

# Effect of the QCD equation of state and strange hadronic resonances on multiparticle correlations in heavy ion collisions ([nucl-th]:1711.05207)

V. Mantovani Sarti<sup>1</sup>, P. Alba<sup>2</sup>, J. Noronha<sup>3</sup>, J. Noronha-Hostler<sup>4</sup>, P. Parotto<sup>5</sup>, I. P. Vazquez<sup>5</sup> and C. Ratti<sup>5</sup>

<sup>1</sup>TUM, <sup>2</sup>FIAS, <sup>3</sup>Universidade de Sao Paulo, <sup>4</sup>Rutgers University, <sup>5</sup>University of Houston

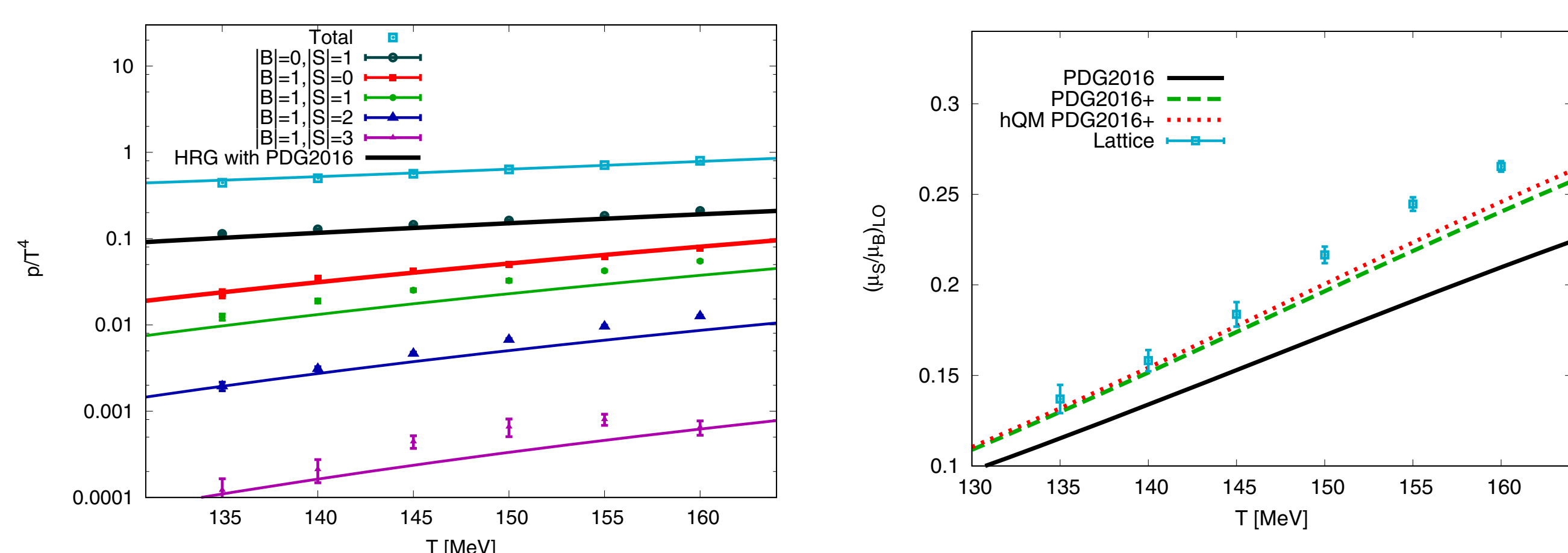
valentina.mantovani-sarti@tum.de



## Why do we need additional strange resonances?

- **Additional resonances are needed in order to reproduce Lattice QCD (LQCD) results on partial pressures  $p/T^4$**  [1, 2];
- in the Hadron Resonance Gas model (HRG) the total pressure is given by the contributions from each hadron family with baryon number B and strangeness content S

$$P(\mu_B, \mu_S) = \sum_{i=\text{reson.}} P_{b_i s_i}^{BS} = P_{00}^{BS} + P_{10}^{BS} \cosh(\mu_B) + P_{01}^{BS} \cosh(-\mu_S) + P_{11}^{BS} \cosh(\mu_B - \mu_S) + P_{12}^{BS} \cosh(\mu_B - 2\mu_S) + P_{13}^{BS} \cosh(\mu_B - 3\mu_S)$$



