

# Shear viscosity and resonance lifetimes in the hadron gas

Jean-Bernard Rose

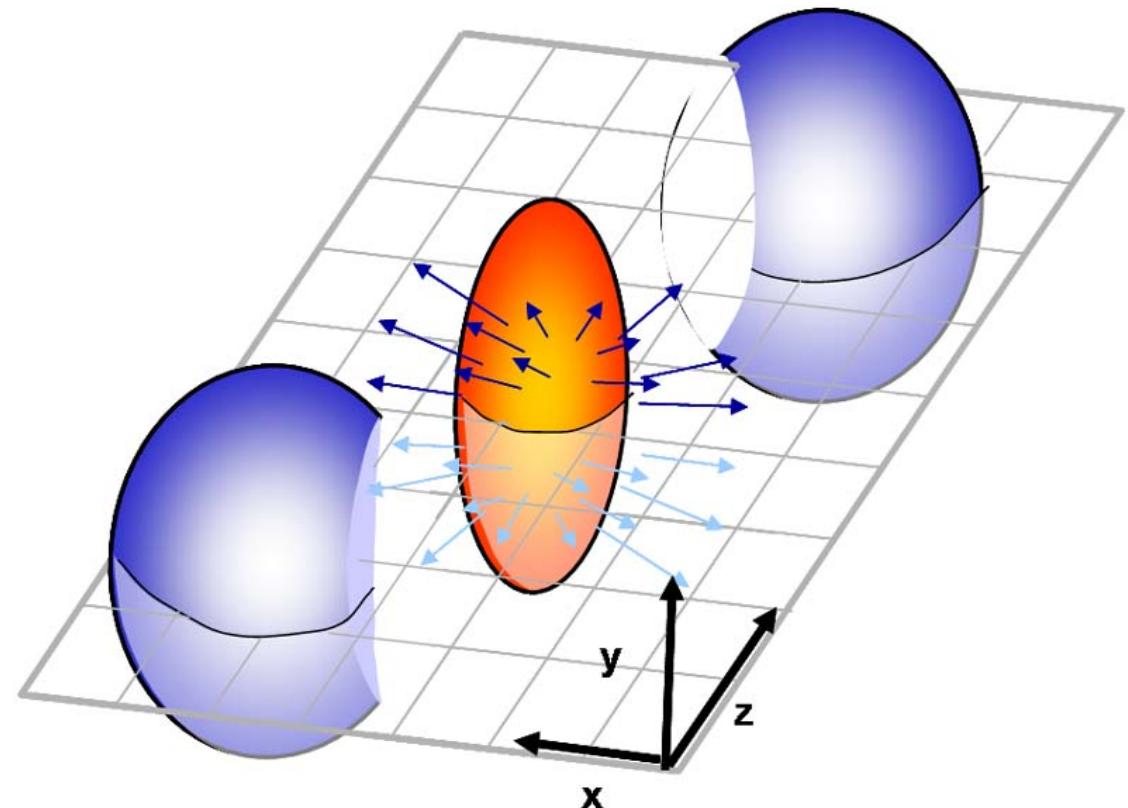
with D. Oliinychenko, J. Torres-Rincon, A. Schäfer,  
J. Hammelmann, H. Petersen

based on Phys Rev C 97.055204 (arXiv:1709.03826) and arXiv:1709.00369



# Viscosity in heavy ion collisions

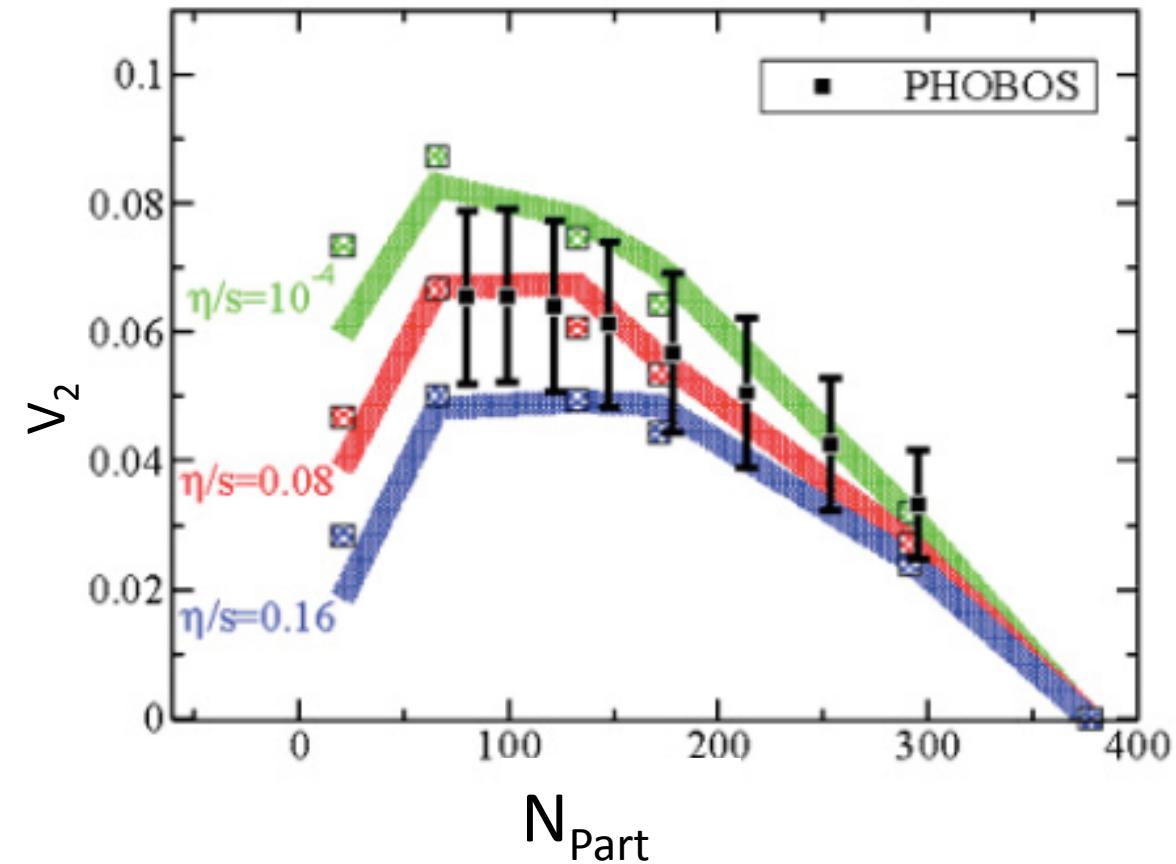
- RHIC and LHC measured large elliptic flow at the high energies corresponding to what is thought to be QGP



<http://www.quantumdiaries.org/wp-content/uploads/2011/02/FlowPr.jpg>

# Viscosity in heavy ion collisions

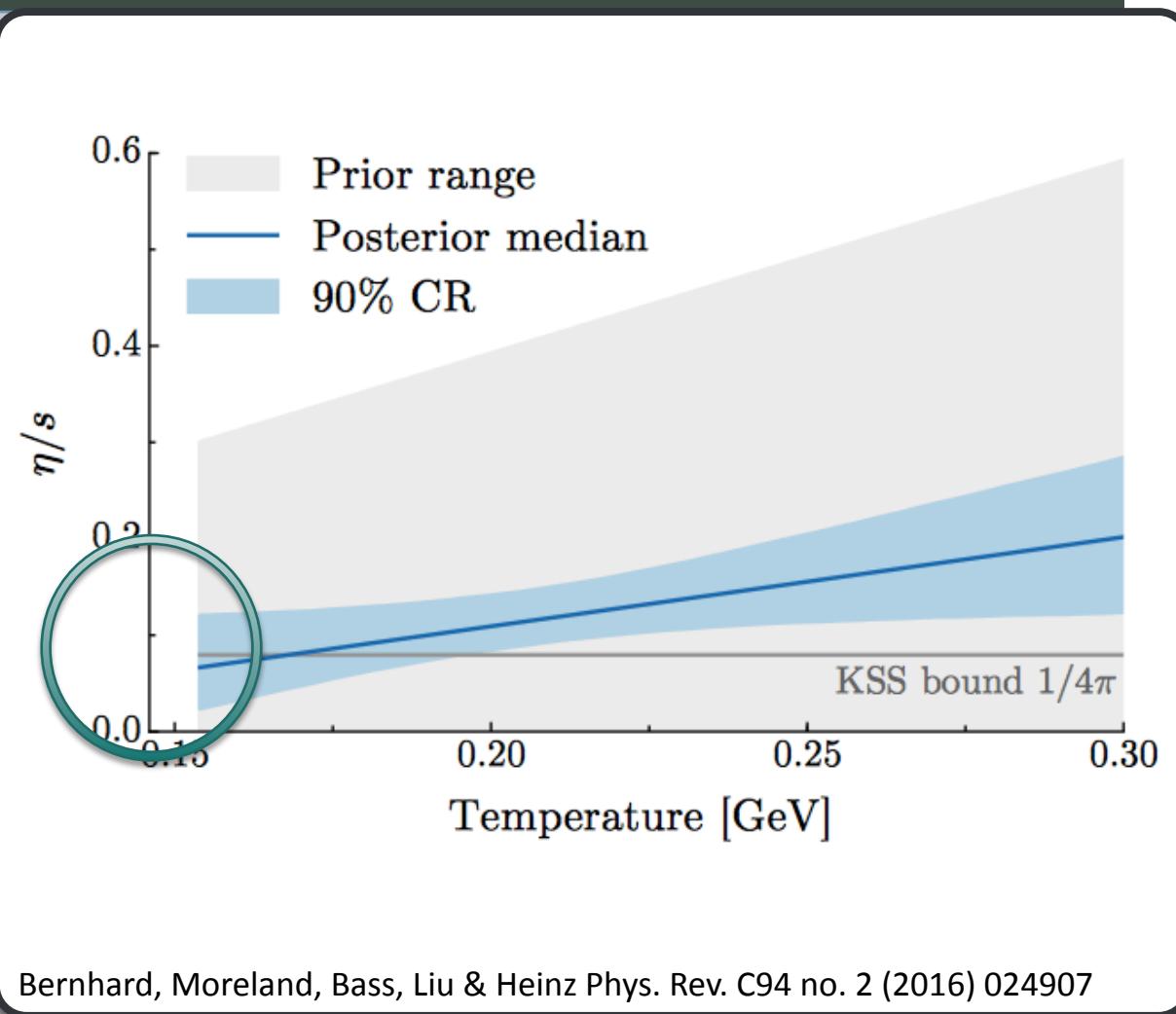
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- Hydrodynamics relatively successful at explaining this with small  $\eta/s$



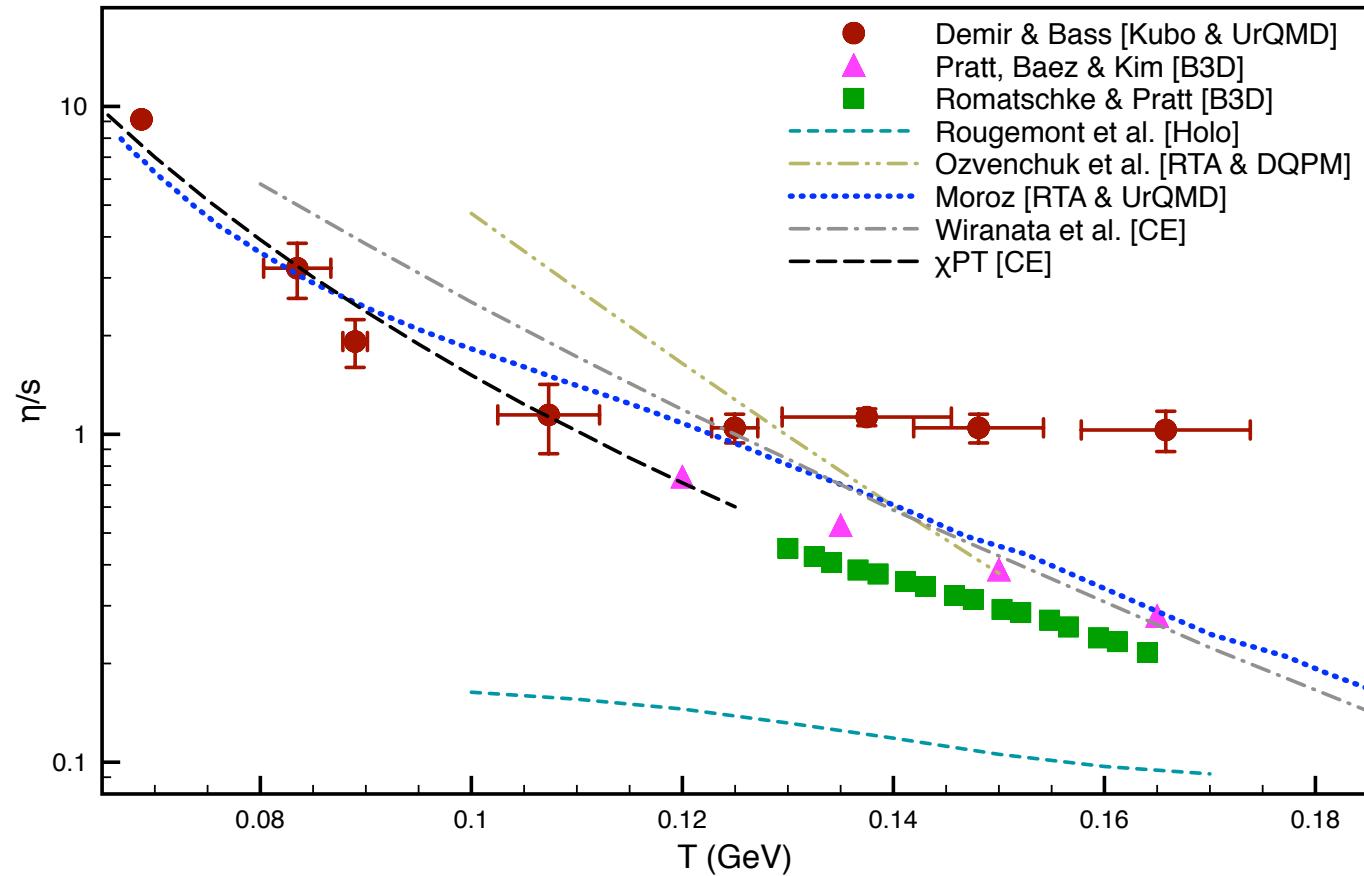
Luzum & Romatschke 10.1103/Phys. Rev. C 78.034915

# Viscosity in heavy ion collisions

- RHIC and LHC measured large elliptic flow at the high energies corresponding to what is thought to be QGP
- Hydrodynamics relatively successful at explaining this with small  $\eta/s$
- Still not clear what the behavior of  $\eta/s$  is at low energies (FAIR, late stage RHIC/LHC)

