

Particle distributions from hydrodynamics

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Introduction

Two analytical solutions for accelerating hydrodynamics are derived to describe the relativistic p+p collisions and heavy ion collisions in Rindler coordinate.

- An exact solution of ideal hydrodynamics is used to estimate the initial thermodynamic density at RHIC and LHC[1].

- A perturbative solution of viscous hydrodynamics (with the Navier-Stokes correction) is obtained. $\sqrt{S_{NN}} = 5.44$ TeV Xe+Xe collisions and $\sqrt{S_{NN}} = 13$ TeV p+p collisions results are predicted[2].

Conservation Law

The conservation laws for temperature profile in Rindler coordinate are:

$$\tau \frac{\partial T}{\partial \tau} + \tanh(\Omega - \eta_s) \frac{\partial T}{\partial \eta_s} + \frac{\Omega'}{\kappa} T = \frac{\Pi_d}{\kappa} \frac{\Omega'^2}{\tau} \cosh(\Omega - \eta_s), \quad (1)$$

$$\tanh(\Omega - \eta_s) \left[\tau \frac{\partial T}{\partial \tau} + T \Omega' \right] + \frac{\partial T}{\partial \eta_s} =$$

$$\frac{\Pi_d}{\tau} (2\Omega'(\Omega' - 1) + \Omega'' \coth(\Omega - \eta_s) \sinh(\Omega - \eta_s)),$$

where Ω is the flow element rapidity, and Π_d is a function of shear viscosity and bulk viscosity.

Final state observations

Final state spectra are yielded from hydrodynamic solutions[1, 2] and Cooper-Frye formula,

$$\frac{d^2 N}{2\pi p_T d p_T dy} = \frac{g}{(2\pi)^3} \int p_\mu d\Sigma^\mu f, \quad (3)$$

The relationship between $\sqrt{S_{NN}}$ and acceleration parameter λ is

$$\lambda = 1 + A \left(\sqrt{S_{NN}} / \sqrt{S_0} \right)^{-B}. \quad (4)$$

The improving initial energy density is:

$$\epsilon_{\text{corr}} = (2\lambda - 1) \left(\frac{\tau_f}{\tau_0} \right)^{\lambda-1} \frac{\langle E \rangle}{S_\perp \tau_0 R^2} \frac{dN}{dy} \Big|_{y=y_0}. \quad (5)$$

Conclusion and summary

- In a limit η_s region, CKCJ [1] model presents new solutions of ideal hydrodynamics.
- CNC [1] model can be using to estimate the initial thermodynamic quantities.
- Accelerating viscous hydrodynamics [2] is an effective theory to study the effect of both longitudinal accelerate and viscosity.

References and Acknowledgements

- [1] M. Csanád, et. al, Universe 3 (2017) no.1, 9.
Z. F. Jiang, et.al, arXiv: 1711.10740.
T. Csörgő, et.al, arXiv: 1805.01427.
- [2] Z. F. Jiang, et.al, in preparation.

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Exact solution of ideal hydrodynamic[1]

Exact solution of ideal hydrodynamic can be obtained in Rindler coordinate. One can extracted **acceleration parameter characteristics** from precise $dN/d\eta$ and obtained a series of **improved initial energy density estimation** (for RHIC and LHC).

The temperature profile is

$$T(\tau, H) = T_0 \left(\frac{\tau_0}{\tau} \right)^{\frac{1+\lambda^*}{\kappa}} \frac{1}{\nu(s)} \left[1 + \frac{\kappa-1}{\lambda-1} \sinh^2(H) \right]^{\frac{1+\lambda^*}{2\kappa}}. \quad (6)$$

where H is a parameter related to η_s .

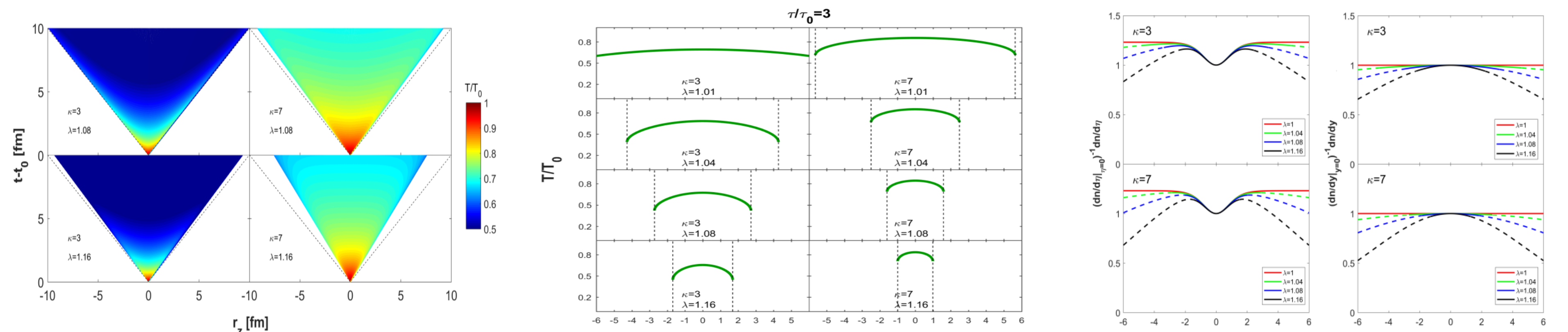


Figure. 1: Left: Heat map of temperature evolution from new CKCJ solutions[1]. Mid: The interesting temperature profile for limit η_s . Right: Effective pseudo-rapidity distribution. Detailed work for CKCJ, arXiv: 1805.01427.

