

# New parton densities with Parton Branching method

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A. Bermudez Martinez (DESY)

in collaboration with

P. Connor, F. Hautmann, H. Jung, A. Lelek, V. Radescu, R. Zlebcik

- Why TMDs are needed
  - TMDs for hadron-hadron collisions
- New developments
  - parton branching algorithm to solve evolution equations
    - benchmark tests
    - advantages for integrated PDFs
  - determination of TMD densities at NLO with xFitter
- Application to DY production and high  $p_T$  dijets

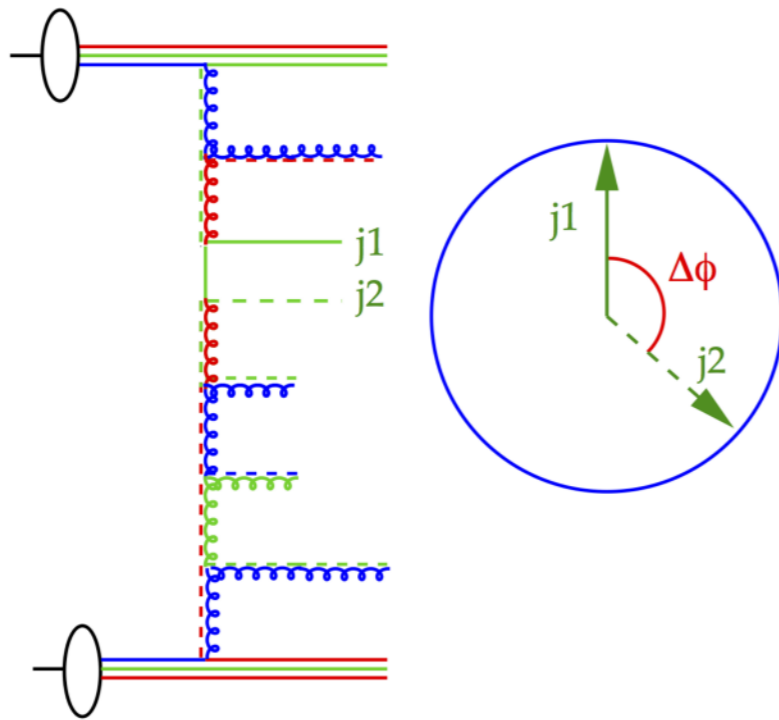
# TMDs – what is it ?

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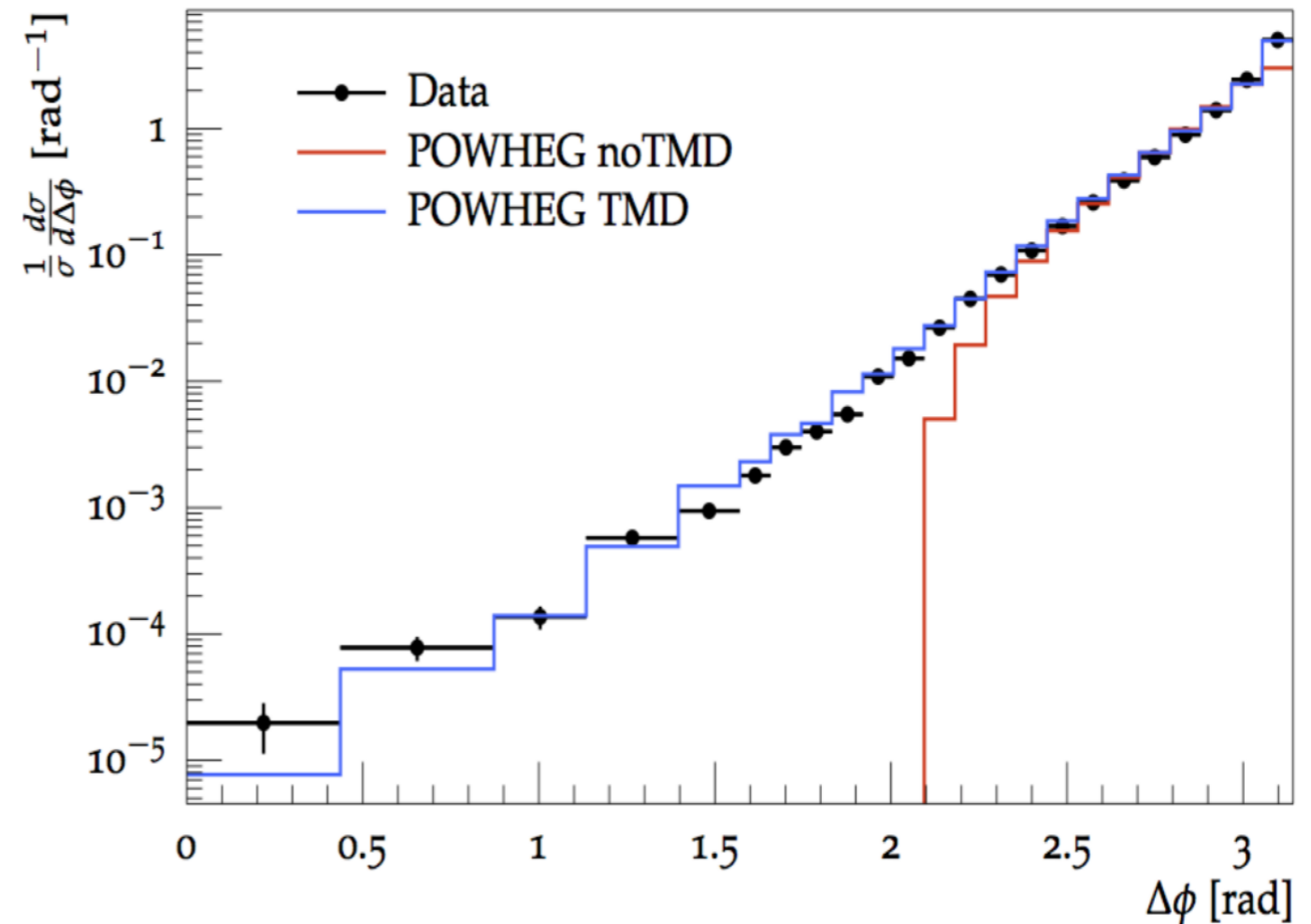
- TMDs (Transverse Momentum Dependent parton distribution)
  - at very small transverse momenta
    - typically for small  $q_t$  in DY production, or semi-inclusive DIS
  - at very small  $x$  – un-integrated PDFs
    - essentially only gluon densities (CCFM, BFKL etc)
- new approach to cover all transverse momenta from small  $k_t$  to large  $k_t$  as well as to cover all  $x$  and all  $\mu^2$ 
  - parton branching method (described here)

# Why TMDs ?

- Measurements with  $p_T > 200$  GeV
  - at least 2 jets



Di-jet azimuthal decorrelation,  $300 < p_T^{\text{leading}} < 400$  GeV



- NLO-dijet (Powheg) w/o PS cannot describe small  $\Delta\phi$
- NLO-dijet (Powheg) with TMDs describes spectrum at small and large  $\Delta\phi$
- Region of higher order emissions described by TMDs

# TMDs – how to determine ?

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- Transverse momentum effects are naturally coming from intrinsic  $k_t$  and parton showers
- TMD effects can be significant in all distributions, even for inclusive (or semi-inclusive) distributions at large  $p_t$
- **New: parton branching method**
  - perform evolution using a parton branching method
  - determine integrated PDF from parton branching solution of evolution eq.
    - check consistency with standard evolution on integrated PDFs
      - at LO, NLO and NNLO
  - determine TMD:
    - since each branching is generated explicitly, energy-momentum conservation is fulfilled and transverse momentum distributions can be obtained

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# How to obtain TMDs – the evolution equation

- [1] F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik. Soft-gluon resolution scale in QCD evolution equations. *Phys. Lett.*, B772:446–451, 2017.
- [2] F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik. Collinear and TMD Quark and Gluon Densities from Parton Branching Solution of QCD Evolution Equations. *JHEP*, 01:070, 2018.
- [3] A. Bermudez-Martinez, P. Connor, F. Hautmann, H. Jung, A. Lelek, V. Radescu, and R. Zlebcik. Collinear and TMD parton densities determined from fits to HERA DIS measurements, DESY-18-042

# DGLAP evolution – solution with parton branching method

- differential form: 
$$\mu^2 \frac{\partial}{\partial \mu^2} f(x, \mu^2) = \int \frac{dz}{z} \frac{\alpha_s}{2\pi} P_+(z) f\left(\frac{x}{z}, \mu^2\right)$$

$$\Delta_s(\mu^2) = \exp\left(-\int^{z_M} dz \int_{\mu_0^2}^{\mu^2} \frac{\alpha_s}{2\pi} \frac{d\mu'^2}{\mu'^2} P^{(R)}(z)\right)$$

- differential form using  $f/\Delta_s$  with

$$\mu^2 \frac{\partial}{\partial \mu^2} \frac{f(x, \mu^2)}{\Delta_s(\mu^2)} = \int \frac{dz}{z} \frac{\alpha_s}{2\pi} \frac{P^{(R)}(z)}{\Delta_s(\mu^2)} f\left(\frac{x}{z}, \mu^2\right)$$

- integral form

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$


  
 no – branching probability from  $\mu_0^2$  to  $\mu^2$

# DGLAP – solution with parton branching method

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$

- solve integral equation via iteration:

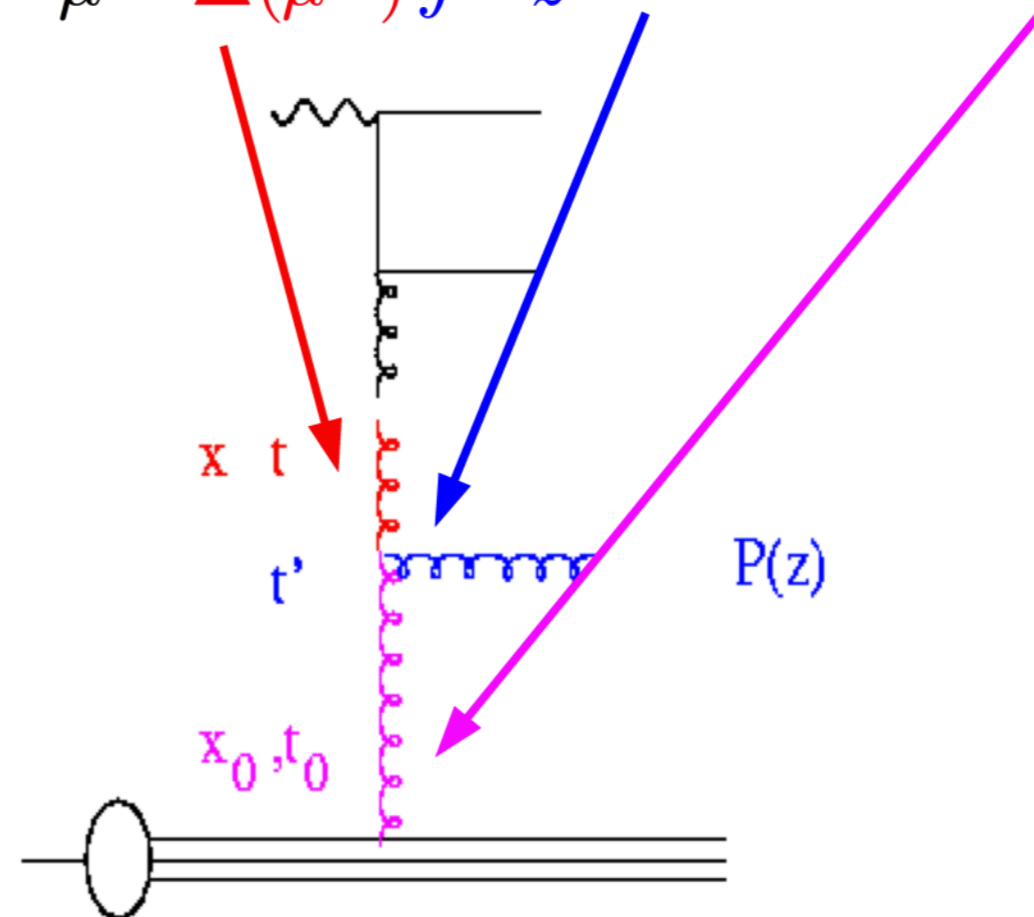
$$f_0(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2)$$

$$f_1(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2) + \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta(\mu^2)}{\Delta(\mu'^2)} \int \frac{dz}{z} P^{(R)}(z) f(x/z, \mu_0^2) \Delta(\mu'^2)$$

from  $t'$  to  $t$   
w/o branching

branching at  $t'$

from  $t_0$  to  $t'$   
w/o branching



# Evolution equation and parton branching method

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- use momentum weighted PDFs:  $x f(x, t)$