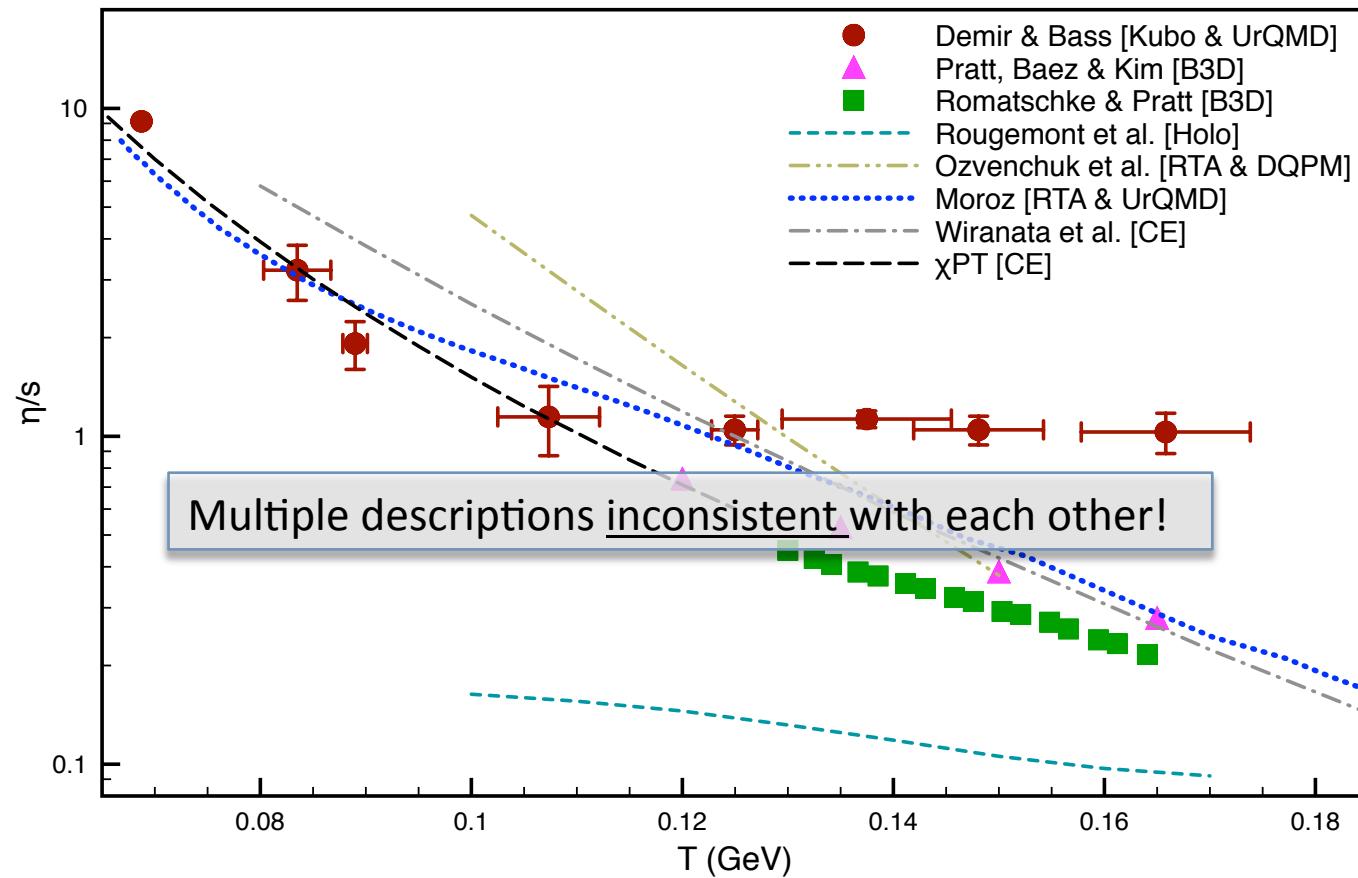


# Previous HG viscosity calculations



- Demir & Bass, Phys. Rev. Lett. 102 (2009) 172302
- Pratt, Baez & Kim, Phys. Rev. C95 (2017) 024901
- Romatschke & Pratt, arXiv:1409.0010v1
- Rougemont et al., arXiv:1704.05558
- Ozvenchuk et al., Phys. Rev. C87 (2013) 064903
- Moroz, arXiv:1301.6670
- Wiranata et al., Phys. Rev. C88 (2013) 044917
- Torres-Rincon, arXiv:1205.0782

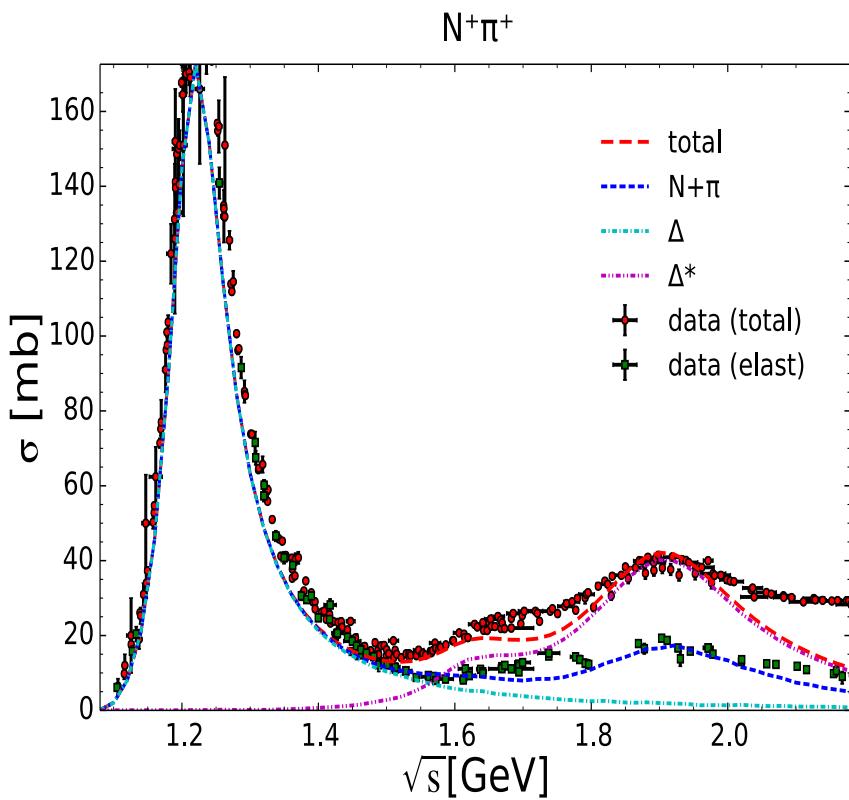
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# Modelling the hadron gas: SMASH

See talk by H. Petersen earlier today at 12:30 in Collective Dynamics session



- **SMASH is a new semi-classical transport approach for the hadron gas**
- Geometric collision criterion:
$$d_{trans} < d_{int} = \sqrt{\frac{\sigma_{tot}}{\pi}}$$
- Spectral functions of resonances are described by relativistic Breit-Wigner functions, with resonance lifetime
$$\tau_{res} = \frac{1}{\Gamma(m)}$$
- Elastic scatterings parameterized for NN; all other elastic scatterings assumed to go through resonances
- Inelastic scatterings, currently include
  - NN↔NR, NN↔ΔR
  - KN↔KN, KN↔πH
  - +antiparticles
- Strings (turned off for detailed balance)

# Green-Kubo Formalism

The shear viscosity  
is calculated from

$$\eta = \frac{V}{T} \int_0^\infty C^{xy}(t) dt$$

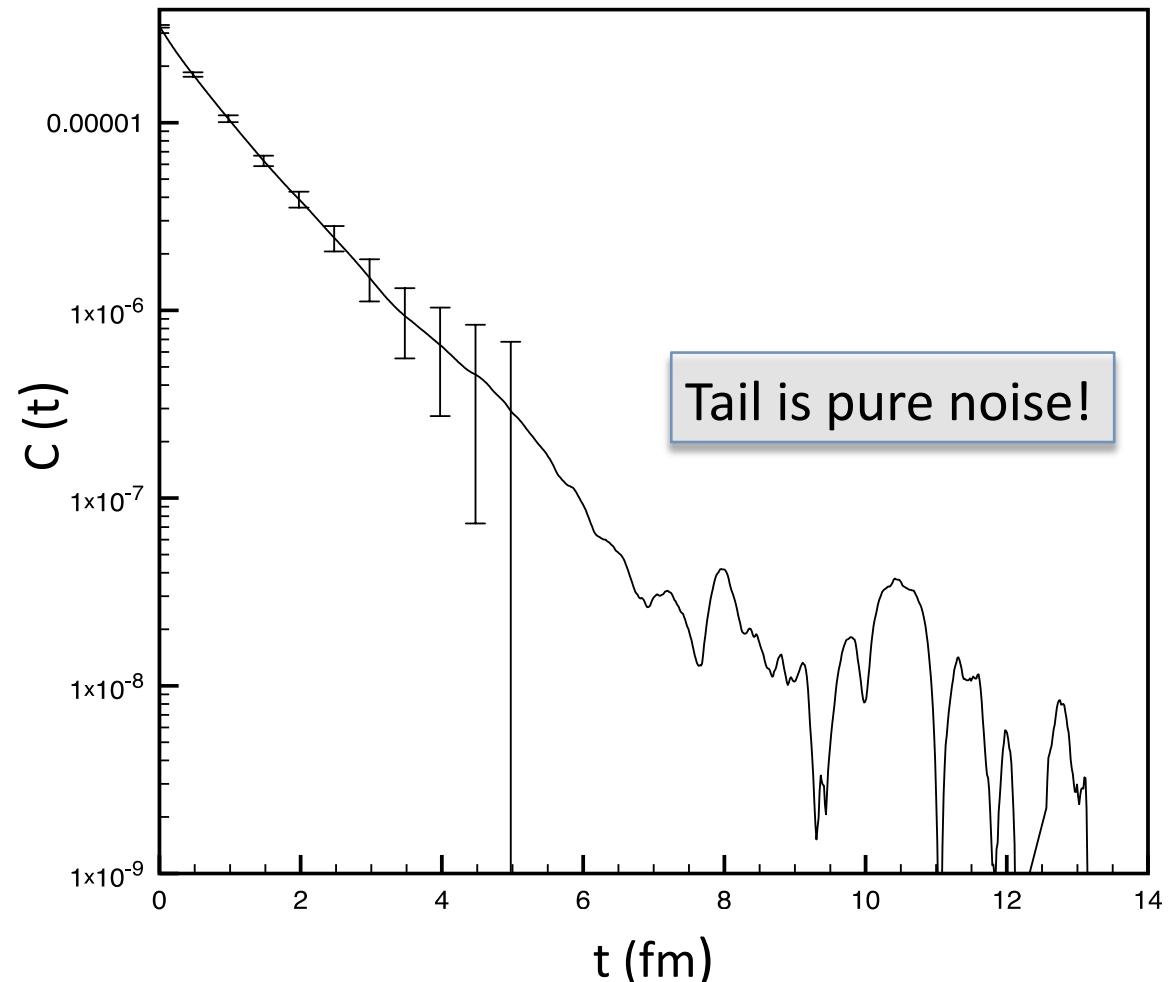
where

$$C^{xy}(t) = \frac{1}{N} \sum_s T^{xy}(s) T^{xy}(s+t)$$

and

$$T^{\mu\nu} = \frac{1}{V} \sum_i^{N_{part}} \frac{p_i^\mu p_i^\nu}{p_i^0}$$

$N$  is the number of time steps, and  $N_{part}$  the number of particles



# Green-Kubo Formalism

It has been shown that the correlation

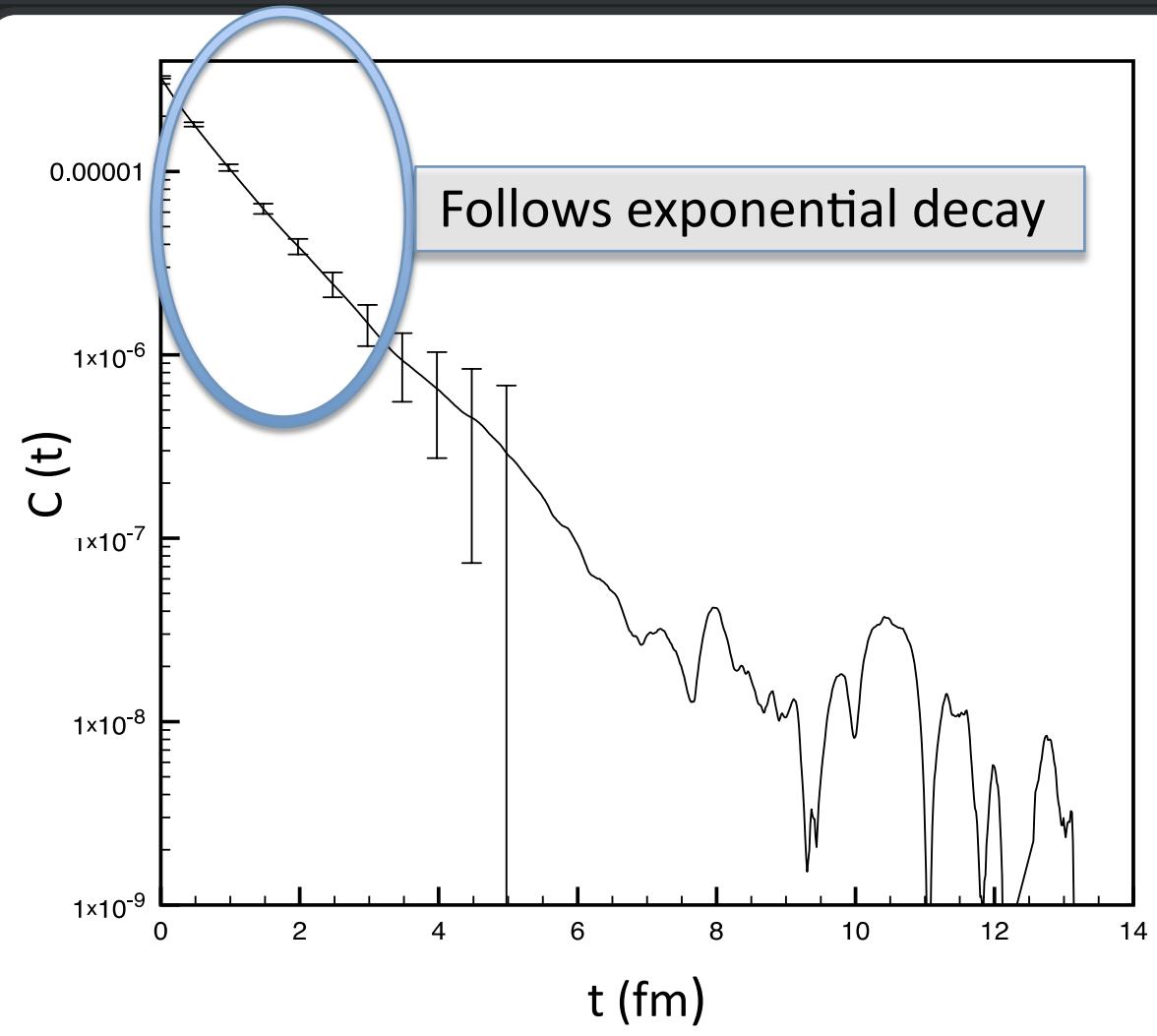
$$\eta = \frac{V}{T} \int_0^{\infty} C^{xy}(t) dt$$

