

# Chiral phase transition of (2+1)-flavor QCD

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for  
HotQCD collaboration



- 1 Introduction
- 2 Scaling analysis : basic definitions and some insights
- 3 Estimation of  $T_c^0$  from generic scaling arguments
- 4 Summary and Outlook

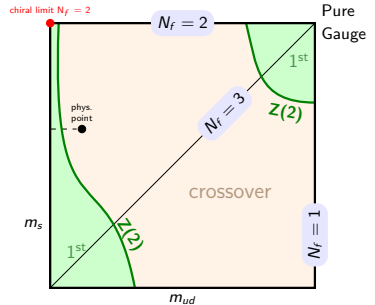
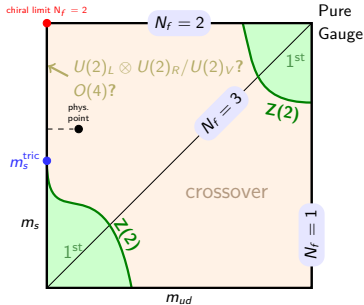
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- In QCD, transition for physical quark masses and vanishing or small baryon density is a crossover.
- $T_{pc}$  for physical mass QCD has been determined with very good accuracy from various observables :  $T_{pc}(\mu_B = 0) = 156.5 \pm 1.5$  MeV  $\Rightarrow$  “QCD transition at zero and non-zero baryon densities” by Patrick Steinbrecher on 16 May 2018, 11:30 for recent HotQCD results.
- Important question : How do singular terms of the chiral phase transition at vanishing quark mass affect observables at physical values of the quark mass?  $\Rightarrow$  How far the physical regime is from chiral limit w.r.t. scaling window?
- Key question : What is the chiral transition temperature,  $T_c^0$ ?

# Introduction?

In this talk, we are trying to address the followings :

- Key question : **What is the chiral transition temperature,  $T_c^0$ ?**
- Possibly another question : **What is the nature of the thermal phase transition in the chiral limit?**
- Two possible scenarios : [O. Philipsen and C. Pinke. Phys. Rev. D93, 114507, 2016.]



- We are trying to pin down the correct one.

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In terms of temperature  $T$  and symmetry breaking field  $H = m_l/m_s$  the scaling variables are defined as :

$$t = \frac{1}{t_0} \frac{T - T_c^0}{T_c^0} \quad \text{and} \quad h = \frac{1}{h_0} \frac{m_l}{m_s} = \frac{1}{h_0} H$$

Scaling variable :

$$z = \frac{t}{h^{\frac{1}{\beta\delta}}} = z_0 \left( \frac{T - T_c^0}{T_c^0} \right) \left( \frac{1}{H^{1/\beta\delta}} \right); \quad z_0 = \frac{h_0^{\frac{1}{\beta\delta}}}{t_0}$$

$$\text{Chiral condensate} : \langle \bar{\psi}\psi \rangle_f = \frac{T}{V} \frac{\partial \ln Z}{\partial m_f}$$

$$\text{Chiral susceptibility} : \chi_m^{fg} = \frac{\partial}{\partial m_g} \langle \bar{\psi}\psi \rangle_f$$

Renormalization group invariant (RGI) definition of order parameter :

$$M = \frac{m_s}{f_K^4} \left( (\langle \bar{\psi}\psi \rangle_u + \langle \bar{\psi}\psi \rangle_d) - \frac{m_u + m_d}{m_s} \langle \bar{\psi}\psi \rangle_s \right)$$

RGI definition of order parameter susceptibility :

$$\chi_M = \frac{T}{V} m_s \left( \frac{\partial}{\partial m_u} + \frac{\partial}{\partial m_d} \right) M$$

Close to chiral limit, singular part behaves as :

$$\begin{aligned} M &= h^{1/\delta} f_G(z) \\ \chi_M &= \frac{1}{h_0} h^{1/\delta-1} f_\chi(z) \end{aligned}$$

$f_G(z)$  and  $f_\chi(z)$  are universal scaling functions which have been precisely determined from various spin models.