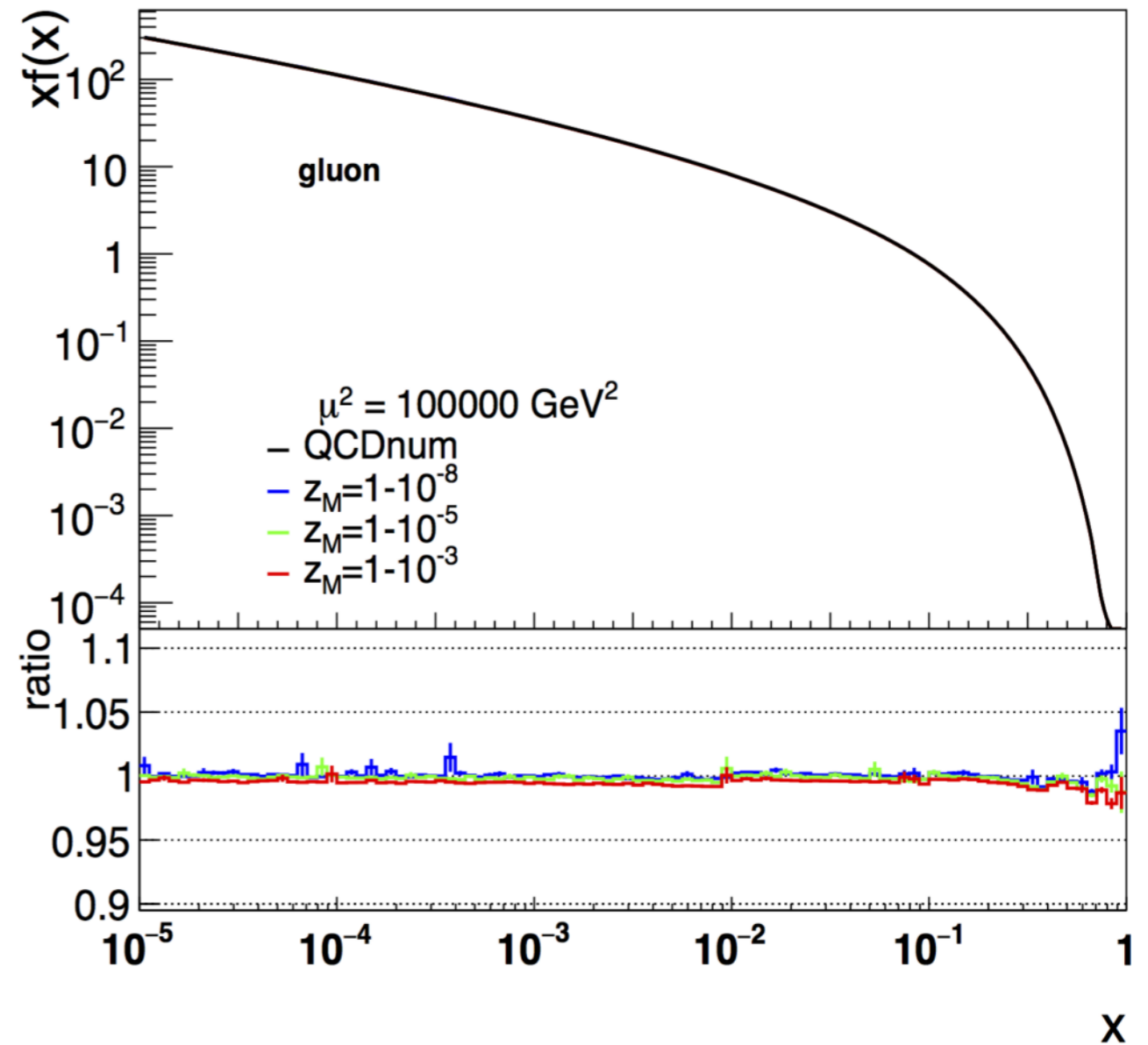
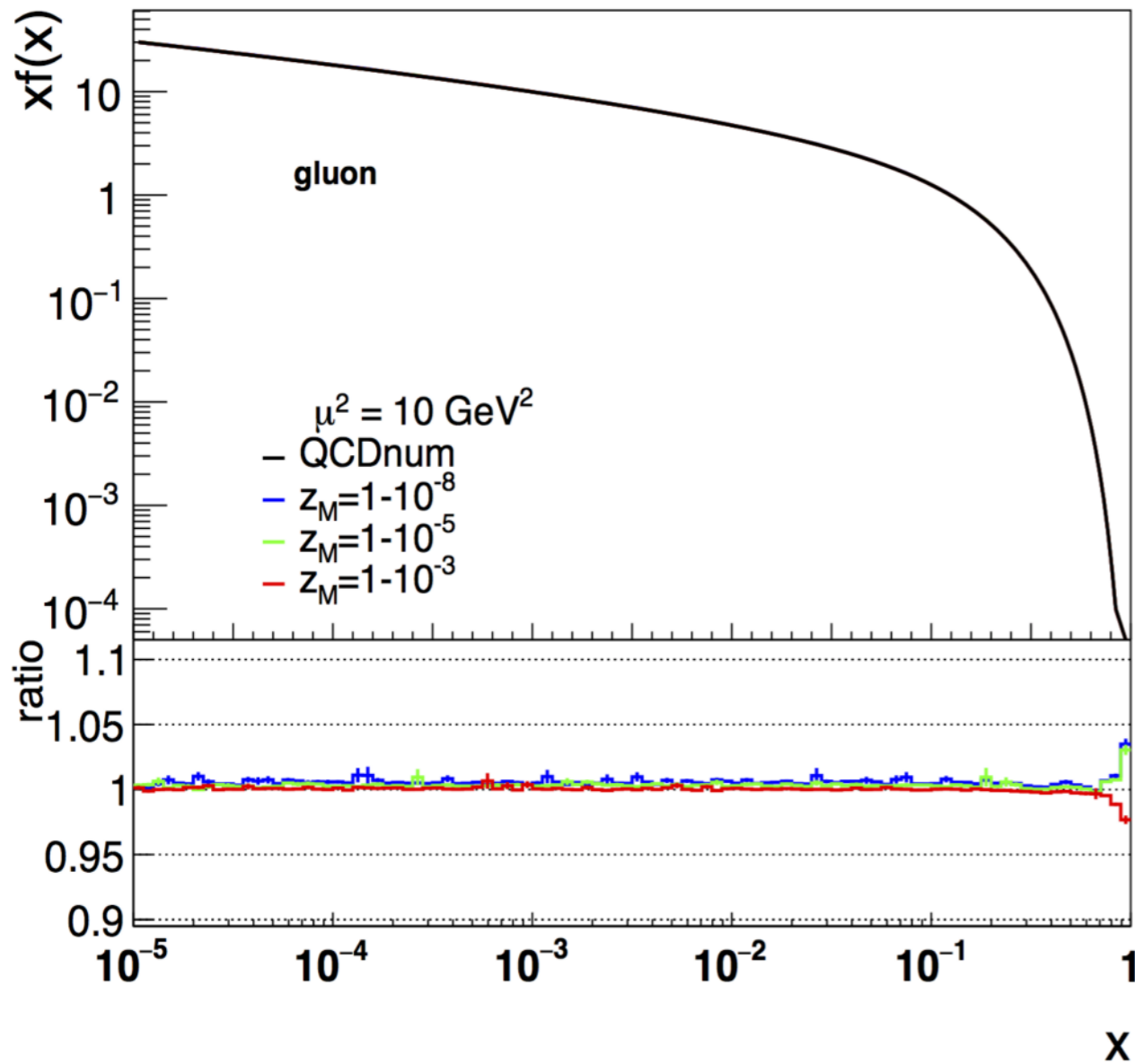
$$x f_a(x, \mu^2) = \Delta_a(\mu^2) x f_a(x, \mu_0^2) + \sum_b \int_{\mu_0}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \int_x^{z_M} dz P_{ab}^{(R)}(\alpha_s, z) \frac{x}{z} f_b\left(\frac{x}{z}, \mu'^2\right)$$

- with  $P_{ab}^{(R)}(\alpha_s(t'), z)$  real emission probability (without virtual terms)
  - $z_M$  introduced to separate real from virtual and non-emission probability
  - reproduces DGLAP up to  $\mathcal{O}(1 - z_M)$
- make use of momentum sum rule to treat virtual corrections
  - use Sudakov form factor to treat non-resolvable and virtual corrections

$$\Delta_a(z_M, \mu^2, \mu_0^2) = \exp \left( - \sum_b \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \int_0^{z_M} dz z P_{ba}^{(R)}(\alpha_s), z \right)$$



# Validation of method with QCDnum at **NLO**

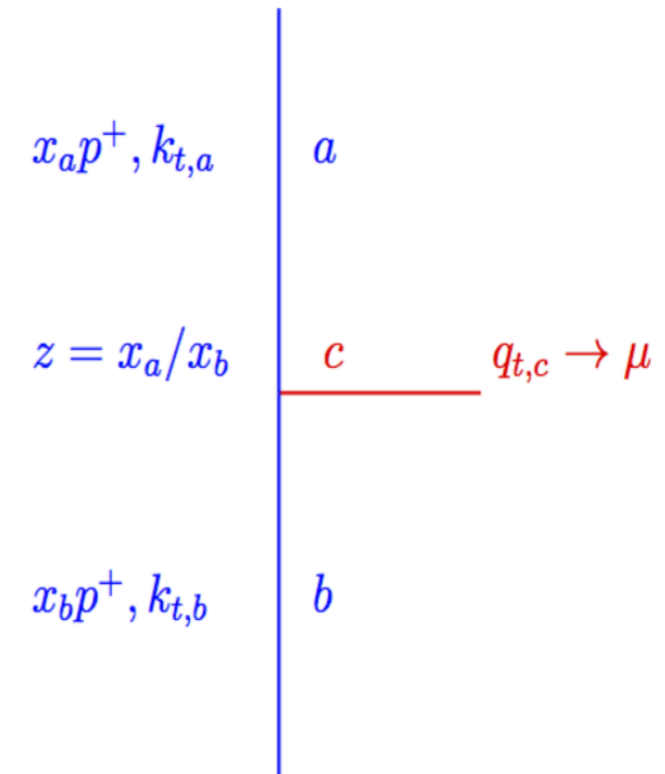


- Very good agreement with **NLO** - QCDnum if  $z_M$  is large enough:
  - approximation is of  $\mathcal{O}(1 - z_M)$

