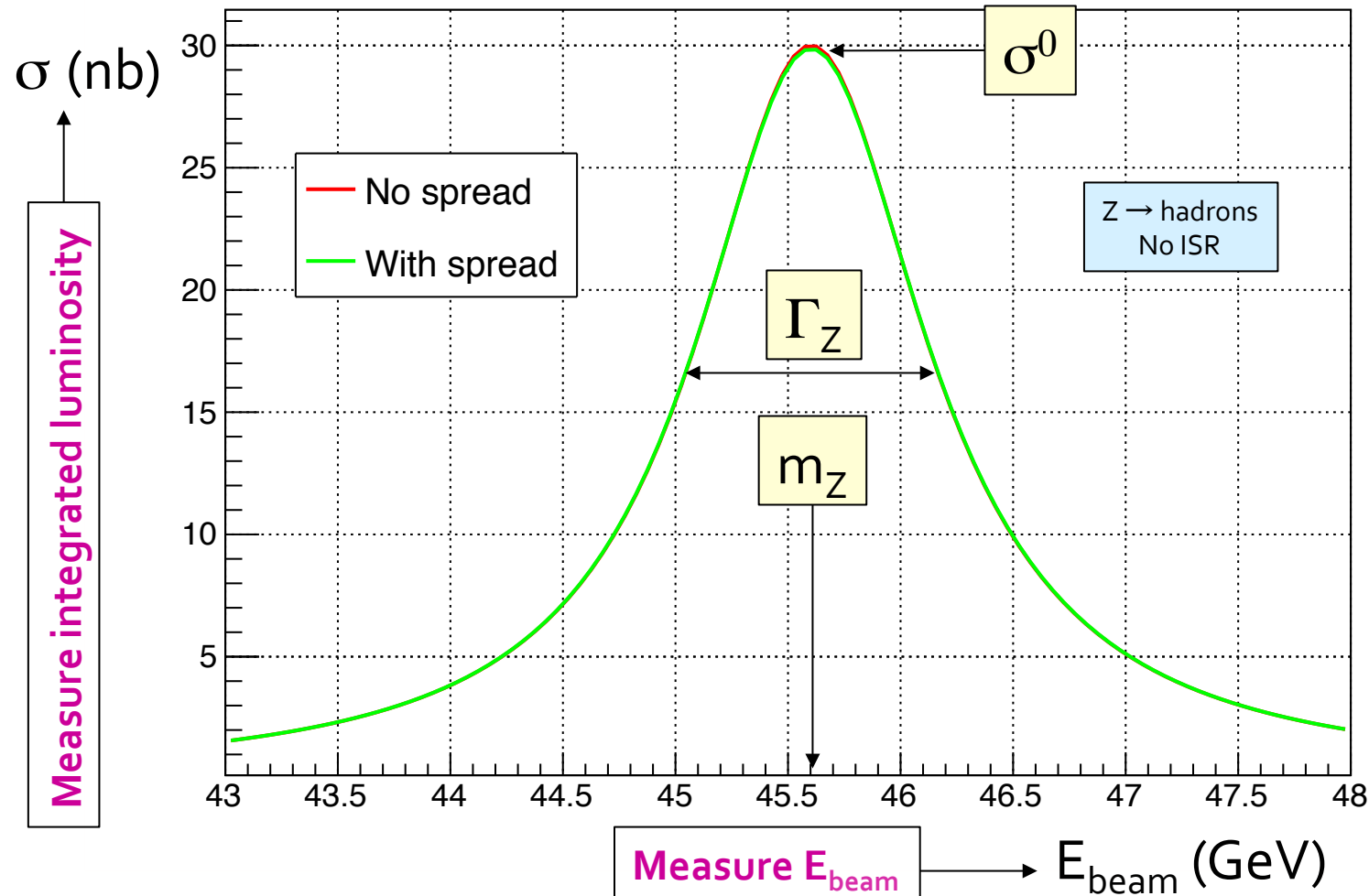


Z pole scan strategy and \sqrt{s} spread measurement

	Z	W	H (ZH)	ttbar
beam energy [GeV]	45.6	80	120	175-182.5
arc cell optics	60/60	60/60	90/90	90/90
emittance hor/vert [nm]/[pm]	0.27/1.0	0.84/1.7	0.63/1.3	1.4/2.8
β^* horiz/vertical [m]/[mm]	0.15/1.8	0.2/1	0.3/1	1/1.6
SR energy loss / turn (GeV)	0.036	0.34	1.72	9.21
total RF voltage [GV]	0.10	0.75	2.0	8.8-10.3
energy acceptance [%]	± 1.3	± 1.3	± 1.7	$\pm 2.4-2.8$
energy spread (SR / BS) [%]	0.038 / 0.132	0.066 / 0.165	0.099 / 0.165	0.15 / 0.20
bunch length (SR / BS) [mm]	3.5 / 12.1	3.0 / 7.5	3.15 / 5.3	2.75 / 3.80
bunch intensity [10^{11}]	1.7	2.3	1.8	3.2-3.35
no. of bunches / beam	16640	1300	328	40-33
beam current [mA]	1390	147	29	6.4-5.4
SR total power [MW]	100	100	100	100
luminosity [$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$]	230	34	8.5	1.9-1.7
luminosity lifetime [min]	70	24	18	25
allowable asymmetry [%]	± 5	± 3	± 3	± 3

Z pole scan strategy: # points ? Energies ?

- At the Z pole, $\delta E_{\text{beam}}/E_{\text{beam}} \approx 0.132\% \Rightarrow \delta E_{\text{beam}} \approx 60 \text{ MeV}$



Z pole scan strategy without energy spread

- **Three parameters to fit : m_Z , Γ_Z , and σ^0**
 - ◆ Need (at least) three centre-of-mass energies at and around the Z pole
 - With a calibration of the beam energy to 50 keV (A. Blondel's talk)
 - With a relative precision on the absolute luminosity of 10^{-4} (M. Dam's talk)
 - ➔ And a point-to-point relative precision of 5×10^{-5}
 - ◆ Precision targets
 - 100 keV on m_Z ; 100 keV on Γ_Z ; dominated by the above.
 - Statistical precision is a few keV with the FCC-ee statistics
 - ◆ Possible beam energies around half-integer spin tunes $\nu = E_{\text{beam}} / 0.4406486$

ν	99.5	100.5	101.5	102.5	103.5
E_{beam}	43.85	44.29	44.72	45.16	45.61
\sqrt{s}	87.69	88.57	89.43	90.31	91.21

ν	103.5	104.5	105.5	106.5	107.5
E_{beam}	45.61	46.04	46.49	46.92	47.36
\sqrt{s}	91.21	92.07	92.97	93.83	94.71

Z pole scan strategy without energy spread

Result of the 3-parameter fit to a Breit-Wigner

- ◆ With 100 ab^{-1} at the peak ($\sqrt{s} = 91.21 \text{ GeV}$)
- ◆ With $30 \text{ ab}^{-1} + 30 \text{ ab}^{-1}$ at peak ± 1 , peak ± 2 , peak ± 3 , peak ± 4 (in spin tune)

Scan	± 1	± 2	± 3	± 4	$(-4,+3)$ or $(-3,+4)$
$\sigma(m_Z)$ [keV]	84	85	91	99	95
$\sigma(\Gamma_Z)$ [keV]	313	153	116	100	107

- Standard at LEP
 - EWPO measurements would have been statistics limited with ± 3 or ± 4
- All these options meet the (100 keV, 100 keV) target at FCC-ee
 - Need another criterion to decide

