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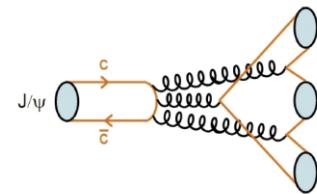
16-20 April 2018 Kobe, Japan

# Study of radiative decays of the $\Upsilon(1S)$ and of three-body decays of the $J/\psi$ at BABAR

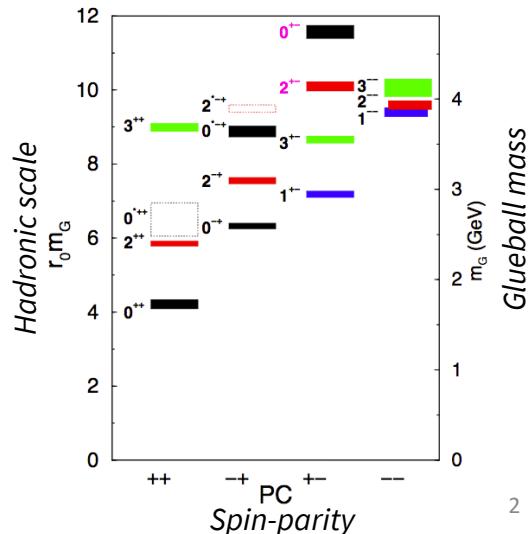
Patrick Robbe, LAL Orsay, for the BABAR Collaboration, 18 April 2018

# Radiative and hadronic decays of quarkonium

- Radiative or hadronic quarkonium decays can be used to perform spectroscopy of light mesons ( $f_J$ ,  $K^*$ , ...).
- They provide also a good laboratory to explore « exotic » QCD states like multiquark states or glueball (bound states of gluons).
- The lowest state  $J^{PC}=0^{++}$  is expected to have a mass around  $1.5 \text{ GeV}/c^2$ , accessible in quarkonium decays.
- One possible candidate for this state could be the  $f_0(1710)$  for which more experimental information is important.
- Good understanding of the decay characteristics is crucial: requires very low background which can be achieved for quarkonium produced in  $e^+e^-$  collisions.

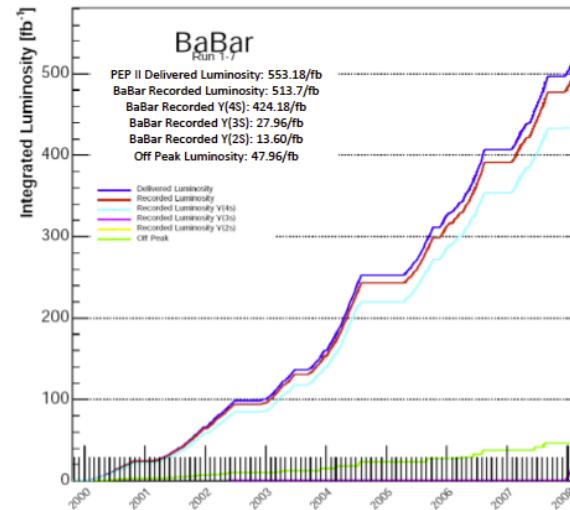
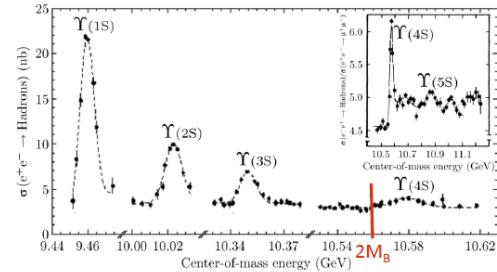
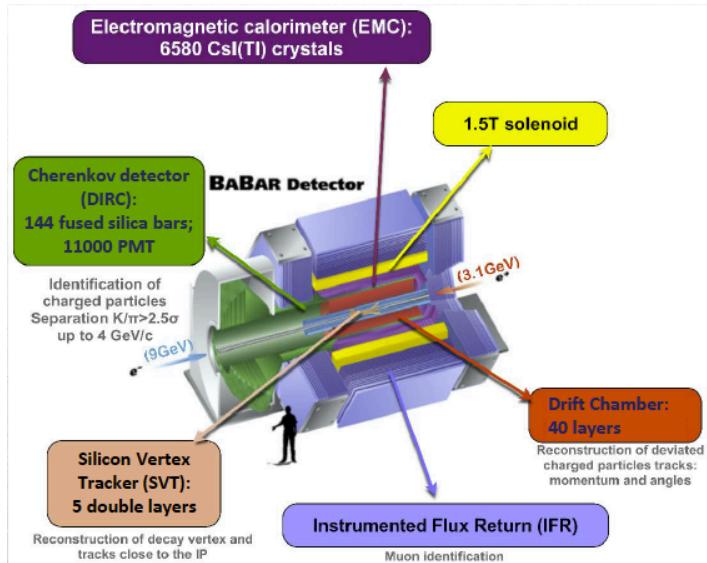


[C. Morningstar, M. Peardon, PRD60 (1999) 034509]



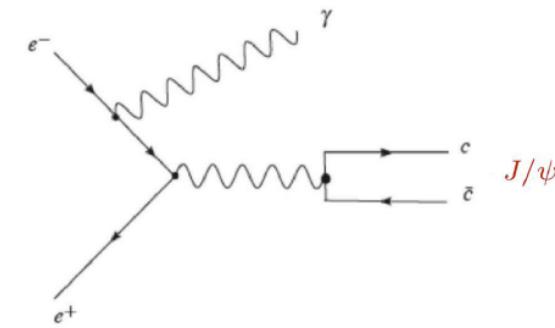
# The BABAR experiment

- Asymmetric  $e^+e^-$  collisions at 10.58 GeV center of mass energy:  $0.5 \text{ ab}^{-1}$  i.e.  $670 \times 10^6 c\bar{c}$  pairs for  $J/\psi$  studies.
- $14 \text{ fb}^{-1}$  and  $28 \text{ fb}^{-1}$  at  $\Upsilon(2S)$  and  $\Upsilon(3S)$  peaks.



