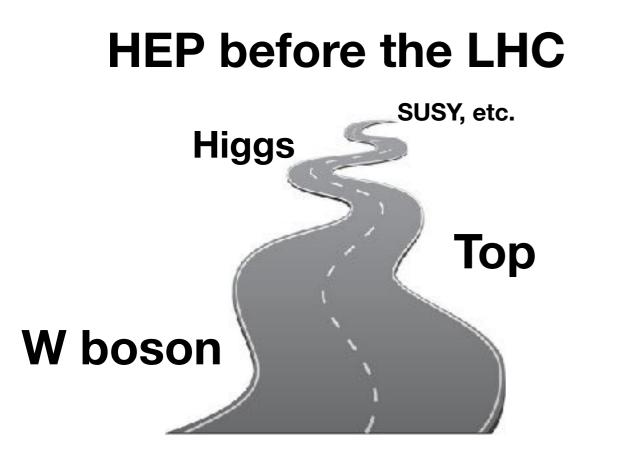
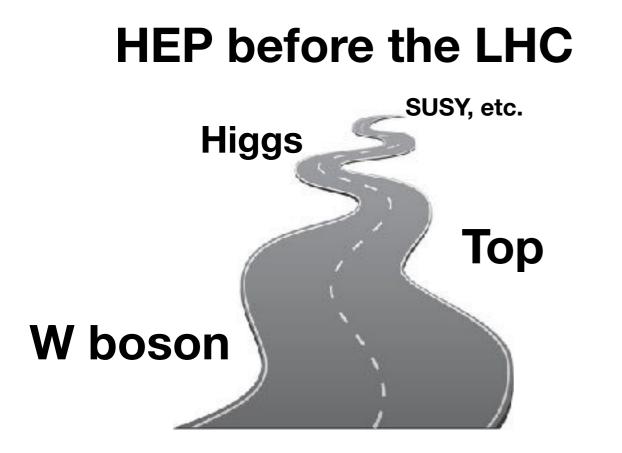
# The CLIC Physics Potential

Andrea Wulzer



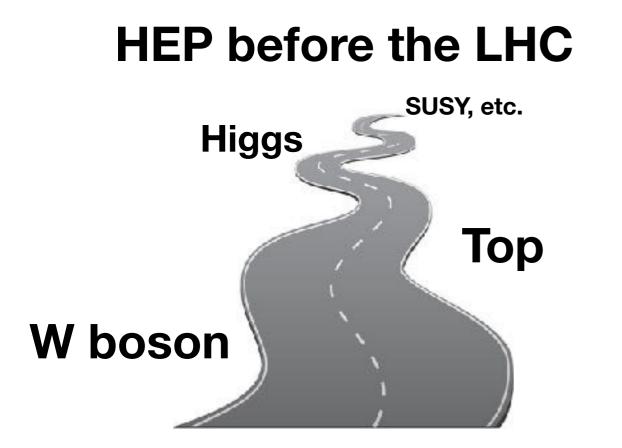


**HEP** before the F.C.



#### **HEP before the F.C.**

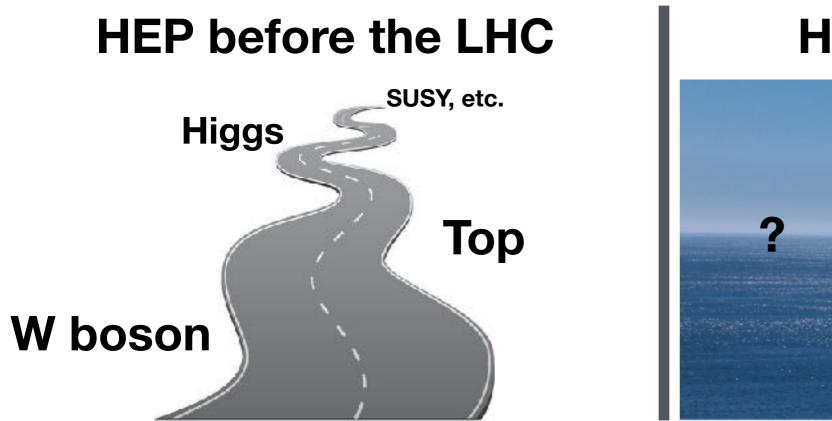




#### **HEP before the F.C.**







#### HEP before the F.C.



# Particle physics is not validation anymore, rather it is exploration of unknown territories \*

\* Not necessarily a bad thing. Columbus left for his trip just because he had no idea of where he was going !!

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy But drawing implications requires BSM.

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy

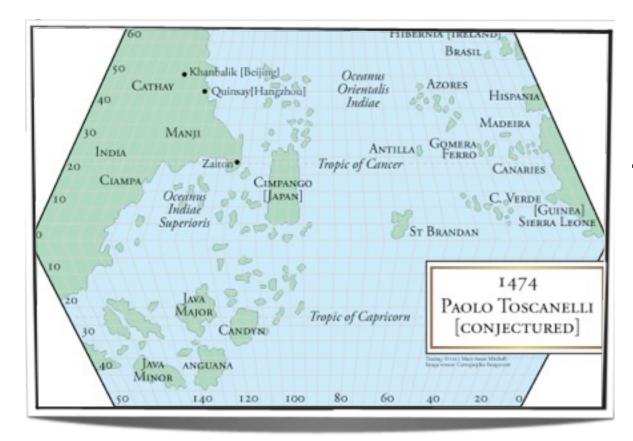
#### But drawing implications requires BSM.

BSM must set a destination, not to go around in circles: BSM = **draw maps** to guide us in F.C. ocean

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy

#### But drawing implications requires BSM.

BSM must set a destination, not to go around in circles: BSM = **draw maps** to guide us in F.C. ocean



Columbus had Toscanelli's map.

It was terribly **wrong**, but **served the purpose** 

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy

#### But drawing implications requires BSM.

BSM must set a destination, not to go around in circles: BSM = **draw maps** to guide us in F.C. ocean

If N.P. is heavy, EFT map:

 $\mathcal{L} = \mathcal{L}_{\rm SM} + \mathcal{L}_{\rm BSM}^{d=6}$ 

operator estimate from structural BSM assumptions. **Different assumptions produce different maps** 

Measuring and comparing with SM predictions is a systematic, BSM-independent exploration strategy

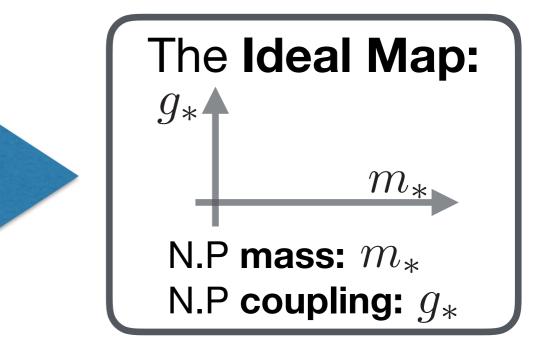
#### But drawing implications requires BSM.

BSM must set a destination, not to go around in circles: BSM = **draw maps** to guide us in F.C. ocean

#### If N.P. is heavy, EFT map:

$$\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \mathcal{L}_{\mathrm{BSM}}^{d=6}$$

operator estimate from structural BSM assumptions. **Different assumptions produce different maps** 

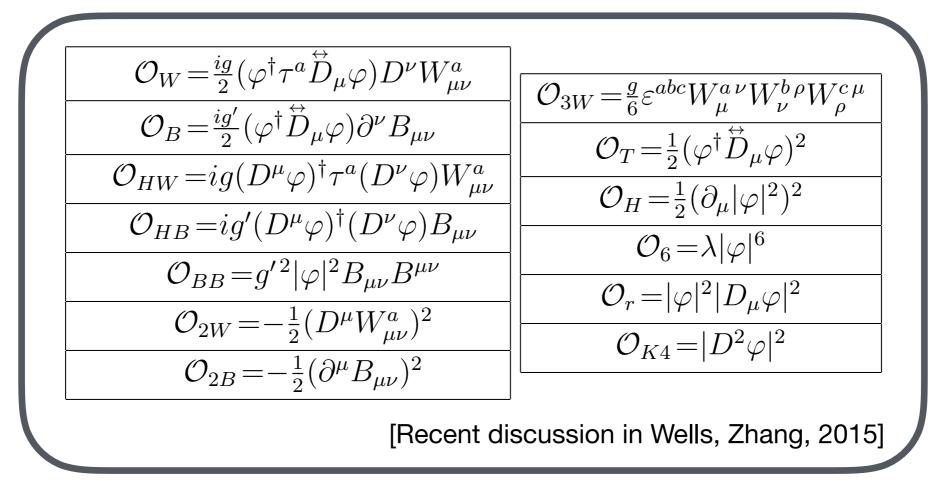


### The EFT Map

#### Dimension-6 operators classified:

#### Universal

BSM only coupled to SM bosons, negligible direct coupling to fermions



## The EFT Map

Dimension-6 operators classified:

#### Universal

BSM only coupled to SM bosons, negligible direct coupling to fermions

#### **Top-philic**

direct BSM coupling to top. Motivated by Naturalness and flavour.

