

Transversity in inclusive DIS and novel sum rules

Alberto Accardi

Hampton U. and Jefferson Lab

DIS 2018

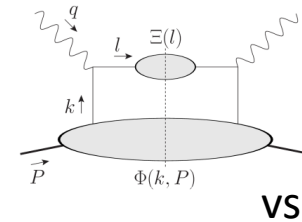
Kobe, Japan, Apr 16th-20th, 2018

*Based on: Accardi, Bacchetta, PLB 773 (2017) 632
Accardi, Signori, work in progress*

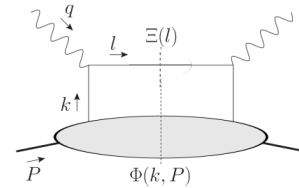
Overview

□ Inclusive DIS with jet correlators

- Quarks are not asymptotic states
- g2: new coupling to transversity...
- Burkhardt-Cottigham sum rule extended



vs.



□ Novel TMD sum rules

- New results, old ones revisited
- Single, and di-hadron FFs

$$\sum_h \int dz \, z \, \text{[diagram with two lines meeting at a point]} = \text{[diagram with a single line]}$$

□ Final thoughts

Inclusive DIS with jet correlators

Accardi, Bacchetta, PLB 773 (2017) 632

TMDs in spin 1/2 targets

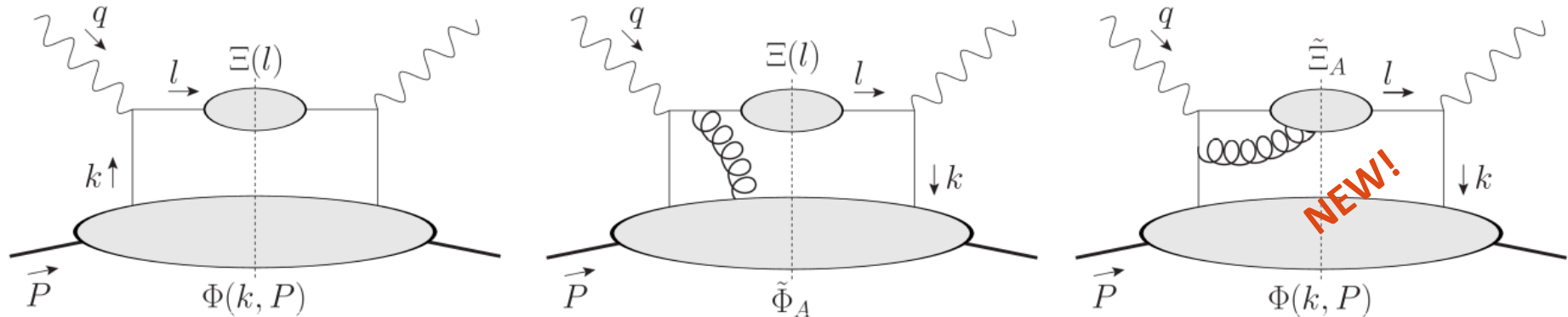
		PARTON SPIN		
TARGET SPIN	QUARKS	γ^+	$\gamma^+\gamma_5$	$\gamma^+\gamma^\alpha\gamma_5$
	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	$h_1^\perp h_{1T}^\perp$

→ H. Gao - Monday

- ❑ Integrated (collinear) correlators: only circled ones survive
- ❑ Christ-Lee theorem (1970): h_1 not observable in inclusive DIS
- ❑ Not quite true:
 - Vacuum fluctuations can flip the spin of the struck quark
 - Large contribution $\sim h_1/x$ pops up in the $g_2 - g_2^{ww}$ structure fn