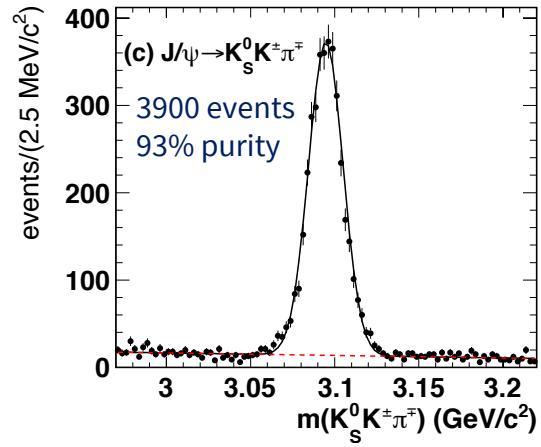
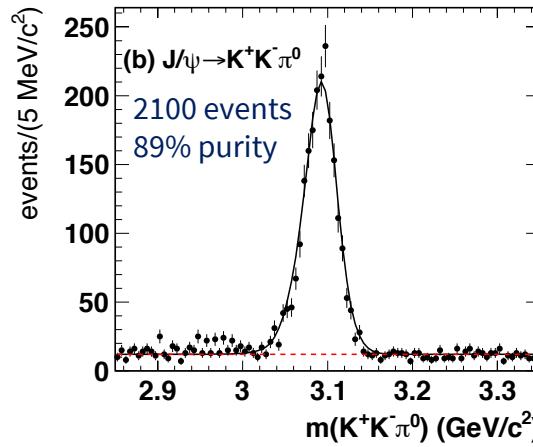
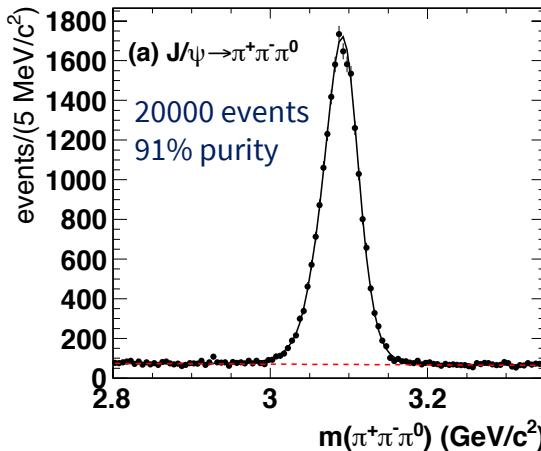
# Dalitz plot analysis of J/ $\psi$ decays

- Results on Dalitz plot analysis of the decays:
  - $J/\psi \rightarrow \pi^+ \pi^- \pi^0$ ,
  - $J/\psi \rightarrow K^+ K^- \pi^0$ ,
  - $J/\psi \rightarrow K_s^0 K^+ \pi^- + \text{charge conjugate}$ .
- Produced via  $e^+e^-$  annihilation with initial-state radiation (ISR):
  - Photon most of the time undetected.
  - Only  $J/\psi$  produced: no background from underlying event.
- Investigate these decays with large statistics:
  - $J/\psi \rightarrow \pi^+ \pi^- \pi^0$  shows complex structure with high mass resonances which are investigated with models implementing Regge asymptotics (Veneziano models).
  - $J/\psi \rightarrow K_s^0 K^+ \pi^-$ : first Dalitz plot analysis.



