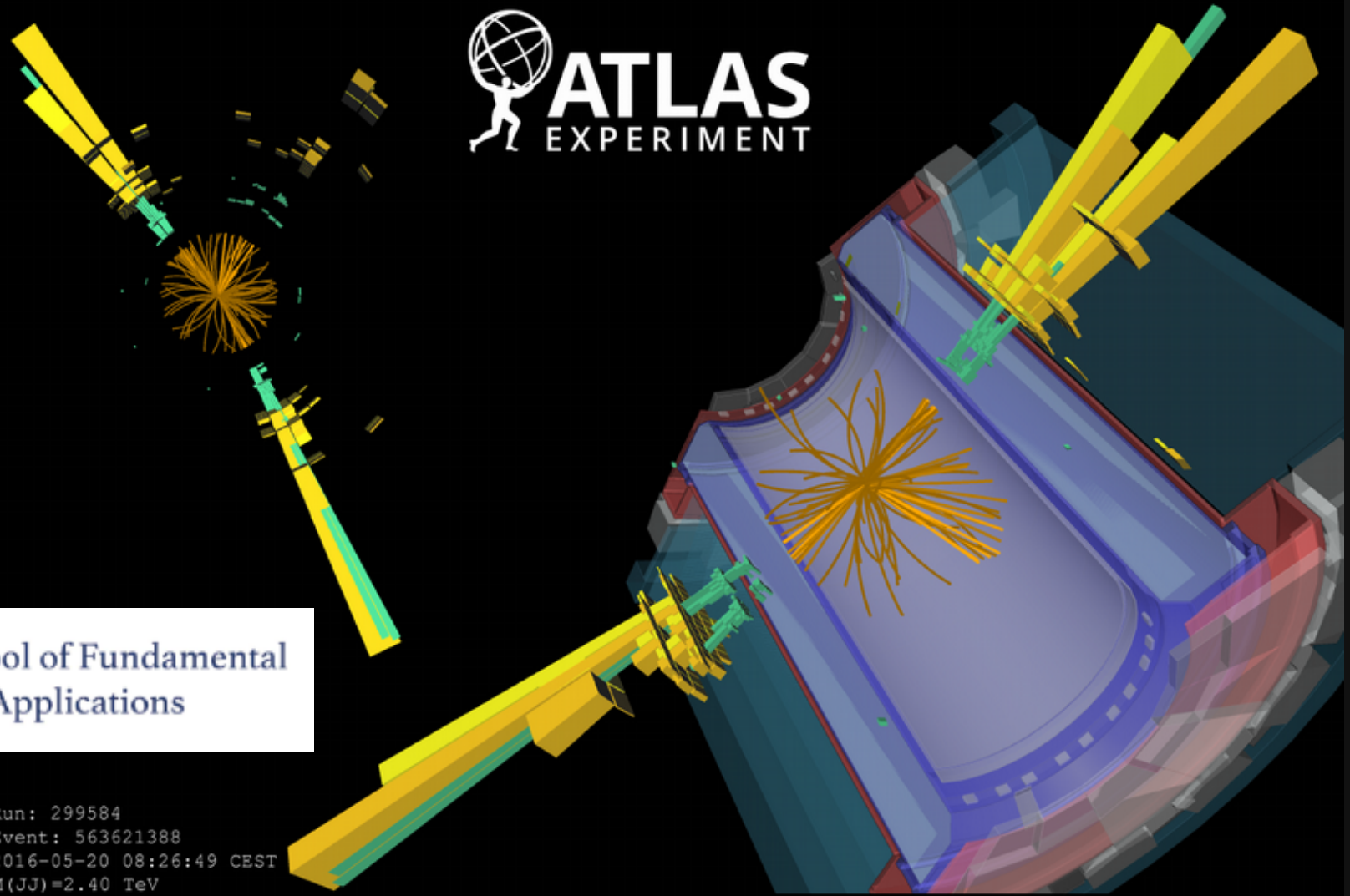


SM Measurements at the LHC

QCD, Electroweak and Top Physics

Dave Charlton (University of Birmingham, UK)
Dave.Charlton@cern.ch

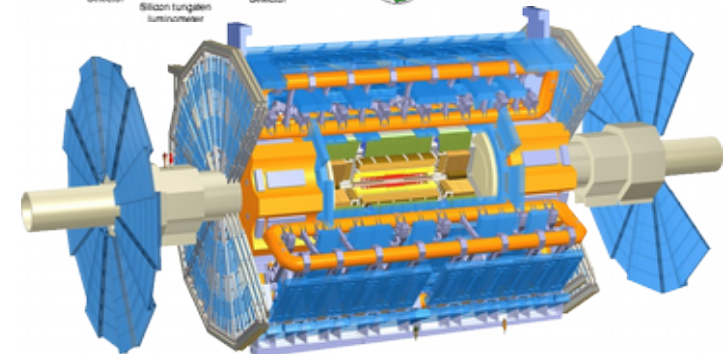
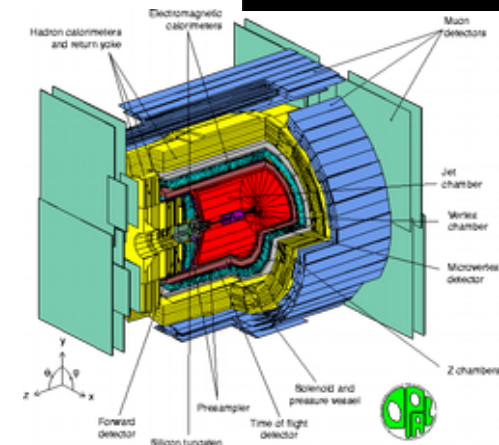
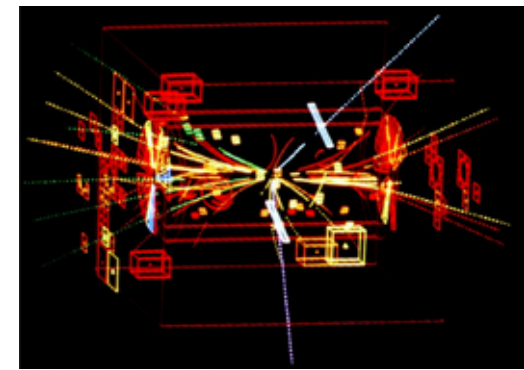


African School of Fundamental
Physics and Applications

Dave Charlton

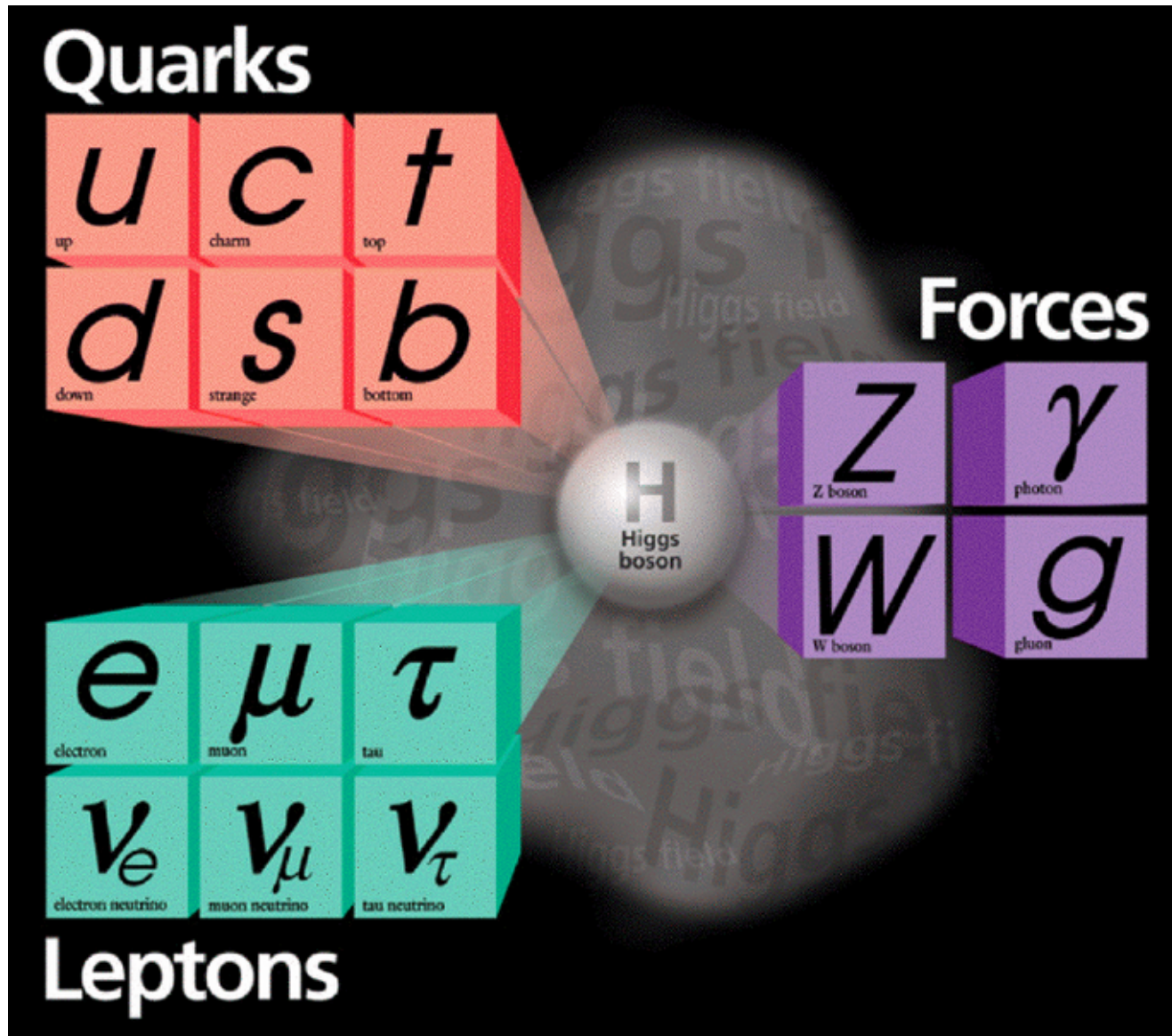
About me:

- PhD student on UA1 experiment 1985-1988 (search for the top quark)
- Moved at start of 1989 to OPAL experiment at LEP, stayed to the end (2000) - electroweak physics with Z and W bosons
- Since 1998, ATLAS experiment at the LHC at CERN
 - Spokesperson (Head) of ATLAS 2013-2017
 - Previously deputy Spokesperson (2009-2013), Physics Coordinator (2008-2009)
 - Worked on calorimeter triggering, silicon tracker construction, analysis of multi-boson production
- Poynting Professor of Physics at the University of Birmingham in the UK since 2017 (I've been with Birmingham since 1994, professor since 2005)



The Standard Model

Matter Particles



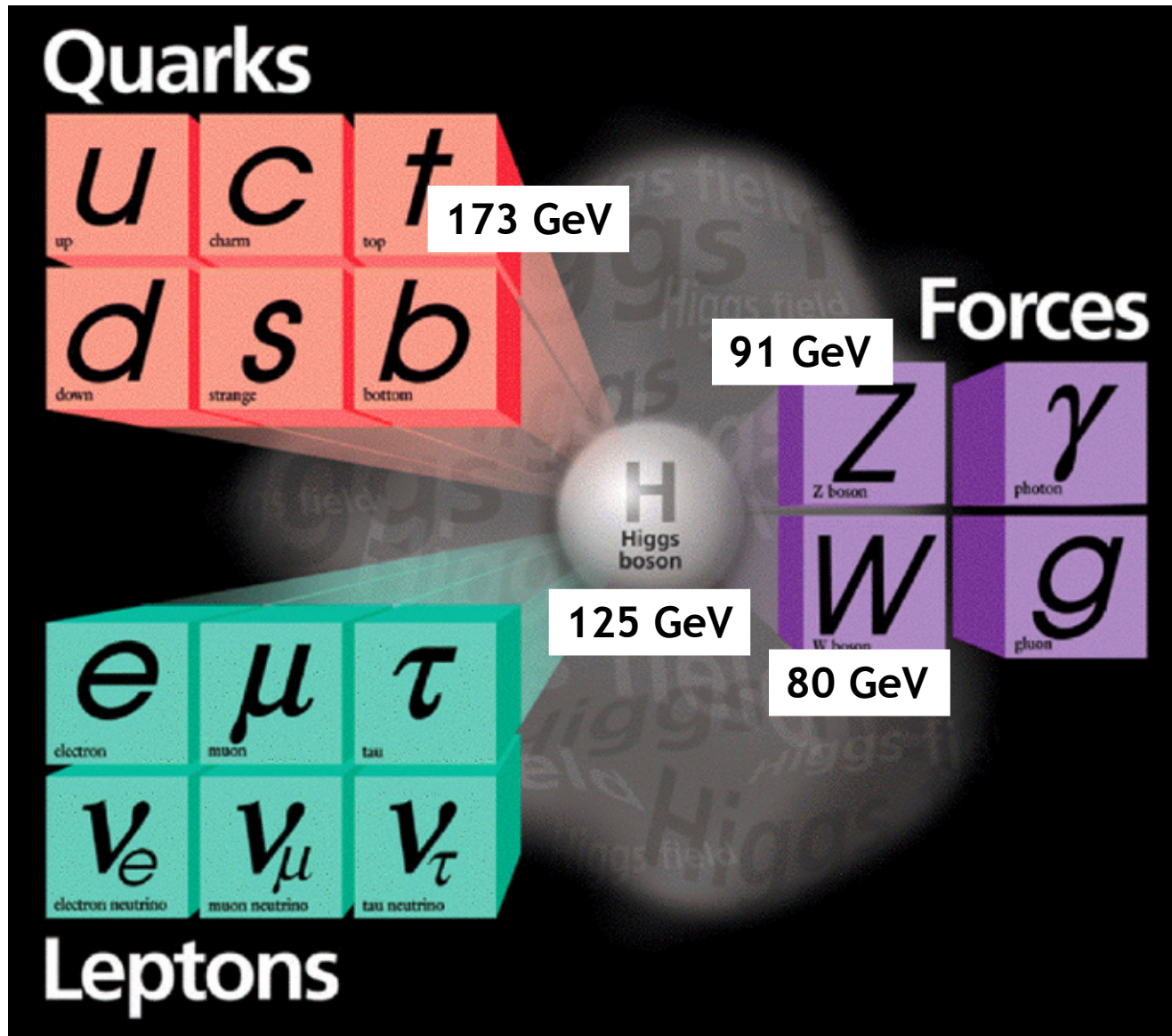
Force-carriers

Fermions

Bosons

The Standard Model

Matter Particles



Force-carriers

Fermions

Bosons

Rates, luminosities and cross-sections

In a collider, the rate, dN_a/dt , of events produced for a given process a is:

$$dN_a/dt = \sigma_a L$$

where

- σ_a is the *cross-section* for the process
 - units of area (1 barn = $10^{-28} \text{ m}^2 = 10^{-28} \text{ cm}^2$)
 - typically mb, μb , nb, pb and fb are (all) met for different processes!
 - it depends on the physics process, eg. $pp \rightarrow W + \text{anything}$ and the centre-of-mass energy \sqrt{s}
- L is the instantaneous luminosity, usually called the luminosity
 - units of inverse-area per unit time (typically $\sim 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ at LHC)
 - Process-independent, depends only on the beam characteristics

Integrated version:

$$N_a = \sigma \int L dt$$

Where $\int L dt$ is the integrated luminosity, typically expressed in fb^{-1}

LHC physics landscape

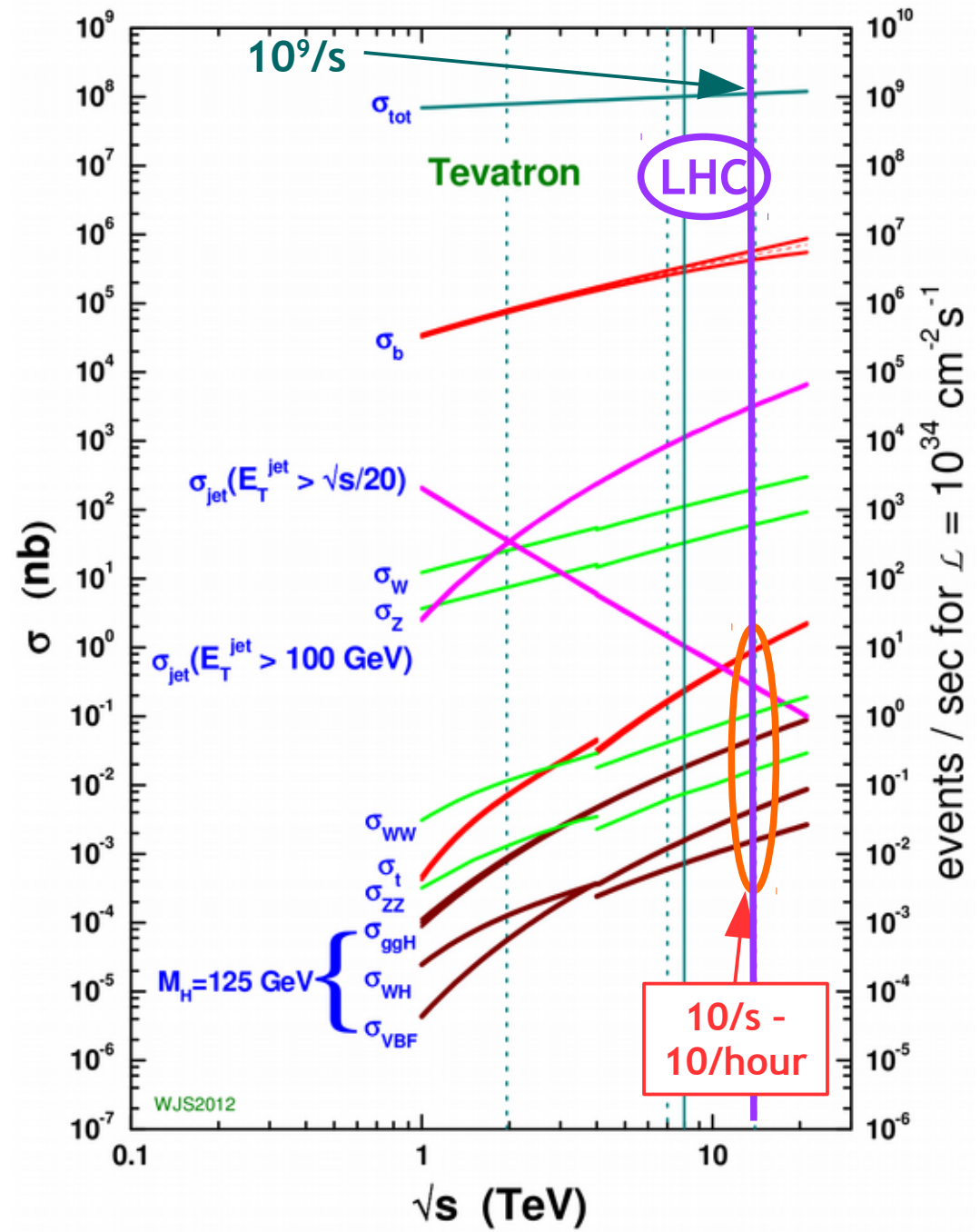
Cross-sections to produce massive particles such as the W, Z, t, (b,) H rise with \sqrt{s}

Range of cross-sections for processes studied, and so of their rates, from ~ 0.1 b to \sim fb i.e. factor $O(10^{14})$

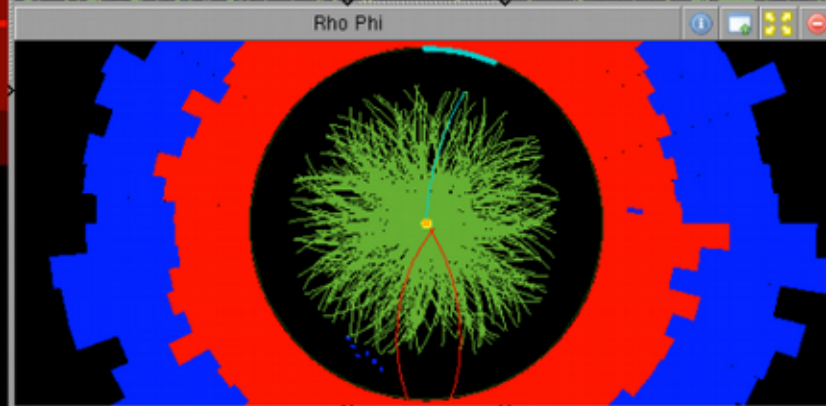
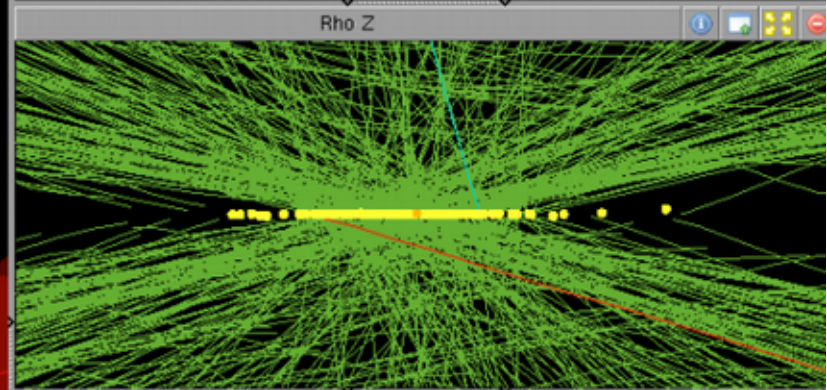
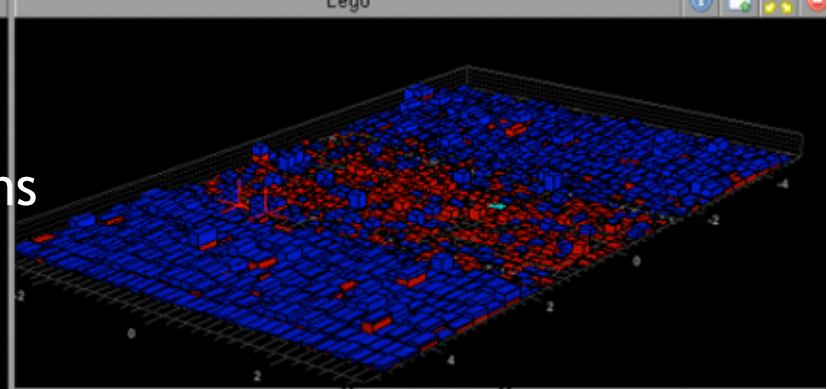
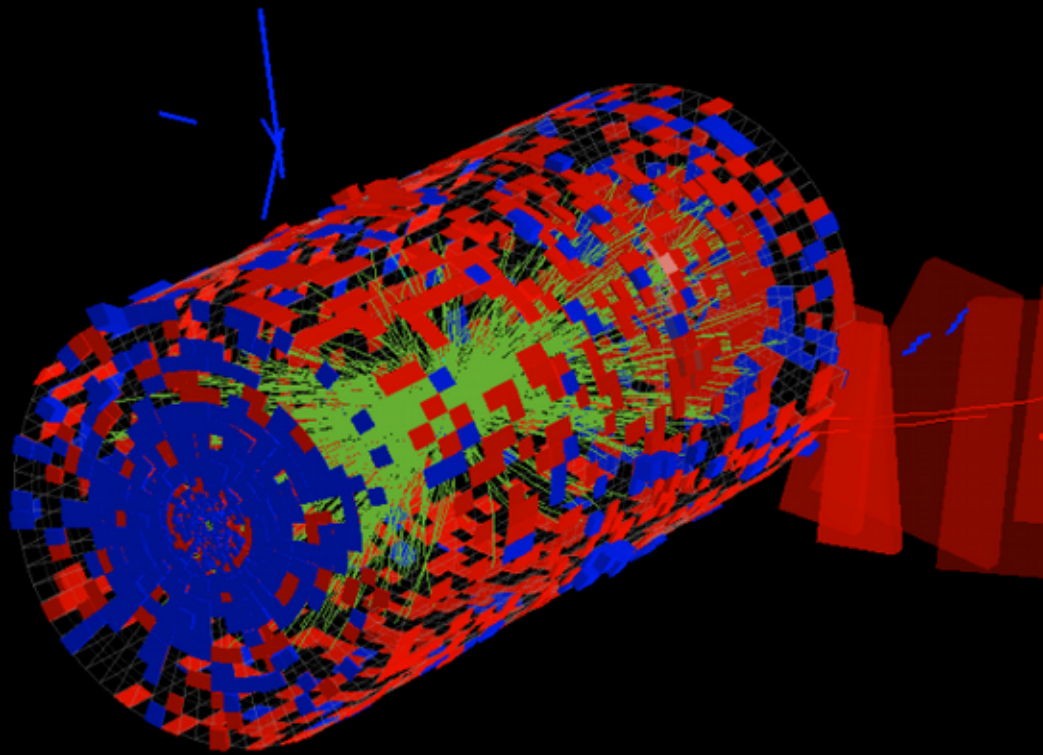
$\sim 10^9$ events per second occur in at most 30 million bunch crossings per second

