



THE 27TH INTERNATIONAL CONFERENCE ON ULTRARELATIVISTIC NUCLEUS-NUCLEUS COLLISIONS  
VENEZIA, ITALY 13-19 MAY 2018

# Directed flow of quarks from the RHIC Beam Energy Scan measured by STAR

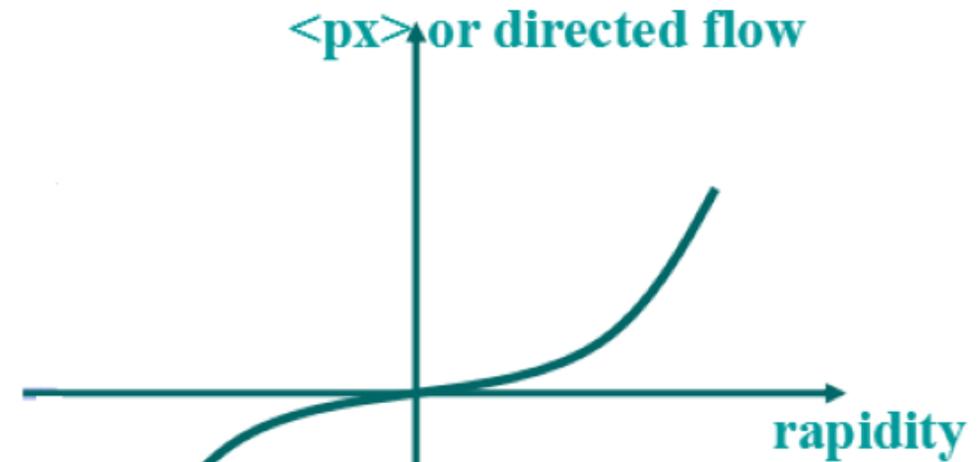
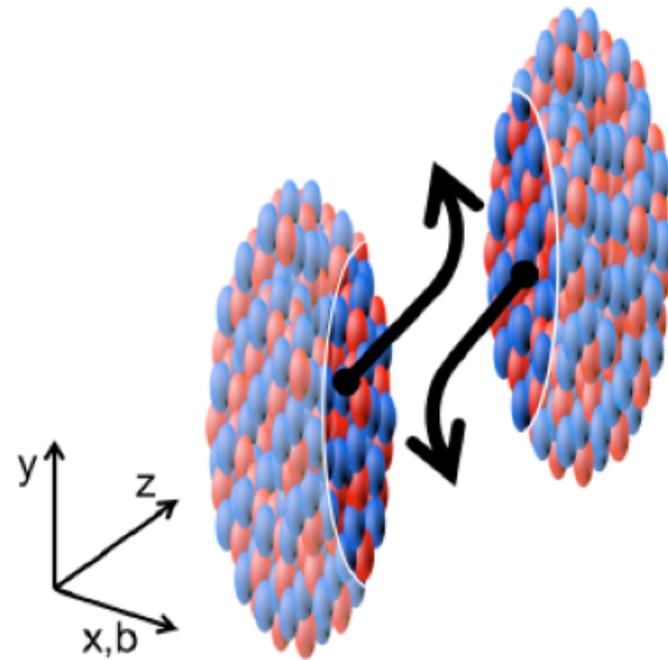
*Gang Wang*  
(for the STAR Collaboration)

UCLA



- ❖ Motivation
- ❖  $v_1$ @BES for 10 particle species
- ❖ Test of the coalescence picture
- ❖  $v_1$  of quarks
- ❖ Summary

# Directed flow ( $v_1$ )

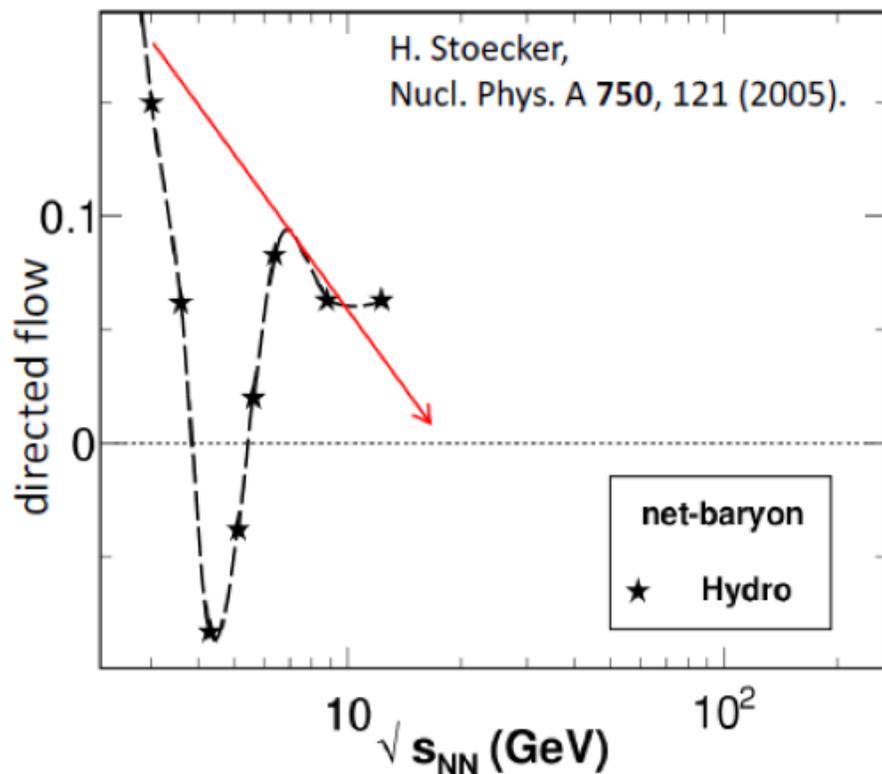


This example mostly applies to baryon transport. Other mechanisms could cause an opposite sign of  $v_1$ .

$$v_1 = \langle \cos(\phi - \Psi_{RP}) \rangle$$

- Directed flow ( $v_1$ ) describes the sideward collective motion of the particles within the reaction plane (x-z plane)
- Generated during the nuclear passage time ( $2R/\gamma$ )
- Probes the very earliest stage of the collision dynamics

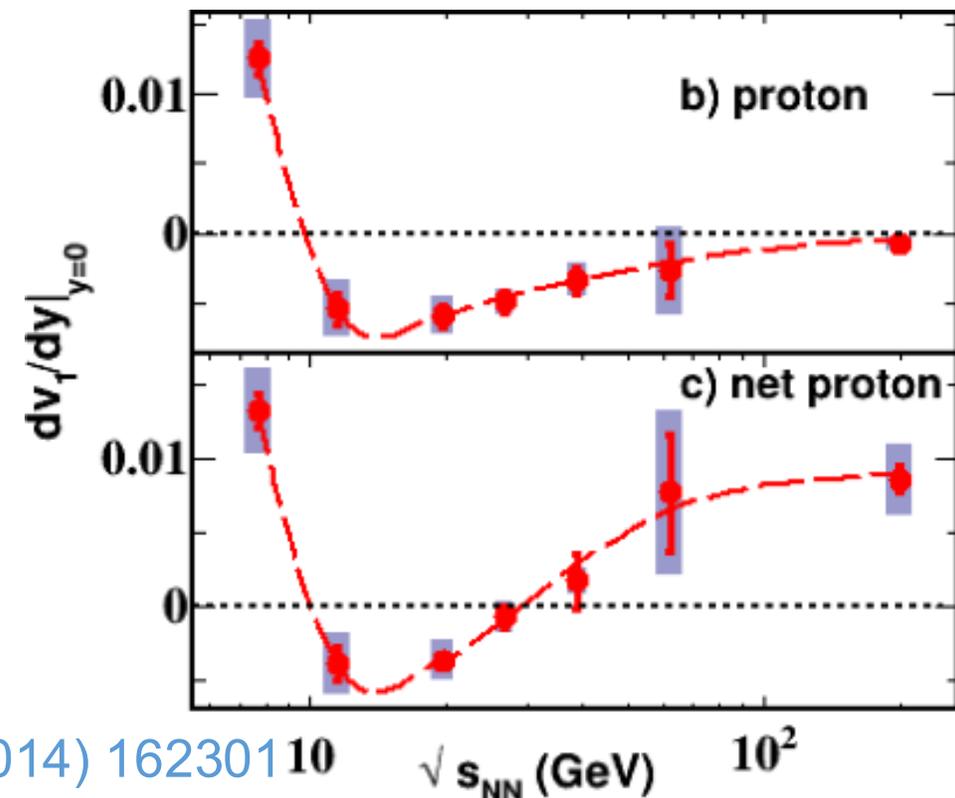
# Softest point



- “Softest Point” in EOS => a minimum in the ratio of pressure to energy density
- Strong softening consistent with the 1<sup>st</sup>-order PT
- Weaker softening is more likely due to crossover

Y. Nara, H. Niemi, J. Steinheimer, and H. Stöcker, Phys. Lett. B769 (2017) 543.  
Yu. B. Ivanov and A. A. Soldatov, Phys. Rev. C91 (2015) 024915.