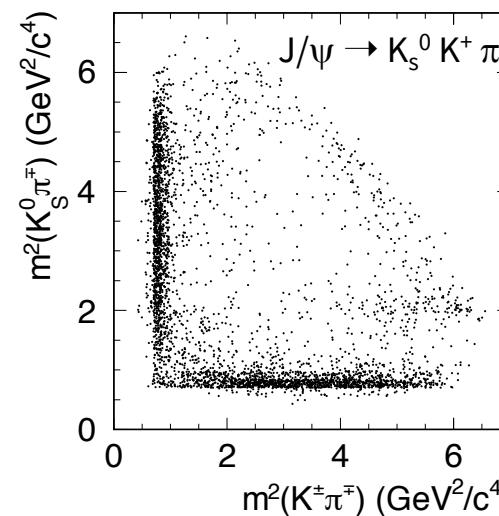
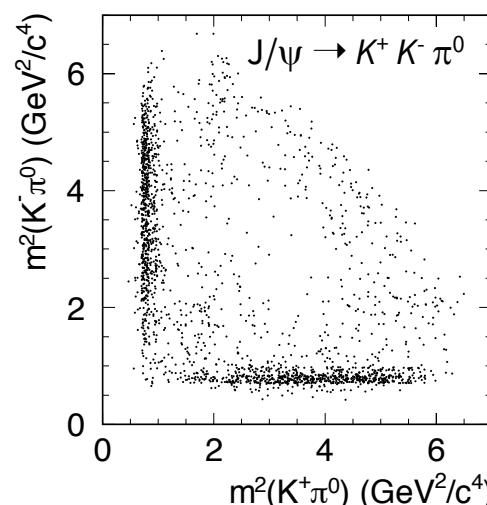
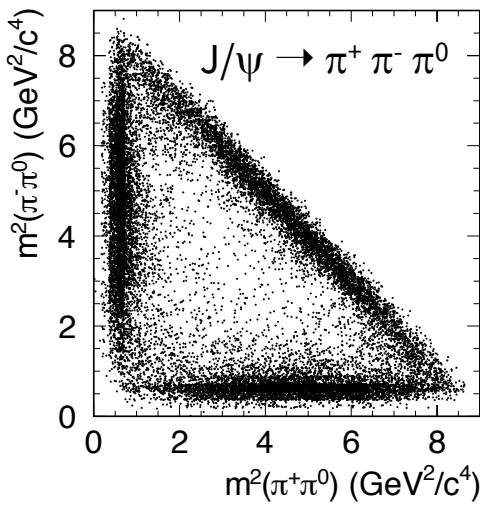
# J/ $\psi$ event selection

- $J/\psi \rightarrow K_s^0 K^+ \pi^-$ ,  $K_s^0 \rightarrow \pi^+ \pi^-$ : 4 tracks and  $K_s^0$  distance to interaction point  $> 2$  mm.
- $J/\psi \rightarrow \pi^+ \pi^- \pi^0$  and  $J/\psi \rightarrow K^+ K^- \pi^0$ : 2 tracks and  $E(\gamma) > 100$  MeV.
- Particle identification for charged  $K$  and  $\pi$ .
- To select  $J/\psi$  coming from ISR events:
  - $| (p_{e^+} + p_{e^-}) - (p_{h^+} + p_{h^-} + p_{\pi^0}) |^2 < 2 \text{ GeV}^2$  or  $| (p_{e^+} + p_{e^-}) - (p_{K_s^0} + p_{K^-} + p_{\pi^+}) |^2 < 1.5 \text{ GeV}^2$ .
  - If ISR photon is in calorimeter acceptance, require measurement compatible with its properties.
- Background from  $e^+ e^- \rightarrow \gamma \pi^+ \pi^-$  rejected with  $|\cos\theta_\pi| < 0.95$  where  $\theta_\pi$  is helicity angle of pion for the  $\pi\pi$  system.



# J/ $\psi$ branching fractions and Dalitz plots

- $\mathcal{R}_1 = \frac{\mathcal{B}(J/\psi \rightarrow K^+ K^- \pi^0)}{\mathcal{B}(J/\psi \rightarrow \pi^+ \pi^- \pi^0)} = 0.120 \pm 0.003 \text{ (stat)} \pm 0.009 \text{ (syst)} \text{ [PDG: } 0.133 \pm 0.038 \text{ [MARK II, PRL51 (1983) 963]}]$
- $\mathcal{R}_2 = \frac{\mathcal{B}(J/\psi \rightarrow K_s^0 K^\pm \pi^\mp)}{\mathcal{B}(J/\psi \rightarrow \pi^+ \pi^- \pi^0)} = 0.265 \pm 0.005 \text{ (stat)} \pm 0.021 \text{ (syst)} \text{ [PDG: } 0.123 \pm 0.033 \text{ [MARK I, PRD15 (1977) 1814]}]$
- With main systematic uncertainties from knowledge of efficiencies.



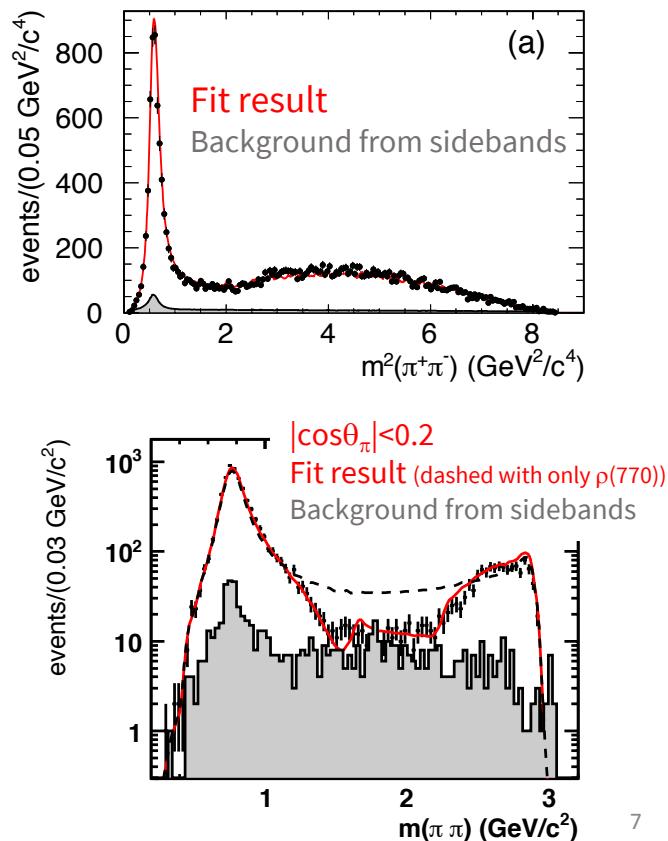
# J/ $\psi$ $\rightarrow \pi^+ \pi^- \pi^0$ isobar model

- Dalitz plot analysis in the J/ $\psi$  signal region, describing resonances by relativistic Breit Wigner shapes. Total of 8 free parameters.

