

Electroweak Physics with the LHeC+FCCeh

Definitions
Measurements
+ Simulations
Couplings
EFT + h.o. Studies

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For the LHeC Study Group



<http://lhec.web.cern.ch>



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Electroweak Physics Probing Nuclear Structure *)

Deep Inelastic Scattering Testing the Electroweak Theory

$$\begin{aligned}
 F_2^\pm &= F_2 + \kappa_Z(-v_e \mp Pa_e) \cdot F_2^{\gamma Z} + \kappa_Z^2(v_e^2 + a_e^2 \pm 2Pv_e a_e) \cdot F_2^Z \\
 xF_3^\pm &= \kappa_Z(\pm a_e + Pv_e) \cdot xF_3^{\gamma Z} + \kappa_Z^2(\mp 2v_e a_e - P(v_e^2 + a_e^2)) \cdot xF_3^Z.
 \end{aligned}$$

$$\kappa_Z(Q^2) = \frac{Q^2}{Q^2 + M_Z^2} \cdot \frac{1}{4 \sin^2 \Theta \cos^2 \Theta} \quad v_f = i_f - e_f 2 \sin^2 \Theta \quad a_f = i_f$$

$$\begin{aligned}
 (F_2, F_2^{\gamma Z}, F_2^Z) &= x \sum (e_q^2, 2e_q v_q, v_q^2 + a_q^2)(q + \bar{q}) \\
 (xF_3^{\gamma Z}, xF_3^Z) &= 2x \sum (e_q a_q, v_q a_q)(q - \bar{q}),
 \end{aligned}$$

	e^2	$2ev$
u	4/9	2/9
d	1/9	2/9

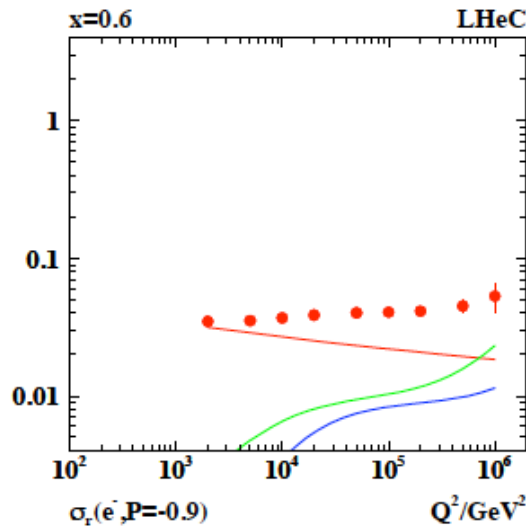
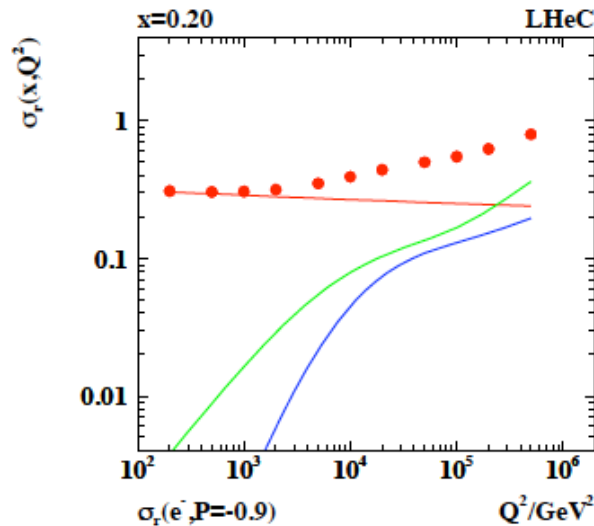
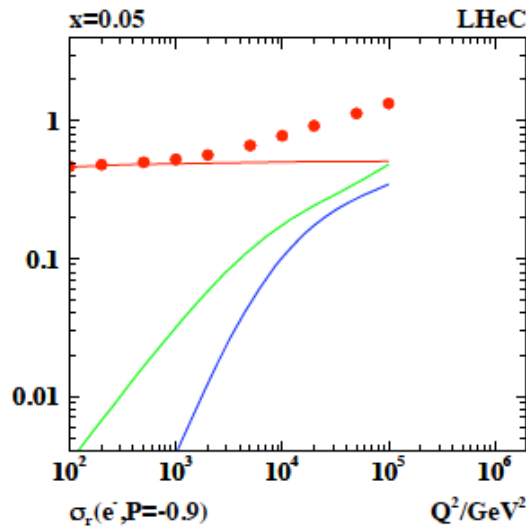
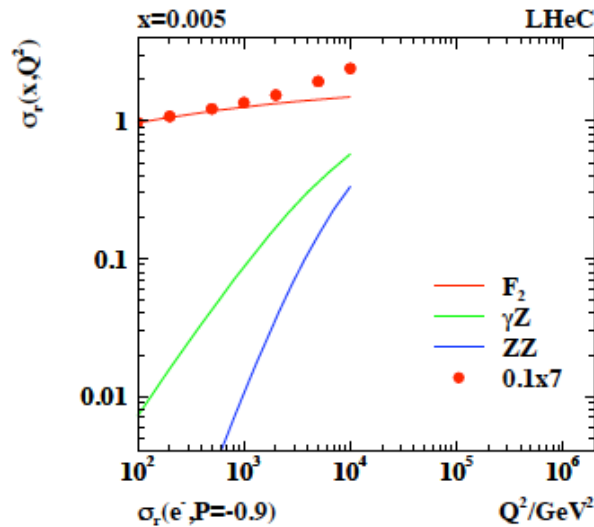
NC: $\gamma\gamma, \gamma Z, ZZ$. Lepton beam helicity P , M_Z , v and a couplings, PV through va
 CC: pure weak cross section (G_F, M_W) \rightarrow 3 independent variables, DIS: OMS

*) M. Klein and T. Riemann¹

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Reduced NC e-p Scattering Cross Section [P=-0.9, 10fb⁻¹]



Huge Z exchange effects at LHeC

At small Q^2 , $\sigma_r = F_2$

Both at $x=0.2$ (the point of Bj scaling) and at larger x (the region of gluon bremsstrahlung which makes F_2 decrease with Q^2), the reduced NC cross section rises

Polarisation Asymmetry and R=NC/CC

$$\frac{2}{P_L - P_R} \cdot A^\pm \simeq \mp \kappa_Z a_e \frac{F_2^{\gamma Z}}{(F_2 + \kappa_Z a_e Y_- x F_3^{\gamma Z} / Y_+)} \simeq \mp \kappa_Z a_e \frac{F_2^{\gamma Z}}{F_2}$$

$$\frac{2}{P_L - P_R} \cdot A^\pm \simeq \pm \kappa \frac{1 + d_v/u_v}{4 + d_v/u_v}$$

Classic asymmetry (Prescott et al, 1978) accesses weak interaction, $F_2^{\gamma Z}$ is a new, direct measure of valence quarks at high x

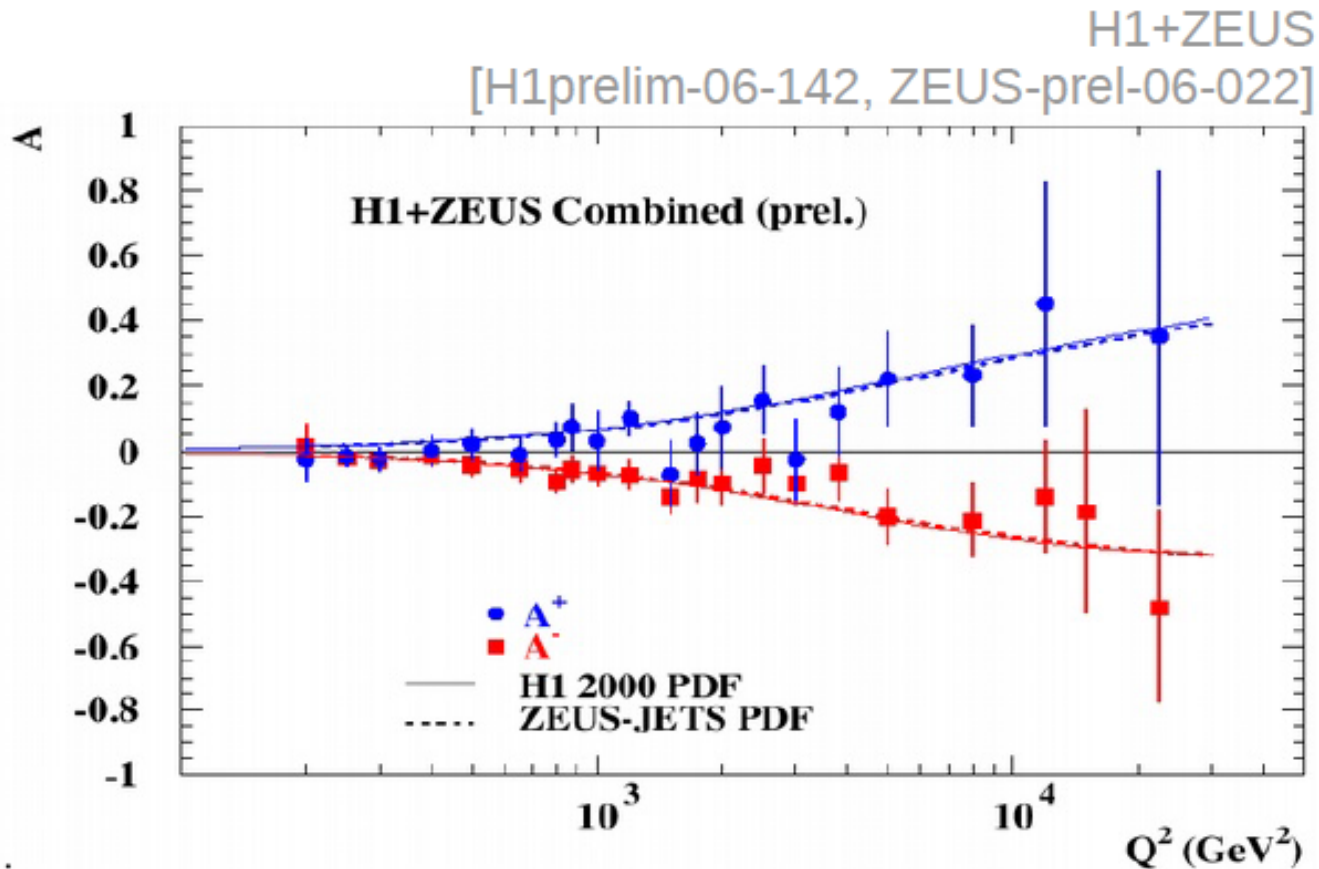
$$R^\pm = \frac{\sigma_{NC}^\pm}{\sigma_{CC}^\pm} = \frac{2}{(1 \pm P)\kappa_W^2} \cdot \frac{\sigma_{r,NC}^\pm}{\sigma_{r,CC}^\pm}$$

$$R^\pm \simeq \frac{2a_e^2}{(1 \pm P)\cos^2 \Theta} \cdot \frac{Y_+ F_2^Z - Y_- P x F_3^Z}{Y_+ W_2^\pm + Y_- x W_3^\pm}$$

R accesses weak interaction and the pure weak structure functions which are best measured at the LHeC/FCC-eh