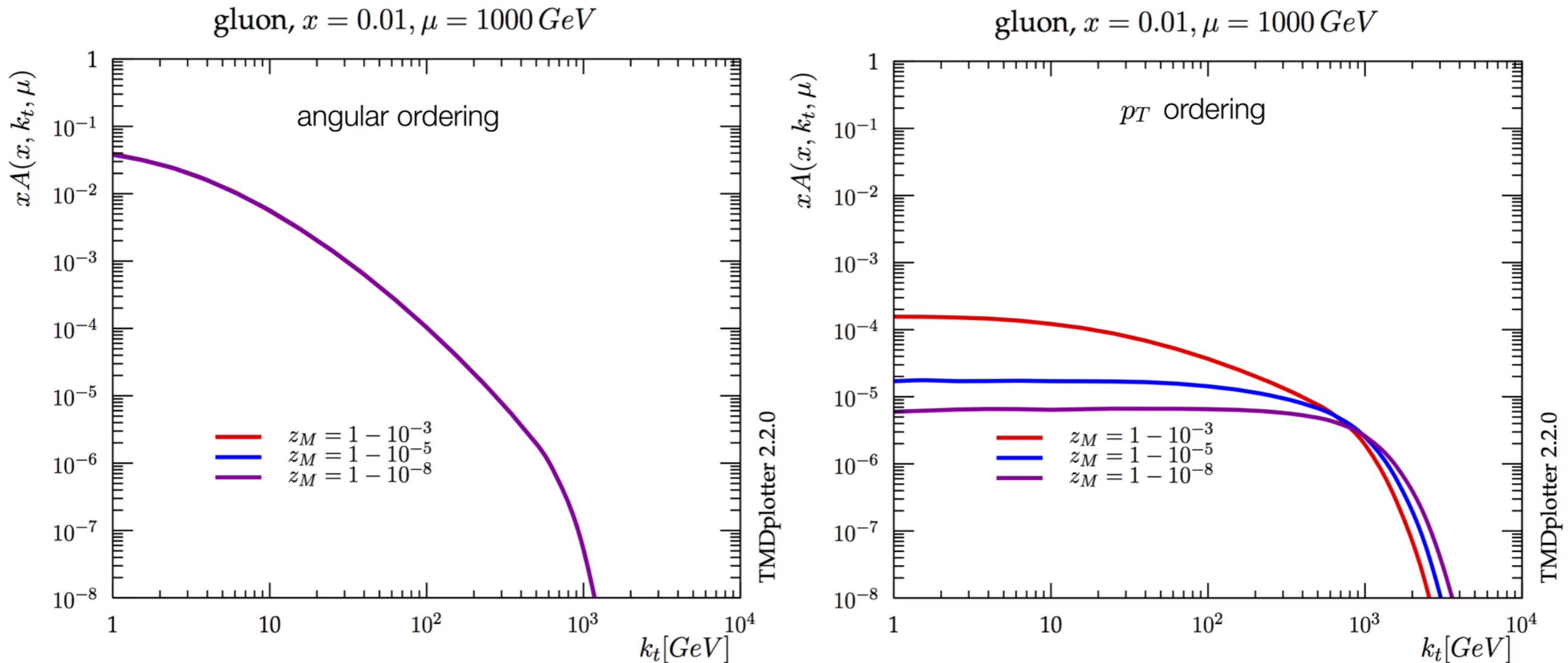
# Transverse Momentum Dependence

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- Parton Branching evolution generates every single branching:
  - kinematics can be calculated at every step
- Give physics interpretation of evolution scale:
  - in high energy limit:  $p_T$  -ordering:
 
$$\mu = q_T$$
  - angular ordering:
 
$$\mu = q_T / (1-z)$$



# Transverse Momentum: dependence on $z_M$



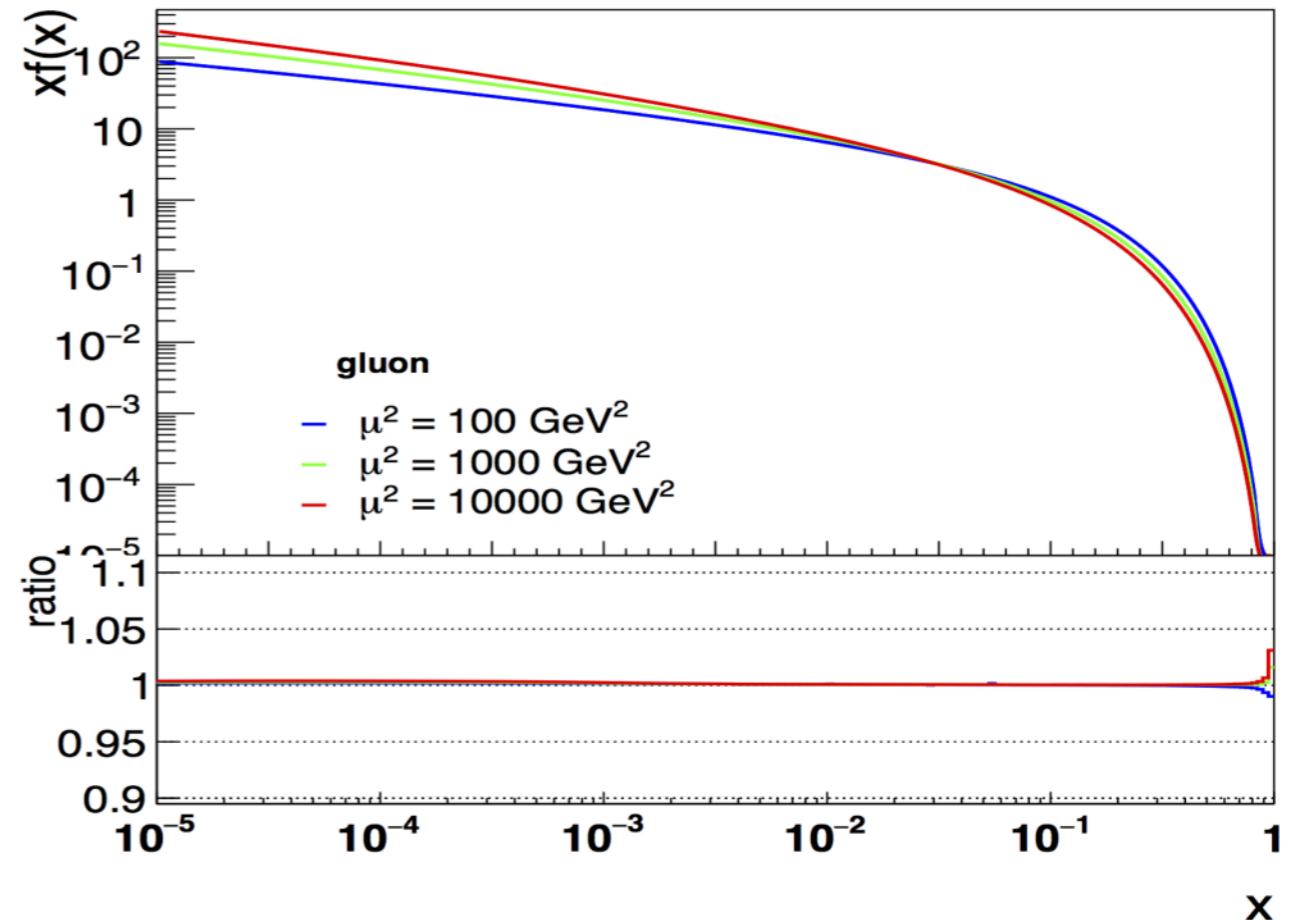
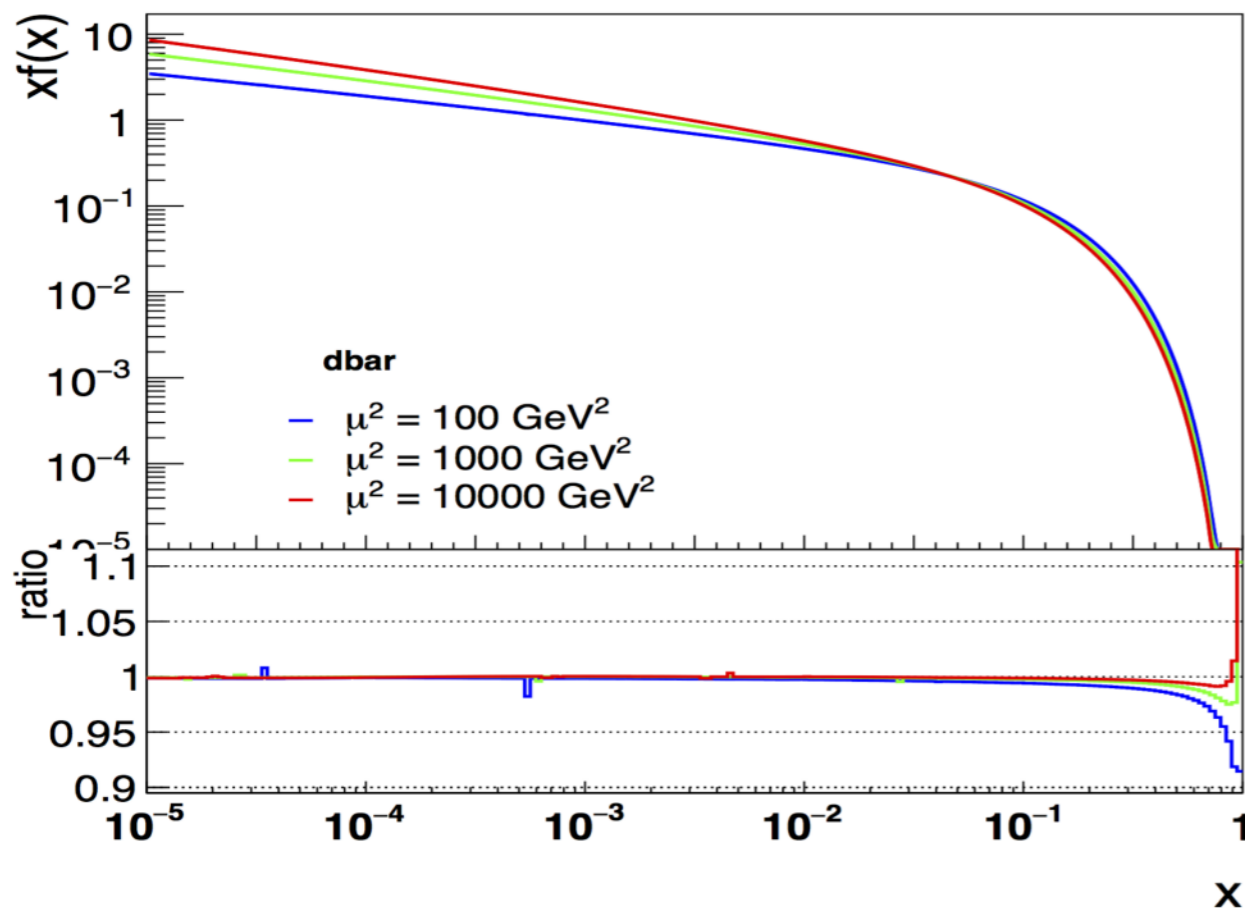
- $p_T$  – ordering ( $\mu = q_T$ ) shows significant dependence on  $z_M$ : unstable result because of soft gluon contribution
- angular ordering ( $\mu = q_T/(1-z)$ ) is independent of  $z_M$ : stable results since soft gluons are suppressed (angular ordering)