Final state	Amplitude	Isobar fraction (%)	Phase (radians)
$\rho(770)\pi$	1.	$114.2 \pm 1.1$	$\pm 2.6$
$\rho(1450)\pi$	$0.513 \pm 0.039$	$10.9 \pm 1.7$	$\pm 2.7$
$\rho(1700)\pi$	$0.067 \pm 0.007$	$0.8 \pm 0.2$	$\pm 0.5$
$\rho(2150)\pi$	$0.042 \pm 0.008$	$0.04 \pm 0.01$	$\pm 0.20$
$\omega(783)\pi^0$	$0.013 \pm 0.002$	$0.08 \pm 0.03$	$\pm 0.02$
$\rho_3(1690)\pi$			
Sum		$127.8 \pm 2.0 \pm 4.3$	
$\chi^2/\nu$		$687/519 = 1.32$	

- When leaving free the  $\rho(1450)$  and  $\rho(1700)$ , the best fit is obtained for:

$$\begin{aligned} m(\rho(1450)) &= 1429 \pm 41 \text{ MeV}/c^2, \\ \Gamma(\rho(1450)) &= 576 \pm 29 \text{ MeV}, \\ m(\rho(1700)) &= 1644 \pm 36 \text{ MeV}/c^2, \\ \Gamma(\rho(1700)) &= 109 \pm 19 \text{ MeV}. \end{aligned}$$



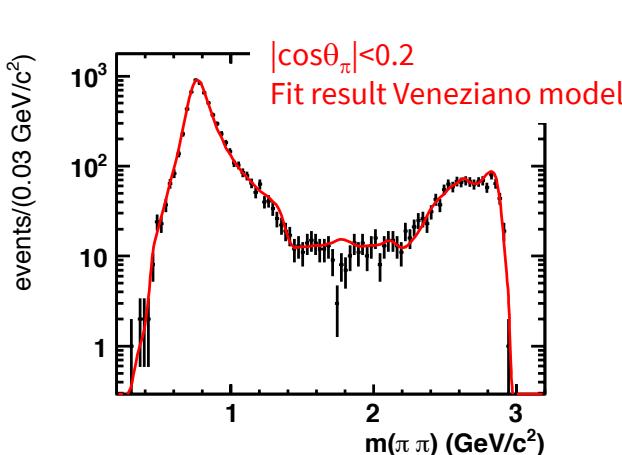
# J/ $\psi$ $\rightarrow \pi^+ \pi^- \pi^0$ Veneziano model

- Veneziano model = alternative model to describe better the high mass region [Szczepaniak, Pennington, PLB737 (2014) 283].

$$A(s, t, \lambda) = \epsilon_{\mu\nu\rho\sigma} p_\mu^+ p_\nu^- p_\sigma^0 \epsilon(p^\psi, \lambda) \sum_{1 \leq m \leq n} \frac{\Gamma(n - \alpha(s)) \Gamma(n - \alpha(t))}{\Gamma(n + m - \alpha(s) - \alpha(t))}$$

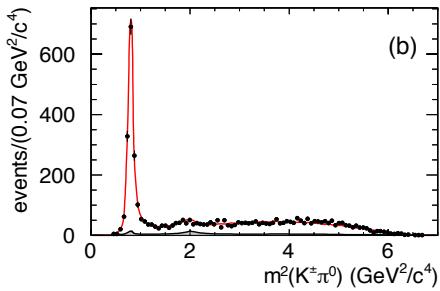
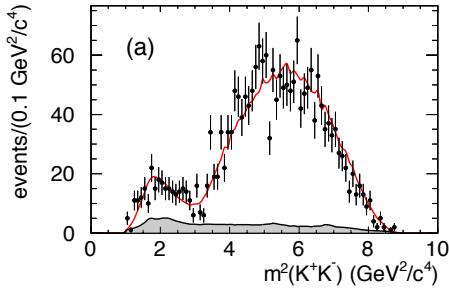
$$\alpha(s) = \alpha_0 + \alpha' s + i\gamma\sqrt{s - s_0}$$

- The model deals with Regge trajectories instead of single resonances. The complexity of the model is related to the number of trajectories,  $n$ , to include in the fit. Here:  $n=7$  i.e. 19 free parameters.