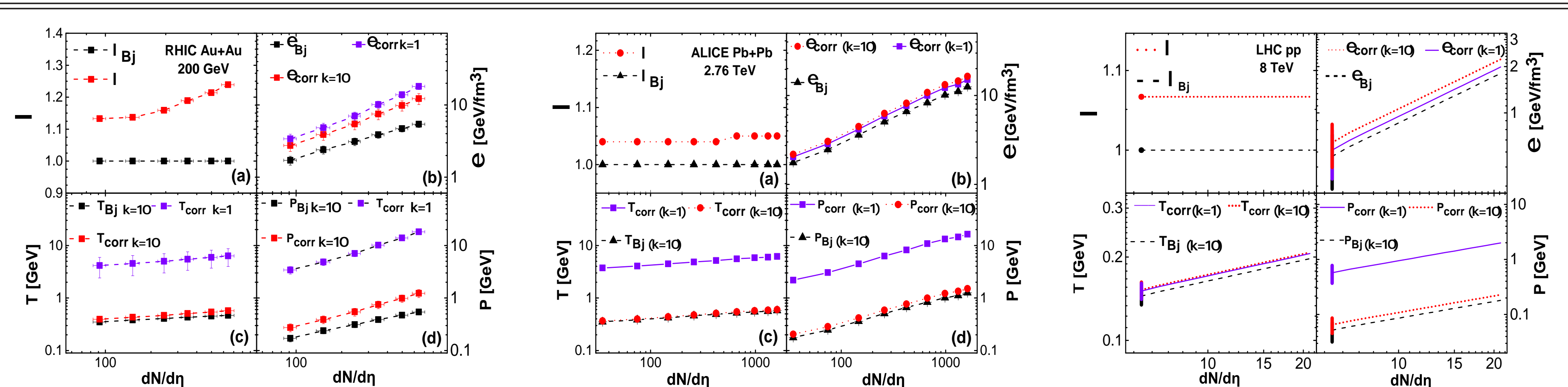


Figure. 2: Accelerate parameter λ evolution, initial energy density, pressure estimation, and temperature profile for RHIC and LHC p+p & A+A collisions. For detailed calculation, see Z. F. Jiang, arXiv: 1711.10740..

Perturbative solution of viscous hydrodynamics[2]

By introducing the Navier-Stokes correction, we obtain a **perturbative solution** which takes acceleration into account the effects of longitudinal acceleration and viscosity.

The new **temperature profile** is

$$T(\tau, \eta_s) = T_0 \left(\frac{\tau_0}{\tau} \right)^{\frac{1+\lambda^*}{\kappa}} \left[\exp\left[-\frac{1}{2}\lambda^* \left(1 - \frac{1}{\kappa}\eta_s^2\right) + \frac{R_0^{-1}}{\kappa-1} \left(2\lambda^* + \exp\left[-\frac{1}{2}\lambda^* \left(1 - \frac{1}{\kappa}\eta_s^2\right) - (2\lambda^* + 1) \left(\frac{\tau_0}{\tau}\right)^{\frac{\kappa-\lambda^*-1}{\kappa}}\right]\right) \right]. \quad (7)$$

where $\lambda^* = \lambda - 1$ the longitudinal acceleration parameter, equation of state $\epsilon = \kappa P$, R_0^{-1} the Reynolds number.