Inclusive DIS with jet correlators

AA, Bacchetta, PLB 773 ('17) 632

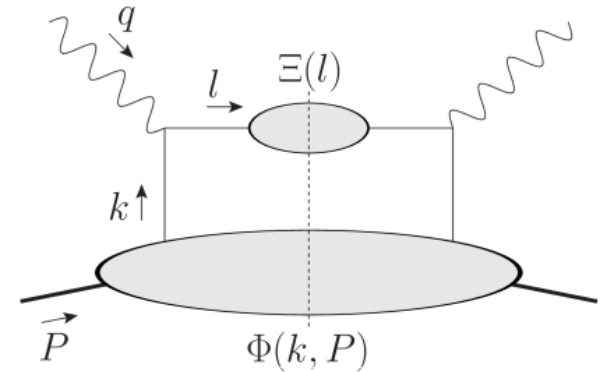


Jet correlators: \rightarrow non-asymptotic quark states

$$\begin{aligned} & \text{Diagram 1: } l \rightarrow \text{blob} \rightarrow \text{line} \quad \Xi_{ij}(l, n_+) = F.T. \langle 0 | \mathcal{U}_{(+\infty, \eta)}^{n_+} \psi_i(\eta) \bar{\psi}_j(0) \mathcal{U}_{(0, +\infty)}^{n_+} | 0 \rangle \\ & \text{Diagram 2: } l \rightarrow \text{blob} \rightarrow \text{line} \text{ with gluon } \quad (\Xi_A^\mu)_{ij} = F.T. \langle 0 | \mathcal{U}_{(+\infty, \eta)}^{n_+} g A^\mu(\eta) \psi_i(\eta) \bar{\psi}_j(0) \mathcal{U}_{(0, +\infty)}^{n_+} | 0 \rangle \end{aligned}$$

Factorization

- At order $1/Q$, neglect k^- compared to q^-
 - The cross section only depends on the **integrated jet correlator**



$$\Xi(l^-, l_T) \equiv \int \frac{dl^2}{2l^-} \Xi(l) = \frac{\Lambda}{2l^-} \xi_1 + \xi_2 \frac{\not{l}_-}{2} + \text{twist-4 terms}$$

- Coefficients interpreted in terms of quark spectral functions $J_{1,2}$:

$$\xi_1 = \int d\mu^2 \frac{\mu}{\Lambda} J_1(\mu^2) \equiv \frac{M_q}{\Lambda} \quad \leftarrow \text{“Current jet” mass} \rightarrow \text{can couple to transversity!}$$

$$\xi_2 = \int d\mu^2 J_2(\mu^2) = 1 \quad \leftarrow \text{Exactly}$$

- Positivity constraints imply

$$0 < M_q < \int d\mu^2 \mu J_2(\mu^2) \implies M_q = O(100 \text{ MeV}) \quad \text{Much larger than } m_q !$$

Jet and TMD sum rules

- Use the jet correlator sum rule: *AA, Bachetta '17*
(see also *Meissner, Metz, Pitonyak '10*)

$$\sum_h \int d^2 p_{hT} \frac{dp_h^-}{2p_h^-} \textcolor{blue}{p}_h^- \Delta^h(l, p_h) = \textcolor{blue}{l}^- \Xi(l)$$

- At TMD level, take Dirac projections:

$$\sum_h \int dz d^2 p_{hT} z D_1^h(z, p_{hT}) = \xi_2 = 1$$

$$\sum_h \int dz d^2 p_{hT} E^h(z, p_{hT}) = \xi_1 = \frac{M_q}{\Lambda}$$

$$\sum_h \int dz d^2 p_{hT} \tilde{E}^{q,h}(z, p_{hT}) = \frac{M_q - m_q}{\Lambda}$$

Novel TMD sum rules

$$M_q^{\text{pert}} = m_q \Rightarrow \text{Old ones}$$

(see later for more...)