$$\begin{aligned} & Q_{\varphi t} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{t} \gamma^{\mu} t) \\ & \mathcal{N}_{tB} = (\bar{t} \gamma^{\mu} t) (\bar{e} \gamma_{\mu} e + \frac{1}{2} \bar{l} \gamma_{\mu} l) \\ & Q_{t\varphi} = (\varphi^{\dagger} \varphi) (\bar{q} t \widetilde{\varphi}) \\ & Q_{tB} = (\bar{q} \sigma^{\mu\nu} t) \widetilde{\varphi} B_{\mu\nu} \\ & Q_{\varphi q}^{(1)} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{q} \gamma^{\mu} q) \\ & Q_{\varphi q}^{(3)} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\bar{q} \tau^{I} \gamma^{\mu} q) \\ & Q_{tW} = (\bar{q} \sigma^{\mu\nu} t) \tau^{I} \widetilde{\varphi} W_{\mu\nu}^{I} \\ & \mathcal{N}_{qB} = (\bar{q} \gamma^{\mu} q) (\bar{e} \gamma_{\mu} e + \frac{1}{2} \bar{l} \gamma_{\mu} l) \\ & \mathcal{N}_{qW} = (\bar{q} \tau^{I} \gamma^{\mu} q) (\bar{l} \tau^{I} \gamma_{\mu} l) \end{aligned}$$

## The EFT Map

Dimension-6 operators classified:

#### Universal

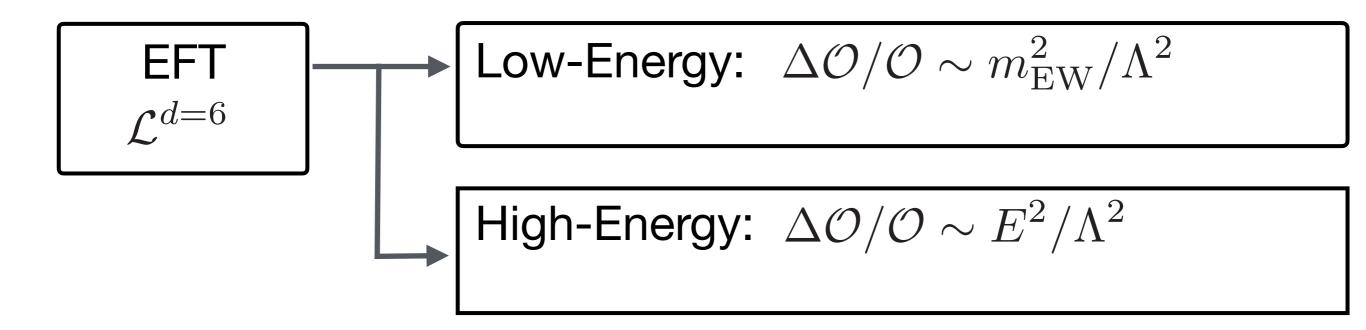
BSM only coupled to SM bosons, negligible direct coupling to fermions

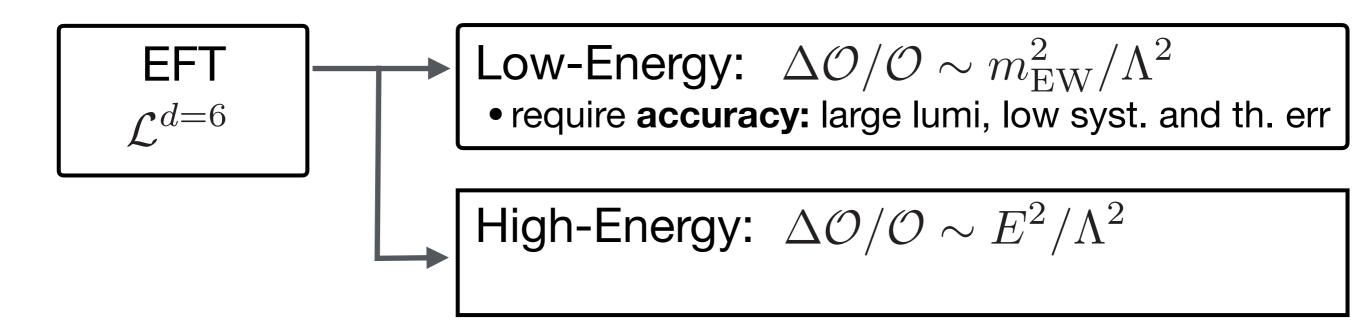
#### **Top-philic**

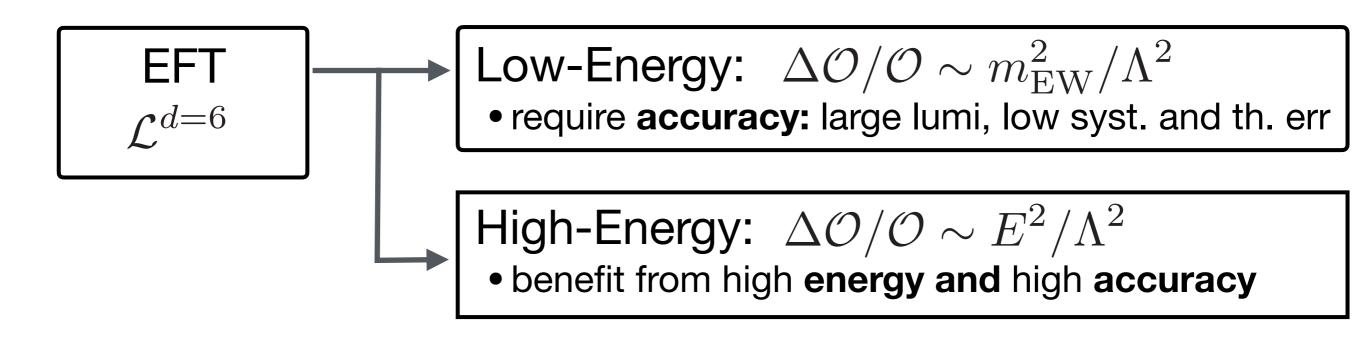
direct BSM coupling to top. Motivated by Naturalness and flavour.