- Equation of State **without** phase transition (PT): a **monotonic** trend
- Equation of State assuming 1<sup>st</sup>-order PT: a dip in  $v_1$  as a function of beam energy



# produced

vs

# transported



- $u, \bar{u}, d, \bar{d}, s$  and  $\bar{s}$

- pair production

- total number not conserved

- different waves of production

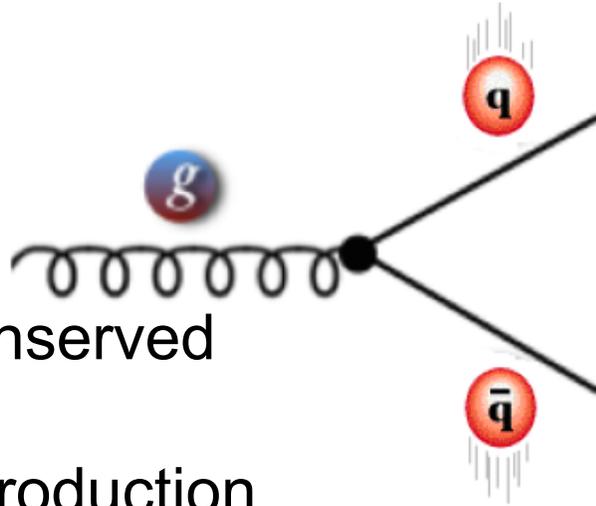
[S. Pratt, PoS CPOD2013 \(2013\) 023](#)

- may or may not be sensitive to the softening of EOS

- dominant at high collision energies

- can be studied via “produced” particles, such as anti-p, anti- $\Lambda$ ,  $K^-$  and  $\varphi$

[whose constituent quarks are all produced](#)



- $u$  and  $d$  only

- from projectile nucleons

- total number conserved

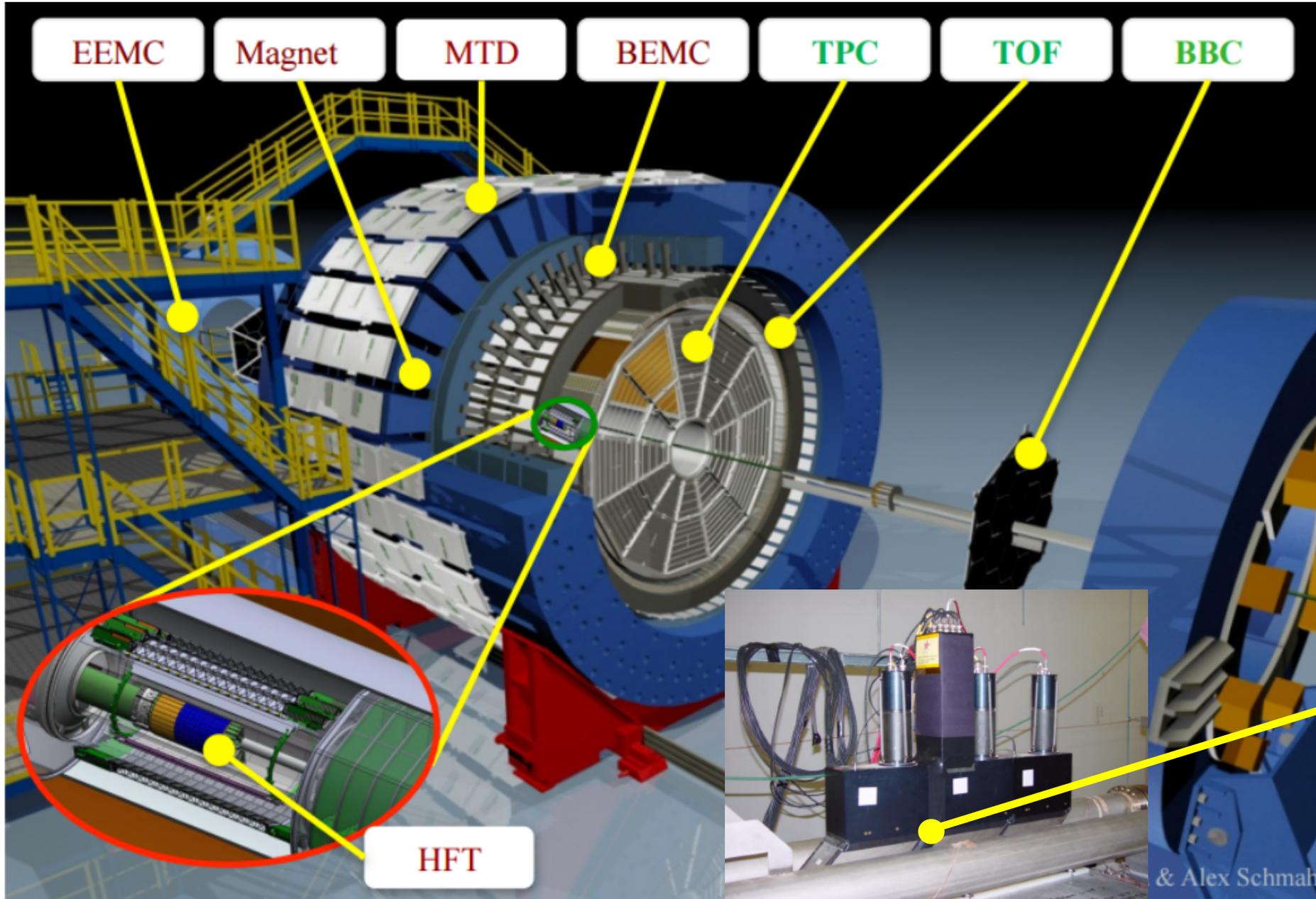
- go thru the whole evolution

- should be sensitive to the softening of EOS, if any

- dominant at low collision energies

- can be studied via net particles, such as net p, net  $\Lambda$  and net K

# STAR detectors



## Detectors used here:

### Time Projection Chamber

- Tracking
- Full azimuthal coverage
- $|\eta| < 1$  coverage
- PID for lower momenta

### Time-Of-Flight

- PID for higher momenta

### Beam-Beam Counter

- 1<sup>st</sup>-order event plane for lower beam energies

### Zero Degree Calorimeter Shower Max Detector

- 1<sup>st</sup>-order event plane for 200 GeV and 62.4 GeV

& Alex Schmah

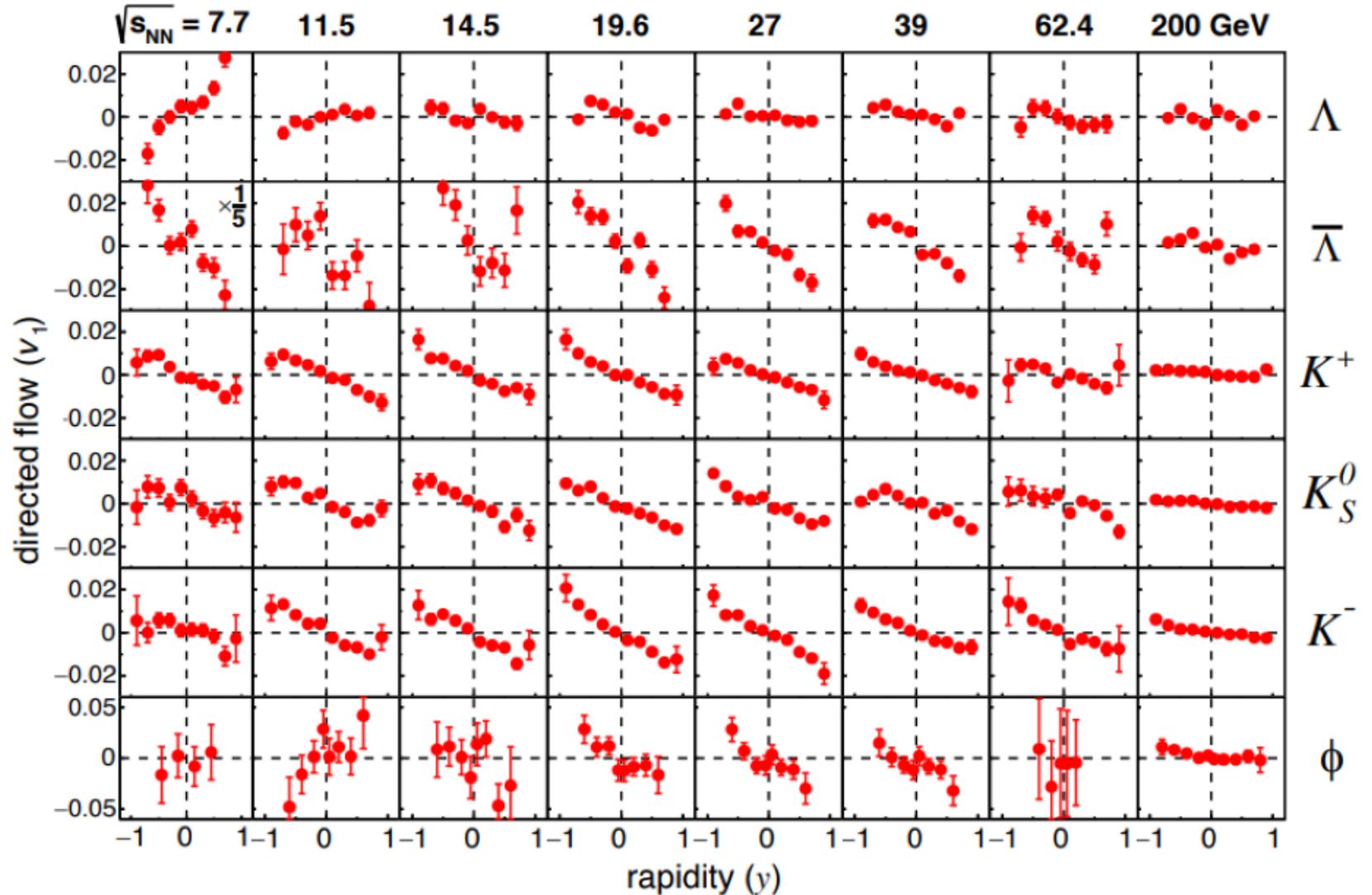
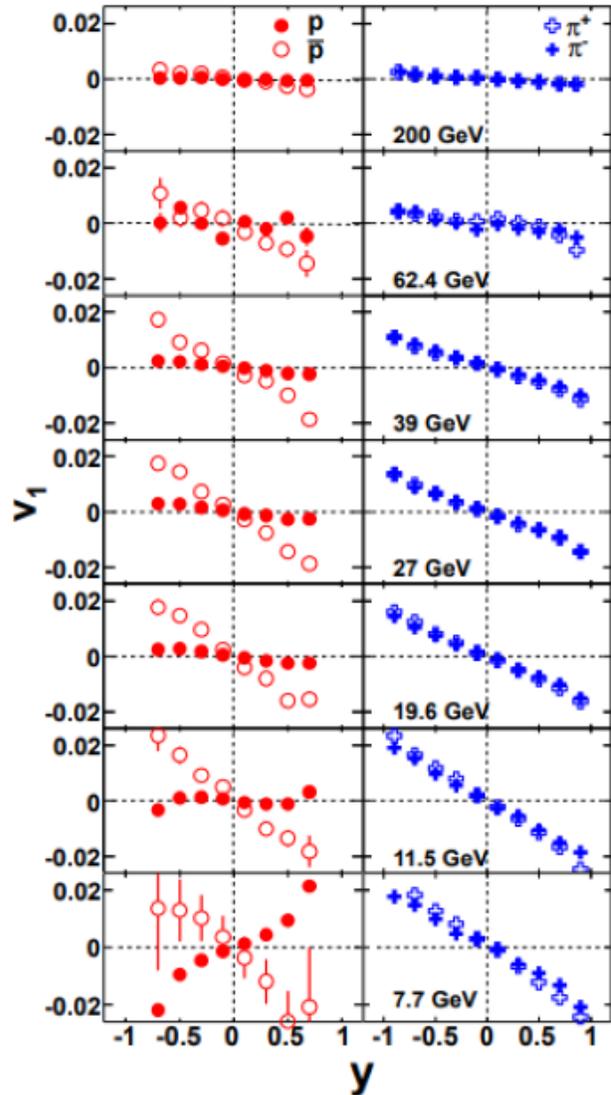
# $v_1(y)$ : 10 species & 8 energies