Finally, the DIS cross section

□ Inclusive DIS

$$\frac{d\sigma}{dx_B dy d\phi_S} \propto \left\{ F_T + \varepsilon F_L + S_{\parallel} \lambda_e \sqrt{1 - \varepsilon^2} F_{LL} + |\mathbf{S}_{\perp}| \lambda_e \sqrt{2\varepsilon(1 - \varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right\}$$

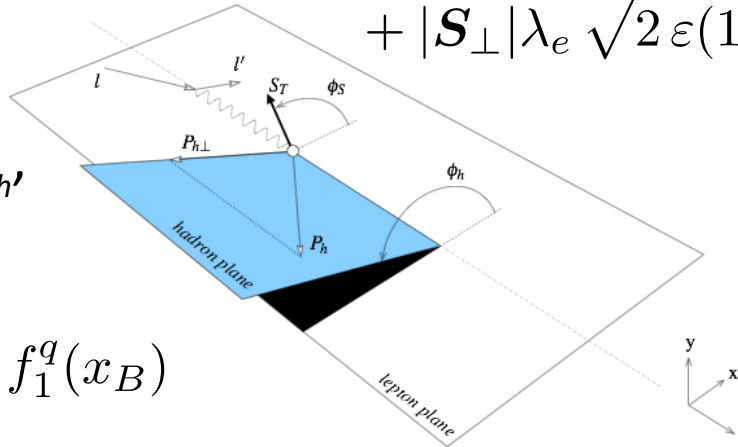
□ Integrate SIDIS over $P_{h'}$
use TMD sum rules:

$$F_T = x_B \sum_q e_q^2 f_1^q(x_B)$$

$$F_L = 0$$

$$F_{LL} = x_B \sum_q e_q^2 g_1^q(x_B)$$

$$F_{LT}^{\cos \phi_S} = -x_B \sum_q e_q^2 \frac{2M}{Q} \left(x_B g_T^q(x_B) + \frac{M_q - m_q}{M} h_1^q(x_B) \right)$$

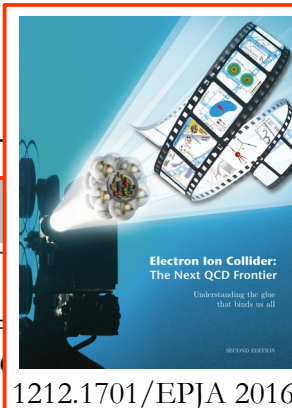


Transversity in inclusive DIS!

Finally, the DIS cross section

Inclusive DIS

$$\frac{d\sigma}{dx_B dy d\phi_S} \propto \left\{ F_T + \epsilon F_L + S_{||} \lambda_e \sqrt{1 - \epsilon} \right\}$$



Deliverables	Observables	What we learn
Sivers & unpolarized TMD quarks and gluon	SIDIS with Transverse polarization; di-hadron (di-jet)	Quantum Interference & Spin-Orbital 3D Imaging of quark's motion: valence + sea 3D Imaging of gluon's motion QCD dynamics in a unprecedented Q^2 (P_{hT}) range
Chiral-odd functions: Transversity; Boer-Mulders	SIDIS with Transverse polarization	3 rd basic quark PDF: valence + sea, tensor charge Novel spin-dependent hadronization effect QCD dynamics in a chiral-odd sector with a wide Q^2 (P_{hT}) coverage

Table 2.2: Science Matrix for TMD: 3D structure in transverse momentum space: (upper) the golden measurements; (lower) the silver measurements.

$$F_{LT}^{\cos \phi_S} = -x_B \sum_q e_q^2 \frac{2M}{Q} \left(x_B g_T^q(x_B) + \frac{M_q - m_q}{M} h_1^q(x_B) \right)$$

Transversity in inclusive DIS!

g₂ structure function revisited

AA, Bacchetta, PLB 773 ('17) 632

□ Using EOM, Lorentz Invariance Relations:

$$\begin{aligned} g_2(x_B) - g_2^{WW}(x_B) &\equiv g_2^{quark} \\ &= \frac{1}{2} \sum_a e_a^2 \left(\underbrace{g_2^{q,tw3}(x_B)}_{\text{Color force distribution}} + \frac{m_q}{M} \left(\frac{h_1^q}{x} \right)^* (x_B) + \boxed{\frac{M_q - m_q}{M} \frac{h_1^q(x_B)}{x_B}} \right) \\ &\equiv g_2^{jet} \end{aligned}$$

Transversity in inclusive DIS!