- comparison of partial pressures  $p/T^4$  from **Wuppertal-Budapest LQCD** collaboration and Hadron Resonance Gas (HRG) model using **PDG2016** with only \*\*\*\*-\*\*\* states [3]  $\Rightarrow$  **need for additional strange states**
- hyperons in the  $S = -1$  (e.g.  $\Lambda$ ) and  $S = -2$  sector (e.g.  $\Xi$ ), are underestimated
- **PDG16** most up-to-date list of all \*\*\*-\*\*\*\* states from the Particle Data Group
- **PDG16+** most up-to-date list of all \*-\*\*\*\* states from the Particle Data Group
- only **strong decays** with non negligible branching ratios ( $\approx 1\%$  or higher)
- if branching ratios do not sum to 100%  $\Rightarrow$  remaining decays  $N_2 \rightarrow N_1 + \gamma$  where  $N_2$  and  $N_1$  are hadrons with the same quantum numbers and  $N_1$  is the next state in descending mass order with parity compatible for such a decay
- if no decay information available  $\Rightarrow$   **$\leq 30\%$  BR hadronic decays,  $\geq 70\%$  radiative decays**

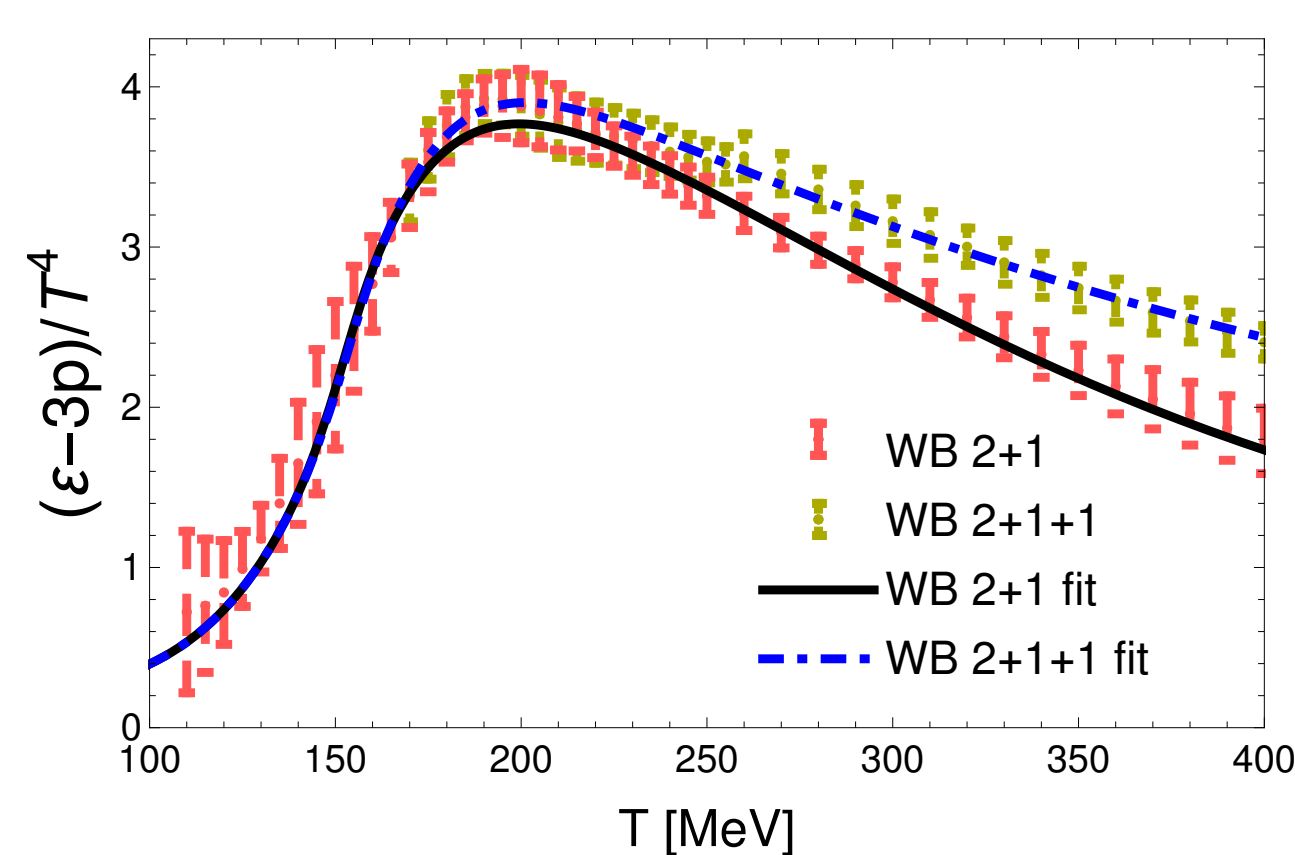
To leading order in the baryochemical potential  $\mu_B$ , the ratio  $\mu_S/\mu_B$  reads

$$\left(\frac{\mu_S}{\mu_B}\right)_{LO} = -\frac{\chi_{11}^{BS}}{\chi_2^S} - \frac{\chi_{11}^{QS}}{\chi_2^S} \frac{\mu_Q}{\mu_B}.$$