Follows

$$C^{xy}(t) = C^{xy}(0) \exp\left(-\frac{t}{\tau}\right)$$

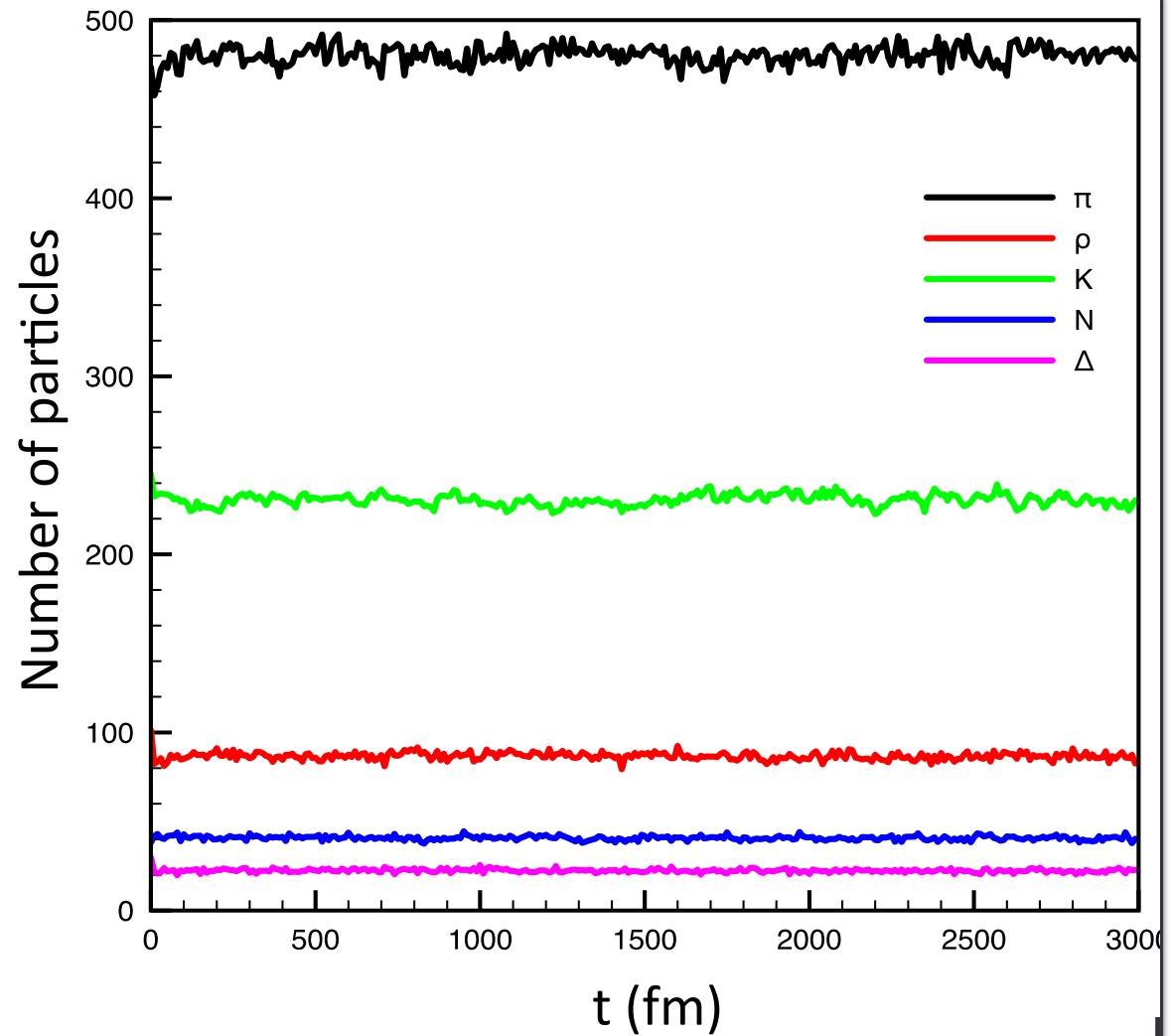
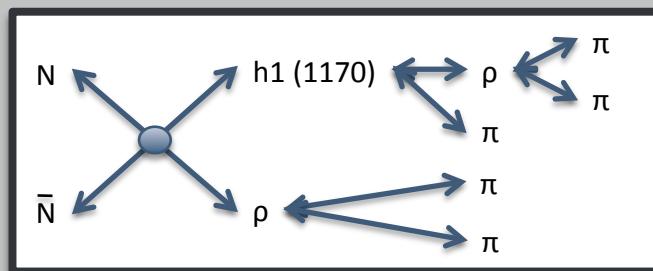
So that

$$\eta = \frac{VC^{xy}(0)\tau}{T}$$

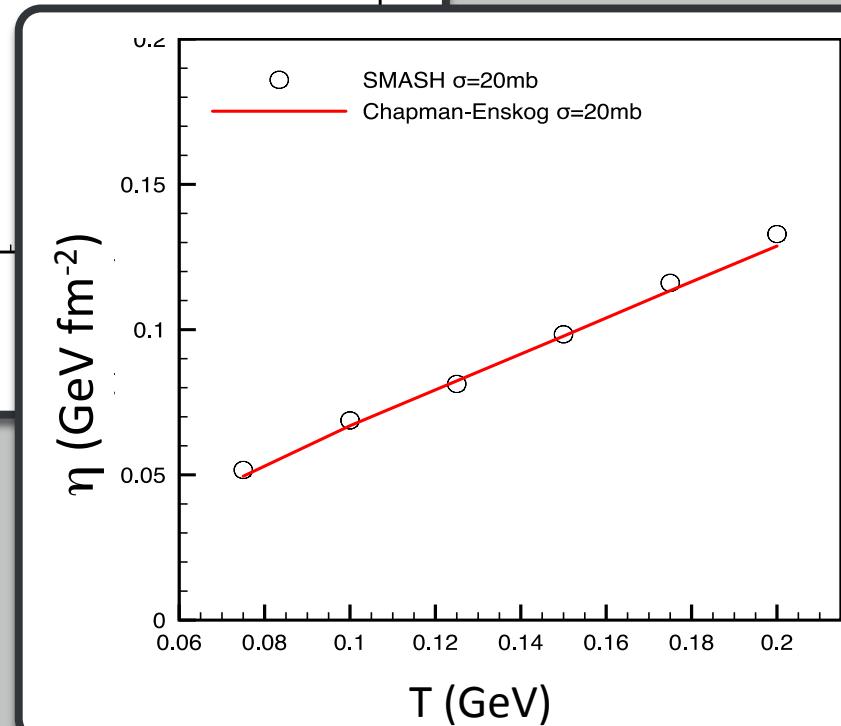
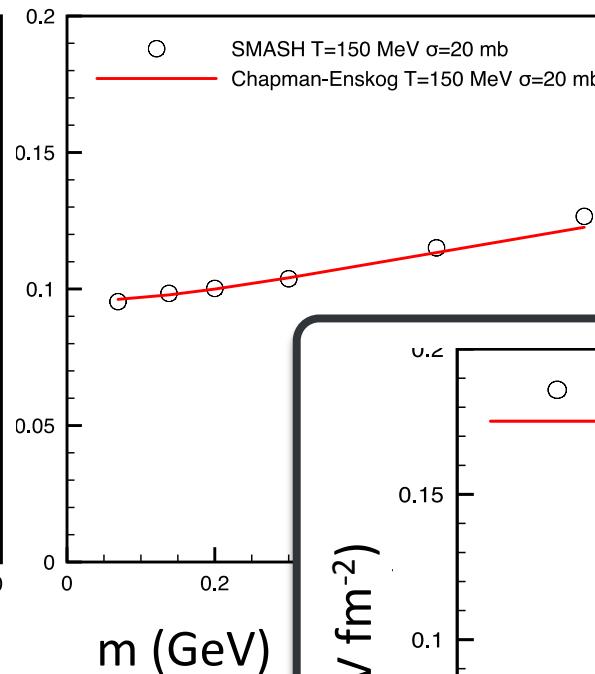
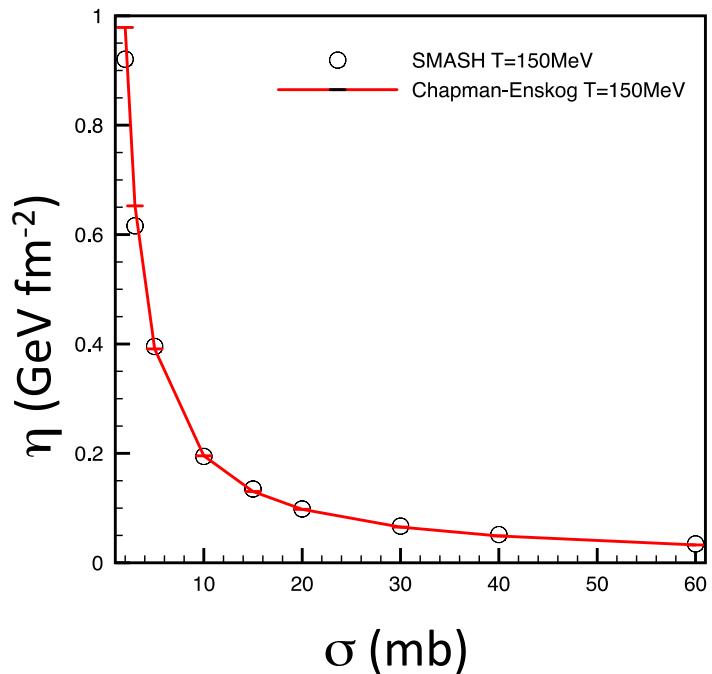


# Equilibrium in SMASH

- Box calculations simulating infinite matter to apply the Green-Kubo procedure
- MUST have thermal & chemical equilibrium
- Baryon/antibaryon annihilation implemented to conserve detailed balance via an average decay to  $5\pi$



# Test case #1: $\pi$ with constant $\sigma$

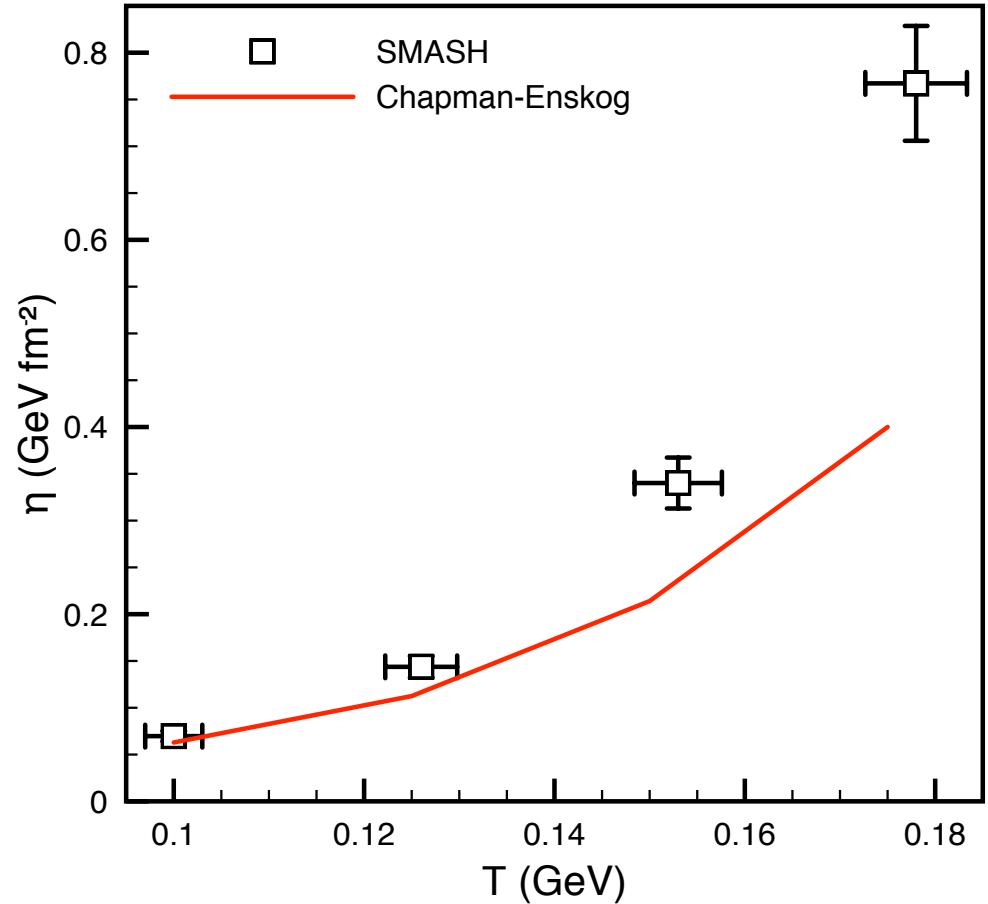


## Main take-away:

The method is relatively insensitive  
to variations of parameters;  
maximum error is less than 10%

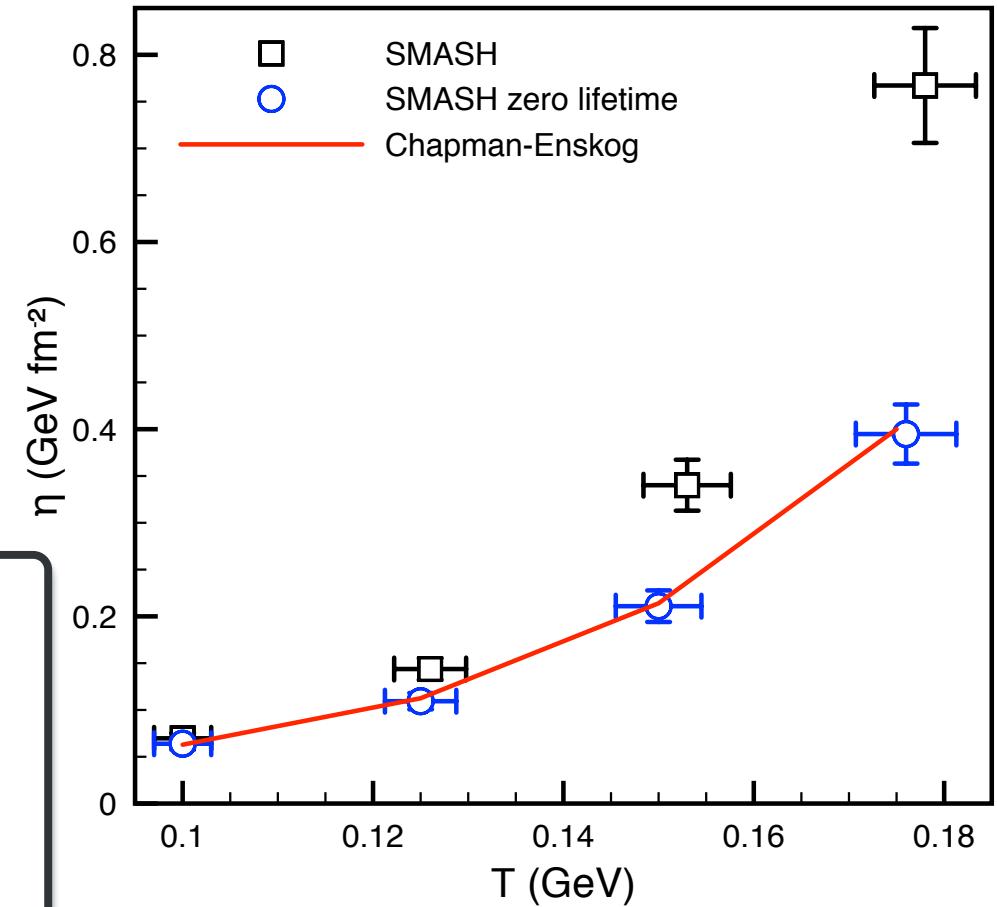
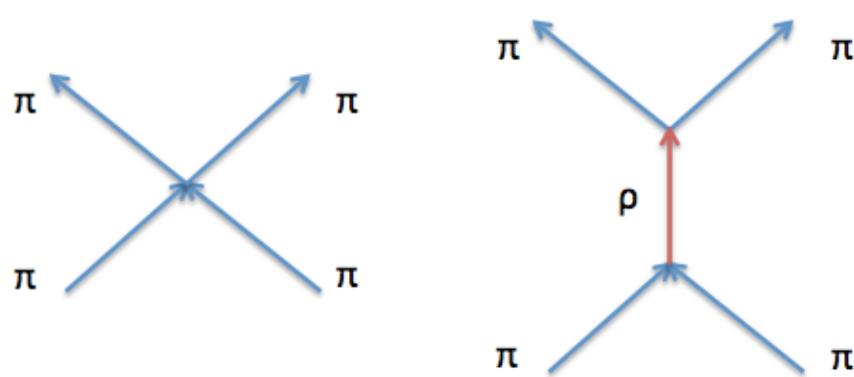
# Test case #2: $\pi$ - $\rho$ gas

