



# INI - 24: Plasmon mass scale in classical nonequilibrium gauge theory in two and three dimensions

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## Introduction

**Classical Yang-Mills (CYM)** calculations have been used to model the **pre-thermal evolution** of the strongly interacting matter created in **ultrarelativistic heavy-ion collisions**. Our aim is to **study the limits of the quasiparticle picture** in real time classical Yang-Mills theory on a **lattice** in two and three spatial dimensions.

## Initial conditions

**3D:** The **initial quasiparticle spectrum** satisfies:

$$f(k, t=0) = \frac{n_0}{g^2} \frac{k}{\Delta} \exp\left(\frac{-k^2}{2\Delta^2}\right).$$

With **momentum scale**  $\Delta$  and **occupation number**  $n_0$ .

**2D:** **Gauge fixing deforms initial quasiparticle spectrum.**

Use similar IC as in 3D and **measure the occupation number**  $n_0^{\text{eff}}$  and **momentum scale**  $Q_{\text{eff}}$  **gauge invariantly**

$$Q_{\text{eff}}^2(t) = \frac{1}{2} \frac{\langle \text{Tr}(\mathbf{D} \times \mathbf{B})^2 \rangle}{\langle \text{Tr}(\mathbf{B}^2) \rangle}. \quad (1)$$

$$n_0^{\text{eff}} = \frac{\pi g^2}{(N_c^2 - 1)} \frac{\epsilon^{2d}}{Q_{\text{eff}}^3}. \quad (2)$$

## Extracting plasmon mass, 3 methods

**DR** **Effective dispersion relation** in the Coulomb gauge

$$\omega^2(k) = \frac{\langle |\dot{E}_i^a(k)|^2 \rangle}{\langle |E_i^a(k)|^2 \rangle},$$

fit as  $\omega^2 = ak^2 + \omega_{pl}^2$ .

**UE** **Add a uniform electric field** at  $t = t_0$ , **measure oscillations** of electric and magnetic energy vs  $t$  [3].

**HTL** **Perturbation theory, Hard Thermal Loop**

$$\omega_{pl}^2 = \frac{4}{3} g^2 N_c \int \frac{d^3k}{(2\pi)^3} \frac{f(k)}{k}.$$

## Conclusions

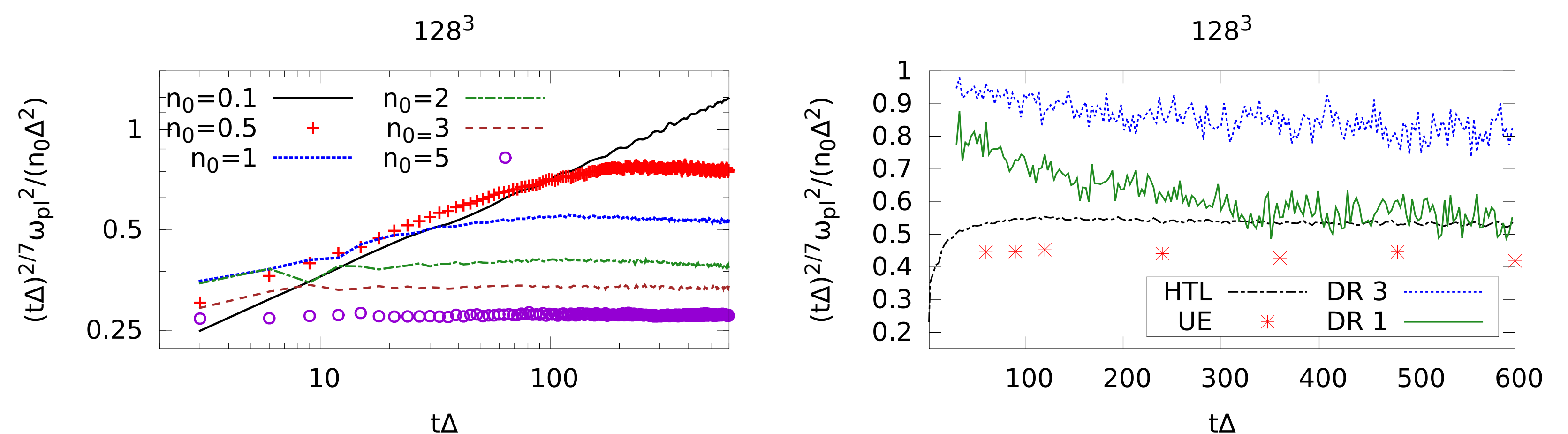
- Study the plasmon mass scale in CYM using 3 methods. The DR agrees with the other methods within a factor of two. **UE and HTL are in rough agreement.**
- The **UE and HTL methods agree in the continuum limit in 3D**. In 2D the difference between the two persists.
- The UE method is insensitive to ultraviolet and infrared cutoffs (3D). In 2D the extracted mass scale increases in the continuum limit for all methods.
- Time-dependence of  $\omega_{pl}^2$  consistent with  $t^{-2/7}$  (3D)  $t^{-1/3}$  (2D) power law.
- More on quasiparticles: **extraction of spectral function** from real time lattice simulations [4], see **poster** by **K. Boguslavski INI - 02**