→ 30+ events per bunch crossing

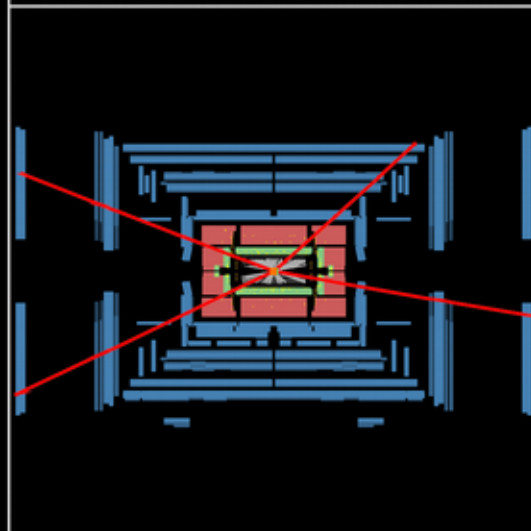
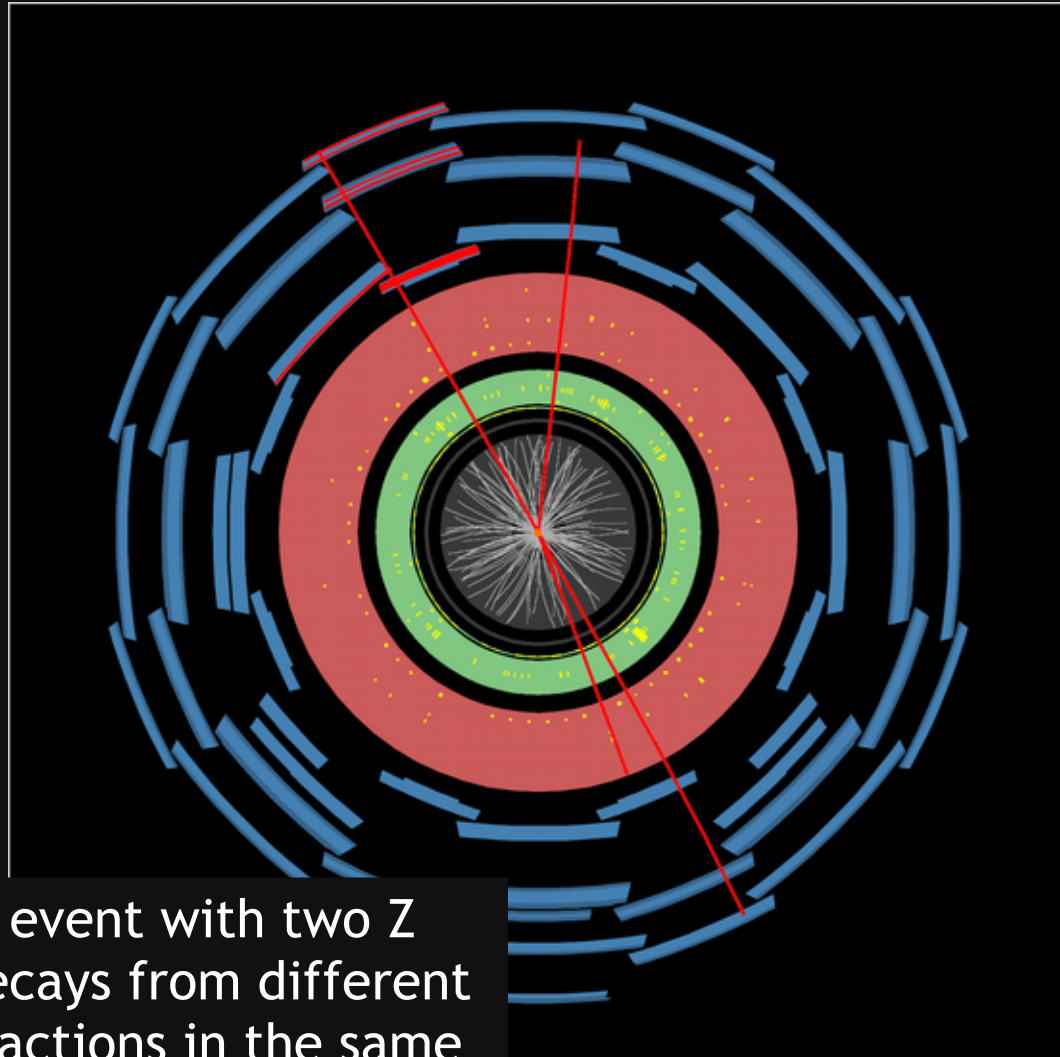
→ “pileup”



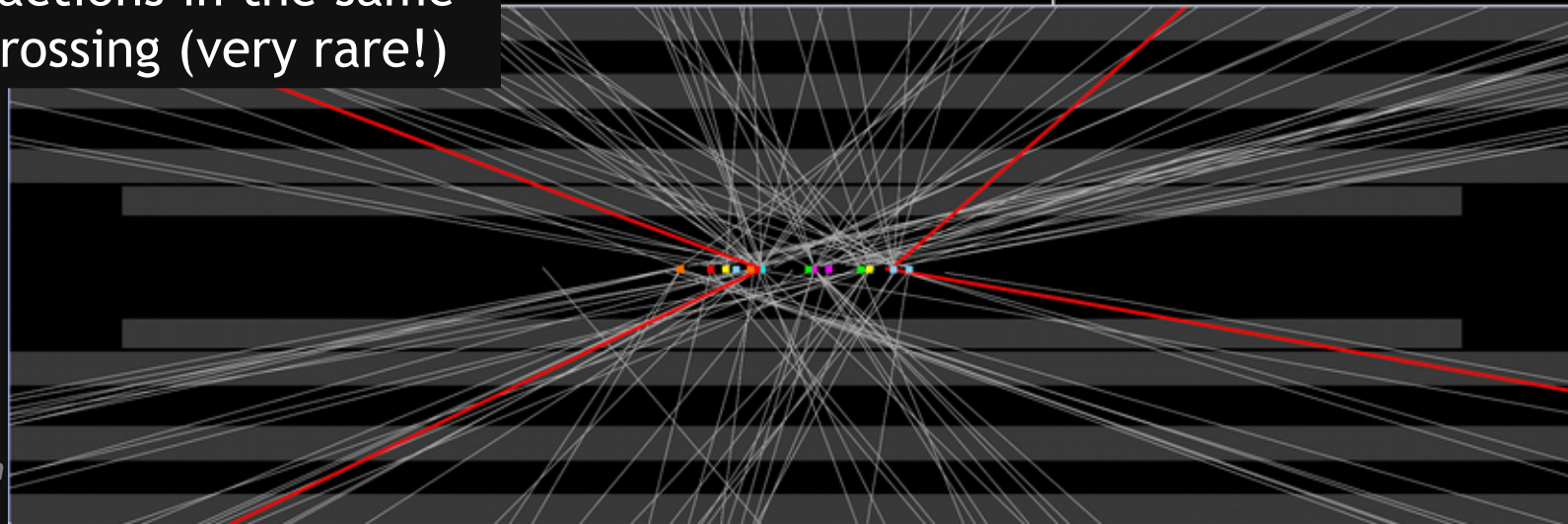
CMS event with 78 reconstructed *pile-up* interactions



Algorithm Name	Result	Bit Number	Prescale
L1_ETM50	0	67	1
L1_ETM70	0	68	1
L1_ETT300	1	69	1
L1_HTT100	0	70	262139
L1_HTT150	0	71	262139
L1_HTT175	1	72	1
L1_HTT200	1	73	1
L1_Mu10er_JetC12_WdEtaPhi1_DoubleJetC_20_12	0	74	262139
L1_Mu10er_JetC32	0	75	262139
L1_DoubleJet64_Central	1	76	1



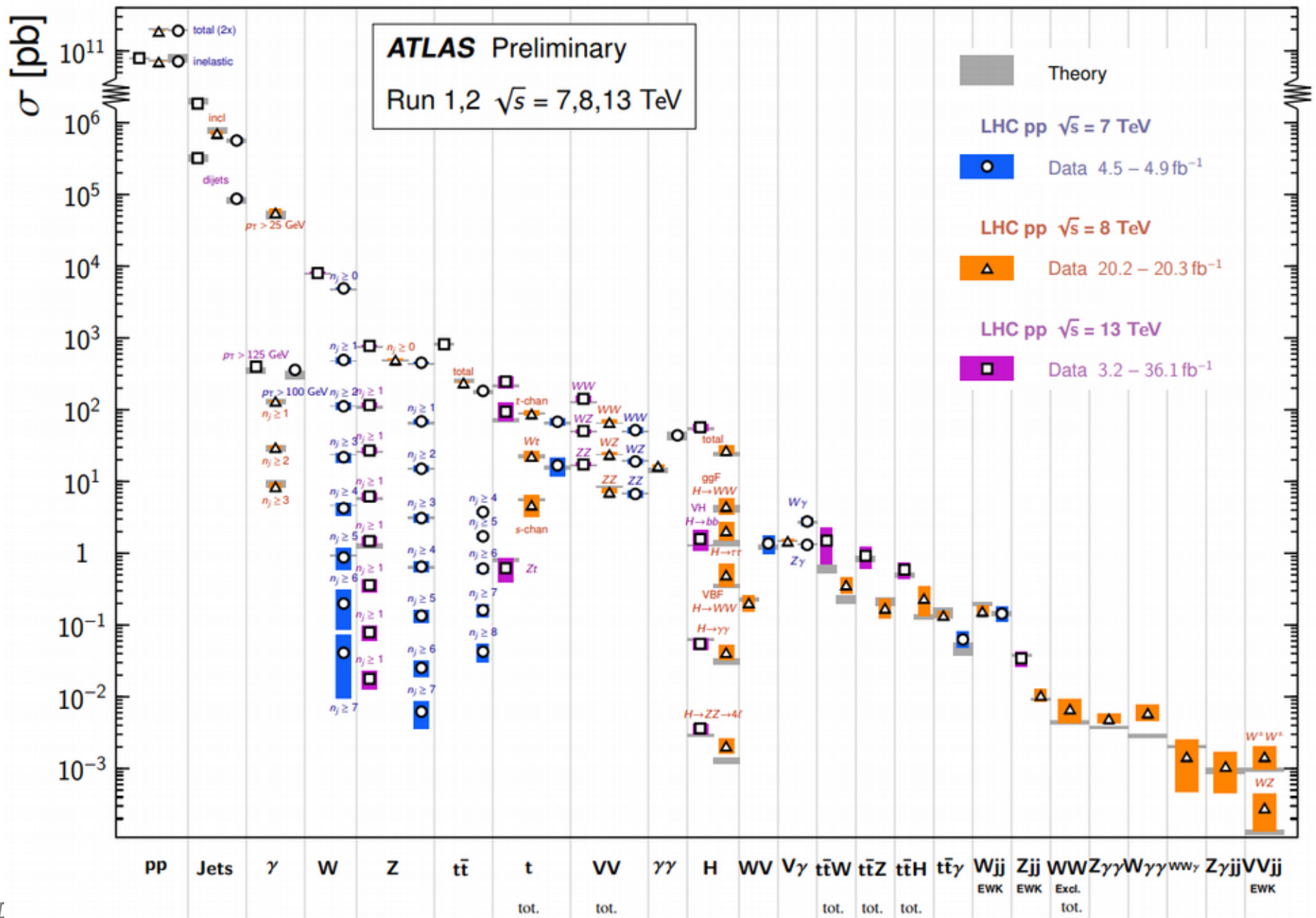
ATLAS event with two Z boson decays from different pp interactions in the same bunch crossing (very rare!)



Measured cross-sections

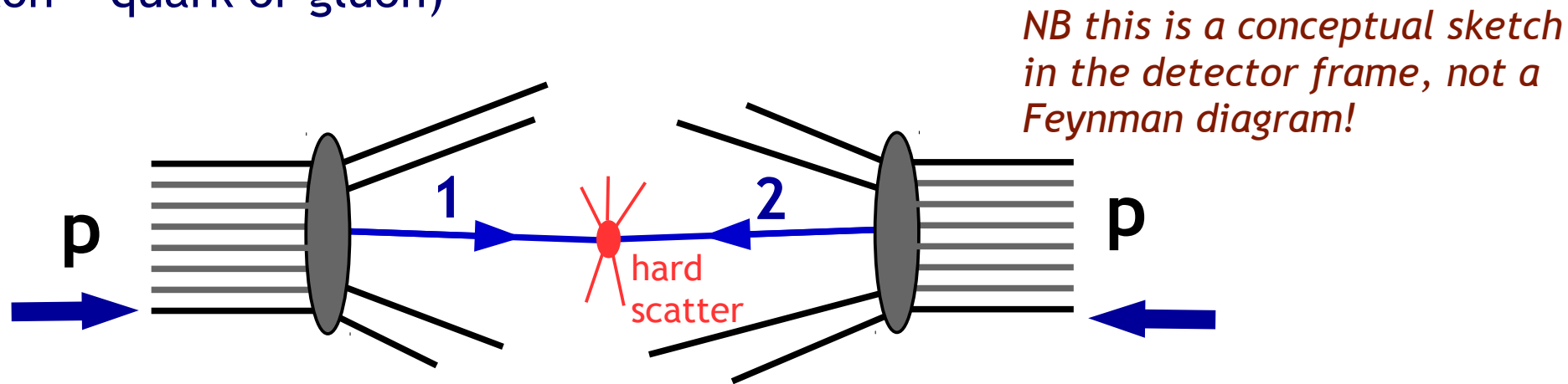
Standard Model Production Cross Section Measurements

Status: March 2018



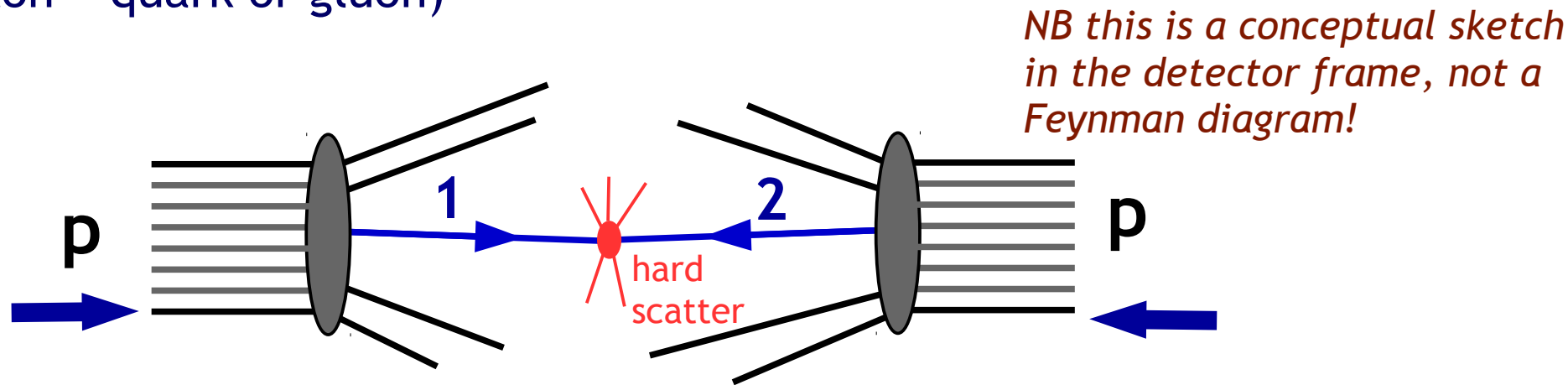
Predicting cross-sections

Although we collide protons in the experiments, at high energy we are really looking at high energy *parton-parton* collisions
(parton = quark or gluon)



Predicting cross-sections

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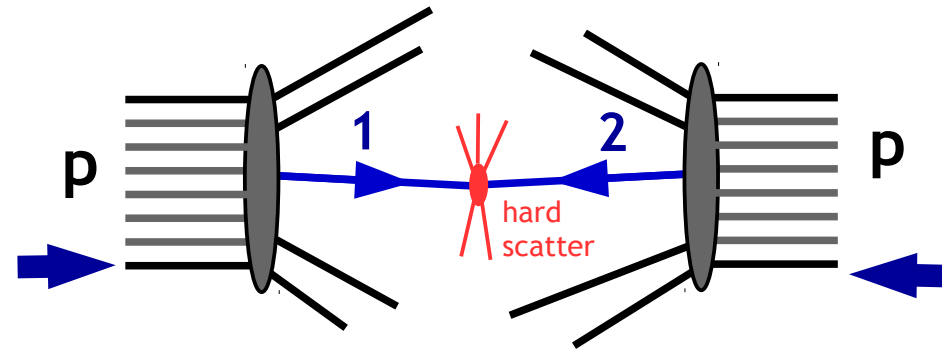
Partons 1 and 2 which collide in the *hard-scattering process* carry fractions x_1 and x_2 of the momentum of their original protons

Reduced (“effective”) centre-of-mass energy of the colliding partons is given by:

$$\sqrt{s_{12}} = \sqrt{(x_1 x_2 s)}$$

Predicting cross-sections (2)

To *predict* the cross-section for a given process, must know cross-section as a function of $\sqrt{s_{12}}$, and the parton density functions (pdfs) f ; then we have:



$$\sigma = \iint \hat{\sigma}(s_{12}) f_1(x_1, Q^2) f_2(x_2, Q^2) dx_1 dx_2$$

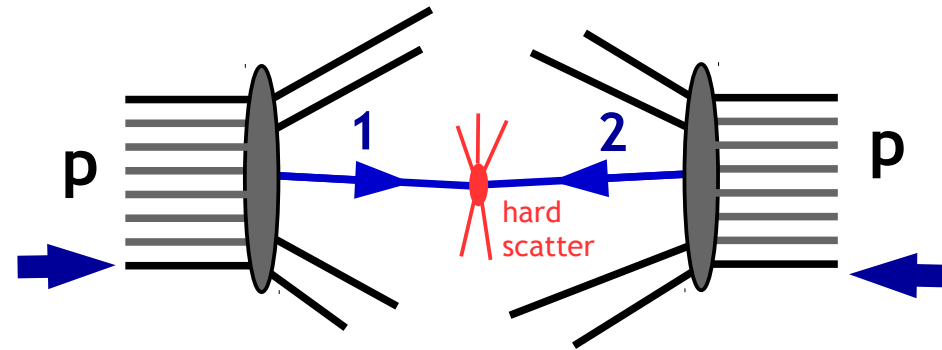
Theorists calculate this using Feynman diagrams and quantum field theory

We measure the pdfs at different experiments, and re-use them here

We measure this, and compare with the prediction

Predicting cross-sections (2)

To *predict* the cross-section for a given process, must know cross-section as a function of $\sqrt{s_{12}}$, and the parton density functions (pdfs) f ; then we have:



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We measure the *total cross-section* σ , or more usually a *fiducial cross-section* σ^{fid} , which is the part of the total cross-section with the final-state particles from the hard-scattering process going into well-defined regions of phase-space (angle, momentum), measurable in the detector

We also measure *differential cross-sections*, which are typically a more finely divided (binned) set of fiducial cross-sections, e.g. we may measure

$$d\sigma/dp_T \quad \text{or} \quad d\sigma/d\eta \quad \text{or} \quad \sigma(N_{\text{jet}})$$

for a specified final-state particle or jet

Parton density functions

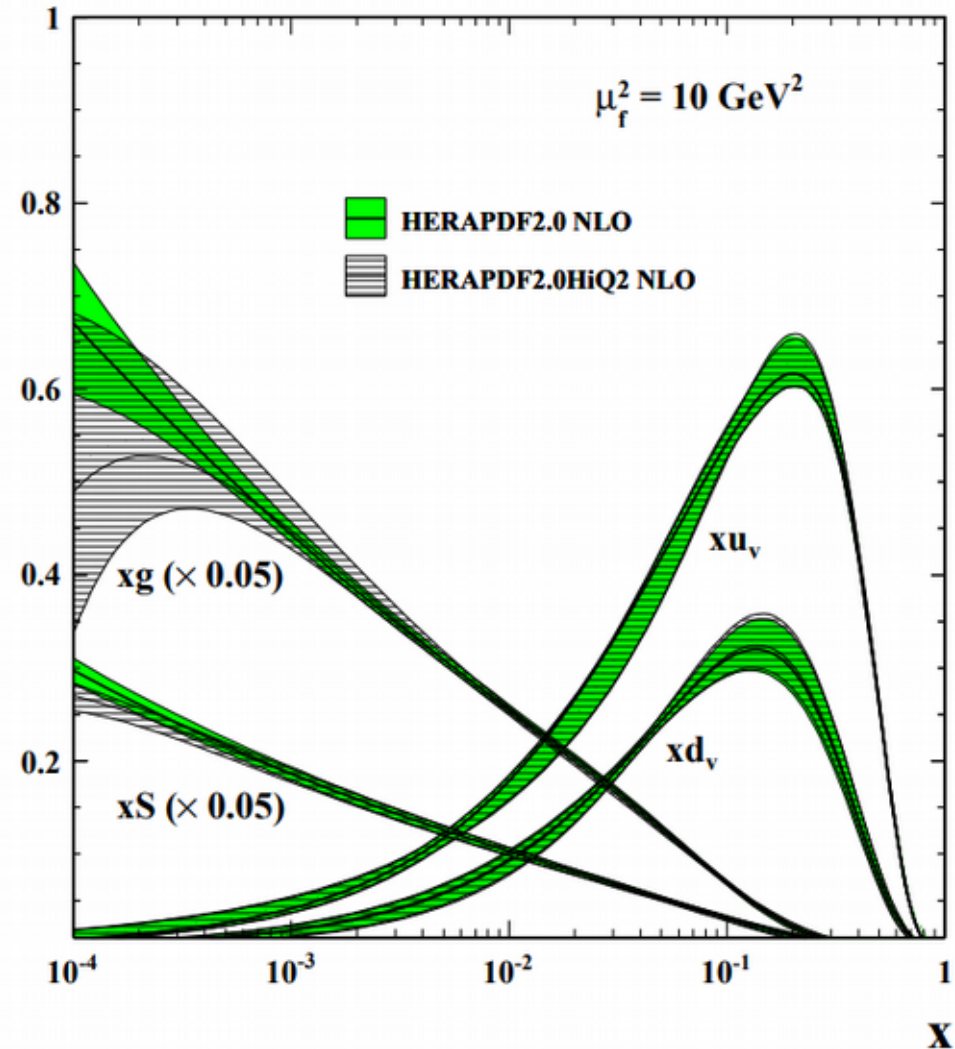
Typical parton density functions

Measured in previous experiments (HERA, Tevatron colliders ...), and we refine them using LHC data

