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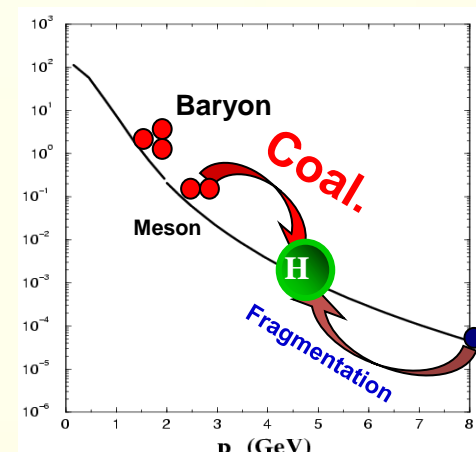



# Heavy hadrons production by coalescence in pp and AA collisions at RHIC and LHC

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We study with a coalescence model the  $p_T$  spectra of charmed hadrons  $D^0$ ,  $\Lambda_c$  and the respective baryon to meson ratio at RHIC and LHC in a range of transverse momentum up to 10 GeV, including both the contribution from decays of heavy hadron resonances and the fragmentation of heavy quarks do not undergo the coalescence process. The  $p_T$  dependence of the heavy baryon/meson ratios is found to be sensitive to the masses of quarks, (in particular we found that the  $\Lambda_c/D^0$  can reach values of about 1-1.5 at  $p_T \sim 3$  GeV) with a quite weak dependence in transverse momentum respect to light hadrons. We also make comparisons with others coalescence models and with predictions from thermal models. In the same framework we finally evaluate the baryon to meson ratios in pp collisions assuming, as suggested in recent works, that at the present LHC energies there can be the formation of QGP matter also in this type of collisions. Applying our model we find a considerable decrease for the peak value and a slightly different behaviour in the  $p_T$  dependence.



Statistical factor colour-spin-isospin

$$\frac{dN_{Hadron}}{d^2p_T dy} = g_H \int \prod_{i=1}^n p_i \cdot d\sigma_i \frac{d^3p_i}{(2\pi)^3 E_i} f_{q_i}(x_i, p_i) f_W(x_1, \dots, x_n; p_1, \dots, p_n) \delta\left(p_T - \sum_i p_{iT}\right)$$

i=2 → mesons ; i=3 → baryons

Parton Distribution function

charm distribution function from parton simulations solving Boltzmann transport equation ( $R_{AA}$  and  $v_2$  see below).

Hadron Wigner function

Fireball Parameters

Constraints from Experiment

RHIC → LHC

dE<sub>T</sub>/dy ~ 740 GeV → 2100 GeV

dN<sub>ch</sub>/dy ≈ 670 → 1600

T<sub>c</sub> ~ 165 MeV

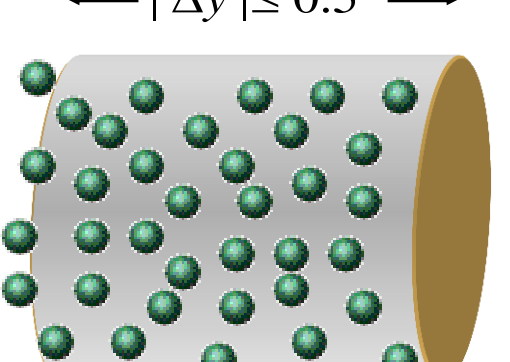
Lifetime and Volume implied

τ ~ 4.5 fm/c → 7.8 fm/c

β<sub>0</sub> = 0.37 → 0.63

V ~ 1100 fm<sup>3</sup> → 2500 fm<sup>3</sup>

In agreement with HBT



COALESCENCE INTEGRAL

LIGHT QUARK

Thermal Distribution (p<sub>T</sub>< 2 GeV)

$$\frac{dN_q}{d^2r_T d^2p_T} = \frac{g_q m_T}{(2\pi)^3} \exp\left(-\frac{\gamma_T(m_T - p_T \cdot \beta_T)}{T}\right)$$

Collective flow β<sub>T</sub> = β<sub>0</sub>  $\frac{r}{R}$

Minijet Distribution (p<sub>T</sub>> 2 GeV)