- Normal SMASH run does not coincide directly with Chapman-Enskog



# Test case #2: $\pi$ - $\rho$ gas

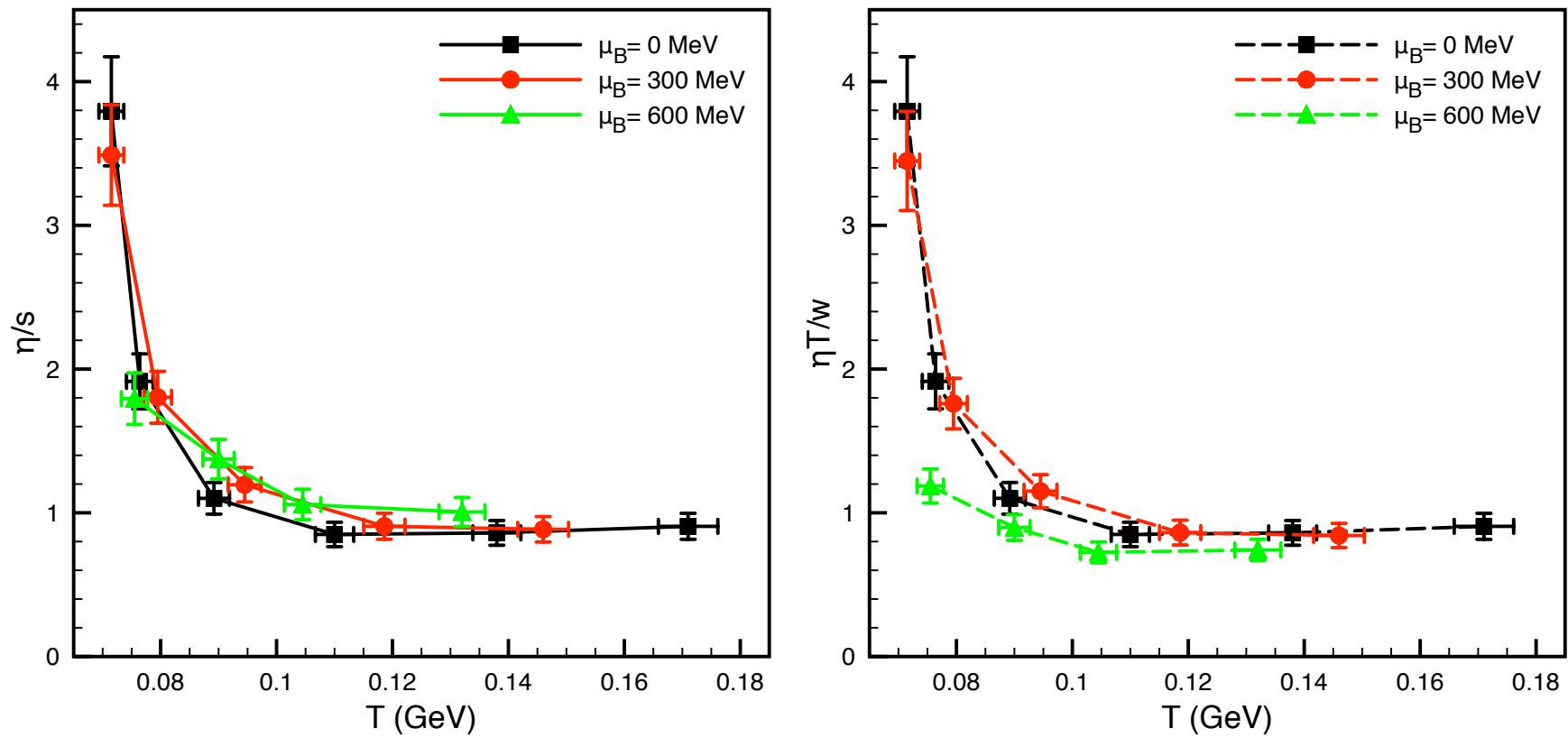
- Normal SMASH run does not coincide directly with Chapman-Enskog
  - Resonance lifetimes



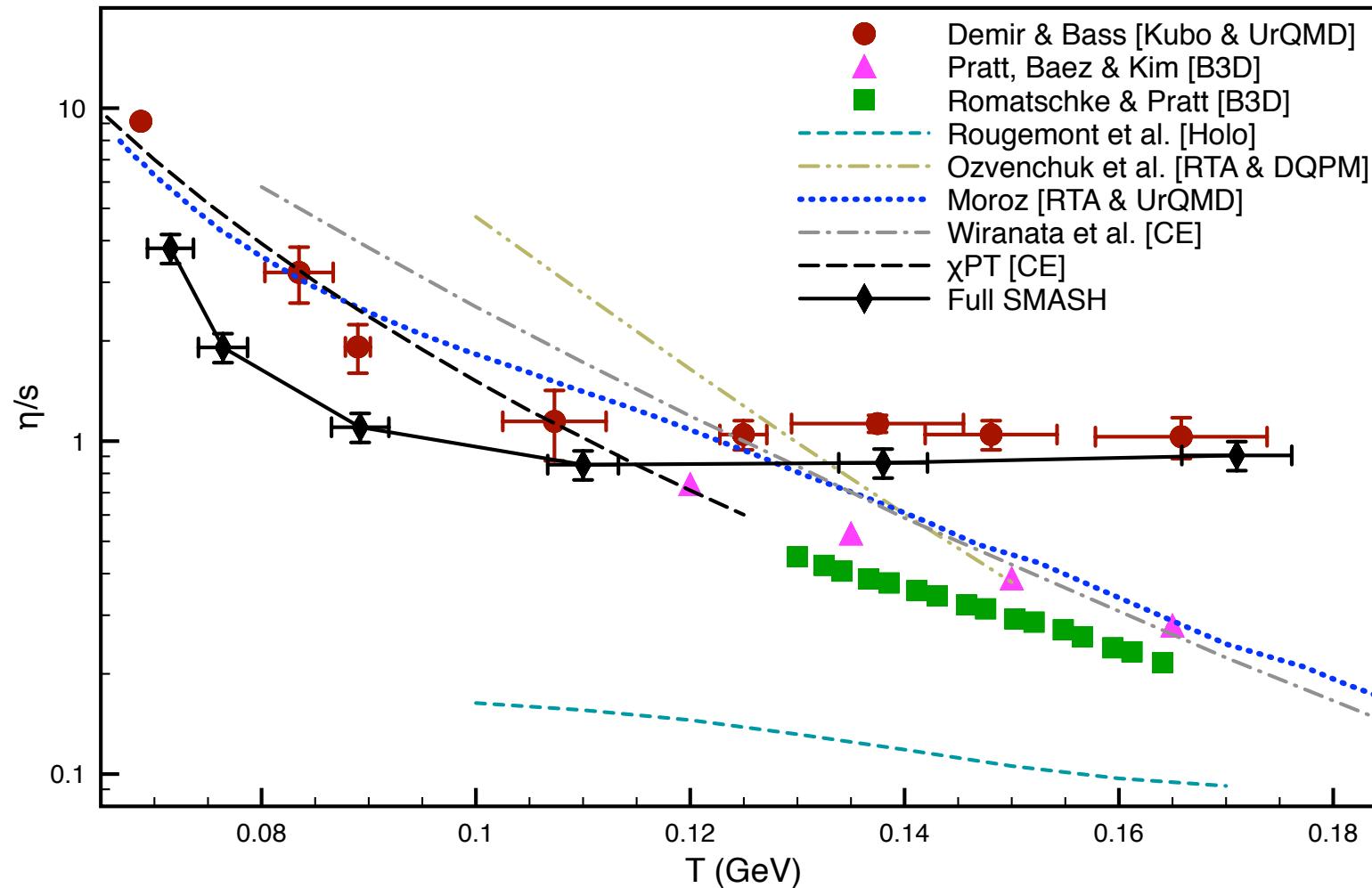
# Hadron Gas: Degrees of freedom

N	$\Delta$	$\Lambda$	$\Sigma$	$\Xi$	$\Omega$	Unflavored					Strange
$N_{938}$	$\Delta_{1232}$	$\Lambda_{1116}$	$\Sigma_{1189}$	$\Xi_{1321}$	$\Omega^{-}_{1672}$	$\pi_{138}$	$f_0\ 980$	$f_2\ 1275$	$\pi_2\ 1670$	$K_{494}$	
$N_{1440}$	$\Delta_{1620}$	$\Lambda_{1405}$	$\Sigma_{1385}$	$\Xi_{1530}$	$\Omega^{-}_{2250}$	$\pi_{1300}$	$f_0\ 1370$	$f_2'\ 1525$		$K^*_{892}$	
$N_{1520}$	$\Delta_{1700}$	$\Lambda_{1520}$	$\Sigma_{1660}$	$\Xi_{1690}$		$\pi_{1800}$	$f_0\ 1500$	$f_2\ 1950$	$\rho_3\ 1690$	$K_1\ 1270$	
$N_{1535}$	$\Delta_{1905}$	$\Lambda_{1600}$	$\Sigma_{1670}$	$\Xi_{1820}$			$f_0\ 1710$	$f_2\ 2010$		$K_1\ 1400$	
$N_{1650}$	$\Delta_{1910}$	$\Lambda_{1670}$	$\Sigma_{1750}$	$\Xi_{1950}$		$\eta_{548}$		$f_2\ 2300$	$\phi_3\ 1850$	$K^*_{1410}$	
$N_{1675}$	$\Delta_{1920}$	$\Lambda_{1690}$	$\Sigma_{1775}$	$\Xi_{2030}$		$\eta'_{958}$	$a_0\ 980$	$f_2\ 2340$		$K_0^*\ 1430$	
$N_{1680}$	$\Delta_{1930}$	$\Lambda_{1800}$	$\Sigma_{1915}$			$\eta_{1295}$	$a_0\ 1450$		$a_4\ 2040$	$K_2^*\ 1430$	
$N_{1700}$	$\Delta_{1950}$	$\Lambda_{1810}$	$\Sigma_{1940}$			$\eta_{1405}$		$f_1\ 1285$		$K^*_{1680}$	
$N_{1710}$		$\Lambda_{1820}$	$\Sigma_{2030}$			$\eta_{1475}$	$\phi_{1019}$	$f_1\ 1420$	$f_4\ 2050$	$K_2\ 1770$	
$N_{1720}$		$\Lambda_{1830}$	$\Sigma_{2250}$				$\phi_{1680}$			$K_3^*\ 1780$	
$N_{1875}$		$\Lambda_{1890}$				$\sigma_{800}$		$a_2\ 1320$		$K_2\ 1820$	
$N_{1900}$		$\Lambda_{2100}$					$h_{1\ 1170}$			$K_4^*\ 2045$	
$N_{1990}$		$\Lambda_{2110}$				$\rho_{776}$		$\pi_1\ 1400$			
$N_{2080}$		$\Lambda_{2350}$				$\rho_{1450}$	$b_{1\ 1235}$	$\pi_1\ 1600$			
$N_{2190}$						$\rho_{1700}$		$a_{1\ 1260}$	$\eta_2\ 1645$		
$N_{2220}$						$\omega_{783}$					
$N_{2250}$						$\omega_{1420}$			$\omega_3\ 1670$		
						$\omega_{1650}$					
<ul style="list-style-type: none"> <li>• Isospin symmetry</li> <li>• Perturbative treatment of non-hadronic particles (photons, dileptons)</li> </ul>											