I've been ignoring Q^2 ($\sim \mu_f^2$ on the plot) so far - this is important, it characterises the momentum-scale (squared) of the hard scattering process

xf

H1 and ZEUS



Parton density functions

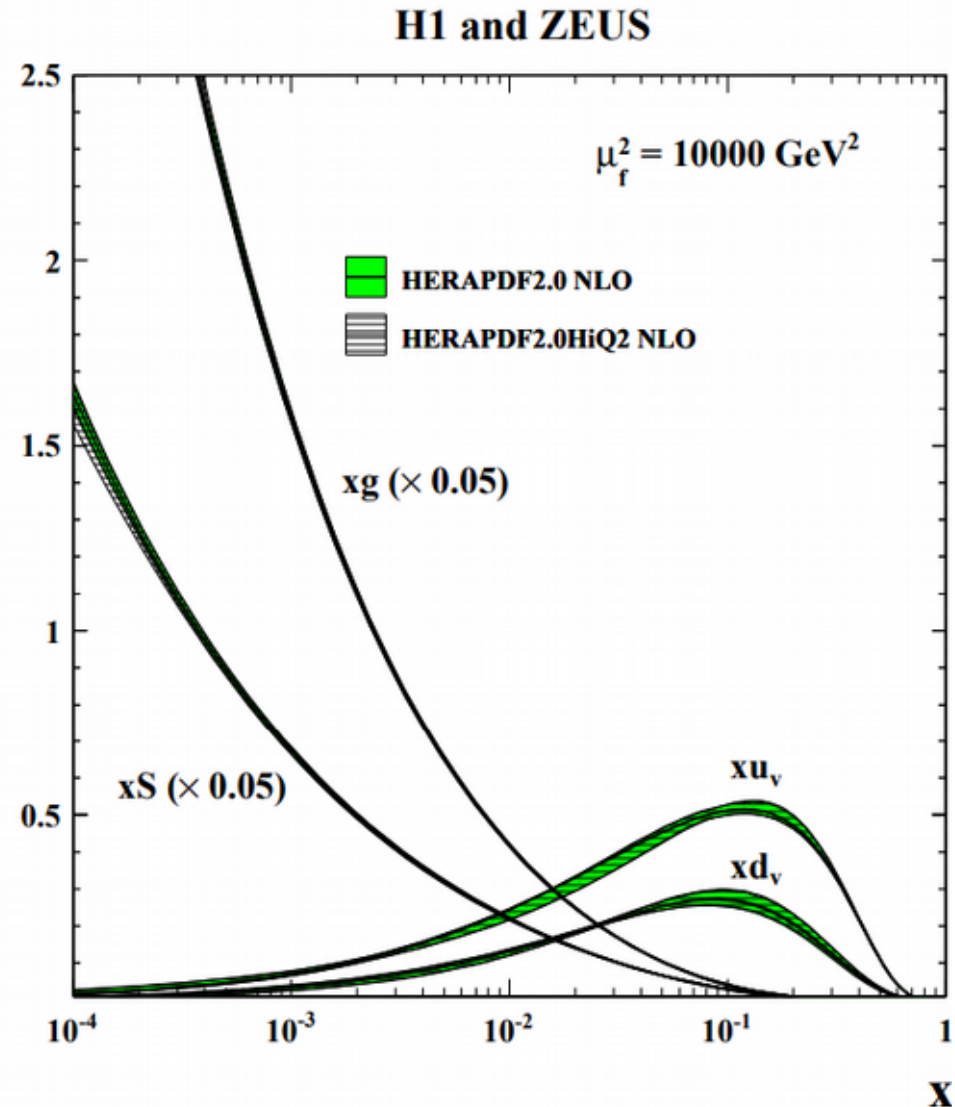
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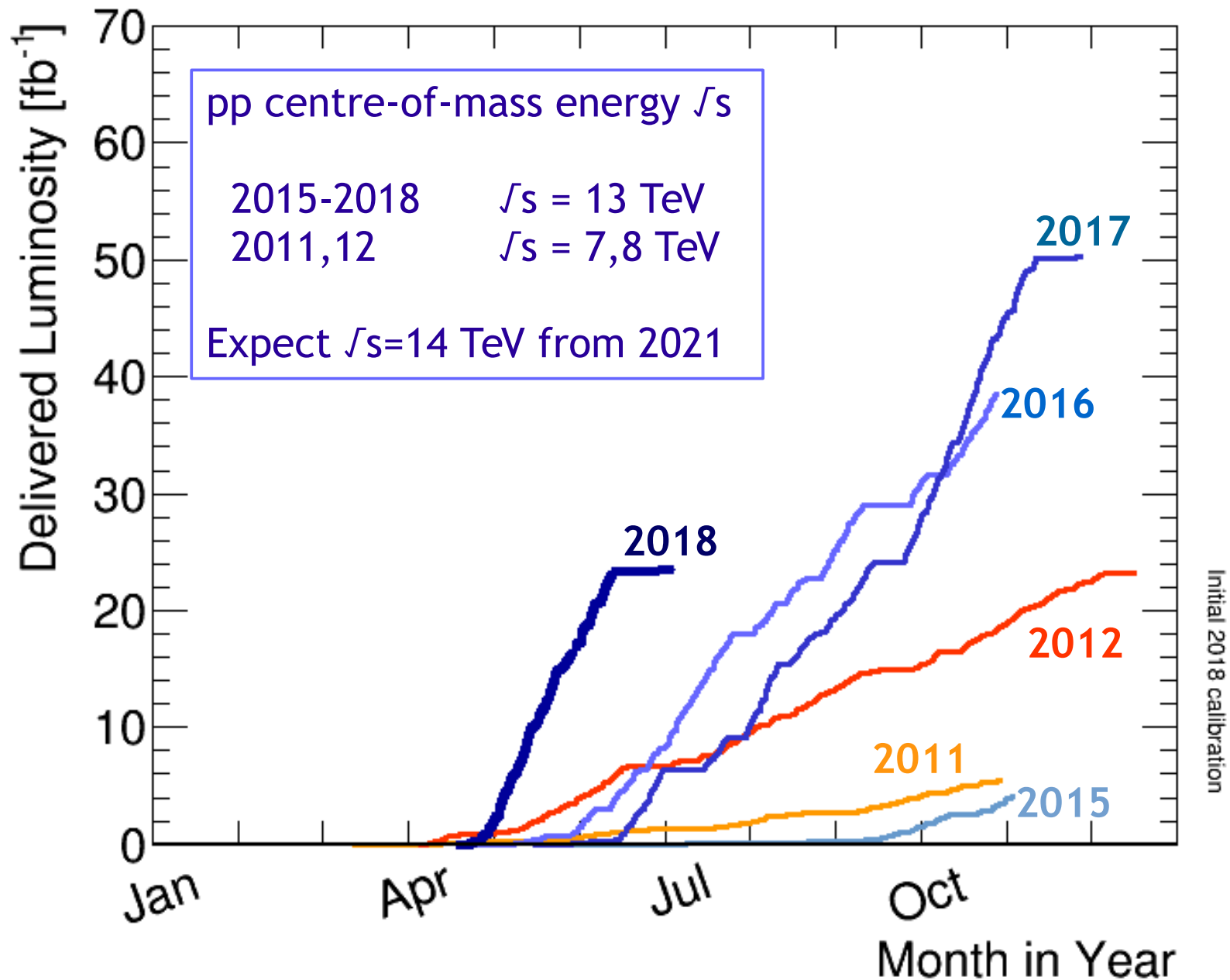
I've been ignoring Q^2 ($\sim \mu_f^2$ on the plot) so far - this is important, it characterises the momentum-scale (squared) of the hard scattering process

pdfs *evolve* with Q^2 , but in a predictable way ("DGLAP")

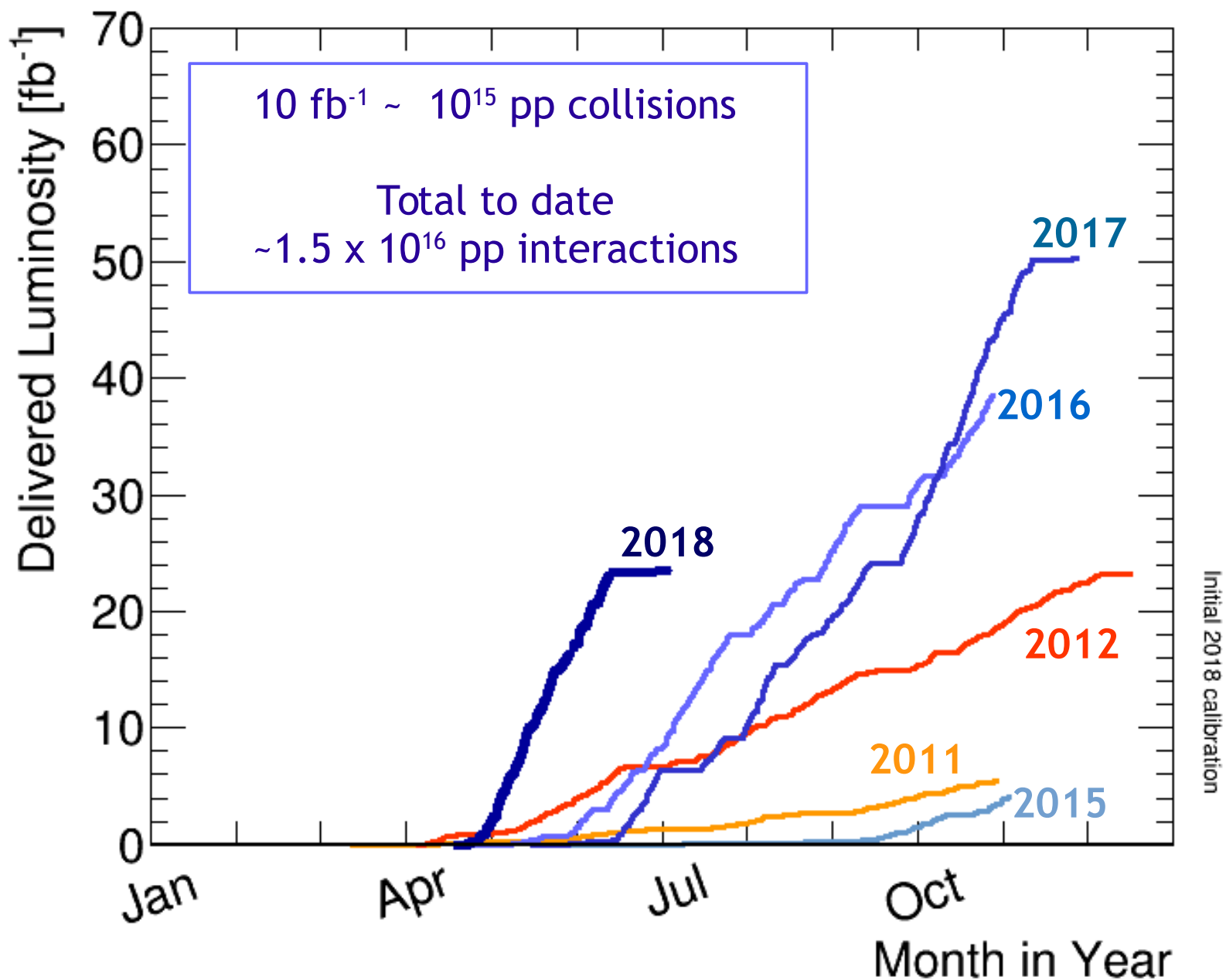
xf



Back to the LHC: pp data samples



LHC pp data samples



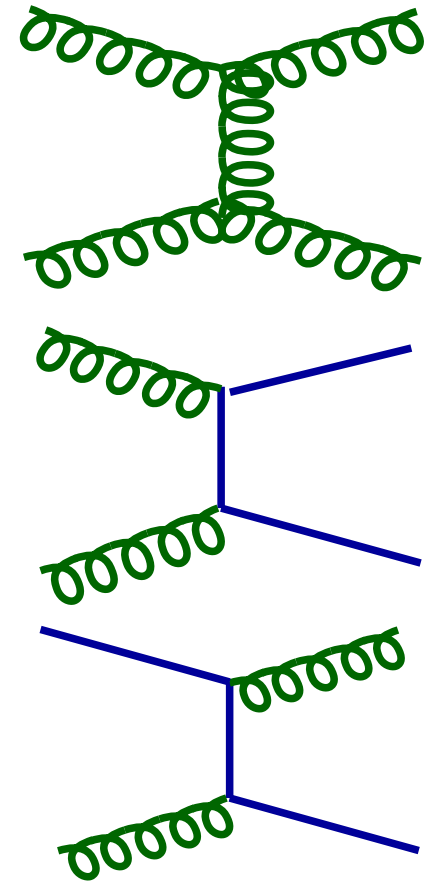
Measuring hadronic jets

Production of hadronic jets has a very high cross-section:

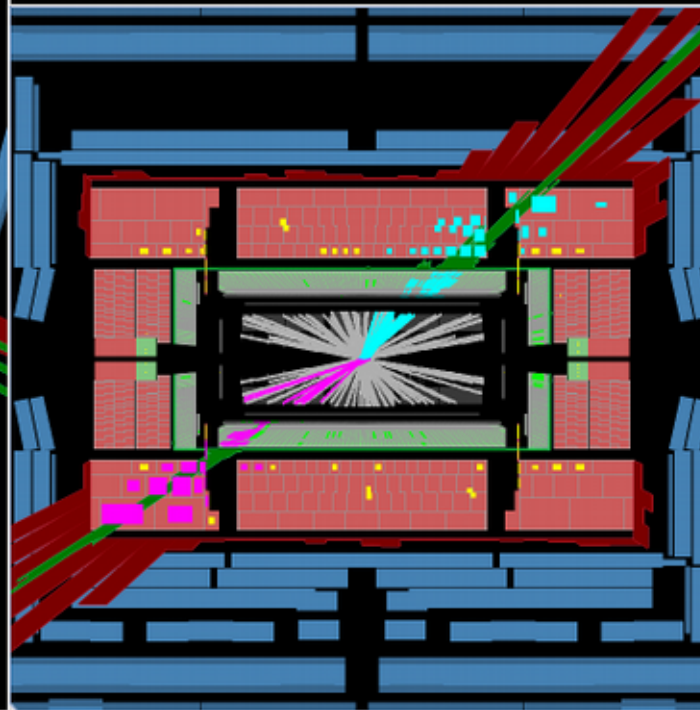
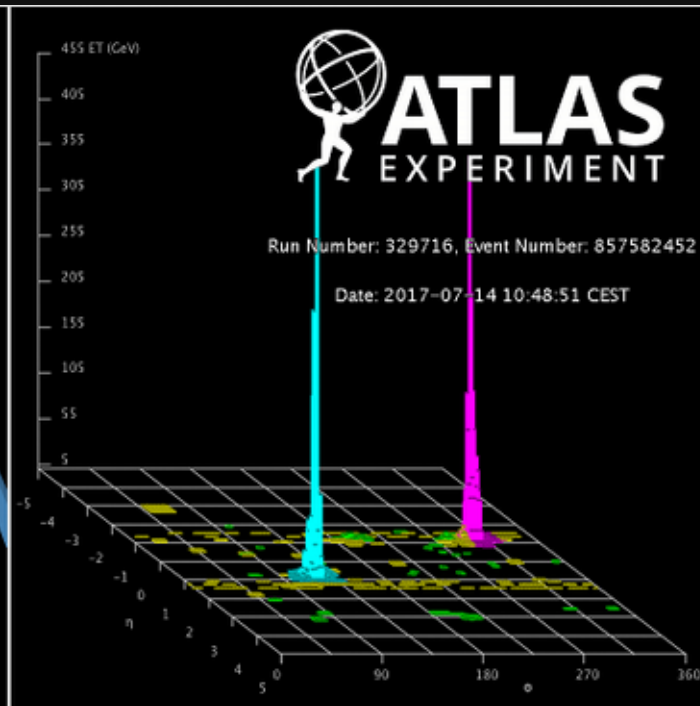
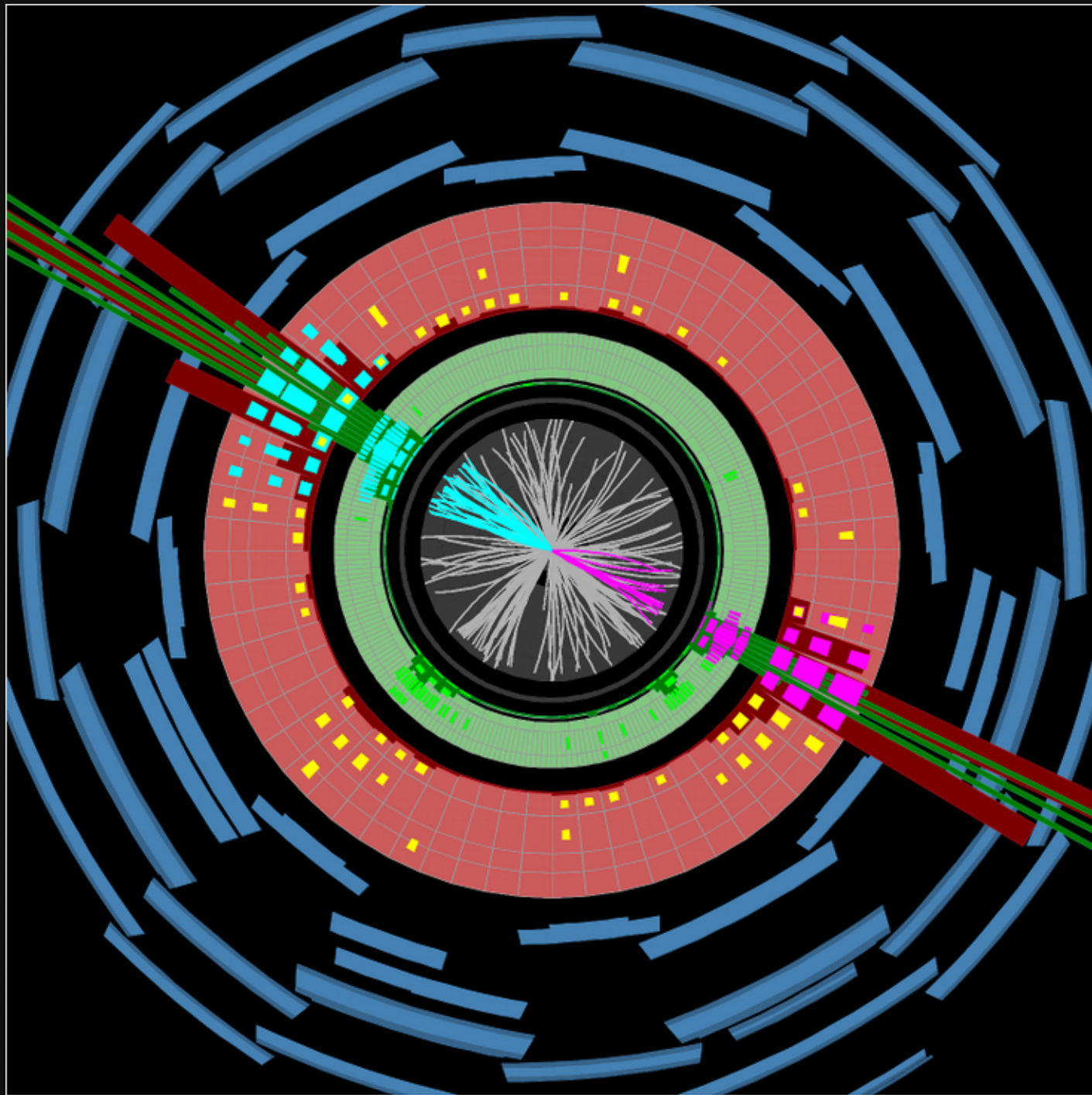
- Strong interaction (QCD) process
- Example parton-level Feynman diagrams, at lowest-order (fewest number of vertices)
- Outgoing partons turn into collimated jets of hadrons within ~ 1 fm



- Reconstruct these "hadronic jets" in the detector to discover the parton configuration using "jet algorithms" (*our favourite one is called "anti- k_t "*)