10-40% Au+Au collisions.

STAR, Phys. Rev. Lett. **112** (2014) 162301; Phys. Rev. Lett. **120** (2018) 62301

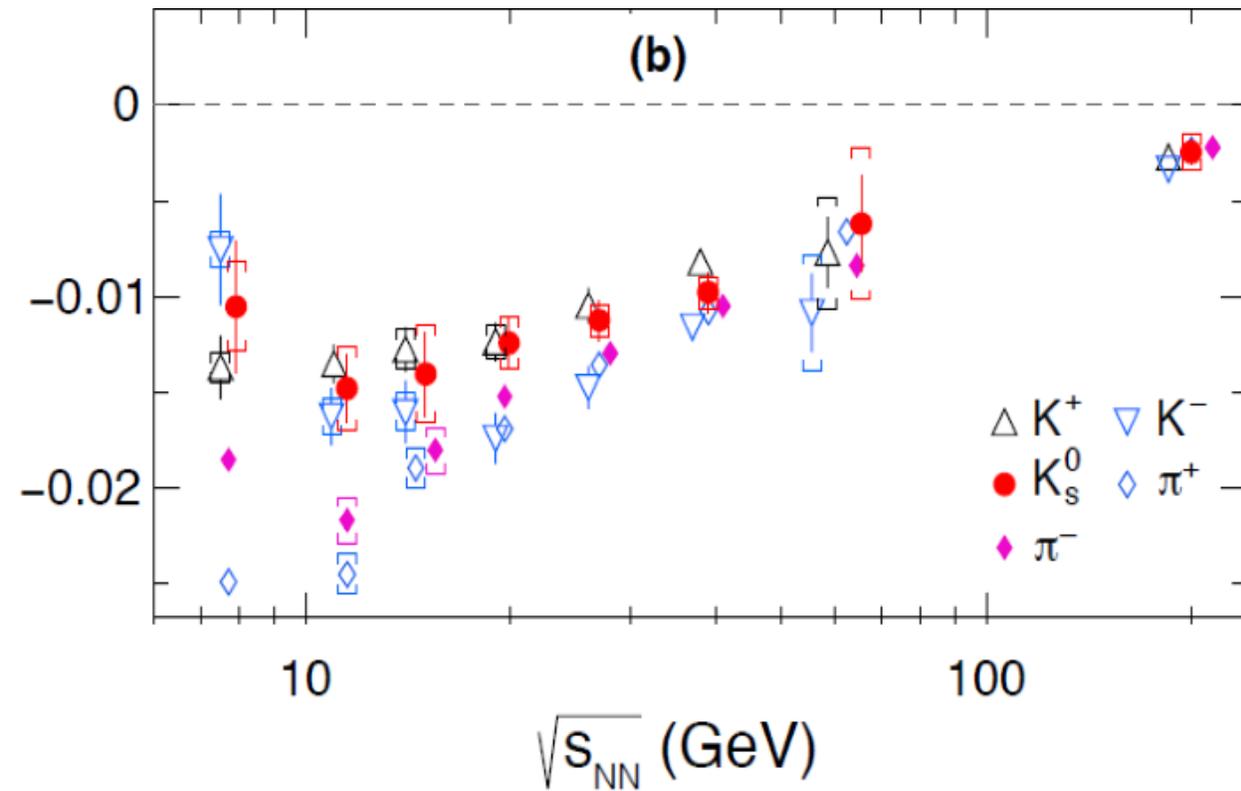
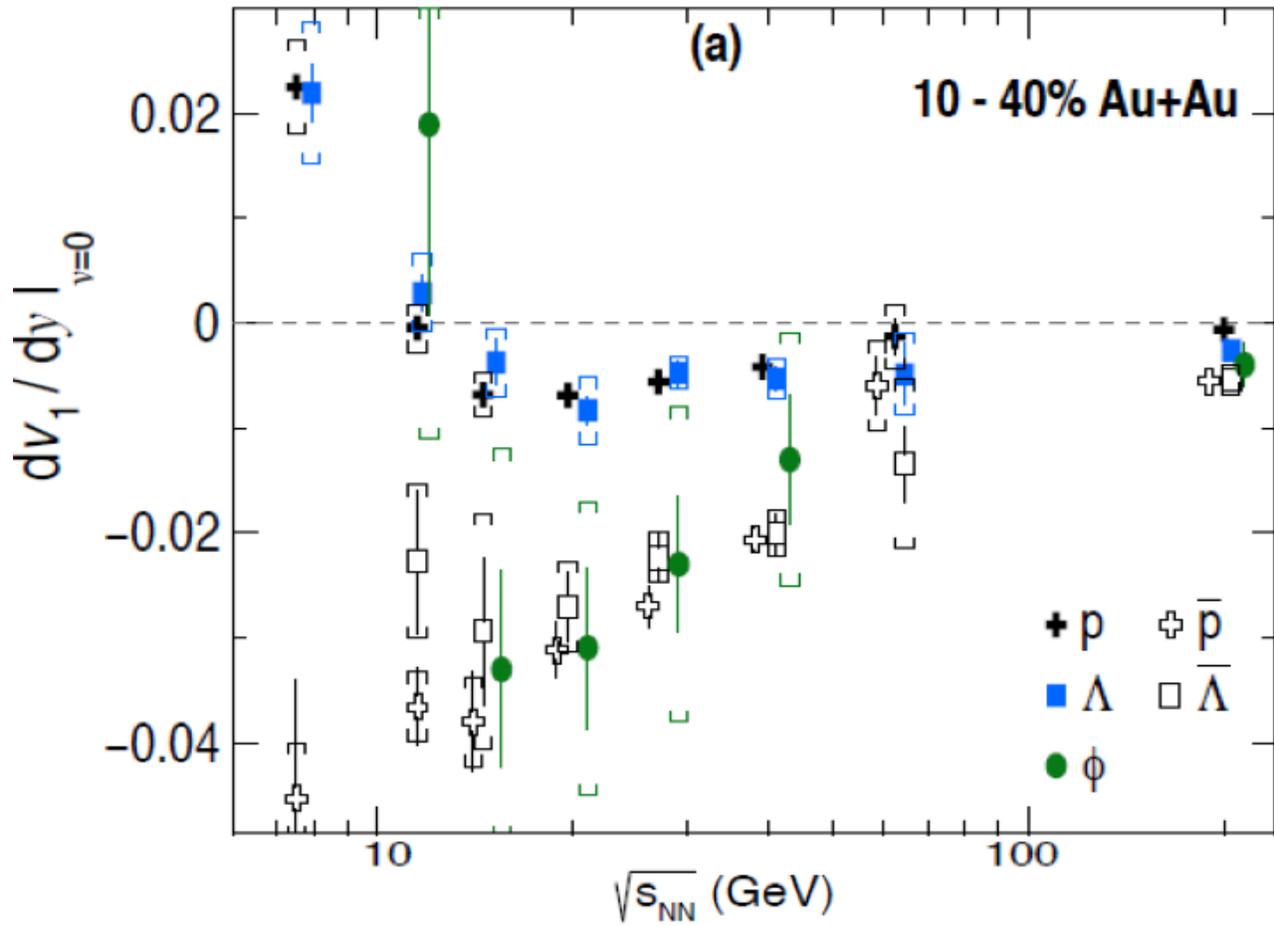
To extract  $v_1$  slope, use linear fit over  $|y| < 0.6$  for  $\Phi$  and over  $|y| < 0.8$  for all other species.