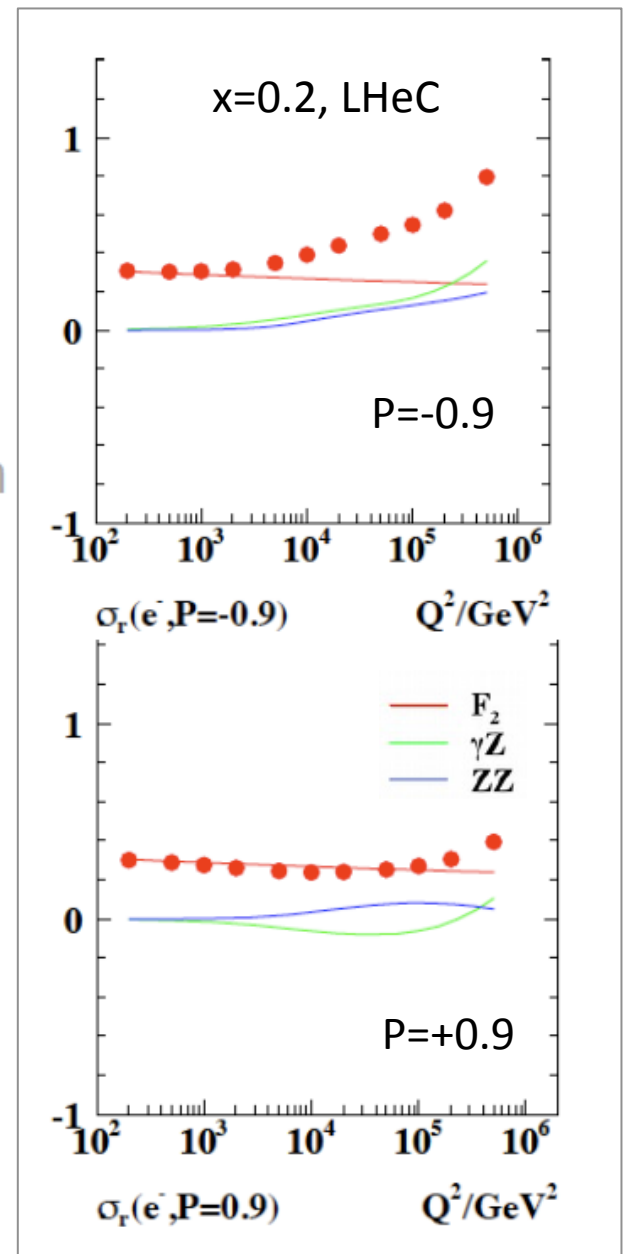
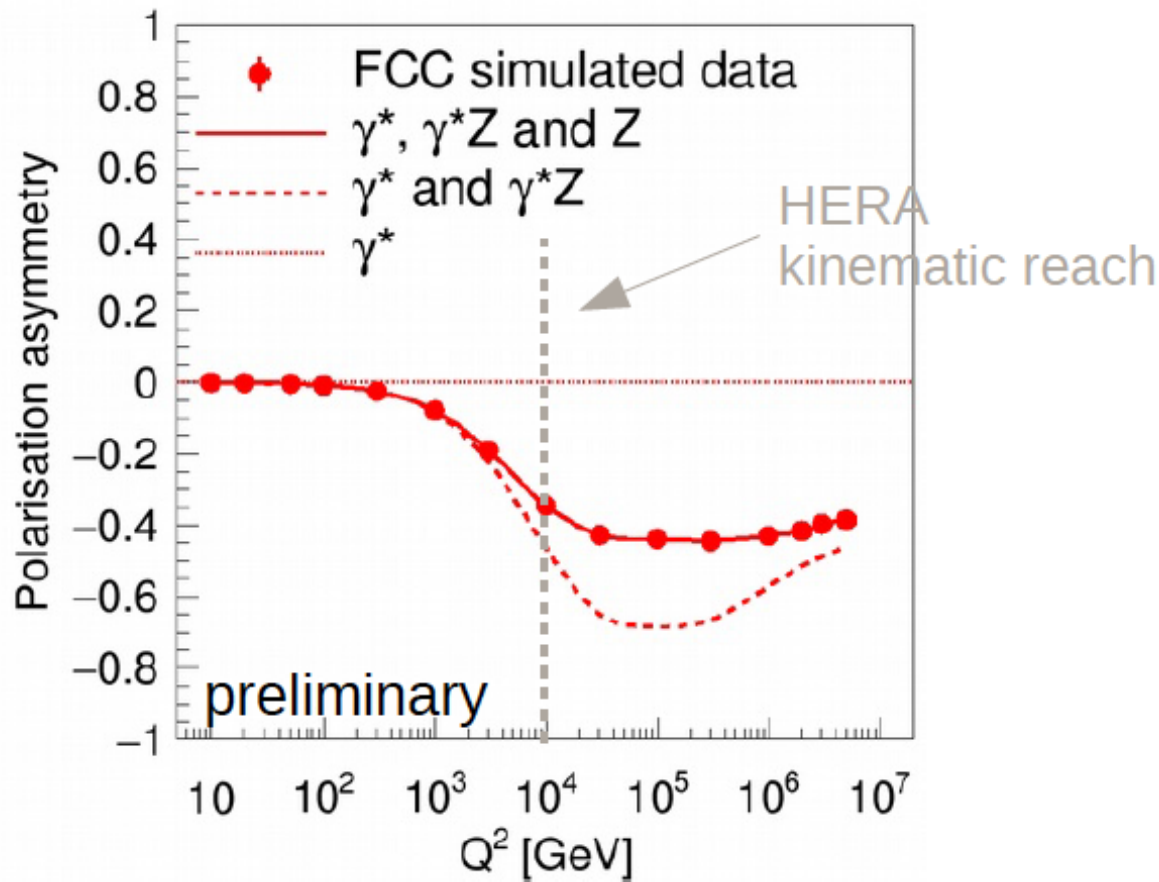
Note that in experiment you would measure the cross sections and determine all correlations which is still more informative than A or R but contains their physics.

PV Asymmetry in NC



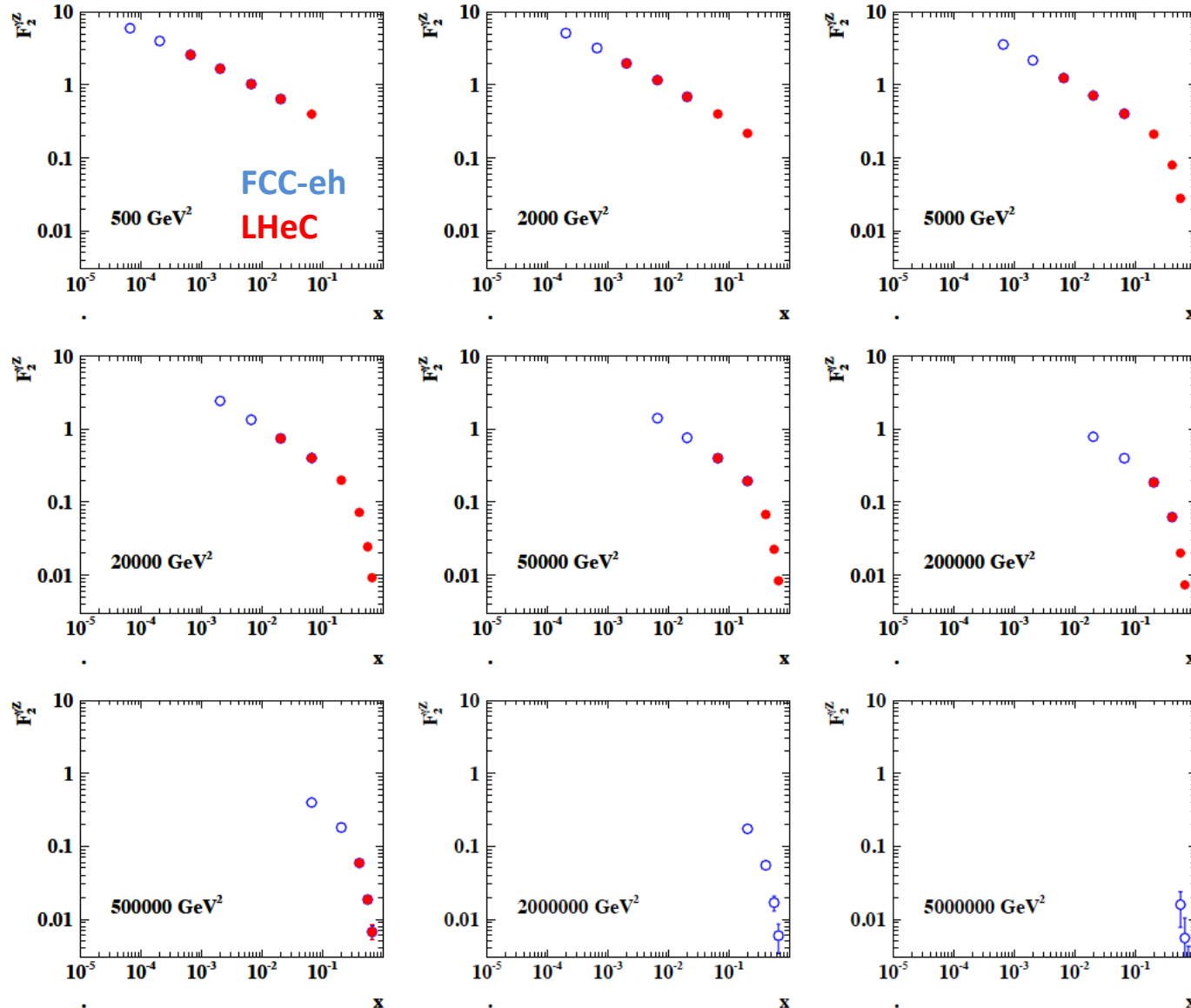
HERA: 20% asymmetry at $Q^2 = 10^4$ GeV² $A^+ = -A^-$

PV Asymmetry in NC

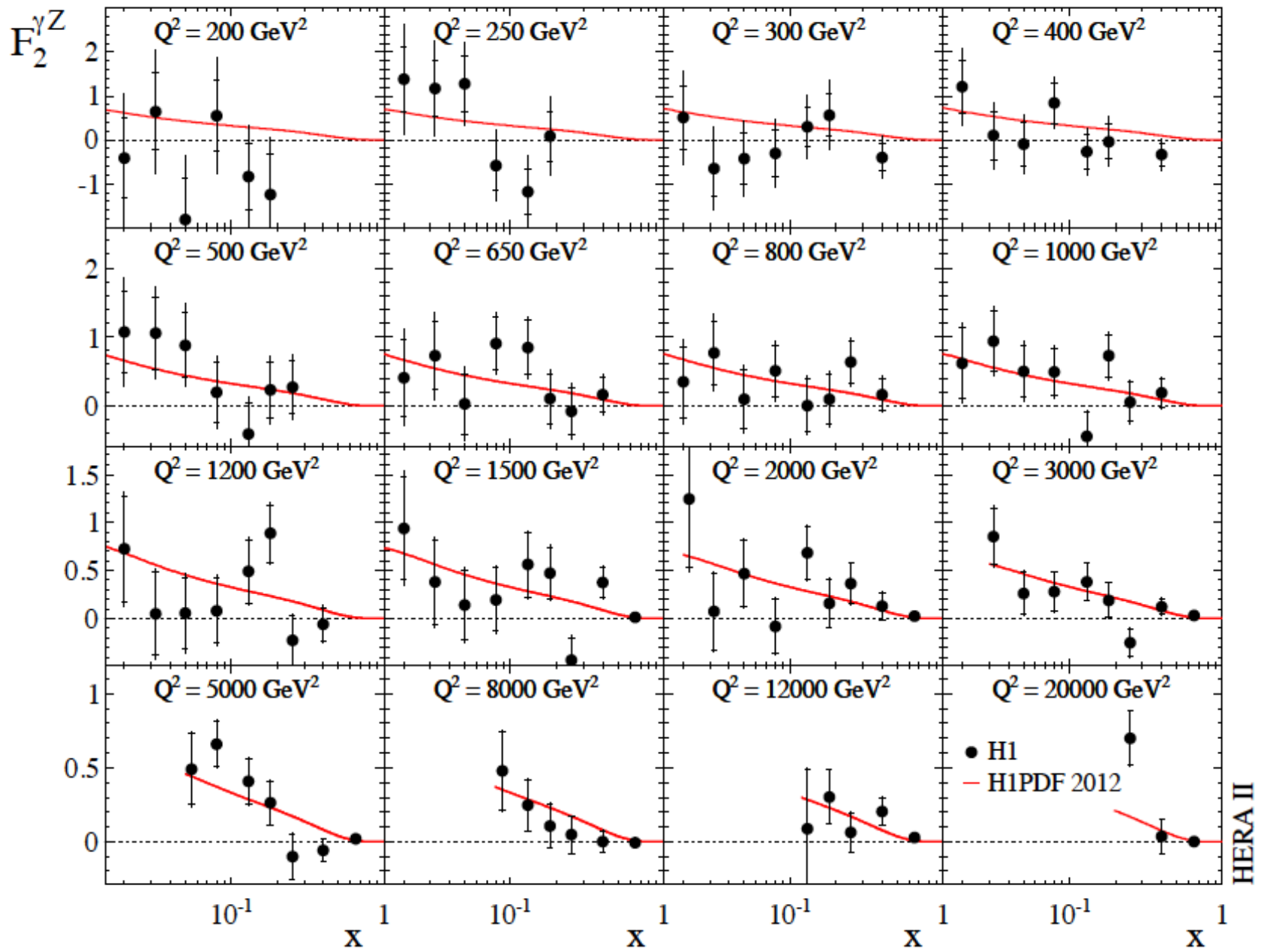


FCCeh: 40% integrated asymmetry for $Q^2 > 10^4 \text{ GeV}^2$, locally (x) much larger