□ Consequences:

- h1 accessible in inclusive DIS! ↔ Potentially large signal
- new background to extraction of qGq effects

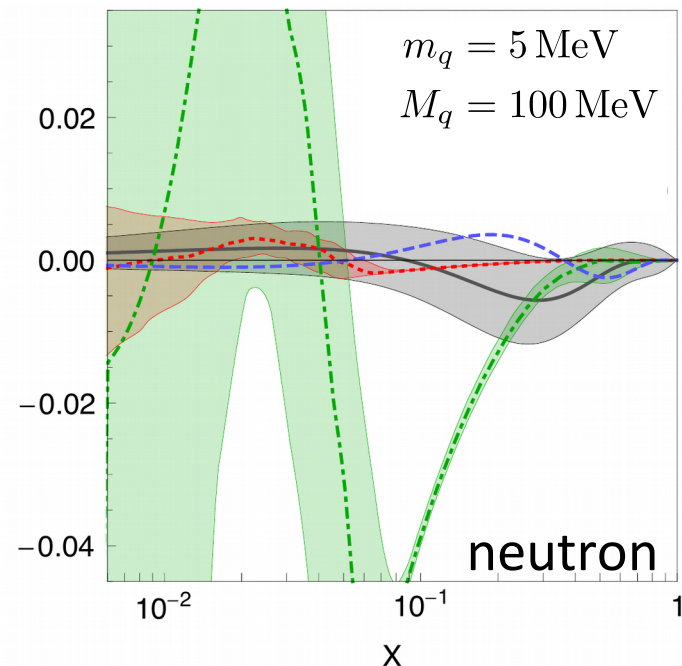
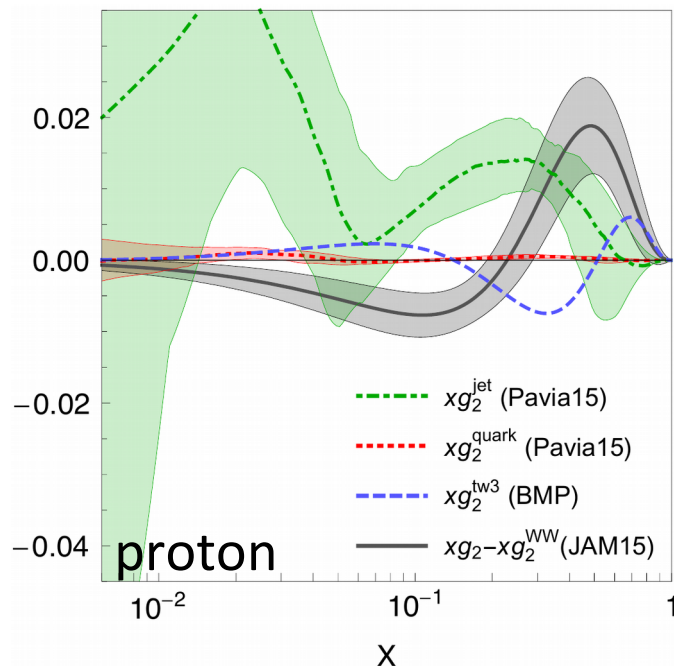
$$f^*(x) = -f(x) + \int_x^1 \frac{dy}{y} f(y)$$

g₂ structure function revisited

AA, Bacchetta, PLB 773 ('17) 632

Using EOM, Lorentz Invariance Relations:

$$\begin{aligned}
 g_2(x_B) - g_2^{WW}(x_B) &\equiv g_2^{quark} \\
 &= \frac{1}{2} \sum_a e_a^2 \left(g_2^{q,tw3}(x_B) + \frac{m_q}{M} \left(\frac{h_1^q}{x} \right)^* (x_B) + \frac{M_q - m_q}{M} \frac{h_1^q(x_B)}{x_B} \right) \equiv g_2^{jet}
 \end{aligned}$$



Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

□ Taking moments of g_2 with $M_u \approx M_d \equiv M_{jet}$

Burkardt-Cottingham

$$\int_0^1 g_2(x) = M_{\text{jet}} \int_0^1 dx \frac{h_1(x)}{x}$$

→ Broken by quark vacuum fluctuations!