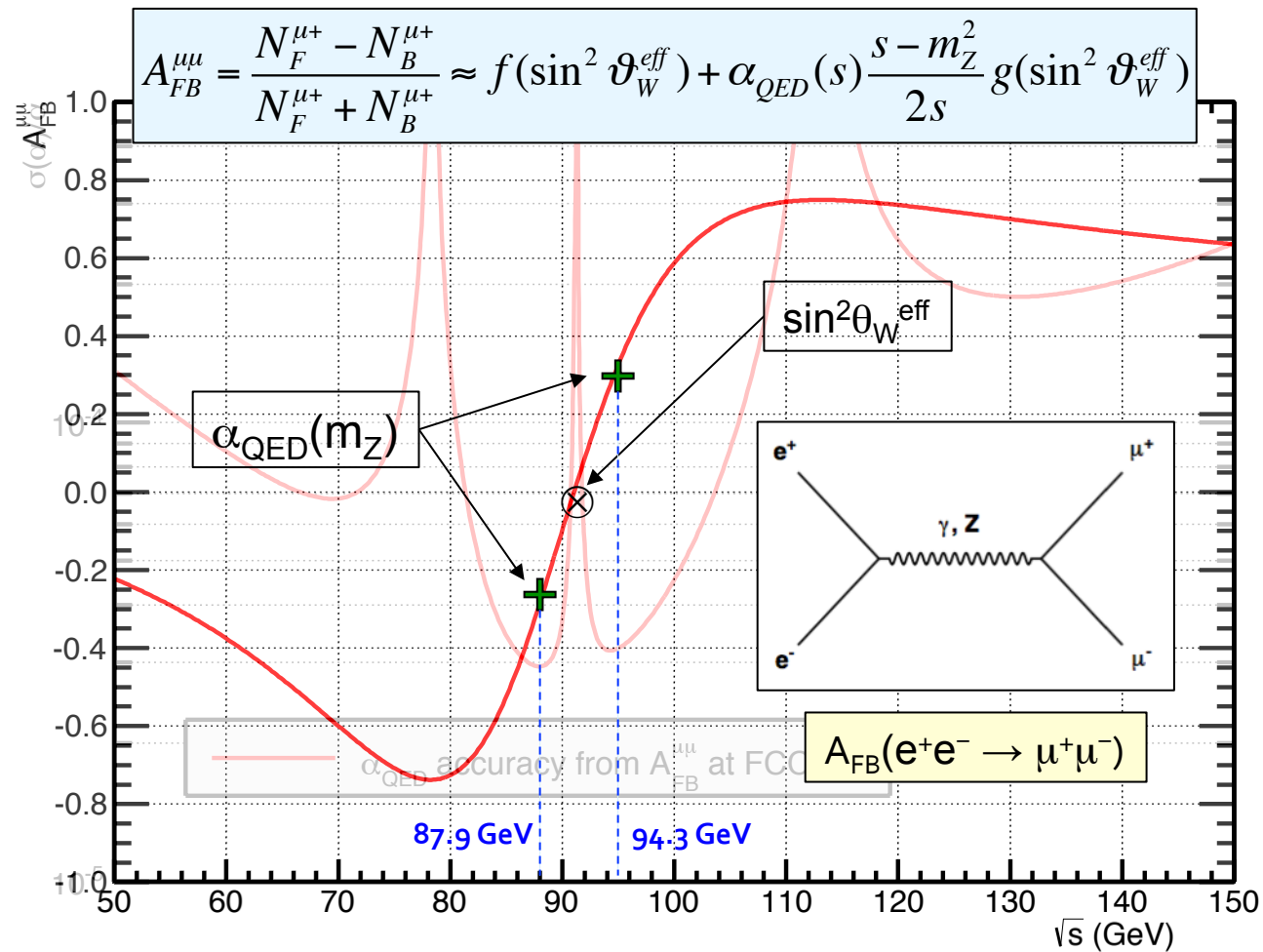
Note: Results of the fit with Peak, Peak ± 2 , Peak ± 4

- $\sigma(m_Z) = 64 \text{ keV}$
- $\sigma(\Gamma_Z) = 97 \text{ keV}$

At the expense of a loss of statistics for rare Z decays

Z pole scan strategy without energy spread

- Choice driven by $\sin^2 \theta_W^{\text{eff}}$ and $\alpha_{\text{QED}}(m_Z)$ determination from $A_{\text{FB}}(\mu\mu)$



Z pole scan strategy without energy spread

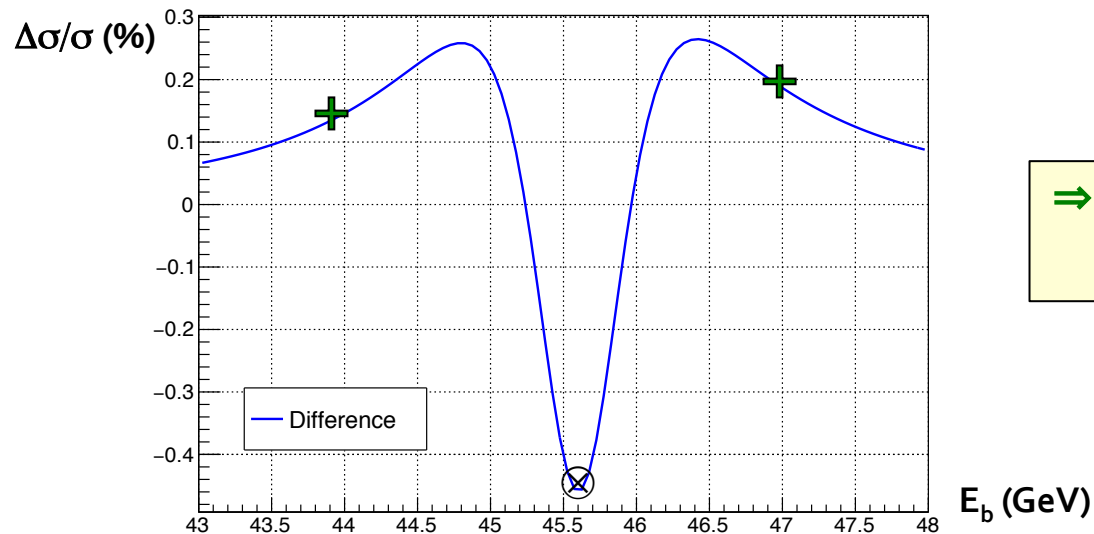
- **Best energies for A_{FB}**
 - ◆ 91.2 GeV for the effective weak mixing angle
 - ◆ 87.9 GeV and 94.3 GeV for the electromagnetic coupling constant
- **Closest half-integer tunes**
 - ◆ $\nu = 103.5$ for the peak
 - $\sqrt{s} = 91.2$ GeV
 - ◆ $\nu = 99.55$ below the peak ($Pk - 4$)
 - $\sqrt{s} = 87.8$ GeV
 - ◆ $\nu = 106.55$ OR $\nu = 107.5$ above the peak ($Pk + 3$ or $Pk + 4$)
 - $\sqrt{s} = 93.9$ GeV OR $\sqrt{s} = 94.7$ GeV (equidistant from 94.3 GeV)
 - ➔ Favour $Pk + 3$ for the sake of statistics (rare decays)
- **Minimal scan strategy within four years and with two detectors**
 - ◆ 100 ab^{-1} at the peak ($\sqrt{s} = 91.2$ GeV)
 - ◆ $30 \text{ ab}^{-1} + 30 \text{ ab}^{-1}$ at peak $- 4$ ($\sqrt{s} = 87.8$ GeV) and peak $+ 3$ ($\sqrt{s} = 93.9$ GeV)