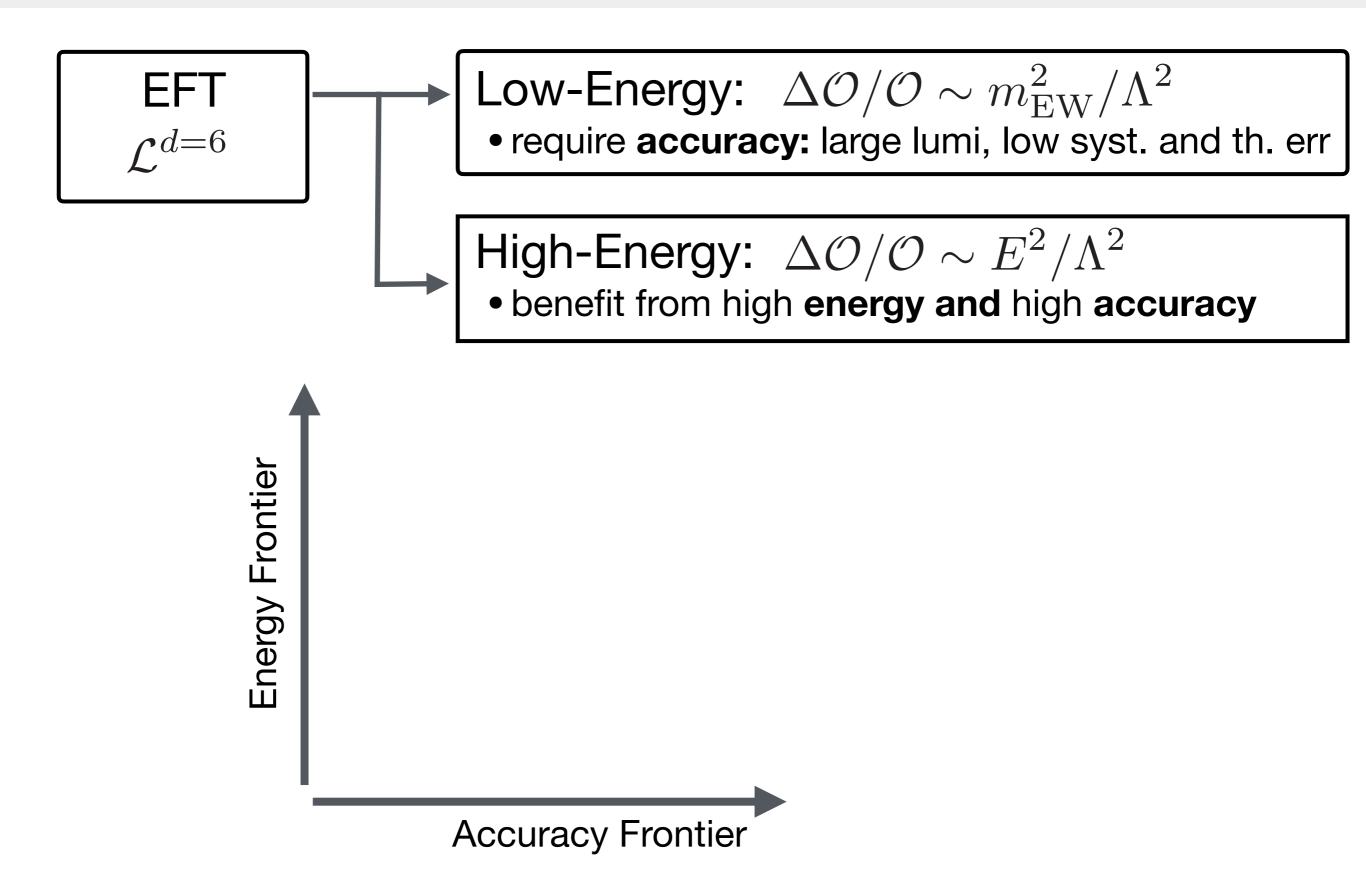
#### **Flavour-breaking**

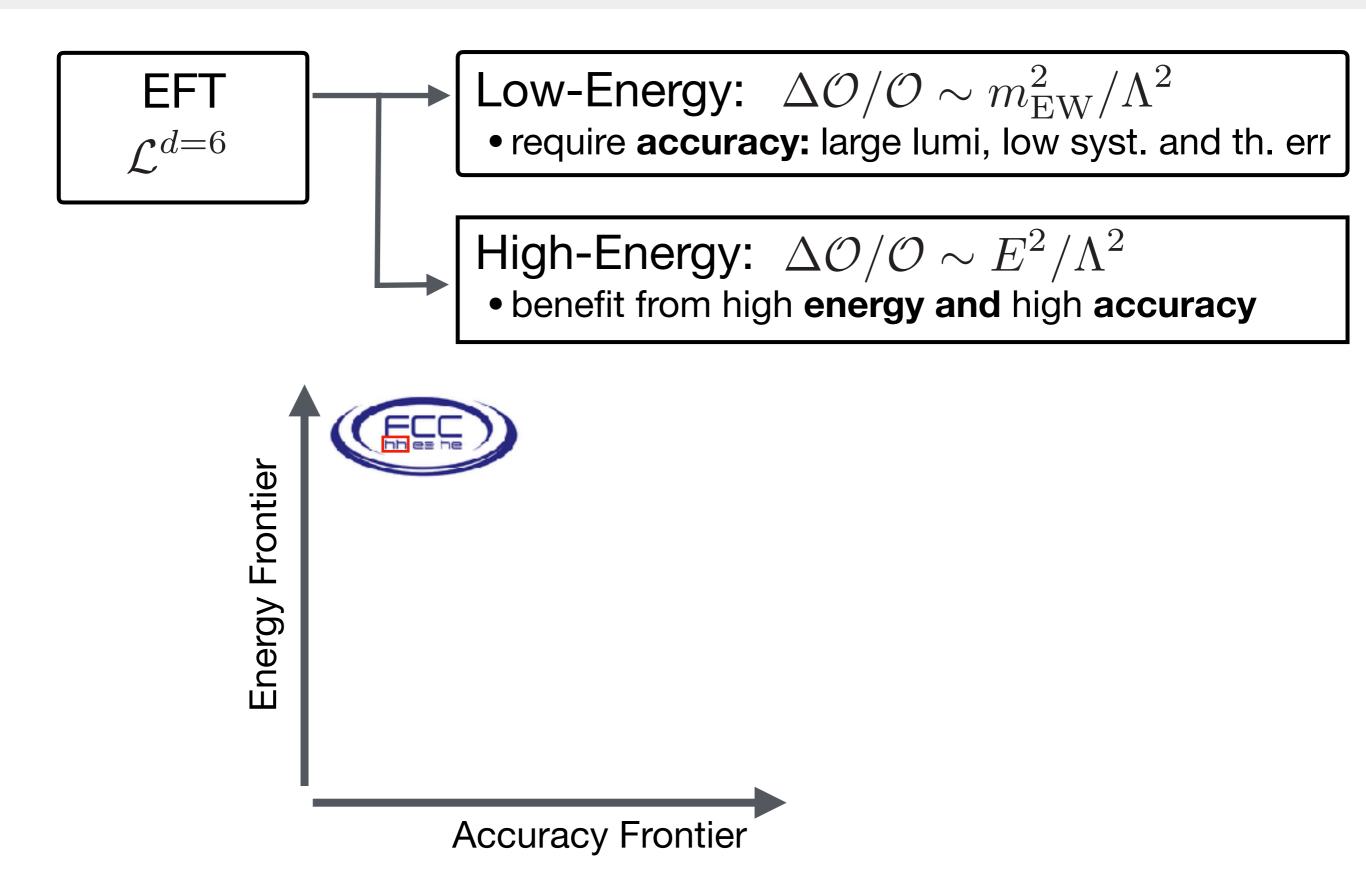
light fermion couplings in BSM, such that not excluded by flavour physics. To be studied by examples, or on the basis of motivated flavour models