Extremely sensitive to the amount of heavy strange and charged particles in the hadronic spectrum. **Additional strange states in PDG2016+  $\uparrow$  the agreement with LQCD up to  $T \approx 145$  MeV.**

## Lattice QCD based Equation of State

The Equation of State (EoS) is the fundamental input to the hydrodynamics evolution that allow us to **test the effect of the assumptions on the number of thermalized quarks.**



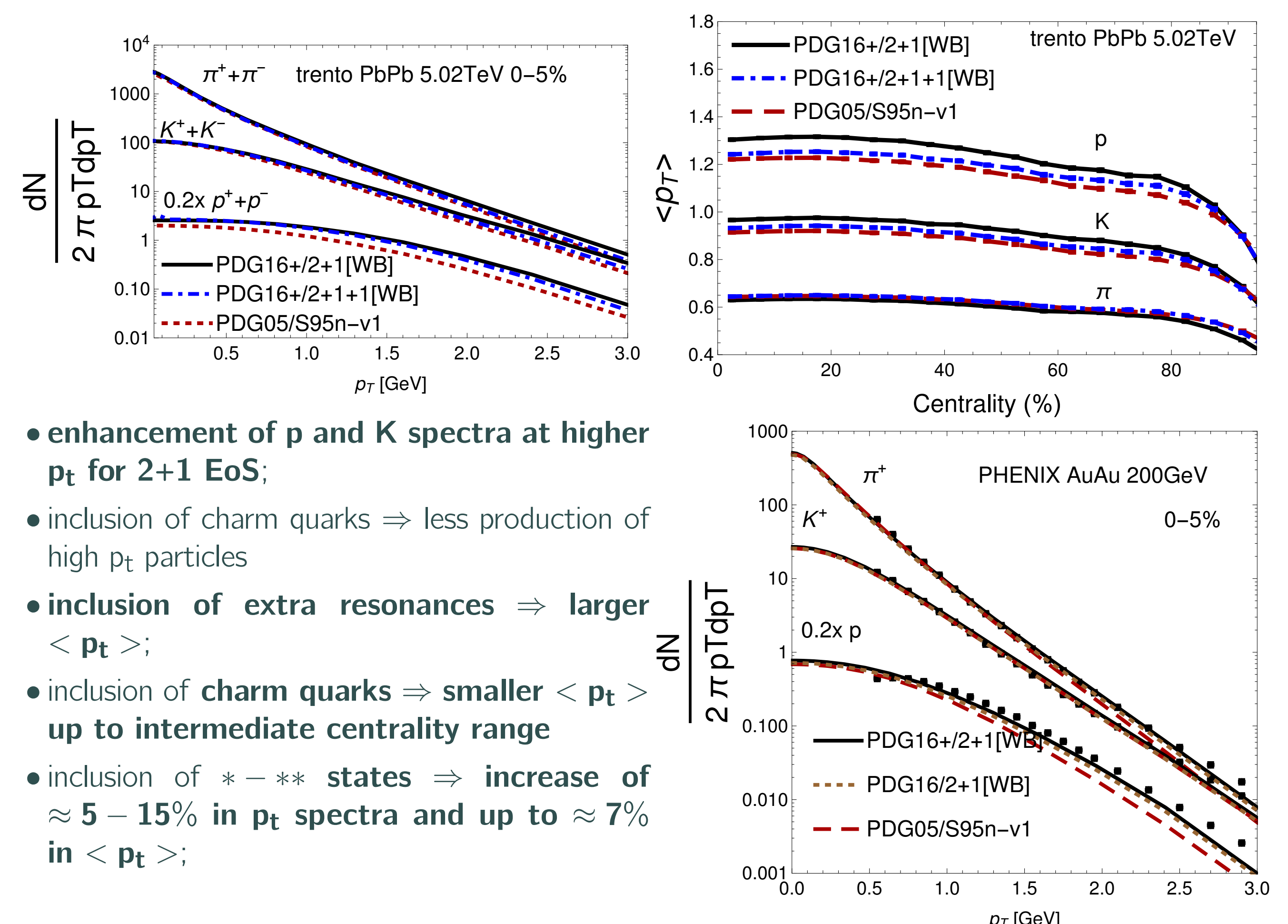
- $T < 153$  MeV  $\Rightarrow$  HRG with PDG16+;
- $T > 153$  MeV  $\Rightarrow$  state-of-the art LQCD fitted EoS for 2+1 [4] and 2+1+1 [5] with thermalized charm degrees of freedom from WB collaboration.