→ *see also caveat in original BC paper*

Perturbatively: $h_1 \sim x$

Small-x asymptotics: → *Sievert – WG6, Thu 11:20*

$$g_1^{NS} \sim x^{\epsilon_g} \quad \epsilon_g = -\sqrt{\alpha_s N_c / \pi} \approx -0.56$$

→ *Kovchegov, Pitonyak, Sievert*
PRD(2017)93

But h_1 is also non-singlet, expect

$$h_1 \sim x^{\epsilon_h} \quad \epsilon_h = \epsilon_g < 0 \quad !!$$

– Is BC badly broken? $1/N_c$ corrections non negligible? Or ...?

Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

□ Taking moments of g_2 with $M_u \approx M_d \equiv M_{jet}$

Burkardt-Cottingham

$$\int_0^1 g_2(x) = M_{\text{jet}} \int_0^1 dx \frac{h_1(x)}{x}$$

→ Broken by quark vacuum fluctuations!

→ **How does spin propagate to small x?**

→ see also caveat in original BC paper

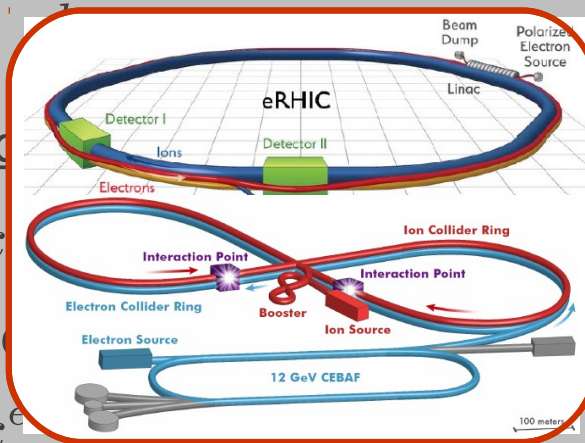
Perturbatively:

Small-x asymptotic

$$g_1^{NS} \sim x$$

But h_1 is also not

$$h_1 \sim x^\epsilon$$



$$/2\pi \approx -0.92$$

→ Kovchegov, Pitonyak, Sievert

PRD(2017)93

— Is BC badly broken? $1/N_c$ corrections non negligible? Or ...?

Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

□ Taking moments of g_2 with $M_u \approx M_d \equiv M_{jet}$

Burkardt-Cottingham

$$\int_0^1 g_2(x) = M_{\text{jet}} \int_0^1 dx \frac{h_1(x)}{x}$$

Efremov-Teryaev-Leader

$$\int_0^1 x g_2^{q-\bar{q}}(x) = 2 M_{\text{jet}} \underbrace{\int_0^1 dx h_1^{q-\bar{q}}(x)}_{\text{Tensor charge } \delta_T}$$

→ Novel way to measure the tensor charge!

