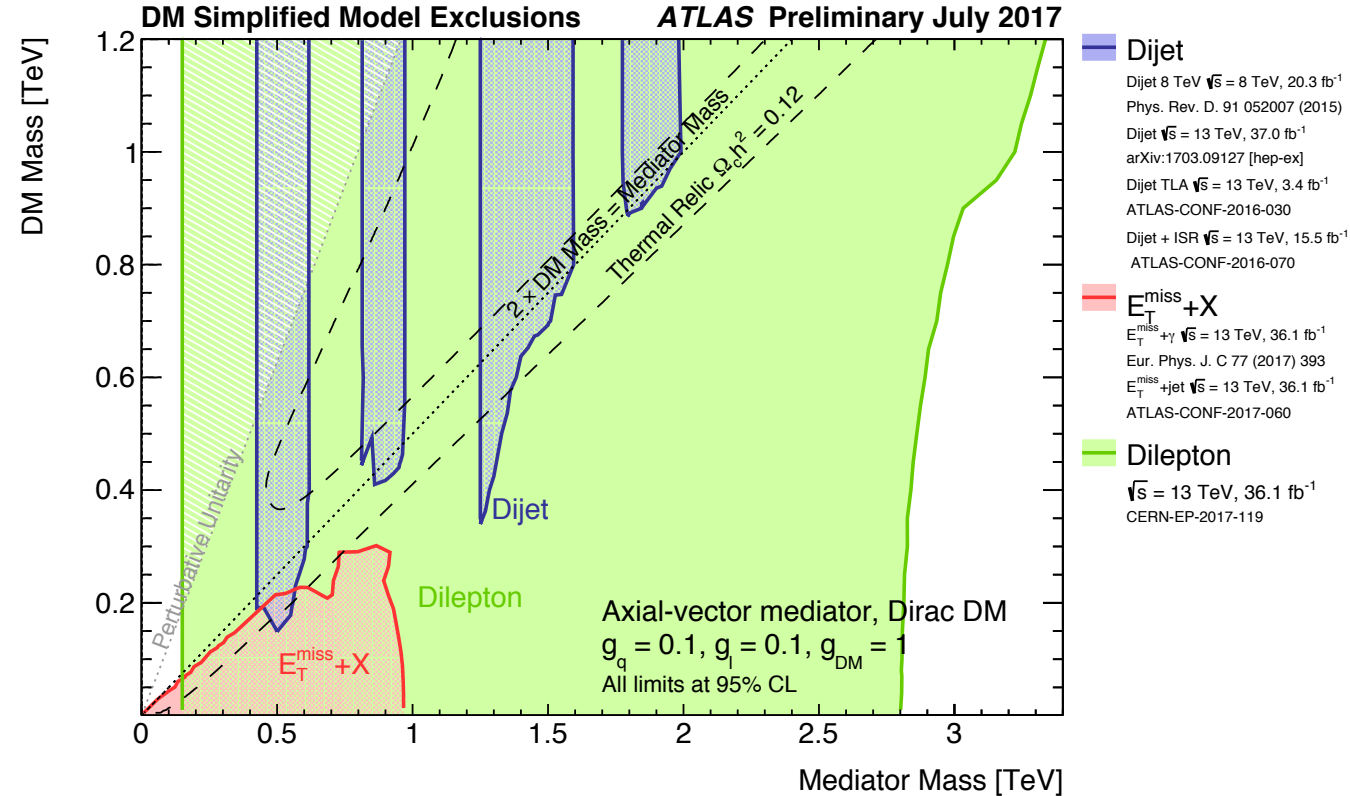
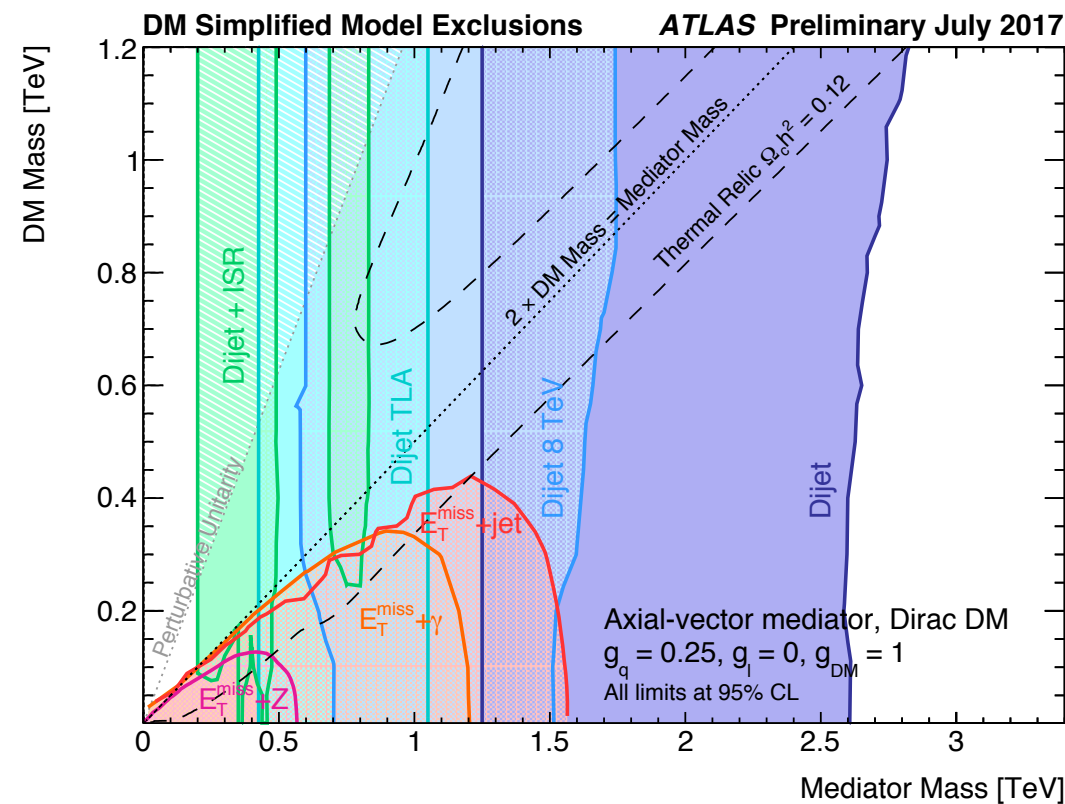


Caveats:

- plots are ~ 1 year old, don't include latest TLA, large-R+ISR or mono-X results
- very model-dependent (eg non-zero lepton coupling causes large changes)

other channels (eg dilepton resonance) cover other model scenarios

Complementarity between DM searches



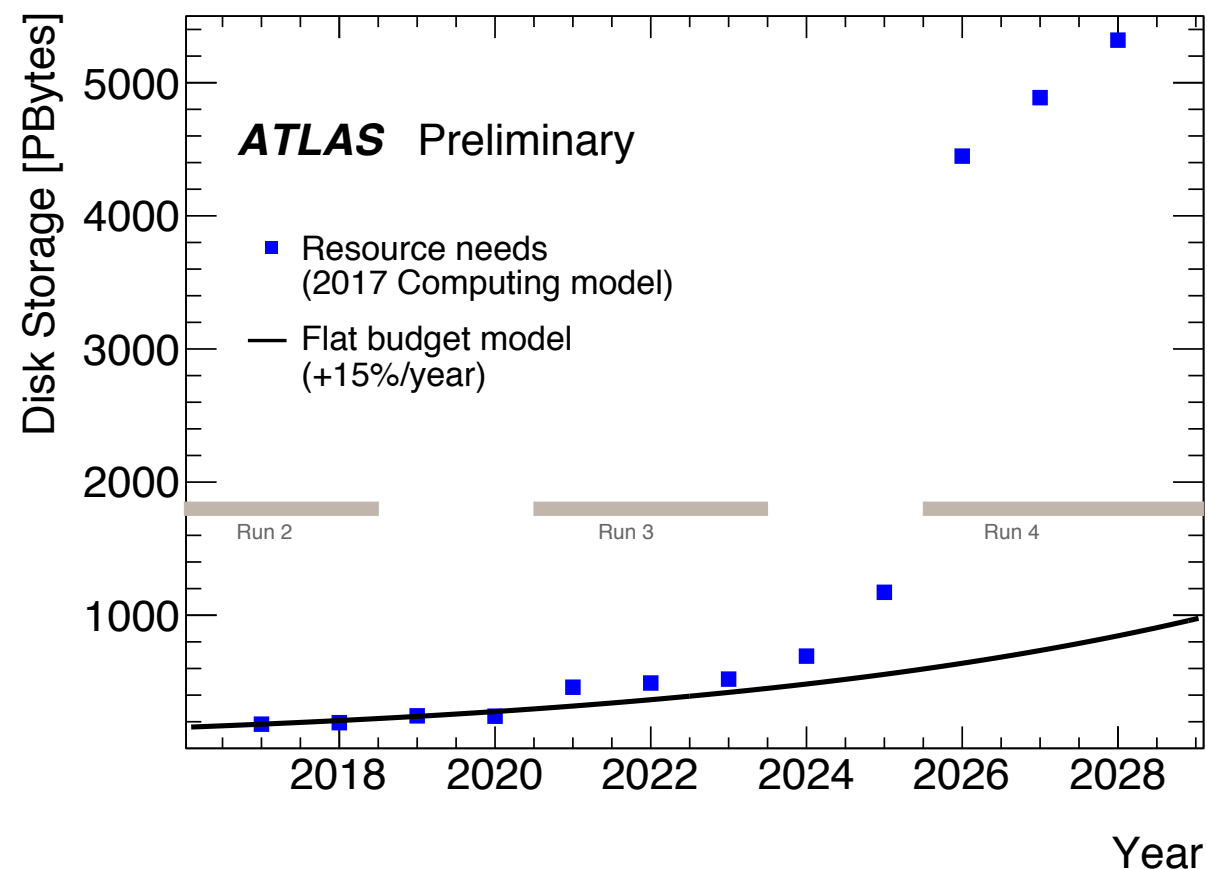
Caveats:

- plots are ~ 1 year old, don't include latest TLA, large-R+ISR or mono-X results
- very model-dependent (eg non-zero lepton coupling causes large changes)
- DD limits 90% CL, collider 95%

**also complementarity
with direct detection**

Conclusions

- Broad set of approaches to searching for Dark Matter with ATLAS
- Various new techniques being exploited to go lower in mass
 - Initial state radiation to evade trigger limitations
 - Substructure to take this into the merged regime
 - Borrowing methods from LHCb and CMS to make the best use of jet trigger system and do a dijet analysis with partial events

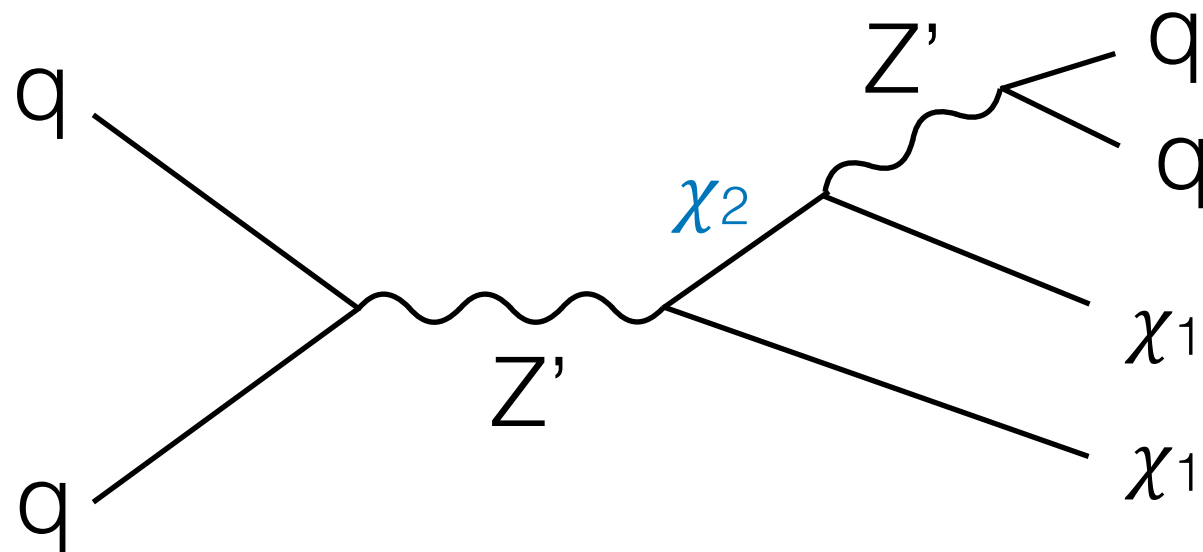


- New methods can all take advantage of LS2 trigger upgrades for sensitivity scaling much better than integrated luminosity alone
- Can also help with significant computing and storage pressures in the future

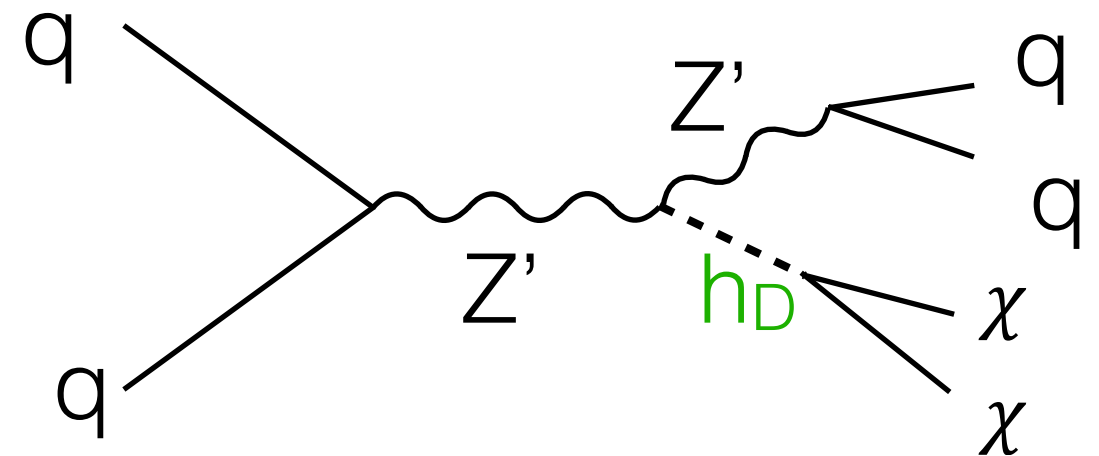
Backup

New: mono- Z'