## Scaling functions : Some intriguing facts

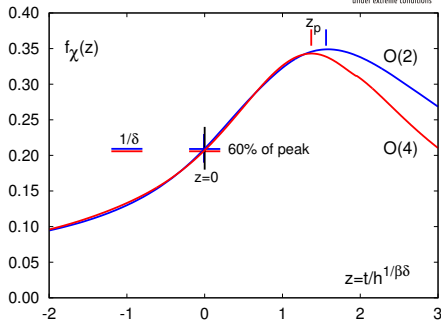
Conventional approach to determine  $T_c^0$  : determine pseudo-critical temperature  $T_{pc}(H)$  from the peak location of  $\chi_M(T, H)$  and

$$T_{pc}(H) = T_c^0 \left( 1 + \frac{z_p}{z_0} H^{1/\beta\delta} \right)$$

Our approach in this work : determine temperature  $T_{60\%}^-(H)$  at 60% of peak height of  $\chi_M(T, H)$  and

$$T_{60\%}^-(H) = T_c^0 \left( 1 + \frac{z_{60\%}^-}{z_0} H^{1/\beta\delta} \right)$$

Dependence on quark mass (H) reduced by two orders of magnitude



Advantage :  $z_{60\%}^- \simeq 0 \Rightarrow$  reduces influence of regular terms, simplifies scaling analysis.

$\Rightarrow$

	$z_p$	$z_{60\%}^-$
$O(2)$	1.56	-0.009
$O(4)$	1.37	-0.01

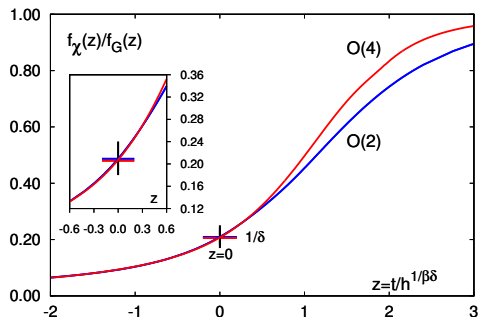
# Scaling functions : Some intriguing facts

- Other approach used for consistency check : Crossing points of

$$\frac{H\chi_M}{M} = \frac{f_\chi(z)}{f_G(z)} + \text{regular terms.}$$

is unique for  $H \rightarrow 0$ .

- Ratio has a curvature in  $z \Rightarrow$  Non-linear part of scaling function is important.

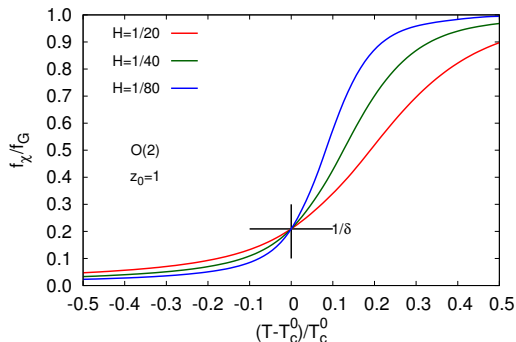


	$1/\delta$
$O(2)$	0.209
$O(4)$	0.207

Difference between  $O(4)$  and  $O(2)$  is negligible in the  $z$  range we are working.

# Scaling functions : Some intriguing facts

$$\frac{f_{\chi}(z)}{f_G(z)} = \left\{ \begin{array}{ll} 0 & , z \rightarrow -\infty \\ 1/\delta & , z = 0 \\ 1 & , z \rightarrow +\infty \end{array} \right\}$$



- Ratio is expected to have a constant value at the crossing point,  $z = 0$ , *i.e.* in chiral limit at  $T_c^0$ .
- Uniqueness of the crossing point get spoiled in presence of regular terms.
- Around  $T_c^0$  :  
slope  $\propto H^{-1/\beta\delta}$  and  
curvature  $\propto H^{-2/\beta\delta}$ .