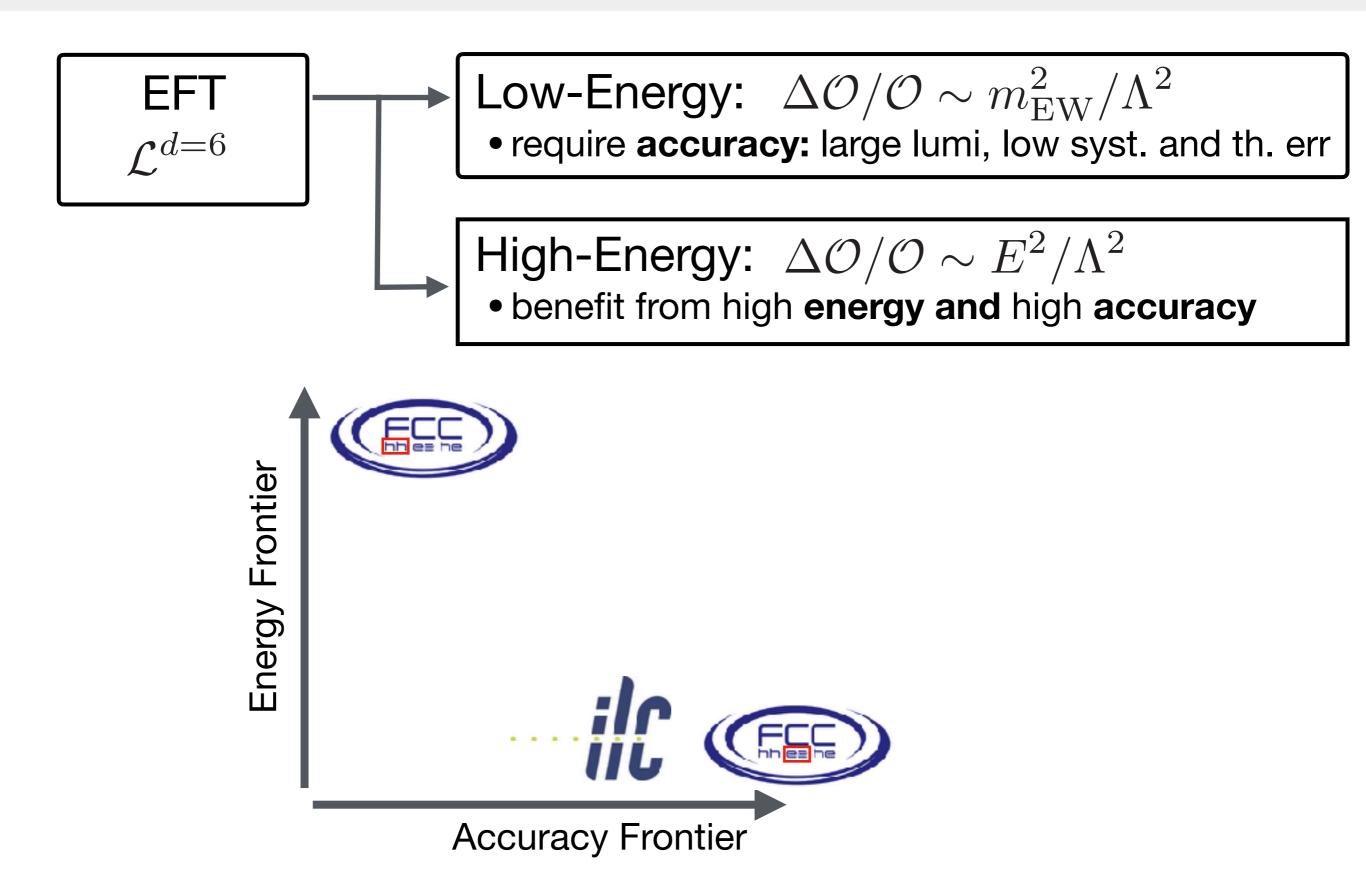


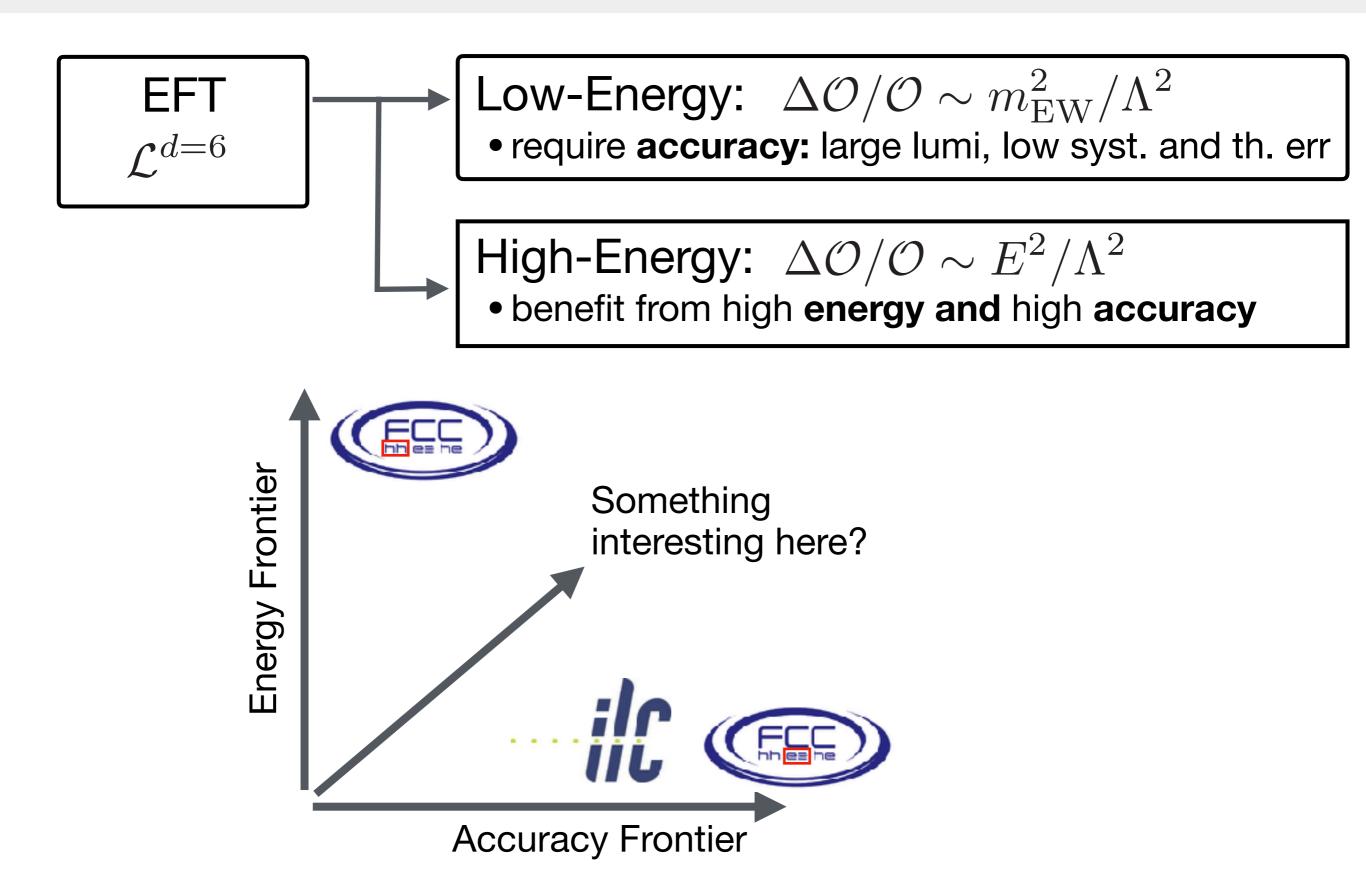


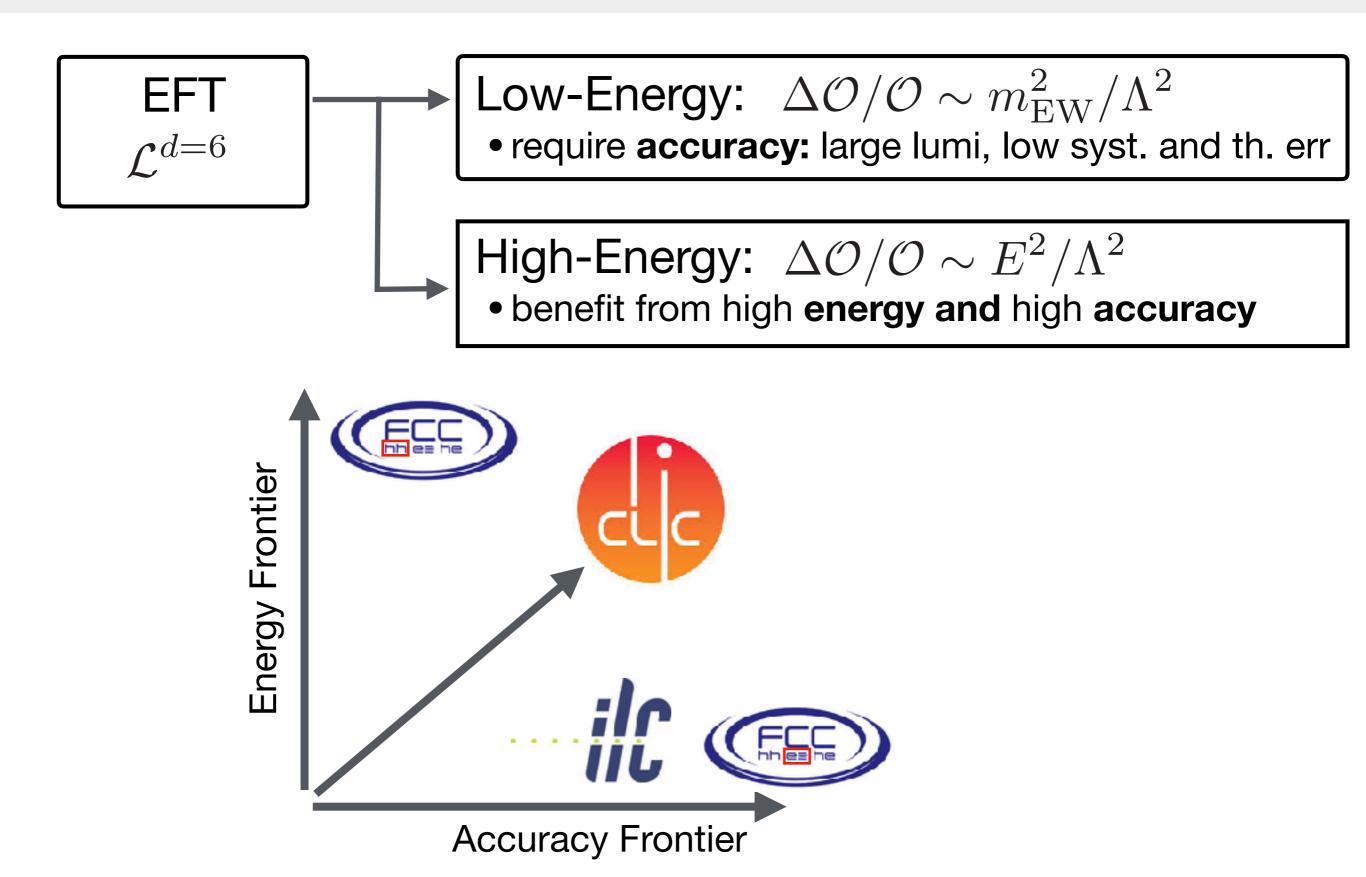


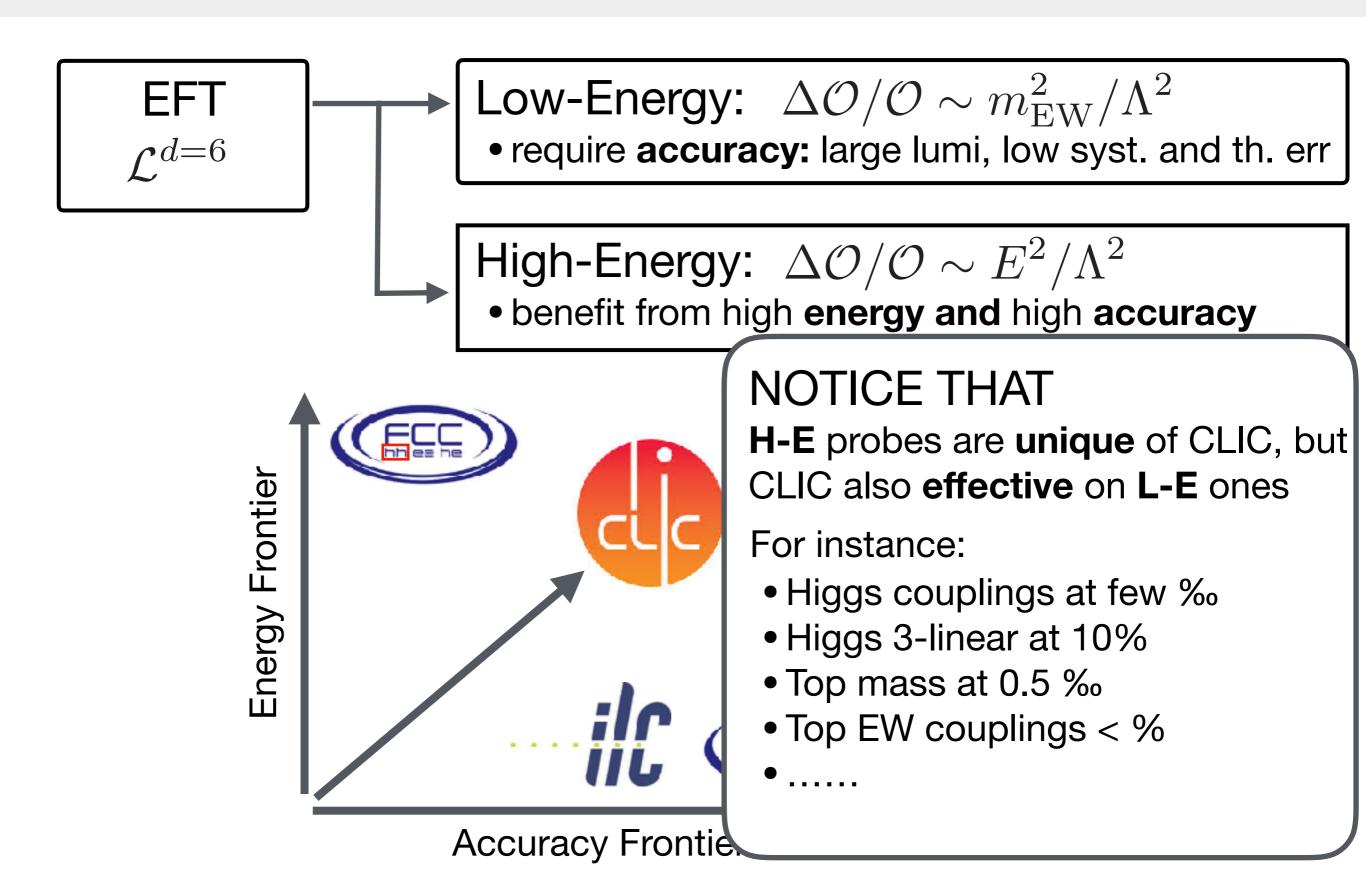




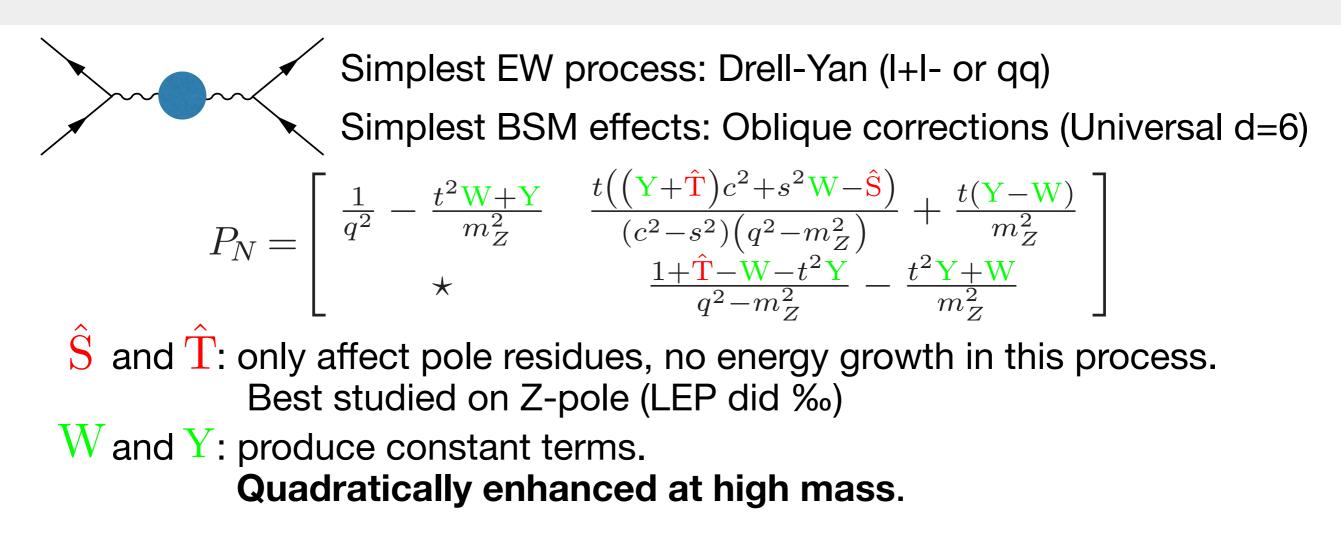




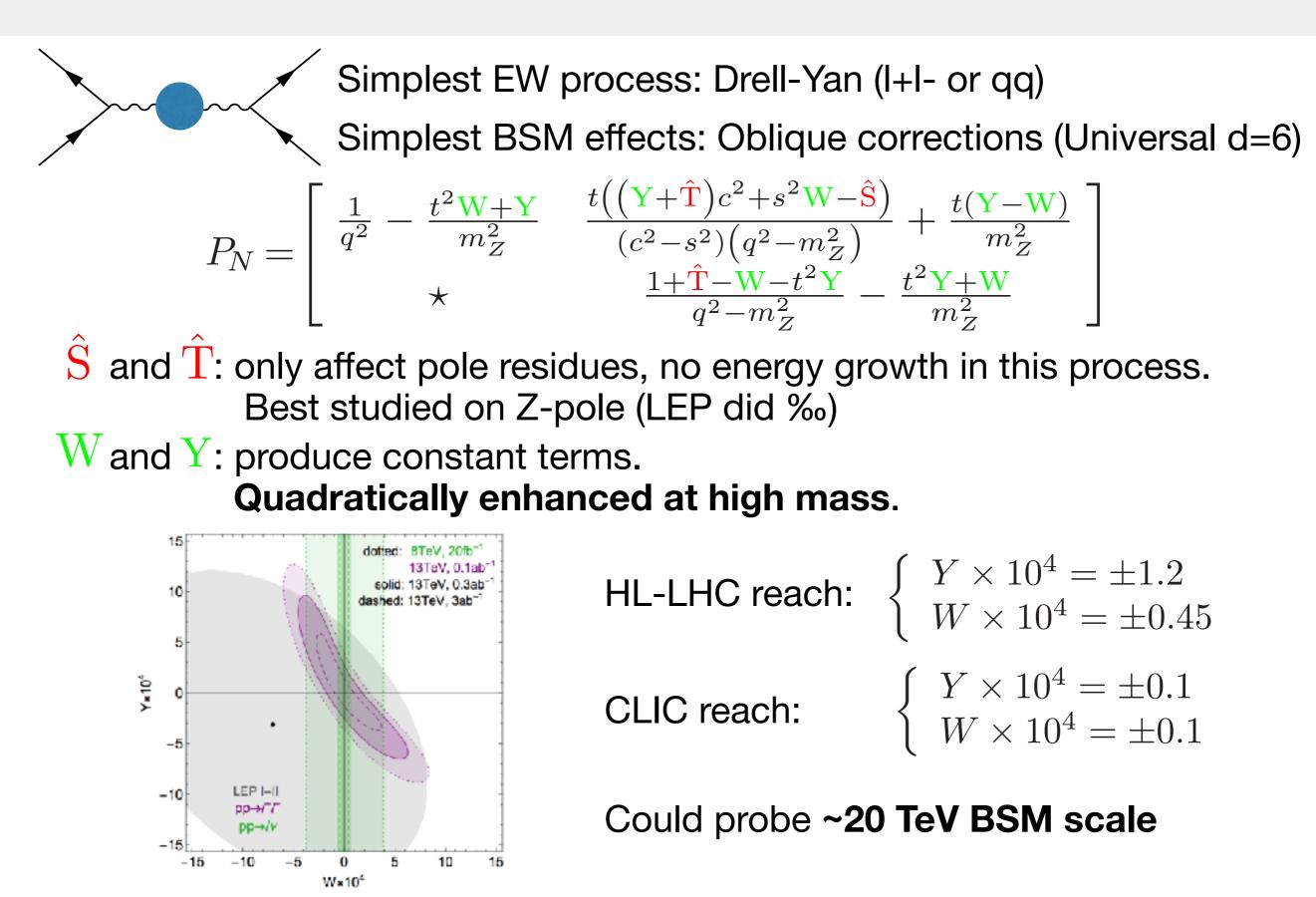




### High-Energy Drell-Yan



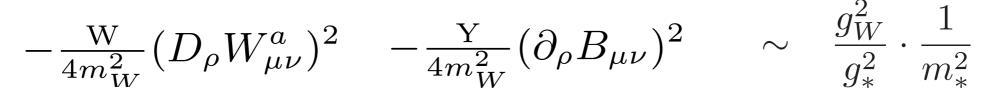