# $dv_1/dy$ vs beam energy



STAR, Phys. Rev. Lett. **120** (2018) 62301



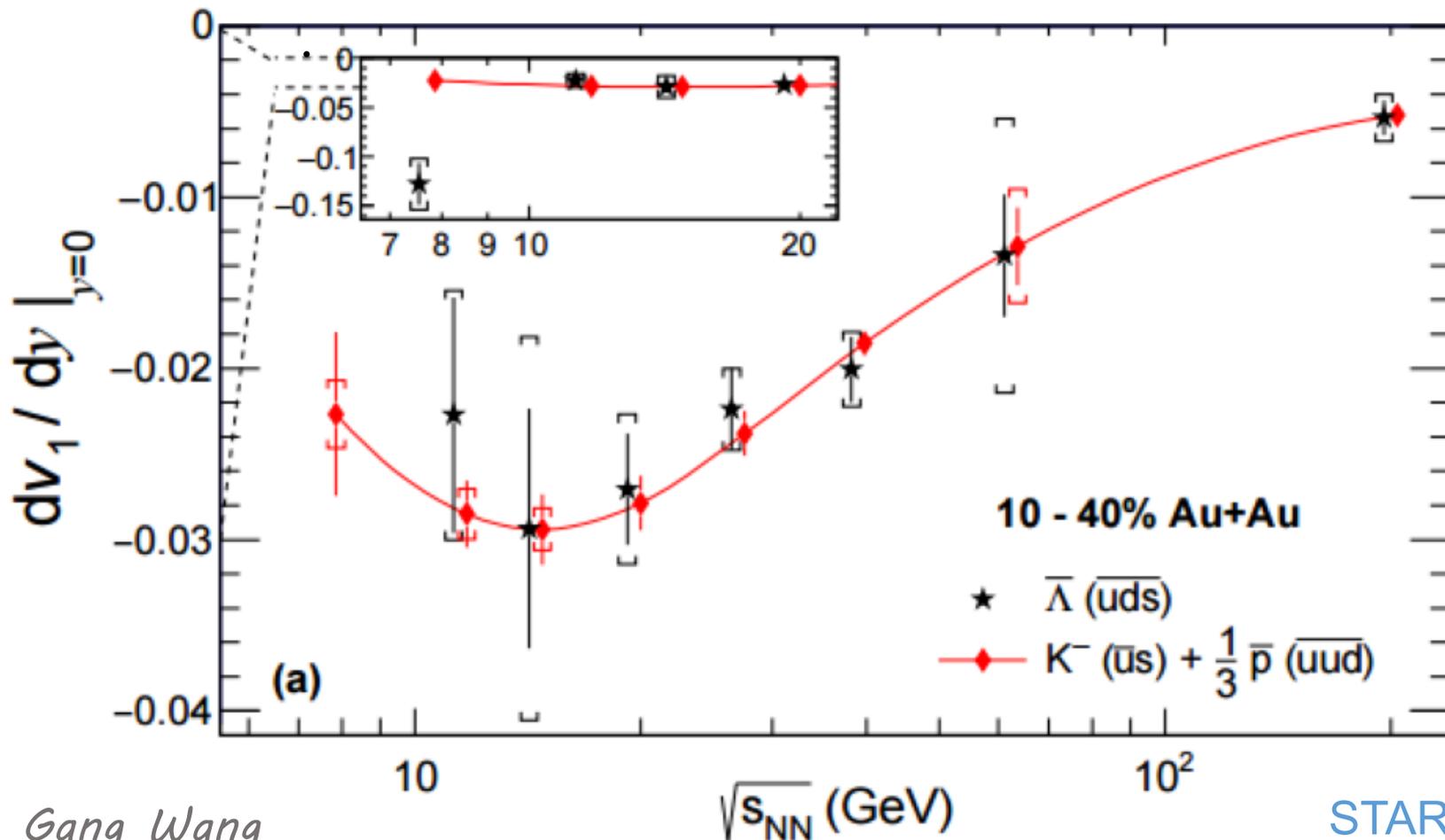
- $dv_1/dy$  for  $\Lambda$  follows proton, and changes sign in the same region  $\sqrt{s_{NN}} < 14.5$  GeV
- Purely produced particles (anti-p, anti- $\Lambda$ ,  $\Phi$ ) show similar behavior above 14.5 GeV
- Mesons show negative  $dv_1/dy$

# Coalescence sum rule: “produced” particles



Assumptions:

- $v_1$  is developed in prehadronic stage
- Hadrons are formed via coalescence:  $(v_n)_{\text{hadron}} = \sum (v_n)_{\text{constituent quarks}}$
- $(v_1)_{\bar{u}} = (v_1)_{\bar{d}}$  and  $(v_1)_s = (v_1)_{\bar{s}}$

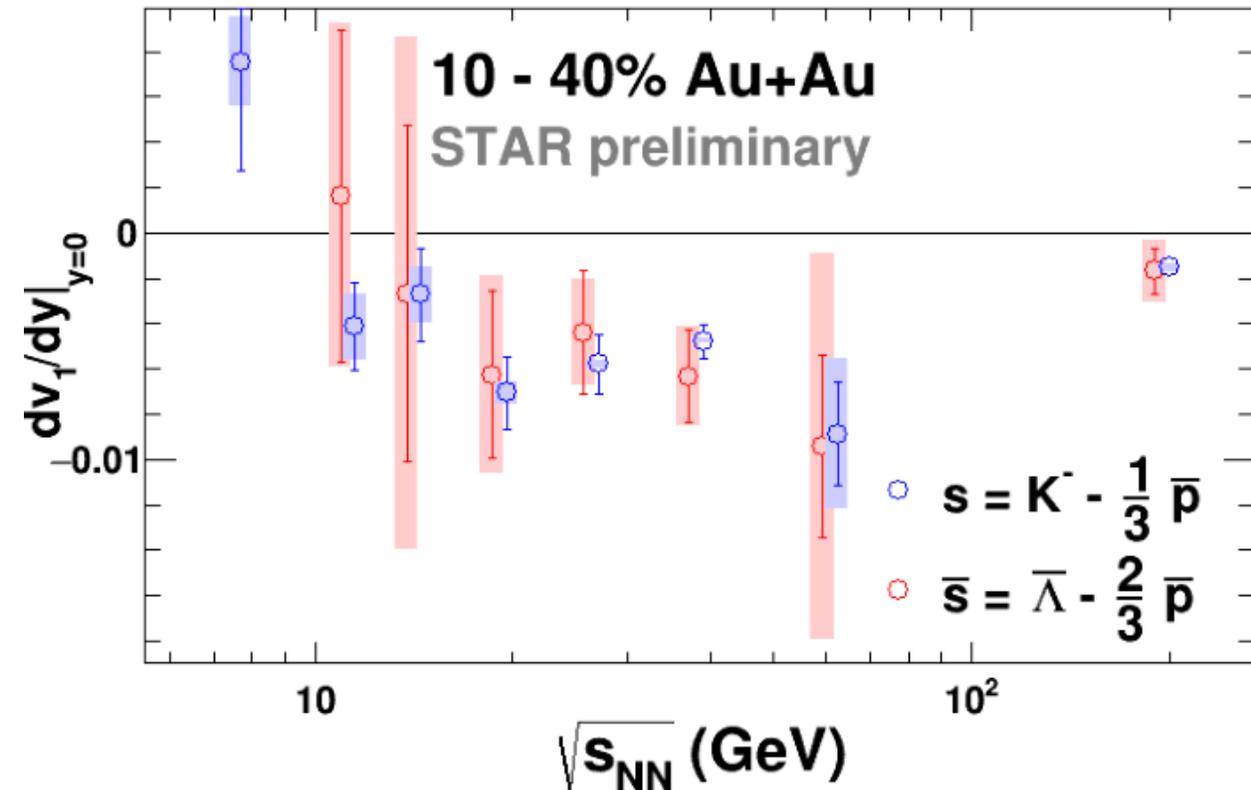
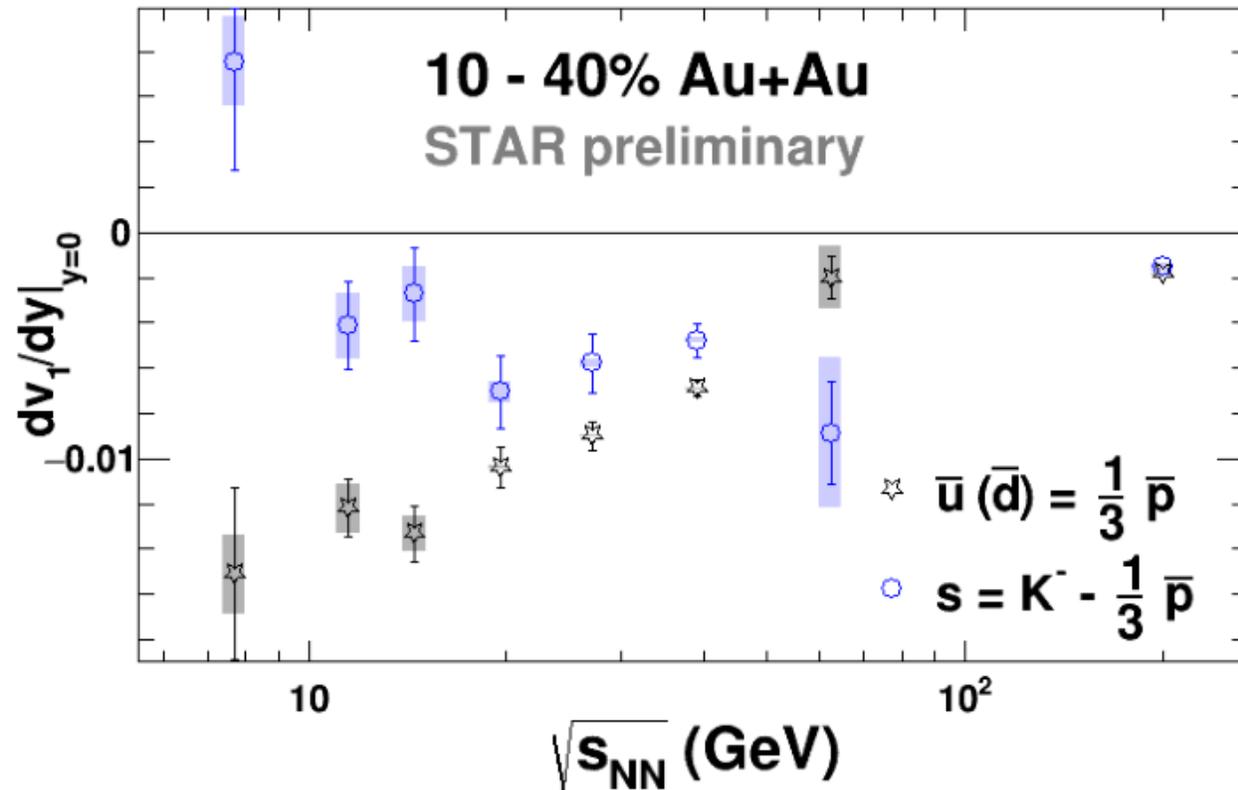


- Constituent quarks of anti-p, anti- $\Lambda$  and  $K^-$  are all produced in the collision.
- For anti- $\Lambda$ s, prediction using coalescence sum rule agrees with measured  $v_1$  above  $\sqrt{s_{NN}} = 11.5$  GeV.
- Disagreement at 7.7 GeV implies the failure of one or more of the assumptions below 11.5 GeV.