HEAVY QUARK (initial state)

Fixed-Order plus Next-to-Leading-Log (FONLL)

M. Cacciari, P. Nason, R. Vogt, PRL 95 (2005) 122001

$$f_M(x_1, x_2; x_1, p_2) = A_W \exp\left(-\frac{x_{r1}^2}{\sigma_r^2} - p_{r1}^2 \sigma_r^2\right)$$

A<sub>W</sub> is a normalization constant fixed to guarantee that in the limit p<sub>T</sub> → 0 we have all the charm hadronizing. While σ<sub>r</sub> is the covariant width parameter, it can be related to the oscillator frequency ω by σ = 1/√μω where μ = (m<sub>1</sub>m<sub>2</sub>)/(m<sub>1</sub> + m<sub>2</sub>) is the reduced mass. The width of f<sub>M</sub> is linked to the size of the hadron and in particular to the root mean square charge radius of the meson by:

$$\langle r^2 \rangle_{ch} = \frac{3}{2} \frac{Q_1 m_2^2 + Q_2 m_1^2}{(m_1 + m_2)^2} \sigma_r^2$$

$\langle r^2 \rangle_{ch}^{D^0} = 0.184 \text{ fm}^2$   
 $\sigma_p^{D^0} = \sigma_{r^{-1}D^0} = 0.283 \text{ GeV}$

$\langle r^2 \rangle_{ch}^{\Lambda_c} = 0.15 \text{ fm}^2$   
 $\sigma_{p1}^{\Lambda_c} = 0.18 \text{ GeV}$   
 $\sigma_{p2}^{\Lambda_c} = 0.342 \text{ GeV}$

## Introduction

The bulk properties of the matter created in the Quark Gluon Plasma are governed by the light quarks and gluons while heavy quarks like charm or bottom quarks are useful probes of the QGP properties. In their final state the charm quarks appear as constituent of charmed hadrons mainly D mesons and  $\Lambda_c$ ,  $\Sigma_c$  baryons. The experimental advances in the study not only of heavy mesons like D mesons but also of  $\Lambda_c$  baryons are important in order to have a new insight in understanding the hadronization mechanism in the QGP. Recent experimental results from STAR collaboration have shown a  $\Lambda_c/D^0 \approx 1.3 \pm 0.5$  for 3- $p_T < 6$  GeV which is a very large enhancement compared to the value predicted by the charm hadron fragmentation ratio or by the PYTHIA for p-p collisions. The idea of the coalescence model comes from the fact that comoving partons in the quark-gluon plasma combine their transverse momentum to produce a final state meson or baryon with higher transverse momentum. In the heavy quark sector, coalescence leads to a modification of the relative abundance of the different heavy hadron species produced. In particular this can manifest in a baryon-to-meson enhancement for charmed hadrons.

## Heavy Baryon to Meson Ratios

Fragmentation fraction for charm quarks

into hadrons →  $\left(\frac{\Lambda_c^*}{D^0}\right)_{e^+e^-} \simeq 0.1$

We use the Peterson fragmentation function

C. Peterson, D. Schallatter, I. Schmitt, P.M. Zerwas PRD 27 (1983) 105

$$D_{f \rightarrow h}(z) \propto \frac{1}{z \left[1 - \frac{1}{z} - \frac{\epsilon}{1-z}\right]^2}$$

The parameter  $\epsilon$  for D meson hadronization fixed by pp collisions data. For baryons we fix it in accordance with e<sup>+</sup>e<sup>-</sup> collisions as done in *S.K. Das et al. PRD94 (2016) no.11, 114039*. Measurement in e<sup>+</sup>e<sup>-</sup> p, pp and e<sup>+</sup>e<sup>-</sup> e<sup>+</sup>e<sup>-</sup> are in agreement within uncertainties: fragmentation at most independent of the specific production process

Comparison with thermal model

• Compared to light baryon/meson ratio the  $\Lambda_c/D^0$  ratio has a larger width (flatter)

Coal with wave function width σ<sub>p</sub> of D<sup>0</sup> and Λ<sub>c</sub> changed to have Λ<sub>c</sub>/D<sup>0</sup>=thermal ratio at p<sub>T</sub>→0

Evolution from RHIC to LHC

Coalescence predict similar baryon/meson ratio of about 1 for both energies. Fragmentation is established from the experimental measured fragmentation fraction into final hadrons channels, and it remains the same changing the collision energy, the ratio from fragmentation fraction is ~ 0.09 .