ATLAS-CONF-2018-005, April 4th



Dark Fermion



Dark Higgs

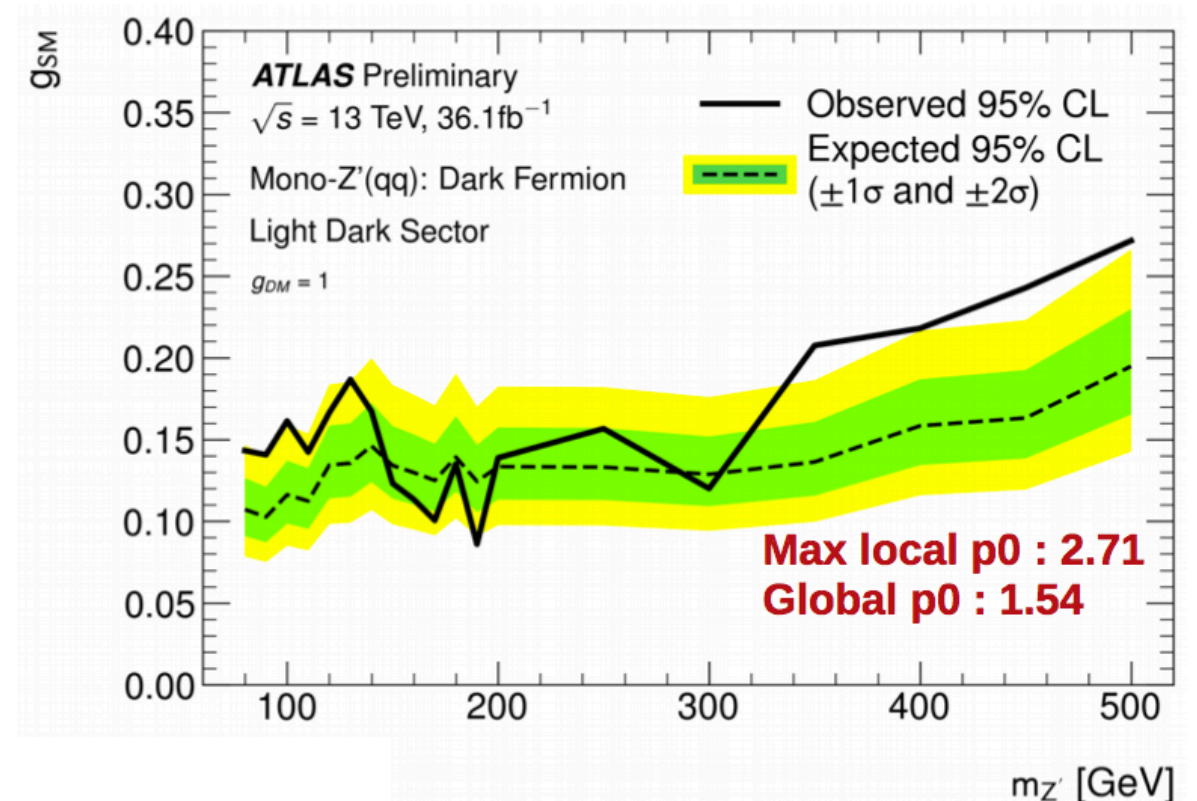
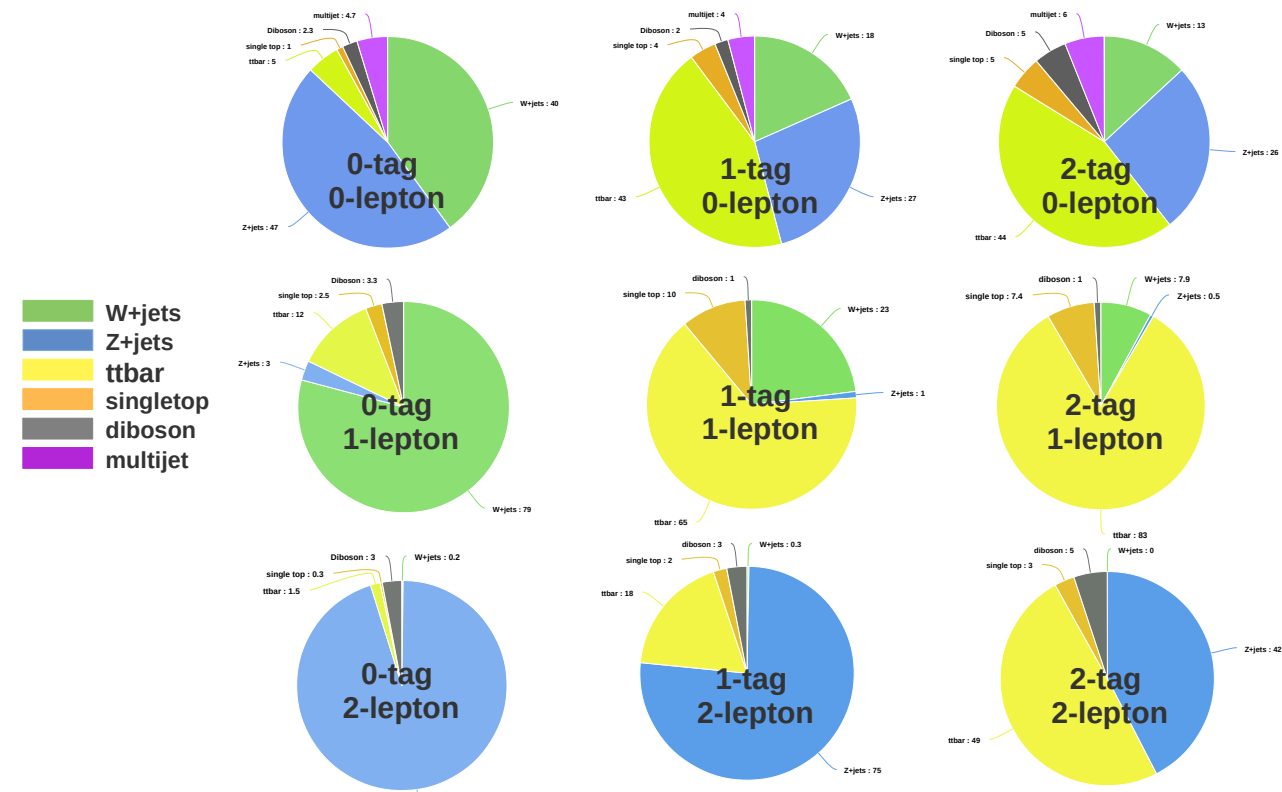
Dijet resonance + MET

Scenario	Dark-fermion model	Dark-Higgs model
	$m_{\chi_1} = 5 \text{ GeV}$	$m_{\chi} = 5 \text{ GeV}$
Light dark sector	$m_{\chi_2} = m_{\chi_1} + m_{Z'} + 25 \text{ GeV}$	$m_{h_D} = \begin{cases} m_{Z'} & , m_{Z'} < 125 \text{ GeV} \\ 125 \text{ GeV} & , m_{Z'} > 125 \text{ GeV} \end{cases}$
	$m_{\chi_1} = m_{Z'}/2$	$m_{\chi} = 5 \text{ GeV}$
Heavy dark sector	$m_{\chi_2} = 2m_{Z'}$	$m_{h_D} = \begin{cases} 125 \text{ GeV} & , m_{Z'} < 125 \text{ GeV} \\ m_{Z'} & , m_{Z'} > 125 \text{ GeV} \end{cases}$

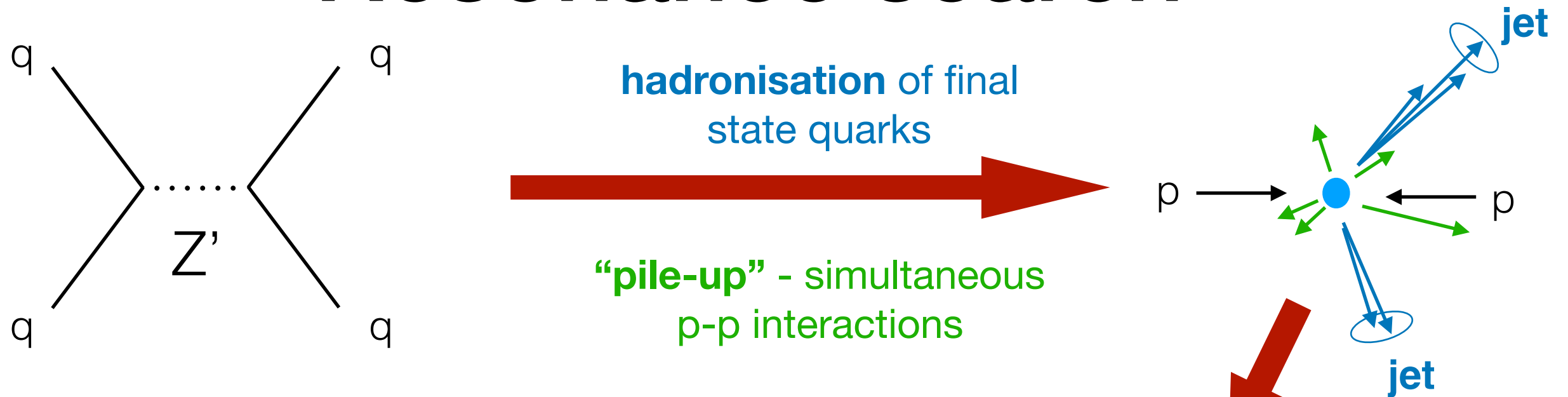
Quick overview: Mono-Z'

- E_T^{miss} trigger
- Merged and resolved jet resonance search
- Use of btagging to enhance sensitivity to $Z' \rightarrow b\bar{b}$
- Combined fit of MC normalisations in 1&2-lepton CRs and 0-lepton SRs
- Limits: heavy dark sector comparable to dijet searches, stronger with light dark sector
- Systematically limited \Rightarrow foresee improvement

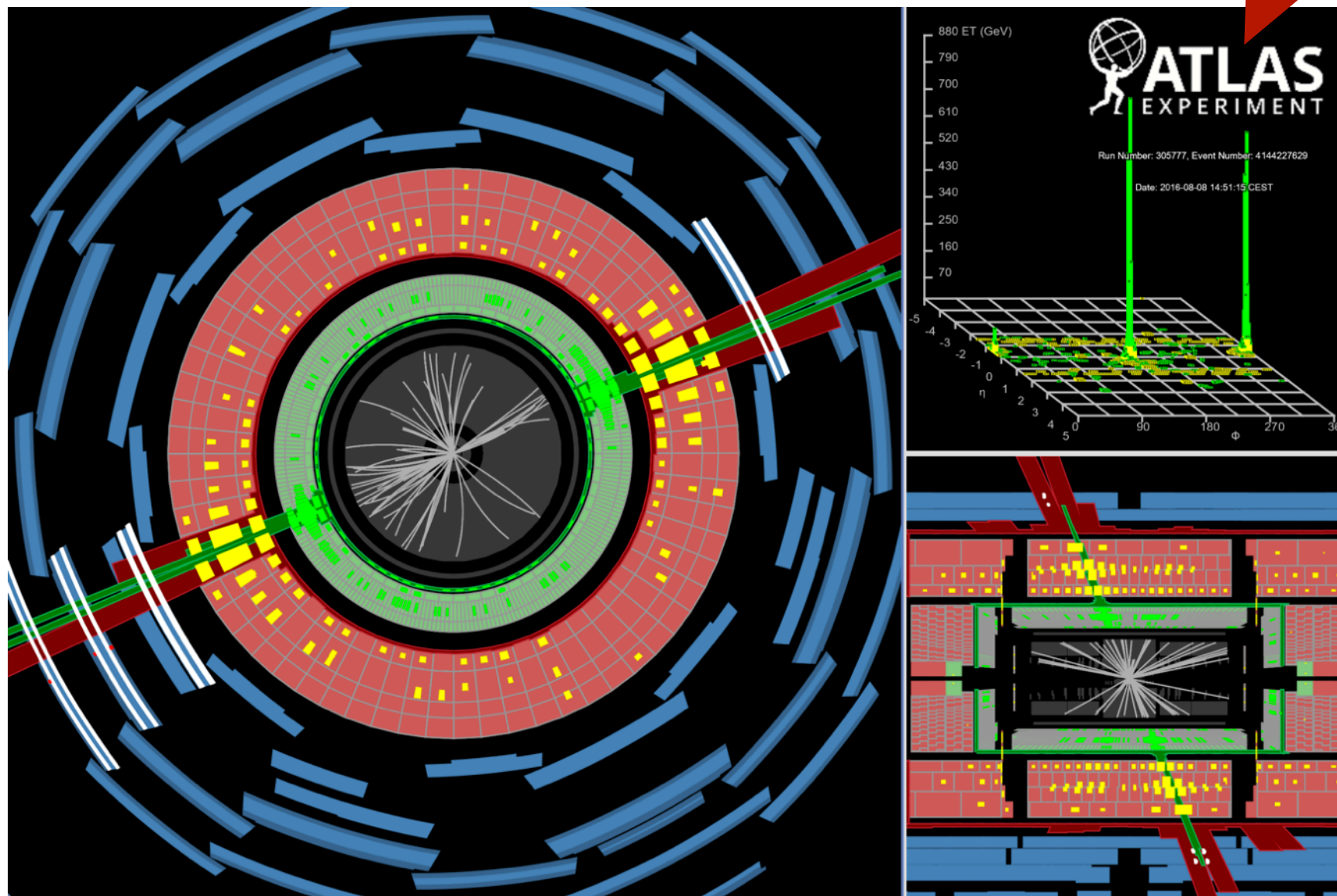
* Example background composition for $m_{Z'} = 90$ GeV, resolved topology.



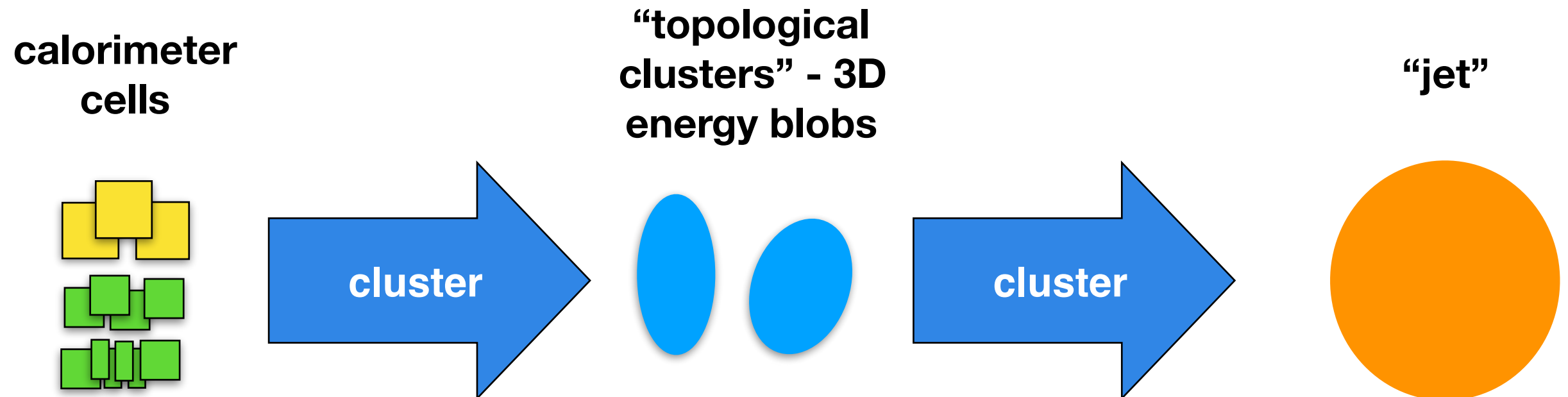
Resonance search



highest-mass
dijet event in
2016
 $p_T(j1, j2) = 3.79$
 $m_{jj} = 8.12 \text{ TeV}$



Jet reconstruction



- Seed from cells with $S/N > 4$
- Grow with cells $S/N > 2$
- Split local maxima (EM calorimeter)

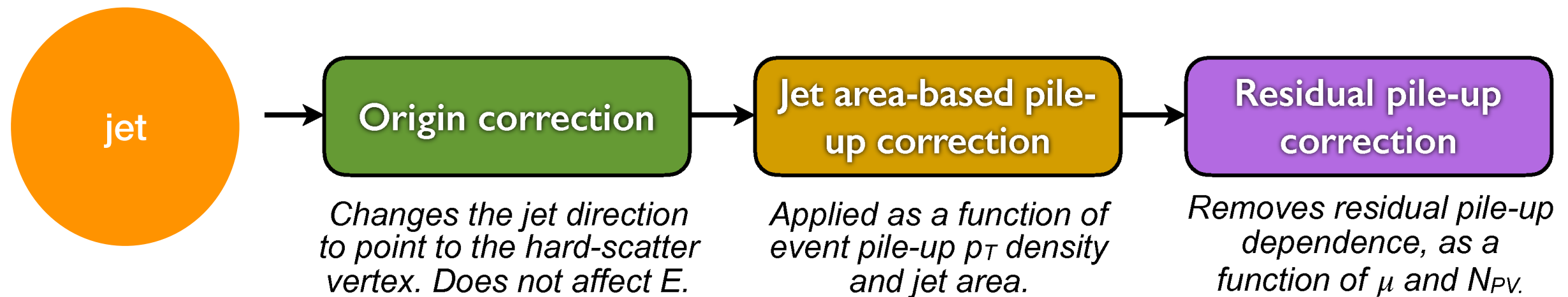
- Sequentially merge topoclusters
- Start from highest E_T
- Size controlled by ‘radius’ parameter, $\Delta R = \Delta\eta \oplus \Delta\phi = 0.4$
- End with a 2D object - \sim circular in η - ϕ (except when touch)

Jet calibration



- Built from raw energy recorded by calorimeter
- **sampling** calorimeters -> don't record all the energy
- Also have energy deposits from other p-p collisions in same event

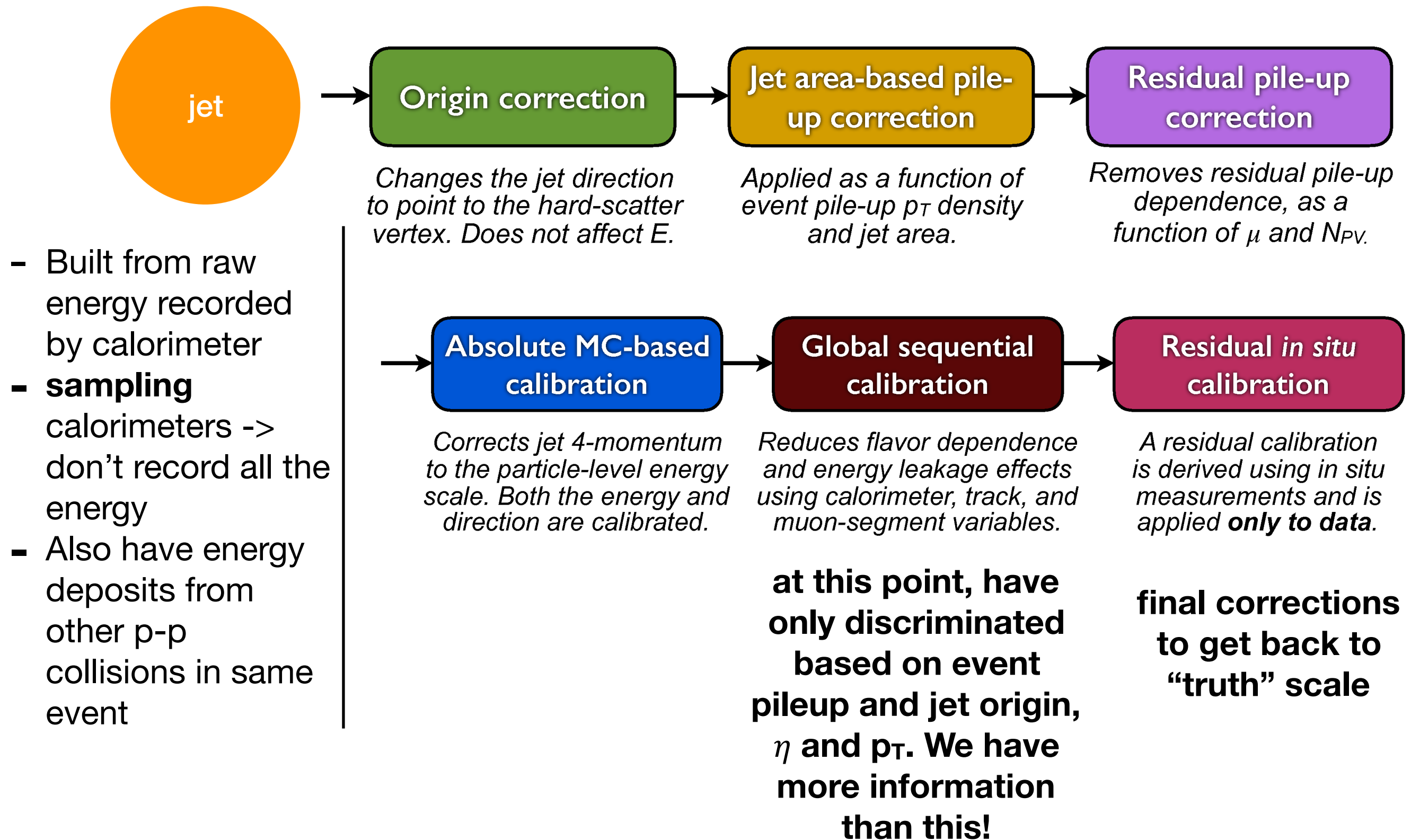
Jet calibration



- Built from raw energy recorded by calorimeter
- **sampling** calorimeters -> don't record all the energy
- Also have energy deposits from other p-p collisions in same event

look at average p_T density of event in the calorimeter, subtract this approximated pileup contribution