Parity Violation Structure Function $F_2^{yZ} = x \sum 2 e_q v_q (q + \bar{q})$



LHeC: 7 TeV, FCC-eh: 50 TeV, $E_e=60$ GeV, integrated L of 100 fb^{-1} for $P= \pm 0.8$



H1: arXiv:1207.7007: much smaller x and Q^2 range, imprecise, but first measurement ever

Charged Currents

$$\frac{d^2\sigma^{\text{CC}}(e^\pm p)}{dx dQ^2} = (1 \pm P_e) \frac{G_F^2}{4\pi x} \left[\frac{m_W^2}{m_W^2 + Q^2} \right]^2 \left(Y_+ W_2^\pm(x, Q^2) \mp Y_- x W_3^\pm(x, Q^2) \right)$$

$$W_2^- = x(U + \bar{D}), \quad xW_3^- = x(U - \bar{D})$$

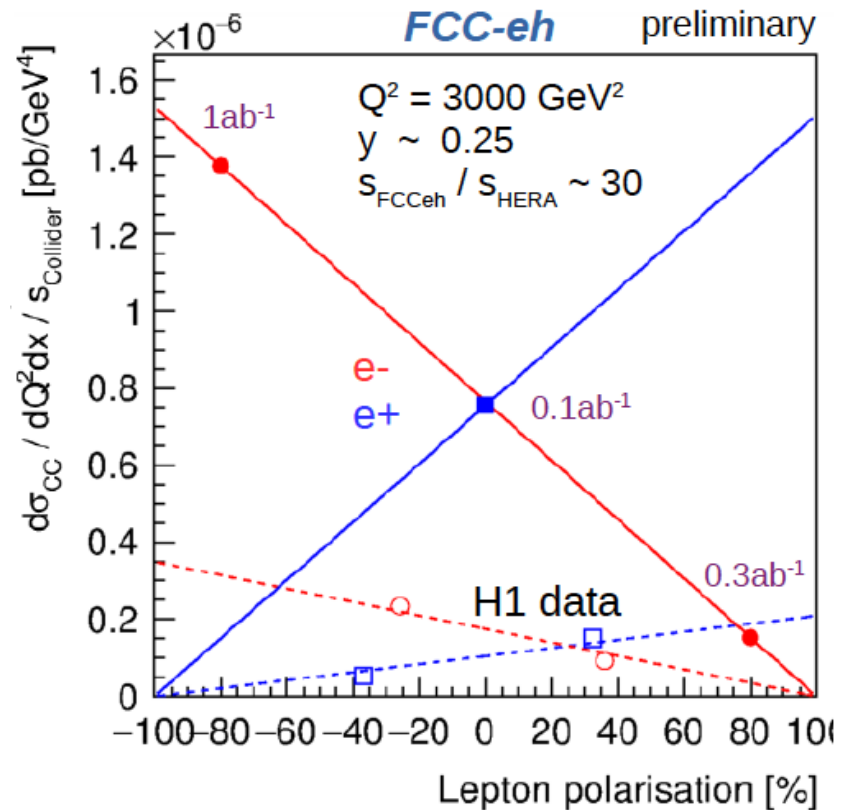
$$W_2^+ = x(\bar{U} + D), \quad xW_3^+ = x(D - \bar{U})$$

Data sets:

1ab⁻¹ P=-0.8 to enlarge rate for WW → H

0.3ab⁻¹ P=+0.8 for eweak physics

0.1ab⁻¹ unpolarised with positrons



Data (NC,CC) Simulation for QCD + el.weak Evaluation

Numerical treatment of correlated and uncorrelated systematic and statistical errors
[based on PHE-1990-02 (J.Blümlein, M.Klein), cross checked with H1 Monte Carlo]

source of uncertainty	error on the source or cross section
scattered electron energy scale $\Delta E'_e/E'_e$	0.1 %
scattered electron polar angle	0.1 mrad
hadronic energy scale $\Delta E_h/E_h$	0.5 %
calorimeter noise (only $y < 0.01$)	1-3 %
radiative corrections	0.5%
photoproduction background (only $y > 0.5$)	1 %
global efficiency error	0.7 %

- Assumptions gauged with H1, probably conservative.
- This approach determines full set of uncorrelated and correlated uncertainties
- ALL PDFs and electroweak fit results presented to this workshop have full systematic error
- This also holds for the LHeC CDR alphas analysis leading to 0.1-0.2% total uncertainty

NC Cross Section Correlated Uncertainties ($Q^2=20000 \text{ GeV}^2$)

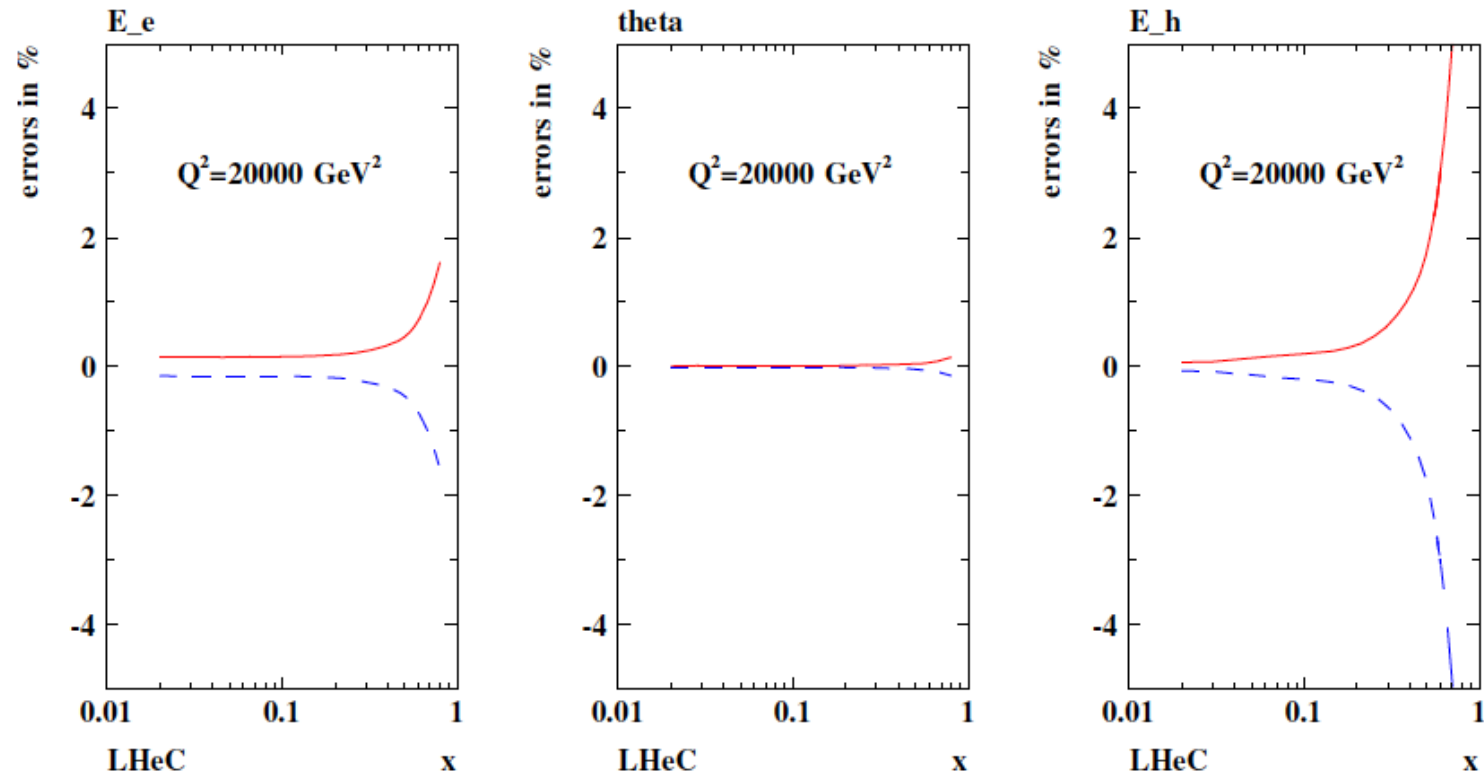


Figure 3.3: Neutral current cross section errors, calculated for $60 \times 7000 \text{ GeV}^2$ unpolarised e^-p scattering, resulting from scale uncertainties of the scattered electron energy $\delta E'_e/E'_e = 0.1\%$, of its polar angle $\delta\theta_e = 0.1 \text{ mrad}$ and the hadronic final state energy $\delta E_h/E_h = 0.5\%$, at large $Q^2 = 20000 \text{ GeV}^2$ and correspondingly large x . Note that the characteristic behaviour of the relative uncertainty at large x , i.e. to diverge $\propto 1/(1-x)$, is independent of Q^2 , i.e. persistently observed at $Q^2 = 200000 \text{ GeV}^2$ for example too.

NC Cross Section Correlated Uncertainties ($Q^2=2 \text{ GeV}^2$)

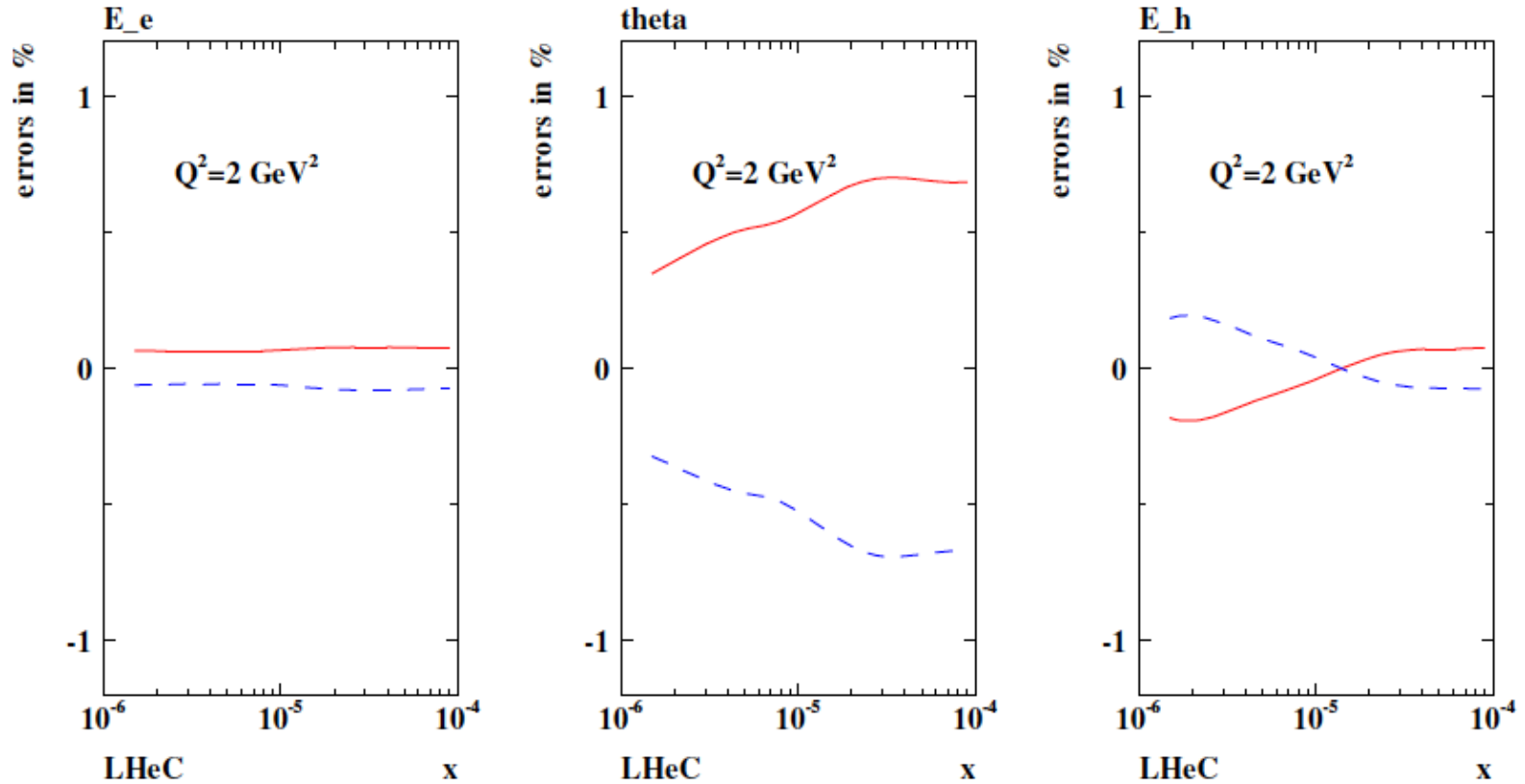


Figure 3.2: Neutral current cross section errors, calculated for $60 \times 7000 \text{ GeV}^2$, resulting from scale uncertainties of the scattered electron energy $\delta E'_e/E'_e = 0.1 \%$, of its polar angle $\delta\theta_e = 0.1 \text{ mrad}$ and the hadronic final state energy $\delta E_h/E_h = 0.5 \%$, at low $Q^2 = 2 \text{ GeV}^2$ and correspondingly low x .