#### High-Energy Drell-Yan



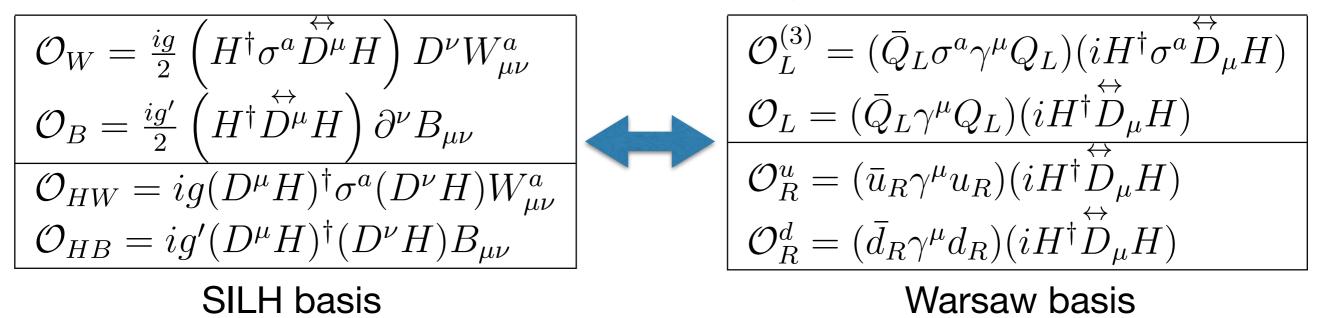
W/Y limits **easily evaded** (e.g., by strongly-coupled SILH):

$$-\frac{W}{4m_W^2} (D_\rho W^a_{\mu\nu})^2 - \frac{Y}{4m_W^2} (\partial_\rho B_{\mu\nu})^2 \sim \frac{g_W^2}{g_*^2} \cdot \frac{1}{m_*^2}$$