# $v_1$ of produced quarks: $u(d/\bar{u}/\bar{d})$ , $s$ and $\bar{s}$



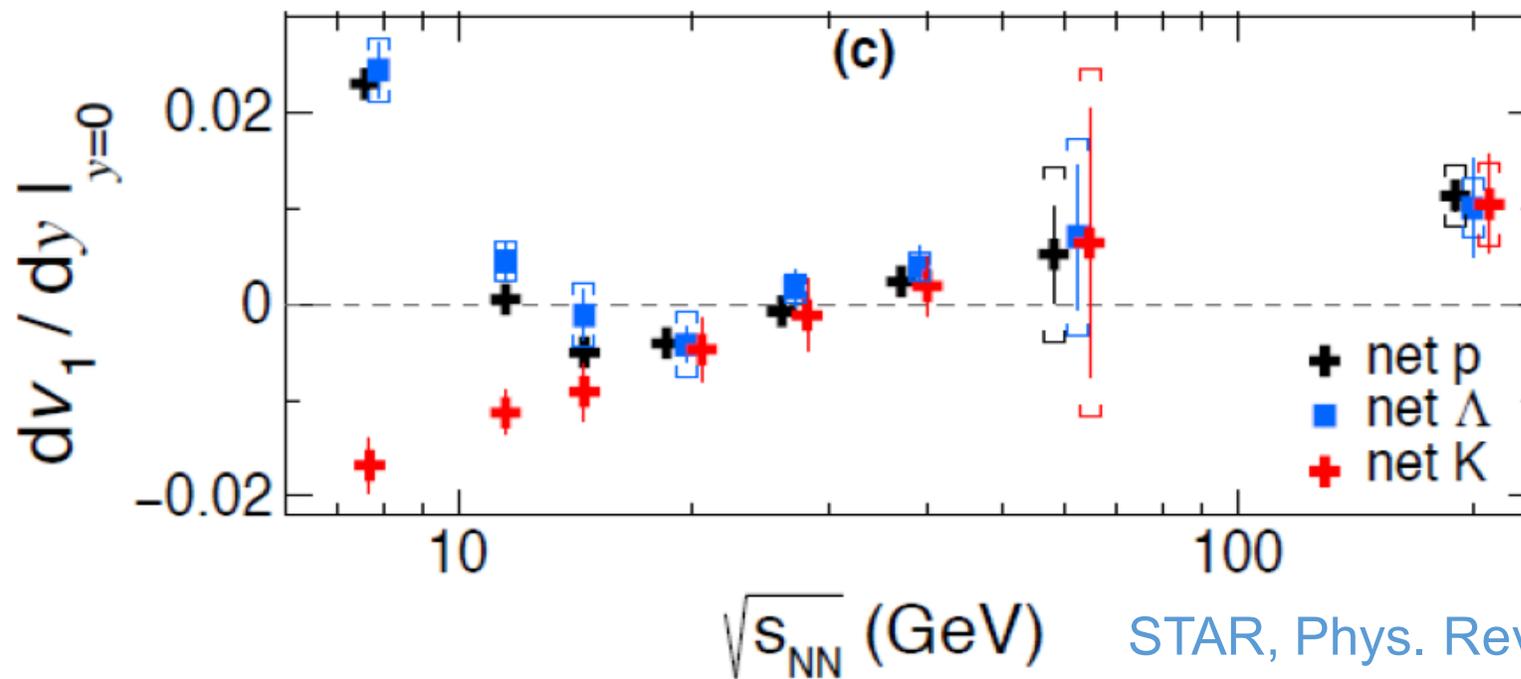
If the coalescence picture works, then...



- $\bar{u}(\bar{d})$  and  $s$  quarks have similar  $dv_1/dy$  at 200 GeV, and deviate at lower energies.
- $s$  and  $\bar{s}$  quarks are consistent with each other, except at the lowest energy.
  - At 7.7 GeV,  $v_1$  slope of  $\bar{s}$  is  $-0.097 \pm 0.023(\text{stat.}) \pm 0.026(\text{syst.})$  (far off the scale).

# $v_1(y)$ slope: net particles

- “Net particle” represents the excess yield of a particle species over its antiparticle.
- Net particles are more directly related to the transported quark number.



STAR, Phys. Rev. Lett. **120** (2018) 62301

$$[v_1(y)]_{net-p} = \frac{[v_1(y)]_p - r(y)[v_1(y)]_{\bar{p}}}{1 - r(y)}$$

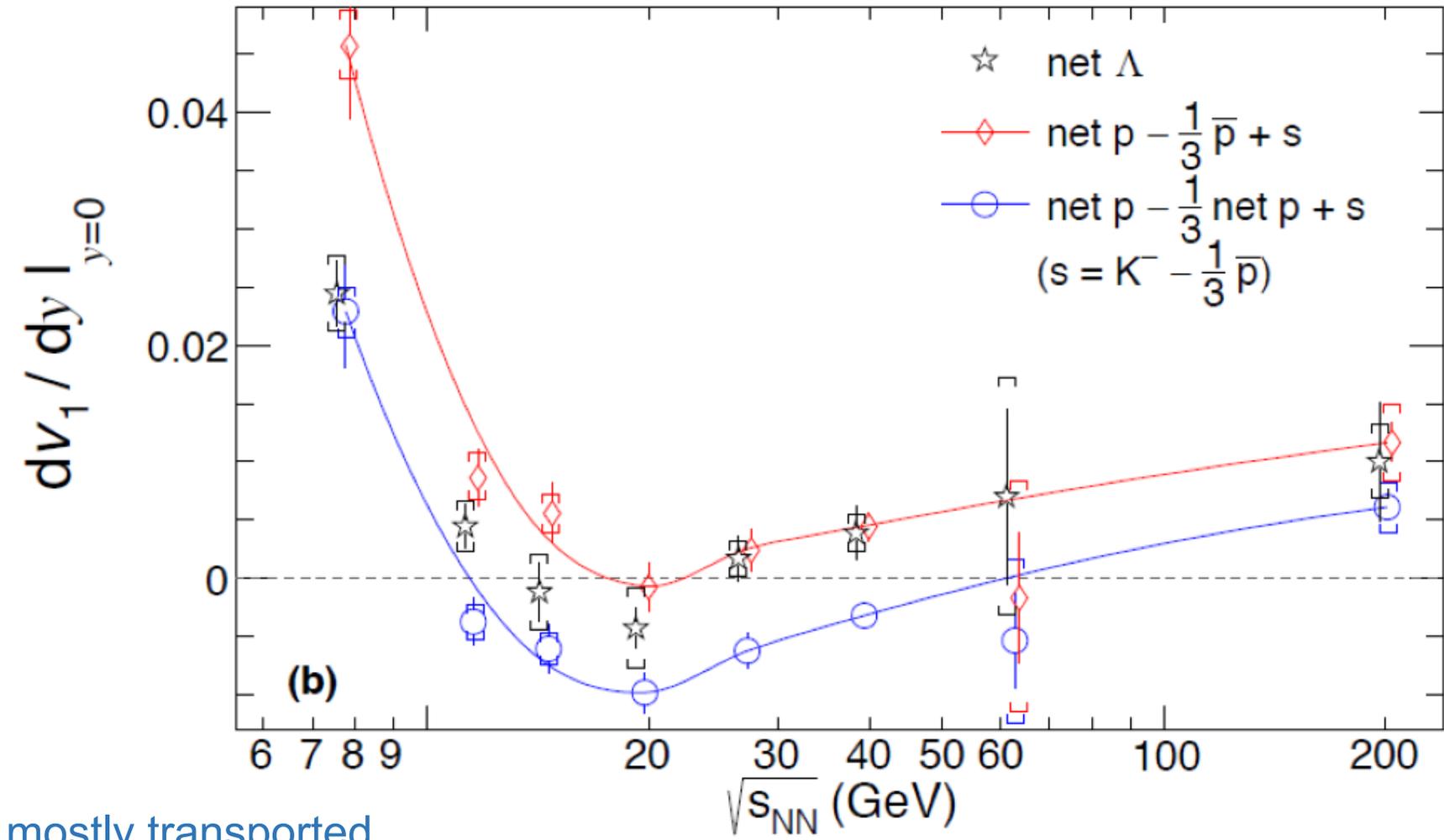
$r(y)$ : the ratio of the observed anti-p to p yield at each collision energy.

For net  $\Lambda$  and net K:

p is replaced with  $\Lambda$  and  $K^+$ , respectively;

anti-p is replaced with anti- $\Lambda$  and  $K^-$ , respectively.

# Coalescence sum rule: net particles



For net particles that contain transported quarks, we replace a **u** quark in net  $p$  with an **s** quark to reproduce net- $\Lambda$  in two scenarios.