# Parton branching method in xFitter

- Convolution of kernel with starting distribution

$$\begin{aligned}
 x f_a(x, \mu^2) &= x \int dx' \int dx'' \mathcal{A}_{0,b}(x') \tilde{\mathcal{A}}_a^b(x'', \mu^2) \delta(x' x'' - x) \\
 &= \int dx' \mathcal{A}_{0,b}(x') \cdot \frac{x}{x'} \tilde{\mathcal{A}}_a^b\left(\frac{x}{x'}, \mu^2\right)
 \end{aligned}$$

- kernel defined on grid (for integrated and TMD distribution)
- validation of method:



# Advantages of parton branching method

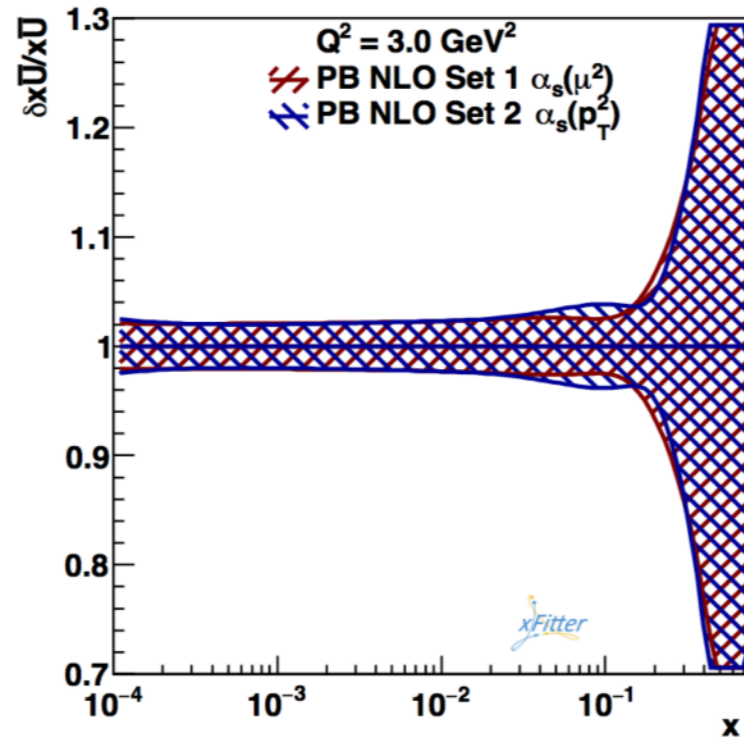
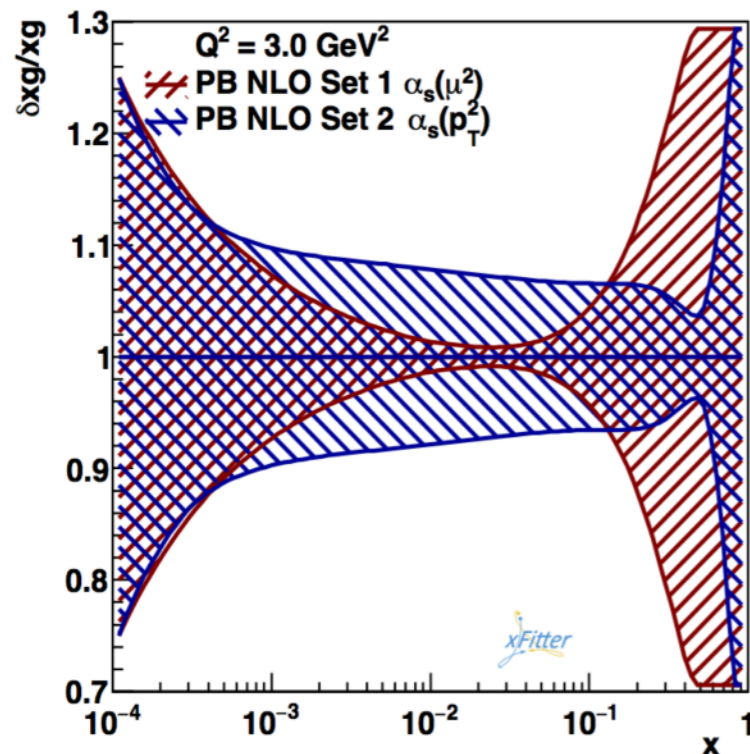
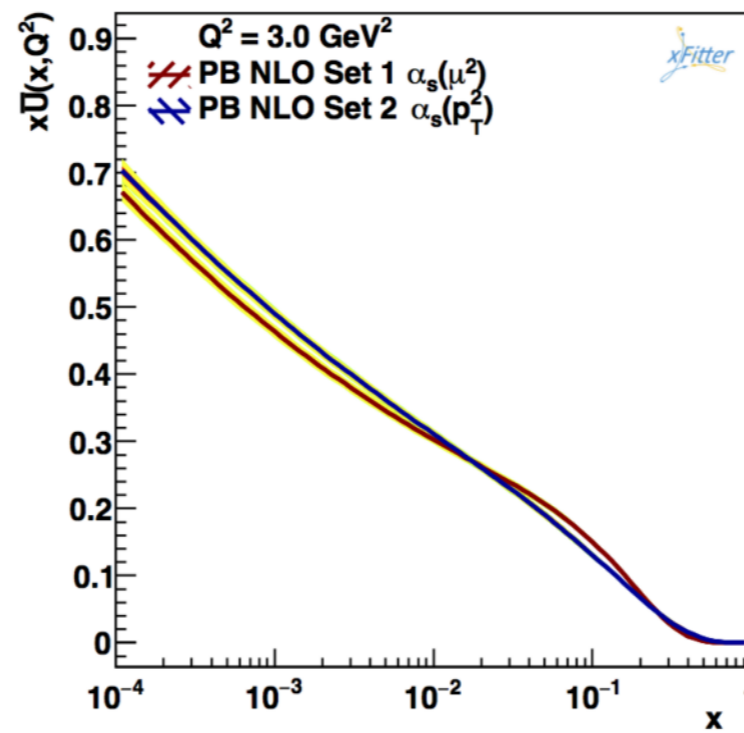
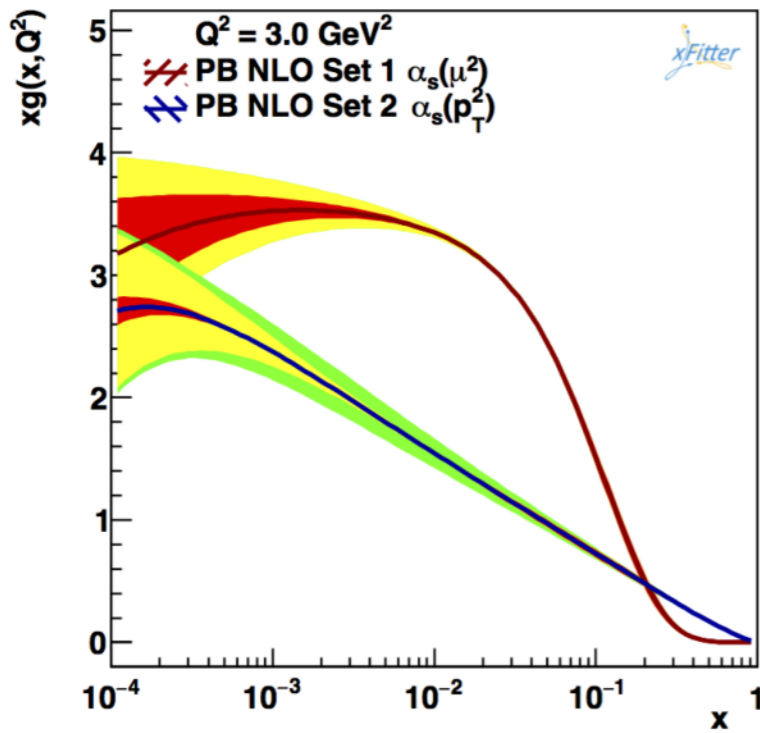
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- DGLAP equation:

$$\mu^2 \frac{\partial}{\partial \mu^2} f(x, \mu^2) = \int \frac{dz}{z} \frac{\alpha_s(\mu_r)}{2\pi} P_+(z) f\left(\frac{x}{z}, \mu^2\right)$$

- Advantages of parton branching method for collinear PDFs:
  - access to all kinematic variables and combinations between them
    - full freedom of choosing:
      - renormalisation scale:  $\alpha_s(\mu_r)$
      - evolution scale:  $\mu_f$
  - studies of different ordering conditions possible for the first time
    - angular ordering with  $\alpha_s(q)$
    - but angular ordering suggests that renormalization scale is  $p_T$  and not angle
      - angular ordering with  $\alpha_s(p_T) \rightarrow \alpha_s(q(1-z))$
      - repeat fits with changed renormalisation scale in pdf (but not yet in coefficient fct)