A high-mass dijet event, $m(jj)=9.3$ TeV

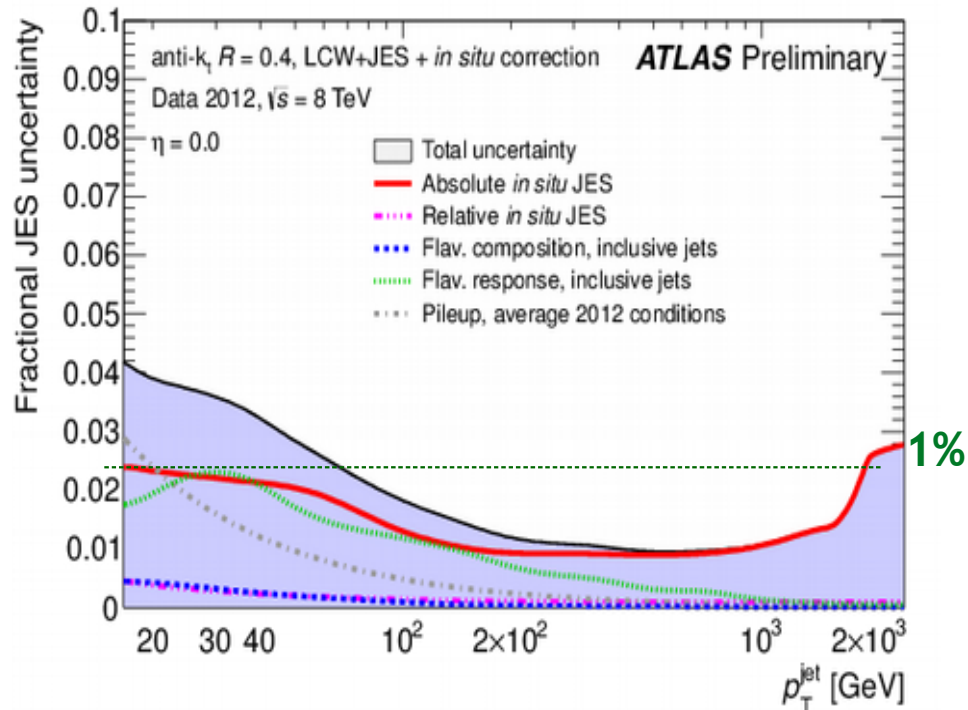


Jet cross-sections

Needs best understanding we can get of the energy of jets

→ Jet energy scale (“JES”) uncertainty

This takes a long time (years), lots of careful work

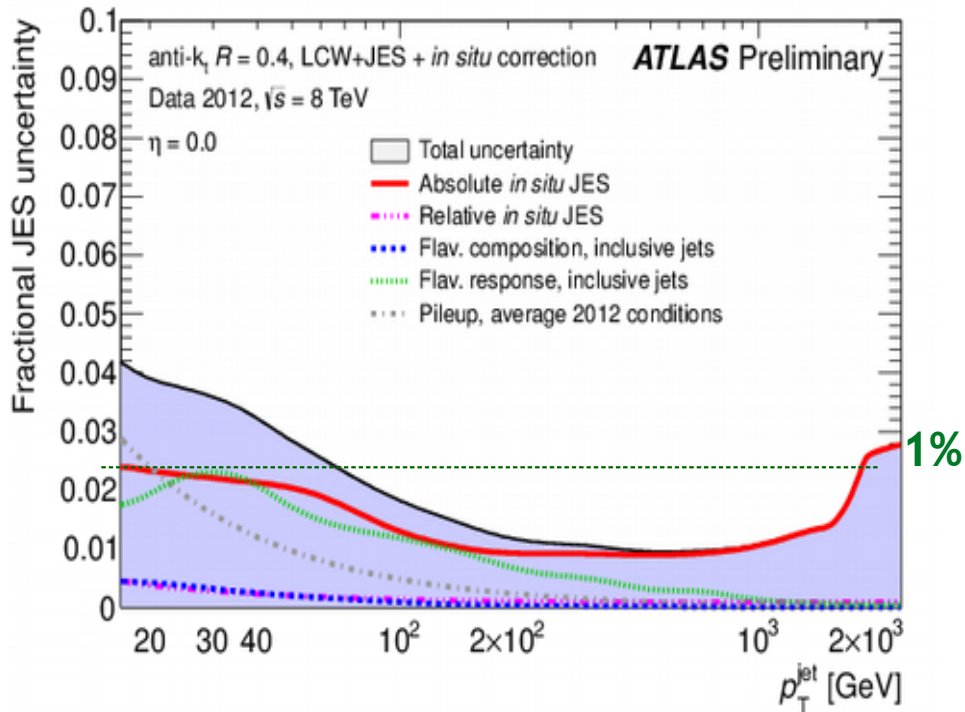


Jet energy scale uncertainty

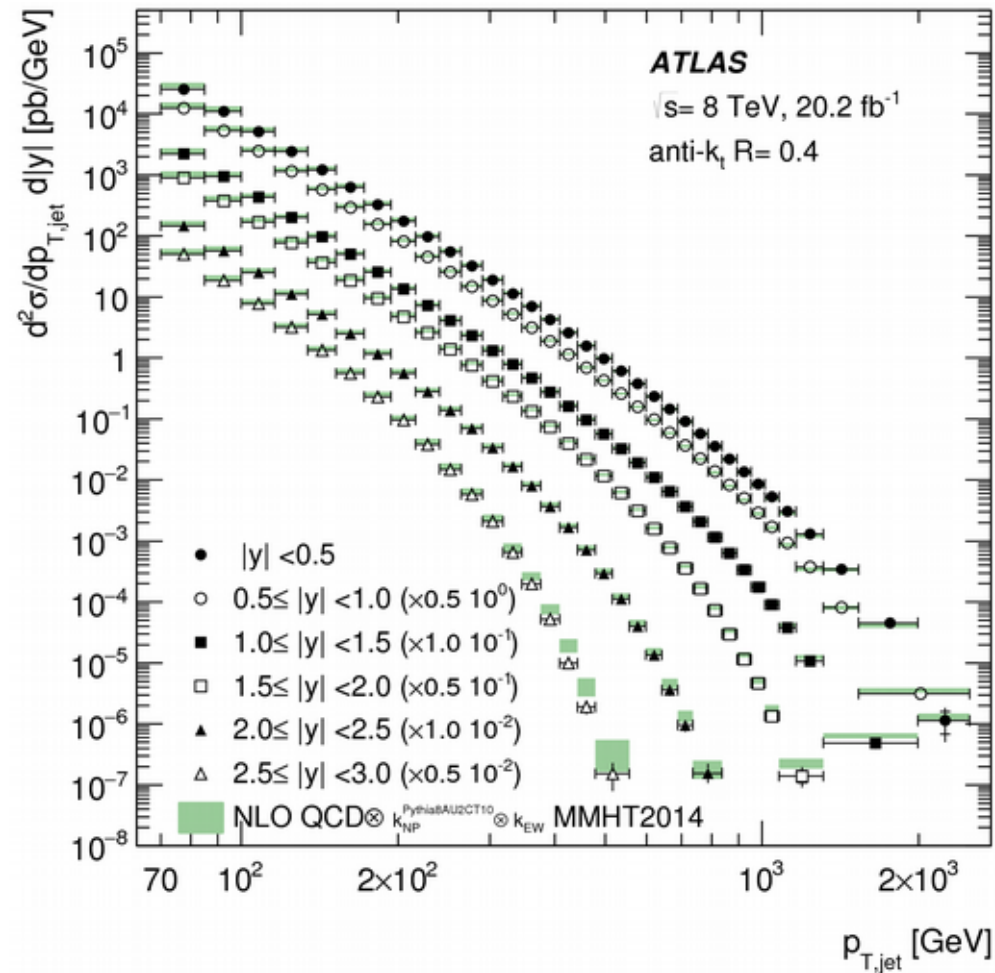
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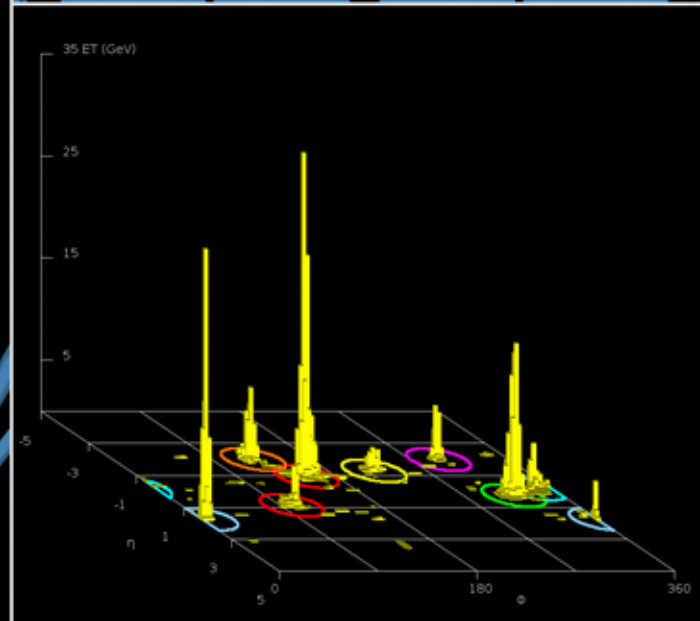
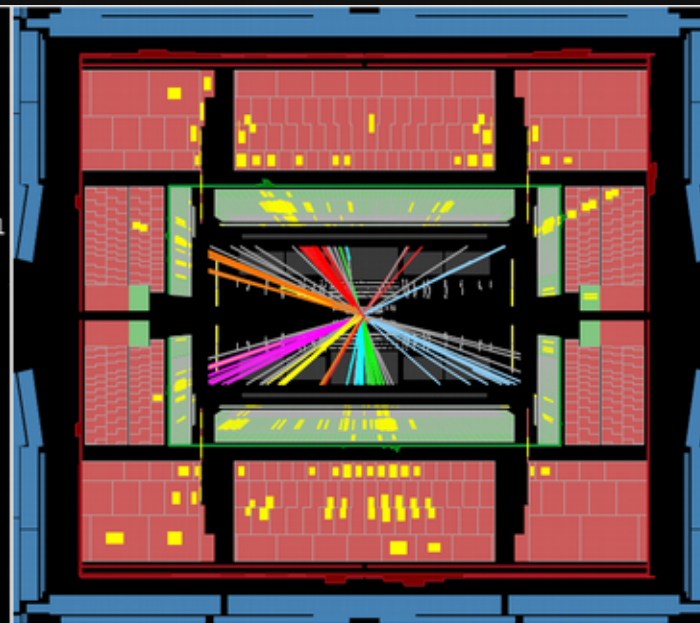
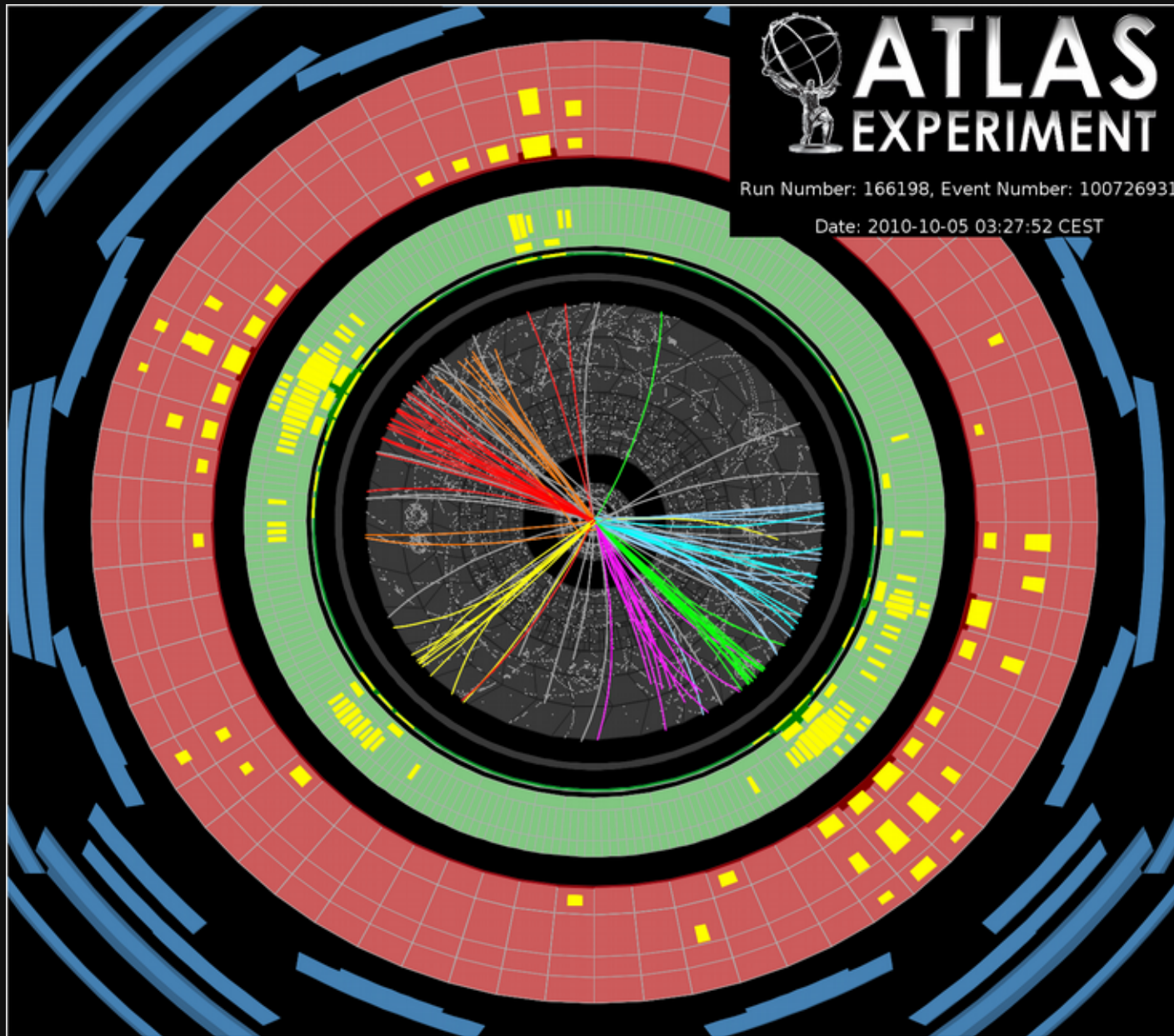
Jet energy scale uncertainty



Single jet cross-section vs jet p_T - spans 10 orders of magnitude

These data will be used to improve the parton density function (pdf) models

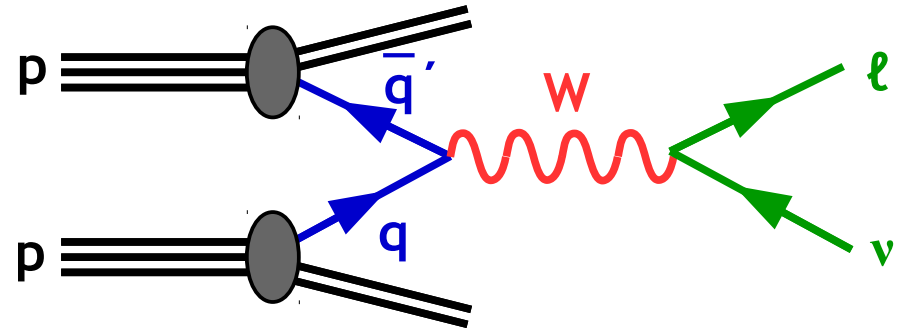
An 8-jet event, with each $p_T(\text{jet}) > 60$ GeV



Measurements of W and Z bosons

Clean experimental signatures and large cross-sections

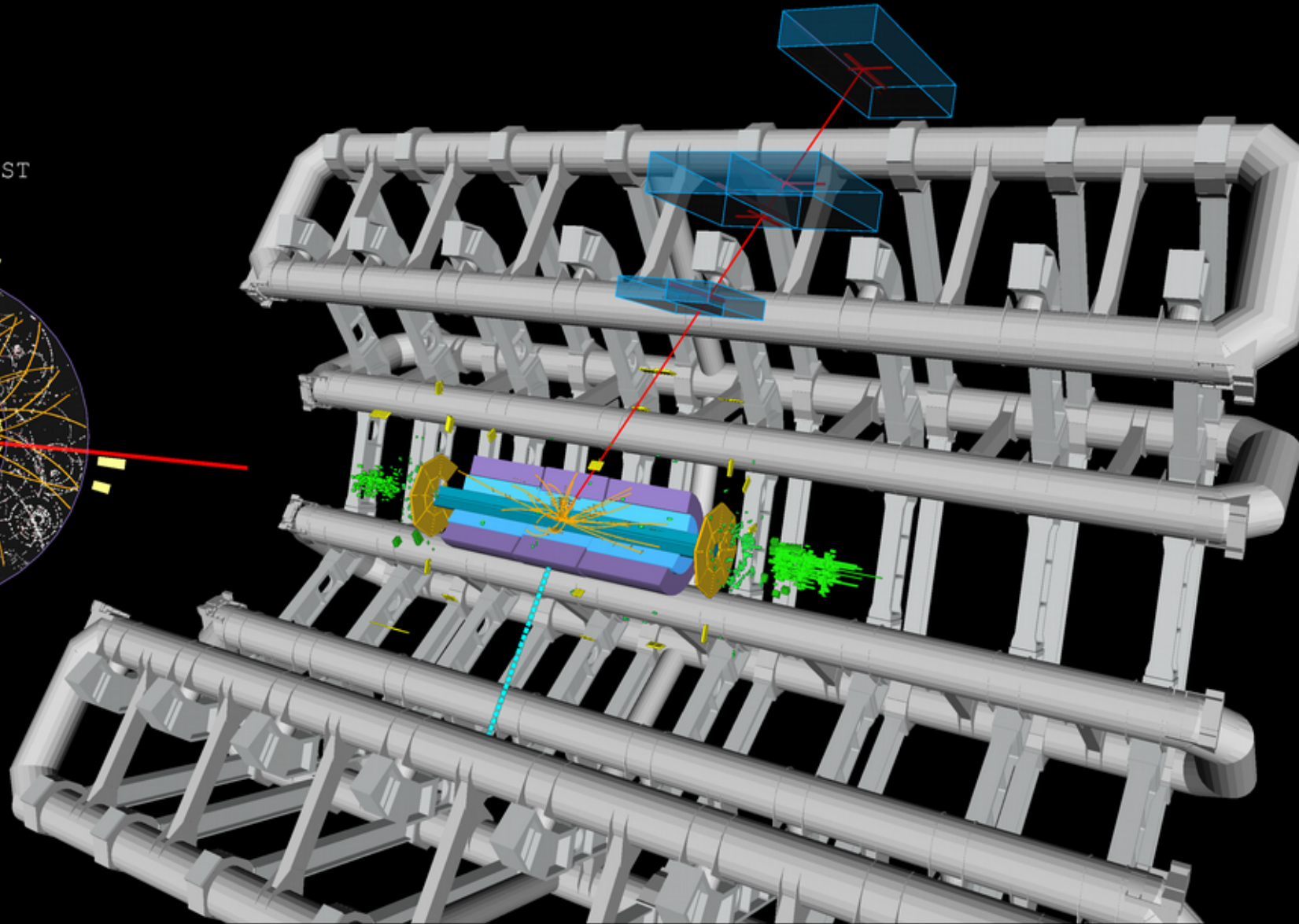
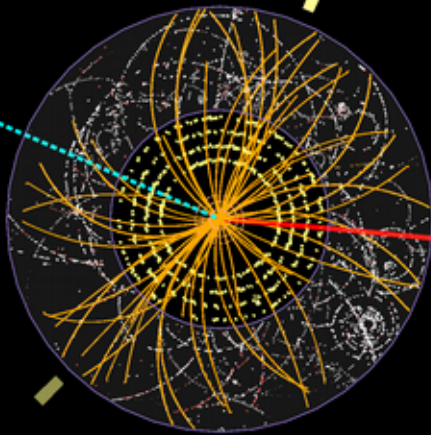
- High precision measurements
- Strong constraints on proton structure
- Tests of consistency of electroweak (EW) sector of SM



$W \rightarrow \mu\nu$ event



Run: 183081
Event: 101291517
2011-06-05 17:09:02 CEST



$M_T = 82.9$ GeV
 p_T muon = 32.8 GeV
 $E_T^{miss} = 52.4$ GeV

$Z \rightarrow ee$ event

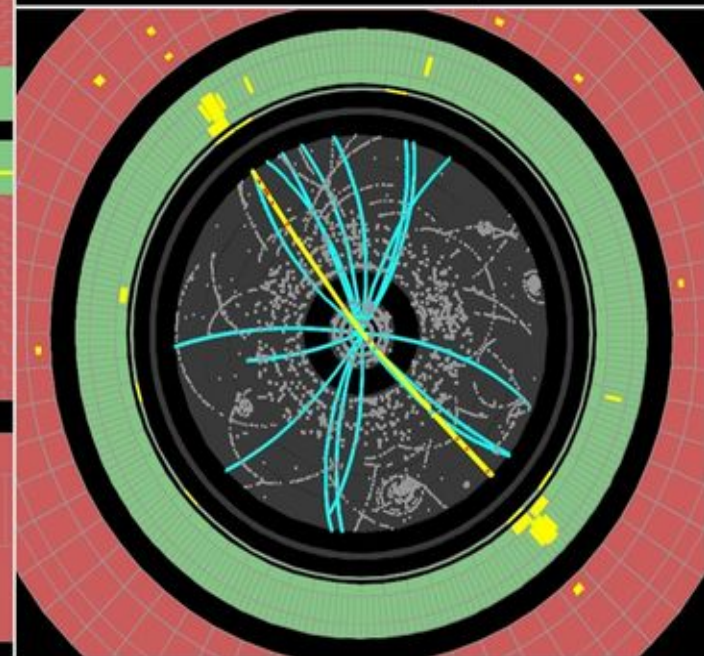
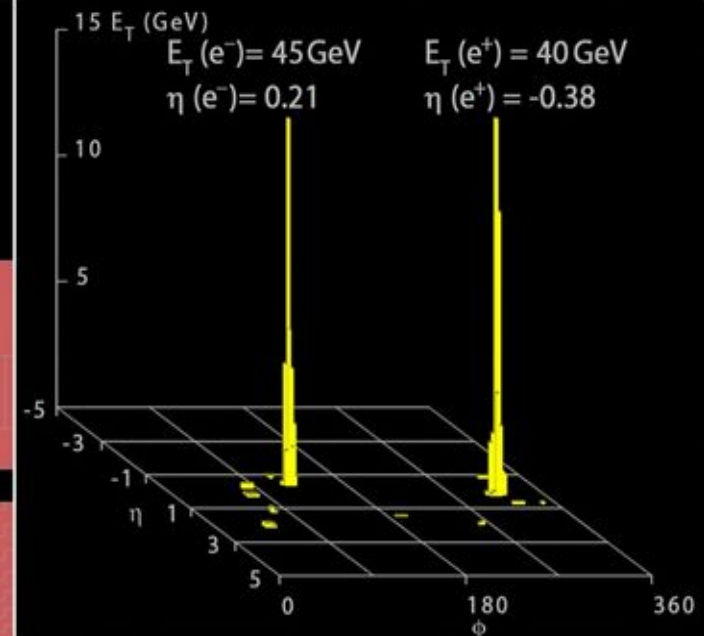


Run Number: 154817, Event Number: 968871

Date: 2010-05-09 09:41:40 CEST

$M_{ee} = 89 \text{ GeV}$

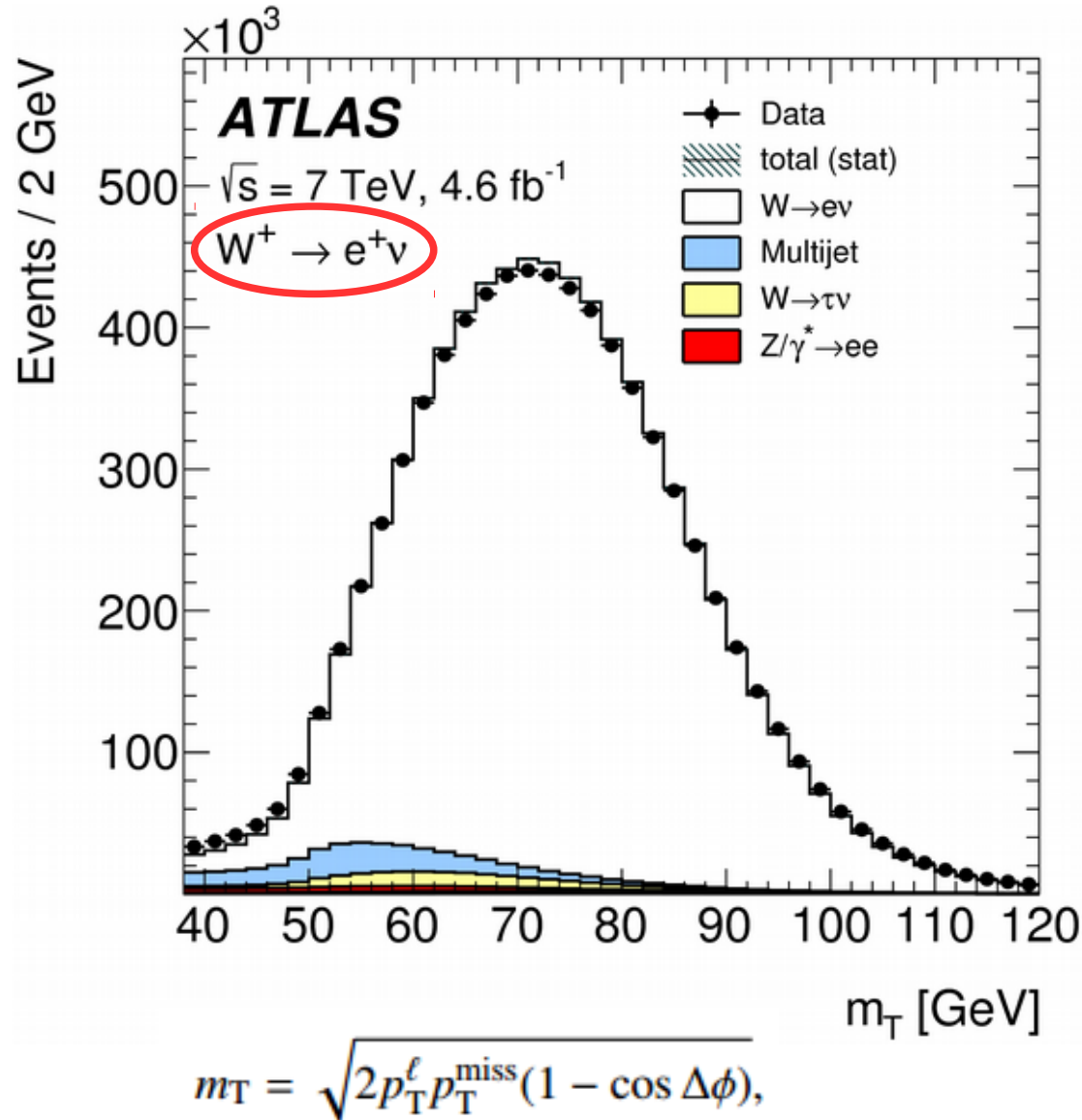
$Z \rightarrow ee$ candidate in 7 TeV collisions