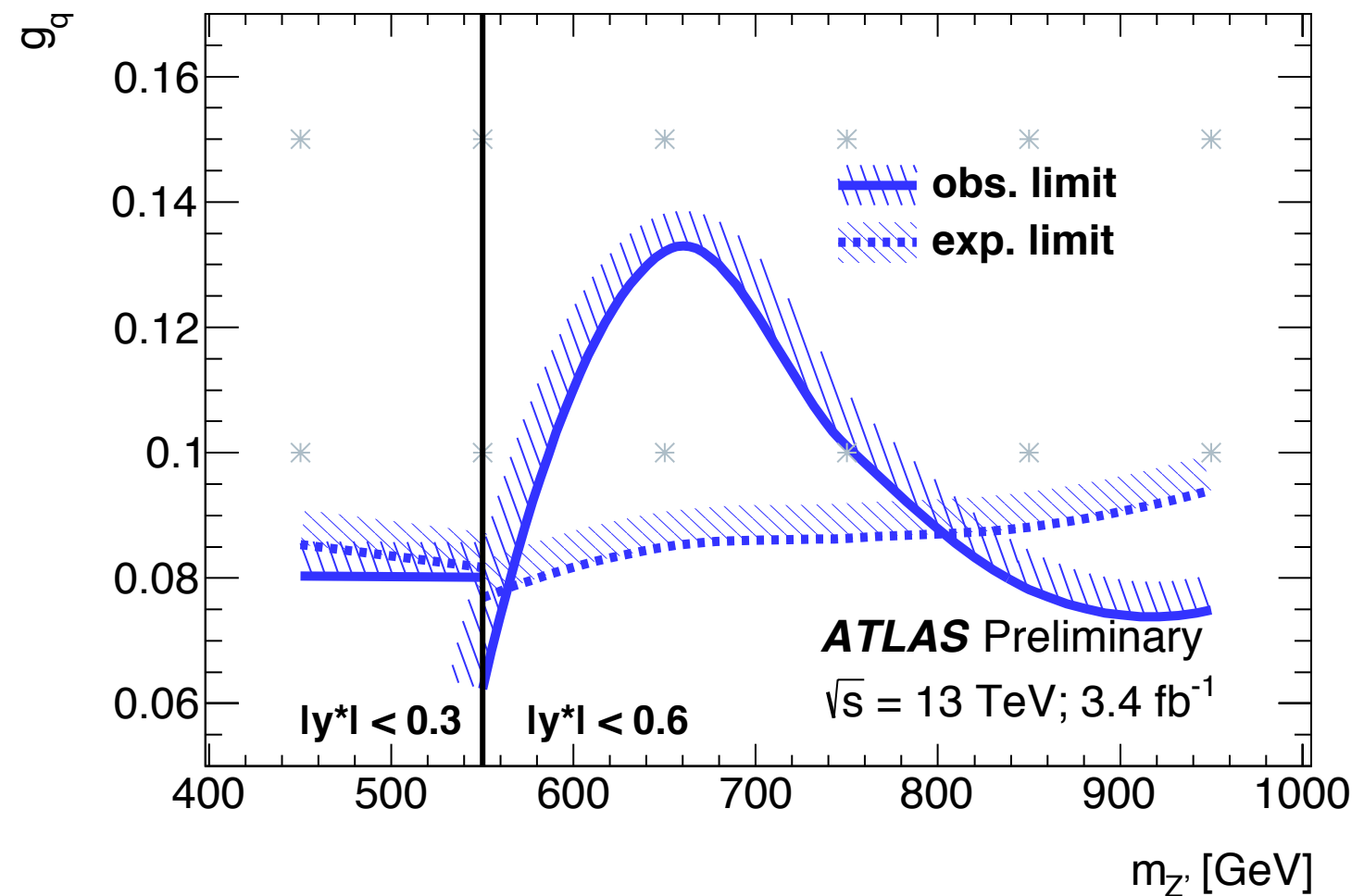
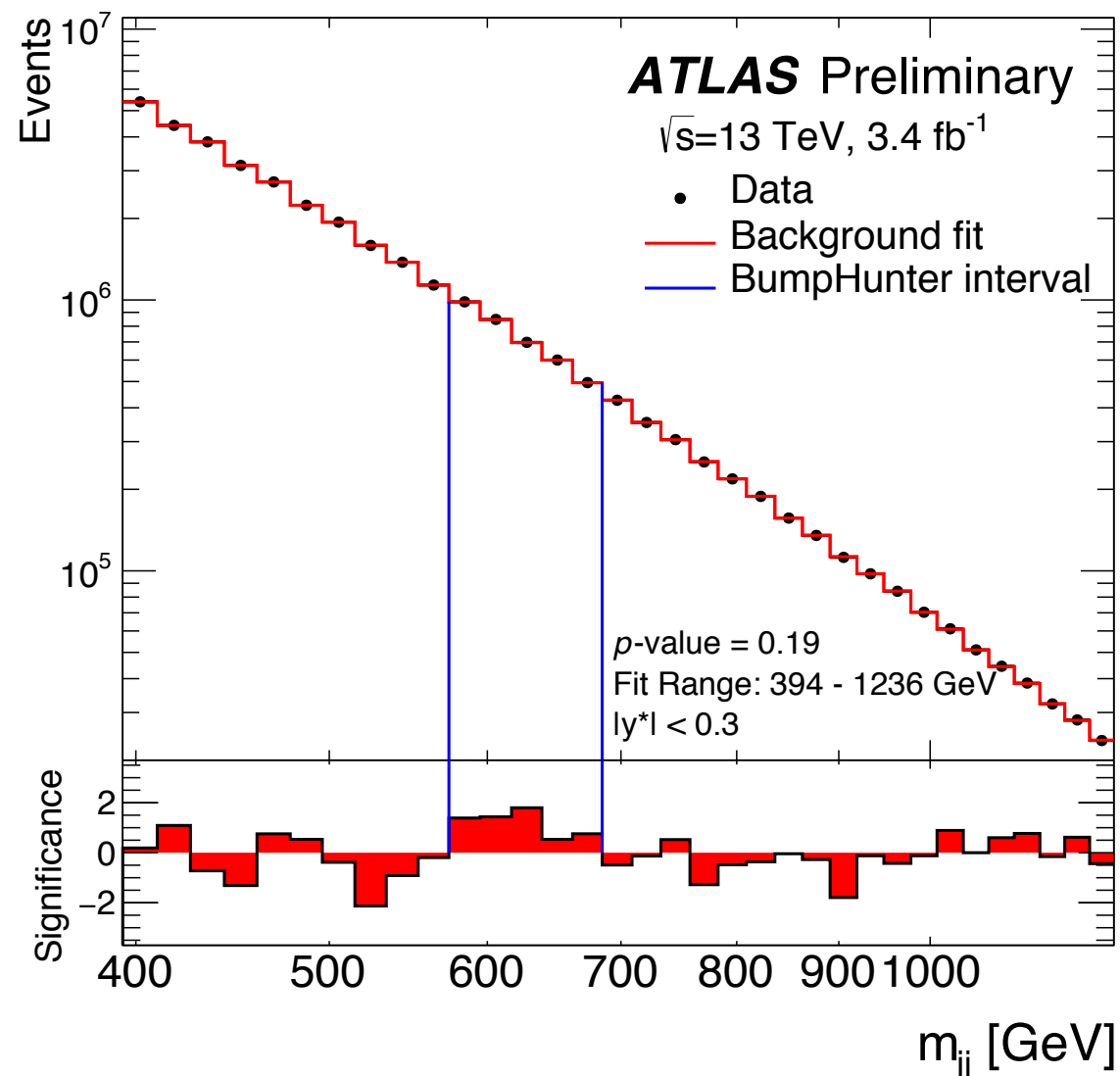
Jet calibration



TLA fitting

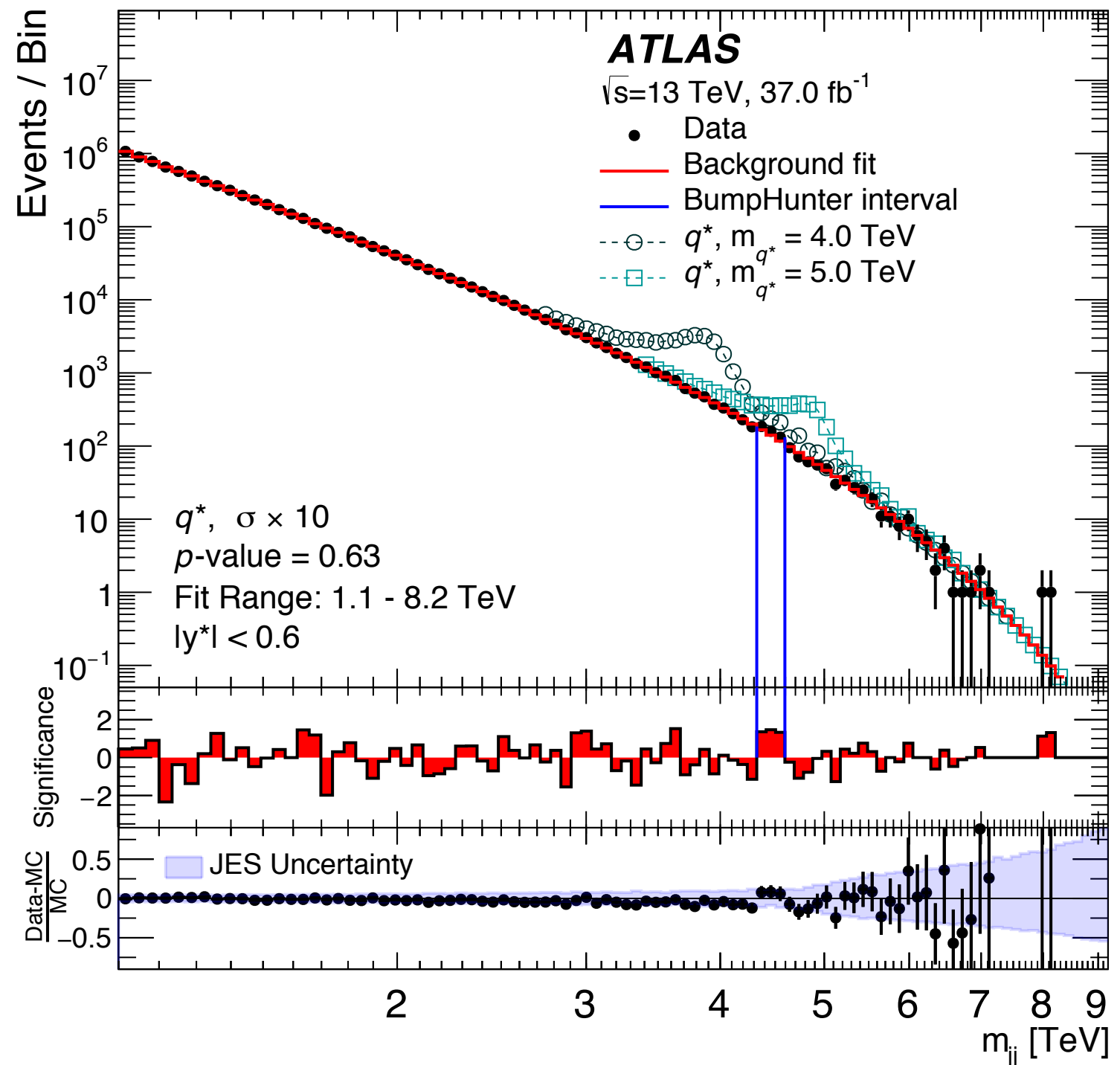
ATLAS-CONF-2016-030

- Very large number of events -> very little scope for QCD to deviate from functional form
- In 2015, could not fit whole m_{jj} range, hence truncated fit at 1250 GeV



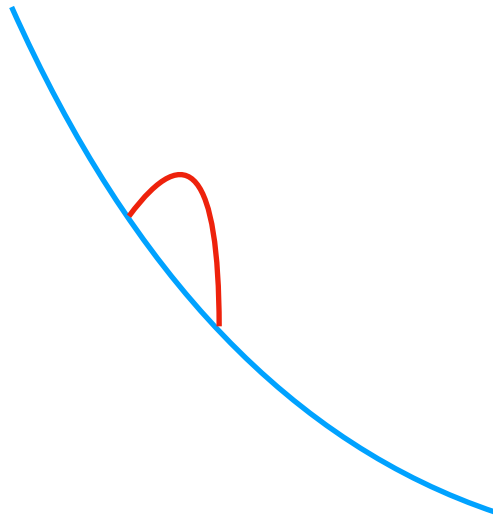
BumpHunter - high-mass dijet

- “BumpHunter” - scans all widths from 1 to $N_{\text{bins}}/2$, finds maximally discrepant interval
- $p\text{-value} < 0.05 \Rightarrow$ there is something there with 95% confidence
- $p\text{-value} > 0.05 \Rightarrow$ there is not something there