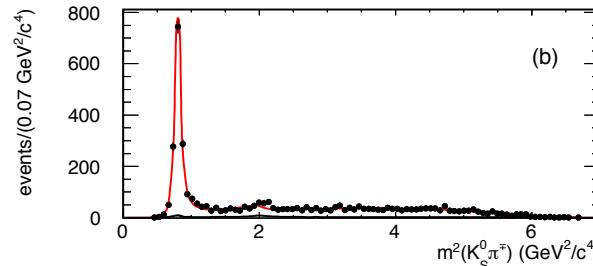
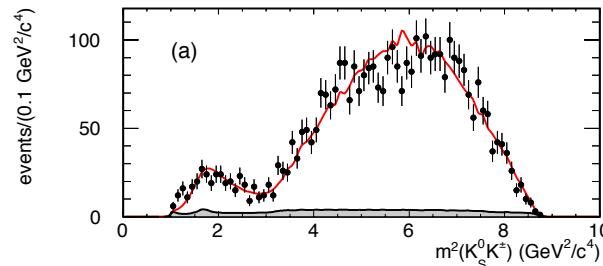
Final state	Isobar fraction (%)	Veneziano fraction (%)
$\rho(770)\pi$	$114.2 \pm 1.1 \pm 2.6$	$133.1 \pm 3.3$
$\rho(1450)\pi$	$10.9 \pm 1.7 \pm 2.7$	$0.80 \pm 0.27$
$\rho(1700)\pi$	$0.8 \pm 0.2 \pm 0.5$	$2.20 \pm 0.60$
$\rho(2150)\pi$	$0.04 \pm 0.01 \pm 0.20$	$6.00 \pm 2.50$
$\omega(783)\pi^0$	$0.08 \pm 0.03 \pm 0.02$	
$\rho_3(1690)\pi$		$0.40 \pm 0.08$
Sum	$127.8 \pm 2.0 \pm 4.3$	$142.5 \pm 2.8$
$\chi^2/\nu$	$687/519 = 1.32$	$596/508 = 1.17$

# $J/\psi \rightarrow K^+ K^- \pi^0$ and $J/\psi \rightarrow K_s^0 K^+ \pi^-$ results



- $J/\psi \rightarrow K^+ K^- \pi^0$  can be used to extract

$$\frac{B(\rho(1450)^0 \rightarrow K^+ K^-)}{B(\rho(1450)^0 \rightarrow \pi^+ \pi^-)} = 0.307 \pm 0.084 \text{ (stat)} \pm 0.092 \text{ (syst).}$$



Final state	fraction (%)	phase (radians)
$K^*(892)^\pm K^\mp$	$92.4 \pm 1.5 \pm 3.4$	0.
$\rho(1450)^0 \pi^0$	$9.3 \pm 2.0 \pm 0.6$	$3.78 \pm 0.28 \pm 0.08$
$K^*(1410)^\pm K^\mp$	$2.3 \pm 1.1 \pm 0.7$	$3.29 \pm 0.26 \pm 0.39$
$K_2^*(1430)^\pm K^\mp$	$3.5 \pm 1.3 \pm 0.9$	$-2.32 \pm 0.22 \pm 0.05$
Total	$107.4 \pm 2.8$	
$\chi^2/\nu$	$132/137 = 0.96$	

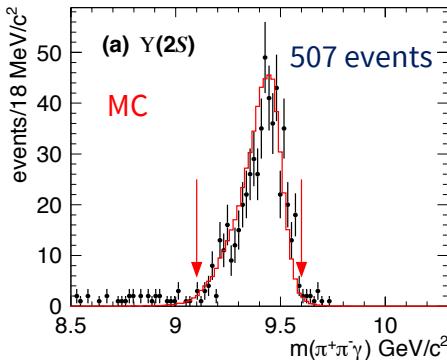
Final state	fraction (%)	phase (radians)
$K^*(892)\bar{K}$	$90.5 \pm 0.9 \pm 3.8$	0.
$\rho(1450)^\pm \pi^\mp$	$6.3 \pm 0.8 \pm 0.6$	$-3.25 \pm 0.13 \pm 0.21$
$K_1^*(1410)\bar{K}$	$1.5 \pm 0.5 \pm 0.9$	$1.42 \pm 0.31 \pm 0.35$
$K_2^*(1430)\bar{K}$	$7.1 \pm 1.3 \pm 1.2$	$-2.54 \pm 0.12 \pm 0.12$
Total	$105.3 \pm 3.1$	
$\chi^2/\nu$	$274/217 = 1.26$	

# Study of Y(1S) radiative decays

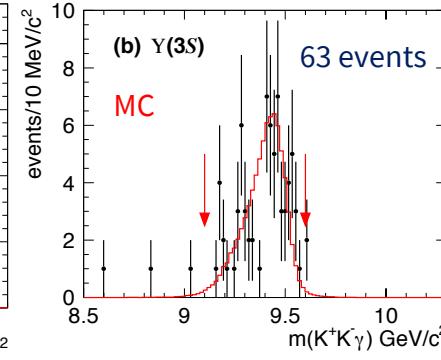
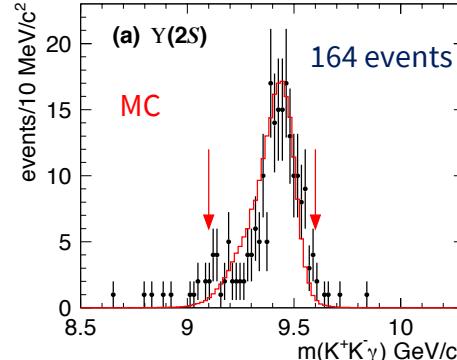
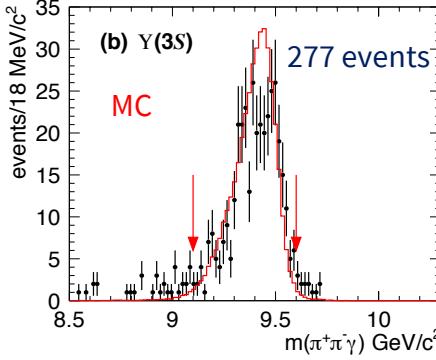
- Radiative decays suppressed in Y system compared to charmonium system: challenging analysis.
- Measurement of branching fractions and resonance structure of:
  - $Y(1S) \rightarrow \pi^+ \pi^- \gamma$ ,
  - $Y(1S) \rightarrow K^+ K^- \gamma$ .
- $Y(1S)$  are produced from  $Y(2S)$  and  $Y(3S)$  decays in the modes  $Y(2S) \rightarrow \pi^+ \pi^- Y(1S)$  and  $Y(3S) \rightarrow \pi^+ \pi^- Y(1S)$  where the two pions are soft: this allows to obtain clean signal using the  $13.6 \text{ fb}^{-1}$  and  $28 \text{ fb}^{-1}$  of  $Y(2S)$  and  $Y(3S)$  on-resonance data.
- Branching fractions are normalized to the  $Y(1S) \rightarrow \mu^+ \mu^-$  decay for which are reconstructed:
  - $435000 Y(2S) \rightarrow \pi^+ \pi^- [Y(1S) \rightarrow \mu^+ \mu^-]$ ,
  - $132000 Y(3S) \rightarrow \pi^+ \pi^- [Y(1S) \rightarrow \mu^+ \mu^-]$ .