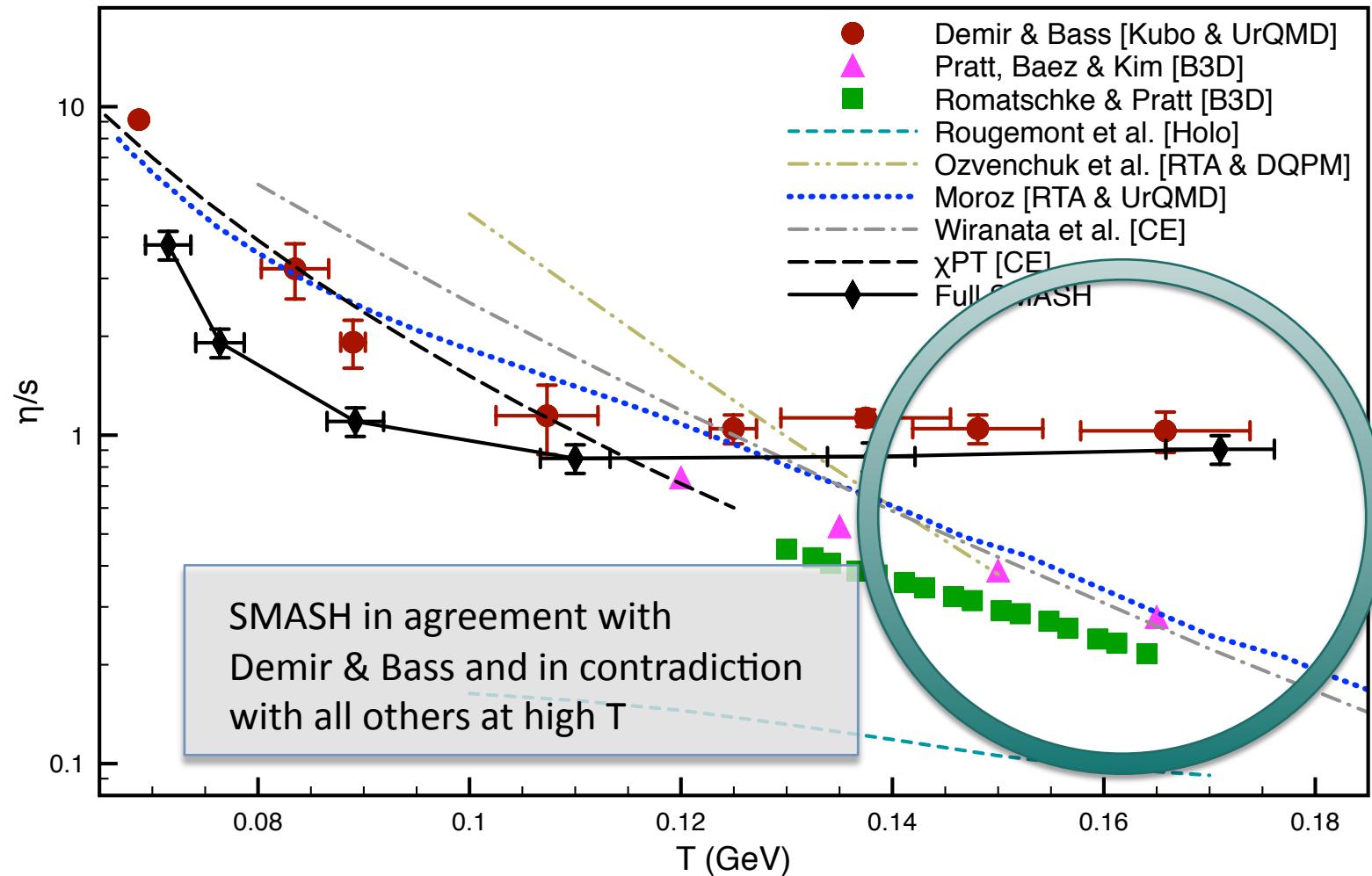
# Hadron Gas: T and $\mu_B$ dependence



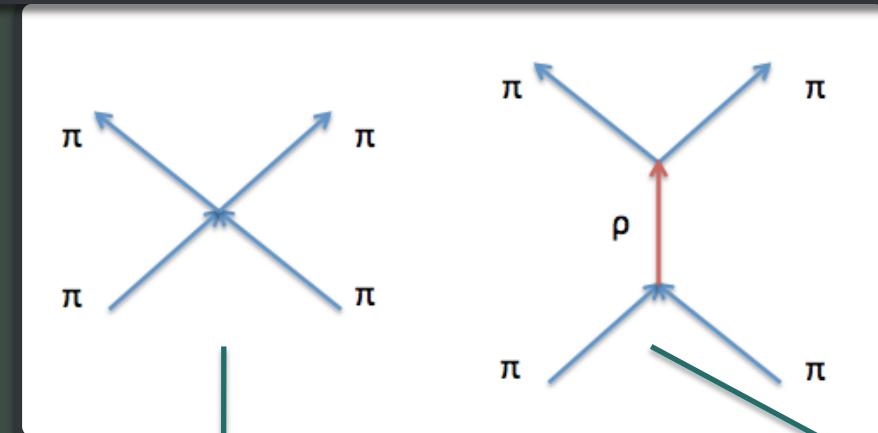
# HG: Viscosity Comparison



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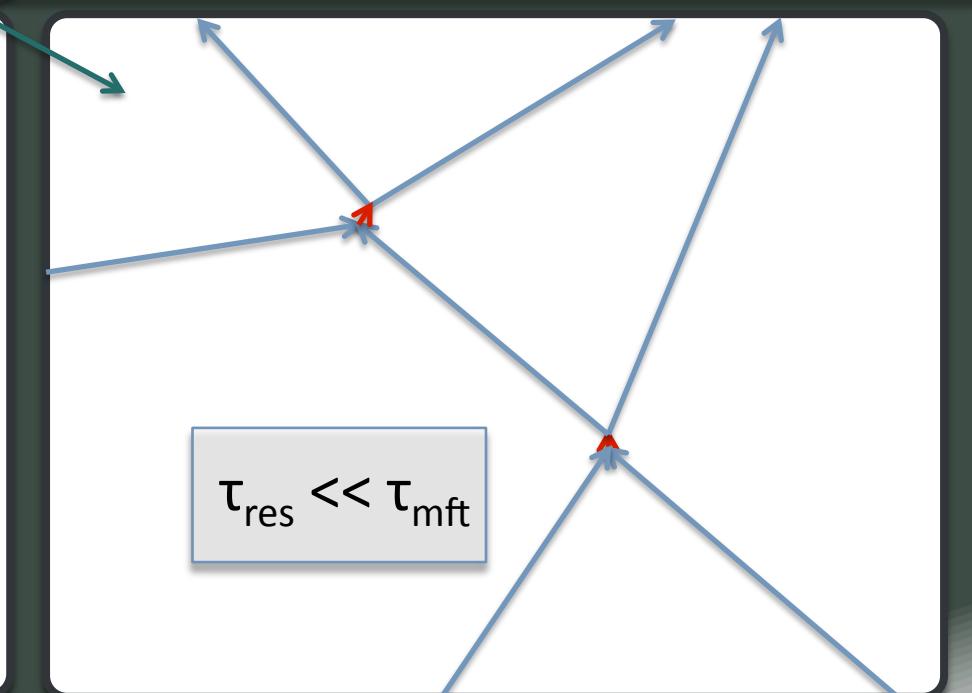
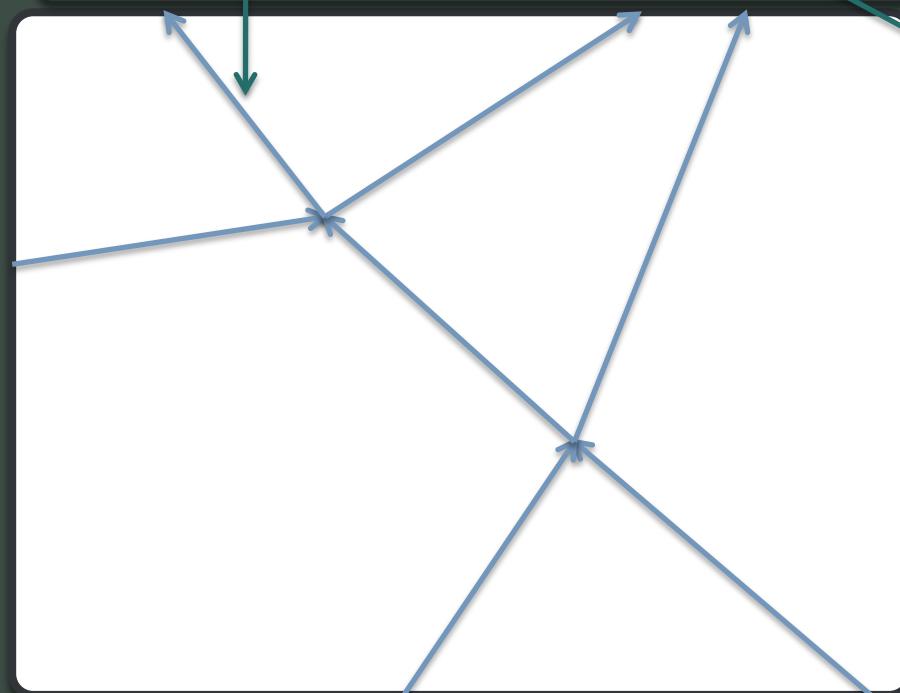


# High temperature $\eta/s$ : Resonance lifetimes

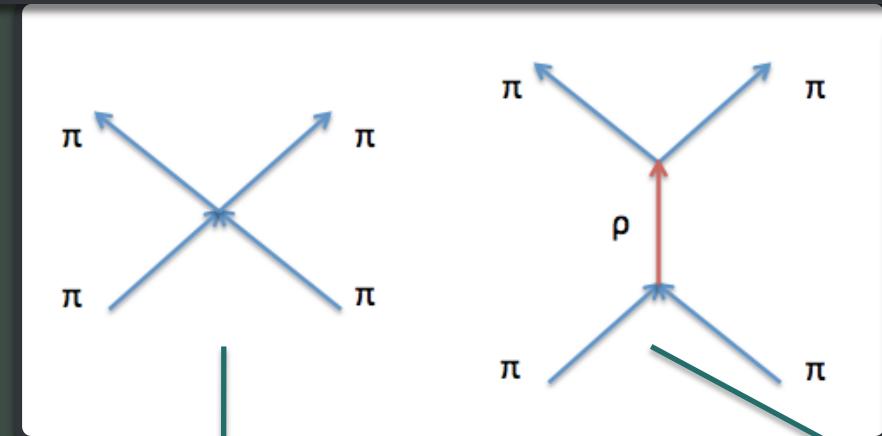


Must look at the microscopic picture from different descriptions  
 $\tau_{\text{res}}$  = resonance lifetime  
 $\tau_{\text{mft}}$  = mean free time

At low T and density:

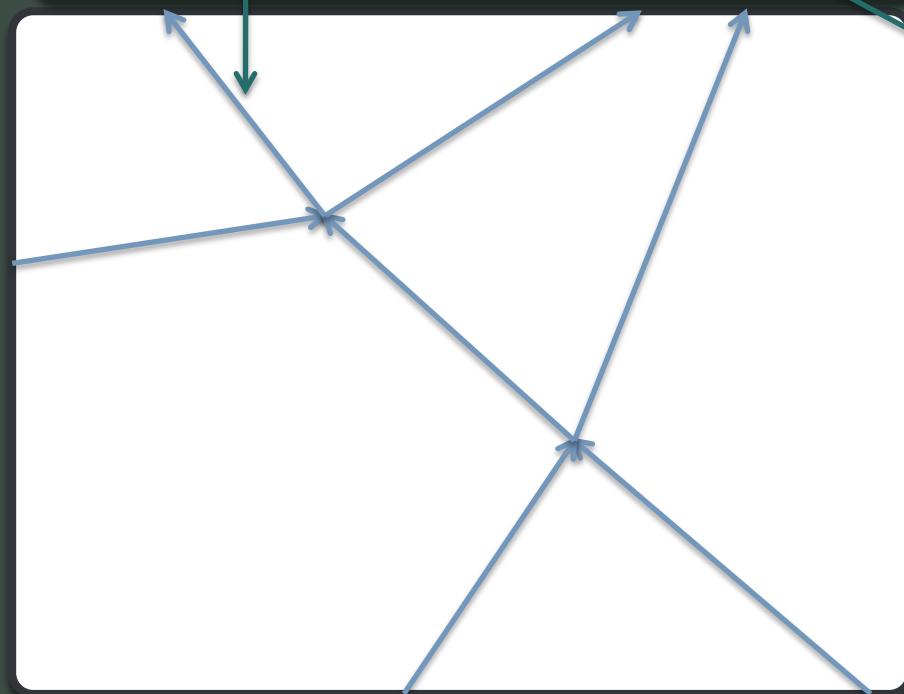


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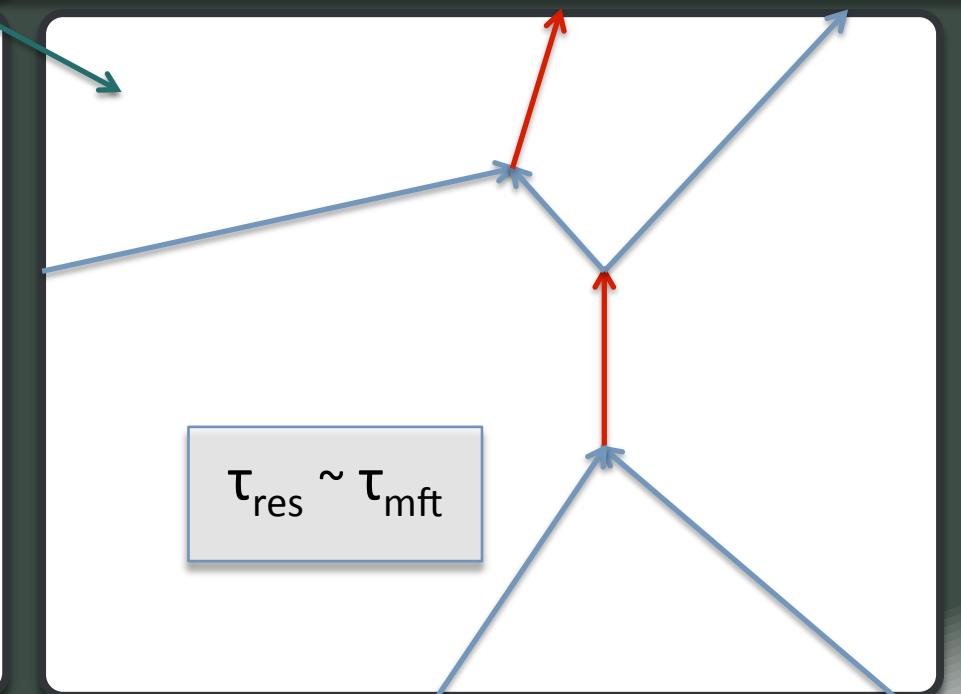


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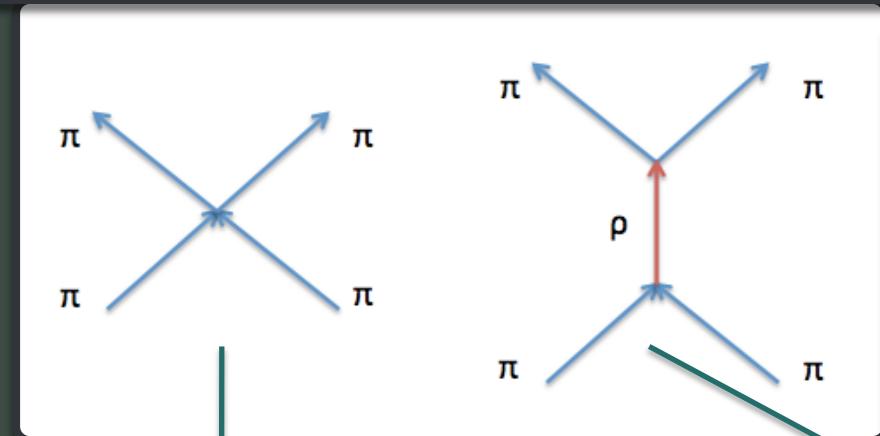
At high T and density:



$$\tau_{\text{res}} \sim \tau_{\text{mft}}$$



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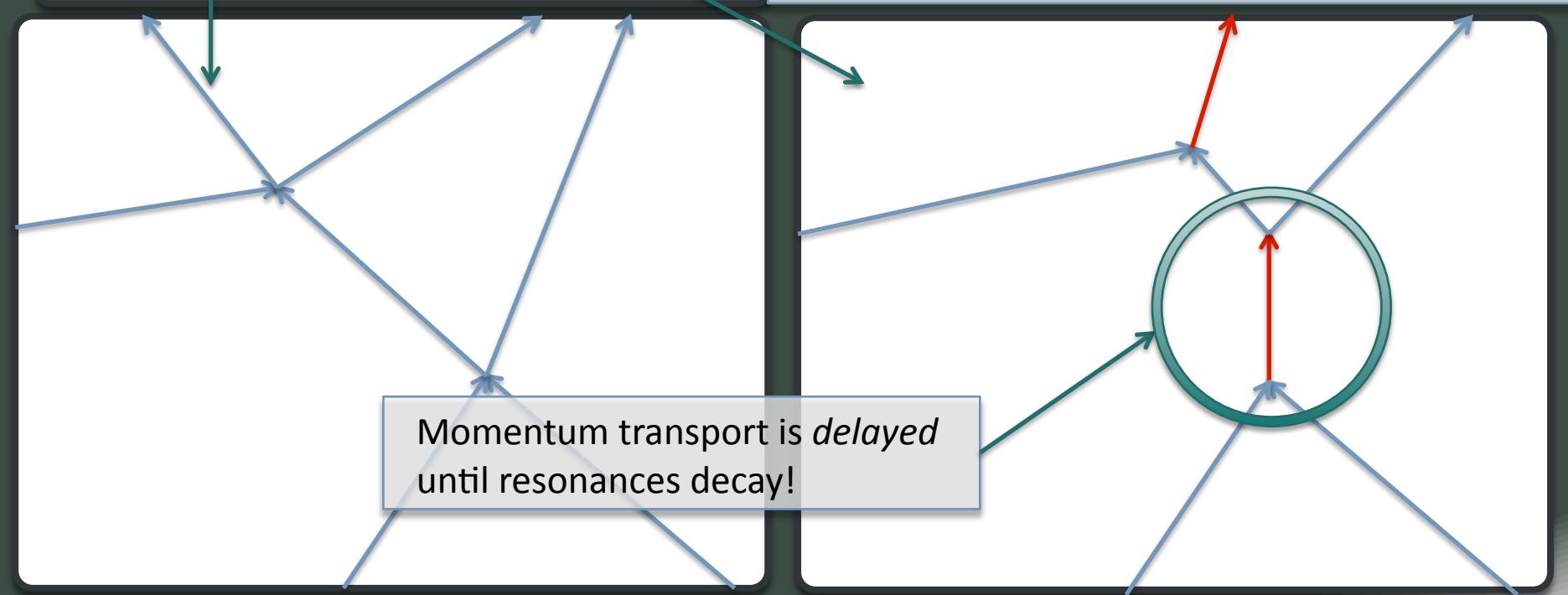