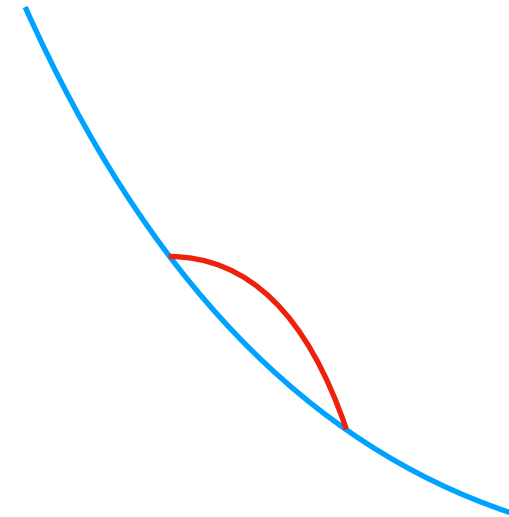


Limits on the limits: m_{jj} resolution

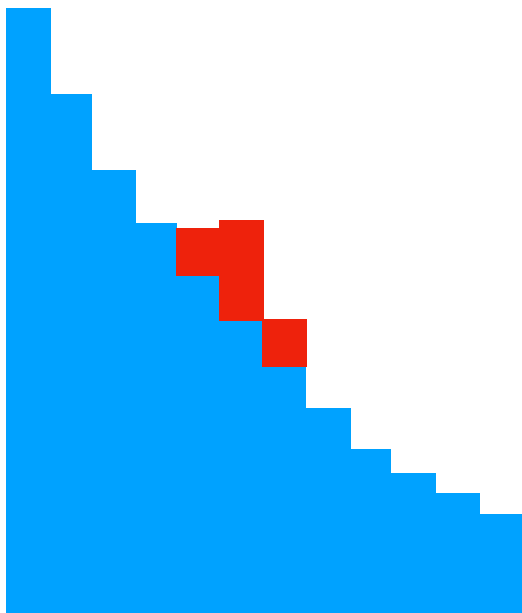
Good resolution



Bad resolution

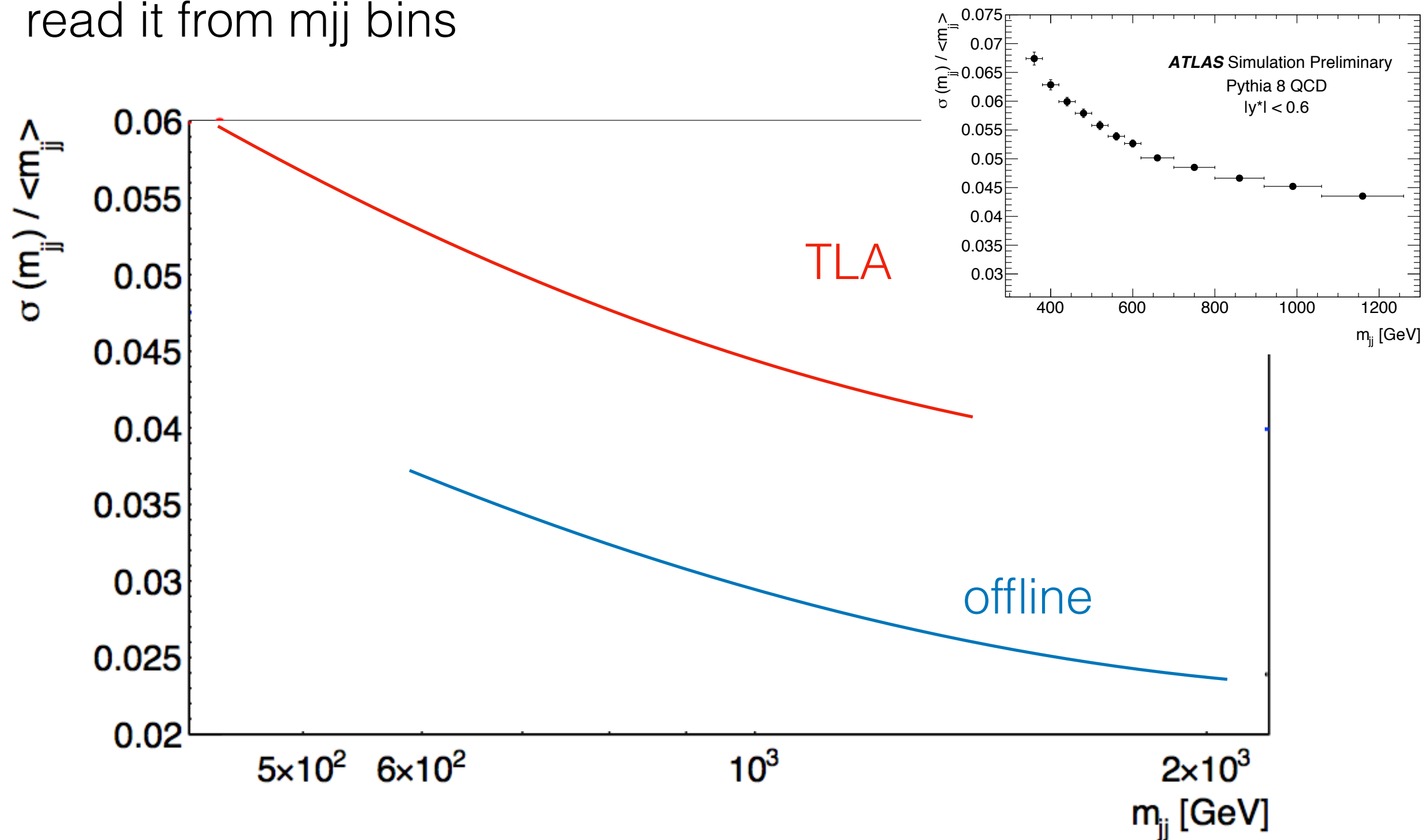


Bad resolution: signal smears out, covers wider m_{jj} range, trying to extract same number of signal events from more background events



m_{jj} resolution

Cartoon because offline plot is internal... but you can read it from m_{jj} bins

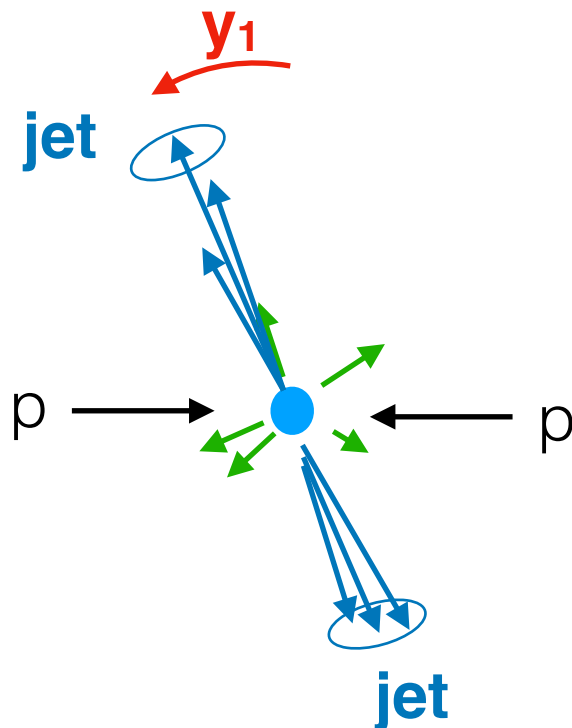


Lower still: exploiting the Kinematics

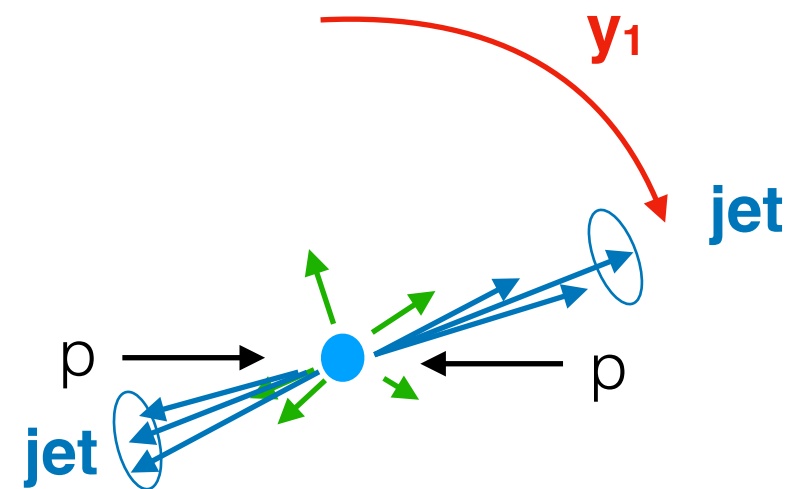
The dijet searches use $|y^*| < 0.6$

$$y^* = \frac{1}{2} (y_1 - y_2)$$

Imagine a centrally produced Z' :
i.e. quarks back to back, $y_1 = -y_2$, $y^* = y_1$



small Δy , large p_T



large Δy , small p_T

TLA: Imposing $|y^*| < 0.3 \Rightarrow$ higher $\langle p_T \rangle$ from given Z' mass \Rightarrow sensitive to lower Z' mass for given p_T (**394 vs 443**)

(signal and background both lose a factor of $\sim 2-3$)

Trigger evolution over time

1. LHC performance increases
2. Decide rate allocation
3. Adjust jet p_T threshold to fit
4. Evaluate performance of this trigger to determine analysis selections

year	L / $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	jet p_T threshold	single jet trigger rate	offline turnon
2015	0.5	260	18	400
2016	1.2	380	38	420
2017	1.7	420	33	435

Jet trigger performance

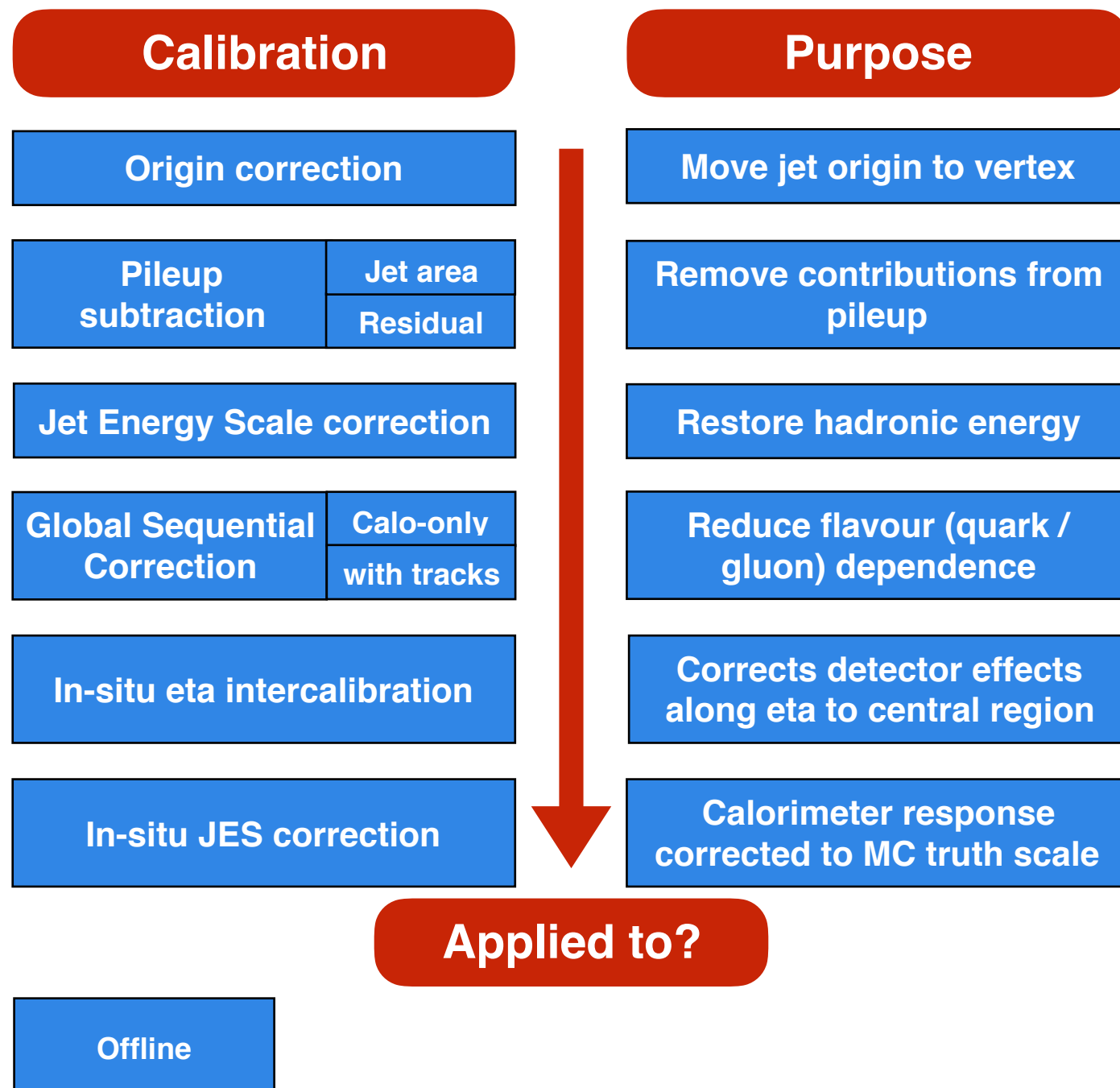
Before: offline - truth resolutions for width of m_{jj} peak

For triggers: trigger - offline resolution, i.e. how good are we at selecting the events we want to analyse?

This is set by how similar we can make trigger jets to offline jets, given:

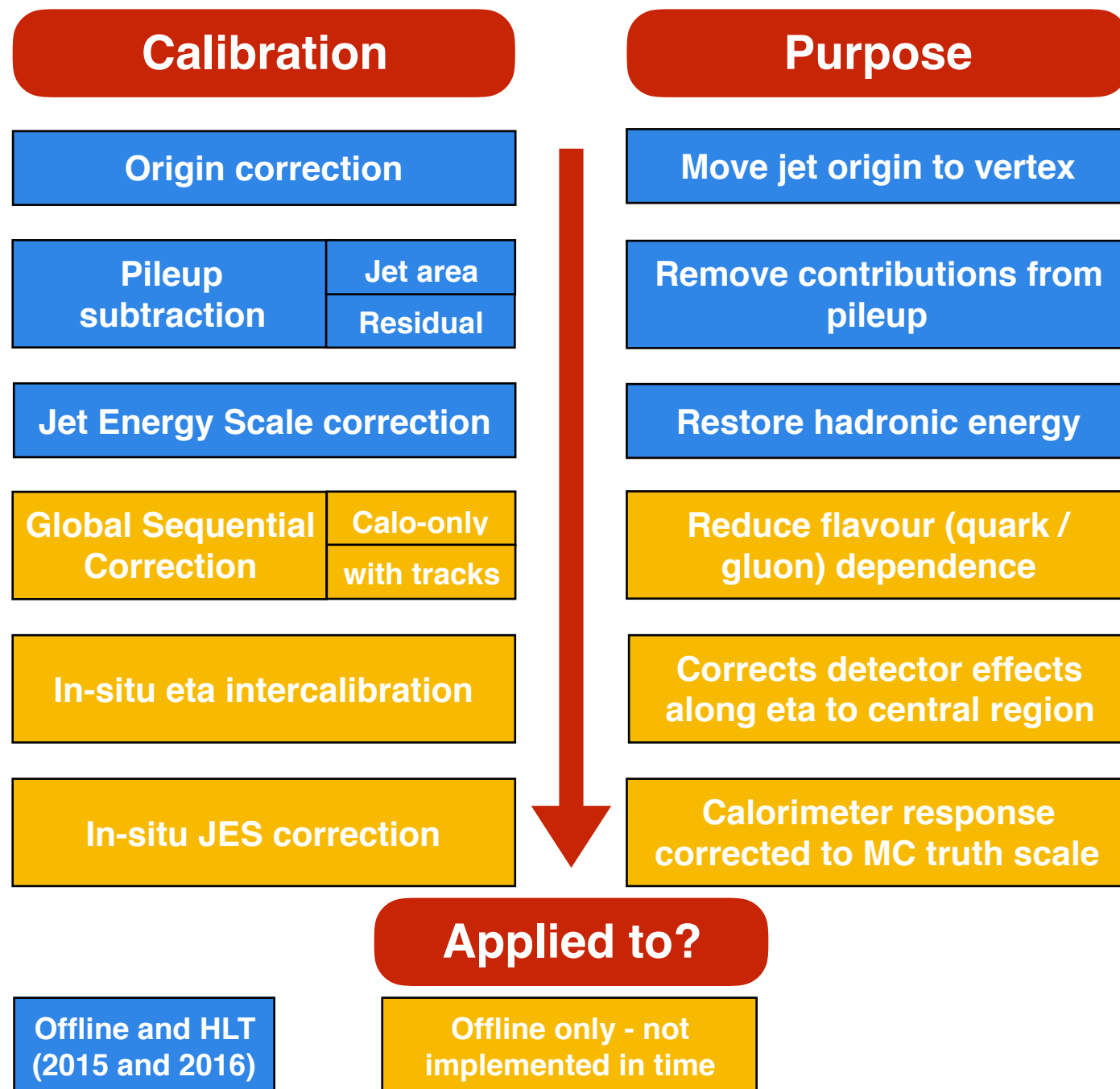
- partial event information (eg restricted / no tracking)
- trigger calibrations determined before data-taking, offline afterwards!

Jet trigger calibration



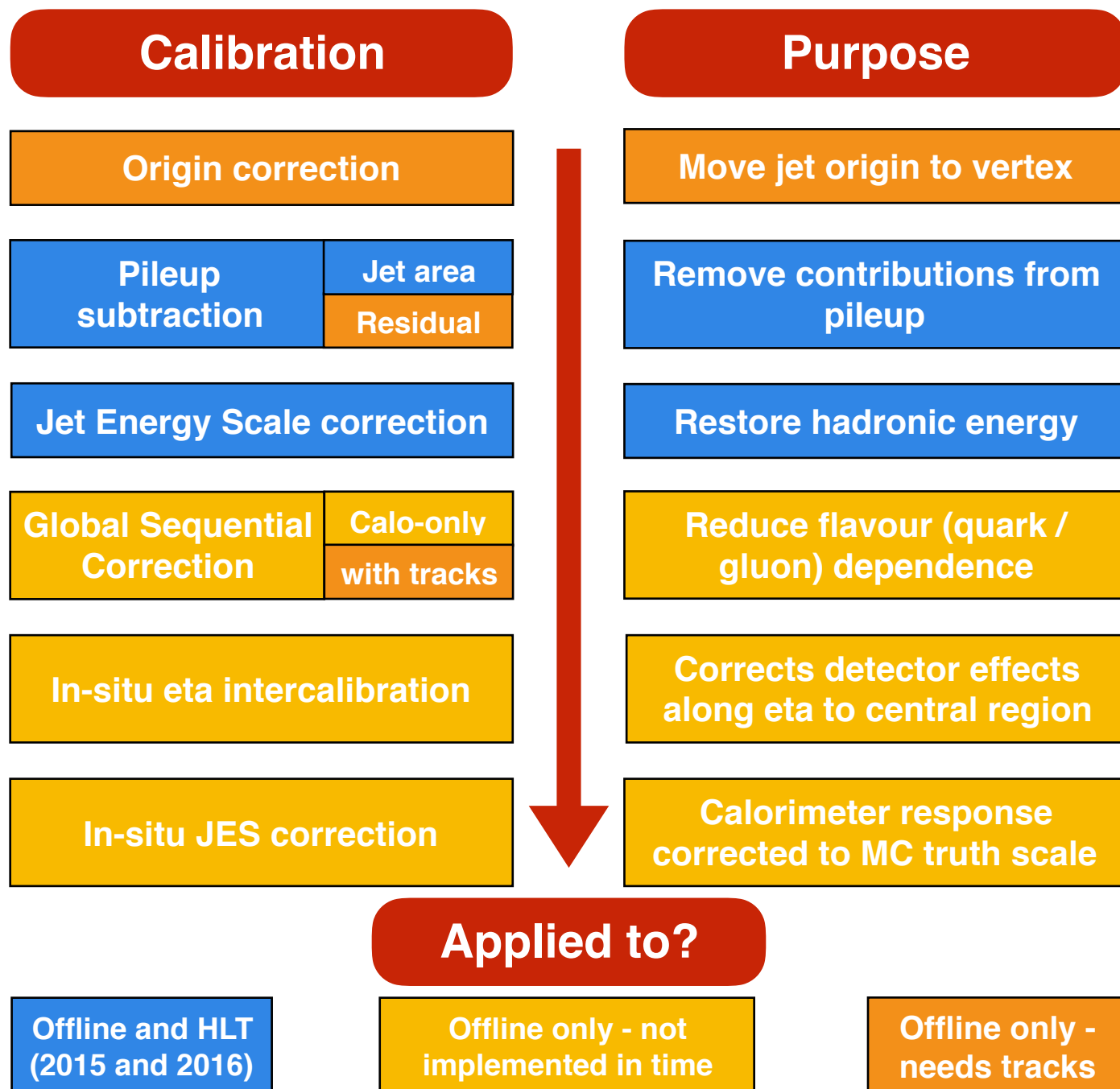
- Start with offline calibration chain

Jet trigger calibration



- Start with offline calibration chain
- No GSC or in-situ in 2015/16 data (developed using 2015 data!)