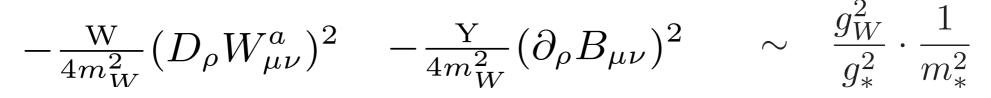
W/Y limits easily evaded (e.g., by strongly-coupled SILH):



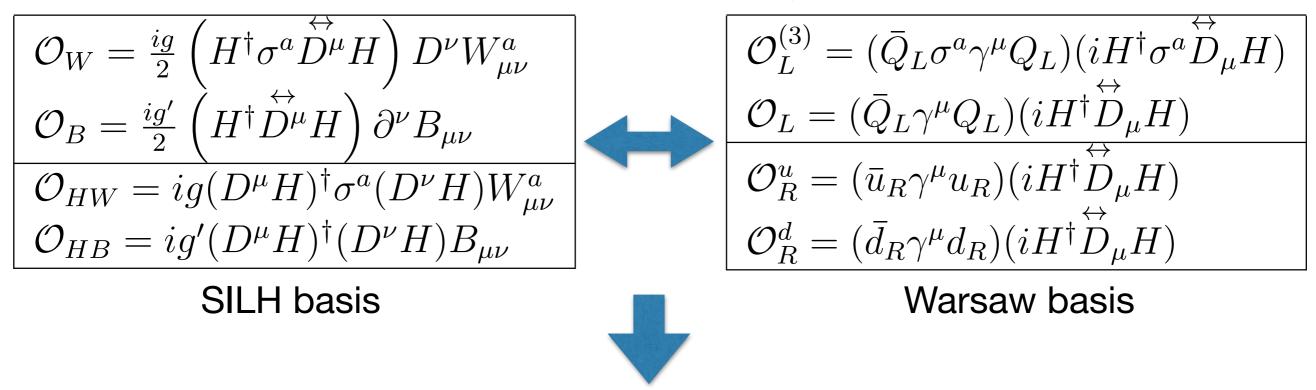
Some un-suppressed Universal operators:  $\sim 1/m_*^2$  (SILH-basis coefficient)



W/Y limits easily evaded (e.g., by strongly-coupled SILH):



Some un-suppressed Universal operators:  $\sim 1/m_*^2$  (SILH-basis coefficient)



Growing-with-energy longitudinal diboson and boson plus Higgs prod.

#### Three growing-with-energy effects (operators).[Franceschini, Panico, Pomarol, Riva, AW]

Amplitude	High-energy primaries	Deviations from SM couplings
$\overline{\bar{u}_L d_L} \to W_L Z_L, W_L h$	$\sqrt{2}a_q^{(3)}$	$\sqrt{2} \frac{g^2 \Lambda^2}{4m_W^2} \left[ c_{\theta_W} (\delta g_{uL}^Z - \delta g_{dL}^Z) / g - c_{\theta_W}^2 \delta g_1^Z \right]$
$\bar{u}_L u_L \to W_L W_L$	$a_q^{(1)} + a_q^{(3)}$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_L t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{u_L} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{dL}^Z} / g \right]$
$\bar{d}_L d_L \to Z_L h$		
$\bar{d}_L d_L \to W_L W_L$	$a_q^{(1)} - a_q^{(3)}$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_L t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{d_L} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{uL}^Z} / g \right]$
$\bar{u}_L u_L \to Z_L h$		
$\bar{f}_R f_R \to W_L W_L, Z_L h$	$a_f$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_{f_R} t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{f_R} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{fR}^Z} / g \right]$

Two of which independent (and same for q and I) for Universal theories

**HL-LHC** has some sensitivity to one of them:

$$a_q^{(3)} \sim \pm 5 \times 10^{-2} \text{TeV}^{-2}$$

#### Three growing-with-energy effects (operators).[Franceschini, Panico, Pomarol, Riva, AW]

Amplitude	High-energy primaries	Deviations from SM couplings
$\bar{u}_L d_L \to W_L Z_L, W_L h$	$\sqrt{2}a_q^{(3)}$	$\sqrt{2} \frac{g^2 \Lambda^2}{4m_W^2} \left[ c_{\theta_W} (\delta g_{uL}^Z - \delta g_{dL}^Z) / g - c_{\theta_W}^2 \delta g_1^Z \right]$
$\bar{u}_L u_L \to W_L W_L$	$a_q^{(1)} + a_q^{(3)}$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_L t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{u_L} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{dL}^Z}/g \right]$
$ \bar{d}_L d_L \to Z_L h$		
 $ \bar{d}_L d_L \to W_L W_L$	$a_q^{(1)} - a_q^{(3)}$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_L t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{d_L} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{uL}^Z} / g \right]$
$\bar{u}_L u_L \to Z_L h$		
 $-\bar{f}_R f_R \to W_L W_L, Z_L h$	$a_f$	$-\frac{g^2\Lambda^2}{2m_W^2} \left[ Y_{f_R} t_{\theta_W}^2 \boldsymbol{\delta \kappa_{\gamma}} + T_Z^{f_R} \boldsymbol{\delta g_1^Z} + c_{\theta_W} \boldsymbol{\delta g_{fR}^Z} / g \right]$

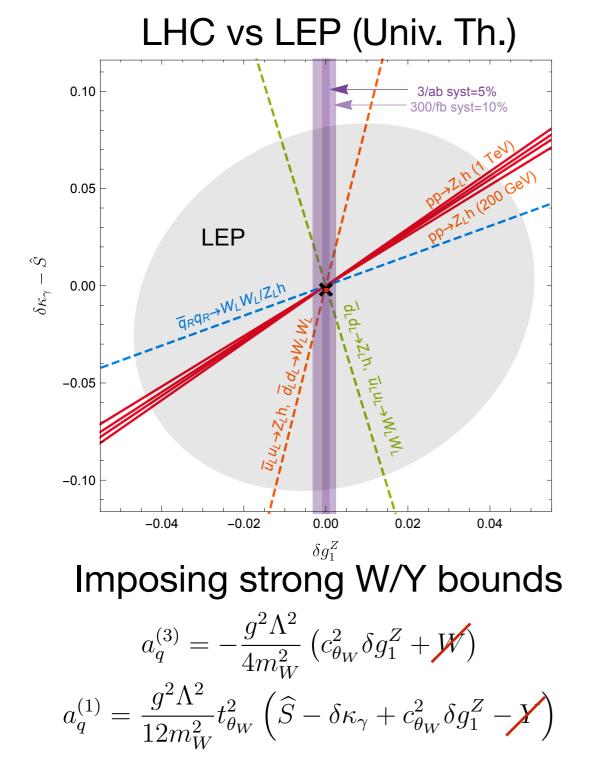
Two of which independent (and same for q and I) for Universal theories

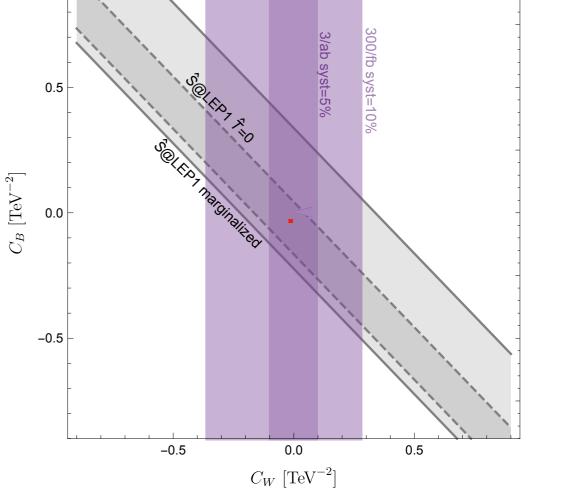
**HL-LHC** has some sensitivity to one of them:

CLIC with polarised beams sensitive to all of them:  

$$a \sim \pm 3 \times 10^{-2} \text{TeV}^{-2}$$
[from Ellis, Roloff, Sanz, You, 2017]

BSM Implications: [CLIC is red point in the middle of the plots]