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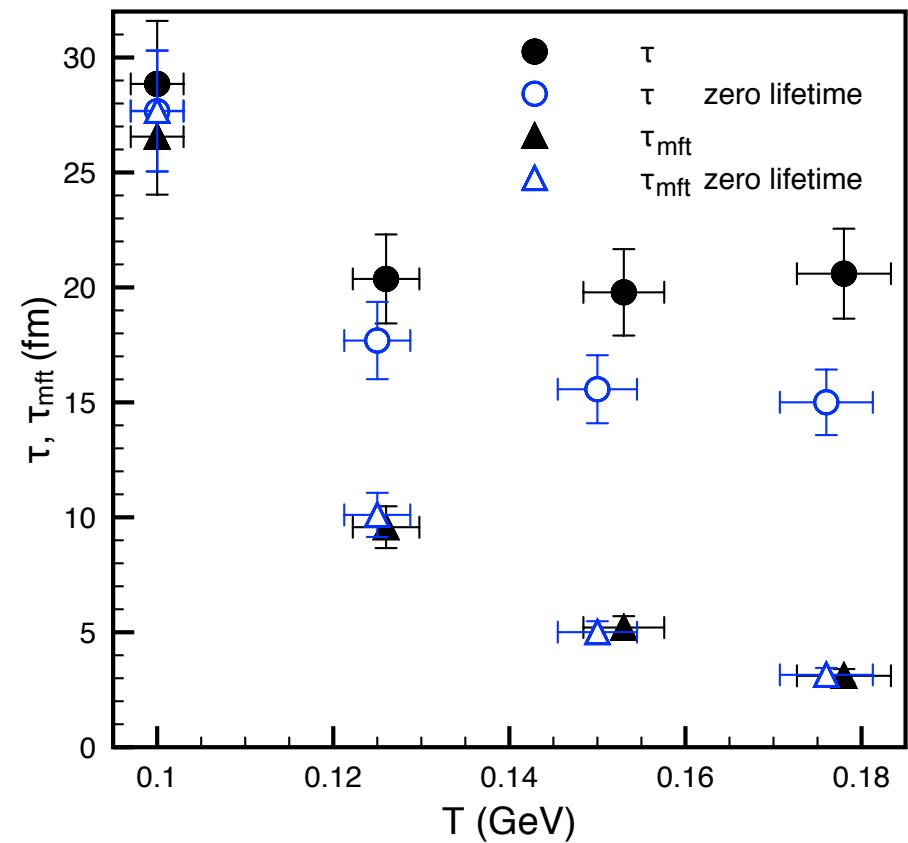
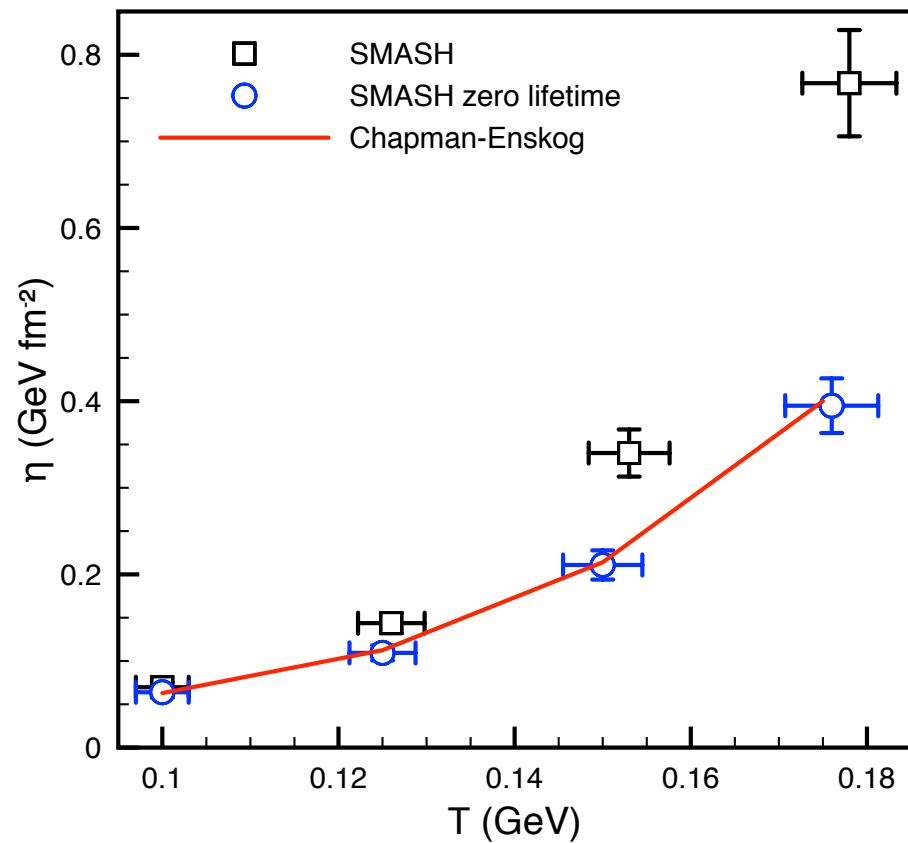
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At high T and density:



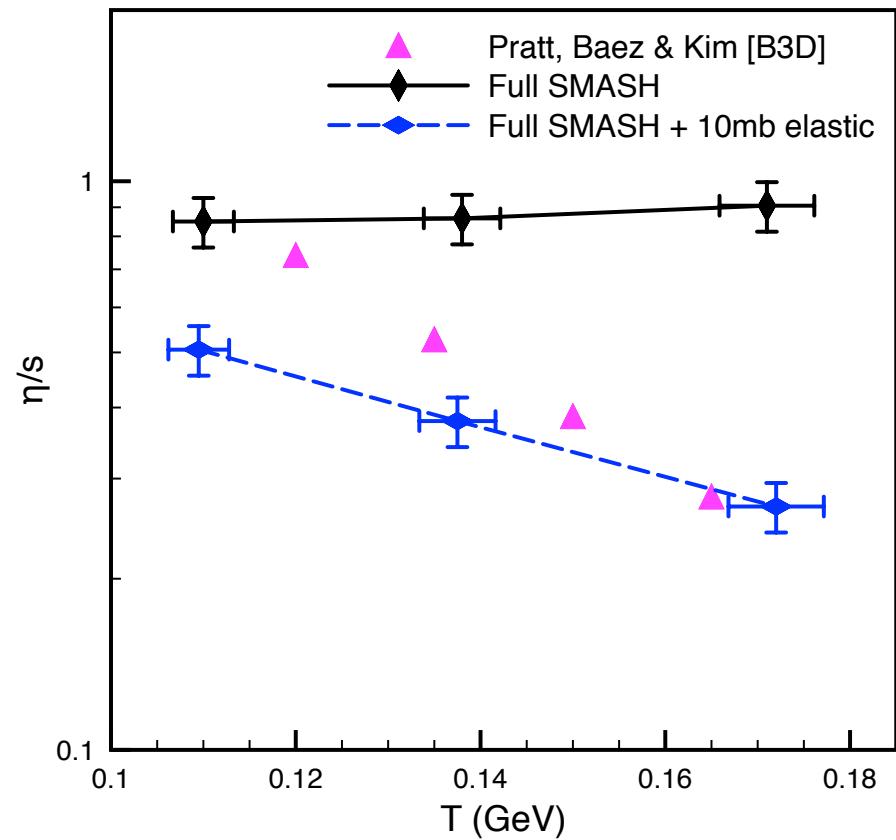
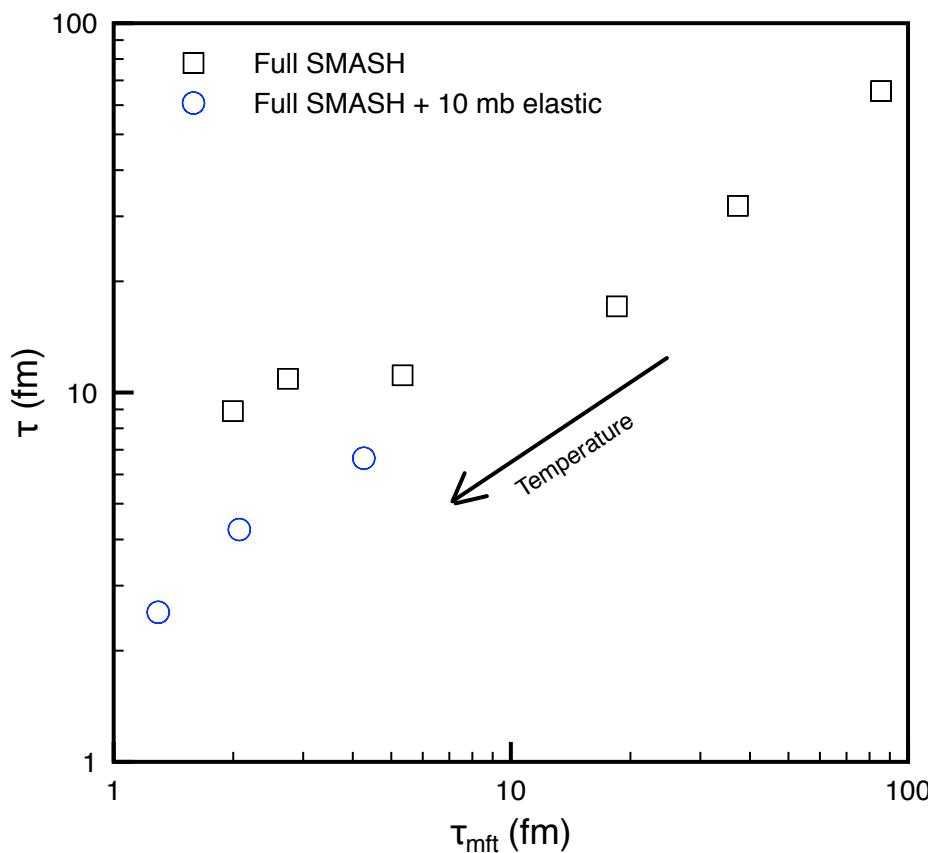
# $\pi$ - $\rho$ : Zero lifetimes vs relaxation time

Large part of the difference explained from eliminating lifetimes



# Effect of many non-resonant interactions

Introduce a constant elastic cross-section between all particles to add many non-resonant interactions

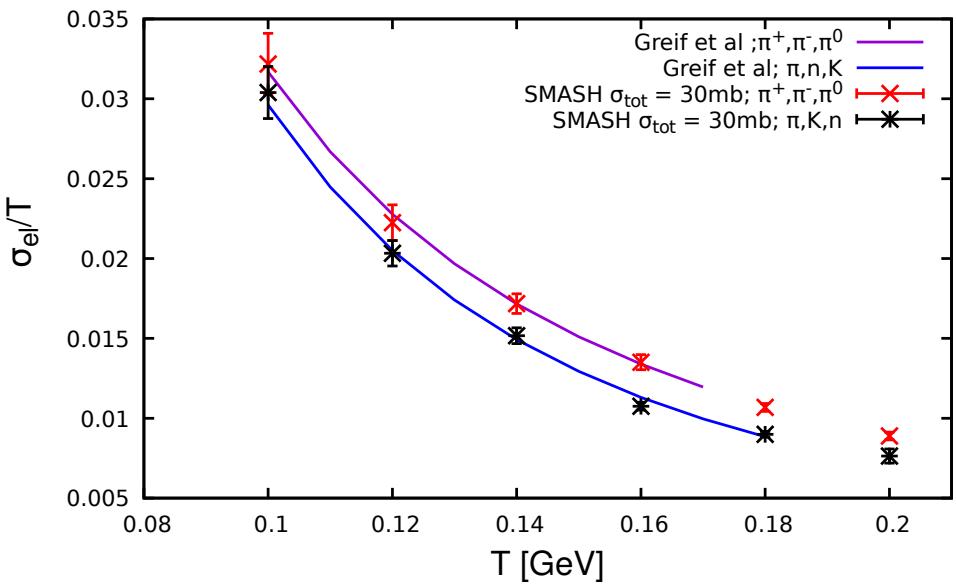


# Summary & Outlook

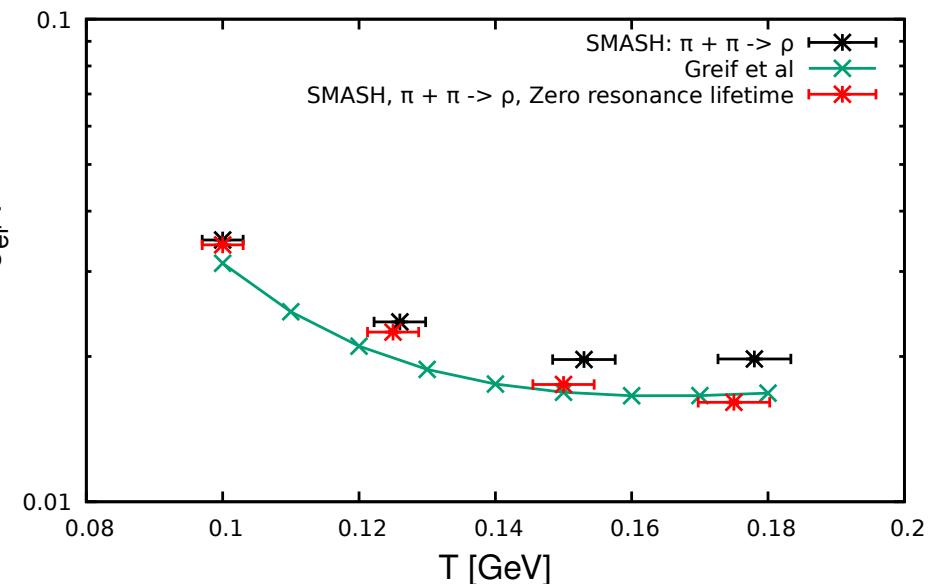
- **Investigated temperature, cross-section and mass dependence of the shear viscosity in an elastic pion box**
  - Very good agreement with Chapman-Enskog approximation (within 10%)
  - Resonance lifetimes need to be considered
- **Full hadron gas  $\eta/s$  calculated at zero and non-zero  $\mu_B$** 
  - High T discrepancy explained by looking at microscopic details of resonance modelling; finite lifetime increases viscosity
  - Could be used to constrain the treatment of resonances
- **Outlook:**
  - Investigation of angular dependent interactions on viscosity
  - At temperatures close to the phase transition, inclusion of multi-particle interaction will play a role, and needs to be investigated
  - Other transport coefficients (electrical conductivity, bulk viscosity, etc.)