- Determine temperature  $T_{\delta}(H)$  which satisfies :

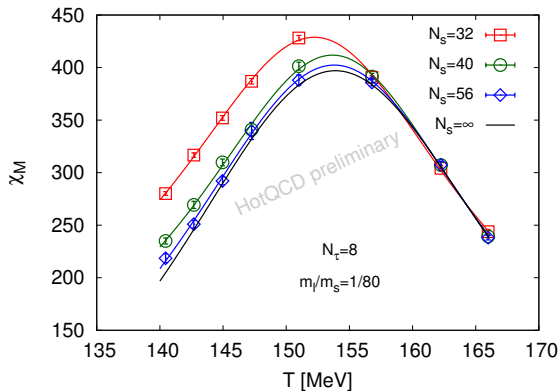
$$\frac{H\chi_M(T_{\delta}, H)}{M(T_{\delta}, H)} = \frac{1}{\delta} \Rightarrow T_c^0 = \lim_{H \rightarrow 0} T_{\delta}(H)$$

- Here we will mainly show results for  $N_\tau = 8$  only.
- Same analyses also carried out for  $N_\tau = 6$ .
- Same analyses for  $N_\tau = 12$  are still at preliminary stage.
- Our goal is to find out the chiral transition temperature,  $T_c^0$ , in continuum limit.
- More lattices are still being generated for low masses and the analyses are not settled yet for those parameter sets.
- We are working on the systematic error estimation for various analyses and most of the results shown here are just indicative.

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# No evidence for 1<sup>st</sup> order transition

- Volume dependence of  $\chi_M$  is studied for  $H = 1/80$  which corresponds to  $m_\pi = 80$  MeV in continuum.
- $\chi_M^{\max}$  is not proportional to volume.



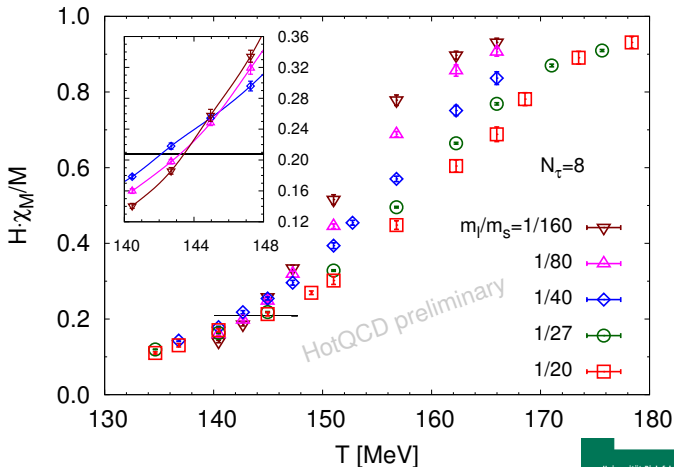
- Possibility for 1<sup>st</sup> order phase transition can be ruled out at  $m_\pi = 80$  MeV for  $N_\tau = 8$ .
- We only focus on  $O(2)$ ,  $O(4)$  and  $Z(2)$ .
- Similar results are also obtained for  $N_\tau = 6$  and 12.
- $N_\tau = 12$  :  $42^3 \times 12$  and  $60^3 \times 12$  lattices have been used for  $m_l/m_s = 1/40$ .

# Estimation of $T_c^0$ from $H\chi_M/M$

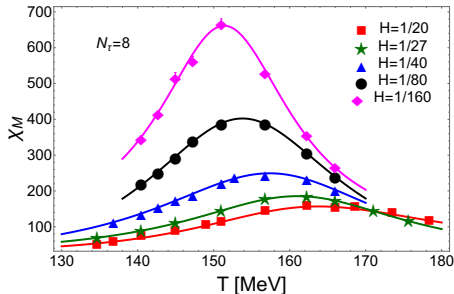
- Behavior of the ratio is like Binder cumulant at critical point.