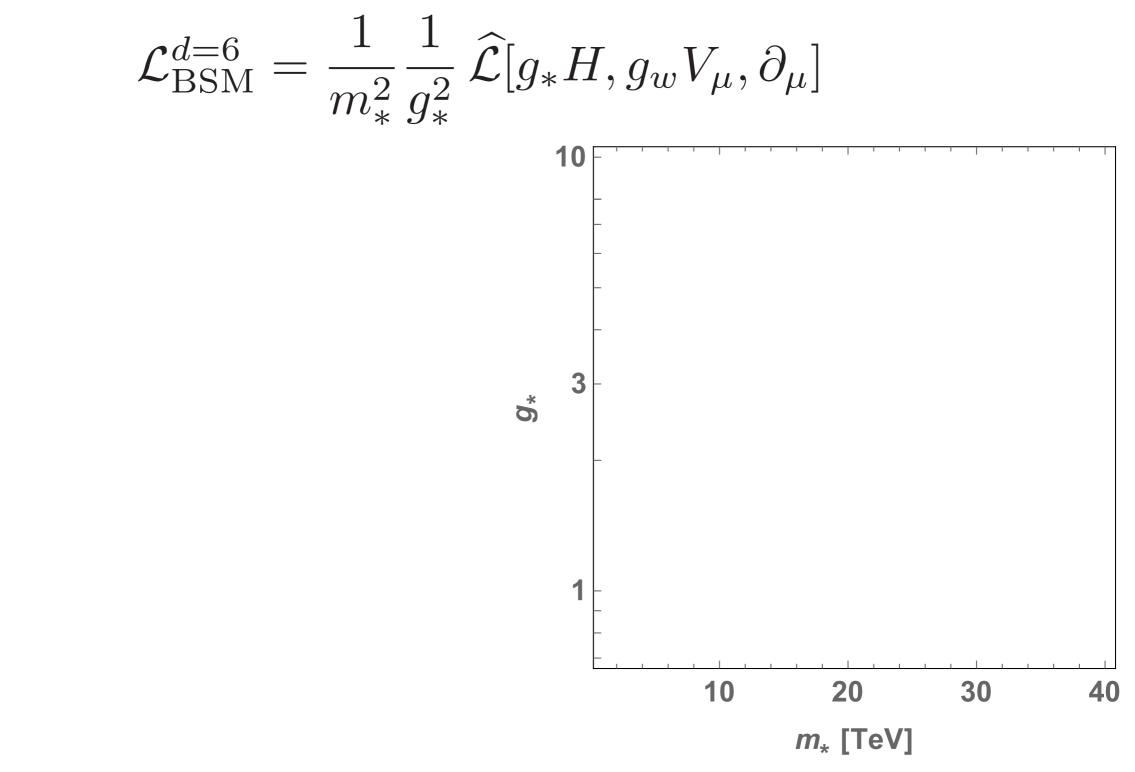




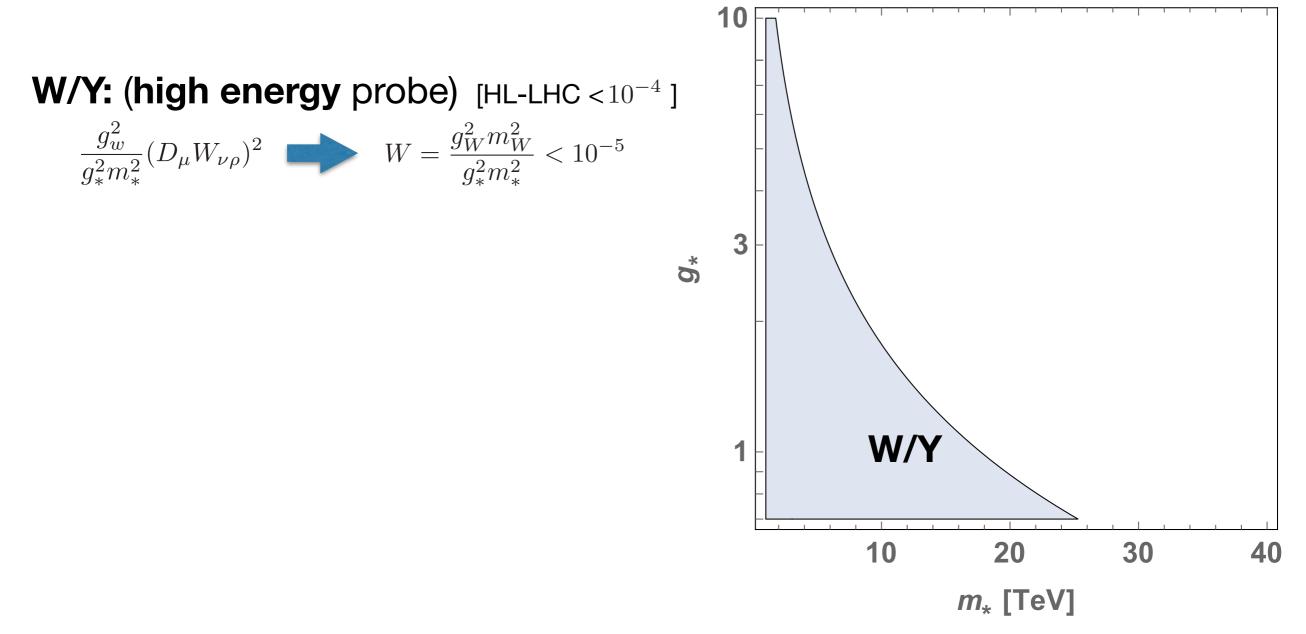
LHC vs LEP (Composite Higgs)

Power-counting + loop suppression

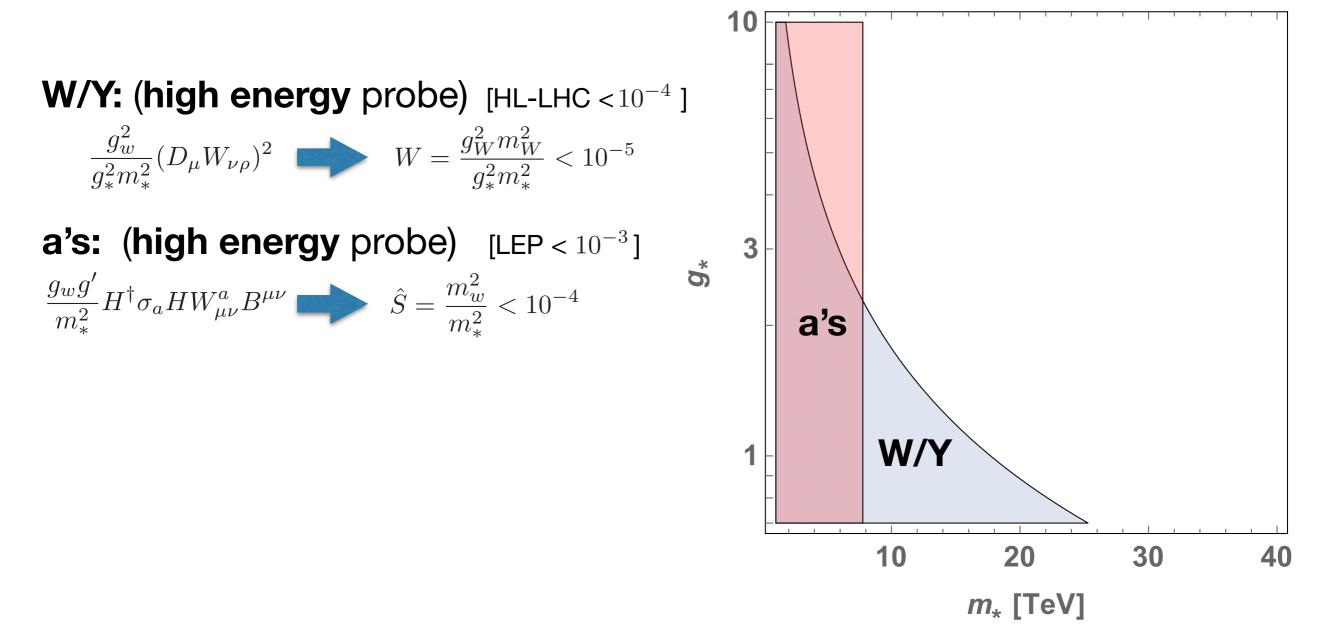
$$a_q^{(3)} = \frac{g^2}{4} (c_W + c_{HW} - c_{ZW})$$
$$\widehat{S} = (c_W + c_B) \frac{m_W^2}{\Lambda^2}$$