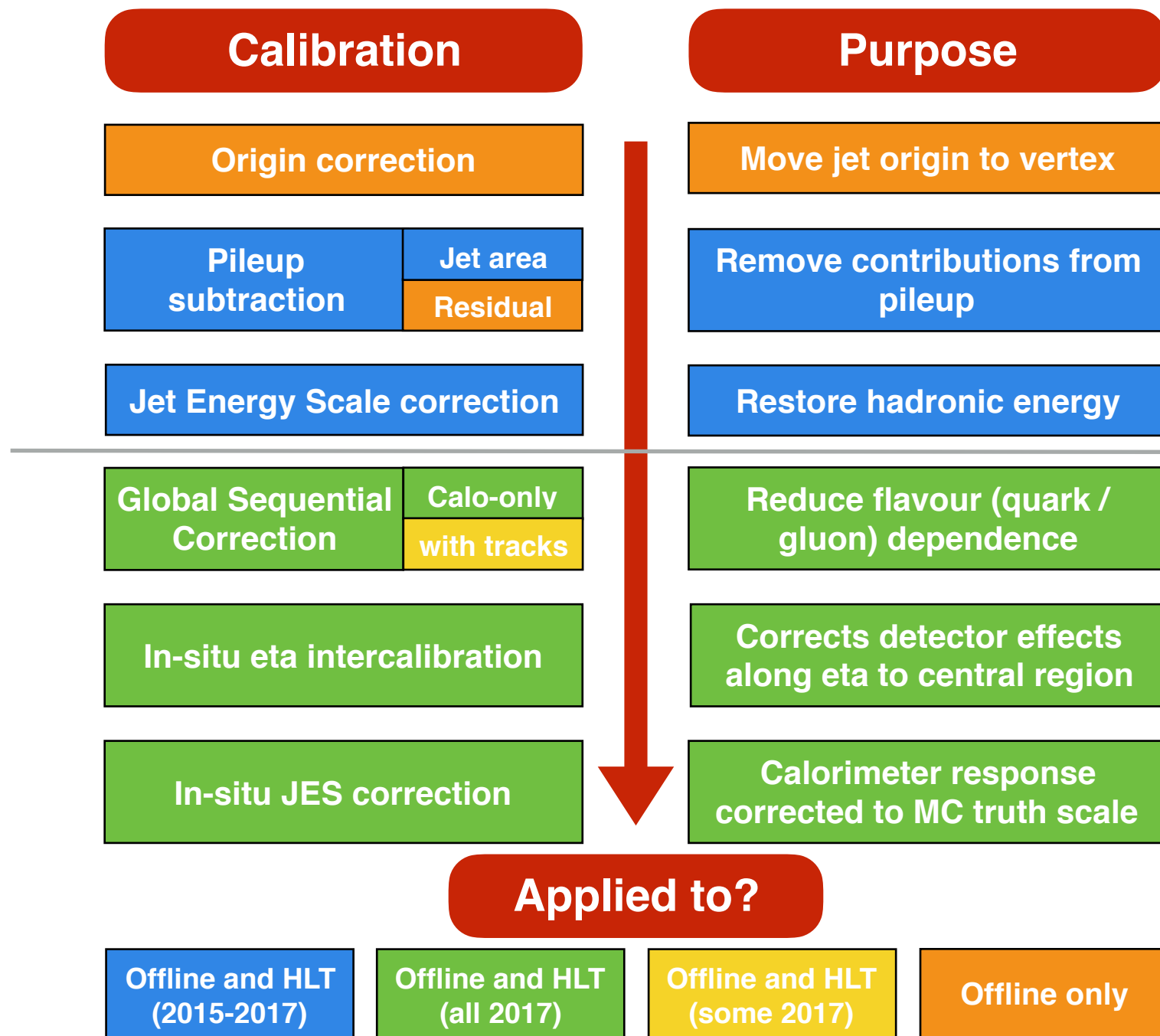
Jet trigger calibration



- Start with offline calibration chain
- No GSC or in-situ in 2015/16 data (developed using 2015 data!)
- Also: no tracks!
- very CPU intensive in ATLAS trigger -> infeasible to run full tracking

Status in 2015 and 2016 data

Jet trigger calibration



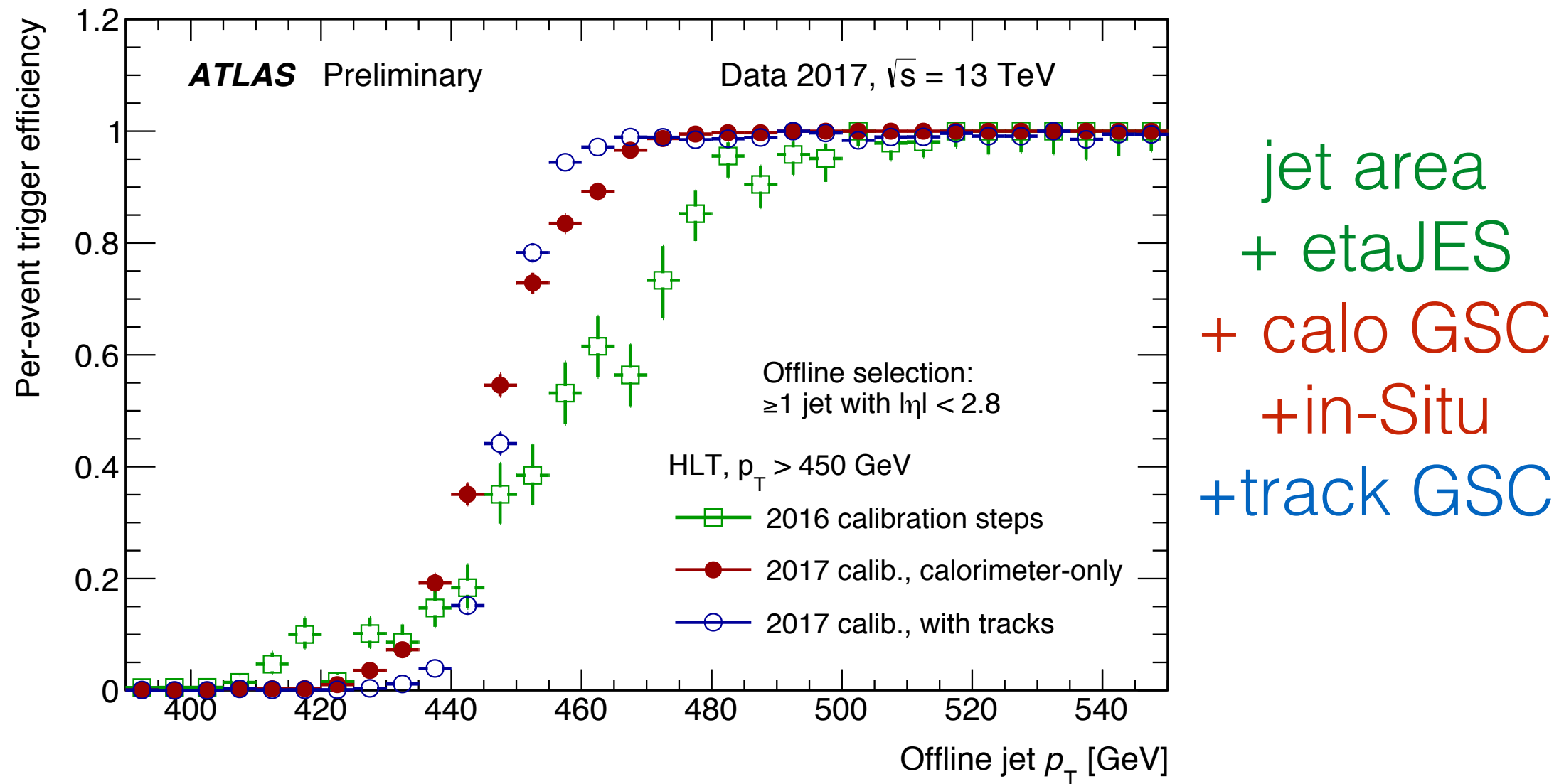
- New in 2017

- Apply partial GSC and in-situ calibrations to all trigger jets

- Some HLT tracking in jets is possible within CPU constraints - can apply GSC to some trigger jets

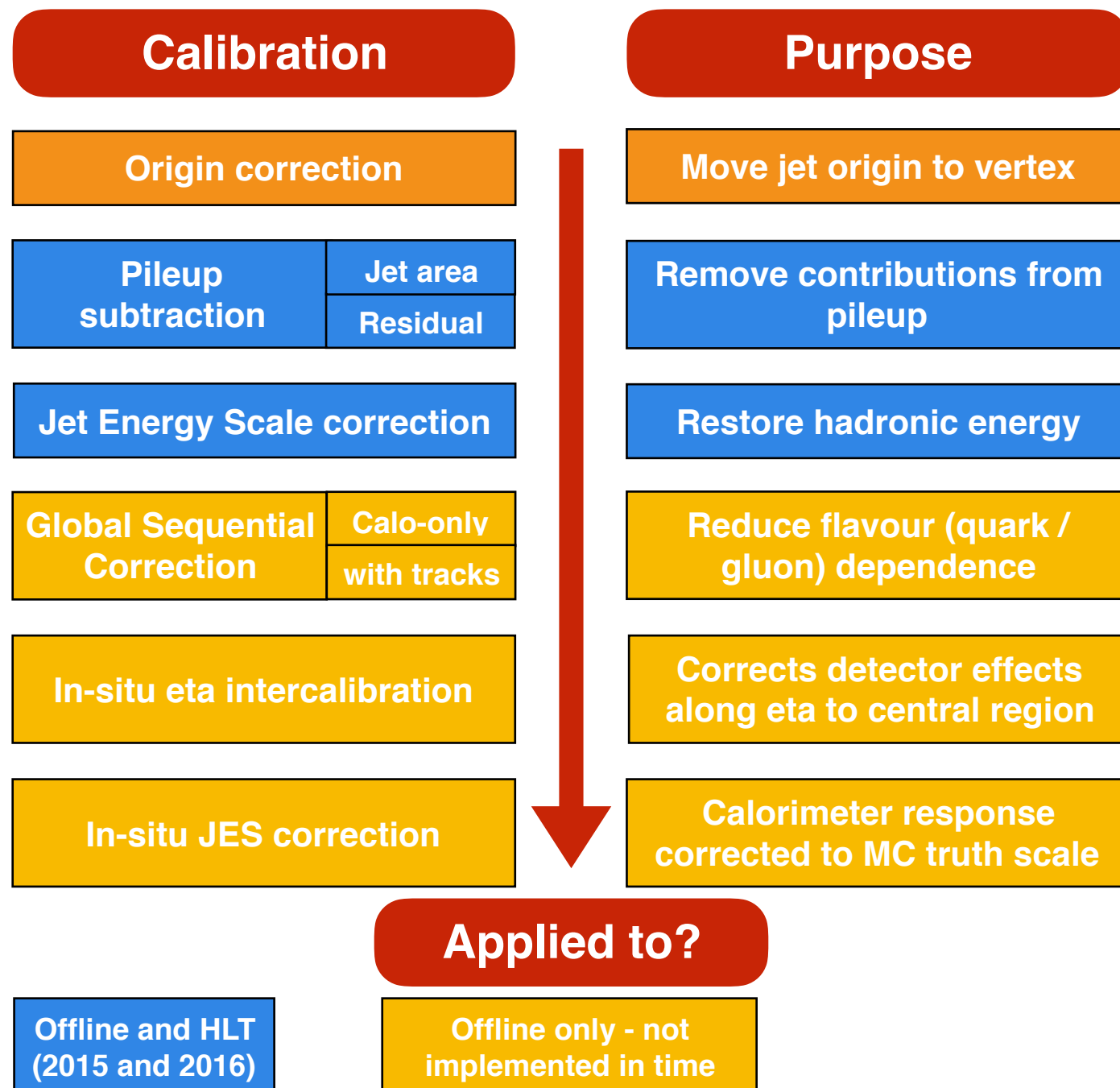
Status in 2017 data

Jet trigger calibration



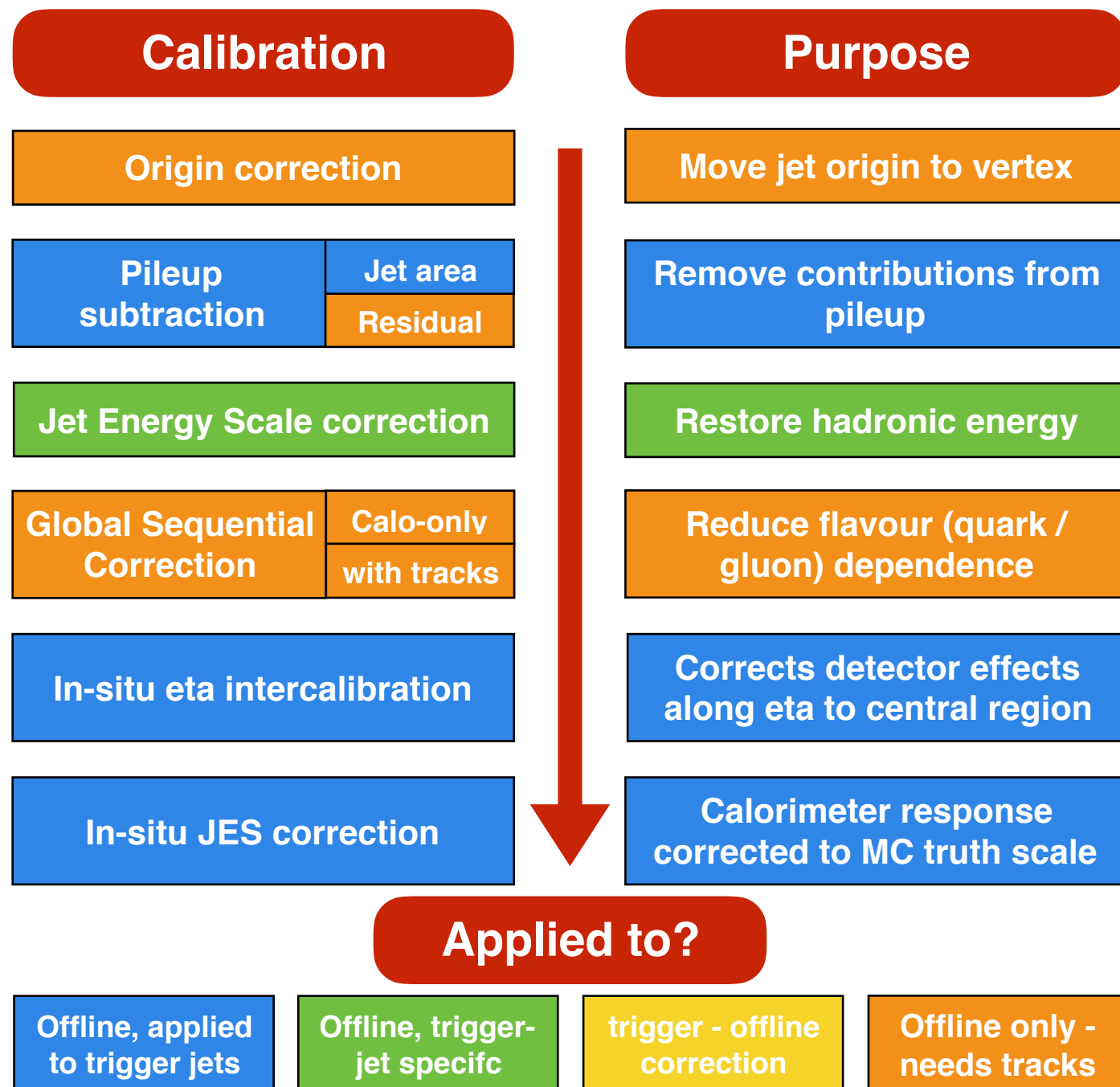
- Application of more steps in calibration chain hugely improves resolution and turnon
- Partially offsets threshold increases required from luminosity increases