Note: more statistics is always welcome
(3 IPs instead of 2; 6 years instead of 4; 5 points instead of 3; etc.)

Z lineshape with energy spread

□ The Z lineshape is slightly modified

- ◆ With $\delta E \sim 6\sigma$ MeV for each beam, $\delta\sigma/\sigma \sim -0.45\%$ at the peak and $+0.15\%$ off peak



⇒ Apparent increase of the Z width

$$\Gamma_Z \rightarrow [\Gamma_Z^2 + 8 \delta E_{\text{peak}}^2]^{1/2}$$

- ◆ If the energy spread δE_{peak} at the peak is known with an uncertainty $\sigma(\delta E_{\text{peak}})$

- $\sigma(\Gamma_Z) \sim 12 \text{ MeV} \times \sigma(\delta E_{\text{peak}})/\delta E_{\text{peak}}$

- ➔ 1% uncertainty of δE_{peak} leads to $> 120 \text{ keV}$ uncertainty on Γ_Z !

- Need to find a way to determine δE_{peak} to $\sim 0.2\%$

- ➔ Relaxed to $\sim 1\%$ for $\sigma(\delta E_{\text{off-peak}})/\delta E_{\text{off-peak}}$

Z lineshape with energy spread

- Can m_Z , Γ_Z , and σ_0 and δE be fit altogether with a five-point scan ?
 - ◆ IFF δE is constant in time, identical at all points, independent of the bunch ...
 - $\sigma(\delta E)/\delta E \sim 1\%$
 - $\sigma(m_Z) = 85 \text{ keV}$
 - $\sigma(\Gamma_Z) = 350 \text{ keV}$
 - ◆ And we certainly do not want to assume that the energy spread
 - Does not depend on time
 - Does not depend on energy
 - Does not depend on the bunch
 - Does not depend on anything else

Better idea(s) required

Note: $A_{\text{FB}}^{\mu\mu}$ with energy spread

- **Convoluting A_{FB} with a Gaussian has almost no effect at the peak**
 - ◆ The sampling of \sqrt{s} around $\sqrt{s_-} = 87.8$ and $\sqrt{s_+} = 93.9$ GeV, however, is not uniform
 - Has to weigh the asymmetry by the production cross section:

$$\Delta A_{\text{FB}}^{\mu\mu}(s_{\pm}) = \frac{\int A_{\text{FB}}^{\mu\mu}(s) \sigma_{\mu\mu}(s) \exp\left[-\frac{(\sqrt{s}-\sqrt{s_{\pm}})^2}{2s_{\pm}\delta^2}\right] d\sqrt{s}}{\int \sigma_{\mu\mu}(s) \exp\left[-\frac{(\sqrt{s}-\sqrt{s_{\pm}})^2}{2s_{\pm}\delta^2}\right] d\sqrt{s}} - A_{\text{FB}}^{\mu\mu}(s_{\pm})$$