# Fit with changed $\alpha_s(p_T)$ : at small $Q^2$

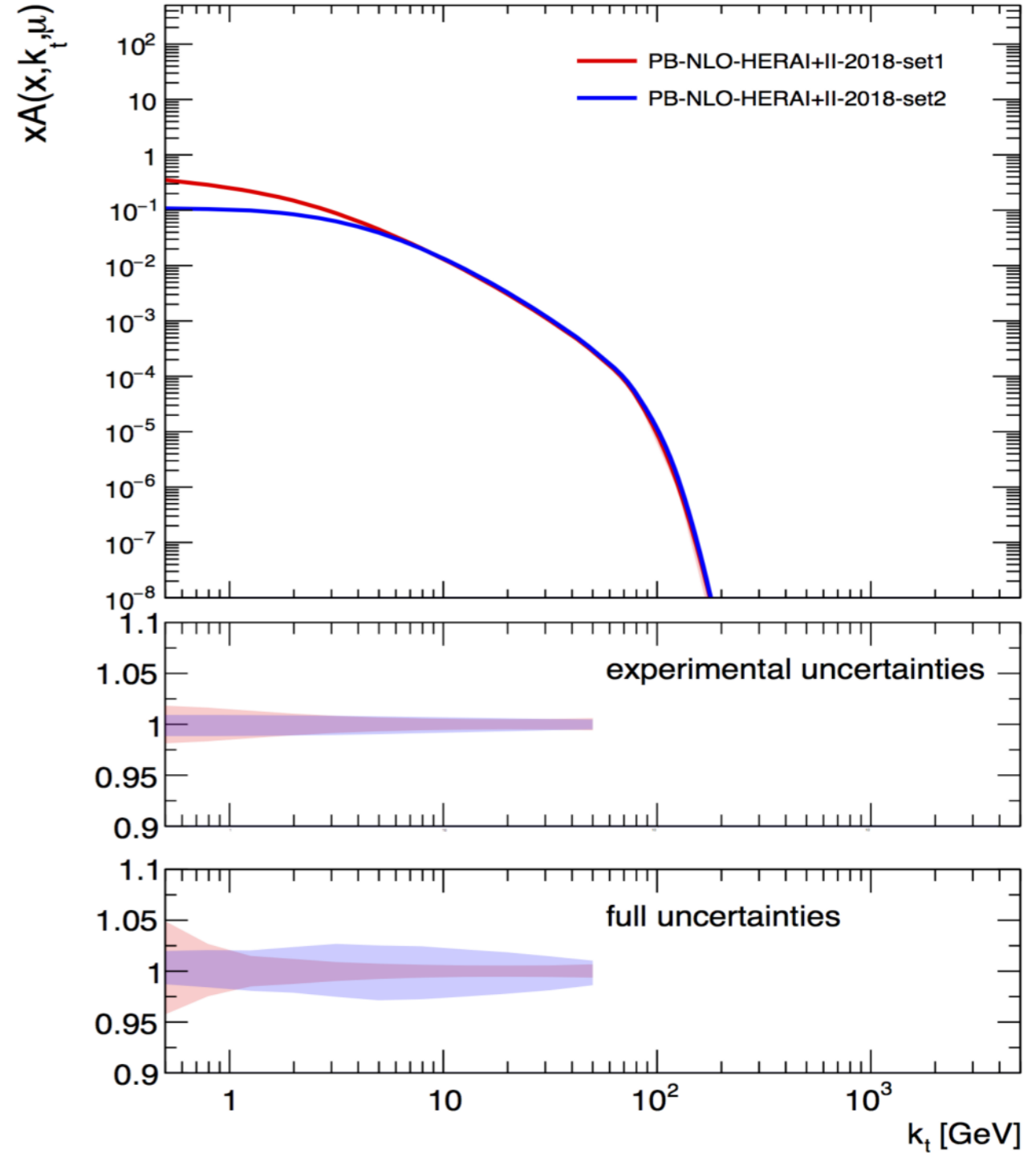
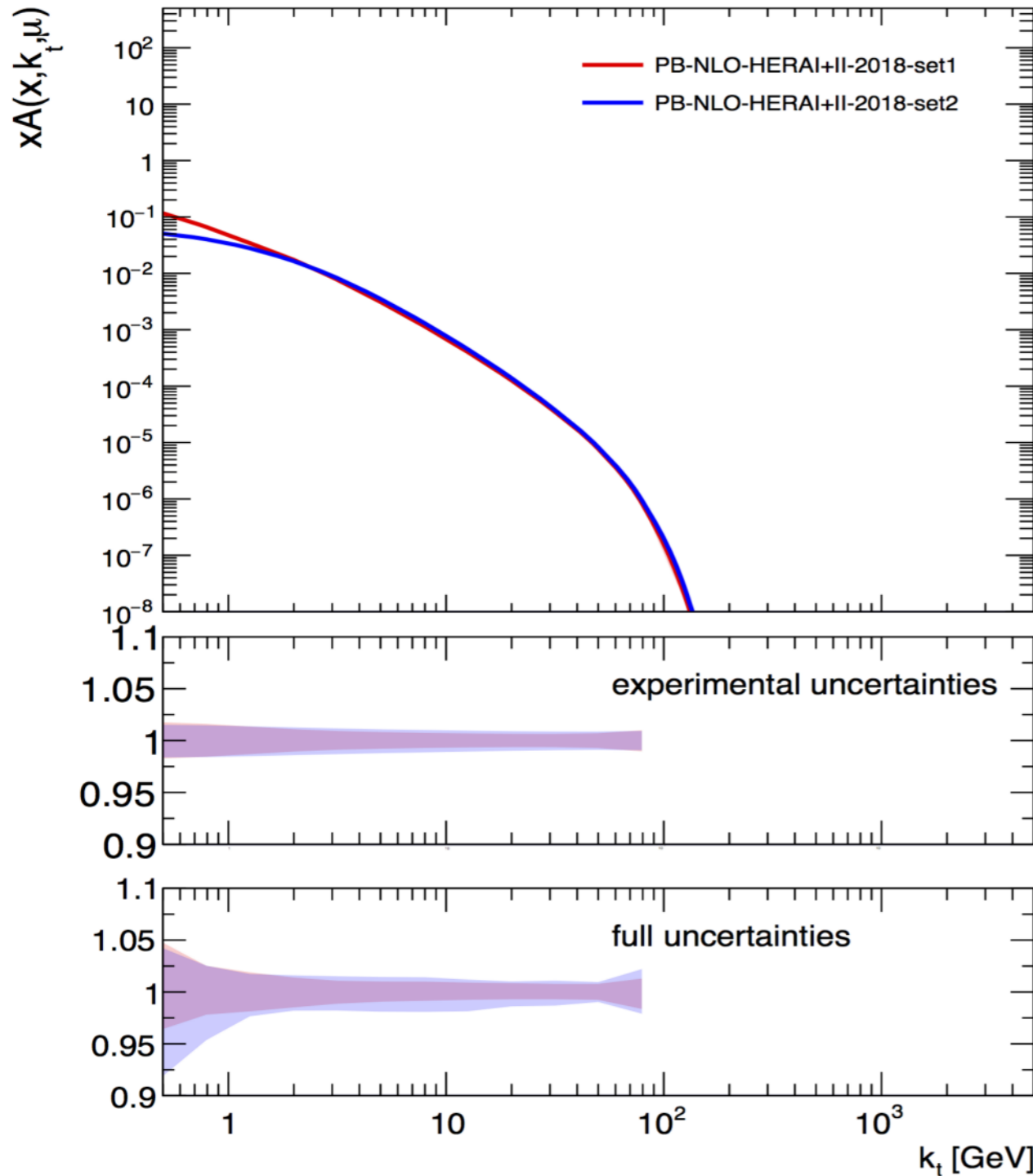


- fit 1 with  $\alpha_s(q)$ 
  - as good as HERAPDF2.0  
 $\chi^2/ndf = 1.2$
- fit 2 with  $\alpha_s(q(1-z))$ 
  - $\chi^2/ndf = 1.21$
- very different gluon distribution obtained at small  $Q^2$

# TMD distributions

anti-up,  $x = 0.01$ ,  $\mu = 100$  GeV

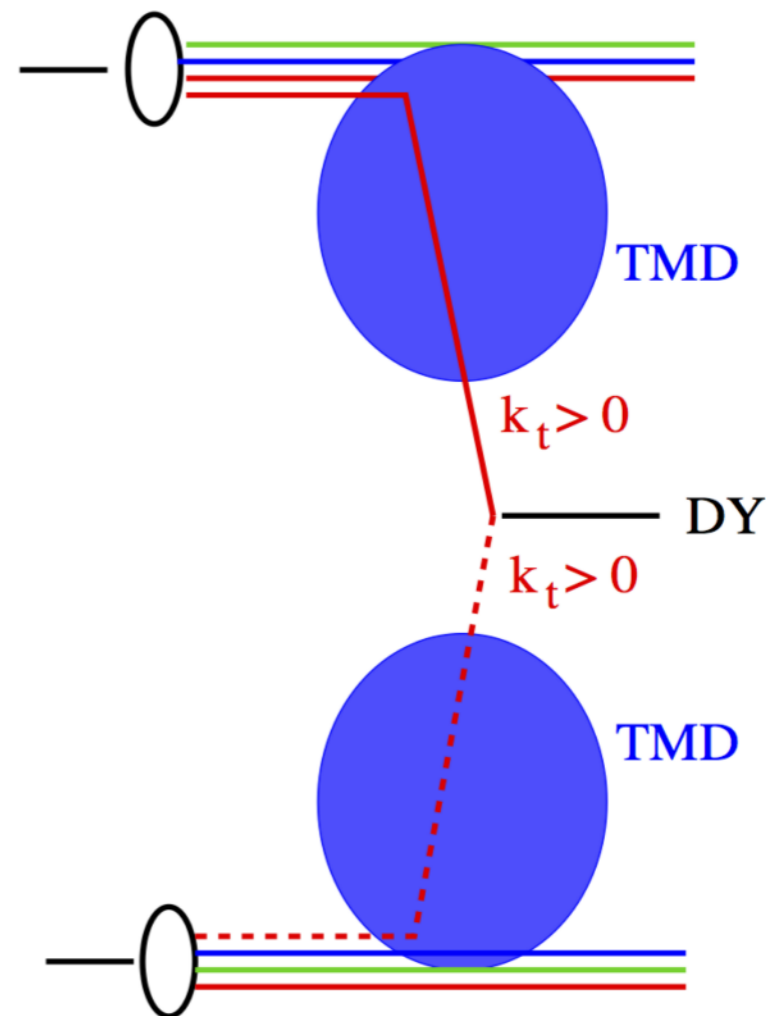
gluon,  $x = 0.01$ ,  $\mu = 100$  GeV



- model dependence larger than experimental uncertainties

# Application to DY $q_T$ - spectrum

- Use LO DY production
  - $q\bar{q} \rightarrow Z_0$
  - add  $k_t$  for each parton as function of  $x$  and  $\mu$  according to TMD
- keep final state mass fixed:
  - $x_1$  and  $x_2$  (light-cone fraction) are different after adding  $k_t$





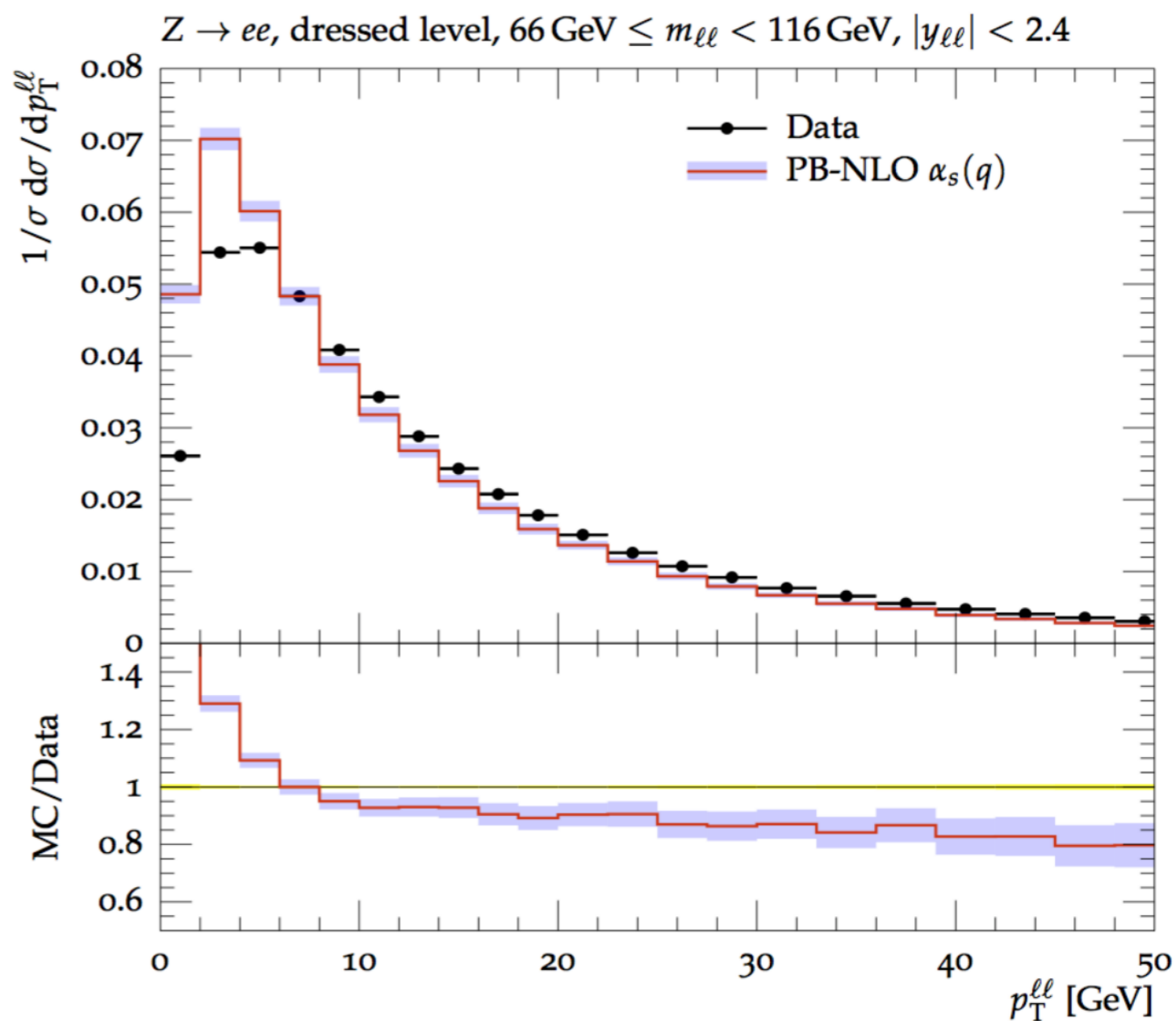
# Application to DY $q_T$ - spectrum

ATLAS Collaboration Eur. Phys. J. C76 (2016), 291  
[arXiv:1512.02192]