Offline trigger jet calibration



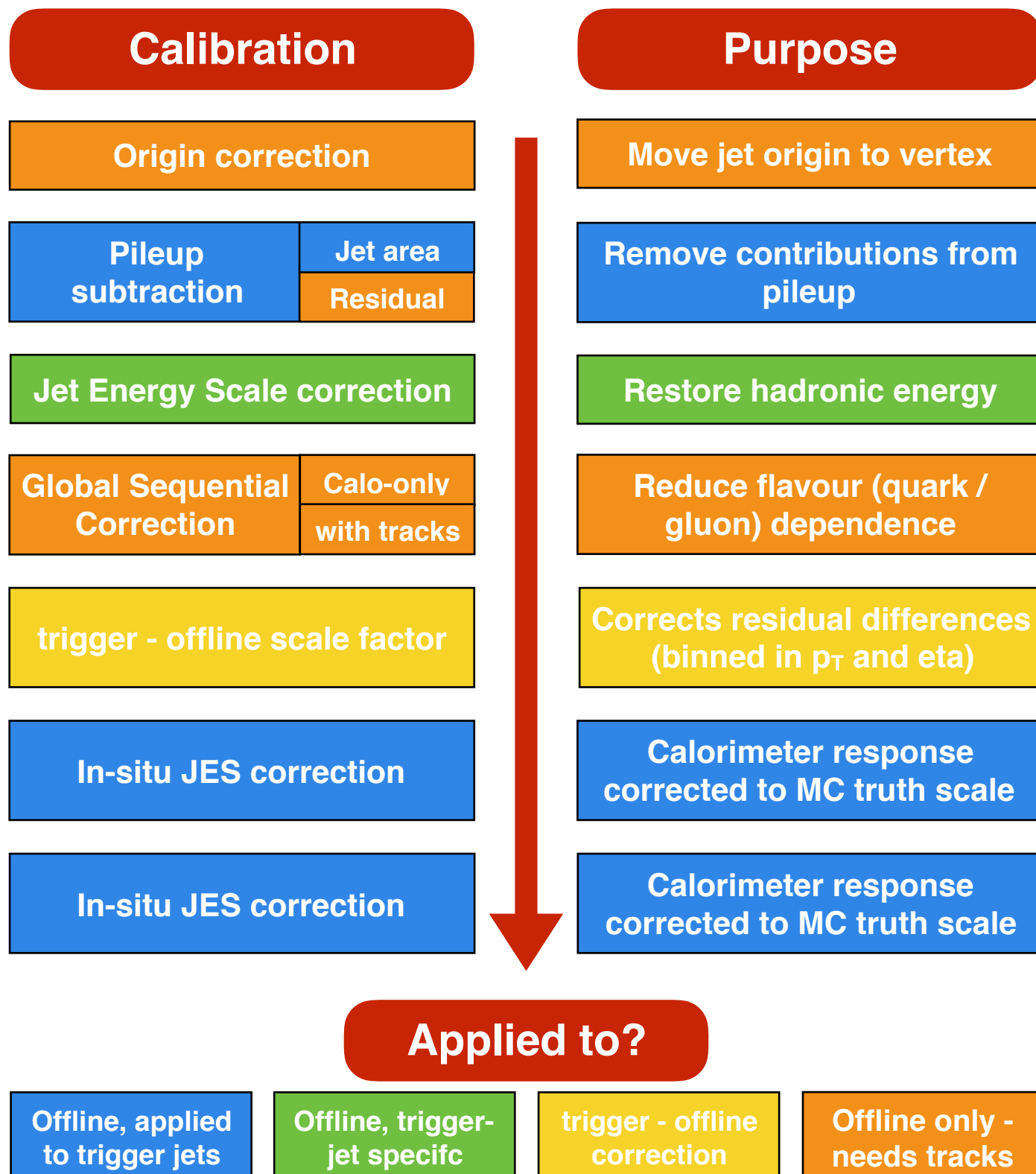
- We save enough information to be able to (re)do most of the calibration offline

Offline trigger jet calibration



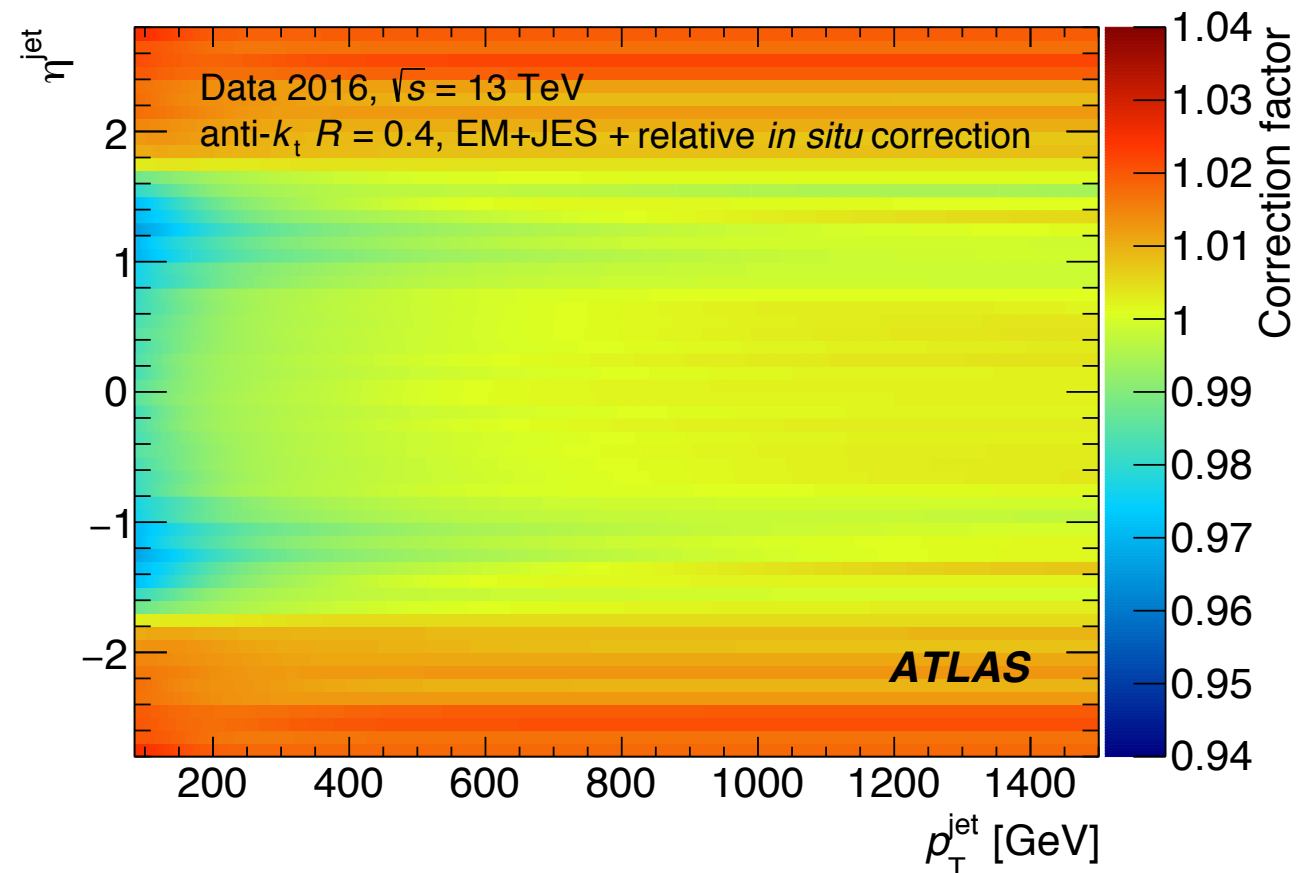
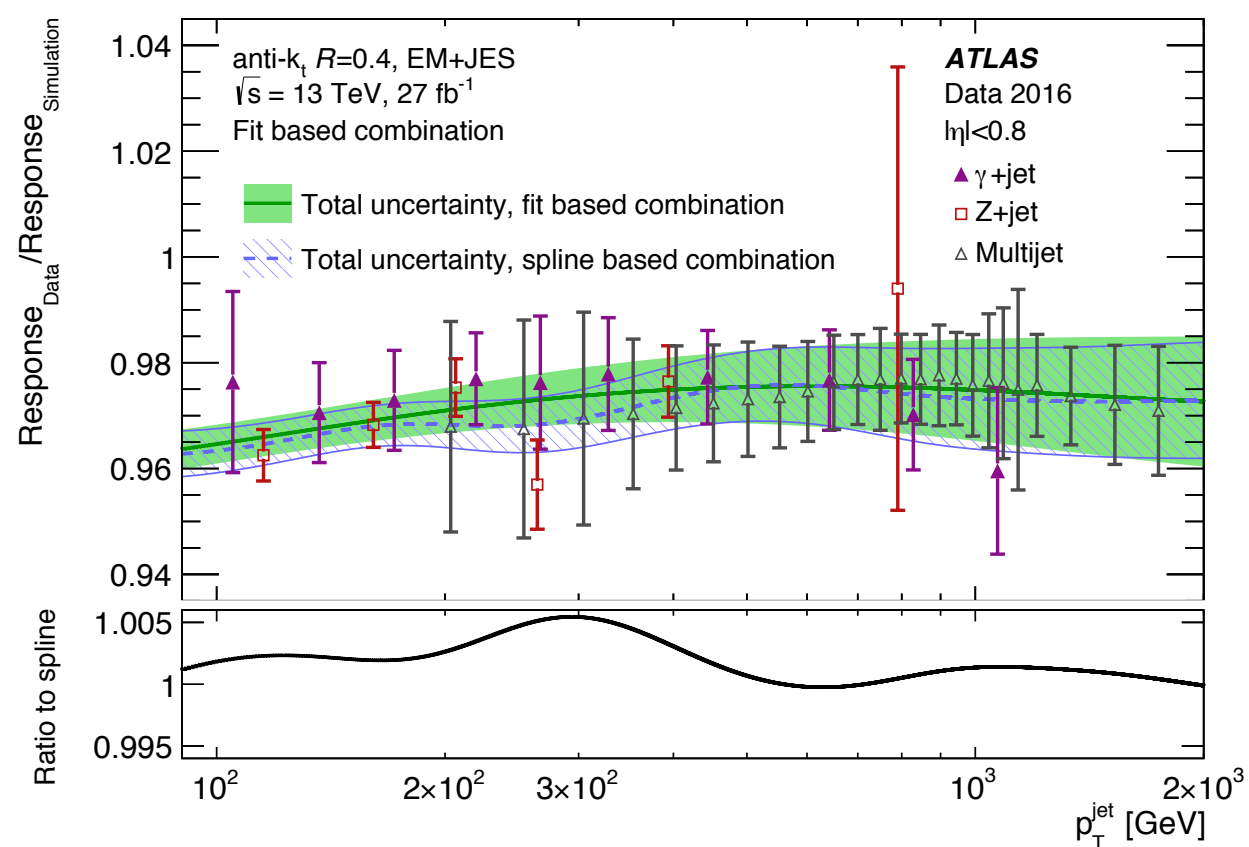
- We save enough information to be able to (re)do most of the calibration offline
- Some parts specifically redefined for trigger jets

Offline trigger jet calibration



- We save enough information to be able to (re)do most of the calibration offline
- Some parts specifically redefined for trigger jets
- Apply scale factor between trigger and offline jets to correct residual differences

TLA trigger jet calibration

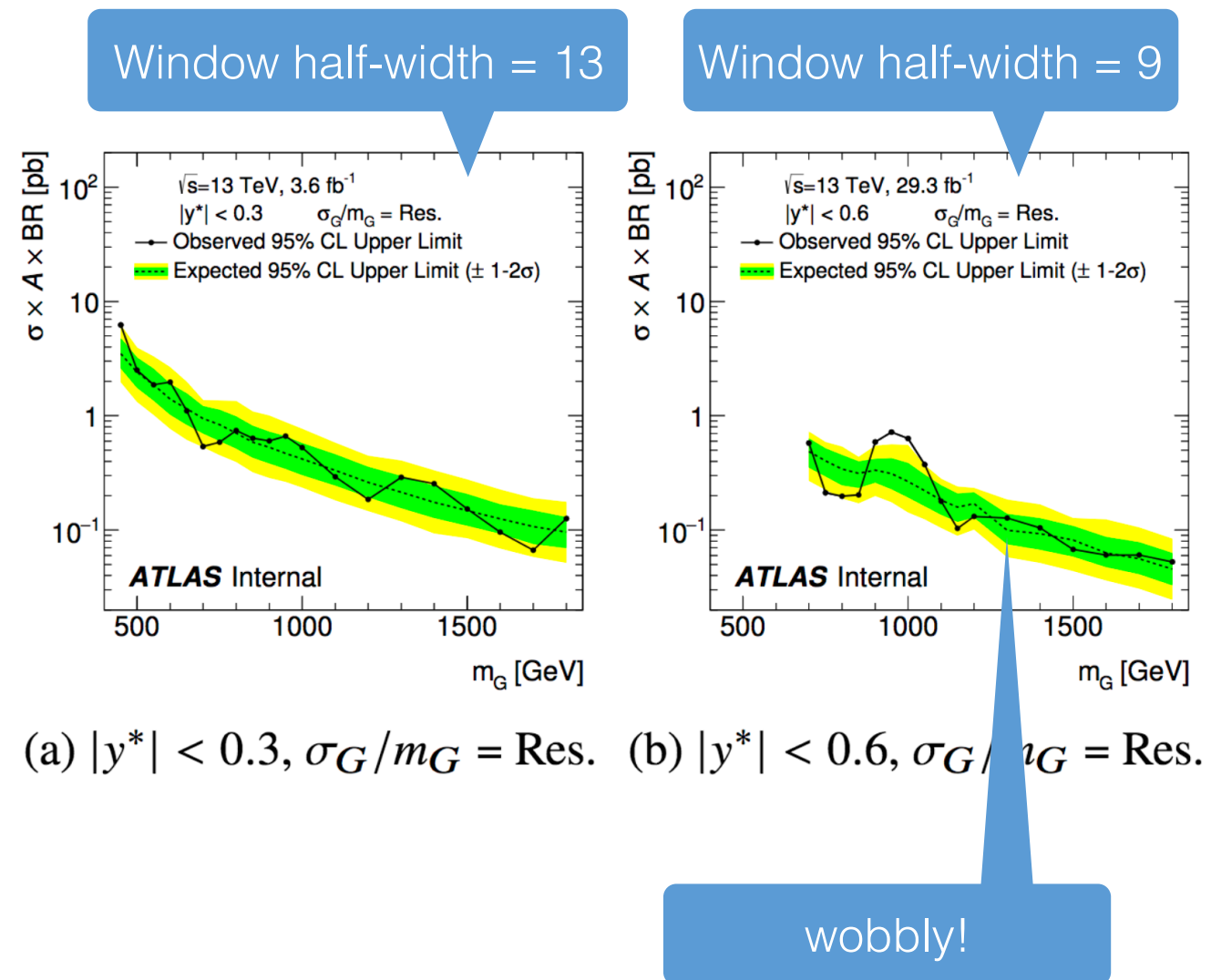


Custom “in-situ” step to ensure smoothness -
 statistical fluctuation in normal spline-based combination leads to bump in p_T and hence m_{jj}