- ◆ Effect is to increase the effective $\sqrt{s_-}$ and decrease the effective $\sqrt{s_+}$
 - And therefore decrease the absolute value of the asymmetry in both points

$$\frac{\Delta A_{\text{FB}}}{A_{\text{FB}}}(s_-) = -0.99 \times 10^{-3} \quad \text{and} \quad \frac{\Delta A_{\text{FB}}}{A_{\text{FB}}}(s_+) = -1.03 \times 10^{-3}$$

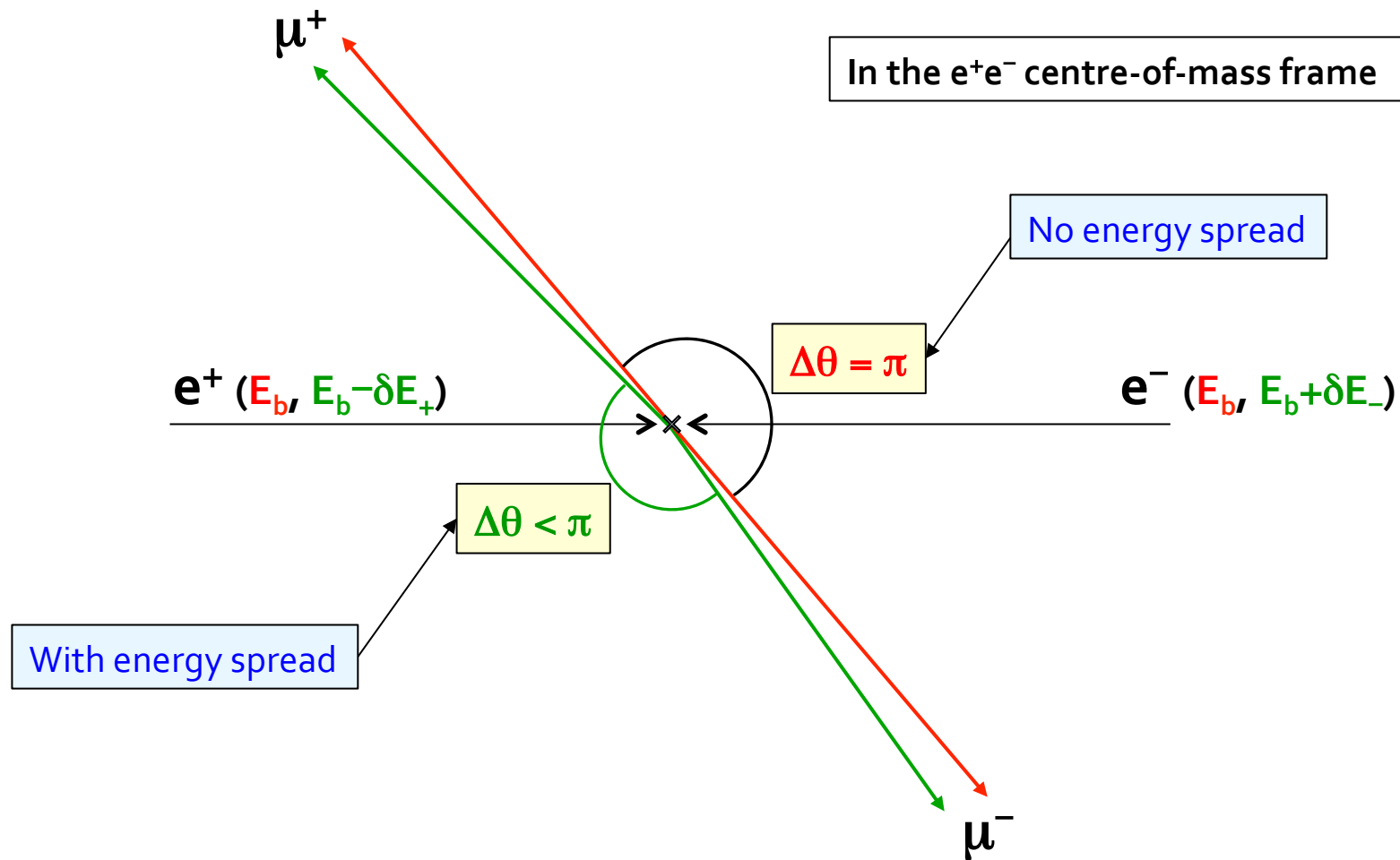
$$\Rightarrow \Delta\alpha_{\text{QED}}/\alpha_{\text{QED}} \sim 10^{-3}$$

Stat: 3×10^{-5}

- **Need to know the energy spread to better than 1% at $\sqrt{s_{\pm}}$**
 - ◆ To limit the effect on α_{QED} to $< 10^{-5}$ (same as the effect from the \sqrt{s} knowledge)
 - Similar constraint as for the Z lineshape