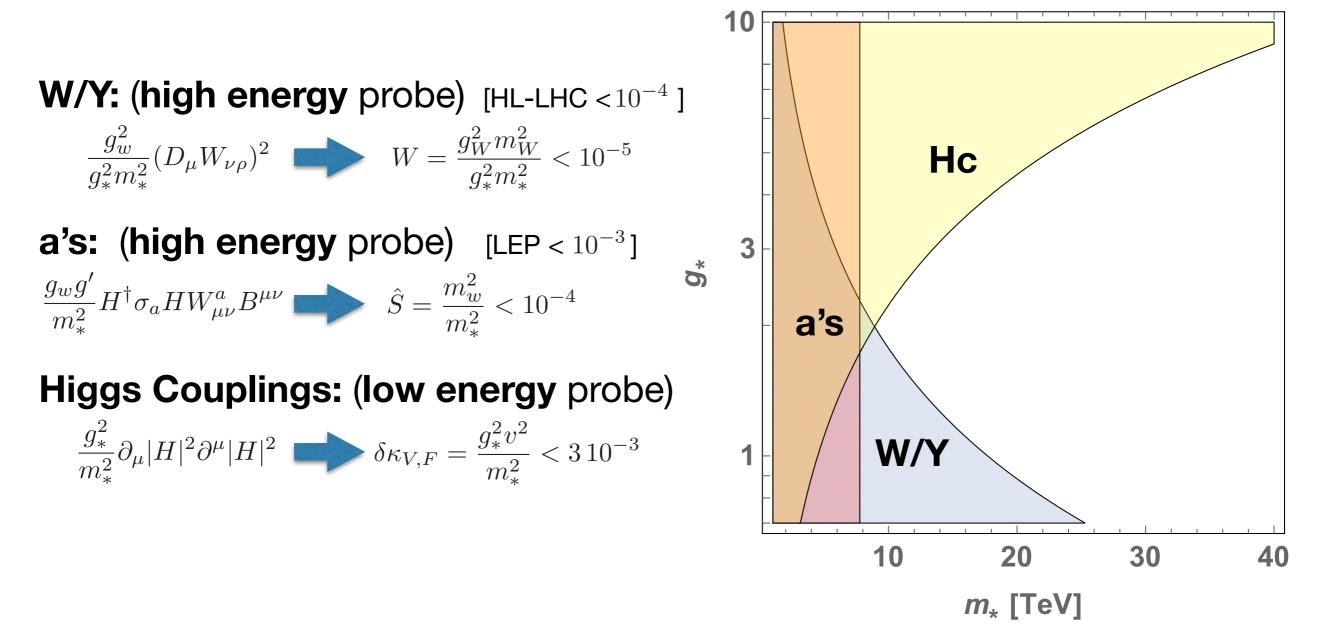
$$\mathcal{L}_{\rm BSM}^{d=6} = \frac{1}{m_*^2} \frac{1}{g_*^2} \,\widehat{\mathcal{L}}[g_*H, g_w V_\mu, \partial_\mu]$$



$$\mathcal{L}_{\text{BSM}}^{d=6} = \frac{1}{m_*^2} \frac{1}{g_*^2} \,\widehat{\mathcal{L}}[g_*H, g_w V_\mu, \partial_\mu]$$



$$\mathcal{L}_{\text{BSM}}^{d=6} = \frac{1}{m_*^2} \frac{1}{g_*^2} \,\widehat{\mathcal{L}}[g_*H, g_w V_\mu, \partial_\mu]$$



### High-Energy Tops

Growing-with-Energy in ee->tt: [Durieux, Perelló, Vos, C.Zhang]		
$Q_{\varphi t} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu} \varphi) (\overline{t} \gamma^{\mu} t)$	quadratic growth	
$\mathcal{N}_{tB} = (\bar{t}\gamma^{\mu}t)(\bar{e}\gamma_{\mu}e + \frac{1}{2}\bar{l}\gamma_{\mu}l)$ $Q_{t\varphi} = (\varphi^{\dagger}\varphi)(\bar{q}t\widetilde{\varphi})$	linear growth (and diff. top decay dist needed)	
$Q_{tB} = (\overline{q}\sigma^{\mu\nu}t)\widetilde{\varphi}B_{\mu\nu}$ $Q_{\varphi q}^{(1)} = (\varphi^{\dagger}i\overleftrightarrow{D}_{\mu}\varphi)(\overline{q}\gamma^{\mu}q)$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	
$Q_{\varphi q}^{(3)} = (\varphi^{\dagger} i \overleftrightarrow{D}_{\mu}^{I} \varphi) (\overline{q} \tau^{I} \gamma^{\mu} q)$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	
$Q_{tW} = (\overline{q}\sigma^{\mu\nu}t)\tau^{I}\widetilde{\varphi}W^{I}_{\mu\nu}$ $\mathcal{N}_{qB} = (\overline{q}\gamma^{\mu}q)(\overline{e}\gamma_{\mu}e + \frac{1}{2}\overline{l}\gamma_{\mu}l)$	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	
$\mathcal{N}_{qW} = (\overline{q}\tau^I \gamma^\mu q) (\overline{l}\tau^I \gamma_\mu l)$		

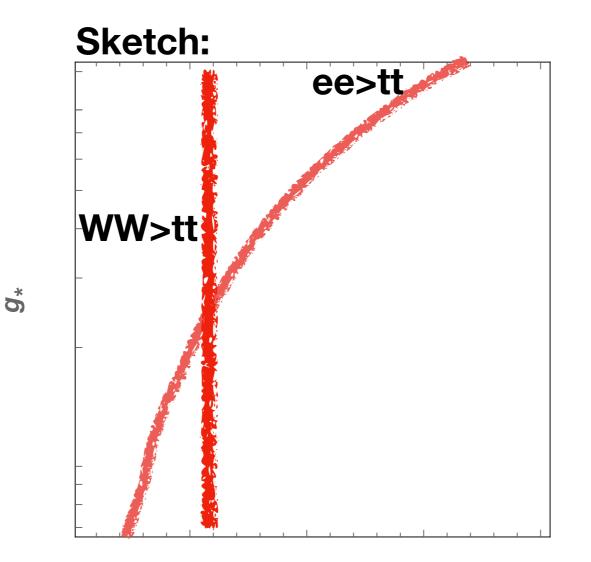
### High-Energy Tops

Growing-with-Energy in **WW->tt**: [Grojean, You, AW, Z.Zhang] quadratic growth  $Q_{\varphi t} = (\varphi^{\dagger} i D_{\mu} \varphi) (\bar{t} \gamma^{\mu} t)$  $\mathcal{N}_{tB} = (\bar{t}\gamma^{\mu}t)(\bar{e}\gamma_{\mu}e + \frac{1}{2}\bar{l}\gamma_{\mu}l)$ linear growth (and diff. top decay dist needed)  $Q_{t\varphi} = (\varphi^{\dagger}\varphi)(\overline{q} t \,\widetilde{\varphi})$ **H-E probe of y**<sub>t</sub>? (no result yet)  $Q_{tB} = (\overline{q}\sigma^{\mu\nu}t)\widetilde{\varphi}B_{\mu\nu}$ 95% CL limits (3TeV CLIC, 3ab<sup>-1</sup>, sys. err. 0, 3%)  $Q_{\varphi q}^{(1)} = (\varphi^{\dagger} i D_{\mu} \varphi) (\overline{q} \gamma^{\mu} q)$ 0.100 PRELIMINARY  $Q^{(3)}_{\varphi q} = (\varphi^{\dagger} i D^{I}_{\mu} \varphi) (\overline{q} \, \tau^{I} \gamma^{\mu} q)$ 0.010  $\blacksquare$   $A_{co} > 0$  (no tt bkg)  $Q_{tW} = (\overline{q}\sigma^{\mu\nu}t)\tau^{I}\widetilde{\varphi}W^{I}_{\mu\nu}$  $A_{co} > 0.1$  (no tt bkg) 0.001  $A_{co} > 0.1$ 0.001  $\blacksquare A_{co} > 0.15$  (no tt bkg)  $\mathcal{N}_{qB} = (\overline{q}\gamma^{\mu}q)(\overline{e}\gamma_{\mu}e + \frac{1}{2}l\gamma_{\mu}l)$  $A_{co} > 0.15$ 0.010  $\mathcal{N}_{qW} = (\overline{q}\tau^I \gamma^\mu q) (\overline{l}\tau^I \gamma_\mu l)$ 0.100  $c_{\phi q}^{(1)}$  $c_{\phi q}^{(3)}$ C<sub>Φt</sub> **C**tW

#### Top-philic EFT

Assuming composite t<sub>R</sub> and H, elementary t<sub>L</sub> and gauge

$$\mathcal{L}_{BSM}^{d=6} = \frac{1}{m_*^2} \frac{1}{g_*^2} \widehat{\mathcal{L}}[g_* t_R, y_t q_L, g_* H, g_w V_\mu, \partial_\mu]$$



#### More Maps

- **High Energy FCNC:** [ee >  $\tau \mu$ , ee > t q, ...] can compete with flavour phys. and/or exotic top dec?
- Light quark Yukawa determination: assessing BSM impact

- - -

- **EW-Charged Particles:** [Higgsino/EW-ikno, Minimal DM] Opportunity: **Millicharged Minimal DM at 1.5 TeV**
- **Exploring Holes in SUSY parameter space.**
- Extra Singlets Production: [for EWBG? related to H<sup>3</sup>?]

### Summary

- Indirect BSM probes of heavy new physics through growing-with-energy effects, exploring the Energy and Accuracy Frontier, are very effective at CLIC.
- Several groups are further exploring CLIC potential in this direction, and assessing BSM implications of the program.
- This adds to, and **complements**, well-studied L-E probes
- Direct search program also to be updated with new ideas, in reaction to LHC non-discovery.
- Discussing Yellow Report report summary by this year.