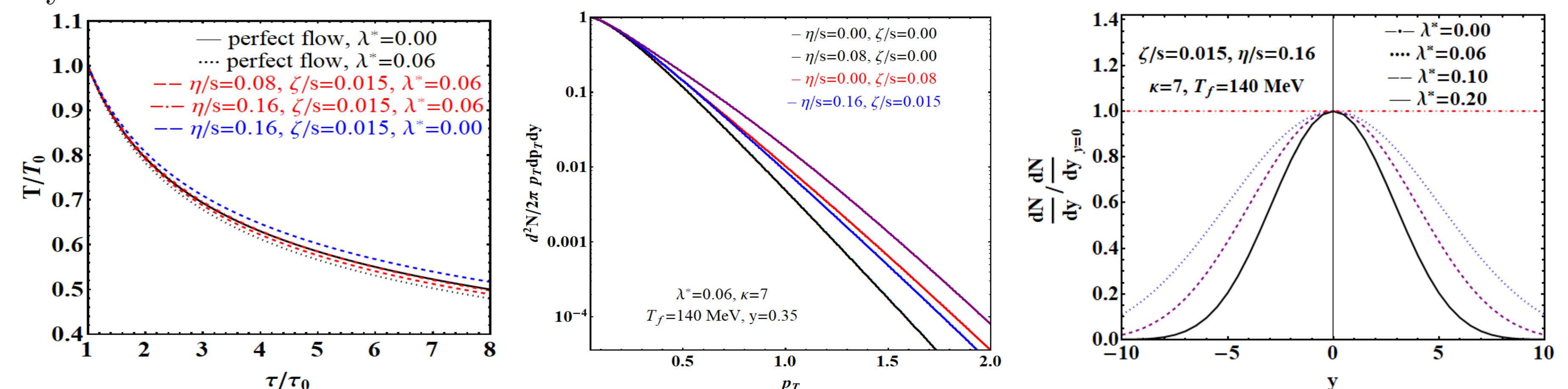


Figure. 3: Left: Temperature profile, Viscosity effect vs. Accelerate effect. Mid: Transverse momentum distribution, viscosity is important! Right: Rapidity distribution for different accelerate (Academic study).

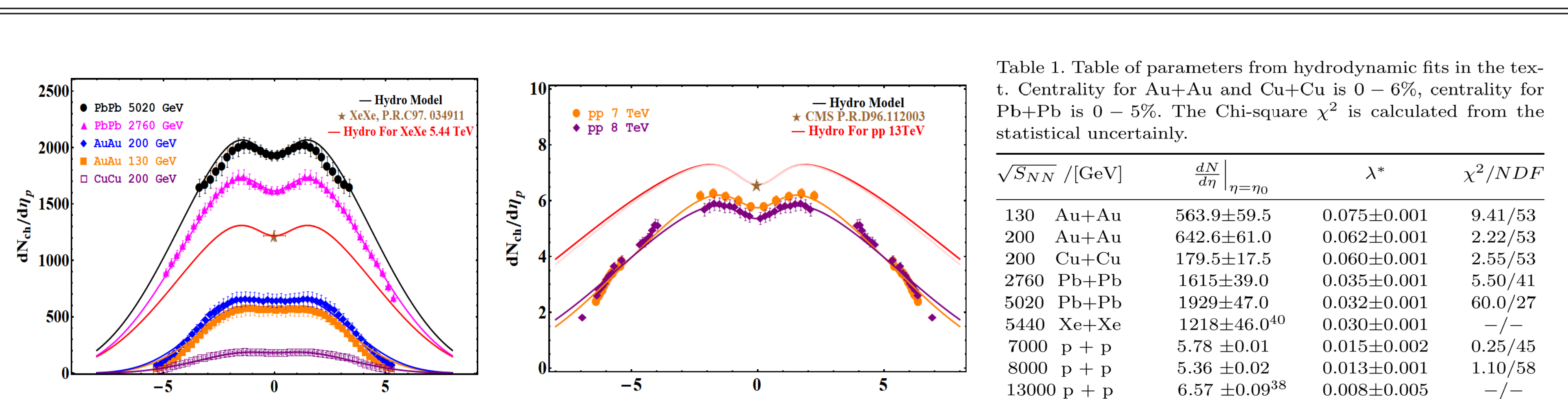


Table 1. Table of parameters from hydrodynamic fits in the text. Centrality for Au+Au and Cu+Cu is 0-6%, centrality for Pb+Pb is 0-5%. The Chi-square χ^2 is calculated from the statistical uncertainty.

$\sqrt{S_{NN}}$ [GeV]	$\frac{dN}{d\eta} \Big _{\eta=0}$	λ^*	χ^2/NDF
130 Au+Au	563.9±59.5	0.075±0.001	9.41/53
200 Au+Au	642.6±61.0	0.062±0.001	2.22/53
200 Cu+Cu	179.5±17.5	0.060±0.001	2.55/53
2760 Pb+Pb	1615±39.0	0.035±0.001	5.50/41
5020 Pb+Pb	1929±47.0	0.032±0.001	60.0/27
5440 Xe+Xe	1218±46.0 ⁴⁰	0.030±0.001	-/-
7000 p+p	5.78±0.01	0.015±0.002	0.25/45
8000 p+p	5.36±0.02	0.013±0.001	1.10/58
13000 p+p	6.57±0.09 ³⁸	0.008±0.005	-/-

Figure. 4: The pseudo-rapidity distributions from our model calculation compared to the RHIC and LHC experimental data. Left: different nucleus-nucleus colliding systems at RHIC & LHC, prediction for $\sqrt{S_{NN}} = 5.44$ TeV Xe+Xe collisions results from hydrodynamics. Right: prediction for 7 TeV and 8 TeV, and $\sqrt{S_{NN}} = 13$ TeV p+p collisions prediction. $A=0.045$ and $B=0.23$ for A+A collisions and $A=0.12$ and $B=1.07$ for p+p collisions, $\sqrt{S_0} = 1$ GeV.