Make use of $e^+e^- \rightarrow \mu^+\mu^-$ events

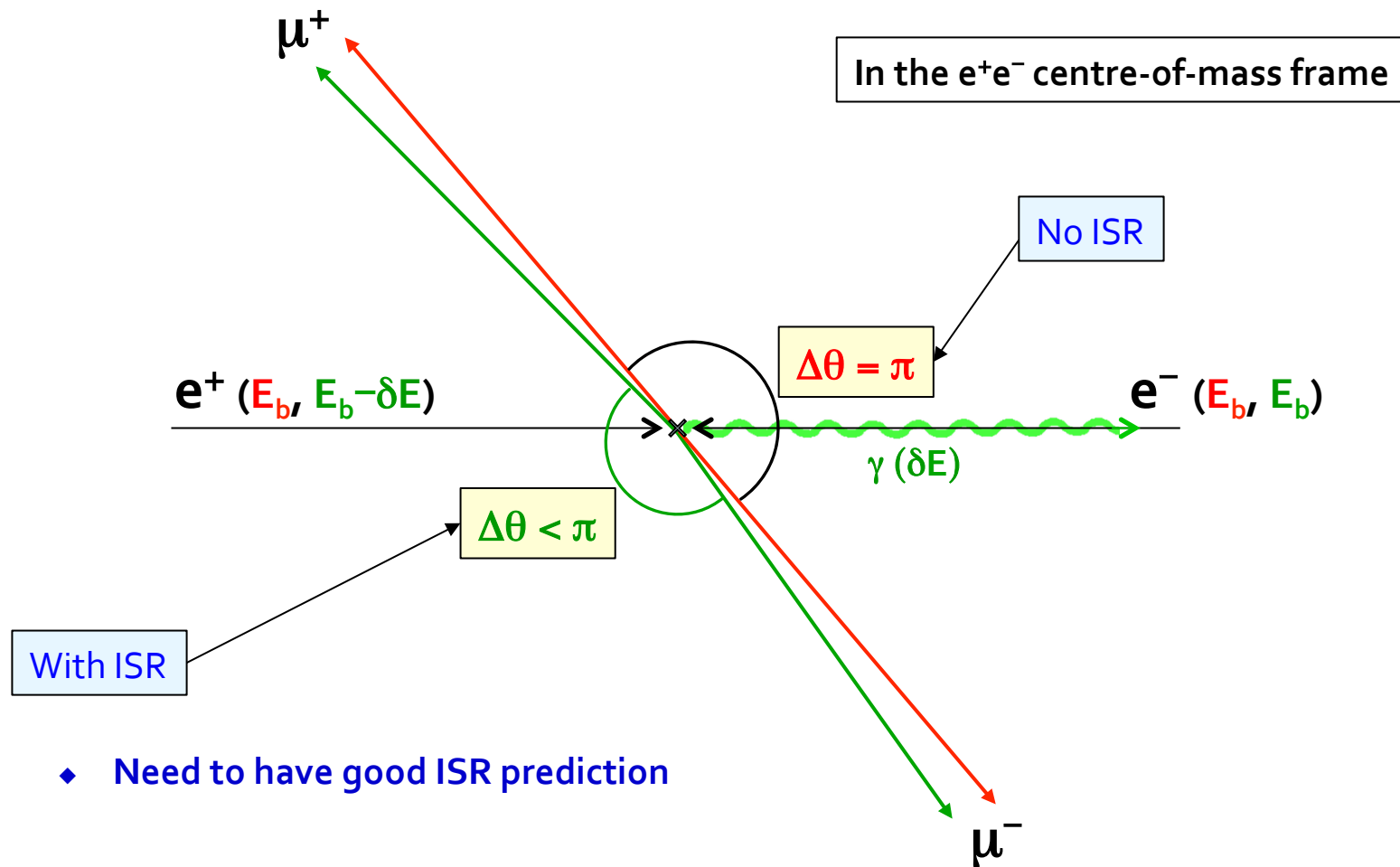
- How are the events modified with energy spread (in the c.o.m.) ?