- 1) the **u** quark (being replaced) was produced: works at higher energies.
- 2) all the quarks in net  $p$  have the same  $v_1$  (mostly transported): works at the lowest energy.

mostly transported  
(low energy limit)

STAR, Phys. Rev. Lett. **120** (2018) 62301

net  $p - \frac{1}{3} \text{net } p + s$

net  $p - \frac{1}{3} \bar{p} + s$

produced  
(high energy limit)

# Number of transported u+d per net proton

Assume  $\mu_{u(d)} = \mu_B / 3$ , and then statistically in a net proton,

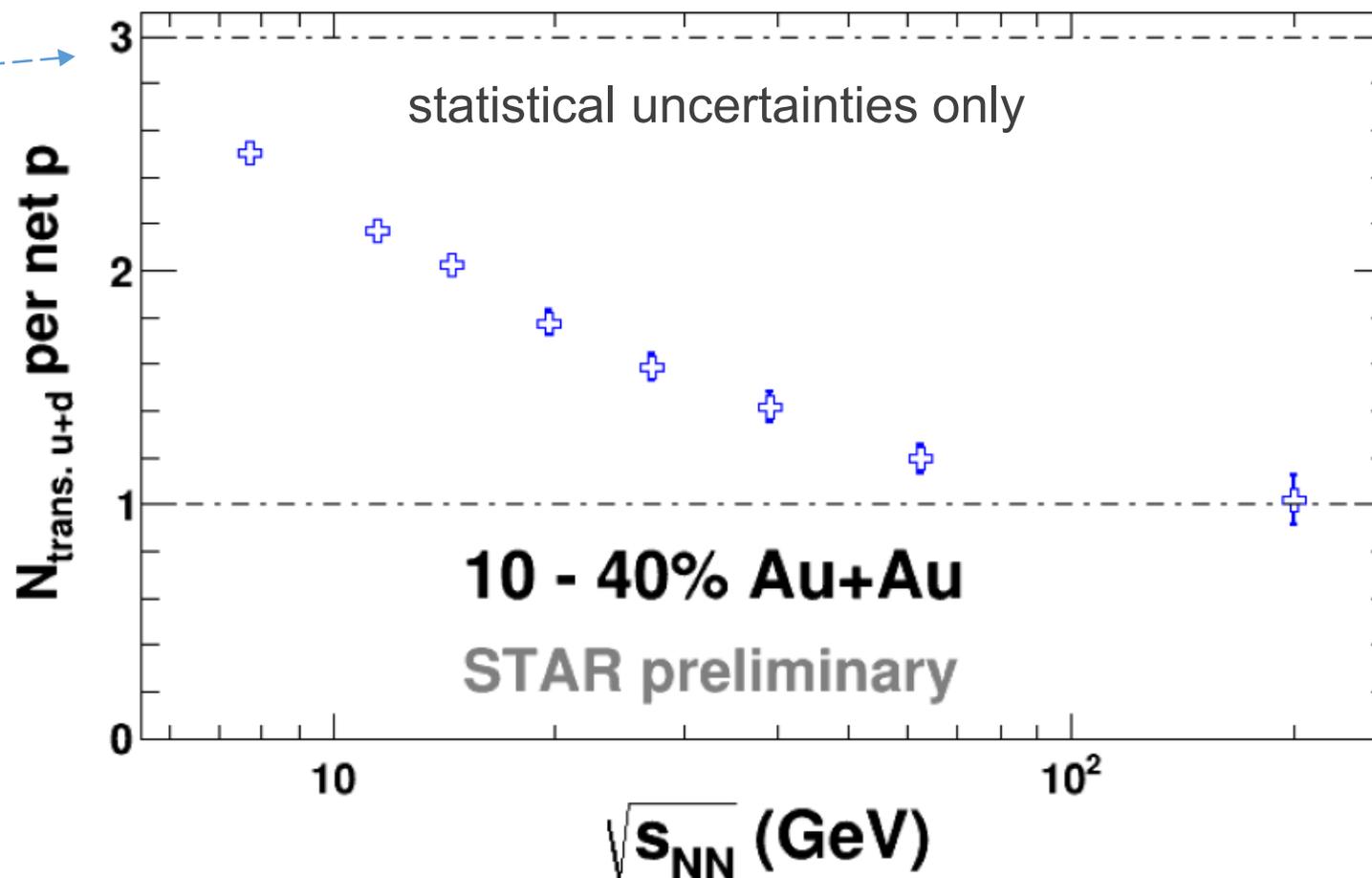
$$N_{\text{trans. } u+d} = 3 * [1 - \exp(- 2\mu_{u(d)} / T_{\text{ch}})] / (1 - r_{\text{anti-p/p}}).$$

$\mu_B$  and  $T_{\text{ch}}$  values have been published by STAR.

$r_{\text{anti-p/p}}$  is the ratio of observed anti-p to proton yield at each energy.

STAR, Phys. Rev. C  
96 (2017) 044904

low energy limit

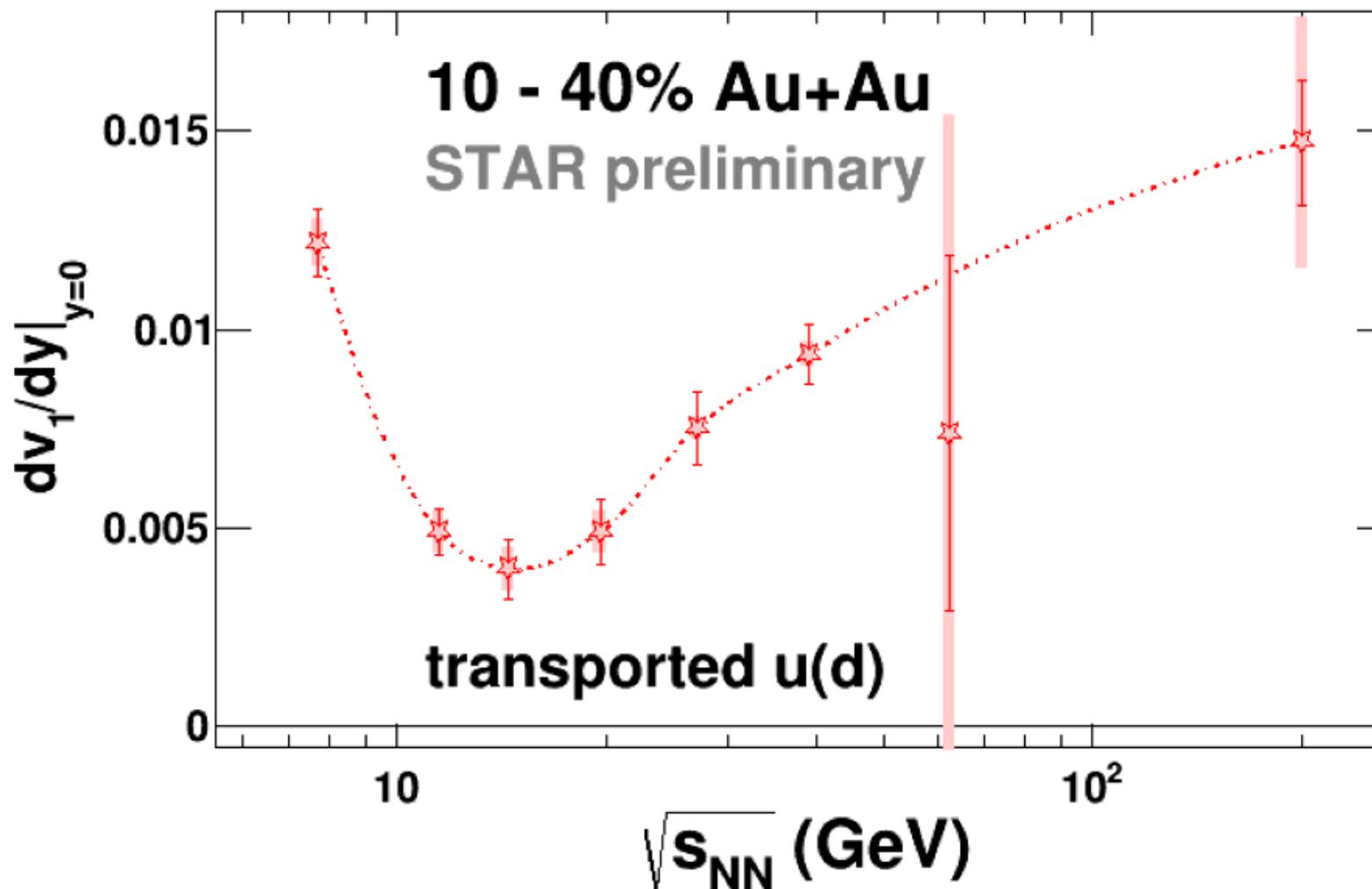


high energy limit

# $v_1$ of transported u(d) quarks

$$(v_1)_{\text{trans. u(d)}} = [(v_1)_{\text{net p}} - (3 - N_{\text{trans. u+d}}) * (v_1)_{\bar{u}(d)}] / N_{\text{trans. u+d}}$$