## References

- [1] T. Lappi, J. Peuron, Phys. Rev. D95, 014025
- [2] T. Lappi, J. Peuron, Phys. Rev. D97, 034017
- [3] A. Kurkela, G. D. Moore, Phys.Rev. D89, 074036
- [4] K. Boguslavski, A. Kurkela, T. Lappi, J. Peuron, arXiv:1804.01966 [hep-ph]

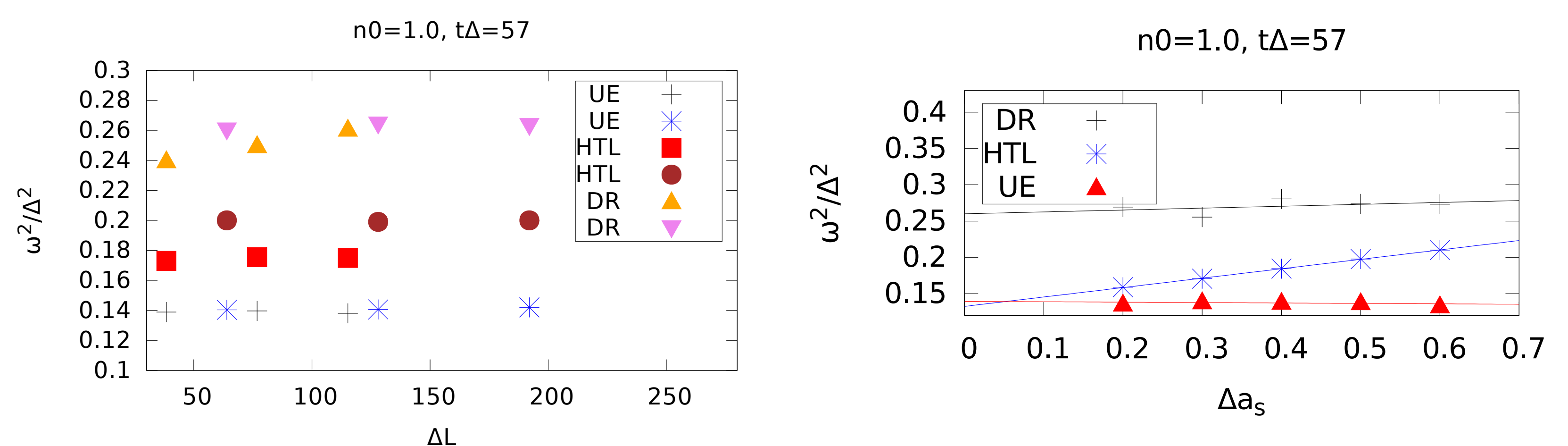
## Results & dependence on the lattice cutoffs, 3D [1]

Dependence on **occupation number** and **time**:



- The **late-time evolution** is **consistent with  $t^{-2/7}$  power law** [3].
- Higher occupation number  $\rightarrow$  asymptotic regime reached faster.
- The DR method depends on maximum  $k^2/\Delta^2$  in fit (DR 1 vs. DR 3).

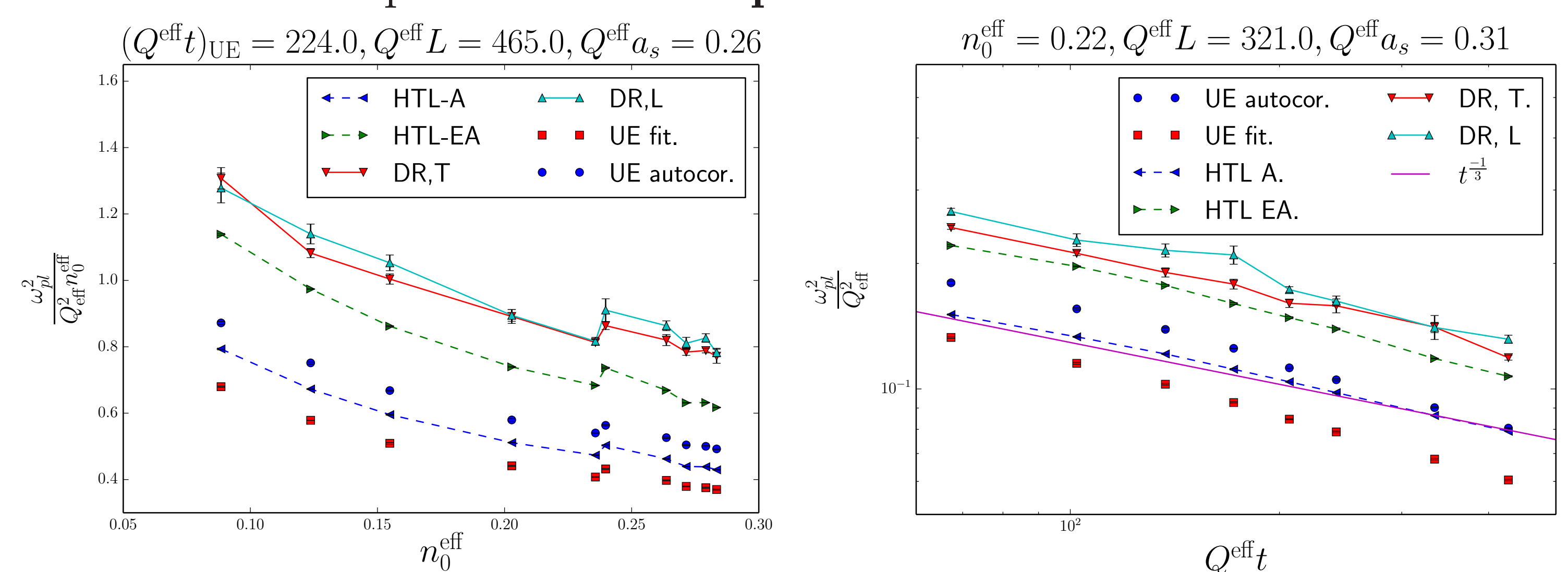
Dependence on **infrared cutoff** and **ultraviolet cutoff**:



- Left: infrared cutoff (lattice size  $L\Delta$ ) dependence with two different ultraviolet cutoffs (lattice spacings  $a_s\Delta = 0.3$  [up],  $a_s\Delta = 0.5$  [down]). **We observe no significant IR-cutoff dependence.**
- Right: UV cutoff dependence. **The HTL and UE agree in the continuum limit.**

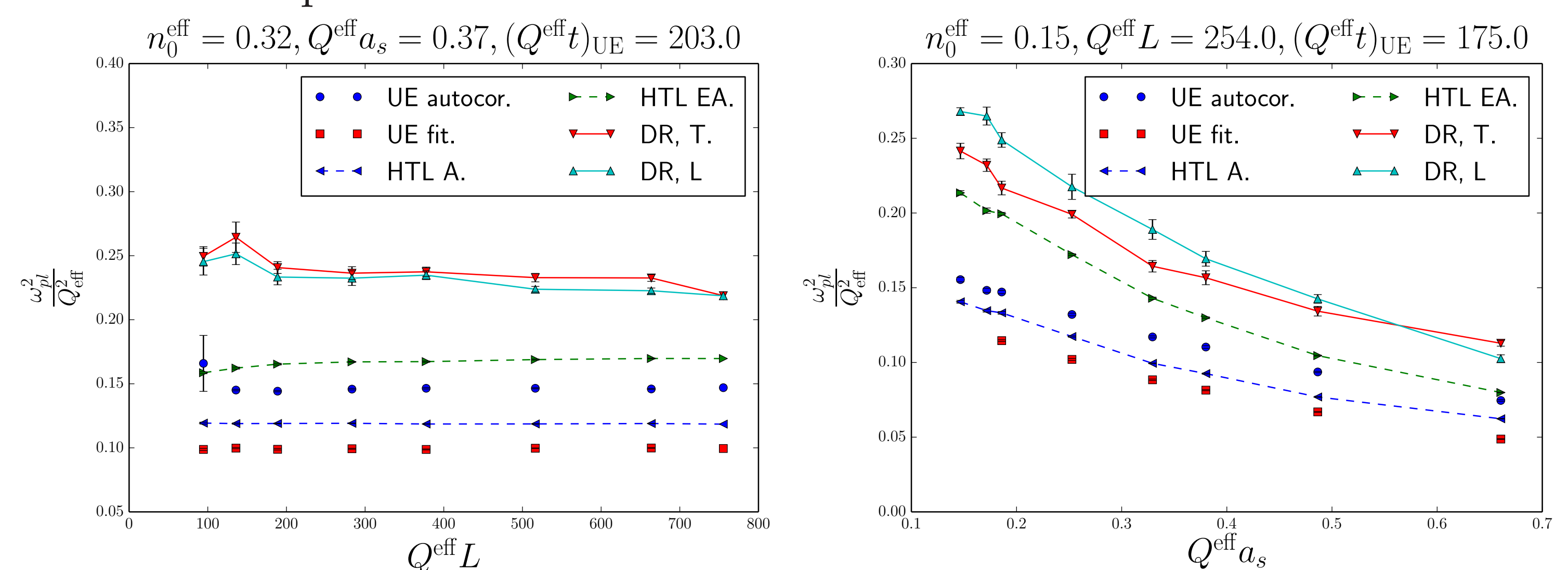
## Results & dependence on the lattice cutoffs, 2D [2]

Dependence on **occupation number** and **time**:



- Right: **Late-time-evolution is consistent with  $t^{-1/3}$  power law.**
- Left: **The occupation number dependence similar to 3D** - faster decrease of the mass scale for higher occupation number

Dependence on **infrared cutoff** and **ultraviolet cutoff**:



- Left: **Infrared cutoff** dependence like in 3D - **No cutoff dependence.**
- Right: **Ultraviolet cutoff** dependence different from 3D: **mass scale increases in the continuum limit.**
- Overall we find a rough agreement between UE and HTL methods. DR method agrees with the other methods within a factor of 2.