# $\Upsilon(1S)$ event selection

- Exactly 4 charged tracks with  $p_T > 100$  MeV/c.
- Exactly 1 photon with  $E > 2.5$  GeV.
- Particle identification for charged  $K$  and  $\pi$ .
- Momentum balance in the event,  $\Delta p_i = p_i^{e^-} + p_i^{e^+} - (p_i^\gamma + p_i^{\pi_s^+} + p_i^{\pi_s^-} + p_i^{h^+} + p_i^{h^-})$ , should be close to 0.
- The recoiling mass to the two soft pions in  $\Upsilon(2S) \rightarrow \pi^+\pi^-\Upsilon(1S)$  and  $\Upsilon(3S) \rightarrow \pi^+\pi^-\Upsilon(1S)$ ,  $M_{rec}^2 = |p_{e^-} + p_{e^+} - p_{\pi^+} - p_{\pi^-}|^2$ , should be close to the  $\Upsilon(1S)$  mass.



$\Upsilon(1S) \rightarrow \pi^+\pi^-\gamma$



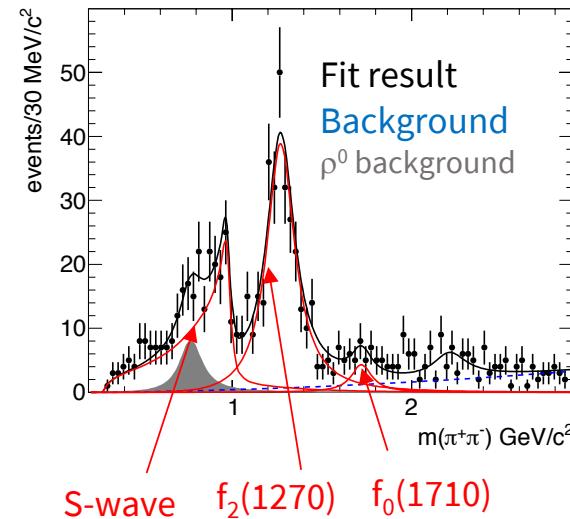
$\Upsilon(1S) \rightarrow K^+K^-\gamma$

# $\pi^+\pi^-$ mass spectrum

- Summed over the Y(2S) and Y(3S) samples.
- Fit with mass function sum of (16 free parameters):
  - s-wave term: sum of  $f_0(500)$  described by relativistic Breit Wigner and  $f_0(980)$  described by Flatté formalism.
  - $f_2(1270)$  and  $f_0(1710)$  as relativistic Breit Wigner.
  - Combinatorial background.
  - $\rho^0$  background for Y(3S) data, as indicated from mass sidebands.

Resonances ( $\pi^+\pi^-$ )	Yield $\Upsilon(2S)$	Yield $\Upsilon(3S)$	Significance ( $\sigma$ )
S-wave	$133 \pm 16 \pm 13$	$87 \pm 13$ (stat only)	12.8
$f_2(1270)$	$255 \pm 19 \pm 8$	$77 \pm 7 \pm 4$	15.9
$f_0(1710)$	$24 \pm 8 \pm 6$	$6 \pm 8 \pm 3$	2.5
$f_0(2100)$	$33 \pm 9$ (stat only)	$8 \pm 15$ (stat only)	
$\rho(770)^0$		$54 \pm 23$ (stat only)	

- Significant s-wave contribution with:
  - $M(f_0(500)) = 0.856 \pm 0.086 \text{ GeV}/c^2$ ,  $\Gamma(f_0(500)) = 1.279 \pm 0.324 \text{ GeV}$ ,
  - fraction of  $f_0(500)$  component of  $(27.7 \pm 3.1)\%$ .