Novel non-perturbative sum rules

AA, Bacchetta, PLB 773 ('17) 632

□ Taking moments of g_2 with $M_u \approx M_d \equiv M_{jet}$

Burkardt-Cottingham

$$\int_0^1 g_2(x) = M_{\text{jet}} \int_0^1 dx \frac{h_1(x)}{x}$$

Efremov-Teryaev-Leader

$$\int_0^1 x g_2^{q-\bar{q}}(x) = 2 M_{\text{jet}} \underbrace{\int_0^1 dx h_1^{q-\bar{q}}(x)}_{\text{Tensor charge } \delta_T}$$

Color polarizability

$$\int_0^1 [3x^2 g_2(x) - 2x^2 g_1(x)] = d_2 + \underbrace{3 M_{\text{jet}} \int_0^1 dx x h_1(x)}_{\text{“background”}} + O(m_q)$$

“pure twist-3” ↗

Color force
– Burkardt $\sim \langle P | \bar{\psi} \gamma^+ F^{+\alpha} \psi | P \rangle$

New momentum sum rules

Accardi, Signori, in preparation

Quark-quark TMD sum rules

□ General jet correlator sum rule: *AA, Signori '18*

$$\sum_h \int d^2 p_{hT} \frac{dp_h^-}{2p_h^-} \mathbf{p}_h^\mu \Delta^h(l, p_h) = \begin{cases} l^\mu \Xi(l) & \mu = - \quad \text{longitudinal} \\ 0 & \mu = 1, 2 \quad \text{transverse} \end{cases}$$

□ For TMDs, take suitable traces:

Longitudinal

Twist-2 $\sum_{h, S_h} \int dz \, z \, D_{1h}(z) = 1$ *Collins-Soper*

Twist-3 $\left\{ \begin{array}{l} \sum_{h, S_h} \int dz \, E_h(z, p_{hT}) = \frac{M_q}{\Lambda} \\ \sum_{h, S_h} \int dz \, H_h(z) = 0 \end{array} \right.$ *NEW!*

Transverse

$\sum_{h, S_h} \int dz \, z \, H_{1h}^{\perp(1)}(z) = 0$ *Schaefer-Teryaev*

$\sum_{h, S_h} \int dz \, D_h^{\perp(1)}(z) = 0$ *NEW!*

$\sum_{h, S_h} \int dz \, G_h^{\perp(1)}(z) = 0$ *NEW!*

$f^{(1)}(z) \equiv \int d^2 P_{hT} P_{hT}^2 f(z, P_{hT})$

Quark-gluon-quark TMD sum rules

□ Using Equation of Motion relations in q-q sum rules:

Longitudinal $\left\{ \begin{array}{l} \sum_{h,S_h} \int dz \tilde{E}_h(z) = \frac{M_q - m_q}{\Lambda} \text{ NEW!} \implies \text{Transversity in DIS!} \\ \sum_{h,S_h} \int dz \tilde{H}_h(z) = 0 \text{ NEW!} \implies \int dz z F_{UT}^{\sin \phi_S}(x, z) = 0 \end{array} \right.$

Diehl-Sapeta

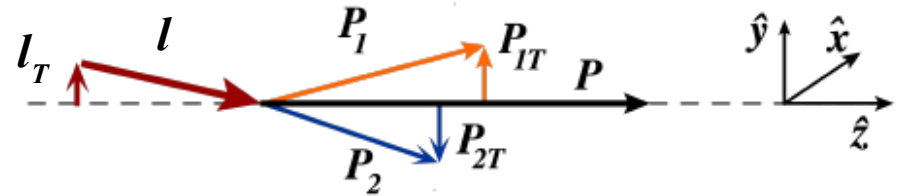
NEW: proven at correlator level

Transverse $\left\{ \begin{array}{l} \sum_{h,S_h} \int dz \tilde{D}^{\perp(1)}(z) = \frac{\langle l_T^2 \rangle}{2\Lambda^2} \text{ NEW!} \\ \sum_{h,S_h} \int dz \tilde{G}^{\perp(1)}(z) = 0 \text{ NEW!} \end{array} \right.$

* in the “hadron frame”, $P_{hT}=0$
 ($\langle p_{hT}/z \rangle$ in “parton frame”)