- Use LO DY production

$$q\bar{q} \rightarrow Z_0$$

- TMD with angular ordering including  $\alpha_s(q)$



# Application to DY $q_T$ - spectrum

ATLAS Collaboration Eur. Phys. J. C76 (2016), 291  
[arXiv:1512.02192]

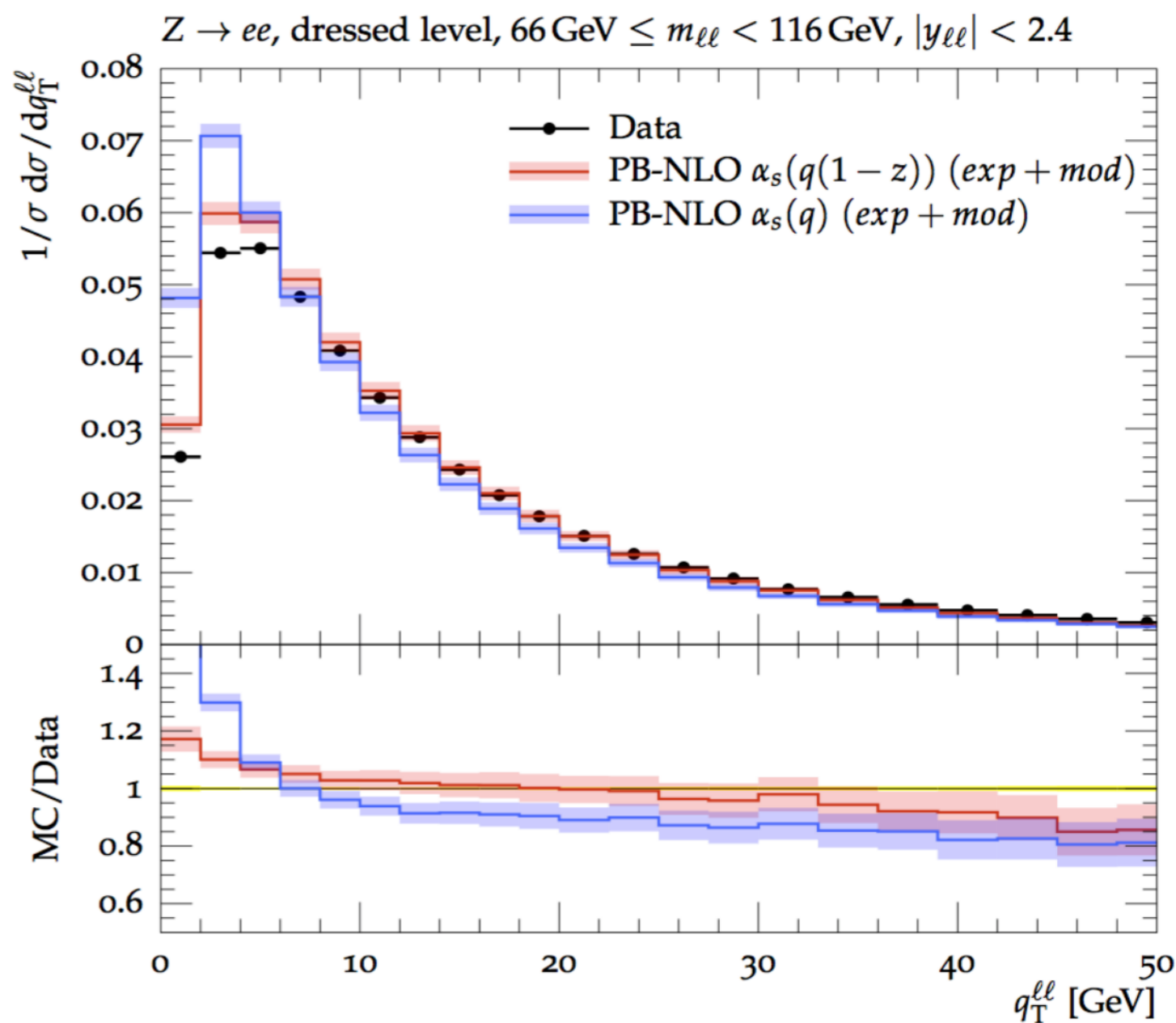
- Use LO DY production

$$q\bar{q} \rightarrow Z_0$$

- TMD with angular ordering including  $\alpha_s(q)$
- TMD with angular ordering including  $\alpha_s(p_T)$ 
  - in low  $p_T$  much better !

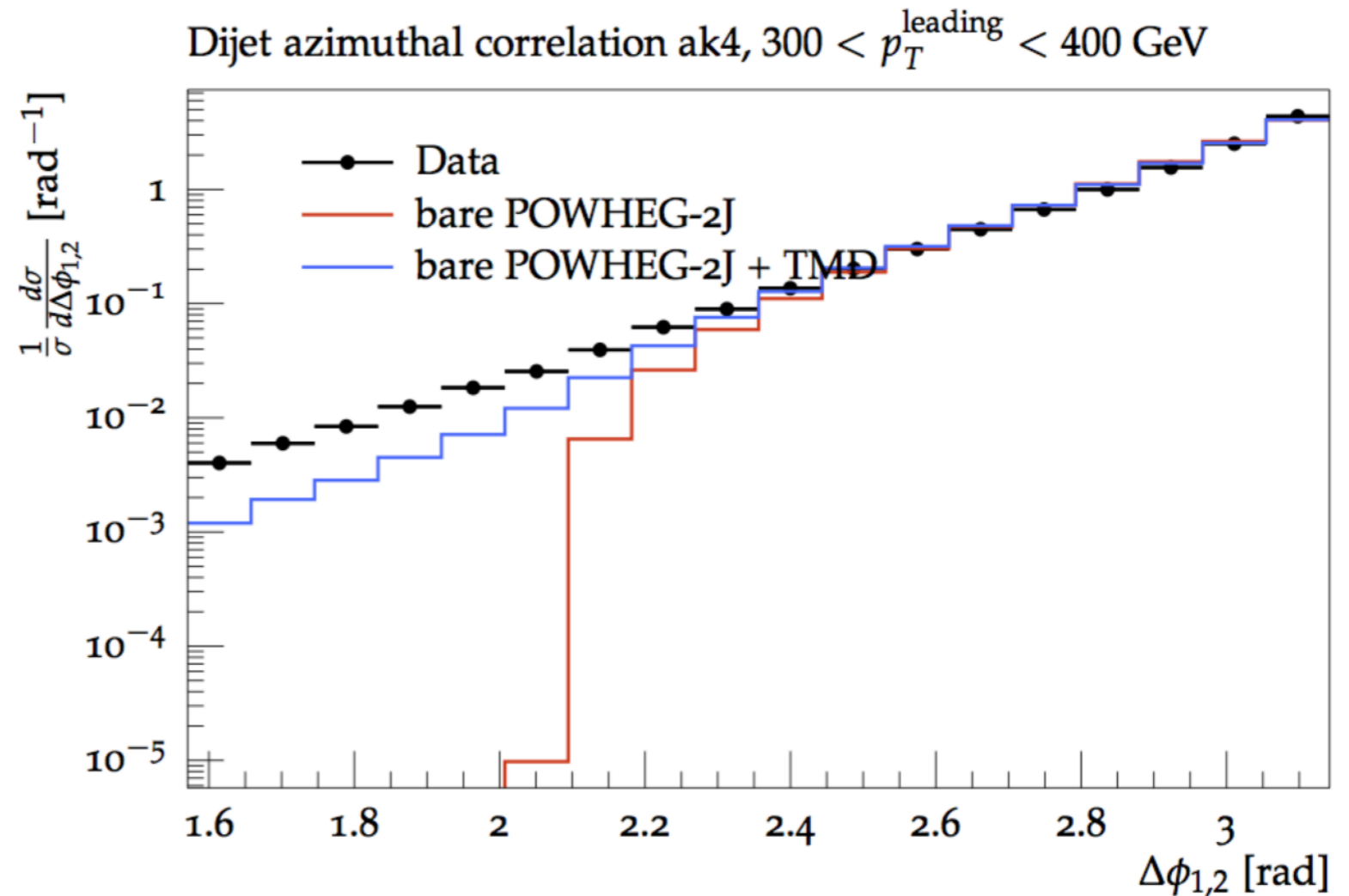
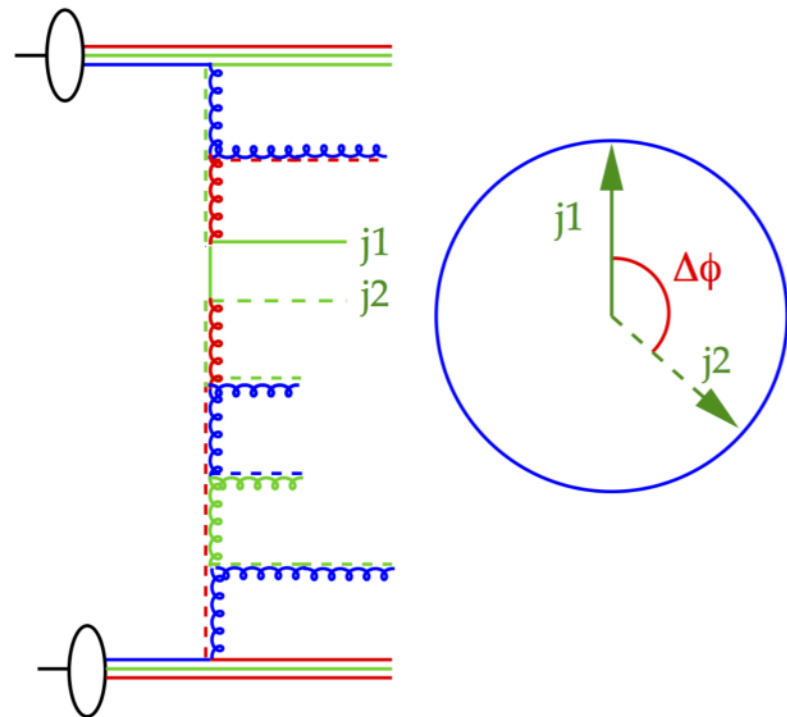
- Additional issues:

- resolvable branching
- freeze  $\alpha_s$
- intrinsic  $k_T$



# Application to high $p_T$ dijets in pp

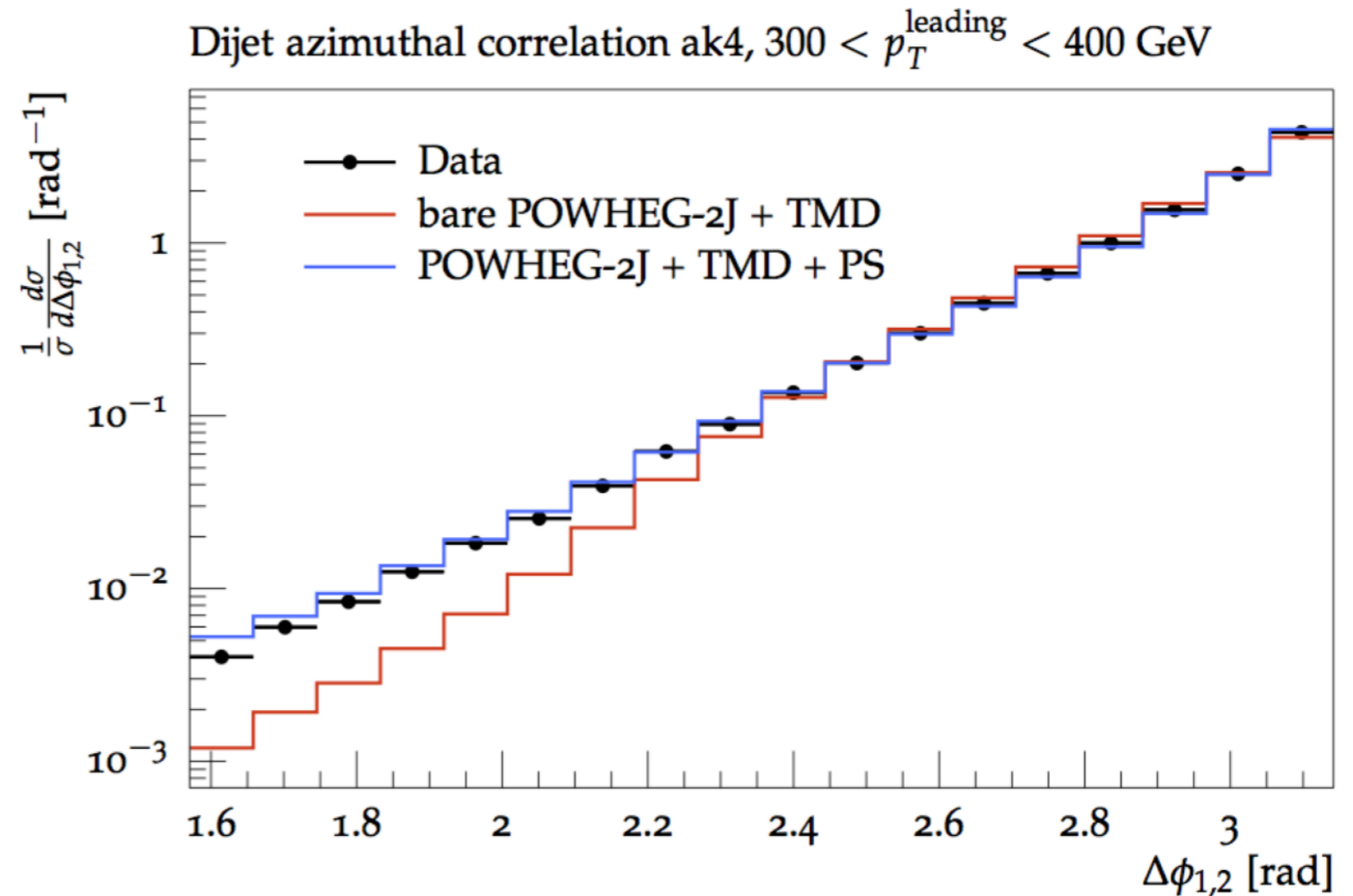
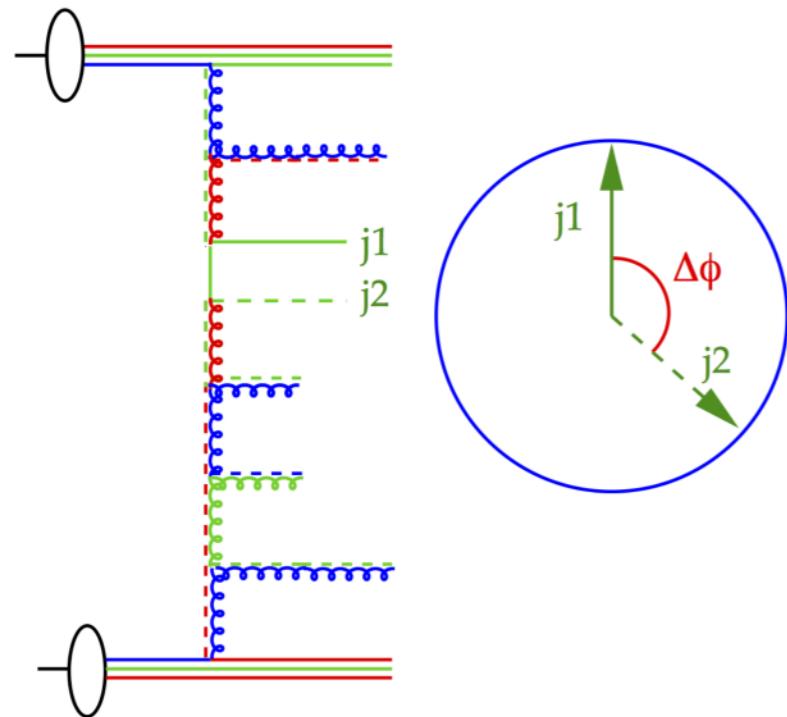
- Dijet production at in pp, a test for TMDs and PS :



- TMDs with NLO dijets get closer to data !

# Application to high $p_T$ dijets in pp

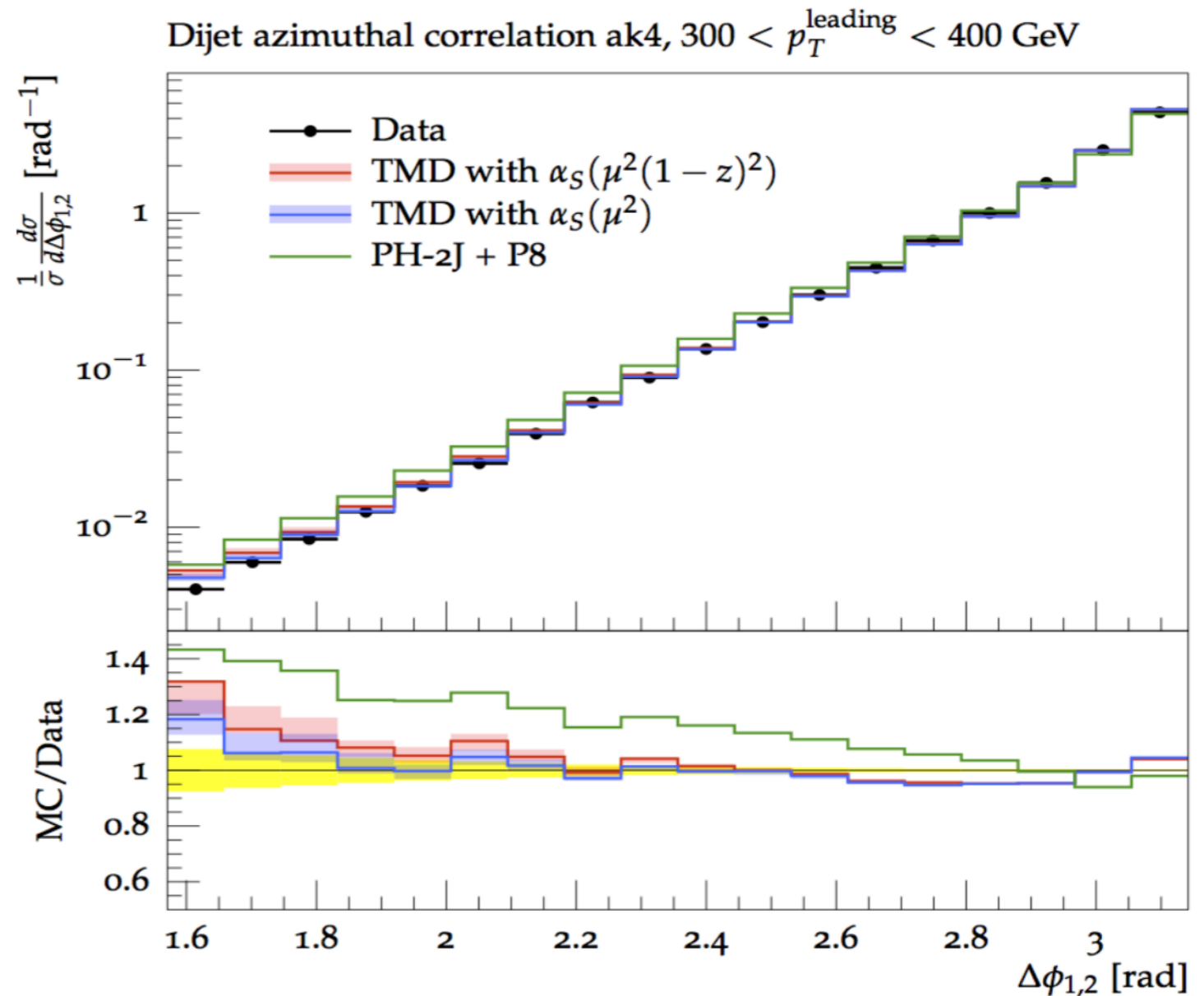
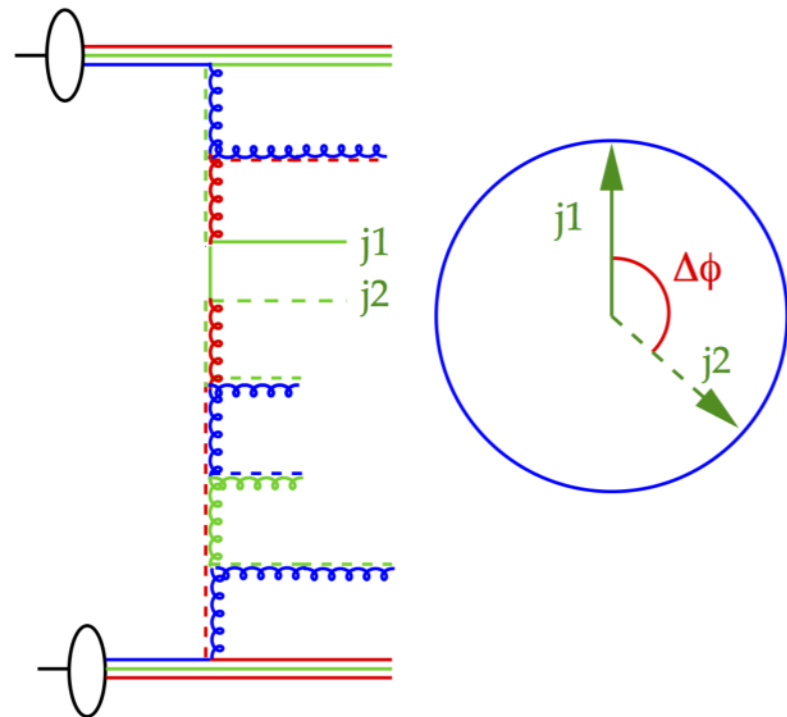
- Dijet production at in pp, a test for TMDs and PS :



- TMDs with NLO dijets + parton shower (following TMD) describes data!