- After efficiency correction:

Resonance	$\Upsilon(2S)$ corrected yield	$\Upsilon(3S)$ corrected yield
S-wave	$541 \pm 65 \pm 53$	
$f_2(1270)$	$1043 \pm 78 \pm 36$	$285 \pm 26 \pm 15$
$f_0(1710)$	$95 \pm 32 \pm 24$	$22 \pm 29 \pm 11$

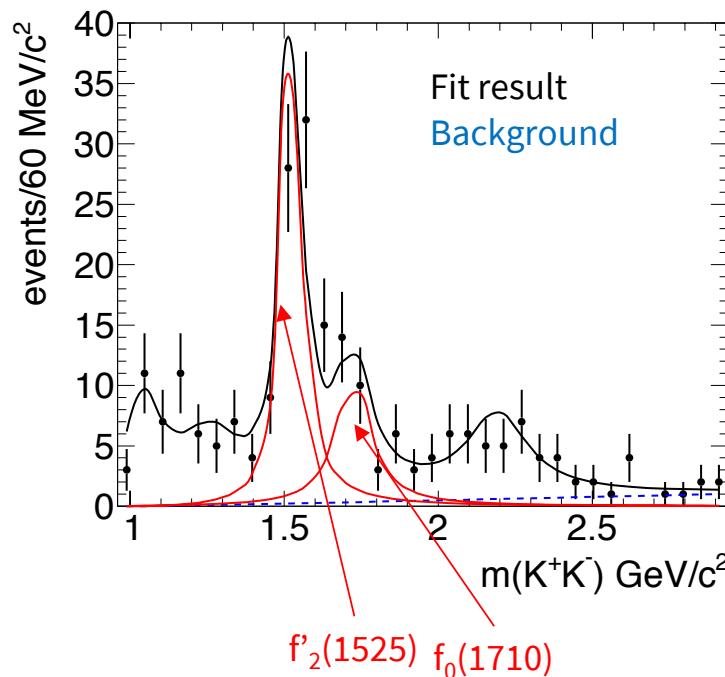
# $K^+K^-$ mass spectrum

- Summed over the Y(2S) and Y(3S) samples.
- Fit with mass function sum of (6 free parameters):
  - $f_0(980)$  described by Flatté formalism.
  - $f_2(1270)$ ,  $f'_2(1525)$ ,  $f_0(1500)$  (sum of both noted as  $f_J(1500)$ ) and  $f_0(1710)$  as relativistic Breit Wigner.
  - Combinatorial background.

Resonances ( $K^+K^-$ )	Yield $\Upsilon(2S) + \Upsilon(3S)$	Significance ( $\sigma$ )
$f_0(980)$	$47 \pm 9$	5.6
$f_J(1500)$	$77 \pm 10 \pm 10$	8.9
$f_0(1710)$	$36 \pm 9 \pm 6$	4.7
$f_2(1270)$	$15 \pm 8$	
$f_0(2200)$	$38 \pm 8$	

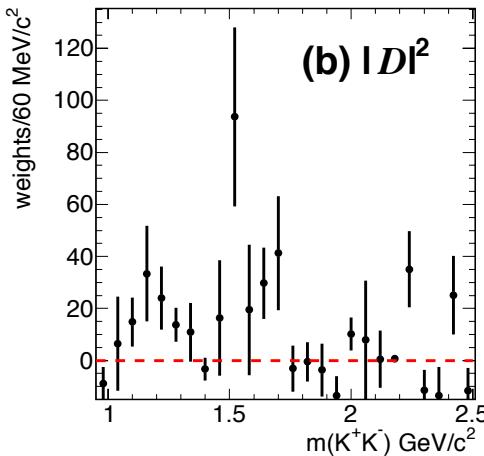
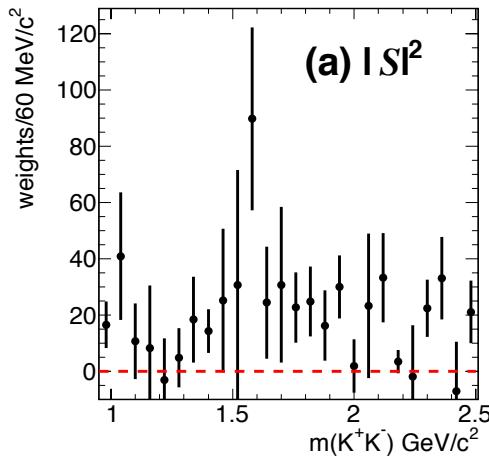
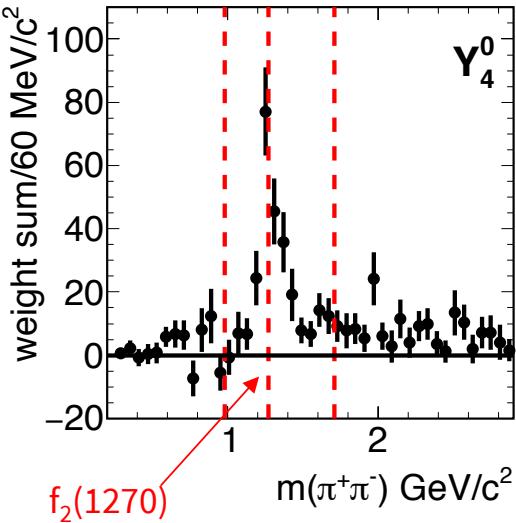
- After efficiency correction:

Resonance	$\Upsilon(2S)/\Upsilon(3S)$ corrected yield
$f_J(1500)$	$281 \pm 37 \pm 38$
$f_0(1710)$	$143 \pm 36 \pm 24$