Di-hadron sum rules

- Two semi-inclusive hadrons
 - Di-hadron frame



→ M.Radici – WG6, Wed 10:10

$$\vec{P}_H = \vec{P}_{h1} + \vec{P}_{h2} \quad z = P_H^- / l^-$$

$$R = \frac{1}{2}(P_{h1} - P_{h2}) \quad \zeta = p_1^- / l^-$$

- Dihadron correlator:

$$\Delta^{2h}(l; P_{2h}, R) = F.T. \sum_X \langle 0 | \psi(\xi) | P_{2h}, R, X \rangle \langle P_{2h}, R, X | \bar{\psi}(\xi) | 0 \rangle$$

- Two correlators → two sets of sum rules
 - **Total P_{2h}** : analogous to previous ones
 - **Relative R** : new!

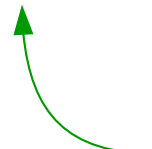
Di-hadron sum rules

□ Relative momentum sum rules:

$$\sum_{h_{1,2}, S_{1,2}} \int dz d\zeta H_1^{\triangleleft(0,1)}(z, \zeta) = 0$$

$$\sum_{h_{1,2}, S_{1,2}} \int dz d\zeta D^{\triangleleft(0,1)}(z, \zeta) = 0$$

$$\sum_{h_{1,2}, S_{1,2}} \int dz d\zeta G^{\triangleleft(0,1)}(z, \zeta) = 0$$


$$f^{(0,1)}(z) \equiv \int d^2 P_{hT} R_T^2 f(z, P_{hT}, R_T)$$

IN PROGRESS

Final thoughts

Final thoughts

□ Jet correlators open up new theory and phenomenology

- Transversity contributes to inclusive g_2
- Extended BC and ETL sum rules
 - New handle on proton tensor charge
- Open question: spin transport to small x !

□ New momentum sum rules

- Quark-quark FF (complete up to twist 3) ; q-g-q (partial at twist 3)
- DiFF : quark-quark only – work in progress

□ How to measure jet correlators?

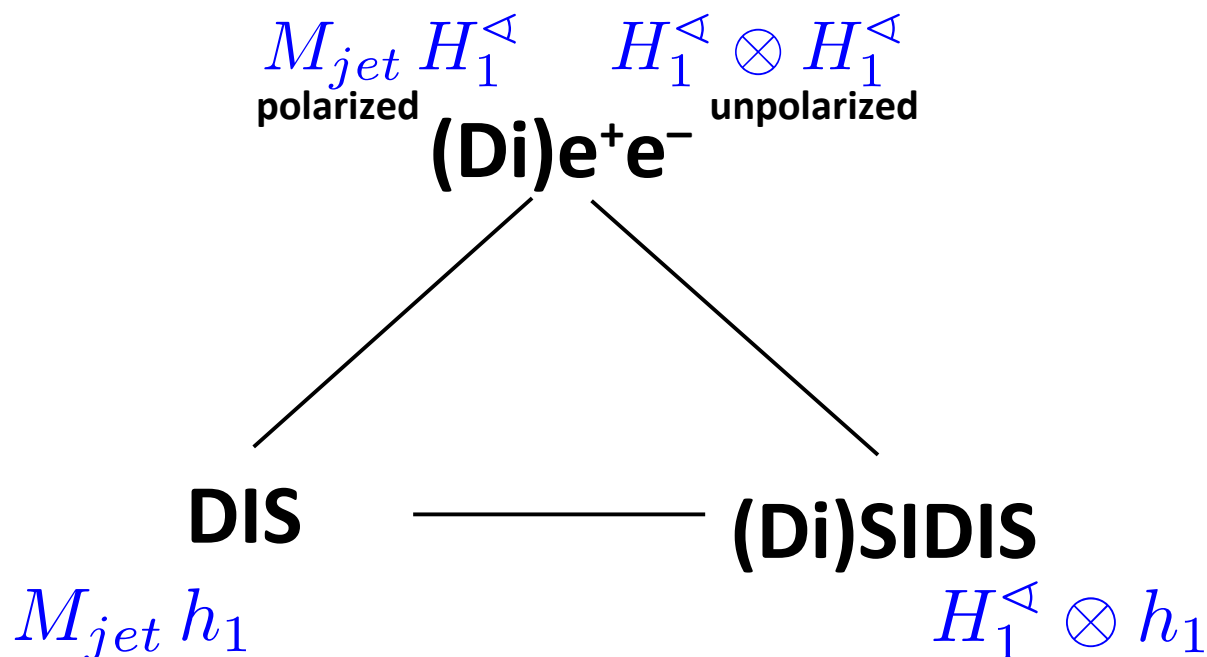
- Need a new “universal fit” of $M_{\text{jet}}, h_1, H_1^\perp$