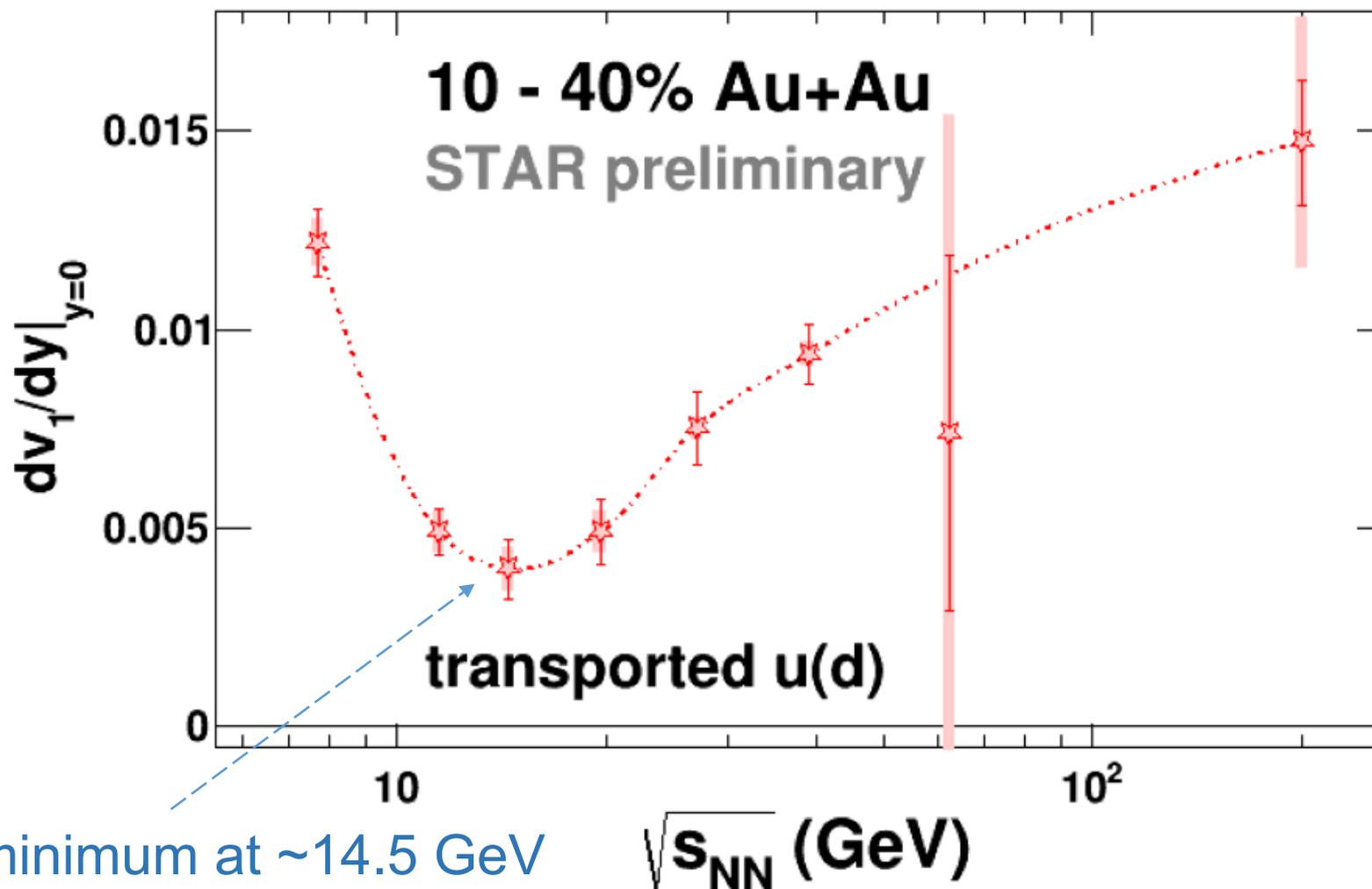
$v_1$  of transported **u(d)** is positive for all beam energies.



# $v_1$ of transported u(d) quarks

$$(v_1)_{\text{trans. u(d)}} = [(v_1)_{\text{net p}} - (3 - N_{\text{trans. u+d}}) * (v_1)_{\bar{\text{u(d)}}}] / N_{\text{trans. u+d}}$$

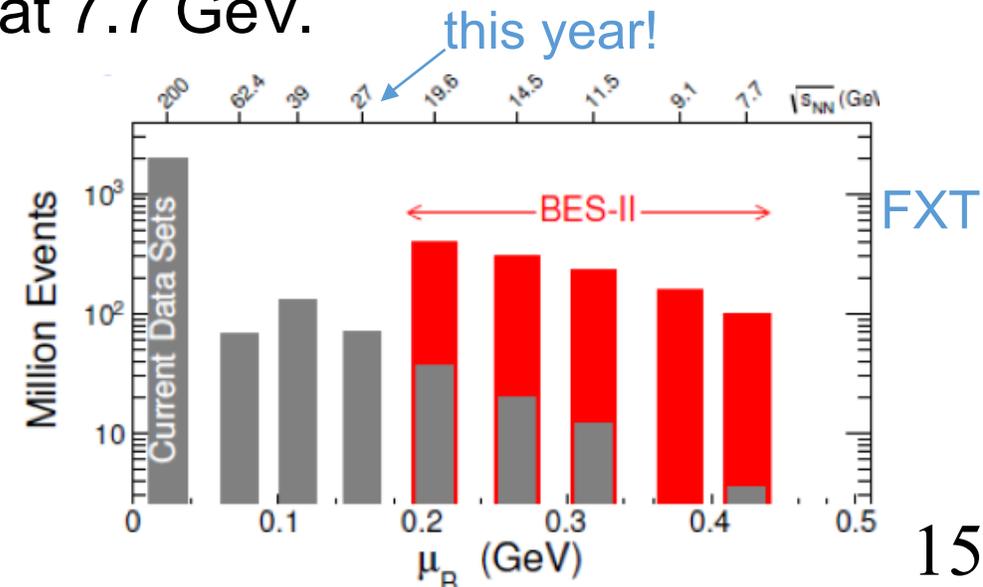
$v_1$  of transported **u(d)** is positive for all beam energies.



# Summary



- 10 species & 8 energies allow a detailed study of constituent-quark  $v_1$ .
- In most cases, the coalescence picture works for both “produced” particles and net particles.
  - the scaling behavior holds for produced quarks at and above 11.5 GeV, but breaks down at 7.7 GeV.  
*whose constituent quarks are all produced*
- $v_1$  of produced quarks
  - $u/\bar{u}/d/\bar{d}$  and  $s$  are close to each other at 200 GeV, and deviate at lower energies.
  - $s$  and  $\bar{s}$  are consistent with each other, except at 7.7 GeV.
- $v_1$  of transported quarks
  - $N_{\text{trans.}u+d}$  per net proton is estimated.
  - a minimum at  $\sim 14.5$  GeV.  
*evidence of the softest point?*
- BES-II & detector upgrades (EPD, eTOF, iTPC)



**Backup slides**

# Comparison: different types of quarks

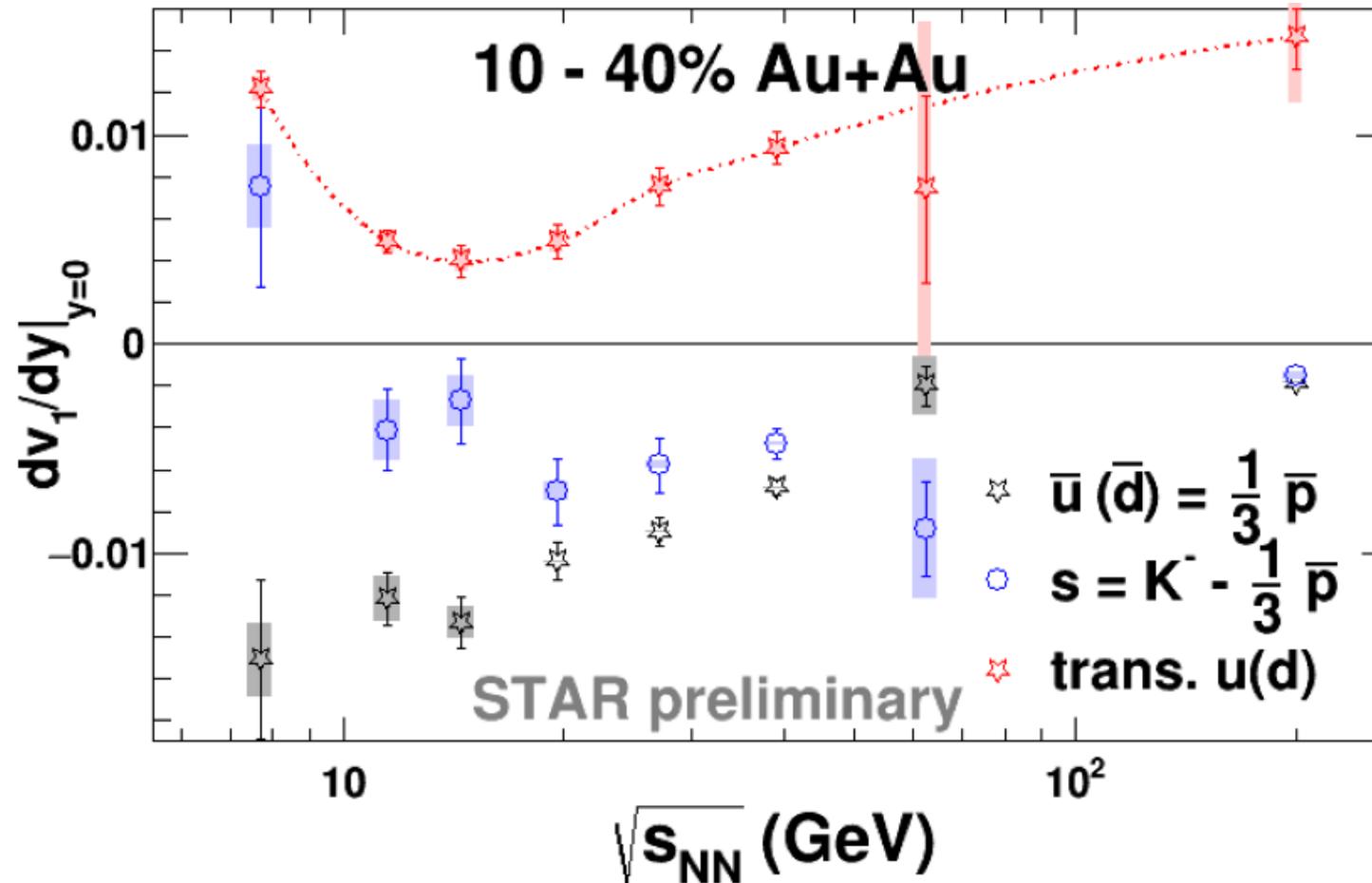


**s** and transported **u(d)** are close to each other at 7.7 GeV (the associated production?)



It looks like the **s** follows a transported **u**.

Then the **s** quark enters other strange hadrons in the medium.



# Coalescence sum rule



We have a set of assumptions, namely that  $v_1$  is imparted while quarks are deconfined, that specific types of quarks have the same  $v_1$  in the QGP, and that the detected hadrons form via statistical coalescence.

The coalescence sum rule should apply to not only  $v_1$ , but also all other  $v_n$ : the  $v_n$  of a particle is the sum of the  $v_n$  of its constituent quarks.