# Legendre polynomial moments

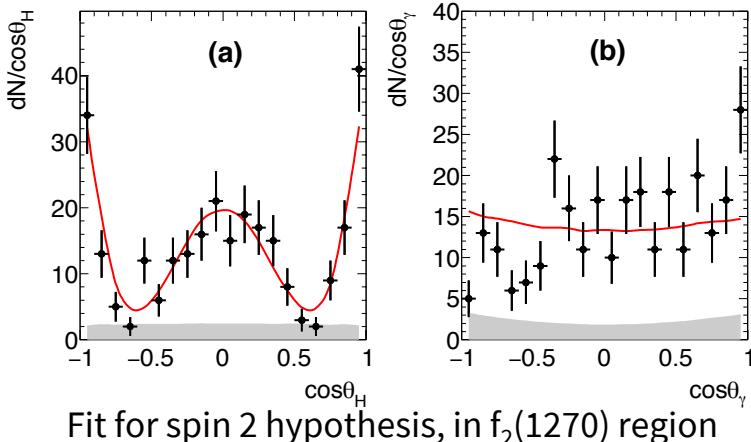
- Information about the spin structure, and proportions of  $S$  and  $D$  wave amplitudes can be obtained from distributions of Legendre polynomial  $Y_n^0$ .
- In the case of the  $K^+K^-$  spectrum, this partial wave analysis shows that the  $f_J(1500)$  component is composed of both spin 0 and 2 resonances.



$$\begin{aligned}\sqrt{4\pi} <Y_0^0> &= S^2 + D^2, \\ \sqrt{4\pi} <Y_2^0> &= 2SD \cos \phi_{SD} + 0.639D^2, \\ \sqrt{4\pi} <Y_4^0> &= 0.857D^2,\end{aligned}$$

# Helicity amplitude fits

- Complete spin-parity analysis is performed fitting the angular distributions expected for the various spin hypotheses, as a function of the decay angles in the  $\Upsilon(2/3S)$  decay chain, with unbinned likelihood fit in each resonance region.
  - $\theta_H$ : helicity angle of  $\pi$  or  $K$ ,
  - $\theta_\gamma$ : photon angle.
- Branching fractions extracted using fractions from these fits and the previous ones.



Spin 2:

$$W_2(\theta_\gamma, \theta_H) = \frac{dU(\theta_\gamma, \theta_H)}{d\cos\theta_\gamma d\cos\theta_H} = \frac{15}{1024} |E_{00}|^2 [6|A_{01}|^2 (22|C_{10}|^2 + 8|C_{11}|^2 + 9|C_{12}|^2) + 2|A_{00}|^2 (22|C_{10}|^2 + 24|C_{11}|^2 + 9|C_{12}|^2) + 24 (|A_{00}|^2 + 3|A_{01}|^2) (2|C_{10}|^2 - |C_{12}|^2) \cos 2\theta_H + 6 (|A_{00}|^2 (6|C_{10}|^2 - 8|C_{11}|^2 + |C_{12}|^2) + |A_{01}|^2 (18|C_{10}|^2 - 8|C_{11}|^2 + 3|C_{12}|^2)) \cos 4\theta_H - 2 (|A_{00}|^2 - |A_{01}|^2) \cos 2\theta_\gamma (22|C_{10}|^2 - 24|C_{11}|^2 + 9|C_{12}|^2) + 12 (2|C_{10}|^2 - |C_{12}|^2) \cos 2\theta_H + 3 (6|C_{10}|^2 + 8|C_{11}|^2 + |C_{12}|^2) \cos 4\theta_H].$$

Spin 0:

$$W_0(\theta_\gamma) = \frac{dU(\theta_\gamma)}{d\cos\theta_\gamma} = \frac{3}{8} |C_{10}|^2 |E_{00}|^2 (|A_{00}|^2 + 3|A_{01}|^2 - (|A_{00}|^2 - |A_{01}|^2) \cos 2\theta_\gamma)$$

Resonance	$\mathcal{B}(10^{-5})$
$\pi\pi$ <i>S</i> -wave	$4.63 \pm 0.56 \pm 0.48$
$f_2(1270)$	$10.15 \pm 0.59 \begin{array}{l} +0.54 \\ -0.43 \end{array}$
$f_0(1710) \rightarrow \pi\pi$	$0.79 \pm 0.26 \pm 0.17$
$f_J(1500) \rightarrow KK$	$3.97 \pm 0.52 \pm 0.55$
$f'_2(1525)$	$2.13 \pm 0.28 \pm 0.72$
$f_0(1500) \rightarrow K\bar{K}$	$2.08 \pm 0.27 \pm 0.65$
$f_0(1710) \rightarrow KK$	$2.02 \pm 0.51 \pm 0.35$

# Conclusions

- Dalitz plot analysis of  $\text{J}/\psi$  decays with high statistics:
  - Provide complementary description of the resonance structures in the  $\text{J}/\psi \rightarrow \pi^+ \pi^- \pi^0$  with Veneziano model.
  - First measurement of the  $\text{J}/\psi \rightarrow K_s^0 K^+ \pi^-$ .
- Studies of radiative  $\Upsilon(1S)$  decays in  $\Upsilon(1S) \rightarrow \pi^+ \pi^- \gamma$  and  $\Upsilon(1S) \rightarrow K^+ K^- \gamma$ .
  - Spin-parity and branching fraction measurements.
  - Observation of the  $f_0(1710)$  state in these decays, which can be used to gain information on glueball.