## The hydrodynamical modeling

The parametrization for switching on/off of the hydrodynamics evolution has been chosen to be as consistent as possible with LQCD results. Event-by-event viscous hydrodynamics **v-USPhydro** [6] with **TRENTO initial conditions** [7], assuming  $T_{FO}^{\text{kin}} = T_{FO}^{\text{chem}}$ . For the simulations both at RHIC Au-Au 200 GeV and ALICE Pb-Pb 5.02 TeV the hydrodynamics description is switched on at  $\tau_0 = 0.6$  fm and the hadrons are formed at  $T_{SW} = 150$  MeV.

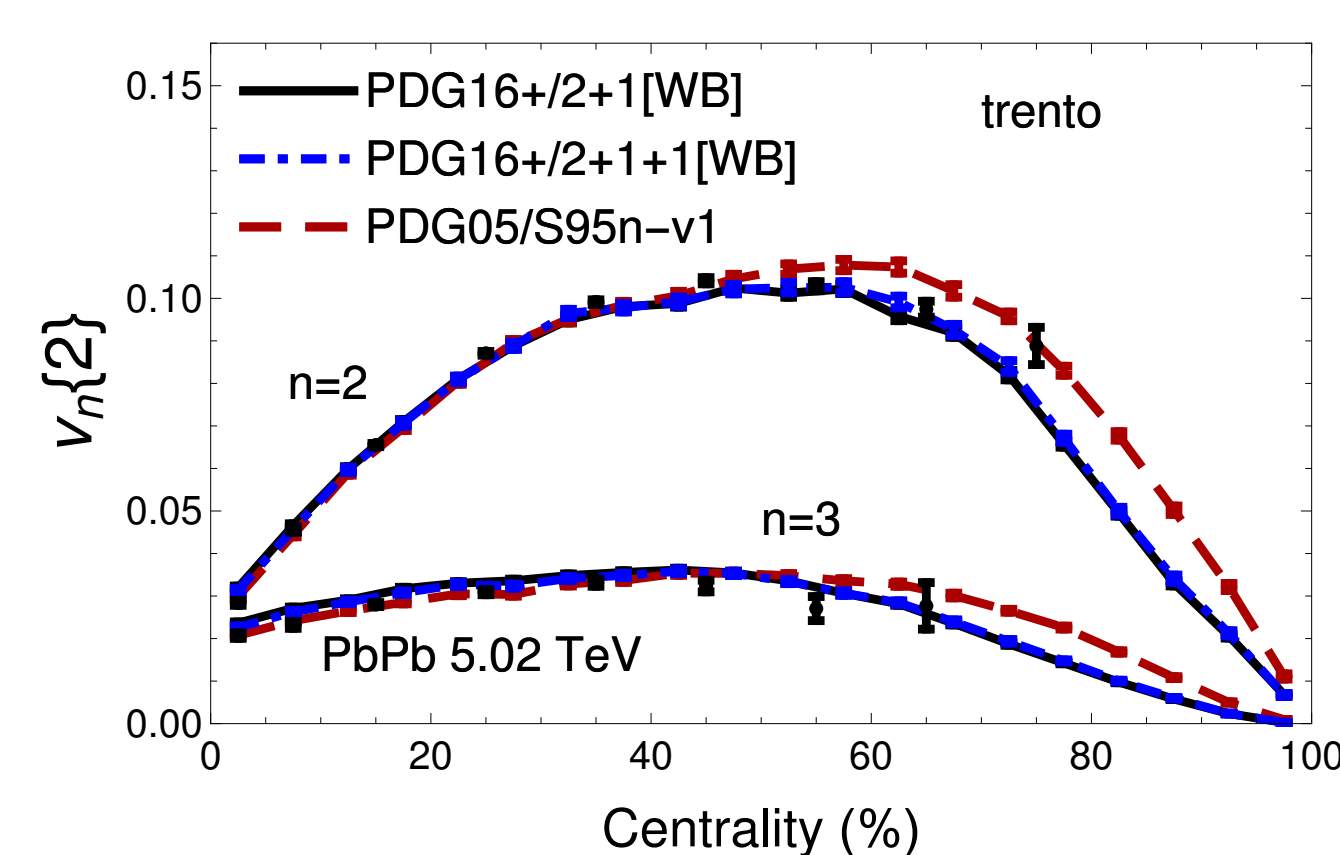
## Results

Effects of the additional resonances on particle spectra and  $\langle p_t \rangle$