Framework and Definitions

PDF+EW-fit

- PDF fit in NNLO precision
- ZM-VFNS using QCDNUM
- 13 free PDF parameters

EW calculations

- 1-loop EW corrections
- On-shell parameters are:
($\alpha_{em}, m_Z, m_W, \Delta r$) with
 $\Delta r = \Delta r(\alpha_{em}, m_W, m_Z, m_t, m_H, \dots)$

$$g_A^q = \sqrt{\rho_{\text{NC},q}} I_{L,q}^3,$$

$$g_V^q = \sqrt{\rho_{\text{NC},q}} \left(I_{L,q}^3 - 2Q_q \kappa_{\text{NC},q} \sin^2 \theta_W \right)$$

$$W_2^- = x \left(\rho_{\text{CC},eq}^2 U + \rho_{\text{CC},e\bar{q}}^2 \bar{D} \right), \quad xW_3^- = x \left(\rho_{\text{CC},eq}^2 U - \rho_{\text{CC},e\bar{q}}^2 \bar{D} \right)$$

$$W_2^+ = x \left(\rho_{\text{CC},eq}^2 \bar{U} + \rho_{\text{CC},e\bar{q}}^2 D \right), \quad xW_3^+ = x \left(\rho_{\text{CC},e\bar{q}}^2 D - \rho_{\text{CC},eq}^2 \bar{U} \right)$$

cf H1 to be published, Z Zhang this conference

MSbar: Spiesberger/Dittmaier, in preparation

- m_t and m_H enter through loop-corrections (Δr)
- $\sin^2 \theta_W$ and g_f are calculated quantities
- More general, also vector and axial-vector couplings are 'free' parameters

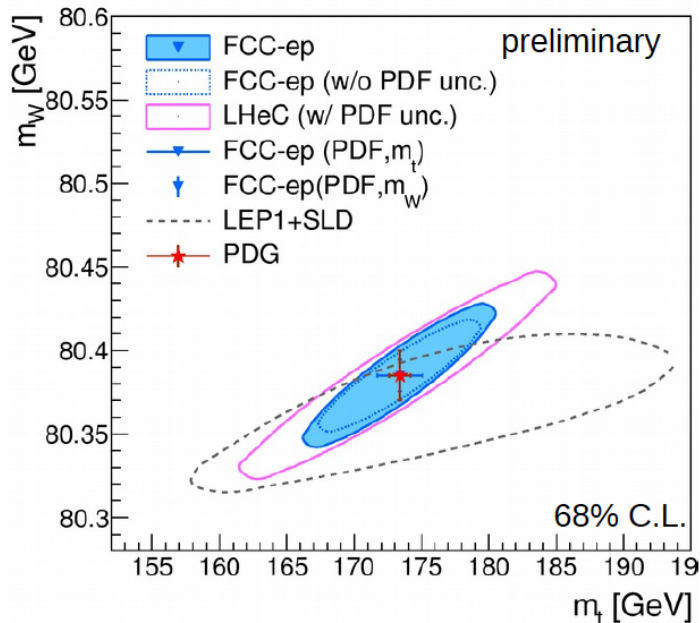
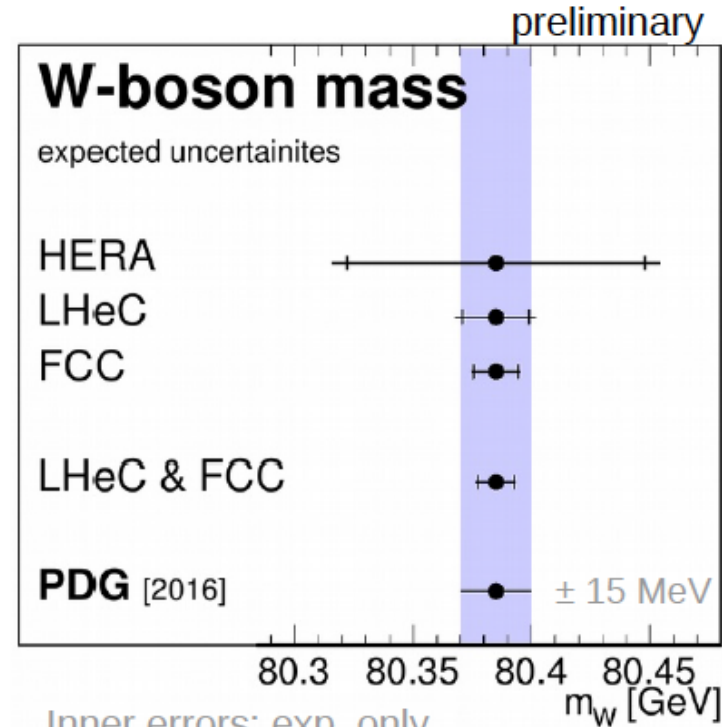
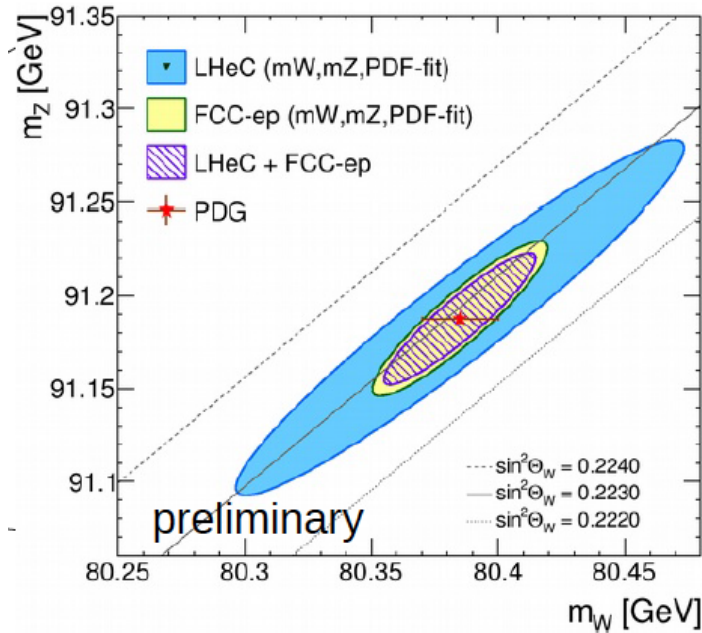
New: parameterise effective h.o. corrections as deviations ρ' and κ'

$$\rho'_{\text{NC}} \rightarrow \rho'_{\text{NC}} \rho_{\text{NC}}$$

$$\kappa'_{\text{NC}} \rightarrow \kappa'_{\text{NC}} \kappa_{\text{NC}}$$

$$\rho'_{\text{CC}} \rightarrow \rho'_{\text{CC}} \rho_{\text{CC}}$$

W,Z, top [loop] Masses [from inclusive data only] at LHeC/FCCeh



- Inner errors: exp. only
 Outer errors: exp. + PDF
- LHeC $\pm 14_{(\text{exp})} 10_{(\text{PDF})}$ MeV
 - FCC $\pm 9_{(\text{exp})} 4_{(\text{PDF})}$ MeV

In ep PDFs not limiting factor
 For pp, LHC predict M_W to 2.8 MeV
 with LHeC PDFs (S Camarda)

MZ similar. H and t from loops for consistency
 measurable directly with much better accuracy.

Masses from inclusive NC+CC Cross Sections

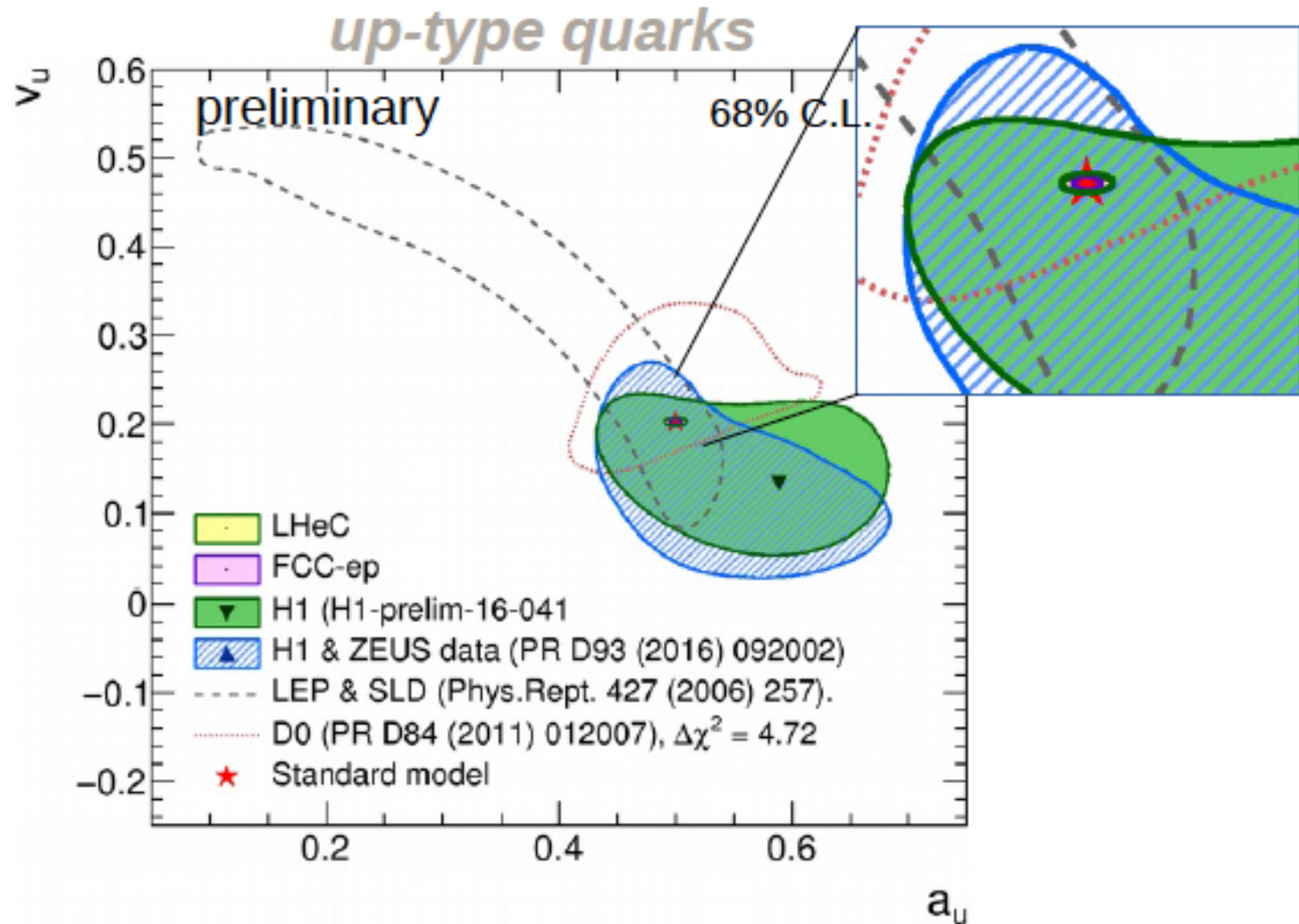
Parameter	HERA	LHeC	FCC-eh
Δm_W [MeV]	$\pm 63_{(\text{exp})} 29_{(\text{pdf})}$	$\pm 14_{(\text{exp})} 10_{(\text{pdf})}$	$\pm 9_{(\text{exp})} 4_{(\text{pdf})}$
Δm_Z [MeV]	$\pm 56_{(\text{exp})} 25_{(\text{pdf})}$	$\pm 16_{(\text{exp})} 10_{(\text{pdf})}$	$\pm 16_{(\text{exp})} 10_{(\text{pdf})}$
Δm_t [GeV]	$\pm 10_{(\text{exp})} 5_{(\text{pdf})}$	$\pm 2.6_{(\text{exp})} 1.7_{(\text{pdf})}$	$\pm 1.7_{(\text{exp})} 0.5_{(\text{pdf})}$
Δm_H [GeV]	$> \mathcal{O}(100 \text{ GeV})$	$\pm 31_{(\text{exp})} 22_{(\text{pdf})}$	$\pm 20_{(\text{exp})} 4_{(\text{pdf})}$

Table 4: Summary of electroweak parameters from HERA-II data and LHeC and FCC-ep simulated data.