# Application to high $p_T$ dijets in pp

- Dijet production at in pp, a test for TMDs and PS :



- TMDs with NLO dijets + parton shower (following TMD) describes data!
  - different TMD sets are very similar
- TMD + NLO dijets + PS  $\rightarrow$  better than conventional treatment !

# Conclusion

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- Parton Branching method to solve DGLAP equation at LO, NLO and NNLO
  - ➔ consistence for collinear (integrated) PDFs shown
  - ➔ advantages of Parton Branching method !
- method directly applicable to determine  $k_t$  distribution (as would be done in PS)
  - ➔ TMD distributions for all flavors determined at LO and NLO, without free parameters
  - ➔ TMD evolution implemented in xFitter – fits to DIS processes at the moment
- Application for pp, ep processes, like DY, jets:
  - ➔ DY  $q_T$  - spectrum without new parameters
  - ➔ TMD initial parton shower:
    - ➔ backward evolution following exactly the TMD density
    - ➔ dijet  $\Delta\phi$  very well described with NLO dijets + TMD + TMD shower

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# Appendix

# Where to find TMDs ? TMDlib and TMDplotter

- TMDlib proposed in 2014 as part of REF workshop and developed since
- combine and collect different ansaetze and approaches:

<http://tmd.hepforge.org/> and  
<http://tmdplotter.desy.de>

- TMDlib: a library of parametrization of different TMDs and uPDFs (similar to LHAPdf)

TMDlib and TMDplotter: library and plotting tools for transverse-momentum-dependent parton distributions, *F. Hautmann et al. arXiv 1408.3015, Eur. Phys. J., C 74(12):3220, 2014.*

- Also integrated pdfs (including photon pdf are available via LHAPDF)

- Feedback and comments from community is needed – just use it !

**Integrated PDF plotter**

Home TMD Plotter Publications HEP Links

**Parameters**

$p^2 = 25$  GeV<sup>2</sup>

$y_{\min} = 1.0E-5$   $y_{\max} = 100$

$x_{\min} = 1.0E-5$   $x_{\max} = 1$

**PDFs**

1. gluon ccfm-JH-2013-set1 x 1
2. gluon NNPDF23\_lo\_as\_0130\_qed x 1
3. photon NNPDF23\_lo\_as\_0130\_qed x 1
4. gluon MRST2004qed\_proton x 1

**Output**

Format: ps

display ratio

display command line

Plot Restore Add PDF field

$p^2 = 25$  GeV<sup>2</sup>

$x f(x, p^2)$

10<sup>2</sup> 10 1 10<sup>-1</sup> 10<sup>-2</sup> 10<sup>-3</sup> 10<sup>-4</sup> 10<sup>-5</sup>

10<sup>-5</sup> 10<sup>-4</sup> 10<sup>-3</sup> 10<sup>-2</sup> 10<sup>-1</sup> 1

x

Contact Imprint

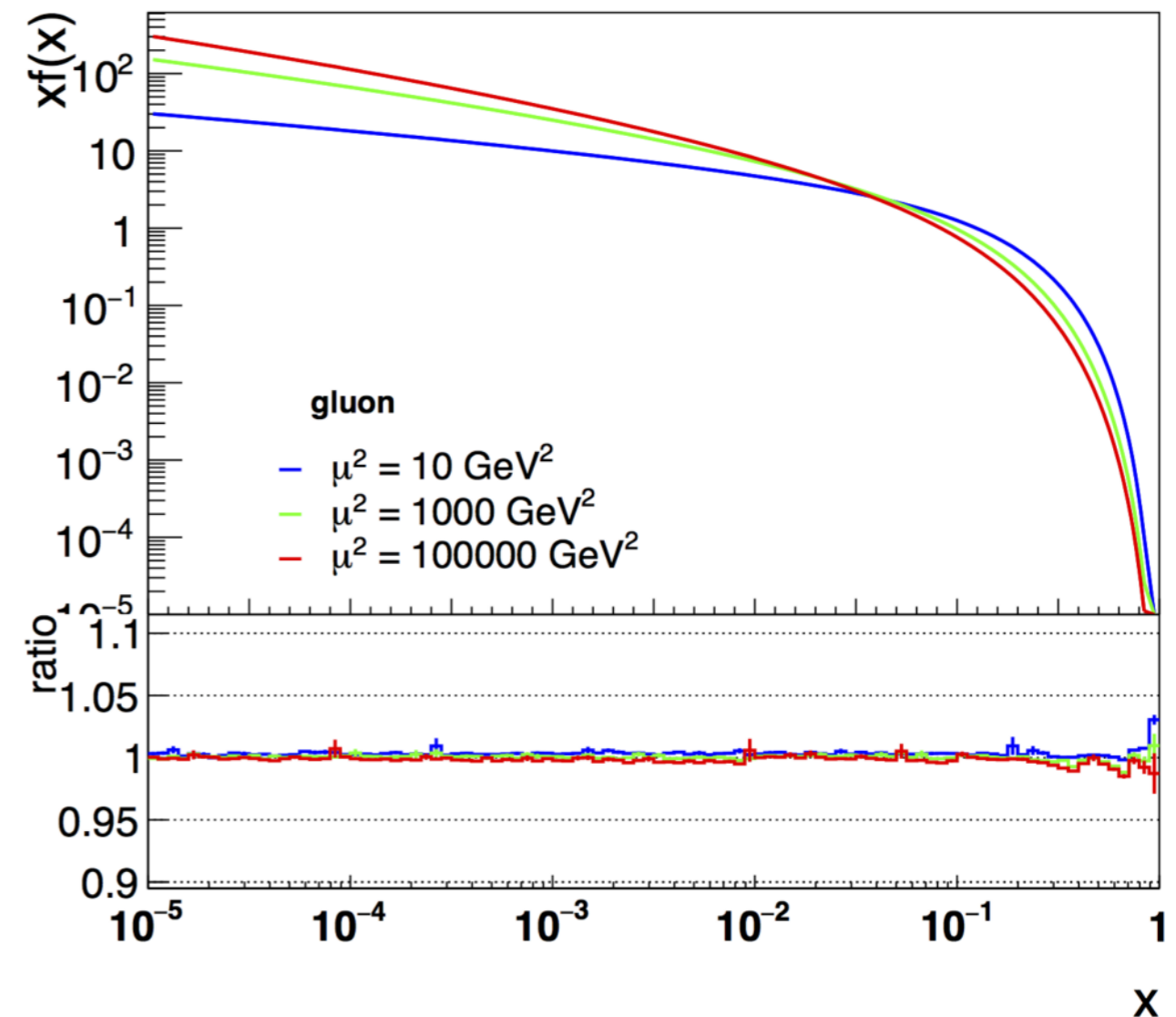
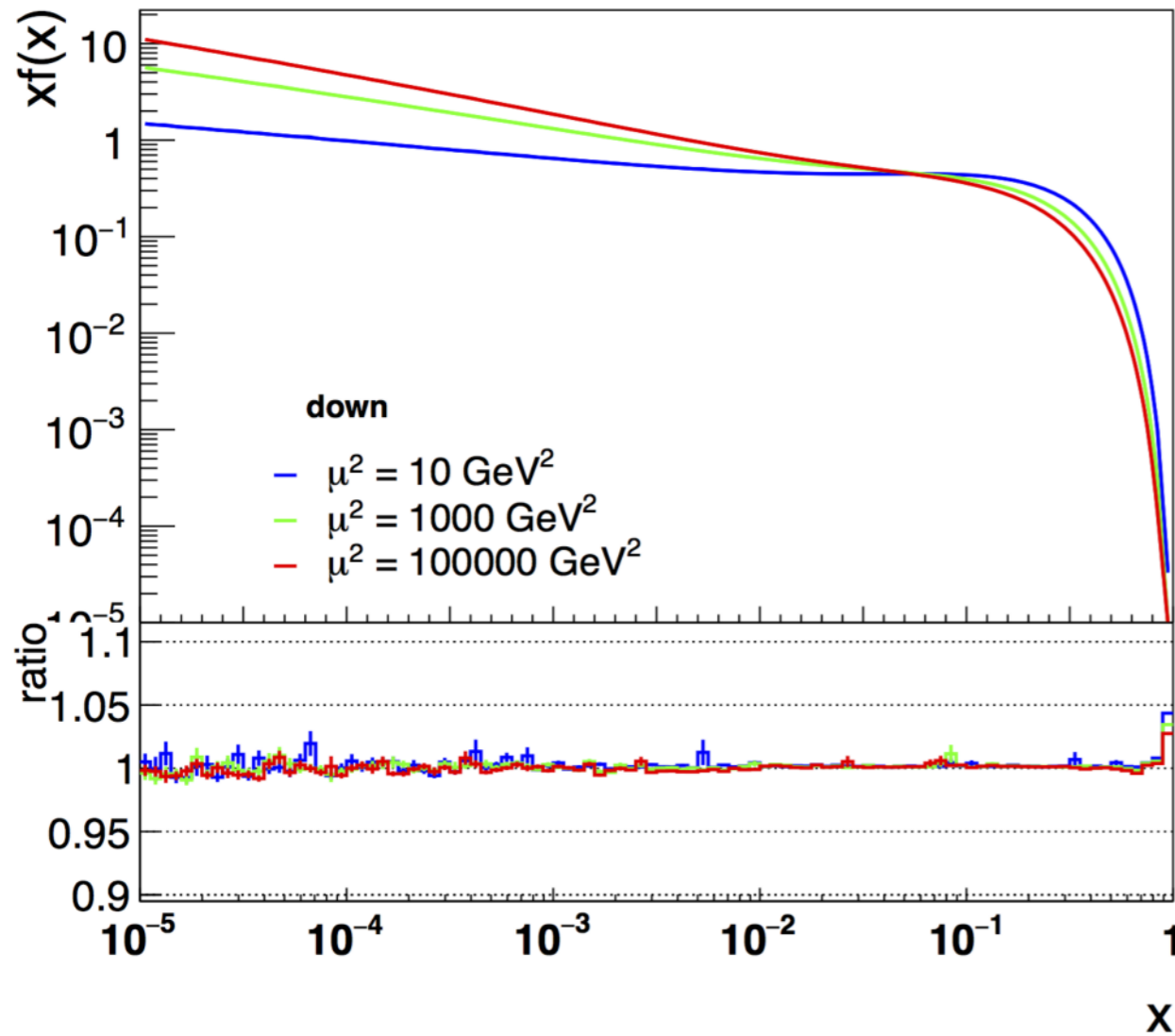
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LHAPDF 6.1.4 and TMDlib 1.0.6

PHYSICS AT THE TERA SCALE Helmholtz Alliance

DESY



# Validation of method with QCDnum at **NLO**



- Very good agreement with **NLO** - QCDnum over all  $x$  and  $\mu^2$ 
  - the same approach work also at NNLO !

# MCEG: TMDs, parton shower

- basic elements are:
  - **Matrix Elements:**
    - on shell/off shell
  - PDFs
    - TMDs
  - Parton Shower
    - following TMDs for initial state !
- Proton remnant and hadronization handled by standard hadronization program, e.g. PYTHIA
- **Parton shower with TMDs follows exactly the evolution of the TMD**
  - no (!) free parameter in shower
  - resolvable branchings and calculation of  $k_T$  defined in TMD
  - no adjustment of kinematics during/after shower