**Excellent
 trigger : offline
 agreement**

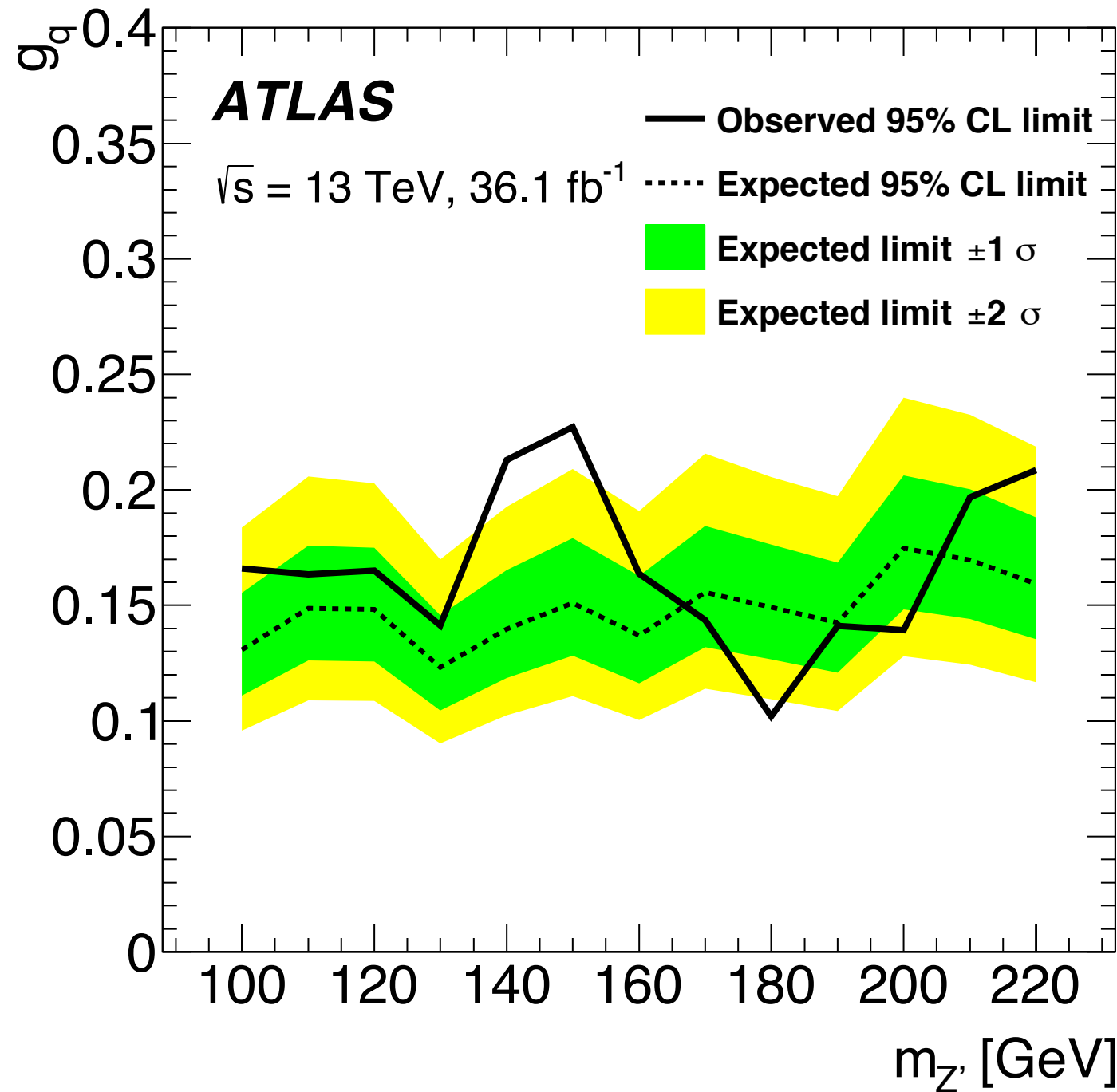
Expected limits fluctuations

- Real signal can exist in data, but expected limits need to represent signal-free background
 - Fit signal+background model for each signal point
 - Set signal component to zero & throw toys for expected limit
- Thus the model used to generate the expected limits is **different for each signal point**, since a different signal is included in each signal+background fit
 - Results in wobbly expected limits
 - **More pronounced the more “flexible” the background estimation is**



Large-R + ISR results

arxiv: 1801.08768, EXOT-2017-01



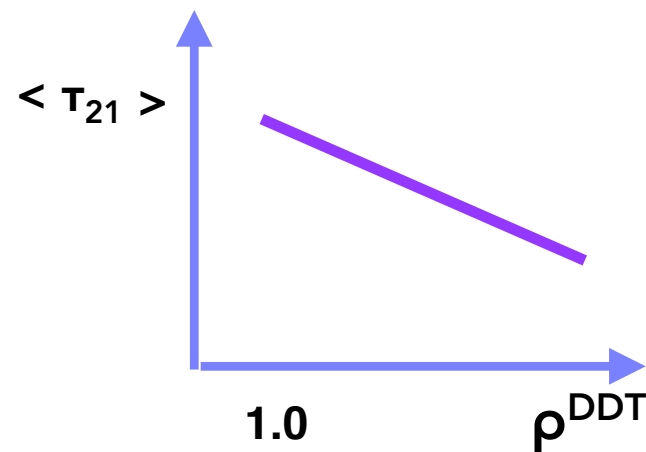
Observed and expected limits at 95% confidence level on the coupling (g_q), for the combination of the ISR jet and ISR γ channels

Large-R + ISR DDT

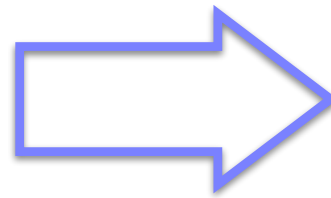
arxiv: 1801.08768, EXOT-2017-01

$$\rho^{\text{DDT}} \equiv \log \left(\frac{m_J^2}{p_T^J \times \mu} \right)$$

Linear relationship between ρ^{DDT} and $\langle \tau_{21} \rangle$ for $\rho^{\text{DDT}} > \sim 1$

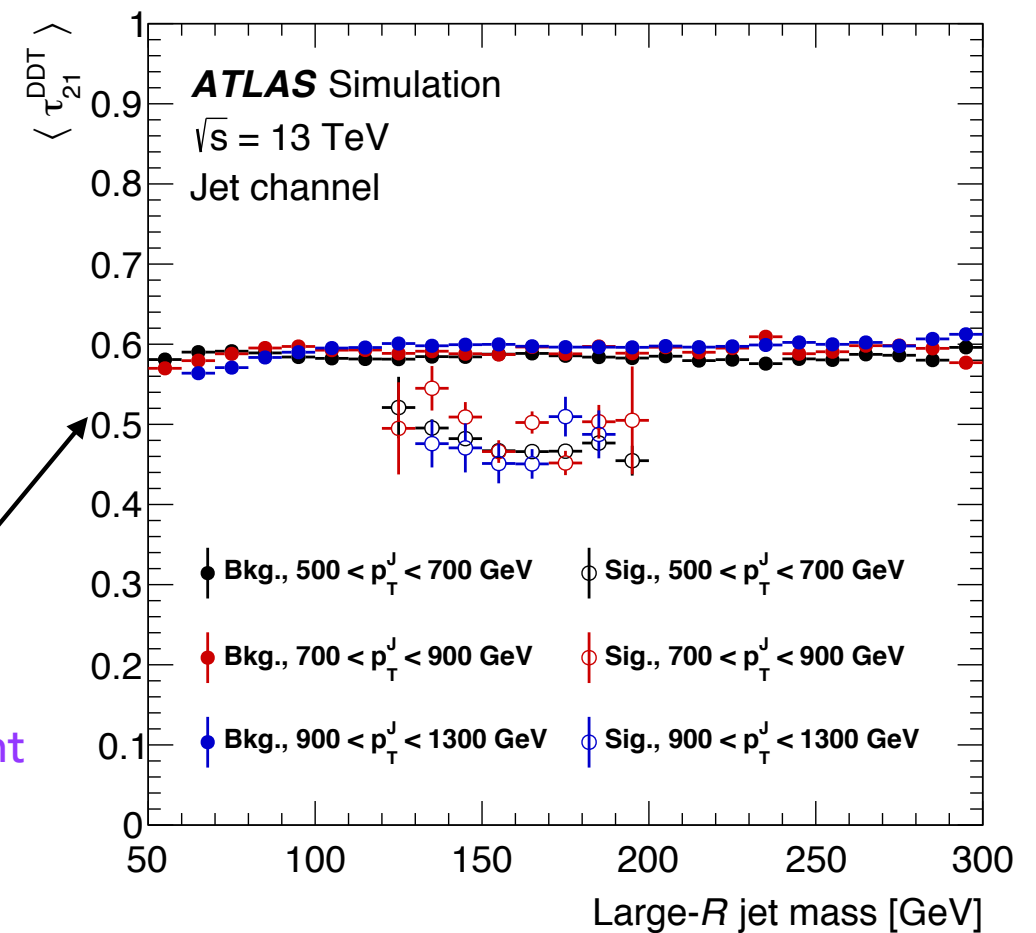


Define τ_{21}^{DDT} : linearly corrected version of τ_{21}

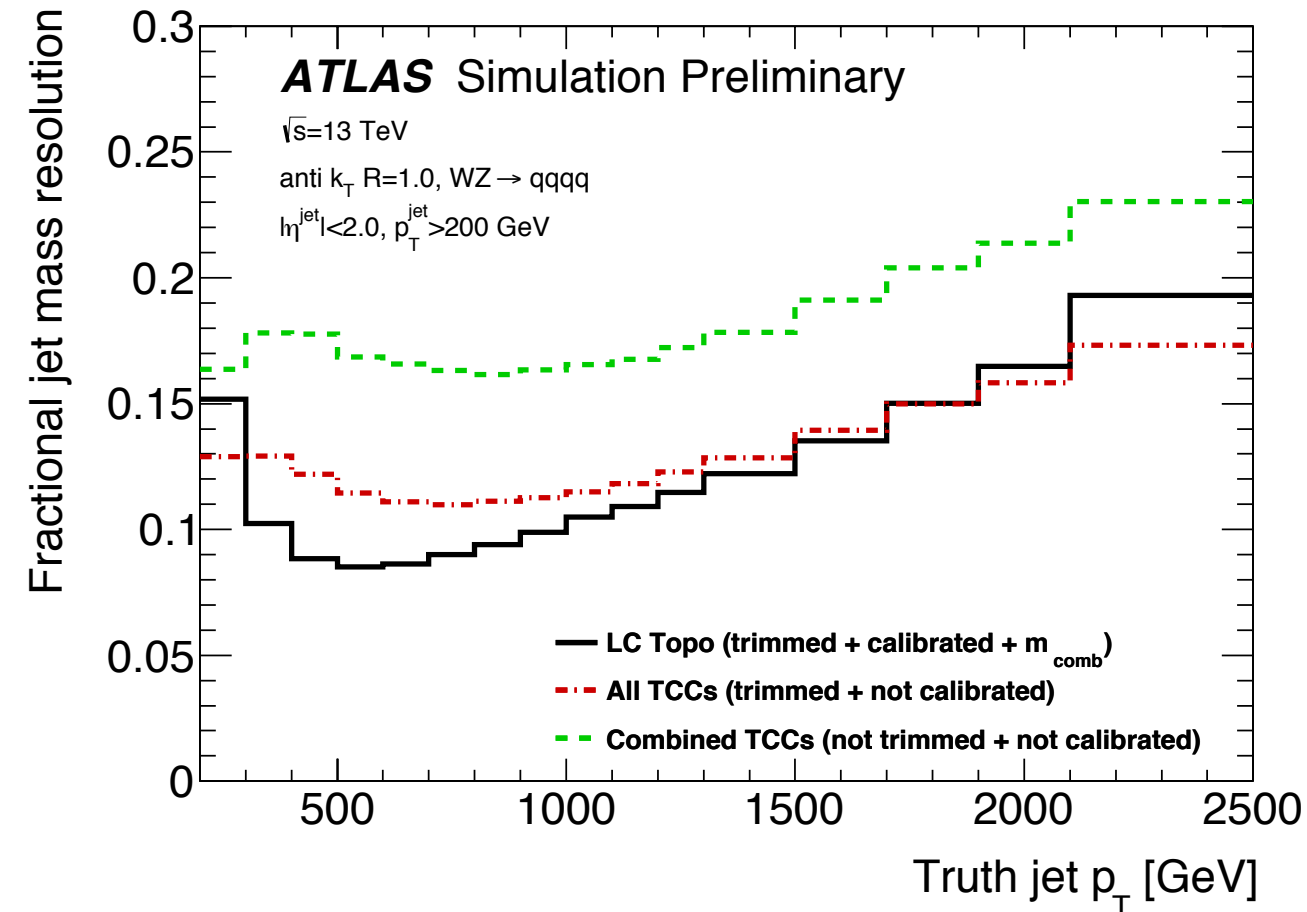
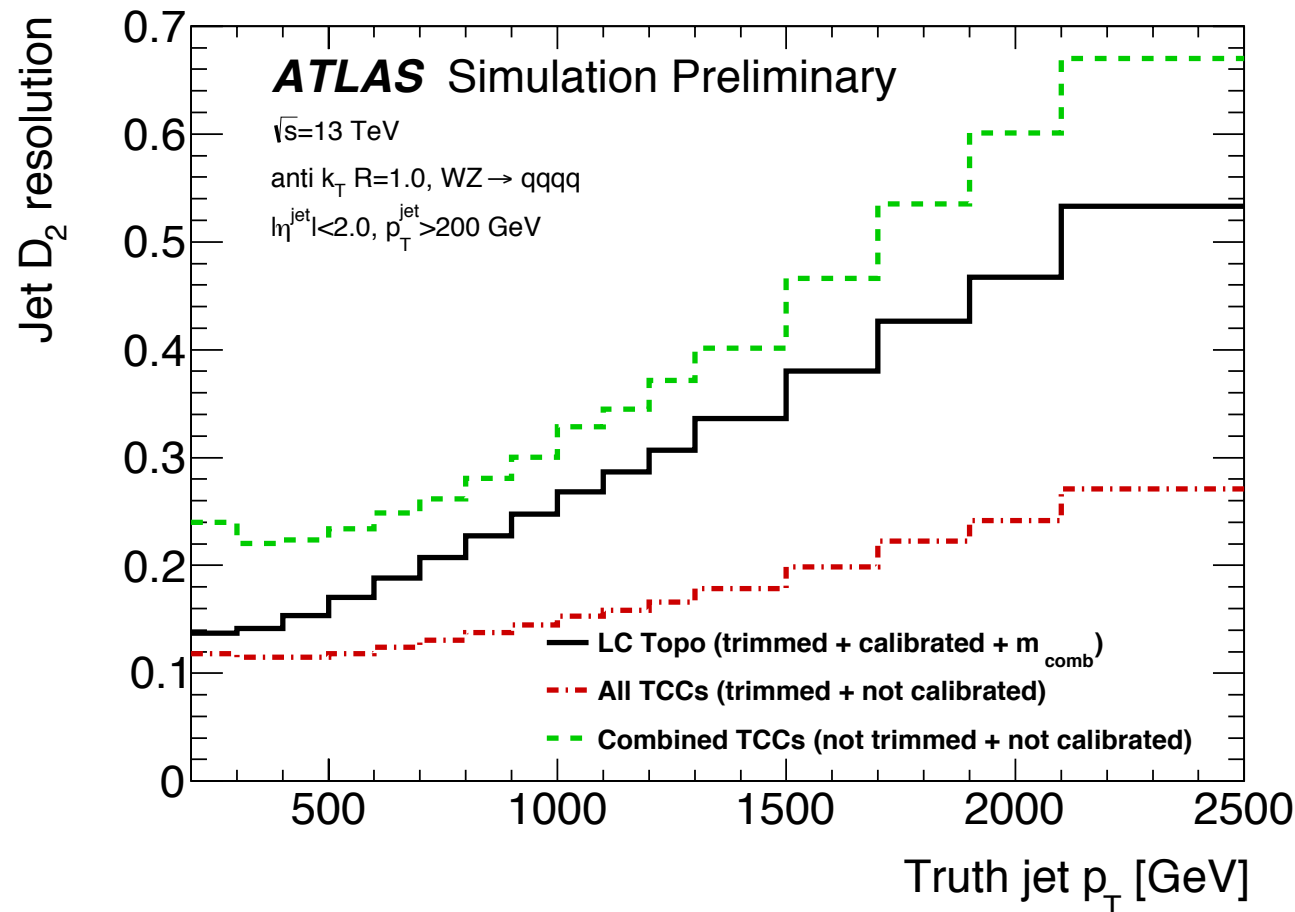