[F. Karsch and E. Laermann. Phys. Rev. D50, 6954, 1994.]

- For  $H = 1/27$  and  $1/20$ , regular contributions may be larger than lowest three masses.
- Crossing points for three lowest masses gives an estimate  $T_c^0 \sim 144$  MeV for  $N_\tau = 8$ .



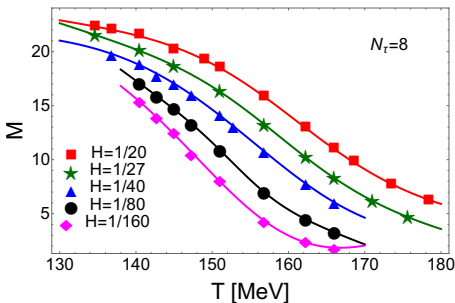
# Fitting $\chi_M^{f_K}$ and $M^{f_K}$ with rational functions



- $\chi_M$  for different  $H$  are fitted with rational functions :

- 1 peak location :  $T_{pc}(H)$ .
- 2 peak height :  $\chi_M^{\max}(H)$ .
- 3  $T_{60\%}^-$ , defined by :

$$\chi_M(T_{60\%}^-) = \frac{60}{100} \chi_M^{\max}$$

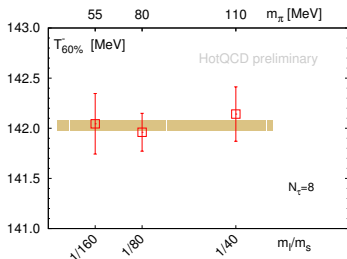


- $M$  at  $T_{pc}$  and  $T_{60\%}^-$  are calculated for different  $H$  by fitting with rational functions.

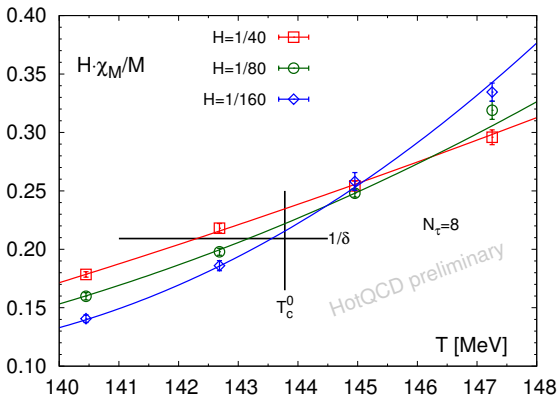


## Sanity checks for $T_c^0$

- $T_{60\%}^-$  for three lowest masses indicates  $T_c^0 \sim 142$  MeV.



- Joint scaling fit of  $M$  and  $\chi_M$  to MEOS with  $O(2)$  exponents also gives  $T_c^0 \sim 145$  MeV for  $N_\tau = 8$ .

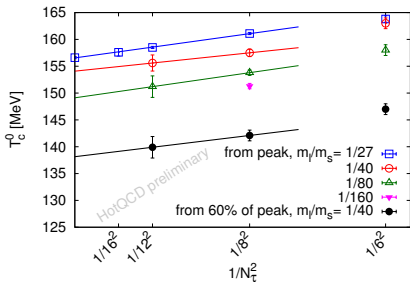
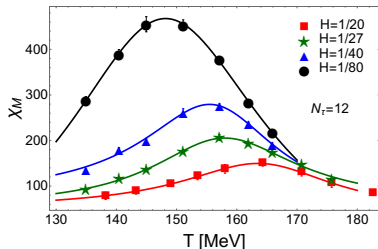


- Simultaneous fit (with regular term) of  $H \cdot \chi_M / M$  to the three lowest masses gives  $T_c^0 \sim 144$  MeV for  $N_\tau = 8$ .