Even if the only coalescence ratio and the only fragmentation ratio remain similar at RHIC and LHC, the **combined ratio is different** because the coalescence over fragmentation ratio at LHC is smaller than at RHIC.

Therefore at LHC the larger contribution in particle production from fragmentation leads to a final ratio that is smaller than at RHIC.

## Heavy Hadron spectrum

In our calculations we take into account main hadronic channels, including the ground states and the first excited states for D and Λ<sub>c</sub>

MESONS		
<ul style="list-style-type: none"> <li>D<sup>+</sup> (l=1/2, J=0)</li> <li>D<sup>0</sup> (l=1/2, J=0)</li> <li>D<sub>s</sub><sup>+</sup> (l=0, J=0)</li> </ul>		
Resonances		
<ul style="list-style-type: none"> <li>D<sup>+</sup>* (l=1/2, J=1) → D<sup>0</sup> π<sup>+</sup> B.R. 68%</li> <li>→ D<sup>+</sup> X B.R. 32%</li> <li>D<sup>0</sup>* (l=1/2, J=1) → D<sup>0</sup> π<sup>0</sup> B.R. 62%</li> <li>→ D<sup>0</sup> γ B.R. 38%</li> <li>D<sub>s</sub><sup>+</sup>* (l=0, J=1) → D<sub>s</sub><sup>+</sup> X B.R. 100%</li> <li>D<sub>so</sub><sup>+</sup>* (l=0, J=0) → D<sub>s</sub><sup>+</sup> X B.R. 100%</li> </ul>		
BARYONS		
<ul style="list-style-type: none"> <li>Λ<sub>c</sub><sup>+</sup> (l=0, J=1/2)</li> </ul>		
Resonances		
<ul style="list-style-type: none"> <li>Λ<sub>c</sub><sup>+</sup>*(2595) (l=0, J=1/2) → Λ<sub>c</sub><sup>+</sup> B.R. 100%</li> <li>Λ<sub>c</sub><sup>+</sup>*(2625) (l=0, J=3/2) → Λ<sub>c</sub><sup>+</sup> B.R. 100%</li> <li>Σ<sub>c</sub><sup>+</sup>*(2455) (l=1, J=1/2) → Λ<sub>c</sub><sup>+</sup> π B.R. 100%</li> <li>Σ<sub>c</sub><sup>+</sup>*(2520) (l=1, J=3/2) → Λ<sub>c</sub><sup>+</sup> π B.R. 100%</li> </ul>		

At RHIC coalescence and fragmentation for D<sup>0</sup> are comparable, at LHC the fragmentation becomes dominant. The same trend also for Λ<sub>c</sub><sup>+</sup>, but the fragmentation fraction of charm into Λ<sub>c</sub> is so small that coalescence is the dominant mechanism of hadronization.

## References

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- Y. Oh, C.M. Ko, S.H. Lee, S. Yasui, *Phys. Rev.* C79,044905 (2009)
- F. Scardina, S.K. Das, V. Minissale, S. Plumari, V. Greco, *Phys.Rev.* C96 (2017) no.4, 044905

## R<sub>AA</sub> and v<sub>2</sub>

Heavy quark hadronization by means of a hybrid model of fragmentation plus coalescence help to increase both the  $R_{AA}$  and  $v_2$  close to the data,

Charm distribution function from parton simulations solving Boltzmann transport equation that give good description of both  $R_{AA}$  and  $v_2(p_T)$  from RHIC to LHC energies

## Prediction: coalescence in pp

Calculation for LHC@5.02 TeV:

- further decrease of  $\Lambda_c/D^0$  ratio in *Pb+Pb* and broadening of the peak momentum dependence
- Disappearance of the peak in *pp*, but still different with respect to only fragmentation