Precise W, Z production measurements

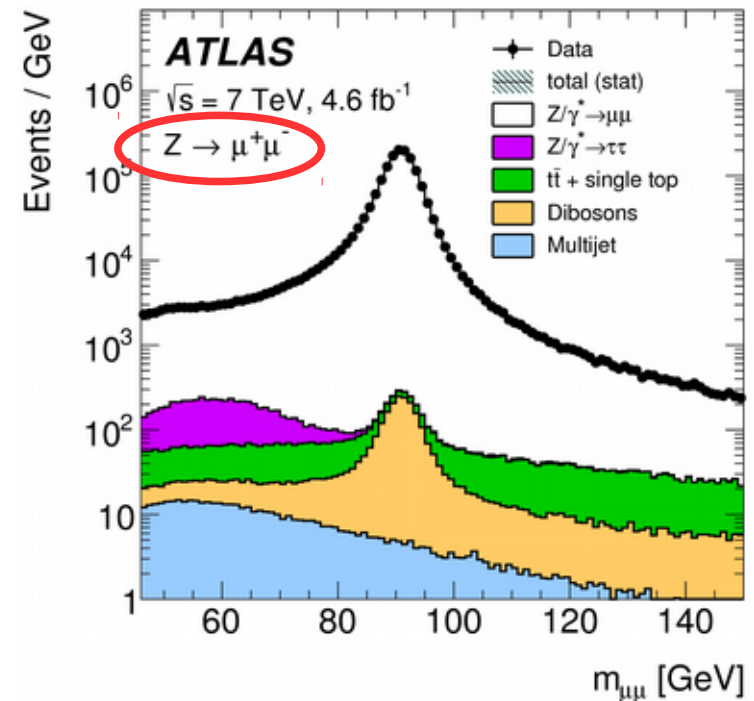
Detailed studies performed with 2011 data at 7 TeV: W^+ , W^- , Z in e, μ decays

arXiv:1612.03016



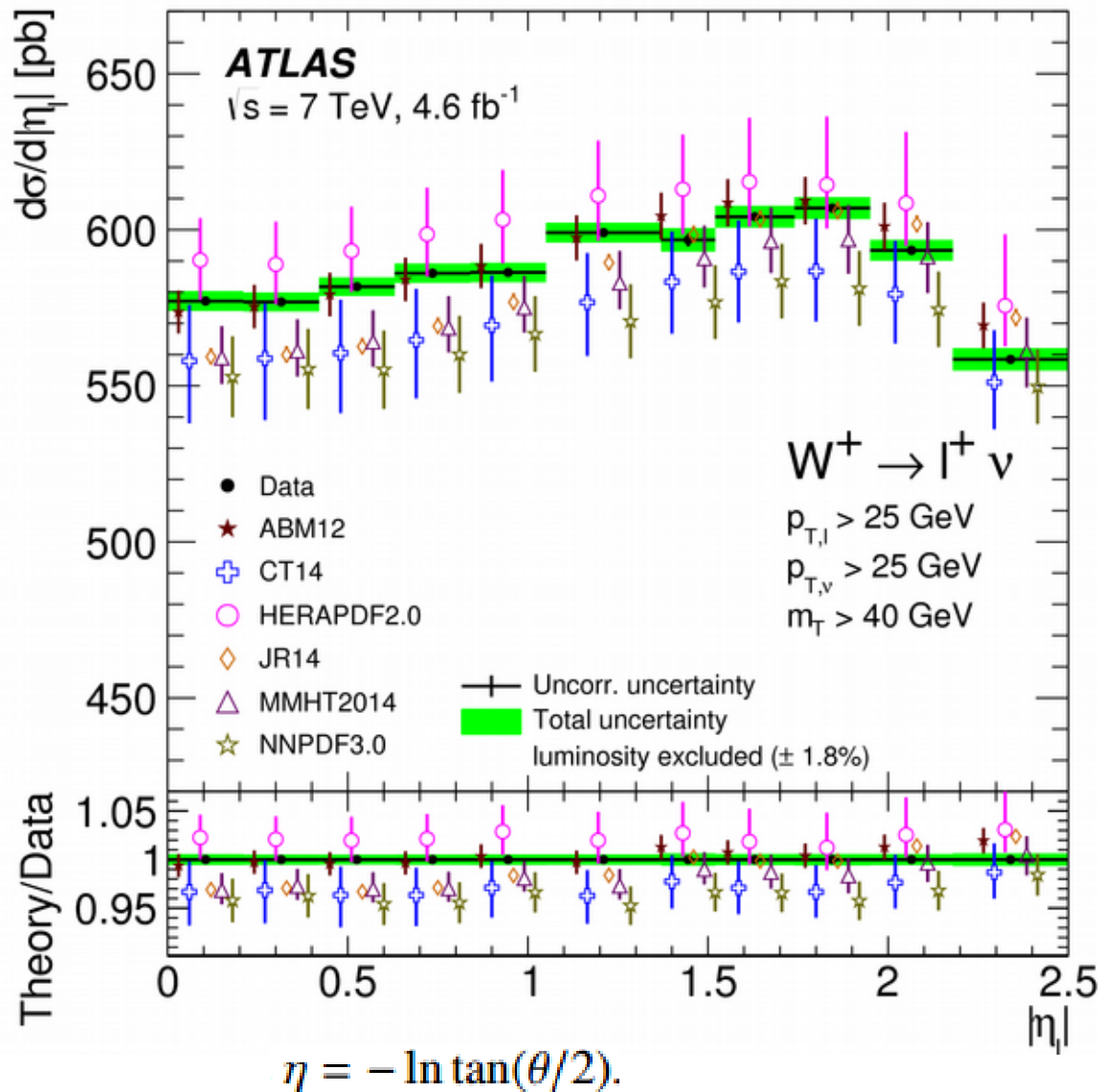
High statistics data well described by simulation

Backgrounds under excellent control



Precise W, Z production measurements

Detailed studies performed with 2011 data at 7 TeV: W^+, W^-, Z in e, μ decays



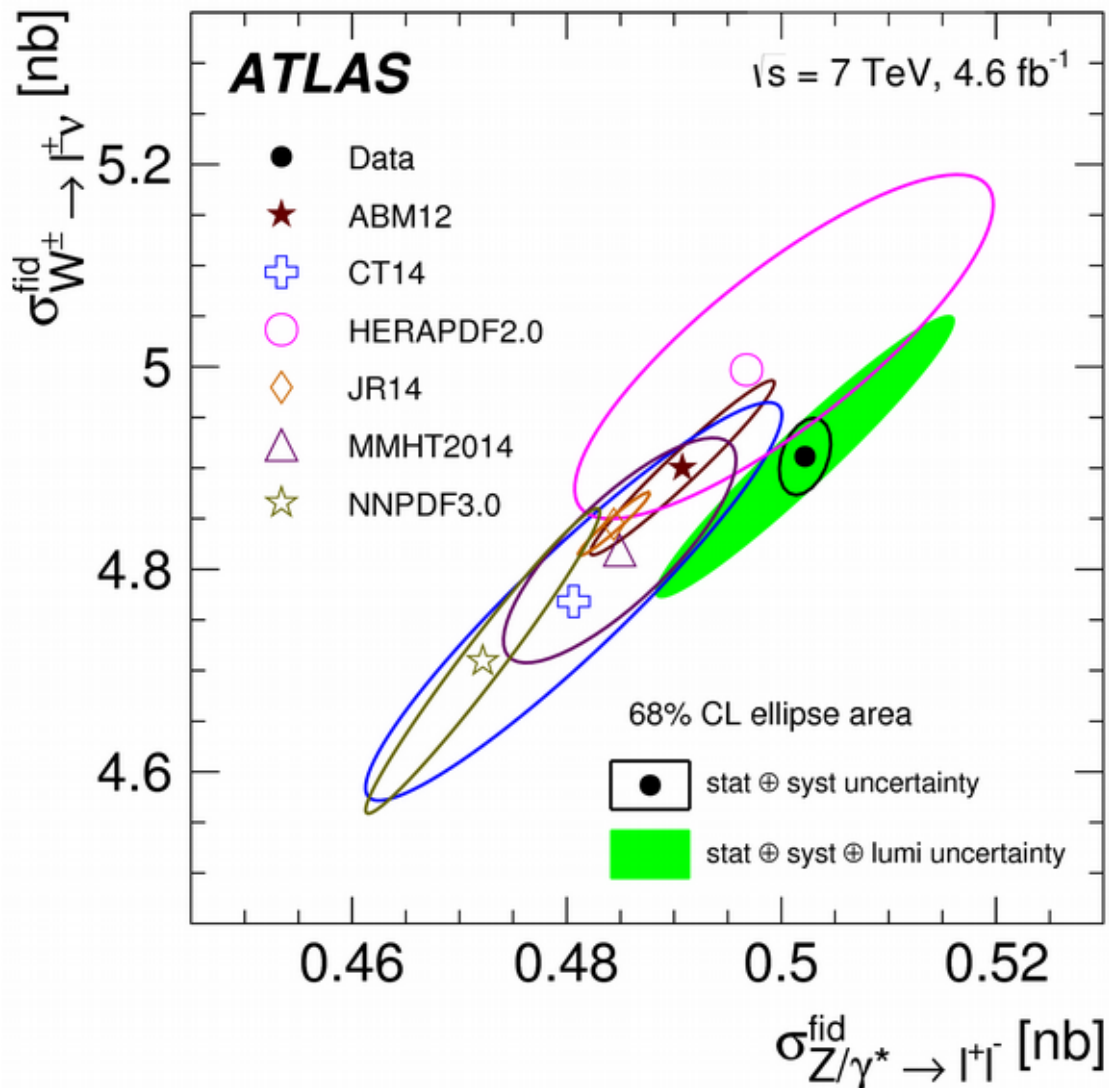
Example: measurement of angular distributions of leptons relative to beam direction, in $W \rightarrow \ell \nu$ decays

Green errors are from the data - errors on predictions from different proton structure (pdf) sets much larger

→ much sensitivity to the pdfs

Precise W, Z production measurements

Detailed studies performed with 2011 data at 7 TeV: W^+, W^-, Z in e, μ decays



Experimental errors better than theoretical/modelling uncertainties

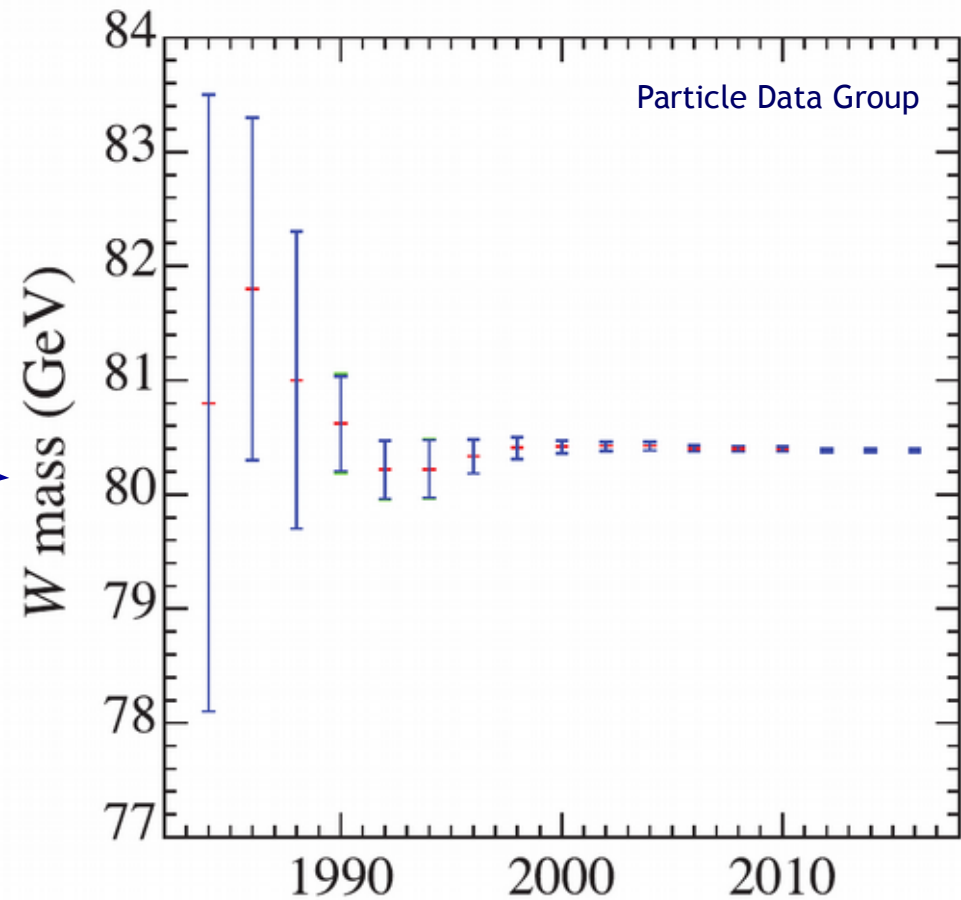
These data are used by ATLAS to make new pdfs

Measuring the W mass

Mass of the W boson is a fundamental parameter of the Standard Model

W mass was first measured directly back in the 1980's soon after it was discovered at CERN

- History of precision



Measuring the W mass

Mass of the W boson is a fundamental parameter of the Standard Model

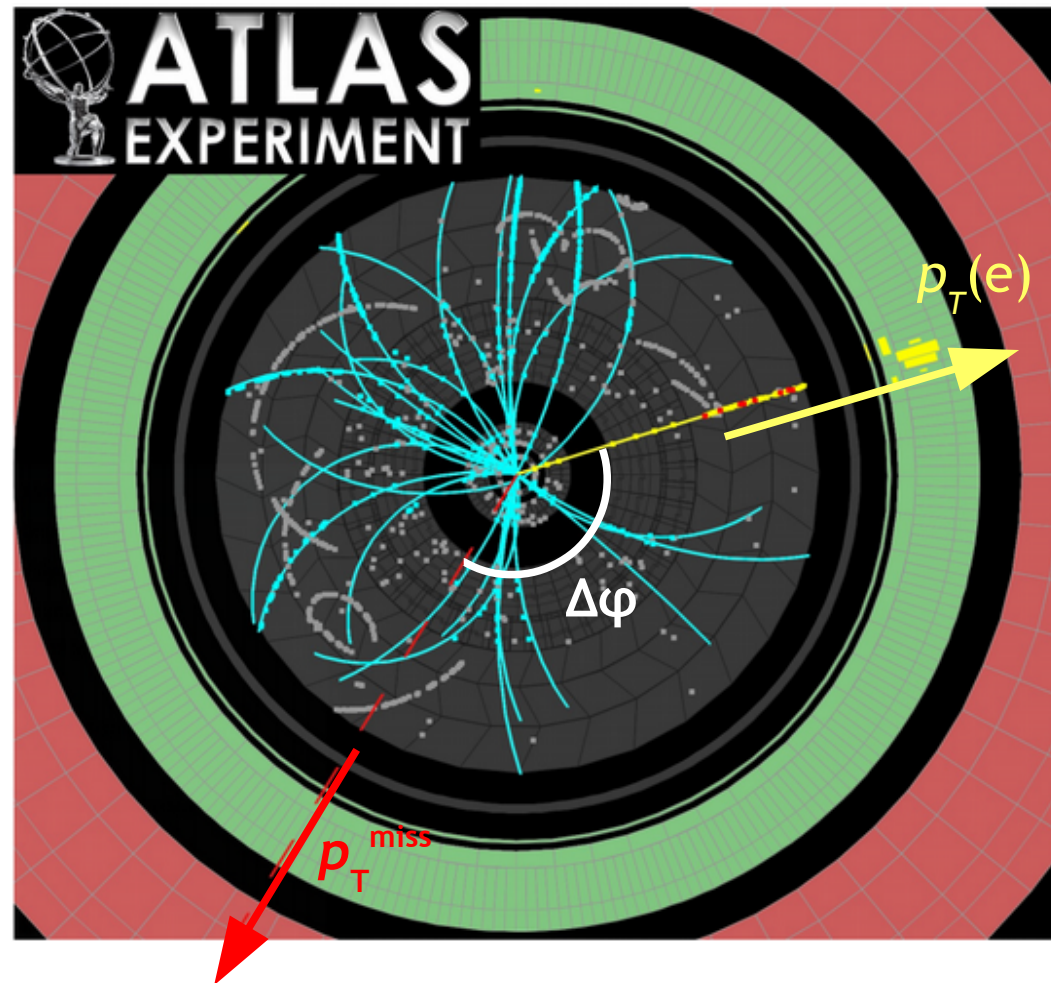
W mass was first measured directly back in the 1980's soon after it was discovered at CERN

- History of precision

A standard method uses “transverse mass”

$$m_T = \sqrt{2p_T^\ell p_T^{\text{miss}}(1 - \cos \Delta\phi)},$$

W → eν event



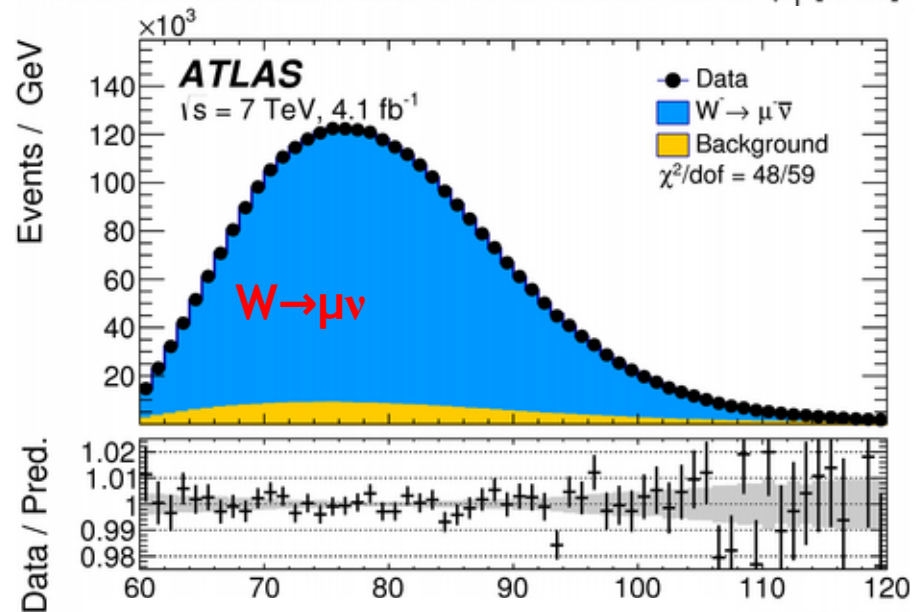
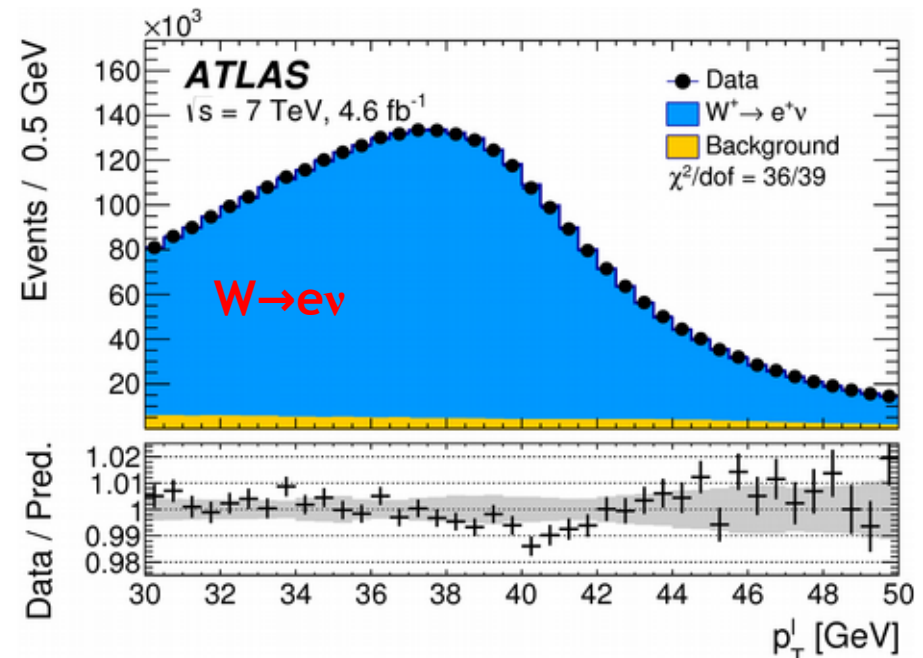
W mass measurement

ATLAS measurement of m_W uses well-understood 2011 data (7 TeV)
 ~15M $W \rightarrow \ell \nu$ decays

Both the lepton transverse momentum [$p_T(\ell)$] distribution, and the transverse mass [m_T] distributions are used - they are both sensitive to the value of m_W

Important experimental features:

- Lepton calibration using high statistics $Z \rightarrow \ell \ell$ sample
- Hadronic recoil ($\rightarrow p_T^{\text{miss}}$) also calibrated against $Z \rightarrow \ell \ell$
- LEP Z mass crucial input (2 MeV error)
- Detailed analysis of modelling uncertainties



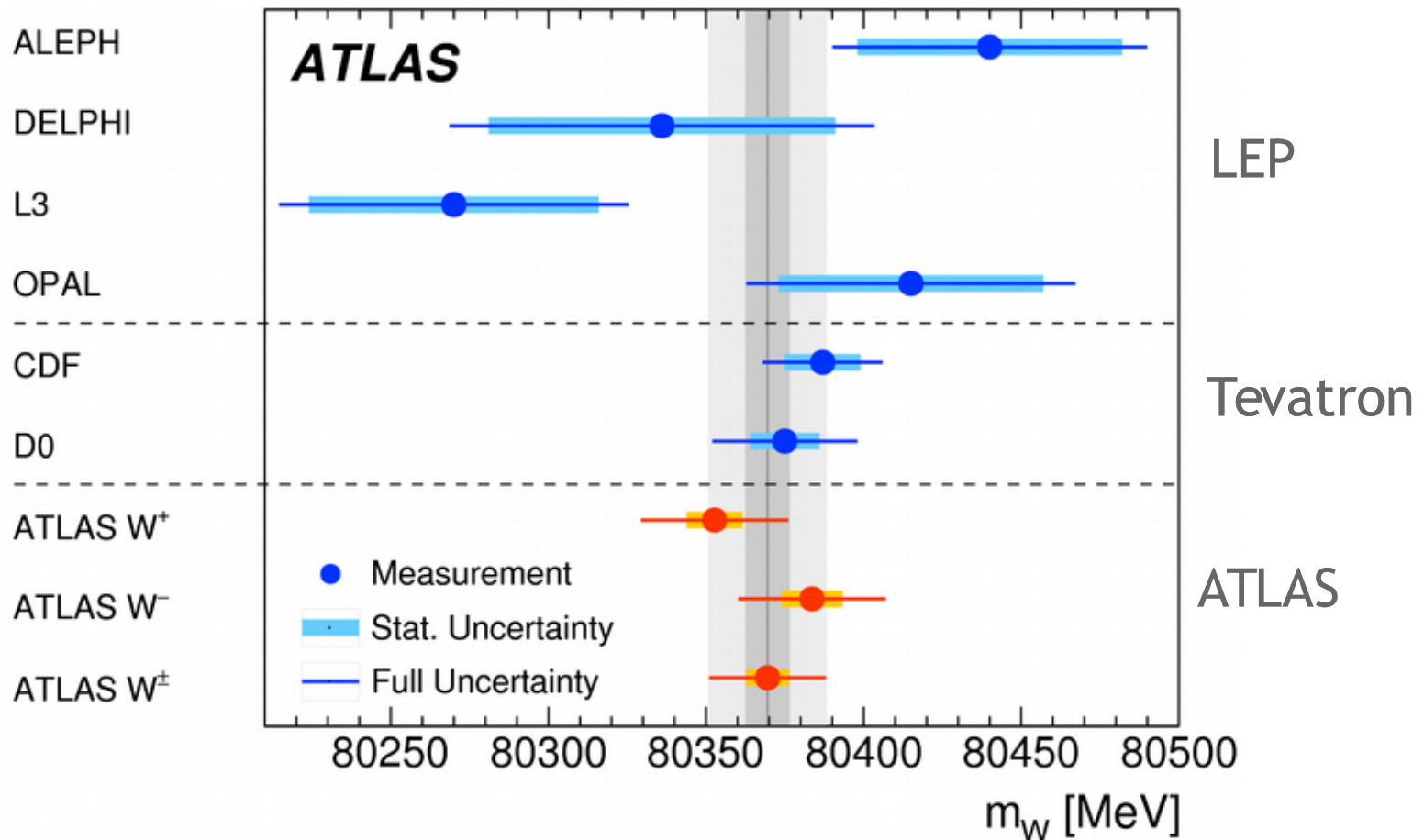
$$m_T = \sqrt{2p_T^\ell p_T^{\text{miss}} (1 - \cos \Delta\phi)}, \quad m_T [\text{GeV}]$$