Current estimate of  $T_c^0$  for  $N_\tau = 8$  is 144(2) MeV.

# Towards continuum $T_c^0$

- Current estimate of  $T_c^0$  is 147(2) MeV for  $N_\tau = 6$  and 138(3) MeV for  $N_\tau = 12$ .
- $T_{60\%}^-$  is approximately independent of  $H$  establishes it as a good estimator of  $T_c^0$ .



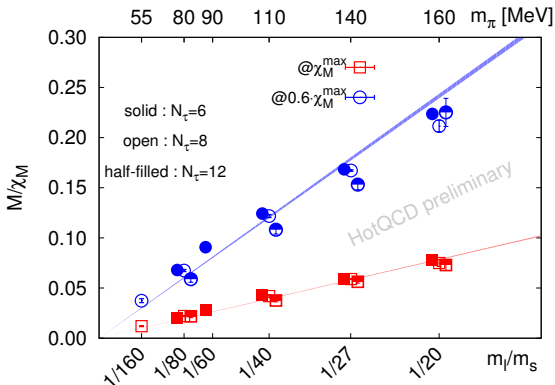
- Calculations at different  $m_l/m_s$  suggest that cut-off effects are only weakly (not) dependent on quark mass.
- Experience from  $m_l/m_s = 1/27$  calculations suggest that extrapolation to continuum limit based on  $N_\tau = 8$  and 12 is safe.

We find  $T_c^0 = 138(5)$  MeV in continuum.

# Order of chiral transition : work in progress....

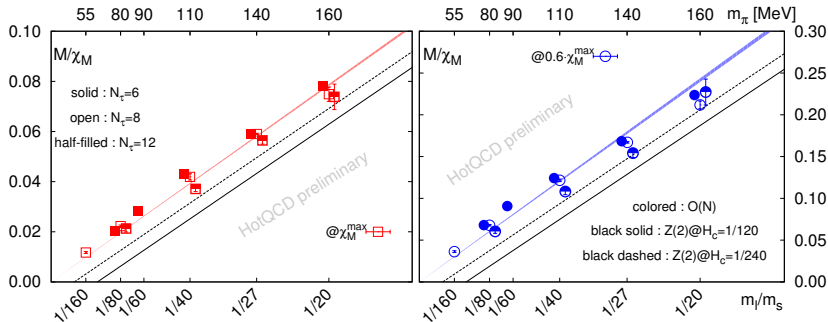
$$\frac{M}{\chi M} = H \frac{f_G(z)}{f_\chi(z)}$$

- For small  $H$  the data seems to be linear.
- Lines are NOT fitted curves rather expectations for  $O(2)$  and  $O(4)$ .



- $Z(2)$  transition, at some finite  $H_c$ , will result into a strong  $V$  dependence and sudden drop in the ratio  $\Rightarrow$  1<sup>st</sup> order transition is unlikely for  $m_\pi > 55$  MeV.
- Additional low  $H$  measurements : slope can be directly determined as fit parameter.

# Order of chiral transition : work in progress....



- $Z(2)$  lines are schematic.
- If  $M$  is not exactly order parameter then the  $Z(2)$  lines will have a curvature.
- Mixing becomes weak as  $H_c$  becomes small.
- Data seems to favor  $O(N)$  compared to  $Z(2)$ .

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- Summary :
  - ① Measurements are according to the expectations even without being specific about the universality class.
  - ② Proposed method can distinguish between  $Z(2)$  and  $O(N)$ .
  - ③ Scaling fits with  $O(2)$  exponents worked reasonably.
  - ④ Current estimate of  $T_c^0$ , in continuum, is 138(5) MeV.
- Outlook :
  - ① Finer lattices are being generated for  $H = 1/40$  and  $1/80$ .
  - ② Calculations at more masses have been planned to have better control on the functional dependence of various quantities.
  - ③ Additional low  $H$  measurements will help us to be confident about the scaling behavior.

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