# Outlook: Electric conductivity

See poster by J. Hammelmann later today in the poster session (COL-15)



Constant  $\sigma$ : good agreement  
with analytical calculation



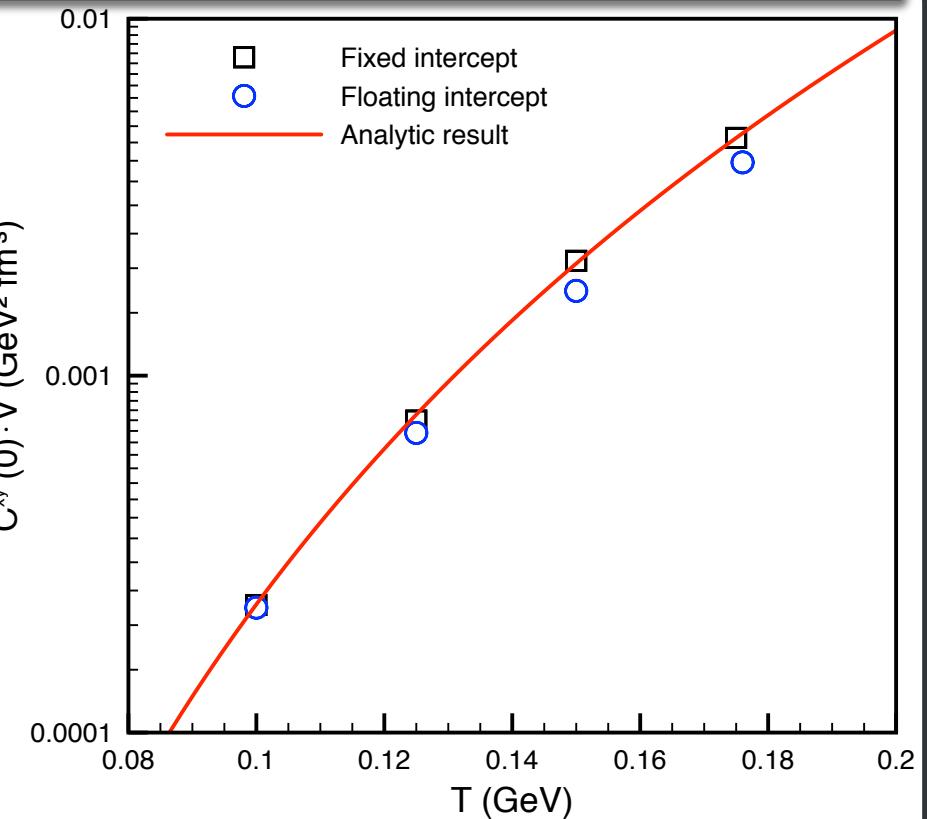
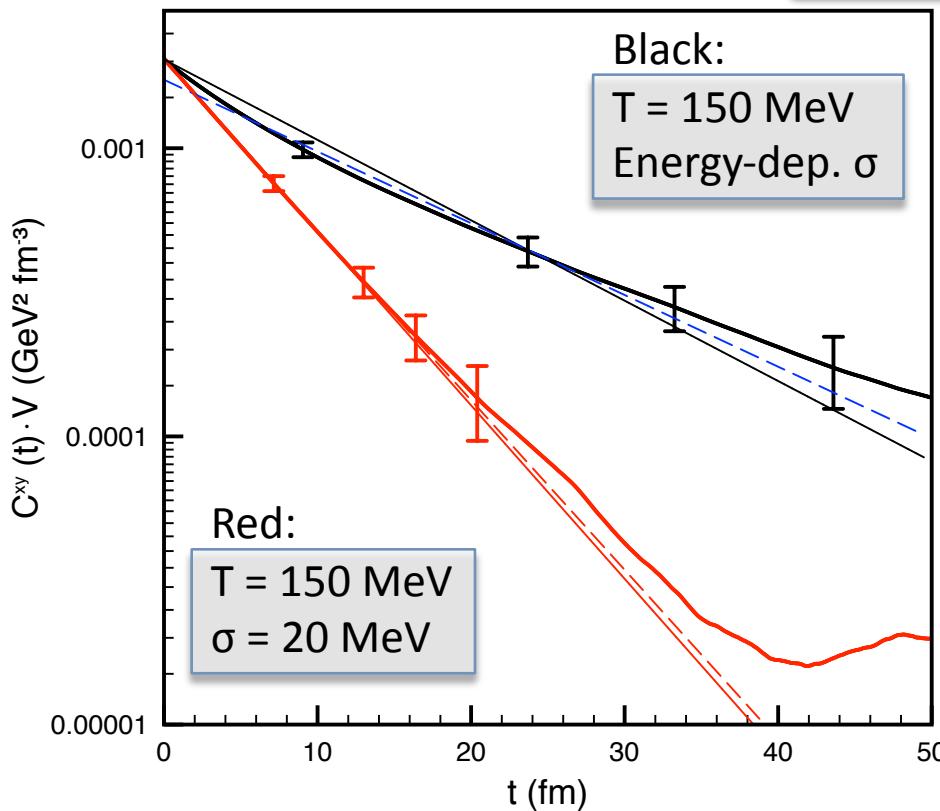
Resonant case also shows sensitivity  
to resonance lifetimes!

M. Greif, C. Greiner, G. S. Denicol, Phys. Rev D93 (2016) 096012, arXiv:1602.05085

# Backup slides

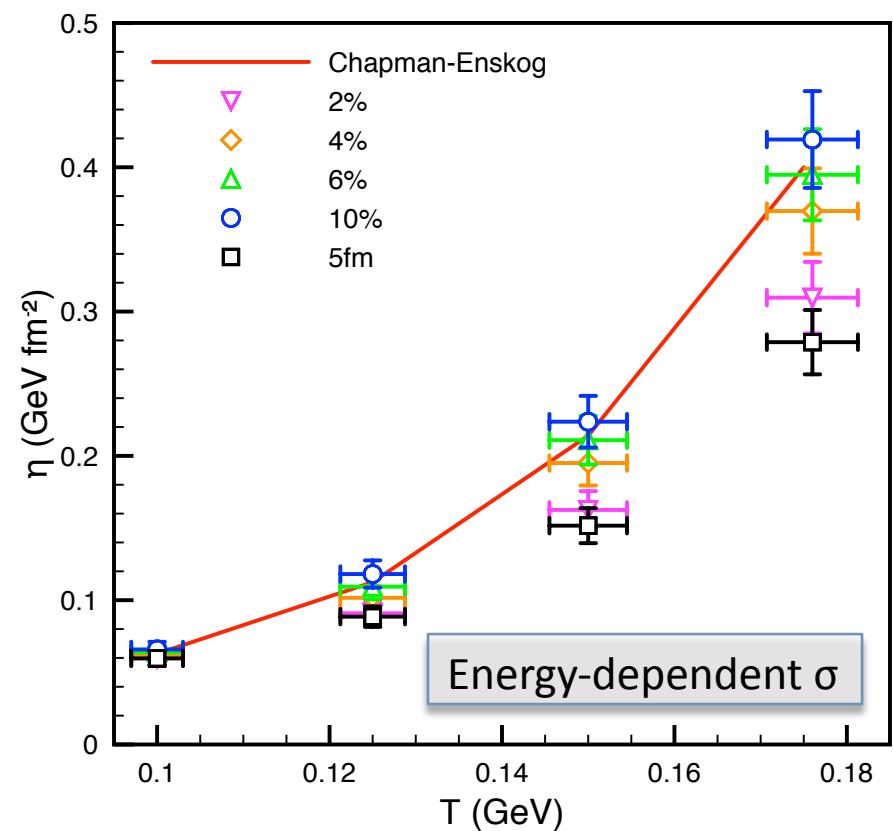
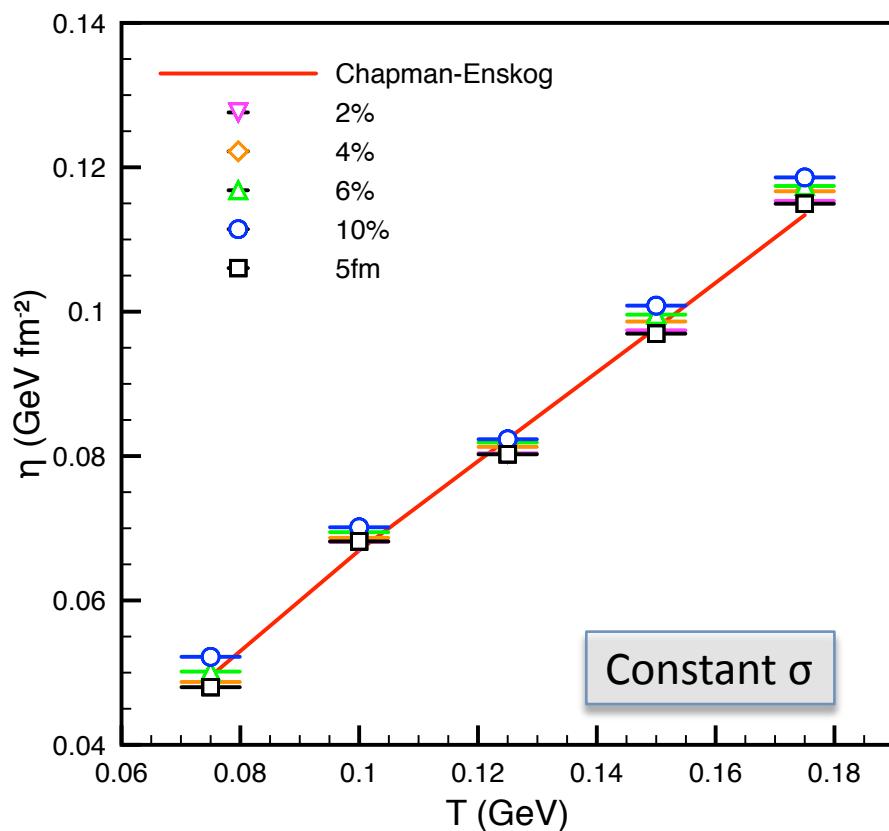
# How to fit?

$$C^{xy}(0) = \frac{g \exp\left(\frac{\mu}{T}\right)}{30\pi^2 V} \int_0^\infty dp \frac{p^6}{m^2 + p^2} \exp\left(-\frac{\sqrt{m^2 + p^2}}{T}\right)$$

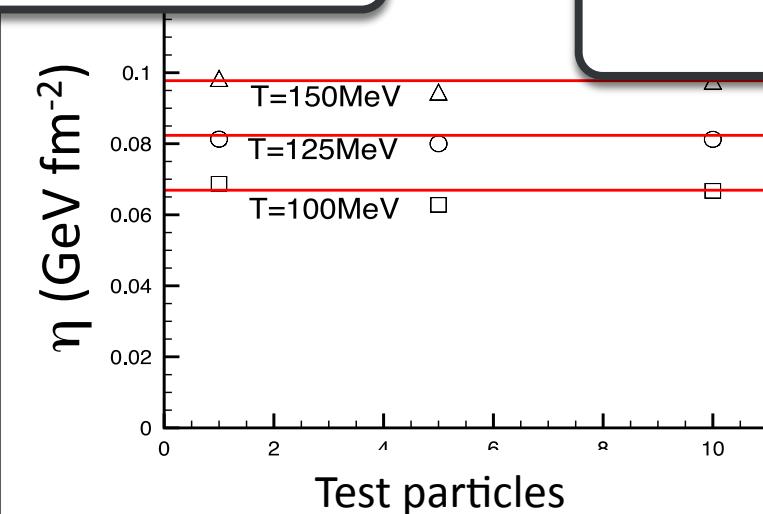
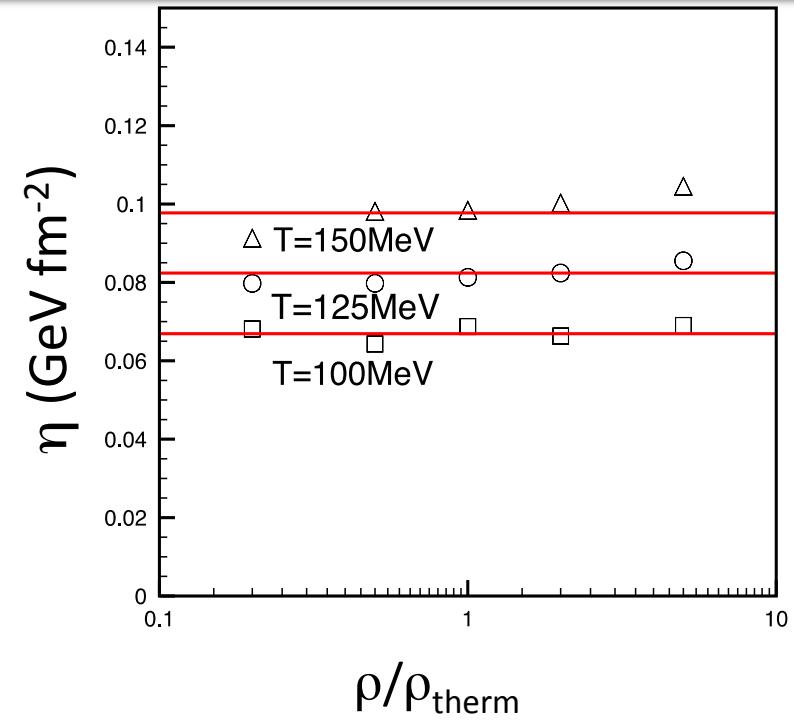
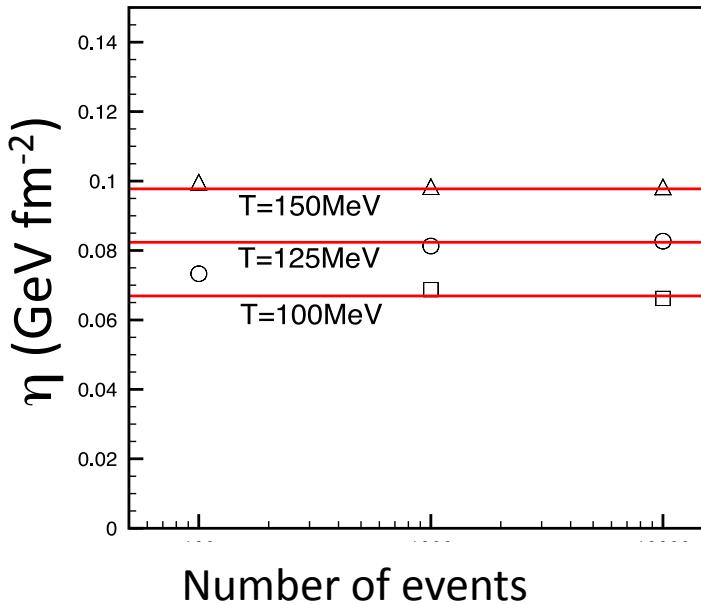


# Where to stop the fitting?

J. Torres-Rincon, PhD dissertation (2012), *Hadronic Transport Coefficients from Effective Field Theories*



# Constant $\sigma$ : Systematics

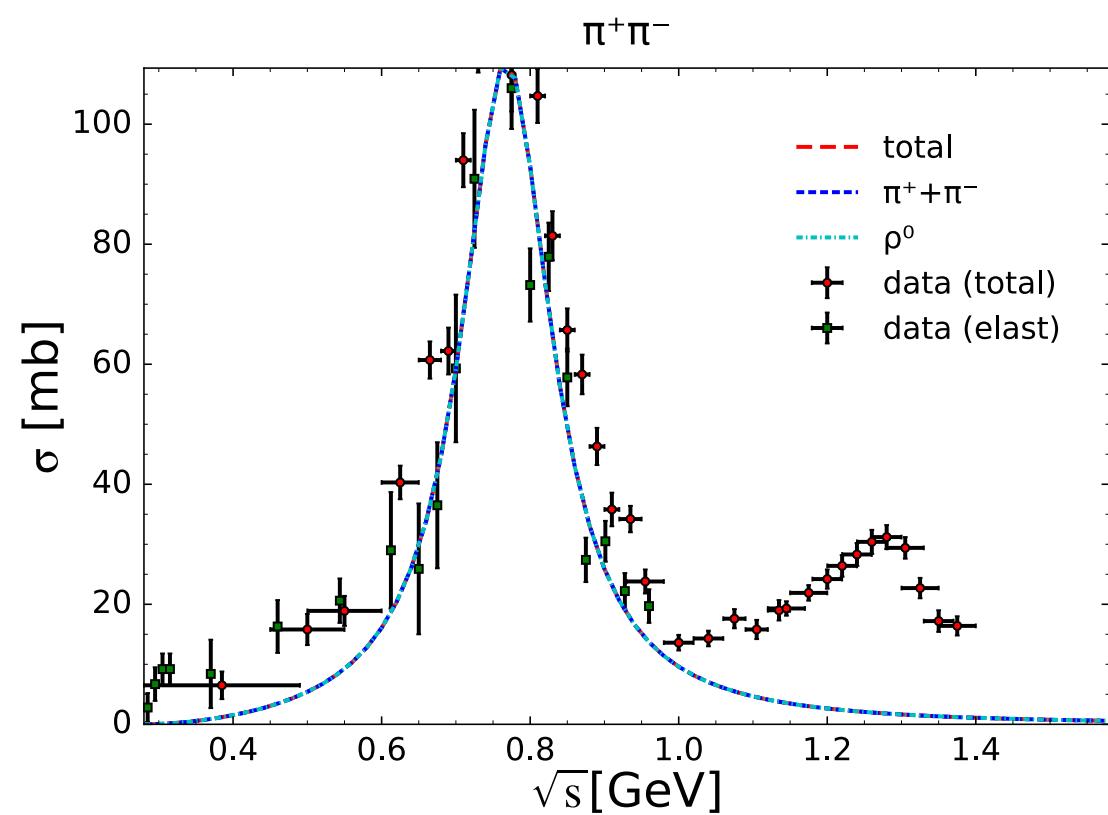


**Main take-away:**

The method is relatively inelastic to variations of most parameters; maximum error is less than 10%