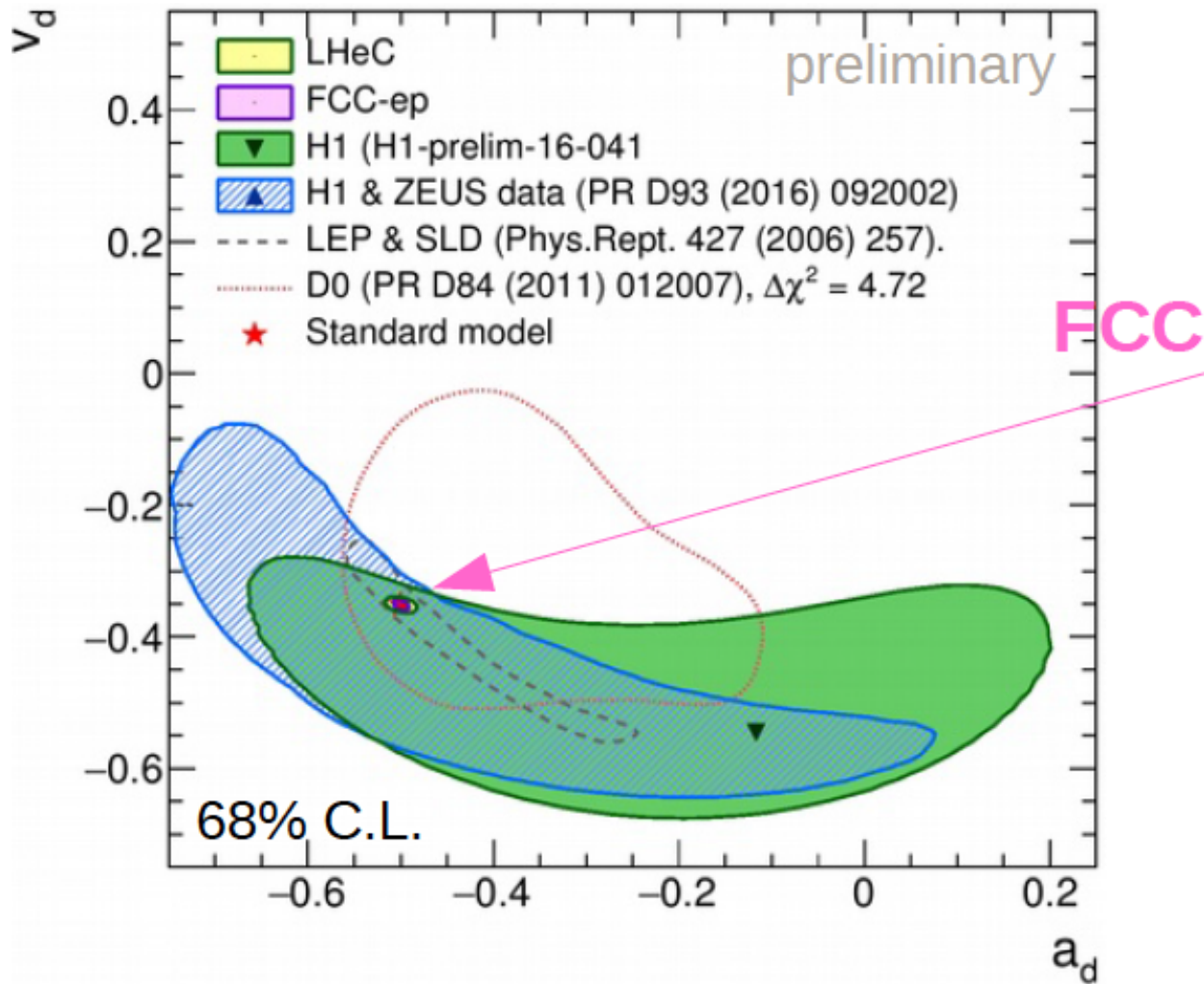
OMS: W,Z direct. top. Higgs through loops

1987 expected 100 MeV for MW [JB, MK, TR] at HERA

Light up Quark NC Couplings



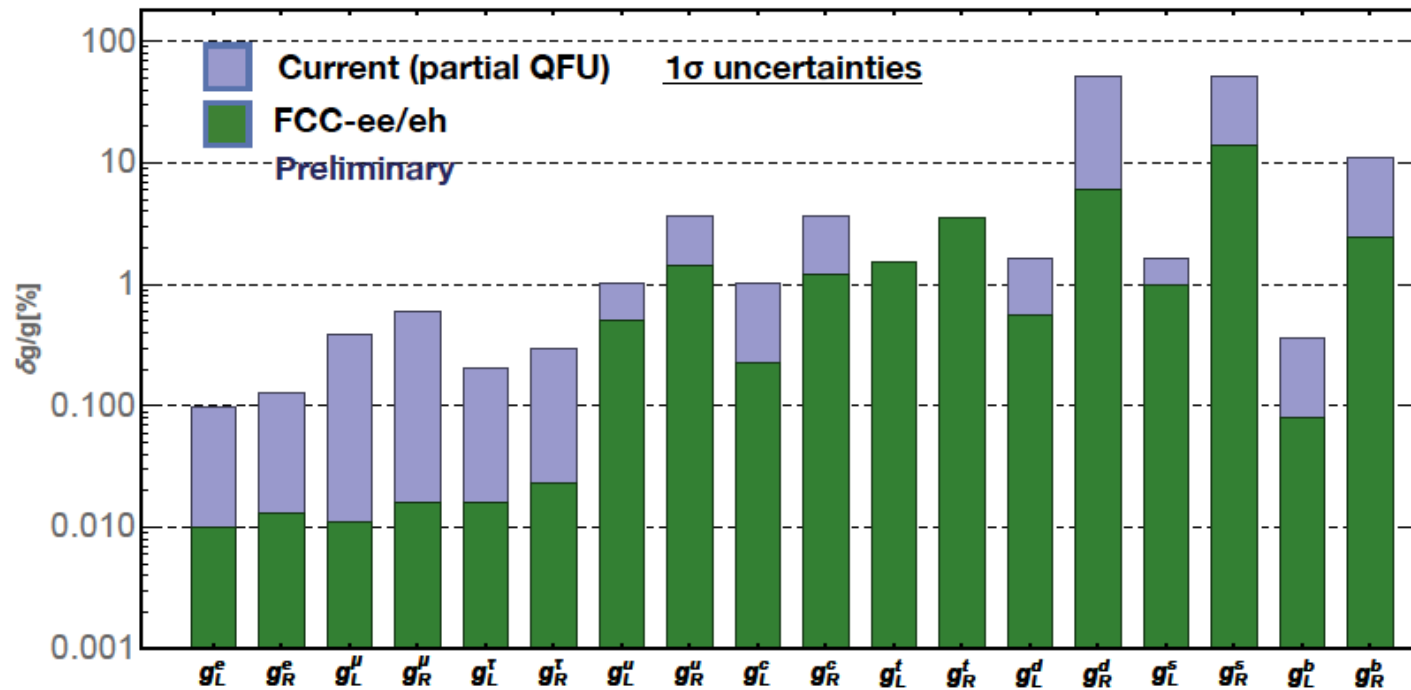
Light down Quark NC Couplings



First preliminary Global Electroweak Analysis FCC ee+eh

J de Blas (Amsterdam FCC week)

- Global fit to electroweak precision measurements at FCC-ee + FCC-eh



No Fermion flavour universality assumed

Independent info about all 3 SM fermion families

FCC-eh (and LHeC) has much more information to provide than up + down NC couplings.

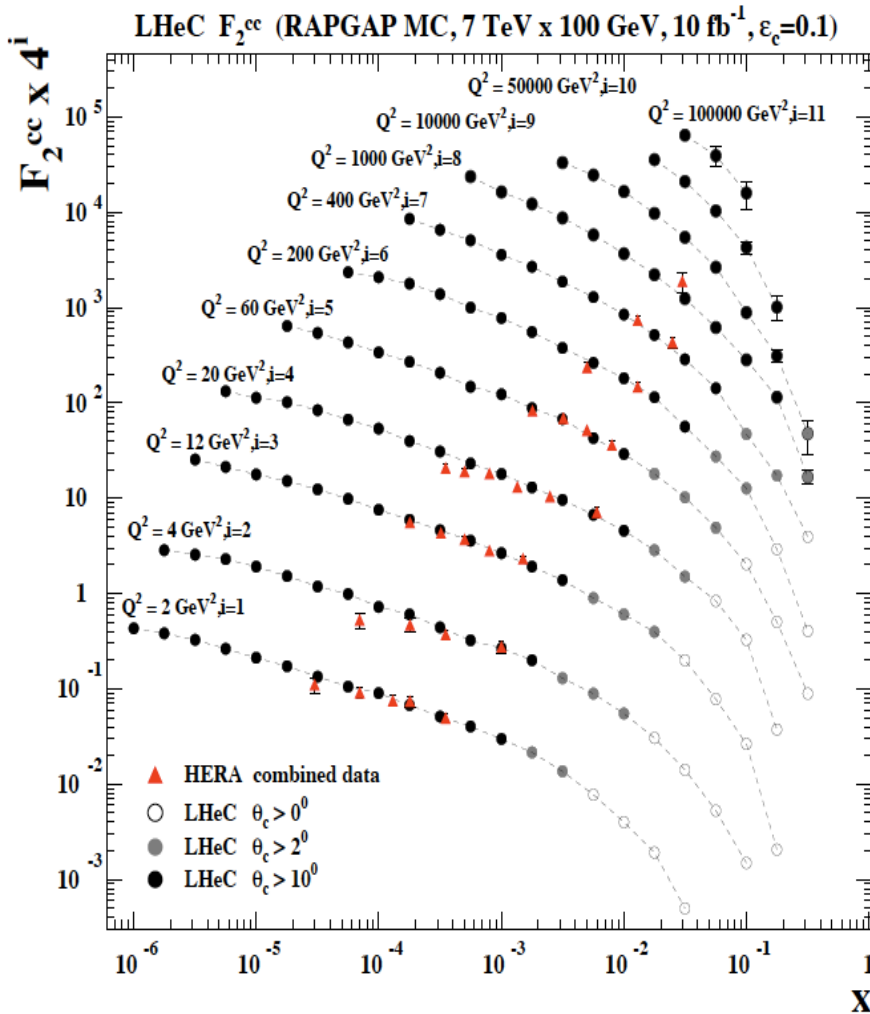
-We can measure the interference parts of F_2^{cc} and F_2^{bb} → get v and a couplings for c, b

-Scale dependence of $\sin^2\Theta_W(Q)$ for $Q \sim 300$ MeV (PERLE) to 1 TeV (LHeC), 3 TeV (FCCeh)

[low scales: elastic lepton-nucleon scattering MK T Riemann, Z Phys C8 (81) 239: Jlab, MESA]

Charm F_2^{cc} and Mass

LHeC CDR arXiv:1206.2913



HERA 0.0005/2.5 .. 0.05/2000 GeV²
 LHeC 0.00001/1 .. 0.2/200000 GeV²

$\epsilon(c)$ assumed 10%, 1% light background, ~3% $\delta(\text{syst})$

Heavy Flavour with LHeC

Beam spot (in xy): 7 μm
 Impact parameter: better than 10 μm
 Modern Silicon detectors, no pile-up
 Higher E, L, Acceptance, ϵ , than at HERA
 → Huge improvements predicted