- Generates a longitudinal boost in the centre-of-mass frame

Make use of $e^+e^- \rightarrow \mu^+\mu^-$ events

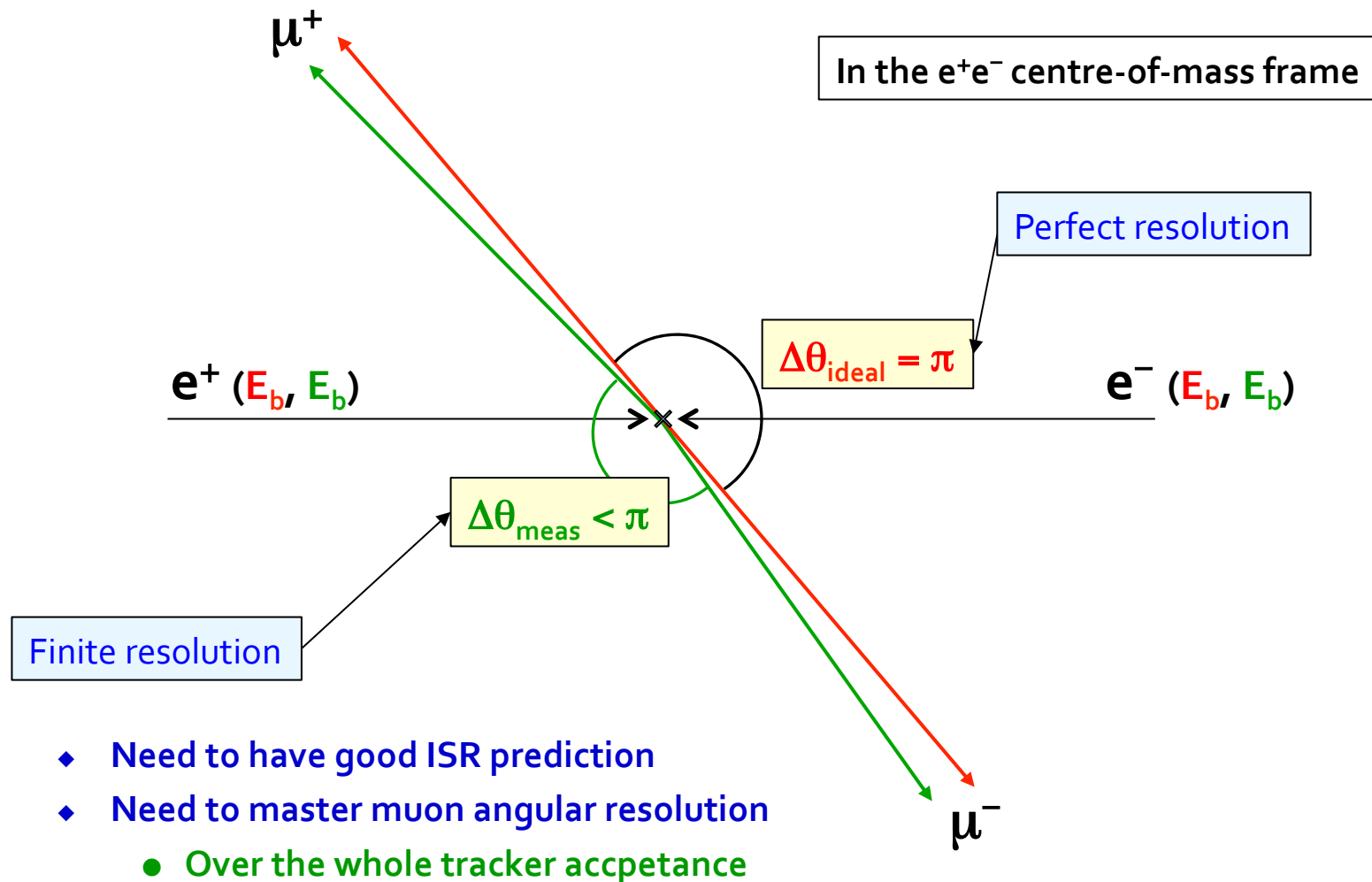
- Competes with initial state radiation (that you cannot get rid of)