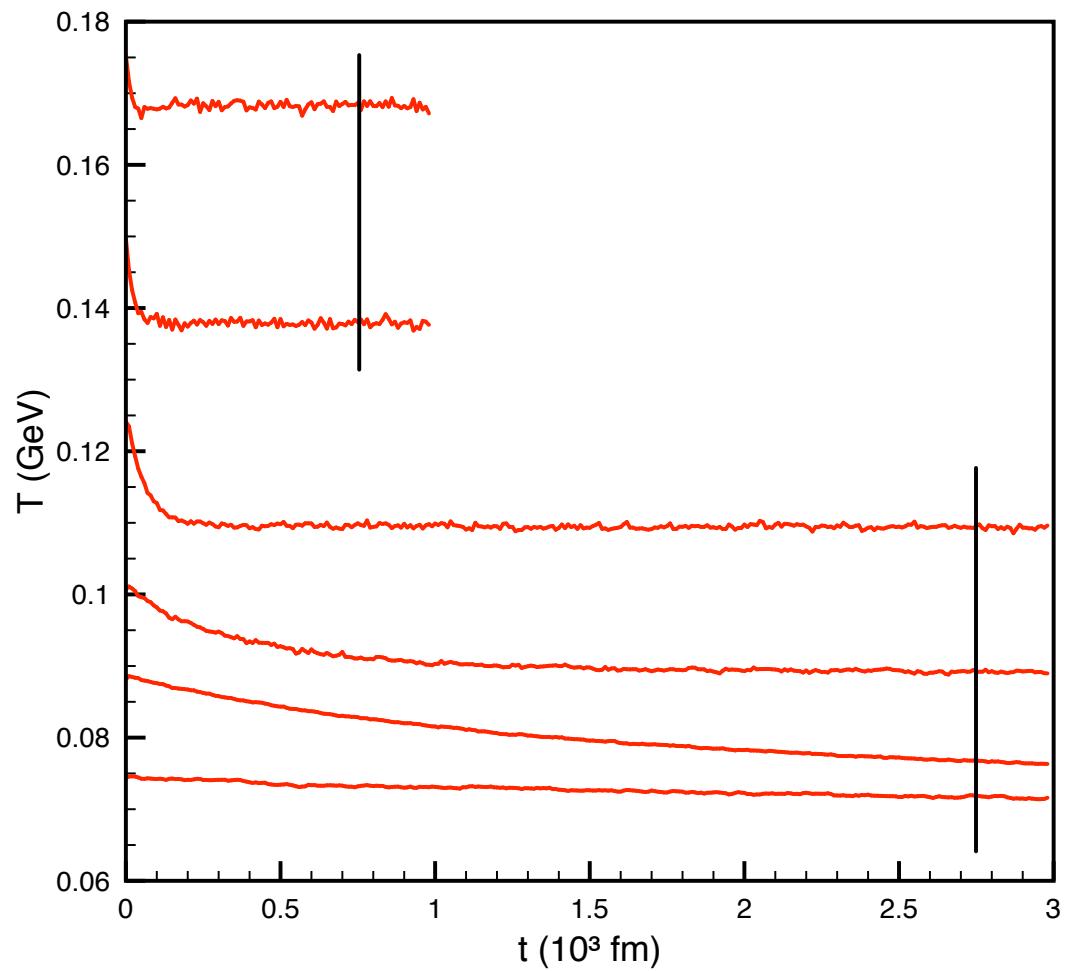
# Test case #2: Energy-dependent $\sigma$

- Pions in a  $(20 \text{ fm})^3$  box simulating infinite matter
- Cross-section uses  $\rho$  resonance
- Runs for  $t_{max} = 200 \text{ fm}$
- Initialized with initial densities consistent with Boltzmann ideal gas



# Hadron Gas

- All particles and resonances initialized to thermal multiplicities (at the pole mass)
- Must wait for equilibration and compute  $T, \mu$  once in equilibrium from most abundant particles
  - $T$  fitted from weighted momentum spectra of  $\pi, K$  &  $N$
  - $\mu_B$  obtained from  $N$  / anti- $N$  ratio



# What about entropy?

The entropy density can be calculated from the Gibbs formula:

$$S = \frac{e + p - \mu n}{T} = \frac{w - \mu n}{T}$$

where the energy density and pressure can be taken from the average shear-stress tensor according to:

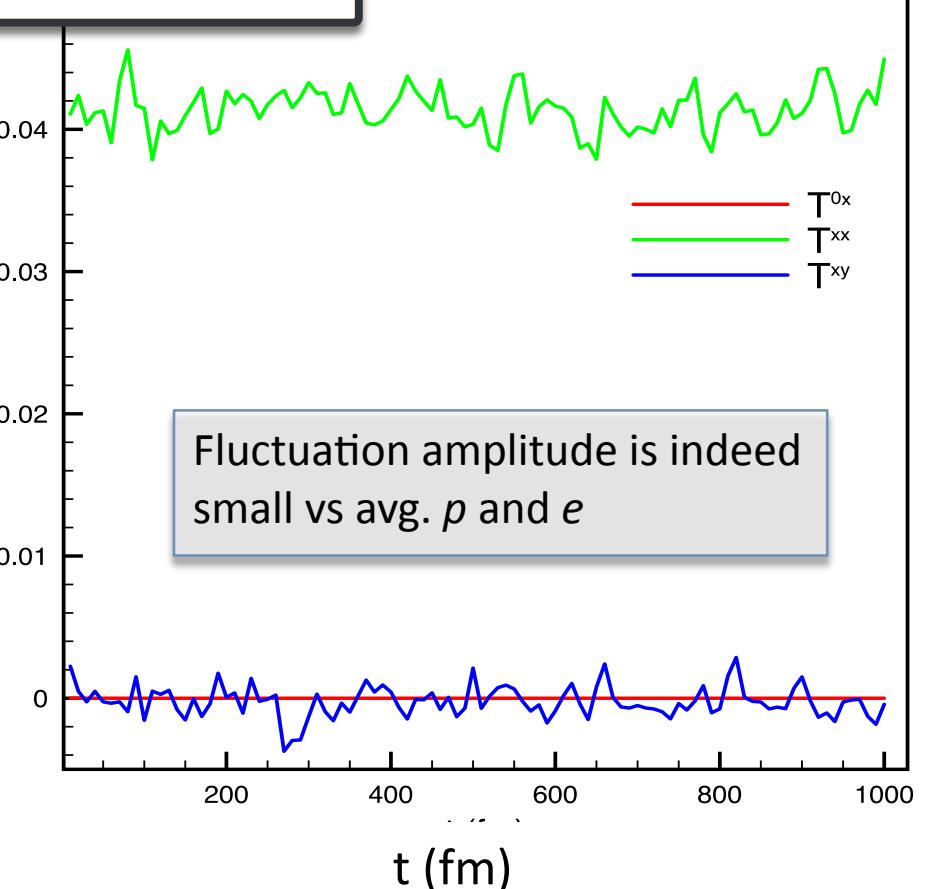
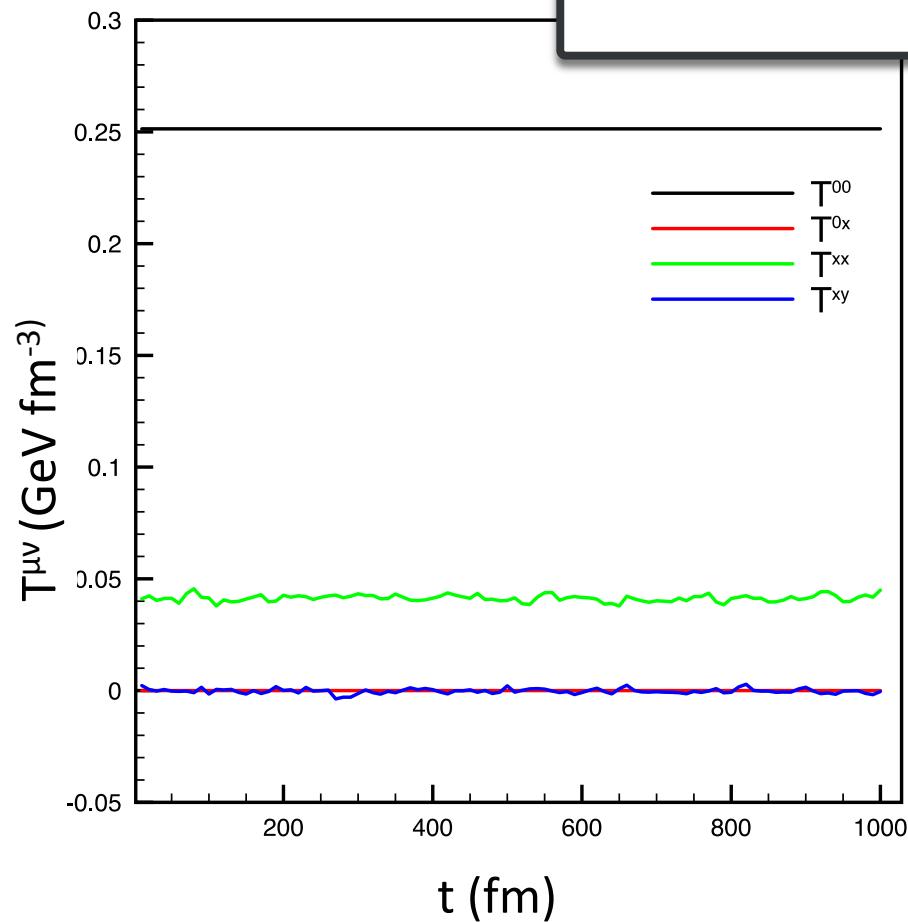
$$T^{\mu\nu} = \text{diag}(e, p, p, p)$$

Assuming a nearly ideal gas, one can fit the temperature and chemical potential with momentum distributions:

$$\frac{dN}{dp} = \frac{g}{2\pi^2} V p^2 \exp\left(-\frac{\sqrt{p^2 + m^2} - \mu}{T}\right)$$

# Energy density and pressure

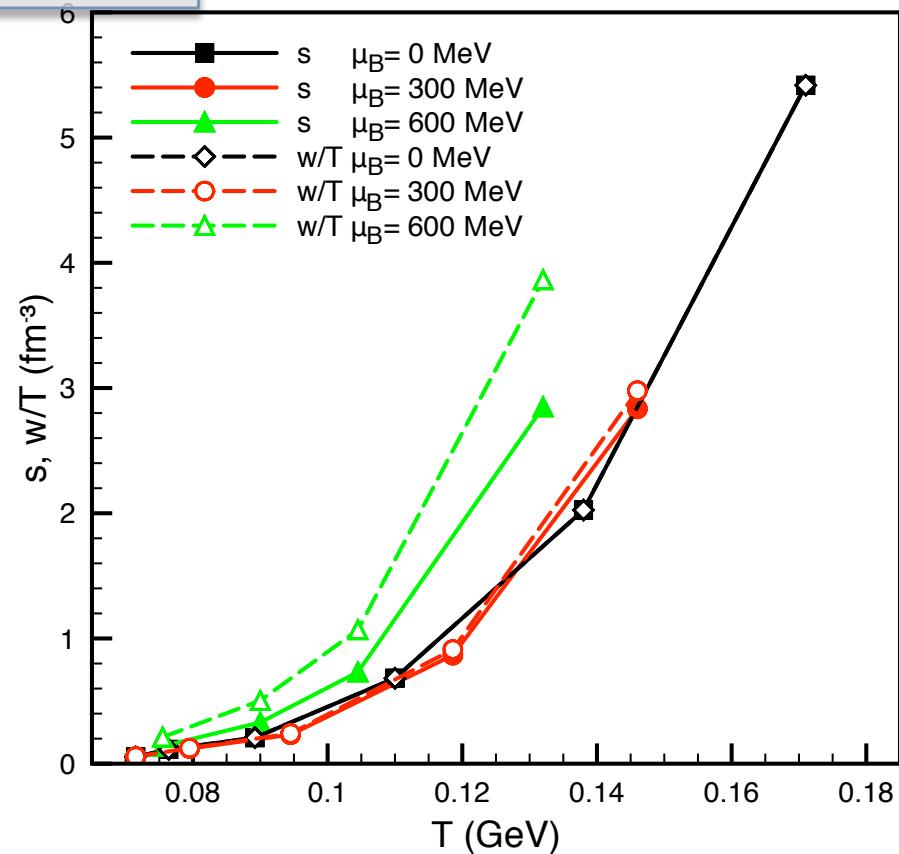
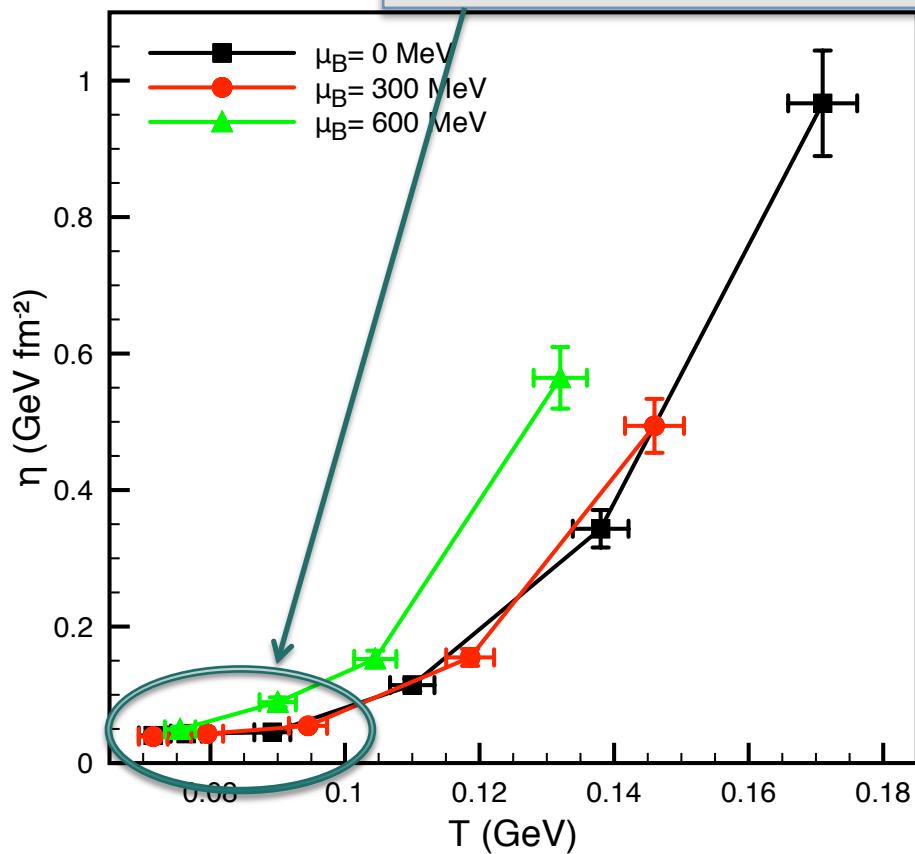
$$T^{\mu\nu} = \text{diag}(e, p, p, p)$$



Fluctuation amplitude is indeed  
small vs avg.  $p$  and  $e$

# Hadron Gas: $\eta$ , $s$ and $w/T$

Viscosity decreases slower at small temperatures; explains rise of  $\eta/s$



# Hadron Gas: Low temperature $\eta/s$

- Low temperature hadron gas is composed almost exclusively of pions
- $\pi\text{-}\pi$  cross-section is then most relevant
  - At very low energy, SMASH much higher than UrQMD/ $\chi$ PT
  - $\chi$ PT includes angular dependence, UrQMD&SMASH don't; increases viscosity by factor up to  $5/3$  for  $\rho$  resonance