	HERA	LHeC
$m_c(m_c)/\text{GeV}$	1.26	?
$\delta(\text{exp})$	0.05	0.003
$\delta(\text{mod})$	0.03	~0.002
$\delta(\text{par})$	0.02	~0.002
$\delta(\alpha_s)$	0.02	0.001

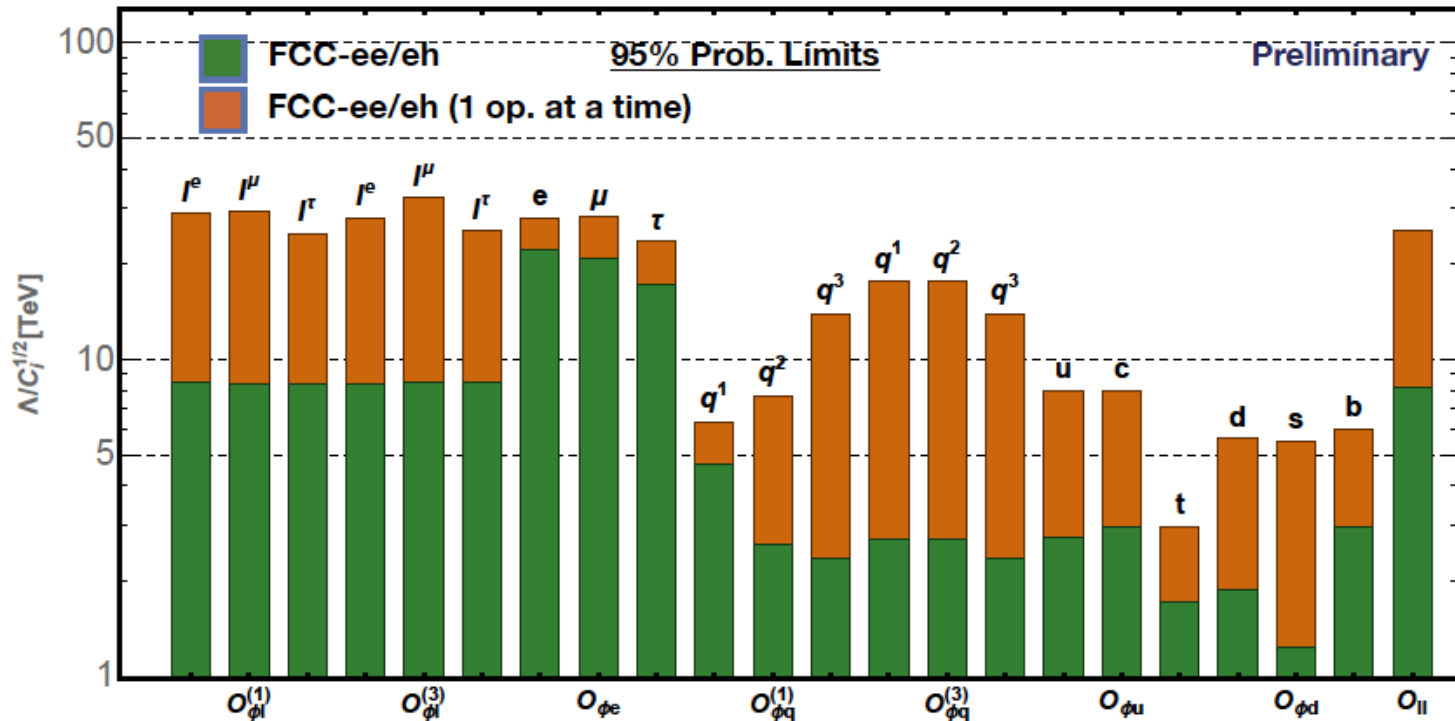
LHeC determines strong coupling to 0.1%
 High precision PDF data will reduce the mod and par errors by a very large amount.

Determination of charm mass to 3 MeV:
 crucial for M_W in pp or $H \rightarrow cc$ in ep
 cf also NNPDF3.1 (arXiv:1706.00428) and refs

First preliminary Global Electroweak Analysis FCC ee+eh

Combination of ee and ep; Test of SM to O(20) TeV; a Model also for LHeC + e⁺e⁻ wherever

- Global fit to electroweak precision measurements at FCC-ee + FCC-eh



No Fermion flavour universality assumed

Independent info about all 3 SM fermion families

$$\mathcal{L}_{\text{NC}} = -\frac{e}{s_c} (1 + \delta^U g_{\text{NC}}) Z_\mu \sum_\psi \bar{\psi}^i \gamma^\mu \left[(g_L^\psi \delta_{ij} + (\delta^D g_L^\psi)_{ij}) P_L + (g_R^\psi \delta_{ij} + (\delta^D g_R^\psi)_{ij}) P_R + \delta^Q g_{\text{NC}} \delta_{ij} \right] \psi^j$$

$$O_{\phi f}^{(1)} = (\phi^\dagger i \overleftrightarrow{D}_\mu \phi) (\bar{f} \gamma^\mu f) \quad O_{\phi f}^{(3)} = (\phi^\dagger i \overleftrightarrow{D}_\mu^a \phi) (\bar{f} \gamma^\mu \sigma_a f) \quad \dots \quad \text{J De Blas}$$

Definitions of ρ' and κ' s

Neutral current

Universal higher-order corrections are taken into account by Q^2 -dependent form factors ρ_{NC} and κ_{NC} . Many extensions of the Standard Model predict modifications of the weak neutral-current couplings. These can be described conveniently by introducing additional parameters ρ'_{NC} and κ'_{NC} , which can be also considered to be Q^2 dependent:

$$g_A^f = \sqrt{\rho_{\text{NC},f} \rho'_{\text{NC},f}} I_{L,f}^3, \quad (1)$$

$$g_V^f = \sqrt{\rho_{\text{NC},f} \rho'_{\text{NC},f}} \left(I_{L,f}^3 - 2Q_f \kappa_{\text{NC},f} \kappa'_{\text{NC},f} \sin^2 \theta_W \right). \quad (2)$$

The estimated relative uncertainties of the ρ'_{NC} or κ'_{NC} achieved with the LHeC or FCC-eh data, can also be interpreted as the relative uncertainty of a direct determination of the ρ_{NC} parameters or $\sin^2 \theta_w^{\text{eff}}$.

Charged current

Higher-order EW corrections to the CC cross sections are collected in form factors $\rho_{\text{CC},eq/e\bar{q}}$. Similarly as for NC, modifications of the SM formalism can be expressed by introducing the additional ρ'_{CC} parameters:

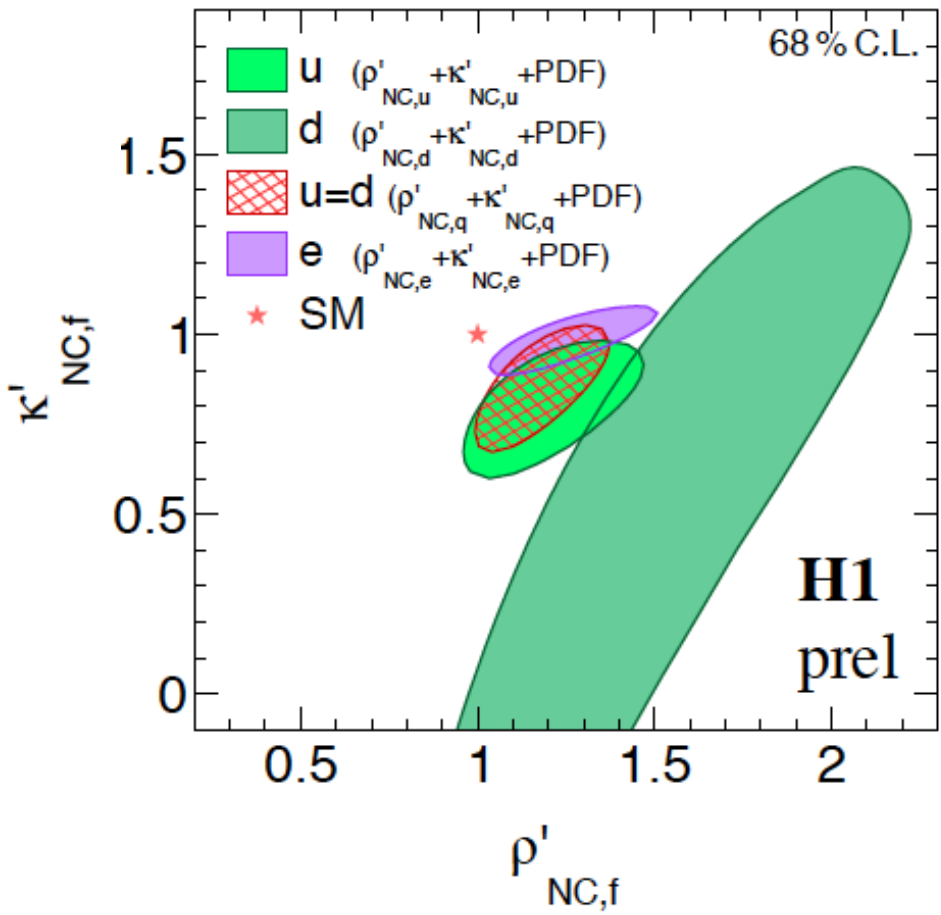
$$W_2^- = x \left((\rho_{\text{CC},eq} \rho'_{\text{CC},eq})^2 U + (\rho_{\text{CC},e\bar{q}} \rho'_{\text{CC},e\bar{q}})^2 \bar{D} \right), \quad (3)$$

$$xW_3^- = x \left((\rho_{\text{CC},eq} \rho'_{\text{CC},eq})^2 U - (\rho_{\text{CC},e\bar{q}} \rho'_{\text{CC},e\bar{q}})^2 \bar{D} \right), \quad (4)$$

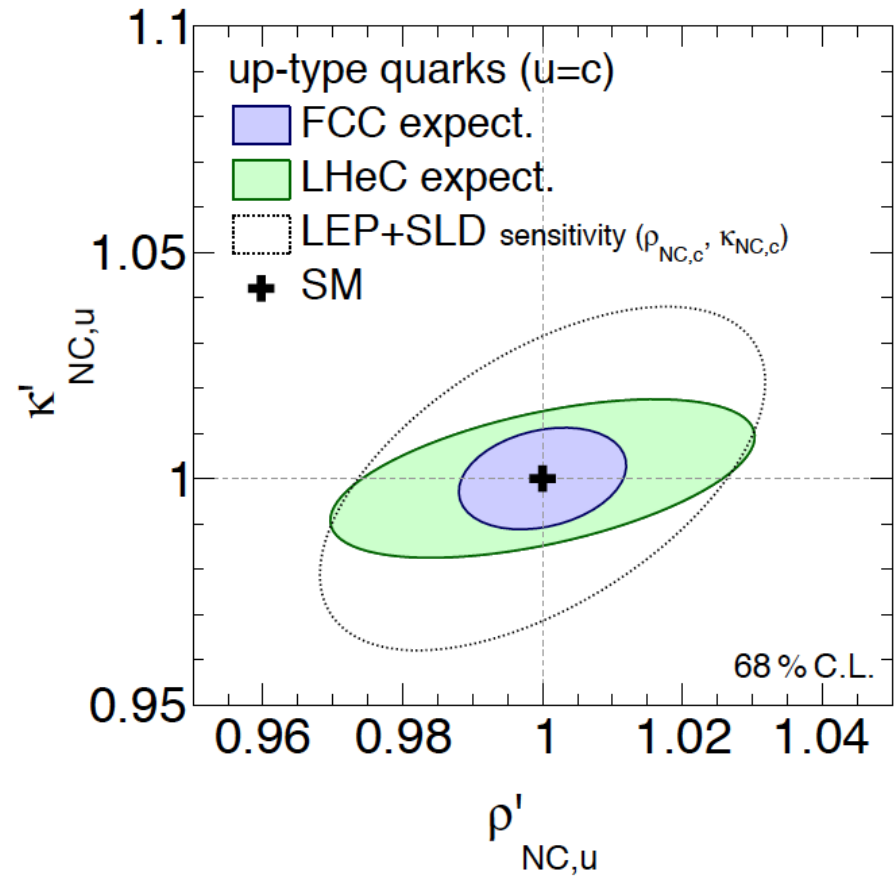
$$W_2^+ = x \left((\rho_{\text{CC},eq} \rho'_{\text{CC},eq})^2 \bar{U} + \rho_{\text{CC},e\bar{q}} \rho'_{\text{CC},e\bar{q}} \bar{D} \right), \quad (5)$$

$$xW_3^+ = x \left((\rho_{\text{CC},e\bar{q}} \rho'_{\text{CC},e\bar{q}})^2 D - \rho_{\text{CC},eq} \rho'_{\text{CC},eq} \bar{U} \right). \quad (6)$$