W mass results

Combining the e and μ channels, charge signs and methods, overall:

$$m_W = 80370 \pm 19 \text{ MeV} \quad \pm 7 \text{ (stat.)} \pm 11 \text{ (exp. syst.)} \pm 14 \text{ (mod. syst.) MeV}$$

Measurement precision of 19 MeV (0.024%) equals best previous uncertainty, from CDF



Electroweak precision test

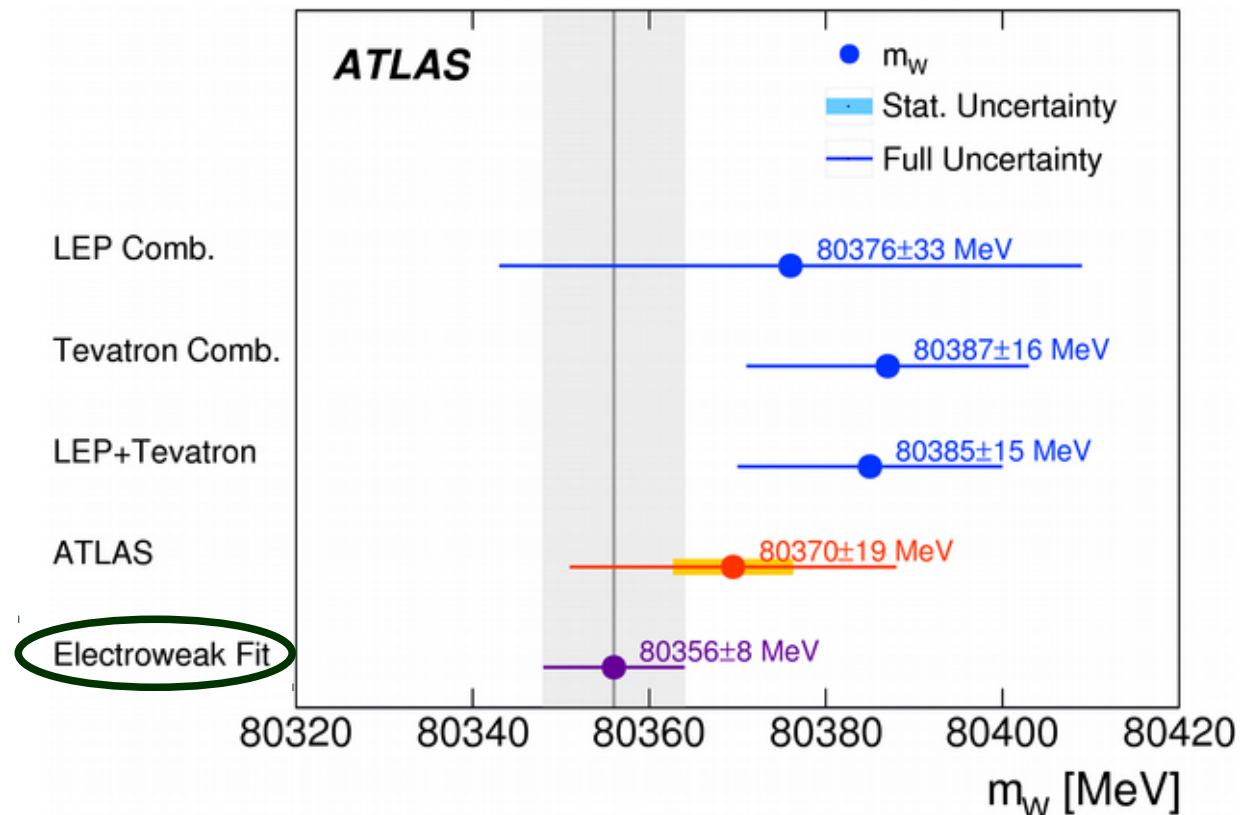
Within the SM framework, m_W is related to other quantities via:

$$m_W^2 \left(1 - \frac{m_W^2}{m_Z^2} \right) = \frac{\pi\alpha}{\sqrt{2}G_\mu} (1 + \Delta r),$$



Δr includes radiative effects (loops), and so depends on m_H and m_{top}

Fits to precision electroweak data from LEP/SLD and others, plus the LHC m_H and Tevatron+LHC m_{top} , provides a prediction of m_W (“indirect measurement in the framework of the SM”)



Precision electroweak fits

Within the SM framework, EW observables can be predicted using just five parameters

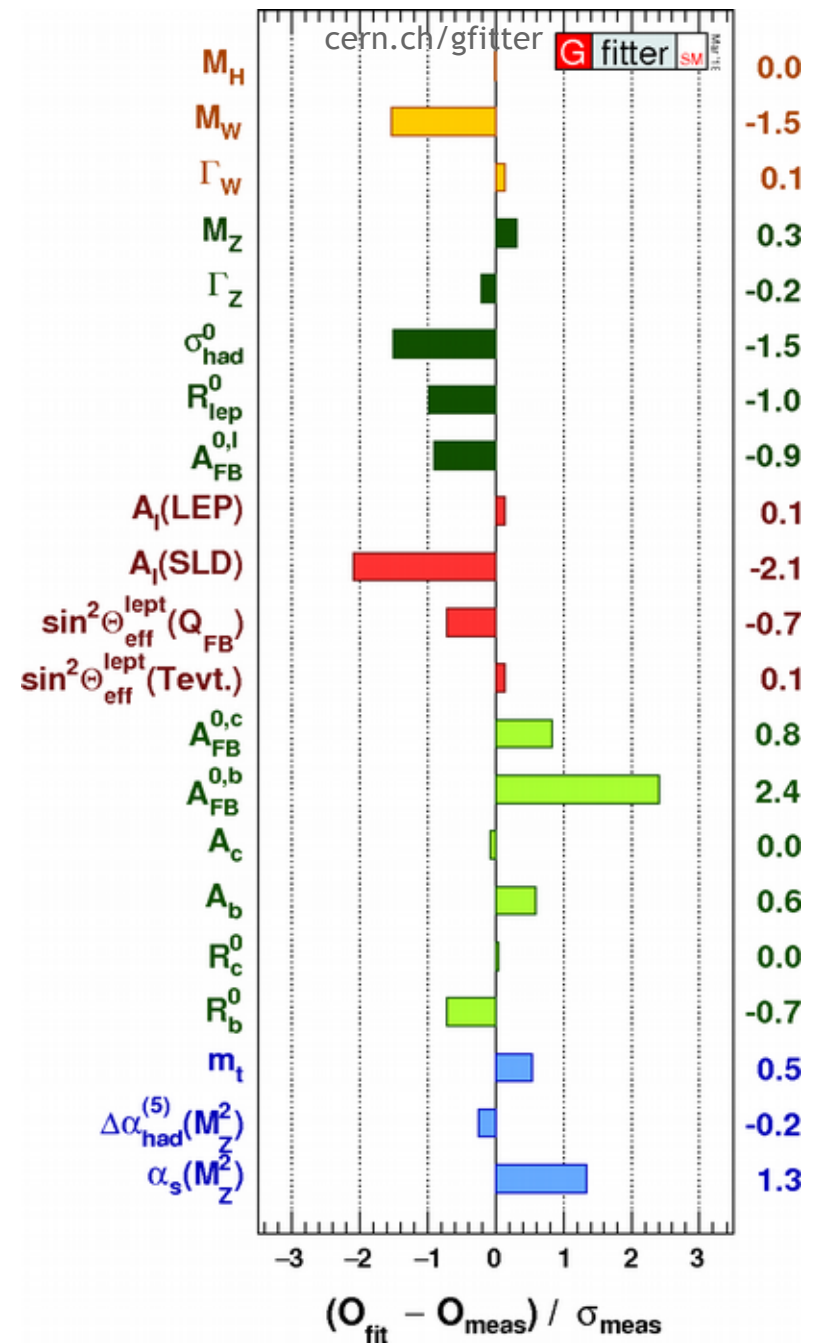
- Many more than five observables have been measured
- Requires theoretical predictions at as high a level as possible (must include loop diagrams!)



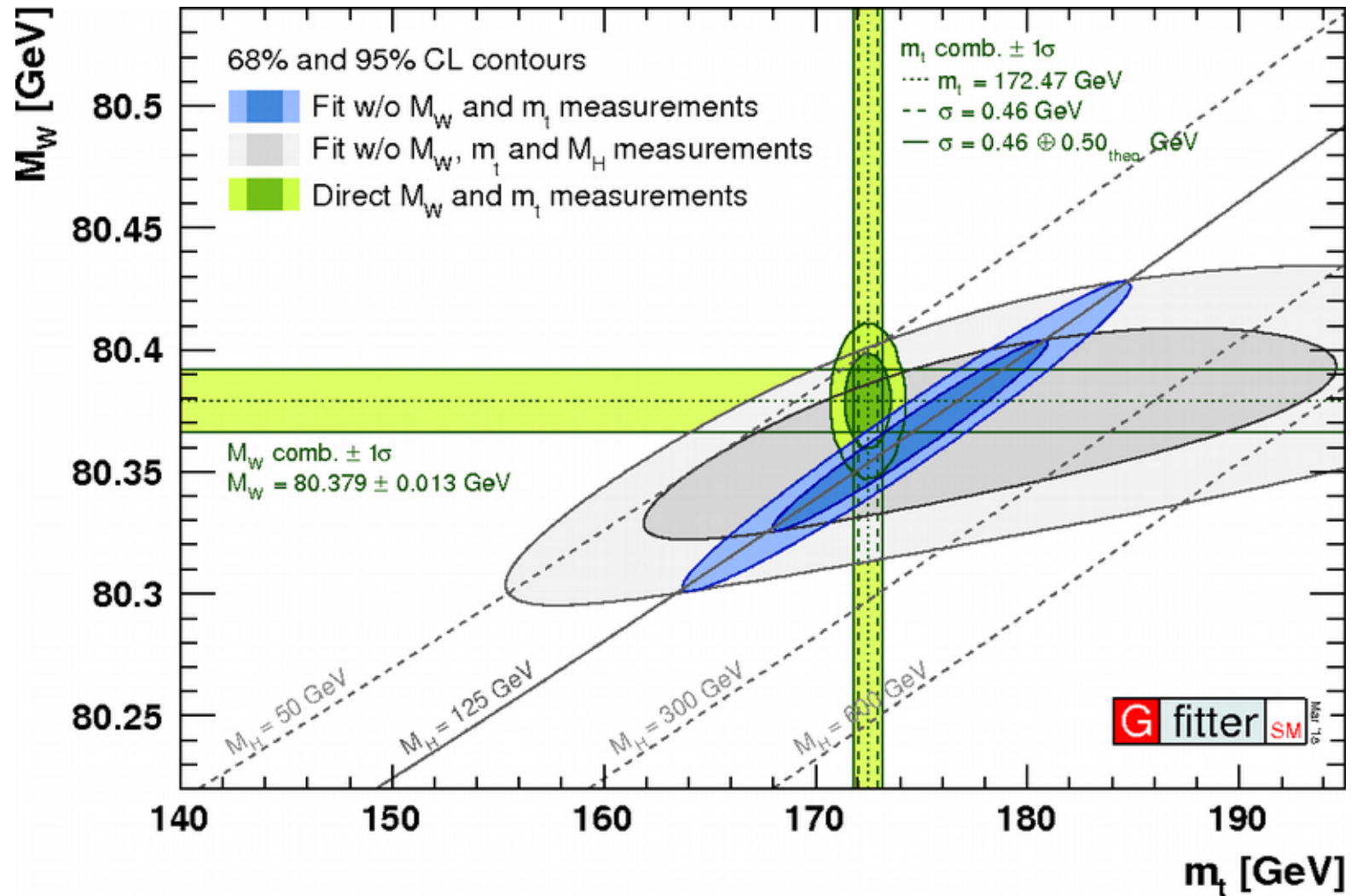
- We can fit all EW measurements for a *global EW precision test*

Latest Gfitter fit: $\chi^2=18.6$ for 15 d-of-f

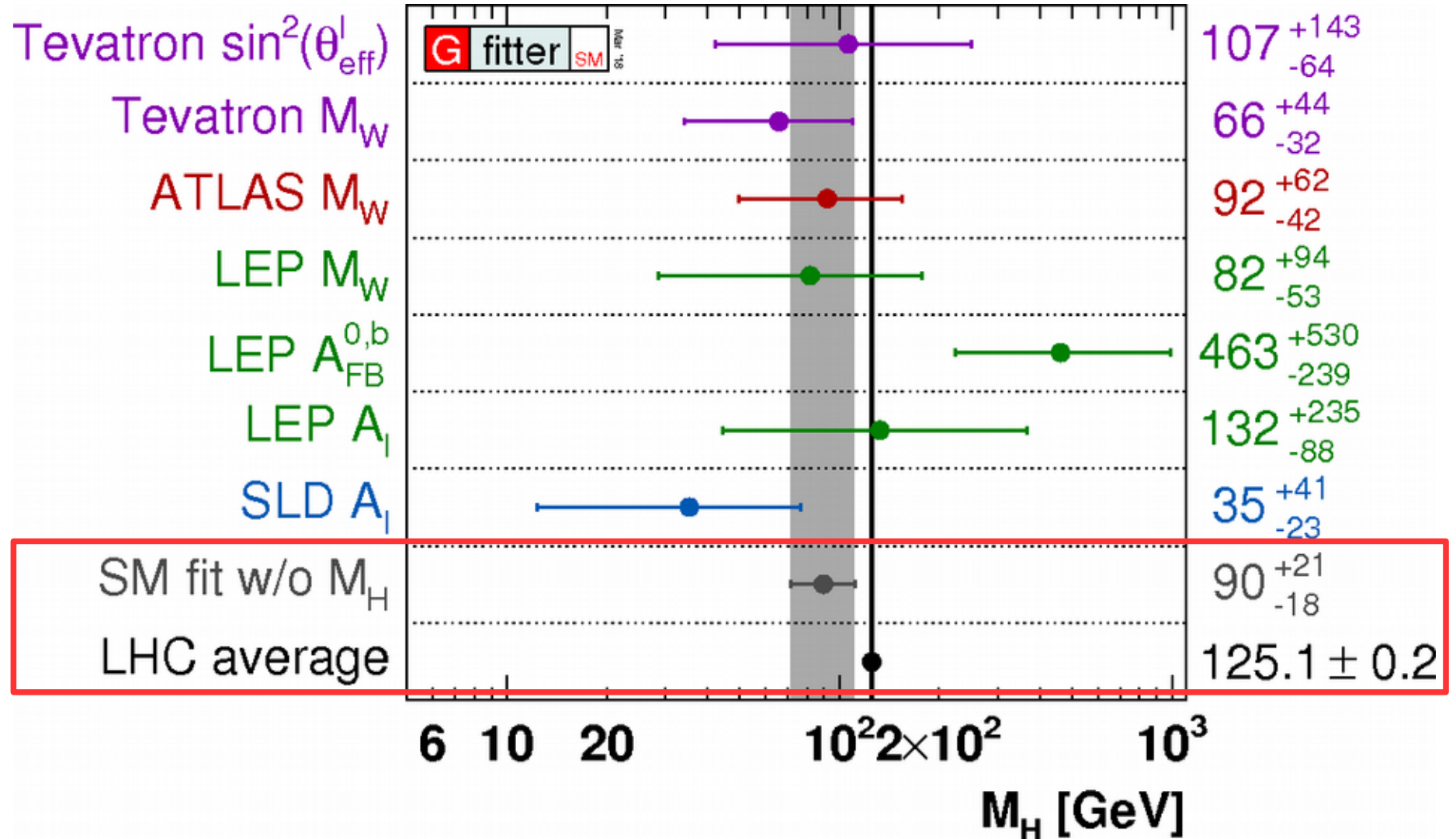
- We can re-interpret the other results into a prediction of m_W and m_{top}
- We can try to predict the Higgs boson mass using all the other measurements



Precision electroweak physics



Precision EW fits: "predicting" m_H



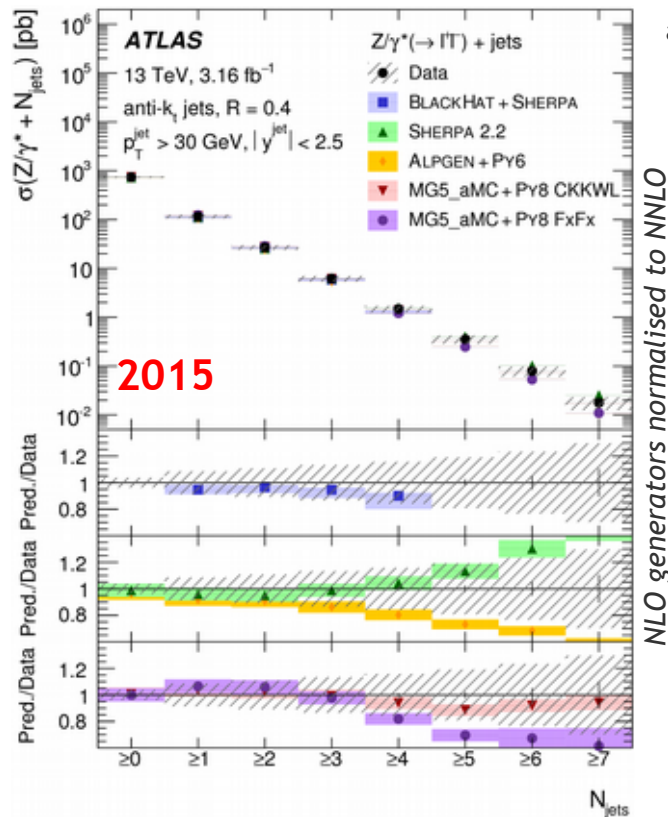
Z+jets at 13 TeV

Access and measure high jet multiplicities in 13 TeV data

- Test NLO predictions on events with high jet multiplicities
- Vector-boson plus jet events are a major background in searches

Fully corrected fiducial and differential cross-sections

- Z with up to 7 additional jets measured

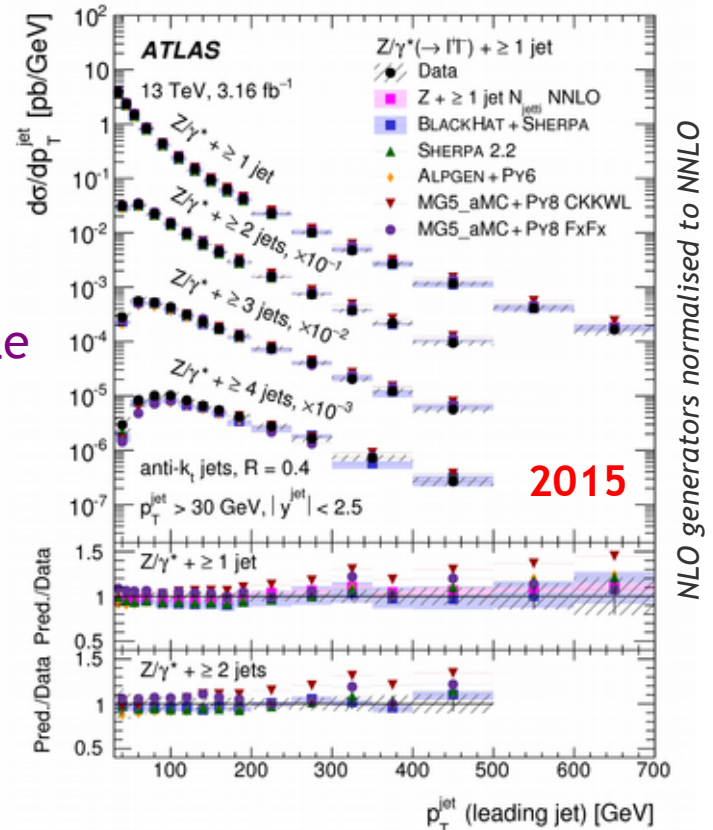


arXiv:1702.05725

~6M $Z \rightarrow \ell\ell$
event sample
(3.2 fb^{-1})

NLO generators normalised to NNLO

Jet multiplicity: main NLO generators do a good job, at least up to 6 jets



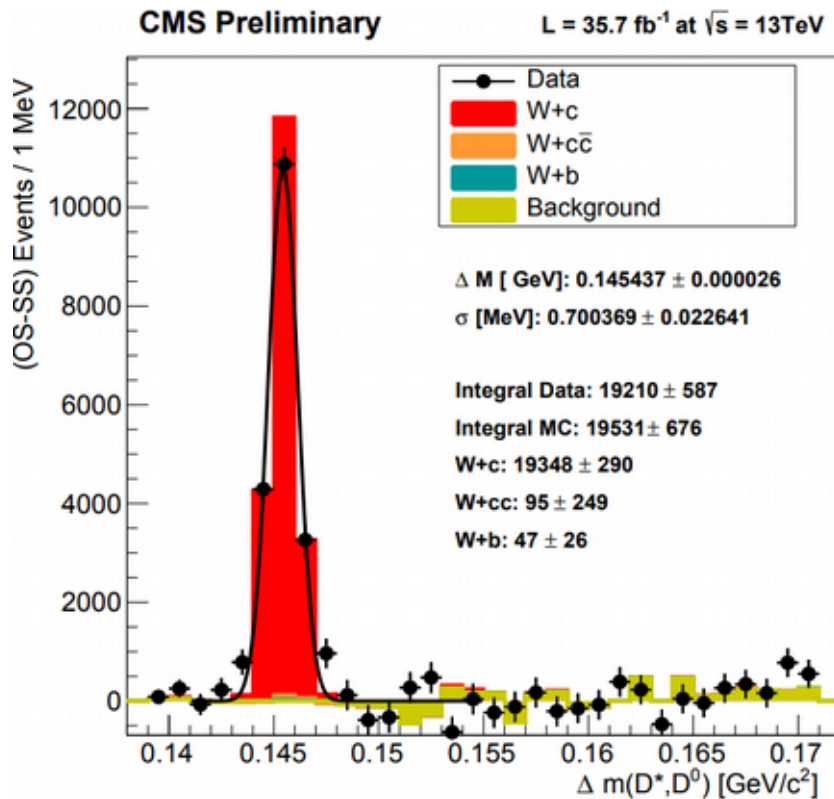
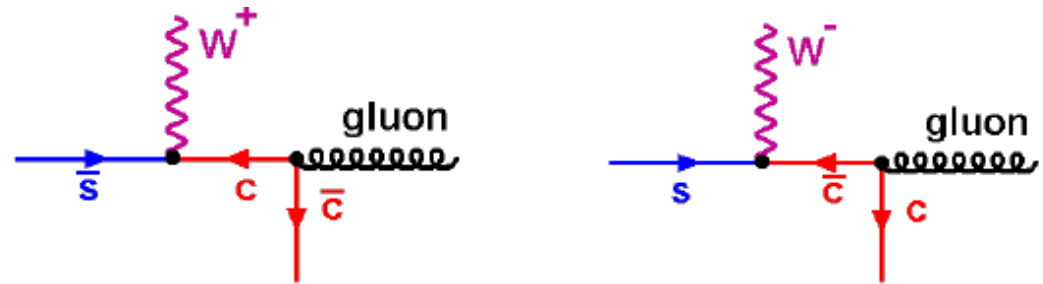
NLO generators normalised to NNLO