Take  $\pi^+(u\bar{d})$  as an example, and ignore the normalization:

$$\begin{aligned} v_n^{\pi^+} &= \int \int \int \cos(n\varphi^{\pi^+}) [1 + 2v_n^u \cos(n\varphi^u)] [1 + 2v_n^{\bar{d}} \cos(n\varphi^{\bar{d}})] \delta(\varphi^{\pi^+} - \varphi^u) \delta(\varphi^{\pi^+} - \varphi^{\bar{d}}) d\varphi^u d\varphi^{\bar{d}} d\varphi^{\pi^+} \\ &= \int \cos(n\varphi^{\pi^+}) [1 + 2v_n^u \cos(n\varphi^{\pi^+})] [1 + 2v_n^{\bar{d}} \cos(n\varphi^{\pi^+})] d\varphi^{\pi^+} \\ &\approx \int \cos(n\varphi^{\pi^+}) [1 + 2(v_n^u + v_n^{\bar{d}}) \cos(n\varphi^{\pi^+})] d\varphi^{\pi^+} \\ &= v_n^u + v_n^{\bar{d}} \end{aligned}$$

The  $\delta$ -functions are due to the coalescence mechanism. The observed  $N_{cq}$  scaling for  $v_2$  assumes all the quarks contribute the same  $v_2$ . In the  $v_1$ @BES study, we have less strict **assumptions**: we assume the produced  $u$ ,  $d$ ,  $\bar{u}$  and  $\bar{d}$  quarks all have the same  $v_1$ , and assume  $s$  and  $\hat{s}$  have the same  $v_1$ . However, the latter may not hold at low energies.

For transported quarks, we assume the transported  $u$  and  $d$  have the same  $v_1$ .

# Estimation with thermal equilibrium



With statistical thermal equilibrium, we have the chemical potentials of quarks:

$$\mu_u = \mu_B / 3 + 2 * \mu_Q / 3$$

$$\mu_d = \mu_B / 3 - \mu_Q / 3$$

$$\mu_s = \mu_B / 3 - \mu_Q / 3 - \mu_S$$

Courtesy of Jinfeng Liao

For each **u** quark, the fraction of the transported **u** quark is

$$f_u = [\exp(\mu_u/T) - \exp(-\mu_u/T)] / \exp(\mu_u/T) = 1 - \exp(-2 * \mu_u/T)$$

Boltzmann statistics is used here. Similar for  $f_d$  and  $f_s$ .

The number of transported quarks per net proton is then

$$(2 * N_p * f_u + N_p * f_d) / (N_p - N_{\text{anti-p}}) = (2 * f_u + f_d) / (1 - r_{\text{anti-p/p}})$$

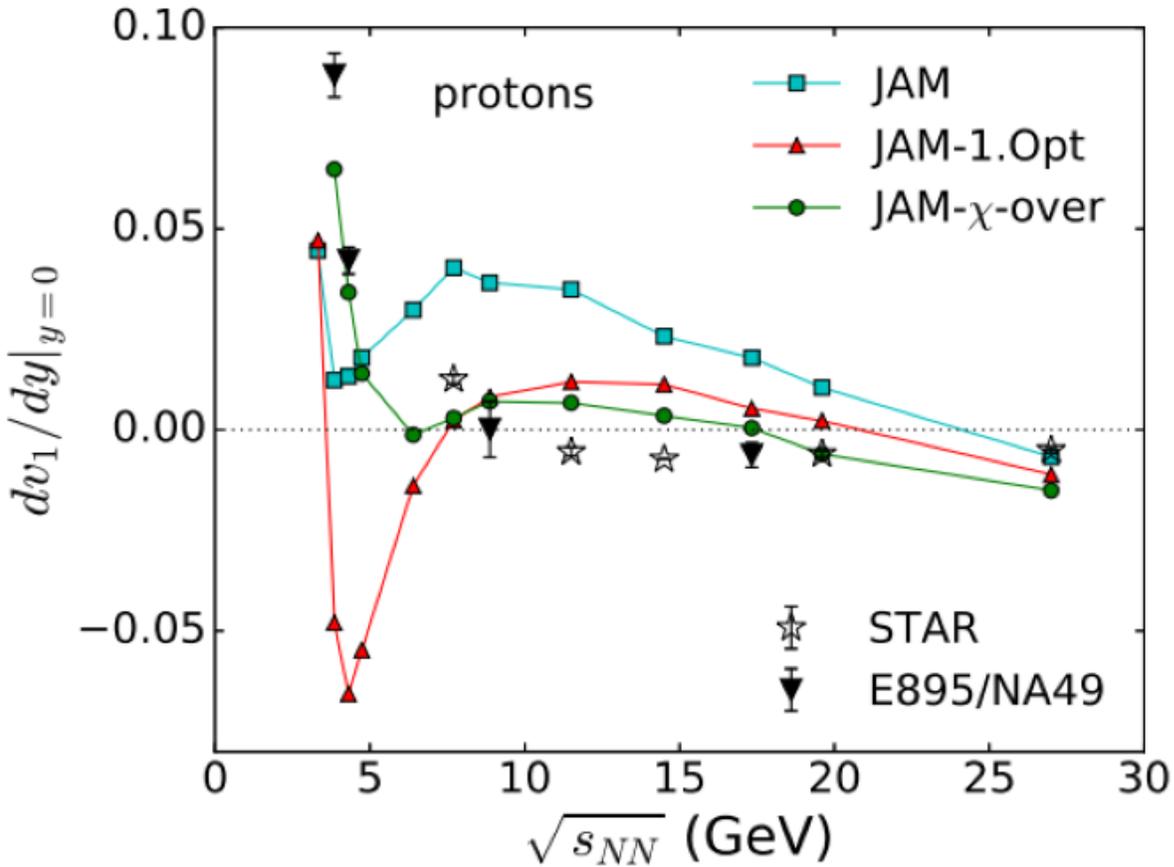
In our current data handling, we ignore the difference between **u** and **d**, so we take

$$\mu_{u(d)} = \mu_B / 3$$

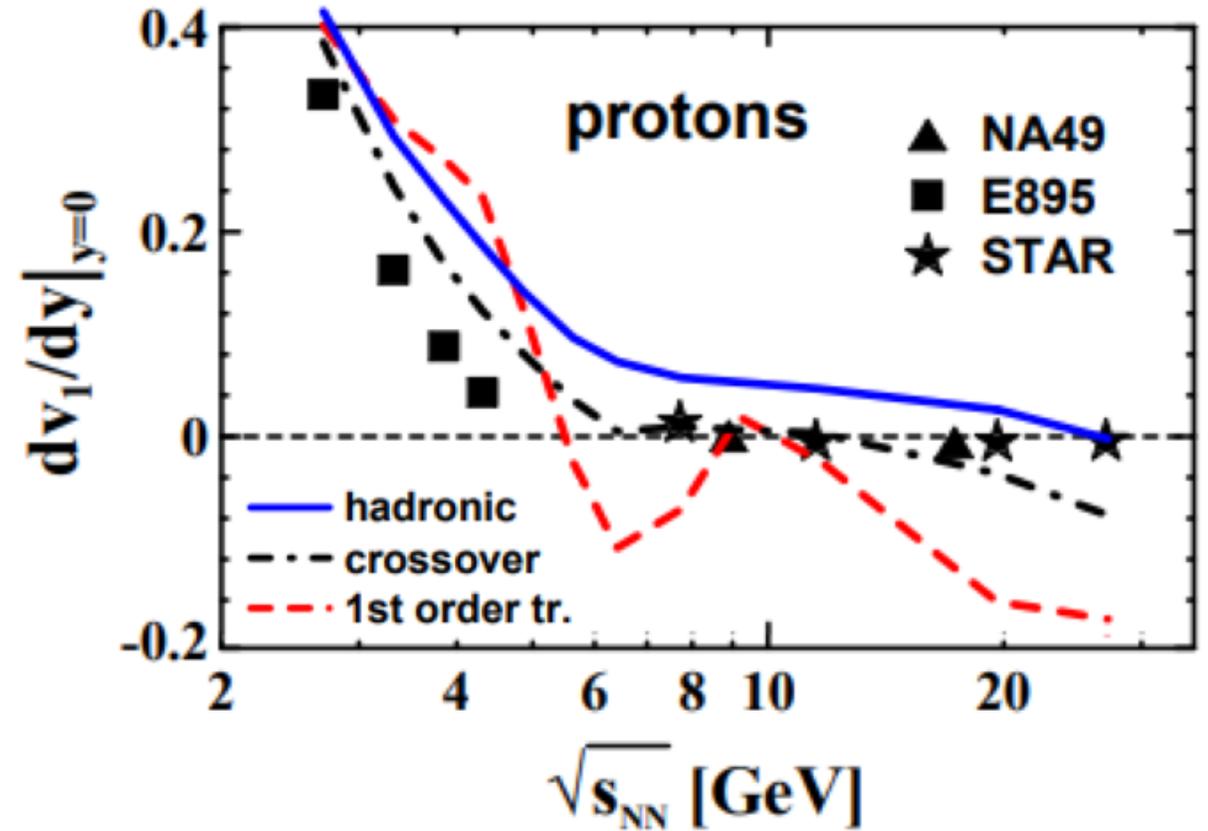
and on average in each net **p**,

$$N_{\text{trans. } u(d)} = 3 * [1 - \exp(-2 * \mu_{u(d)} / T_{\text{ch}})] / (1 - r_{\text{anti-p/p}})$$

$r_{\text{anti-p/p}}$  is the ratio of observed anti-**p** to proton yield at each energy.



Yasushi Nara, Harri Niemi, Jan Steinheimer, and Horst Stöcker, Phys. Lett. B769 (2017) 543. arXiv:1611.08023



3FD: Yu. B. Ivanov and A. A. Soldatov, Phys. Rev. C91 (2015) 024915; arXiv:1412.1669