LHeC/FCCeh h.o. Tests of Electroweak SM

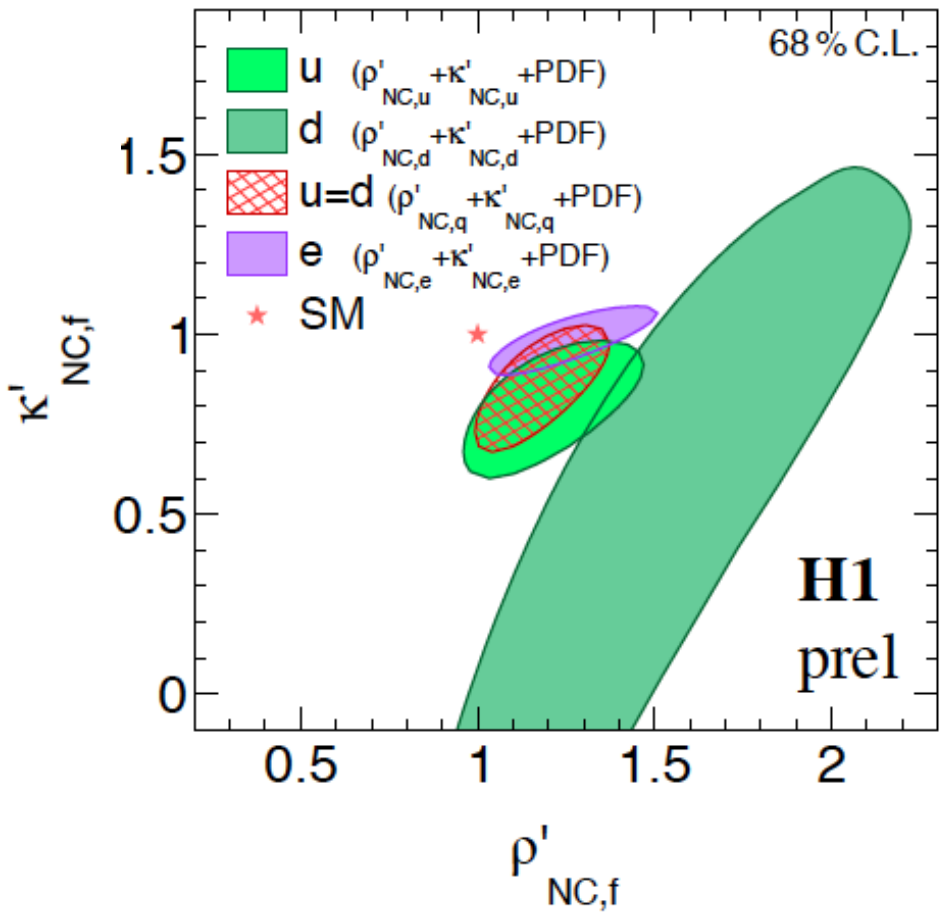


Up Quarks



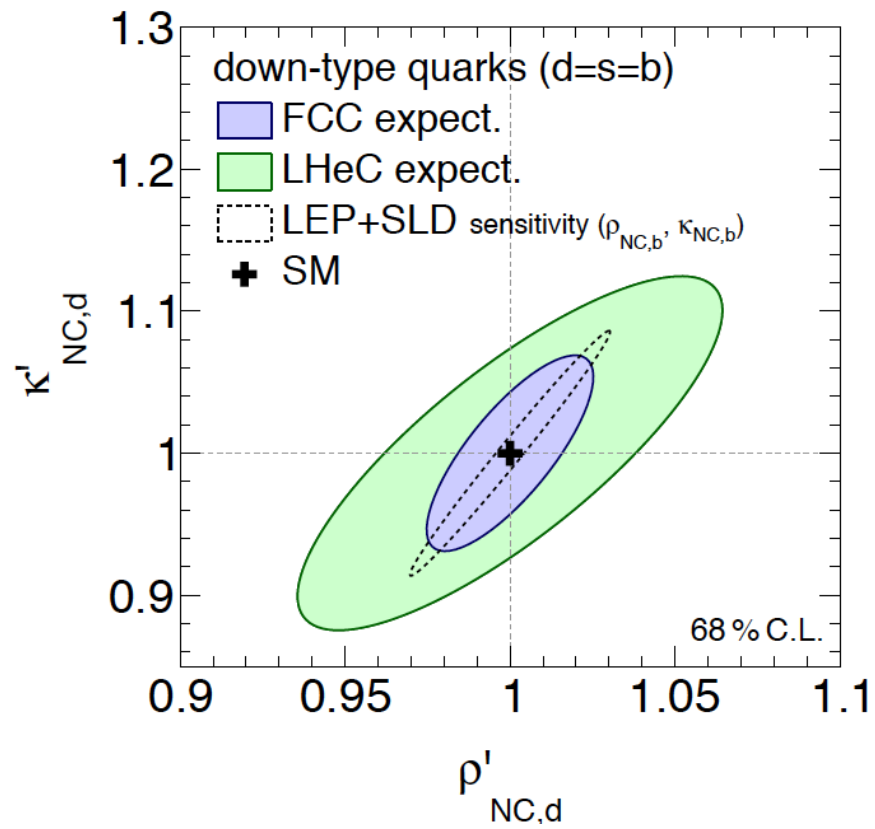
Z Zhang, H1, This Conference.

LHeC/FCCeh h.o. Tests of Electroweak SM



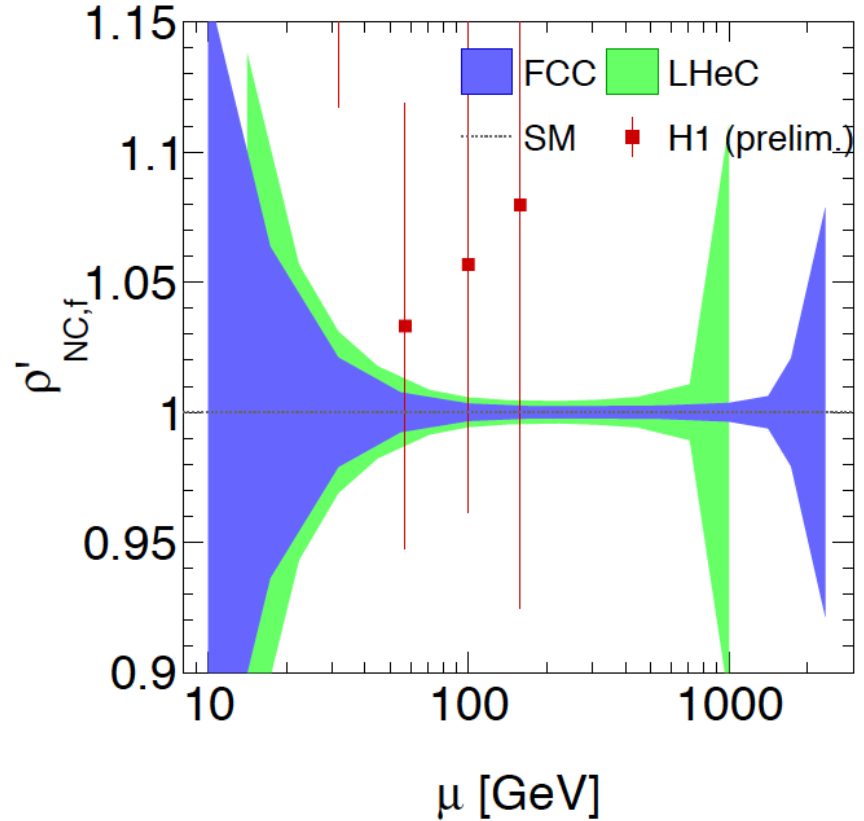
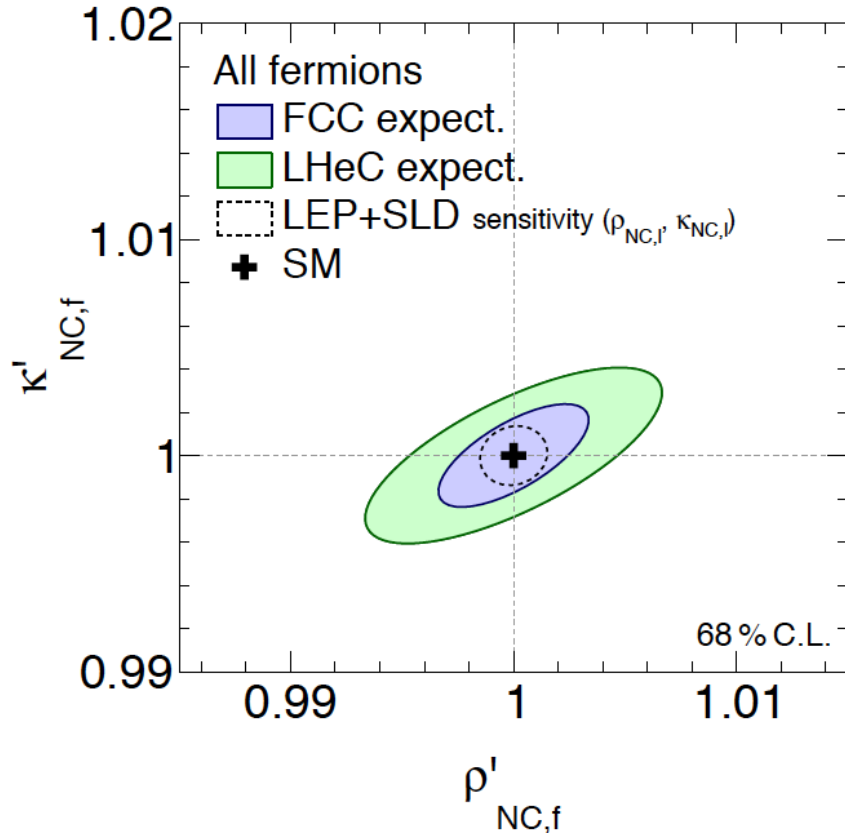
Z Zhang, H1, This Conference.