Leading jet p_T spectra:

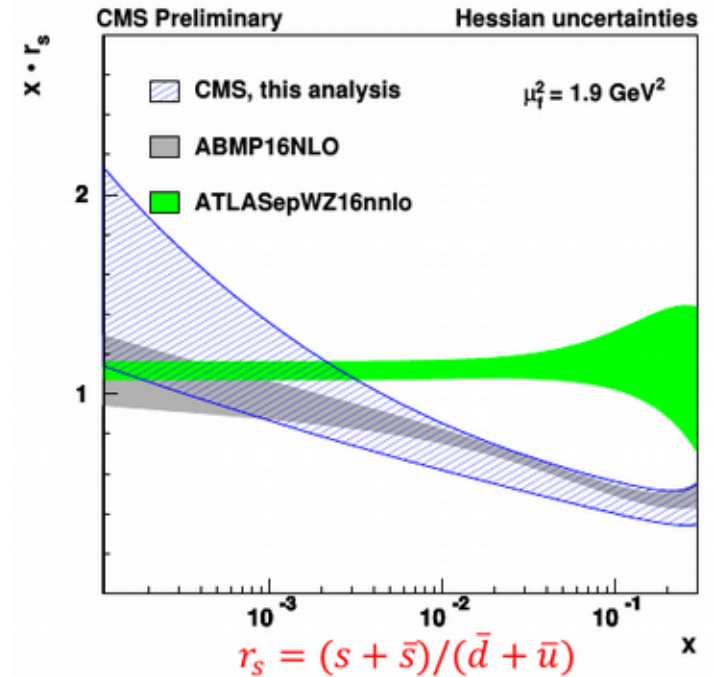
LO generators over-predict high- p_T tail
NLO generators provide better description 37

W + charm

A very new example:
 New result from CMS on W+c at 13 TeV →
 probes strange quark pdf - uses D*-tag



Sensitive to
strange sea



CMS-PAS-SMP-17-014

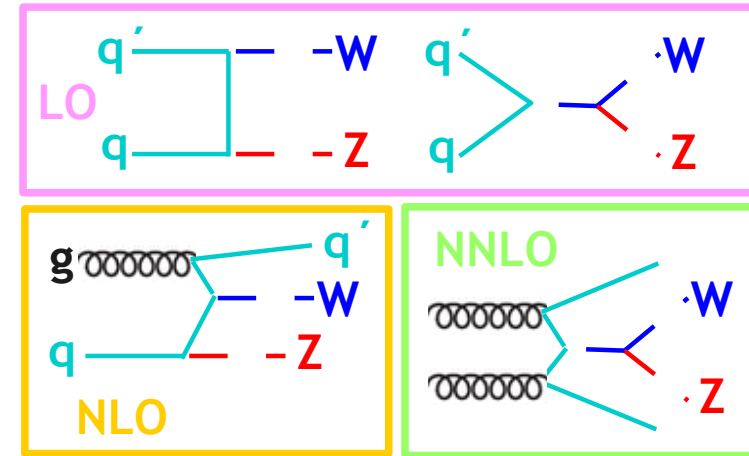
Tension with s-quark pdf from ATLAS - derived
 from inclusive W/Z production
Open question how this will be resolved!

Massive diboson production

Run-1 puzzle to describe inclusive diboson cross-sections

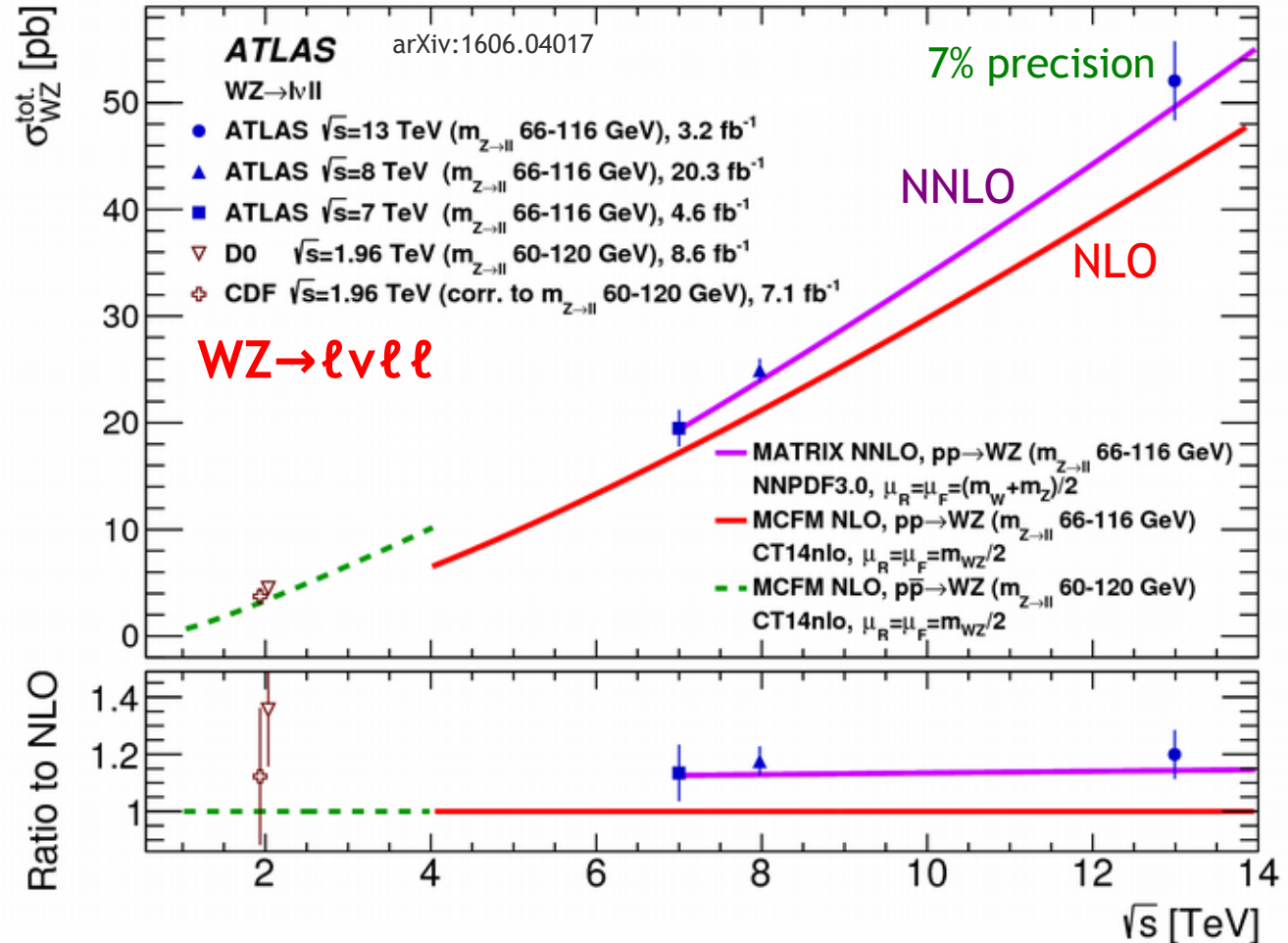
- Measurements tended to lie above next-to-leading order (NLO) calculations

NNLO calculations \rightarrow $\sim +20\%$ corrections and better agreement



Example:
WZ leptonic decays
NNLO calculations
describe data much
better than NLO

This run-1 puzzle
appears to be solved!



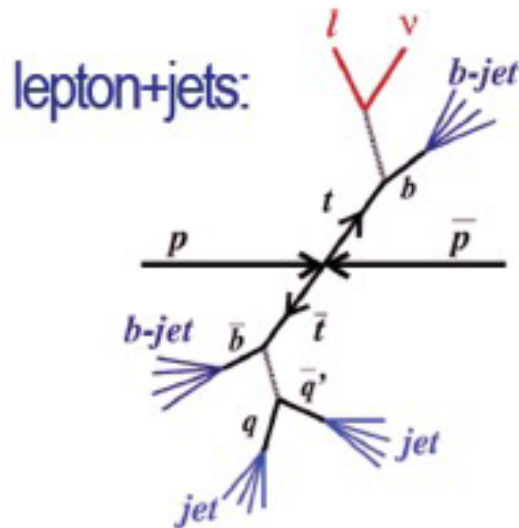
Precision top quark physics

To date tens of millions of $t\bar{t}$ pairs produced at the LHC (cf $\sim 75k$ at Tevatron, where the top quark was discovered)

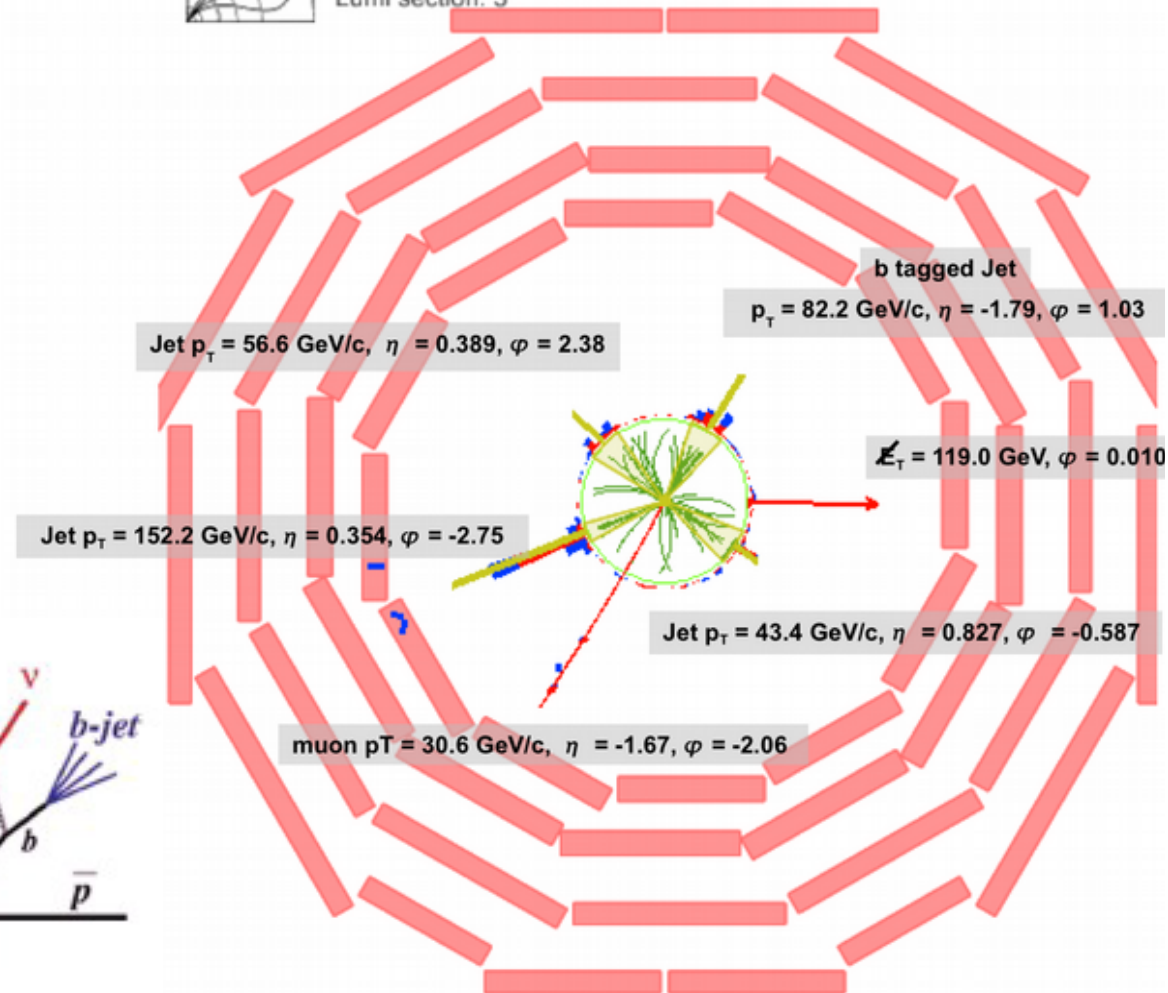
Are top quarks “special” objects?

- The coupling y_t of the $t\bar{t}H$ vertex has a predicted strength $y_t \sim 1$

→ Big programme to measure top production, properties and decays precisely



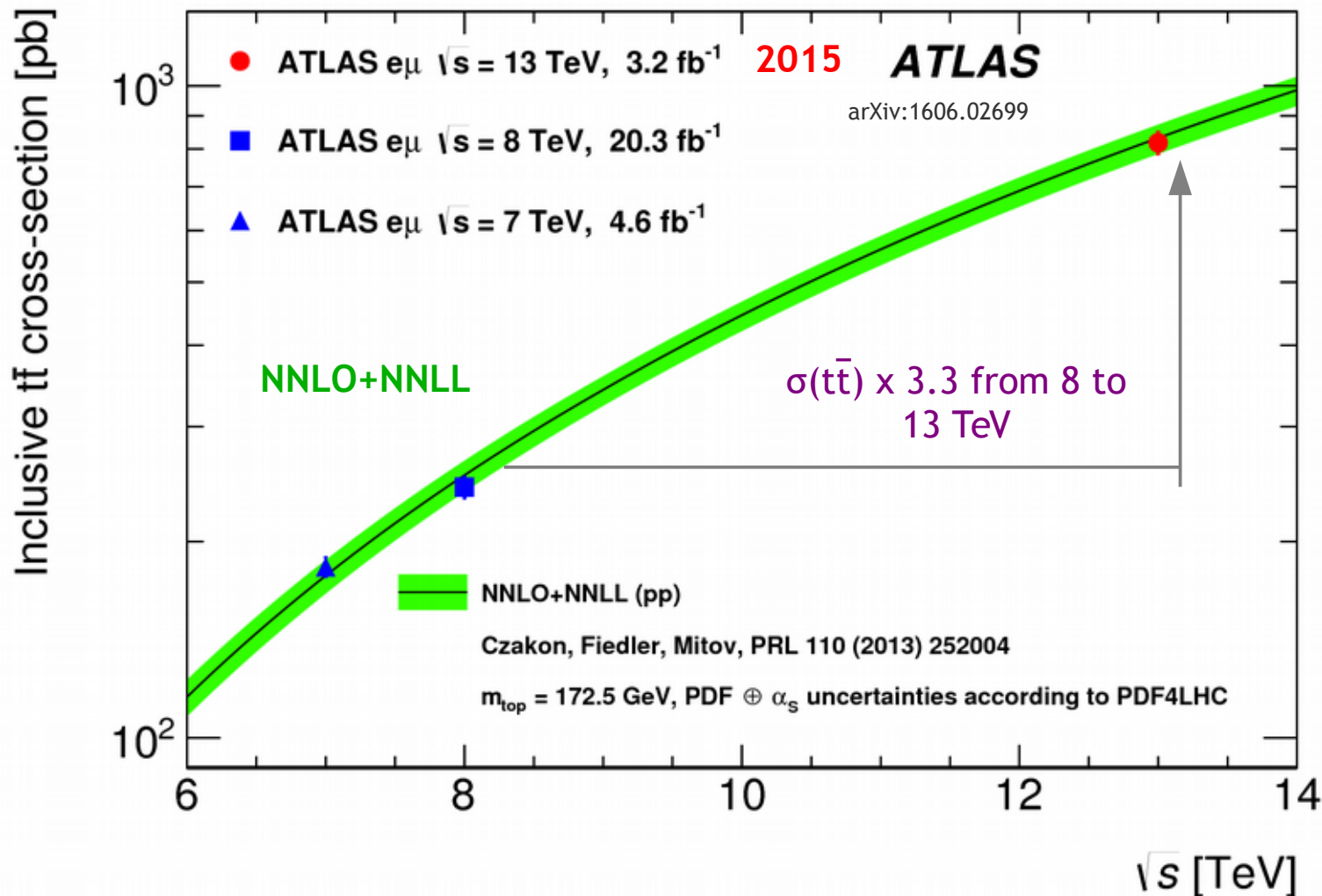
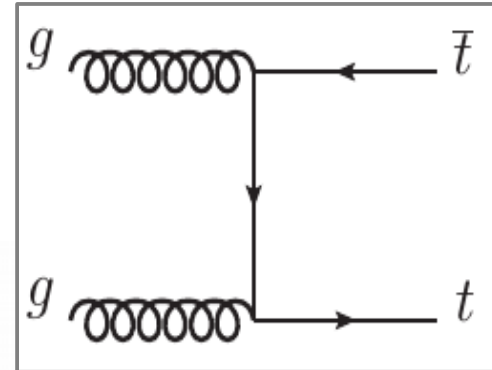
CMS Experiment at LHC, CERN
 Data recorded: Wed Jul 14 03:32:41 2010 CEST
 Run/Event: 140124 / 1749068
 Lumi section: 3



$t\bar{t}$ production

Single and double b-tagged $t\bar{t} \rightarrow bev\bar{b}\mu\bar{\nu}$ events allow to measure $t\bar{t}$ cross-section and b-tagging efficiency simultaneously

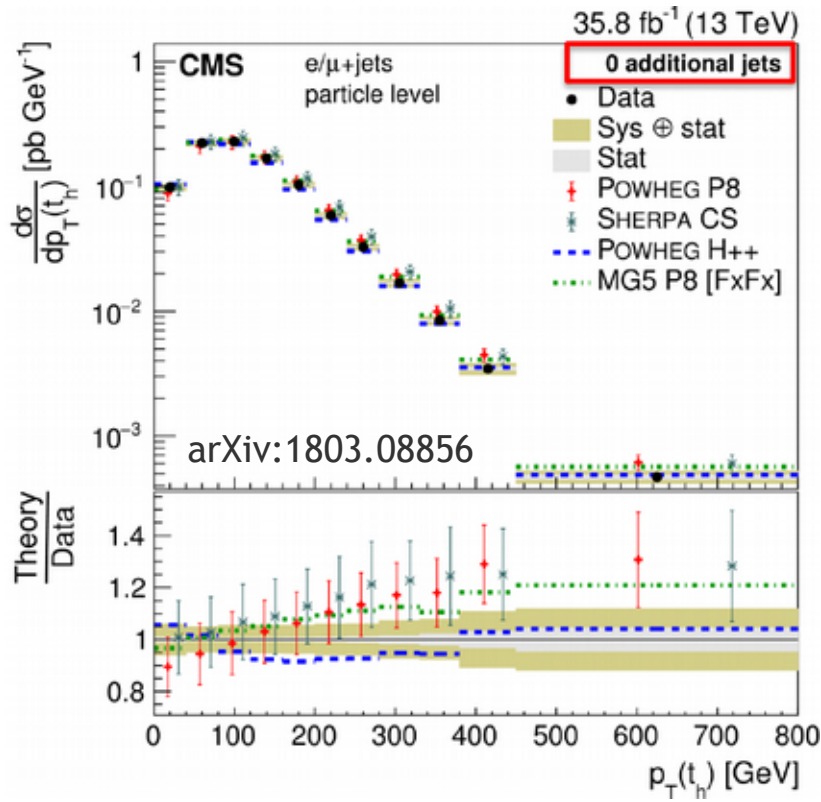
Precision $\pm(3.9-4.4)\%$ (7-13 TeV) better than NNLO+NNLL predictions ($\sim 5\%$)



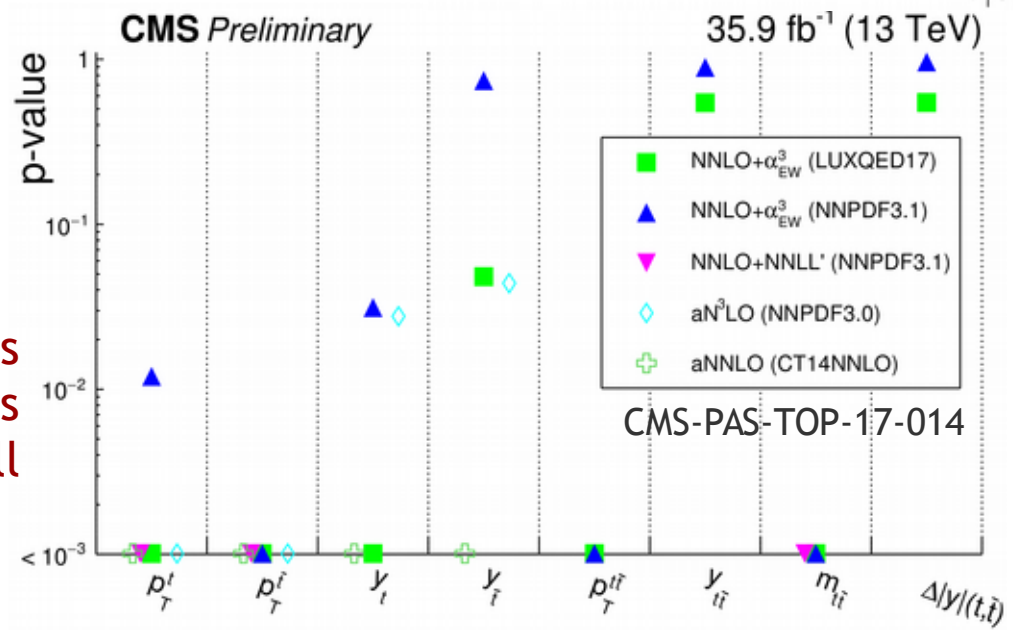
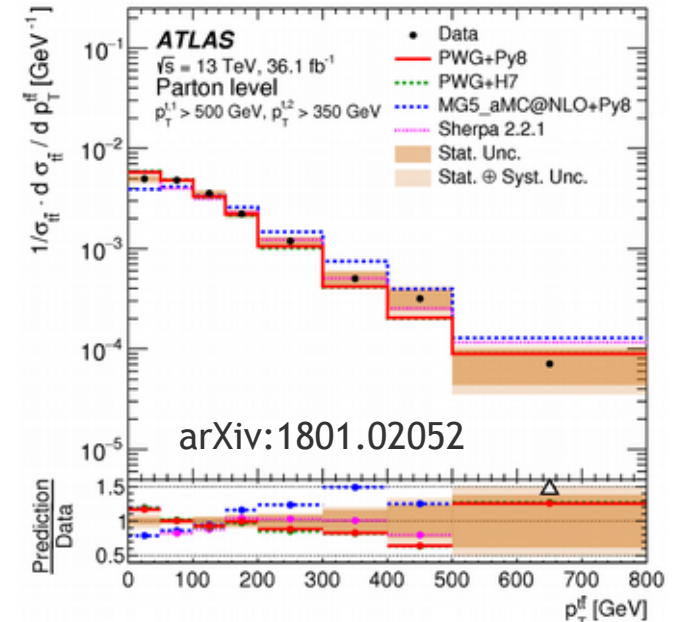
30 fb^{-1} @ 13 TeV
 $\sim 25\text{M } t\bar{t}$ produced