- **enhancement of p and K spectra at higher  $p_t$  for 2+1 EoS;**
- inclusion of charm quarks  $\Rightarrow$  less production of high  $p_t$  particles
- **inclusion of extra resonances  $\Rightarrow$  larger  $\langle p_t \rangle$ ;**
- inclusion of **charm quarks  $\Rightarrow$  smaller  $\langle p_t \rangle$  up to intermediate centrality range**
- inclusion of \*-\*\* states  $\Rightarrow$  **increase of  $\approx 5 - 15\%$  in  $p_t$  spectra and up to  $\approx 7\%$  in  $\langle p_t \rangle$ ;**

Effects of the additional resonances on  $\eta/s$  and on flow harmonics



The shear viscosity  $\eta/s$  has been extracted from the comparison of our theoretical results obtained within different EoS to flow harmonics  $v_2\{2\}$ ,  $v_3\{2\}$  obtained from **STAR in Au-Au 200 GeV** [8] and **LHC run2 Pb-Pb 5.02 TeV** [9].

- at RHIC energies all three EoS describe the data quite well  $\Rightarrow$  **no dependence on the chosen EoS, agreement with previous results based on Bayesian analysis [10];**
- thermalized charmed quarks  $\Rightarrow$  smaller  $\eta/s$  value with respect to the 2+1 EoS.
- **higher temperatures probed at LHC run 2 (up to  $T$  600 MeV)  $\Rightarrow$  splitting between the 2+1 and 2+1+1  $\Rightarrow$  different values of  $\eta/s$**

EoS	Au-Au 200 GeV	Pb-Pb 5.02 TeV
PDG05/S95n - v1 [11]	0.5	0.025
PDG16 + /2 + 1 [WB]	0.5	0.047
PDG16 + /2 + 1 + 1 [WB]	0.5	0.04

## Conclusions

- EoS obtained by matching the **state-of-the-art LQCD calculations for 2+1/2+1+1 quark flavors** to a HRG model based **PDG2016+** spectrum containing extra resonances relevant for reproducing lattice data on partial pressures;
- **the inclusion of additional \*-\*\* states:**
  - **increase the agreement with LQCD data up to  $T \approx 145$  MeV, close to the crossover region;**
  - **enhance the production of particles at higher  $p_t$  and leads to a higher  $\langle p_t \rangle$ ;**
  - results for all three EoS obtained at RHIC energies are in agreement with previous Bayesian analysis;
  - **$\eta/s$  ratio at LHC run2 energies depends on the EoS** and there is roughly a 15% difference in 2+1 and 2+1+1 EoS results which should increase at higher temperatures.

## References

- [1] A. Bazavov et al. *Phys. Rev. Lett.*, 113(7):072001, 2014.
- [2] Paolo Alba et al. *Phys. Rev.*, D96(3):034517, 2017.
- [3] C. Patrignani et al. *Chin. Phys.*, C40(10):100001, 2016.
- [4] Szabocs Borsanyi, Zoltan Fodor, Christian Hoelbling, Sandor D. Katz, Stefan Krieg, and Kalman K. Szabo. *Phys. Lett.*, B730:99–104, 2014.
- [5] Sz. Borsanyi et al. *Nature*, 539(7627):69–71, 2016.
- [6] J. Noronha-Hostler, J. Noronha, and F. Grassi. *Phys. Rev.*, C90(3):034907, 2014.
- [7] J. Scott Moreland, Jonah E. Bernhard, and Steffen A. Bass. *Phys. Rev.*, C92(1):011901, 2015.
- [8] L. Adamczyk et al. 2017.
- [9] Jaroslav Adam et al. *Phys. Rev. Lett.*, 116(13):132302, 2016.
- [10] Jonah E. Bernhard, J. Scott Moreland, Steffen A. Bass, Jia Liu, and Ulrich Heinz. *Phys. Rev.*, C94(2):024907, 2016.
- [11] Pasi Huovinen and Pter Petreczky. *Nucl. Phys.*, A837:26–53, 2010.