Down Quarks



LHeC/FCCeh h.o. Tests of Electroweak SM

Neutral Currents

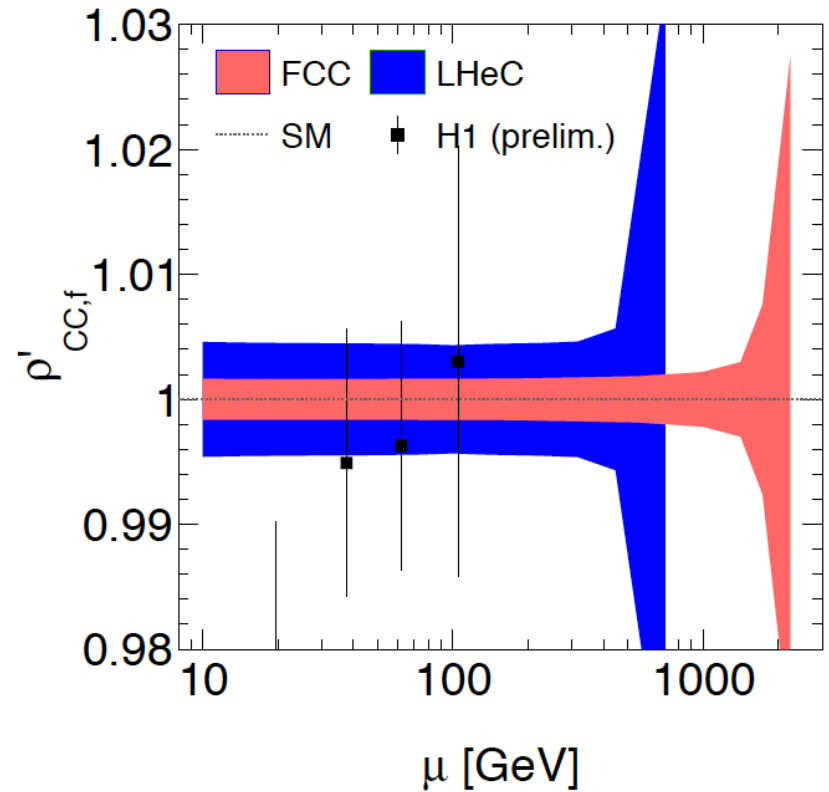
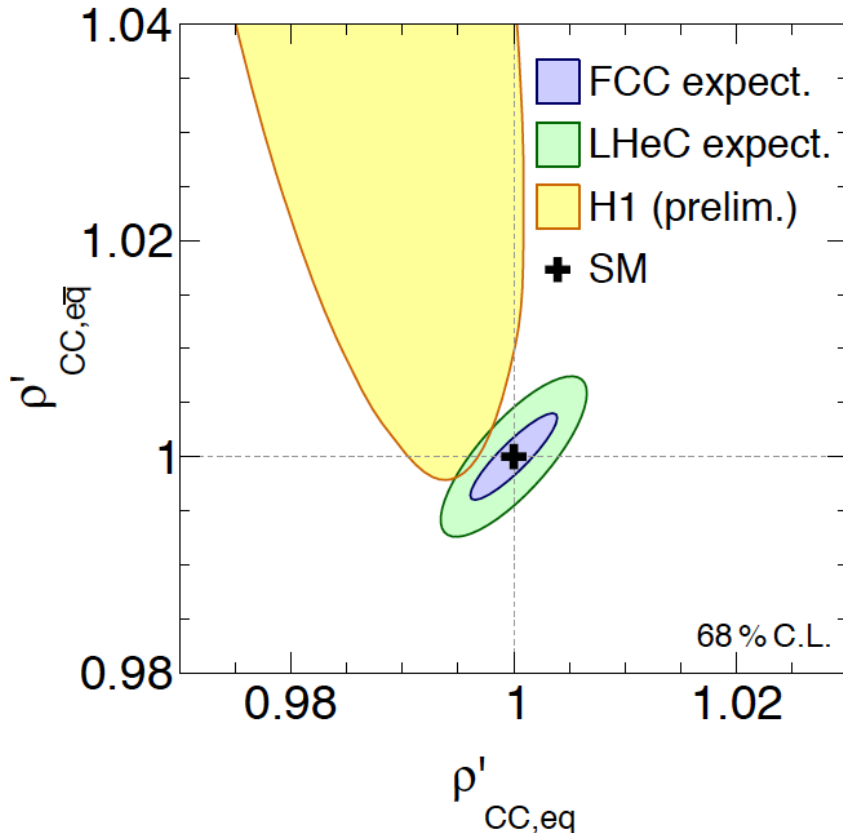


Test of eweak SM in NC to permille level. κ uncertainty describes $\sin^2\theta$ sensitivity.

Test of eweak theory in space like configuration. Precision vs scale $\mu = \sqrt{Q^2}$

LHeC/FCCeh h.o. Tests of Electroweak SM

Charged Currents



Test of eweak SM in CC to few permille level. Quark-Antiquark distinction.

Note that CC at H1 is large x dominated \rightarrow reduced sensitivity to eqbar

Test of eweak theory in space like configuration. Precision vs scale $\mu=VQ^2$

Summary

Study of electroweak effects in NC and CC inclusive cross sections performed.

Full consideration of experimental, syst+stat uncertainties [as in α_s analysis].

Joint QCD (PDF) and electroweak analysis. PDFs do not dominate eweak tests.

$s=Q_{\max}^2 = 4E_e E_p = 1.7 \text{ TeV}^2$ (LHeC) and $12 \text{ TeV}^2 \gg M_{W,Z}^2$. Very large luminosity \rightarrow

High precision measurements \rightarrow New laboratory for testing EW SM at new scales.

Initial determination of light quark couplings done, to 1% precision.

Novel parameterisation of h.o. effects in NC and CC couplings, including $\sin^2\theta$.

Measurement of scale dependence with unique and unprecedented precision.

First joint EFT and coupling fit analysis done for future ee and ep colliders (FCC).

Next: Formulation of DIS in $\overline{\text{MS}}$, Tests for c,b (e) couplings. LHeC & e^+e^- , ...

backup

PDFs and their effect on electroweak physics

FCC-eh and LHeC: Input: high precision (stat+syst) data on
Neutral Current ($x: 10^{-6}-1; Q^2: 1-10^6$) Charged Current ($10^{-4}-1; 100-10^6$)
Tagging of Charm and Beauty with high precision and coverage. ep (eD)

Completely new PDF Programme

Determine ALL pdfs in a coherent way + the strong coupling to 0.1% accuracy
No higher twists, no nuclear corrections, no symmetry assumptions, N³LO

→ \bar{u} , \bar{v} , \bar{d} , \bar{v} , s , c , b , t , xg and α_s

This essentially removes the PDF uncertainties on the electroweak variables, in ep but as well for pp.

For the Higgs this means that ep can turn pp into a precision Higgs facility

Definitions (J De Blas, Amsterdam FCC week)

- EWPO sensitive to modifications of NC couplings**

$$\mathcal{L}_{\text{NC}} = -\frac{e}{sc} (1 + \delta^U g_{\text{NC}}) Z_\mu \sum_\psi \bar{\psi}^i \gamma^\mu \left[\left(g_L^\psi \delta_{ij} + (\delta^D g_L^\psi)_{ij} \right) P_L + \left(g_R^\psi \delta_{ij} + (\delta^D g_R^\psi)_{ij} \right) P_R + \delta^Q g_{\text{NC}} \delta_{ij} \right] \psi^j$$

Flavor non-universal contributions

$$\begin{aligned} \delta^D g_L^\nu &= -\frac{1}{2} \left(C_{\phi l}^{(1)} \mp C_{\phi l}^{(3)} \right) \frac{v^2}{\Lambda^2}, & \delta^D g_R^e &= -\frac{1}{2} C_{\phi e}^{(1)} \frac{v^2}{\Lambda^2} \\ \delta^D g_L^u &= -\frac{1}{2} \left(C_{\phi q}^{(1)} \mp C_{\phi q}^{(3)} \right) \frac{v^2}{\Lambda^2}, & \delta^D g_R^d &= -\frac{1}{4} C_{\phi d}^{(1)} \frac{v^2}{\Lambda^2} \end{aligned}$$

Flavor-universal contributions

$$\begin{aligned} \delta^U g_{\text{NC}} &= -\frac{1}{2} \left[\Delta_{GF} + \frac{C_{\phi D}}{2} \right] \frac{v^2}{\Lambda^2} \\ \delta^Q g_{\text{NC}} &= -Q \left(\frac{sc}{c^2 - s^2} C_{\phi WB} + \frac{s^2 c^2}{c^2 - s^2} \left[\Delta_{GF} + \frac{C_{\phi D}}{2} \right] \right) \frac{v^2}{\Lambda^2} \end{aligned}$$

Indirect effect associated to modifications in μ decay (G_F)

$$\Delta_{GF} = \left(C_{\phi l}^{(3)} \right)_{22} + \left(C_{\phi l}^{(3)} \right)_{11} - (C_U)_{1221}$$

10 Operators

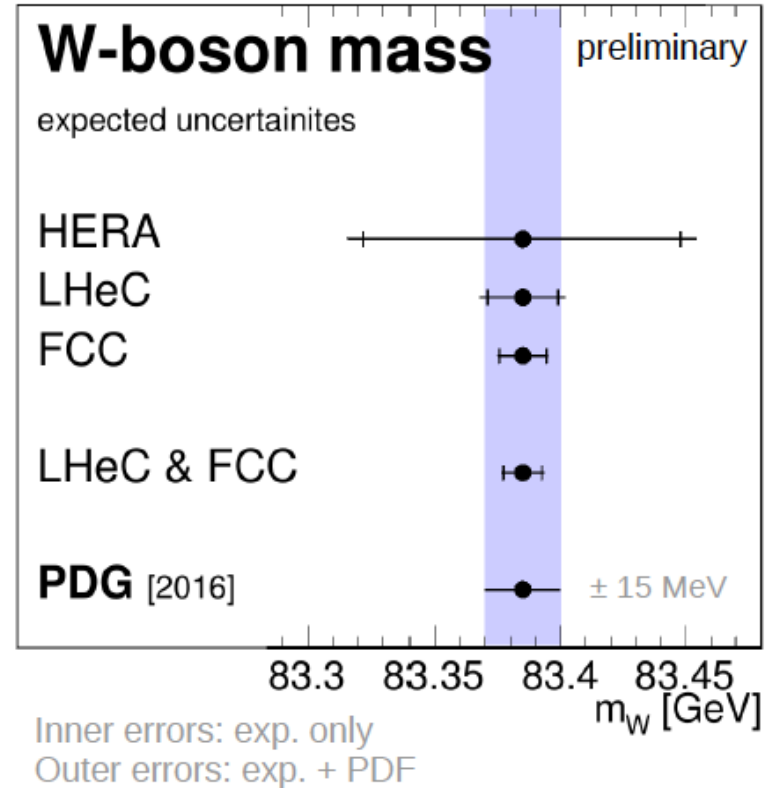
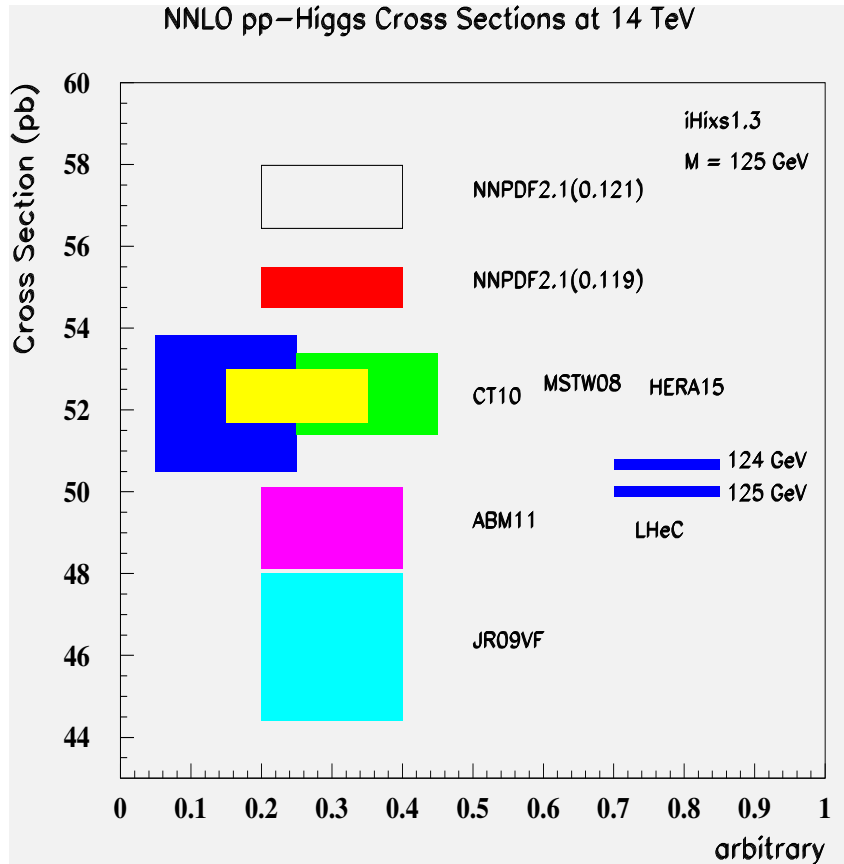
$$\begin{aligned} \mathcal{O}_{\phi f}^{(1)} &= (\phi^\dagger i \overleftrightarrow{D}_\mu \phi) (\bar{f} \gamma^\mu f) \\ \mathcal{O}_{\phi f}^{(3)} &= (\phi^\dagger i \overleftrightarrow{D}_\mu^a \phi) (\bar{f} \gamma^\mu \sigma_a f) \end{aligned}$$

$$\mathcal{O}_{\phi D} = |\phi^\dagger i D_\mu \phi|^2$$

$$\mathcal{O}_{\phi WB} = (\phi^\dagger \sigma_a \phi) W_{\mu\nu}^a B^{\mu\nu}$$

$$\mathcal{O}_l = (\bar{l} \gamma_\mu l) (\bar{l} \gamma^\mu l)$$

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