Top pair production

Long-standing (Run-1) difficulty to describe the observed $p_T(\text{top})$ spectrum, although NLO-based MC models much improved for Run-2

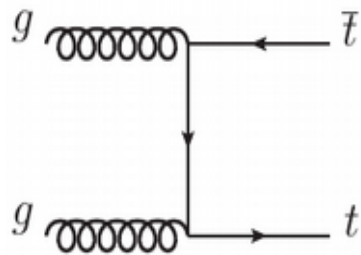


Still none of the predictions describe all of the observables well

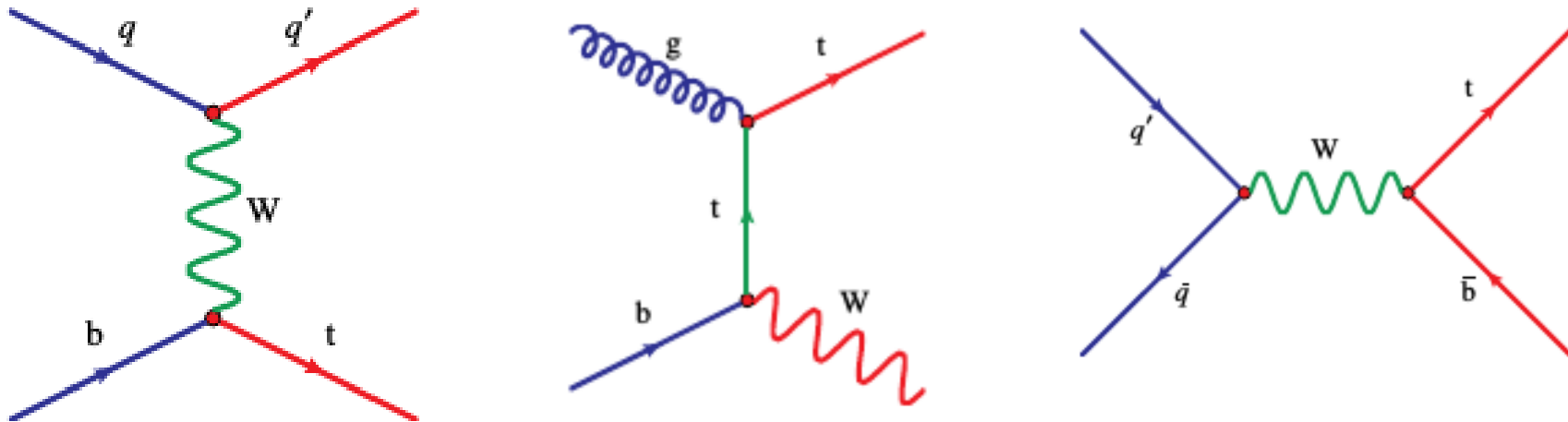


Single top

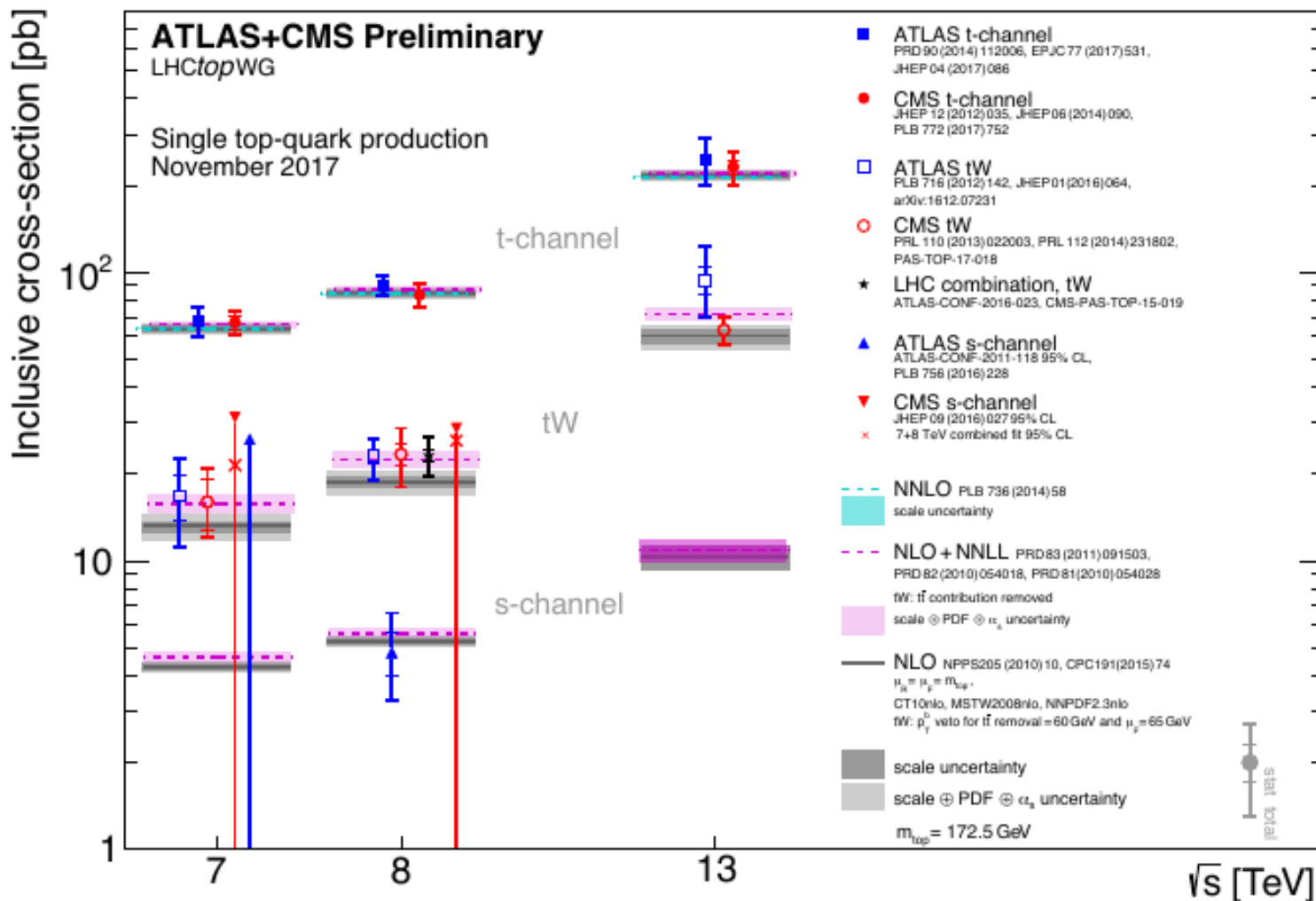
Top-anti-top quark pair-production is a strong interaction process with a high cross-section, as we just saw



We can produce top quarks in other ways too at the LHC, for example singly, via electroweak processes with a W involved



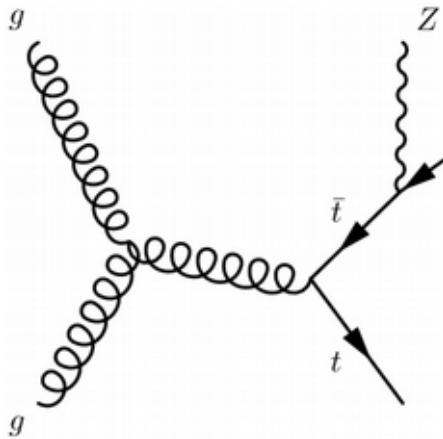
Single top cross-sections



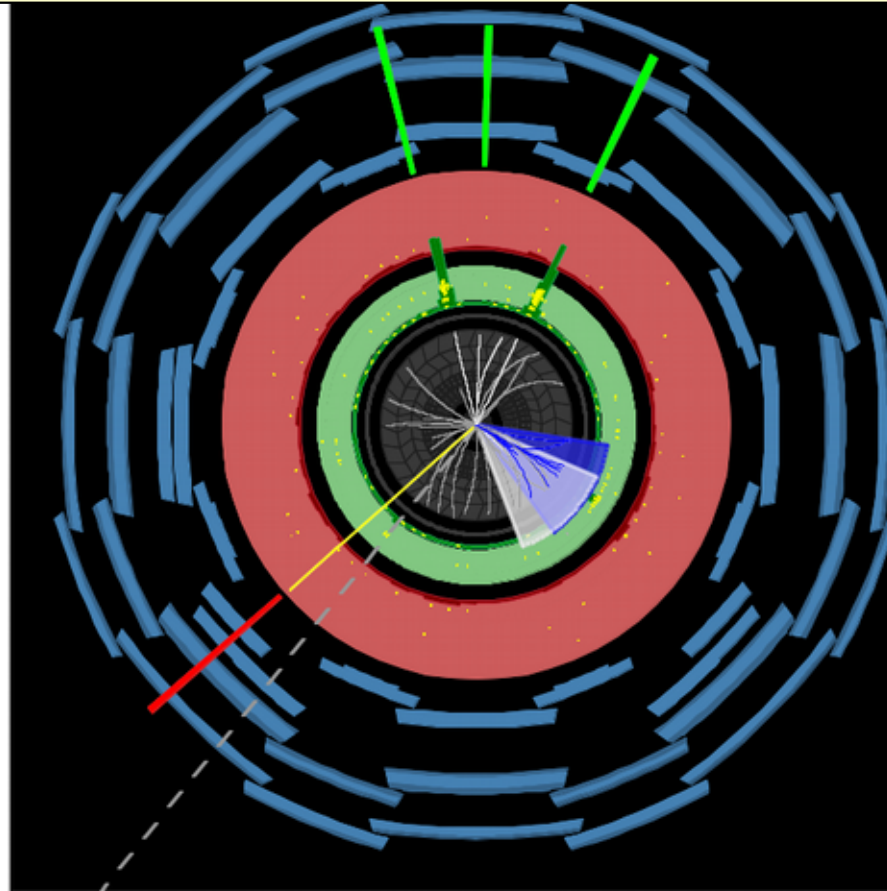
t-channel and Wt production measured differentially
s-channel still unobserved at LHC - observed at Tevatron
Other associated production channels tZq and tq̄q should be seen soon

Two tops and a Z boson!

Three very massive particles produced together - example diagram:



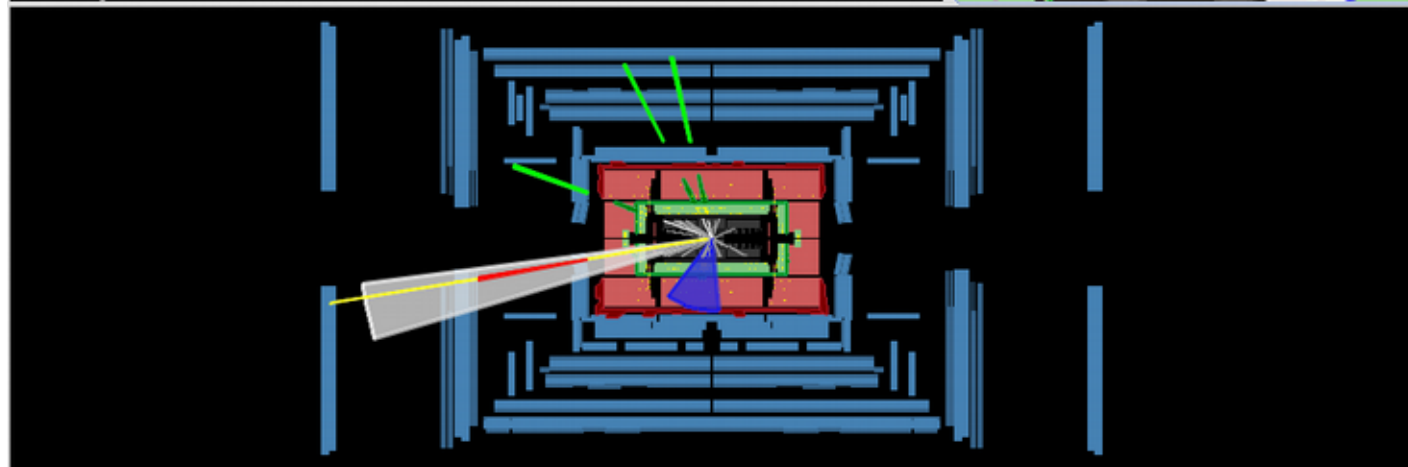
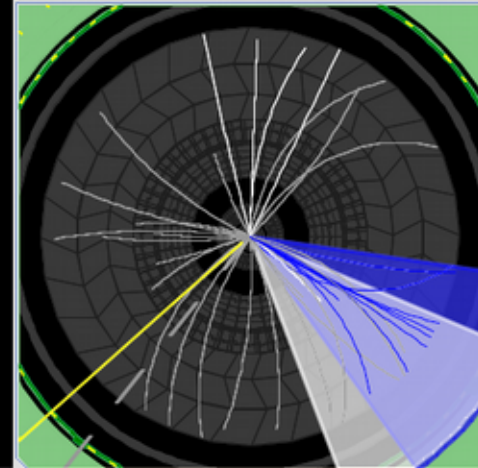
In the event shown, both top quarks decay to Wb , and the W decays to lepton plus neutrino \rightarrow total of four charged leptons (3e, 1 μ) plus 2 b-jets



 **ATLAS**
EXPERIMENT

Run Number: 212172, Event Number: 70154285

Date: 2012-10-06 08:03:38 CEST



Summary

- Computational technology to predict cross-sections of Standard Model process at the LHC is now pretty sophisticated (NLO, NNLO ...)
- Many processes have been measured, and generally are well described by the Standard Model
 - Measurements now often more precise than the predictions
 - Work for the theorists!!! (and experimenters, e.g. to constrain better the pdfs)
- Only a small part (<5%) of the LHC data sample has been collected - there is much more to explore with precision measurements, and advancing our understanding of QCD and electroweak physics - even ignoring the Higgs (next talk) and possible new physics (this morning)!

All the best to all of you with your studies and research - we hope to see some of you working at the LHC in future!

Measuring the luminosity

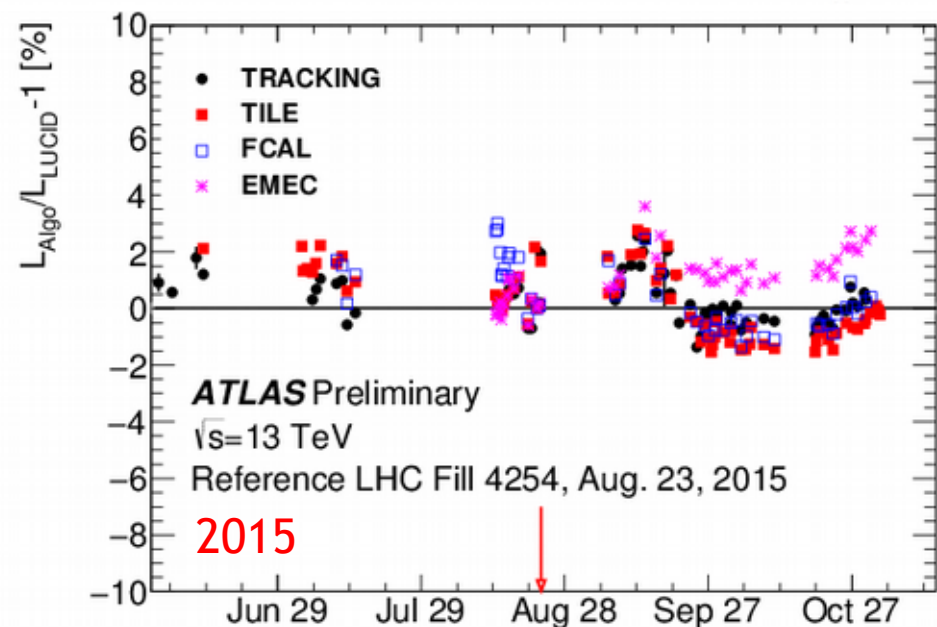
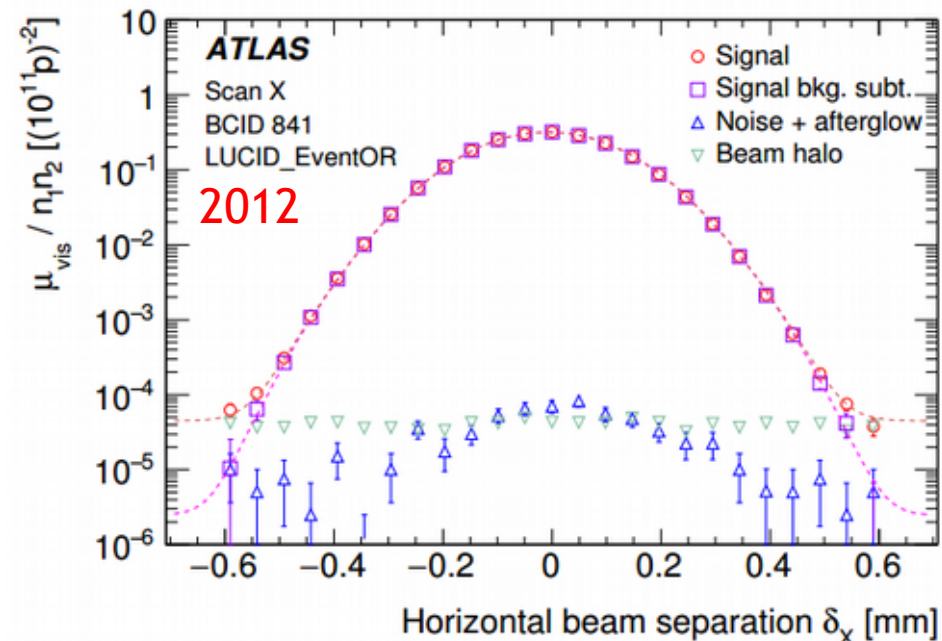
Van Der Meer (VDM) beam-beam separation scans allow to determine the absolute bunch luminosity from measured beam parameters:

$$\mathcal{L}_b = \frac{f_r n_1 n_2}{2\pi \Sigma_x \Sigma_y}$$

f_r revolution frequency; n_i number of p/bunch
 $\Sigma_{x,y}$ convolved beam sizes, from the scans

Multiple detectors allow to map between VDM-scan conditions and normal operations

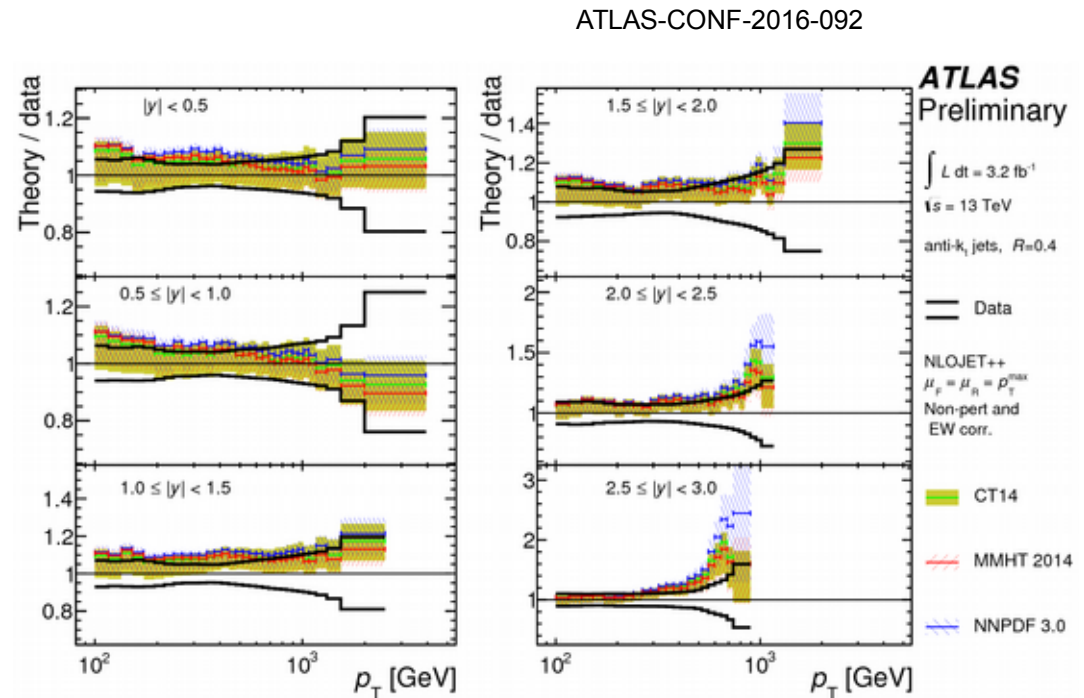
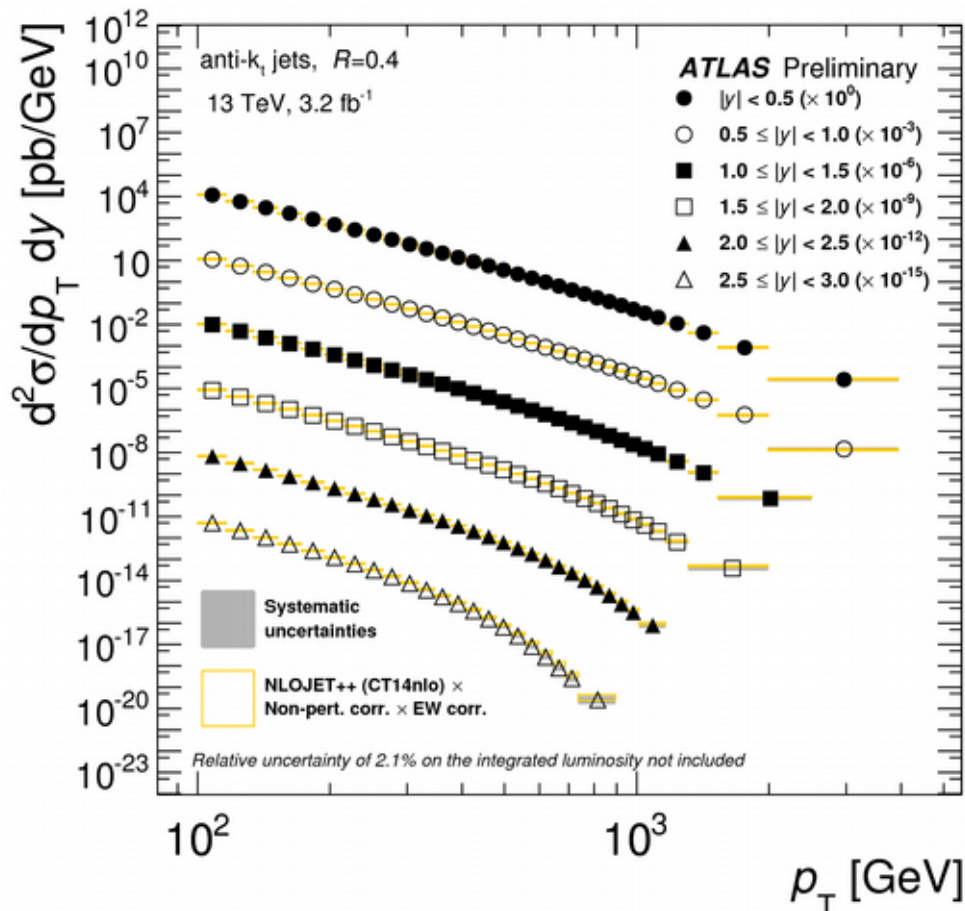
Luminosity uncertainty $\sim \pm 2\%$



Dijet cross-sections at 13 TeV

Measurements in progress on 13 TeV data - preliminary results released with 2015 data in August last year

- Range of corrected measurements extends to jet transverse momentum $p_T \sim 3.2$ TeV



Ratios of measured spectrum compared to models using different pdfs