A new “universal” fits

□ “Universal fits”

- Different PDF, FF “sectors” fitted simultaneously
- e.g. polarized PDFs, and FFs → *N.Sato – WG1, Thu 10:20*

□ Chiral-odd collinear case:



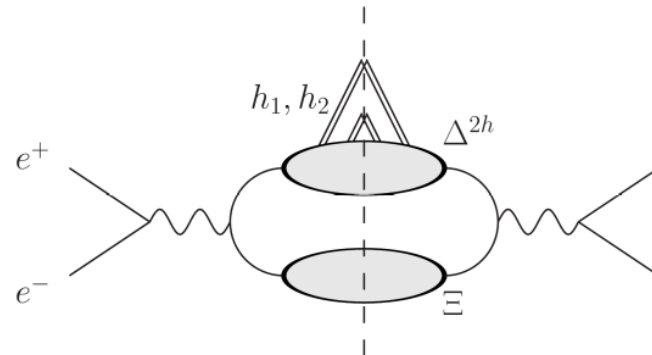
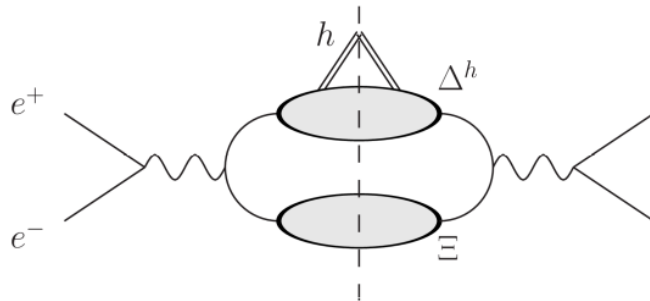
Thank you!

Measuring the jet correlator

Accardi, Bacchetta, Radici, in prep.

Related to confinement, mass generation [*C.Roberts*]

□ Jet mass M_{jet} can be measured in polarized $e^+ + e^-$:



— Needs **LT asymmetry** in semi-inclusive **Lambda** production

$$\frac{d\sigma^L(e^+e^- \rightarrow \Lambda X)}{d\Omega dz} = \frac{3\alpha^2}{Q^2} \lambda_e \sum_a e_a^2 \left\{ \frac{C(y)}{2} \lambda_h G_1 + D(y) |S_T| \cos(\phi_S) \frac{2M_h}{Q} \left(\frac{G_T}{z} + \frac{M_q - m_q}{M_h} H_1 \right) \right\}$$

— Similarly a **LU asymmetry** in unpolarized **dihadron** production

Need polarized e+e- colliders!

❑ Are existing facilities enough?

adapted from Particle Data Book

	BEPC	HIEPA	super KEKB	ILC	JLab/eEIC ??
E beam [GeV]	1.9	symmetric	4 (e ⁻) 7 (e ⁻)	250	?
√s [GeV]	3 – 5	2 – 7	10	500	?
polarization	? (beam self- polarization)	One beam	maybe	80% e ⁻ 60% e ⁺	YES!

❑ Can we get a (polarized) e+ e- collider at JLab / BNL?

– At JLab12 ? JLEIC ? eRHIC?

❑ What else is interesting to study?

– Factorization tests for FFs (low s, unpol), ...

← Ideas?