τ_{21}^{DDT} independent of jet mass

arXiv:1801.08769

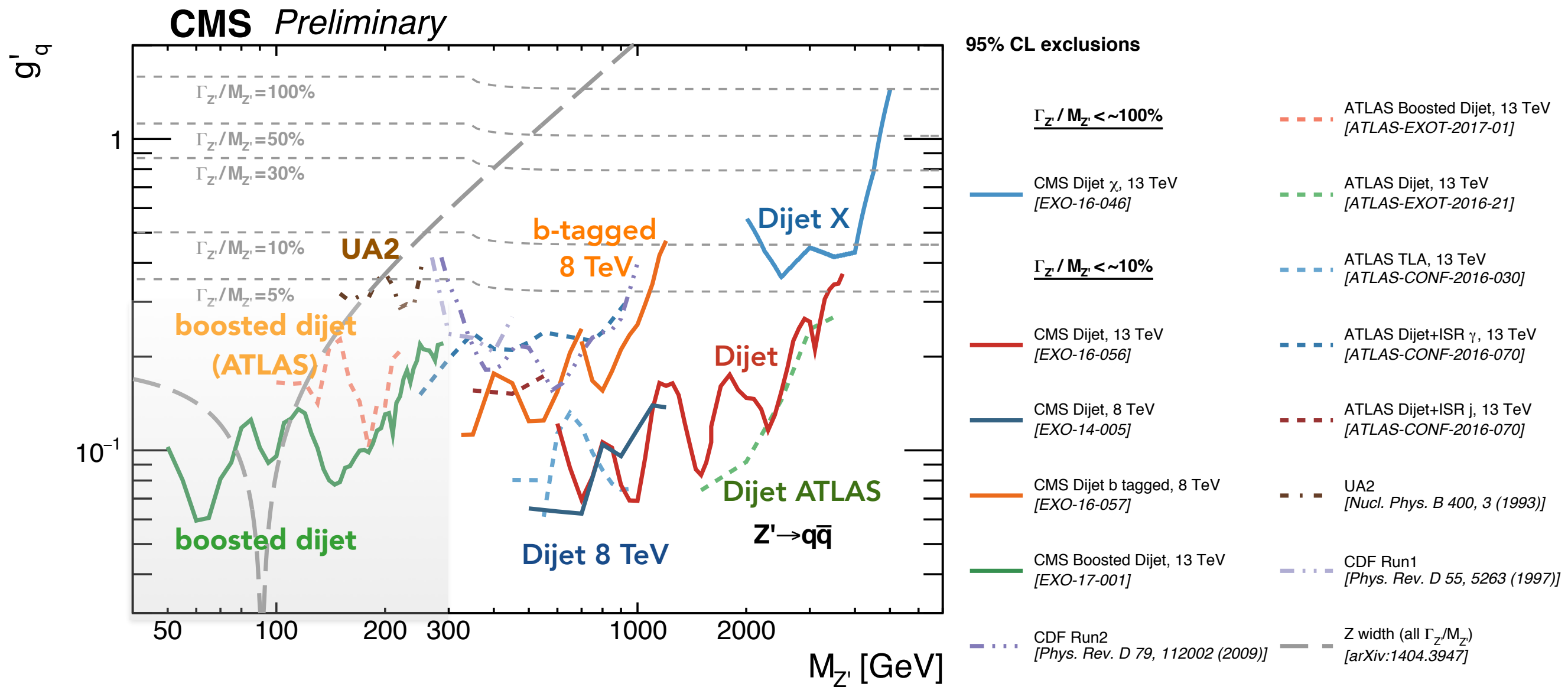


Tracking in CaloClusters



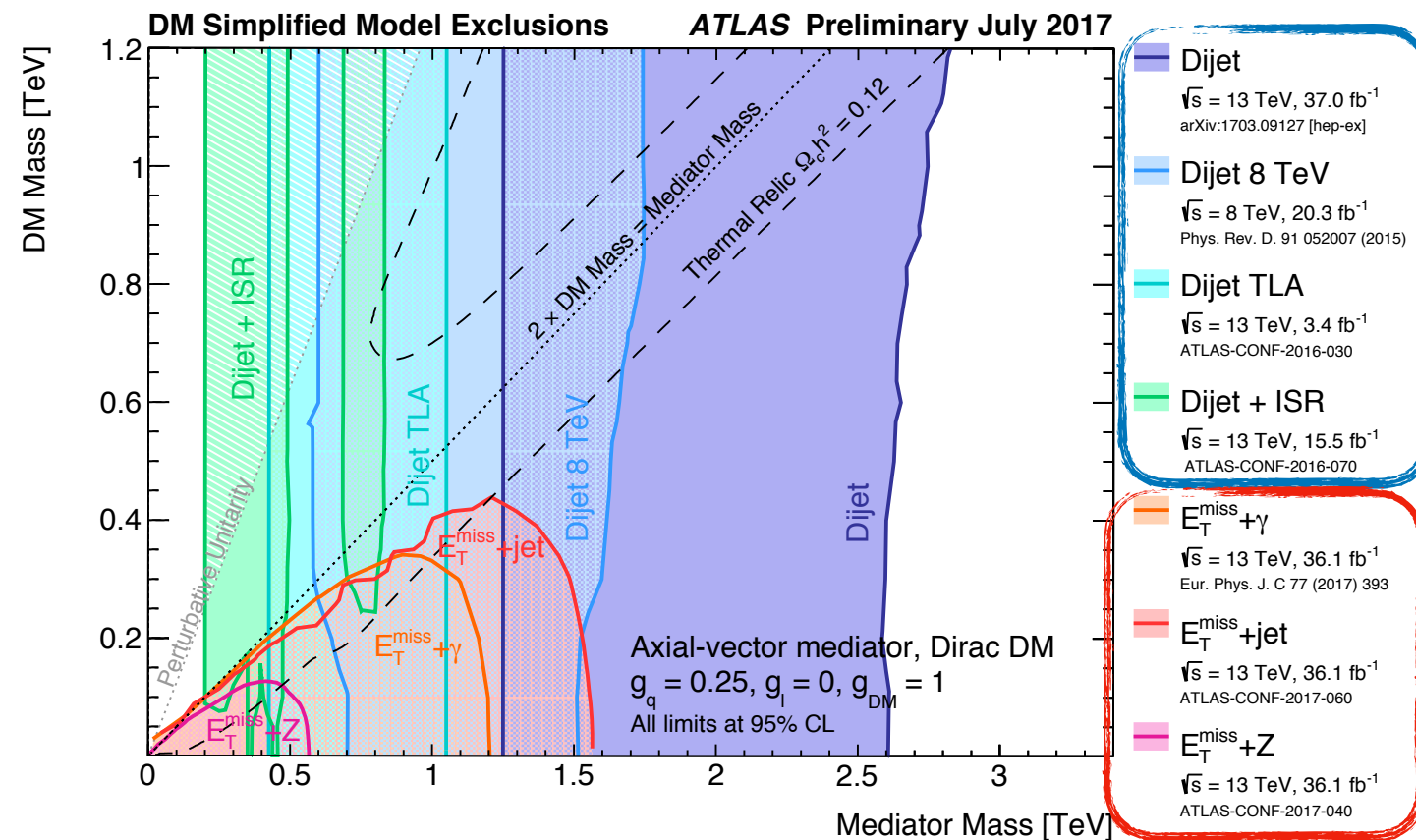
- Improvements in jet substructure resolution thanks to track information in jet reconstruction inputs ATL-PHYS-PUB-2017-015
- Black -> Red
 - Mostly low p_T -> improvement in D_2 , degradation in mass

CMS and ATLAS limits



ATLAS TLA updated since this plot

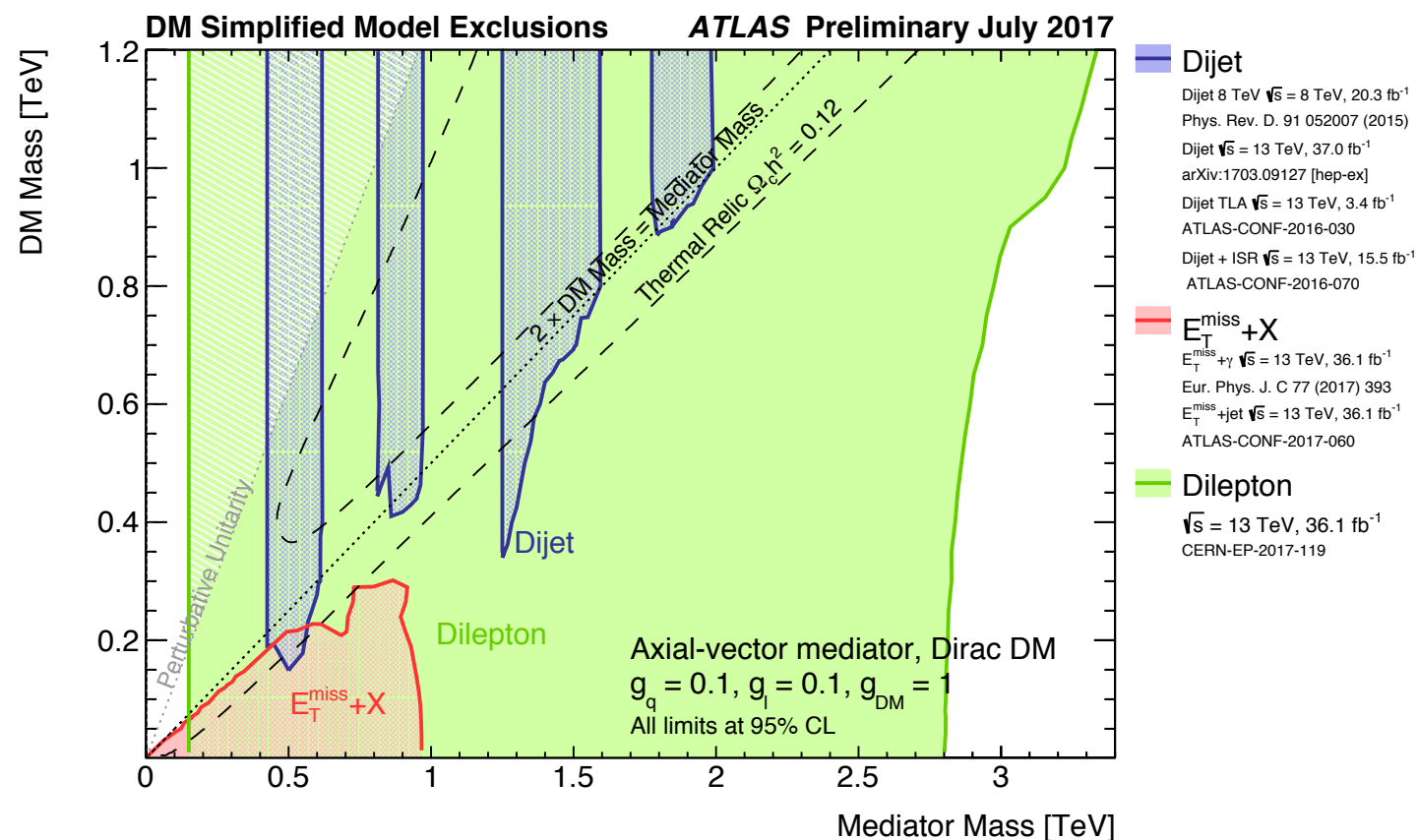
Wider context



dijet

mono-X

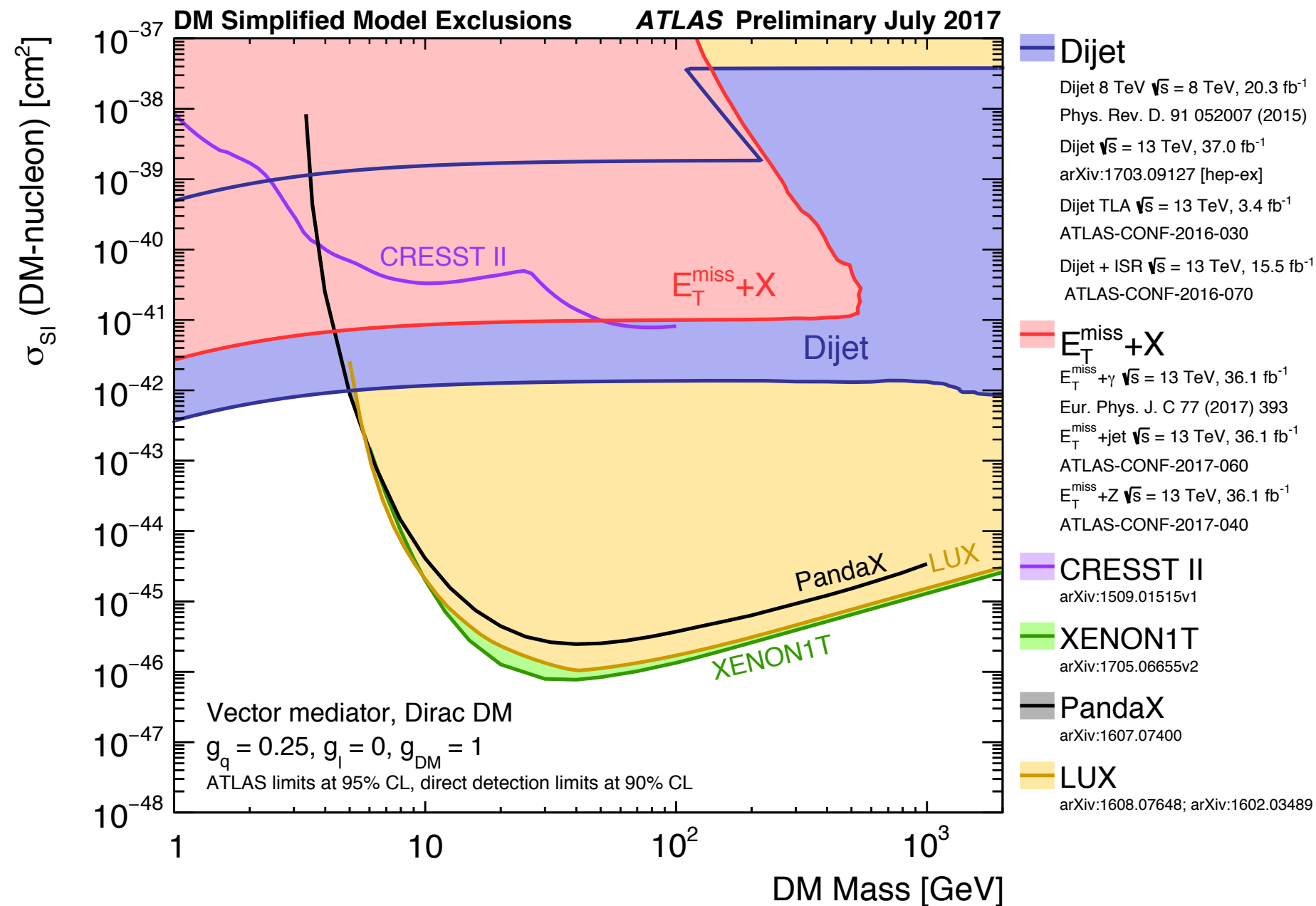
Sensitivity decreases
 as DM mass decreases
 (Z' branching ratio to
 dijets decreases)
 -> covered by mono-X
 searches



Interpretation is very
 model-dependent

Sensitivity decreases
 as lepton coupling g_l
 increases and quark
 coupling g_q decreases
 -> covered by dilepton
 resonance searches

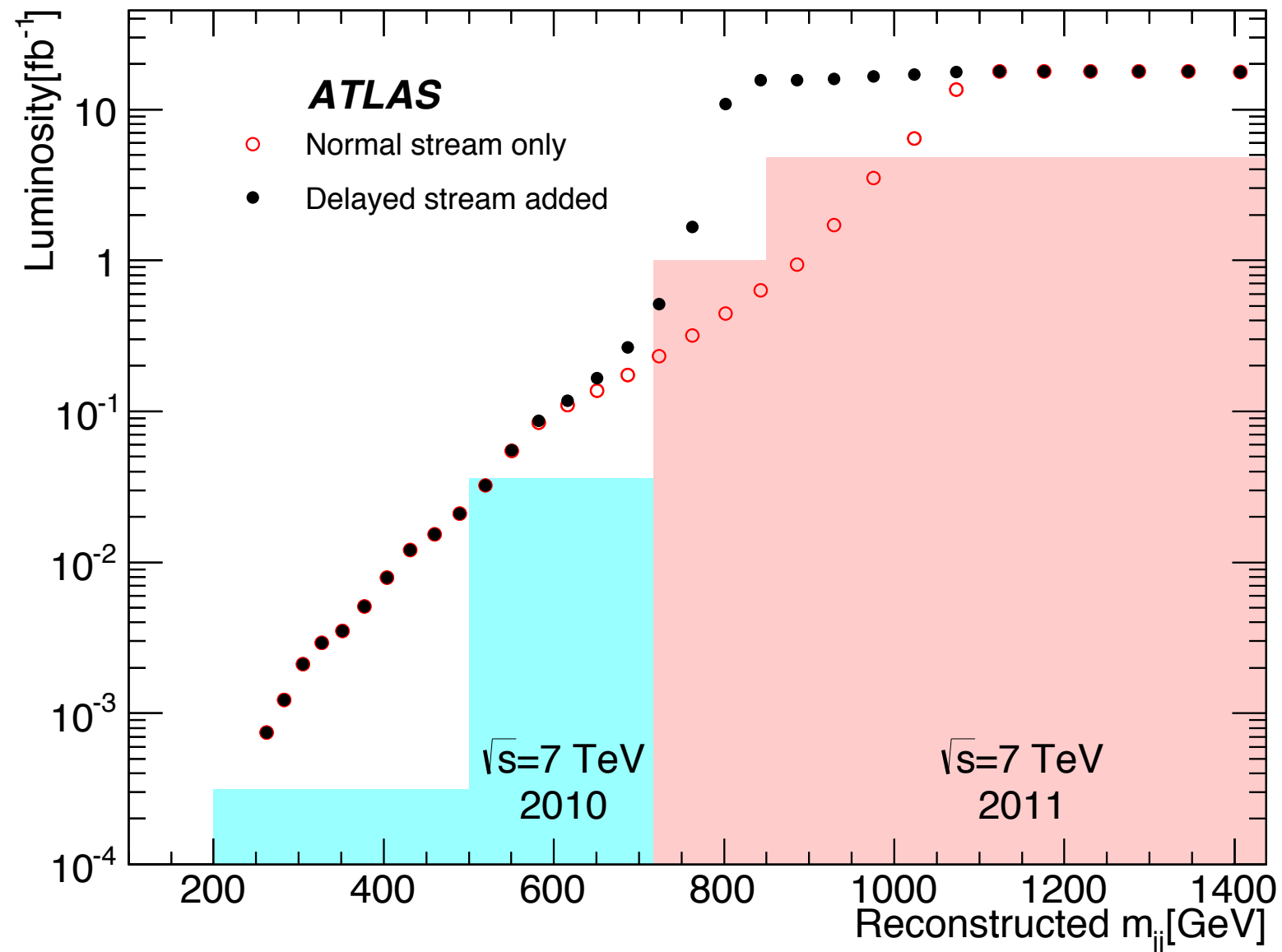
Even wider context



**Interpretation is
even more
model-
dependent**

Nice complementarity between direct detection,
collider production with mono-X and “indirect
searches” with dijet resonances

8 TeV 20.3 fb⁻¹ triggers



prescaled single jet triggers
plus delayed stream