- ◆ Need to have good ISR prediction

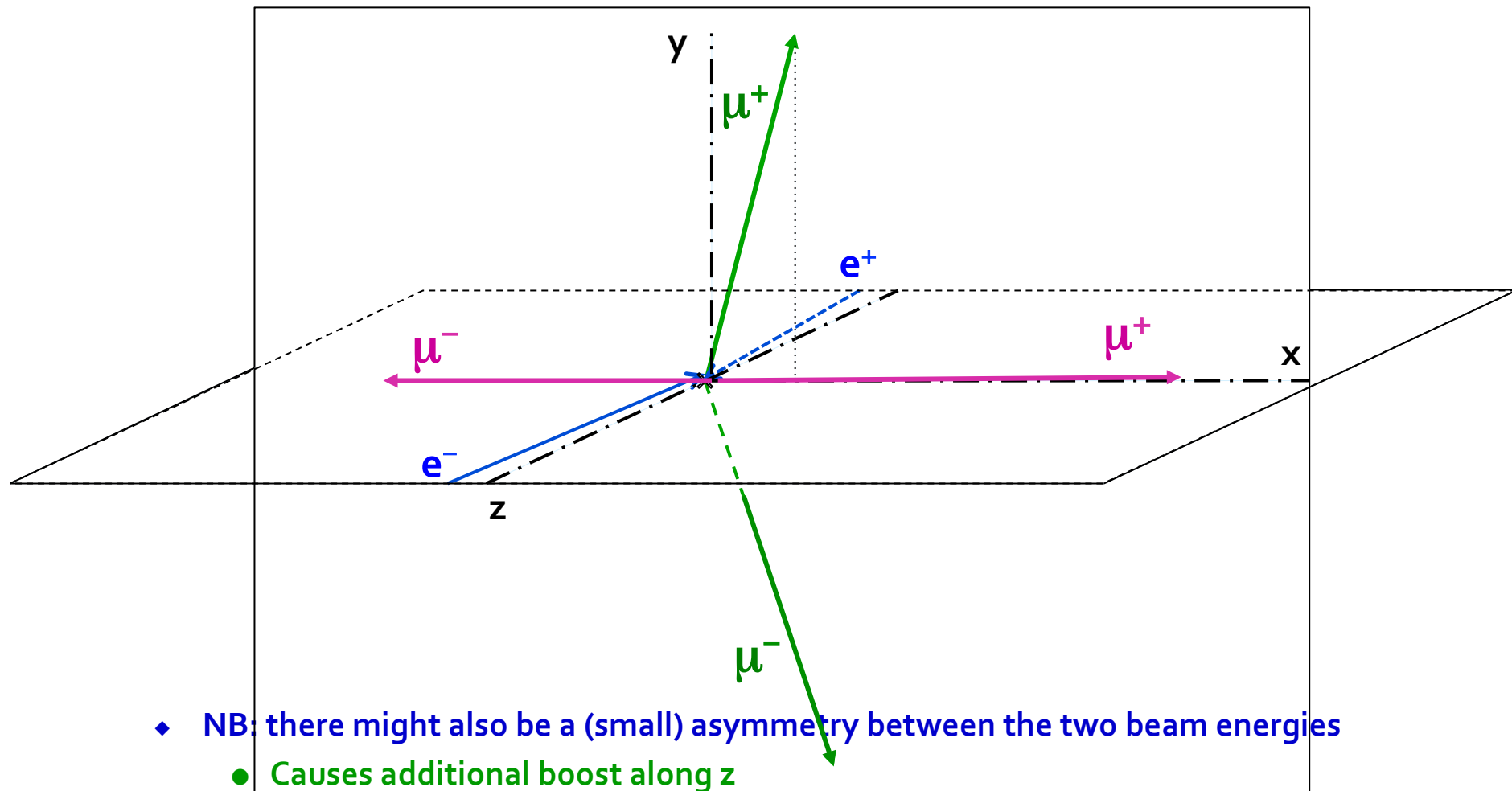
Make use of $e^+e^- \rightarrow \mu^+\mu^-$ events

- Also competes with muon angular resolution ...



The FCC-ee frame is not the e^+e^- c.o.m frame

- An $e^+e^- \rightarrow \mu^+\mu^-$ event produced at $\theta=\pi/2$, with $\phi=0$ or $\phi=\pi/2$



Total energy-momentum conservation

- Assuming one photon emitted along one of the two beams

$$\begin{aligned} E^+ \sin \theta^+ \cos \varphi^+ + E^- \sin \theta^- \cos \varphi^- + |p_z^\gamma| \tan \alpha/2 &= \sqrt{s} \tan \alpha/2, \\ E^+ \sin \theta^+ \sin \varphi^+ + E^- \sin \theta^- \sin \varphi^- &= 0, \\ E^+ \cos \theta^+ + E^- \cos \theta^- + p_z^\gamma &= 0, \\ E^+ + E^- + |p_z^\gamma| / \cos \alpha/2 &= \sqrt{s} / \cos \alpha/2, \end{aligned}$$

- Where E^\pm are the measured energies of the μ^\pm
- Where α is the beam crossing angle (nominal : 30 mrad),
- Where the z axis is the bisector of the two beam axes,
- Where the two beam axes form the (x,z) plane,
- Where θ^\pm are measured with respect to the z axis in the FCC-ee frame,
- Where φ^\pm are measured with to the x axis in the plane transverse to the z axis,
- Where \sqrt{s} is the centre-of-mass energy of the collision

Longitudinal boost $x_\gamma = p_z(\gamma)/\sqrt{s}$

- Solve (E,p) conservation equations for x_γ

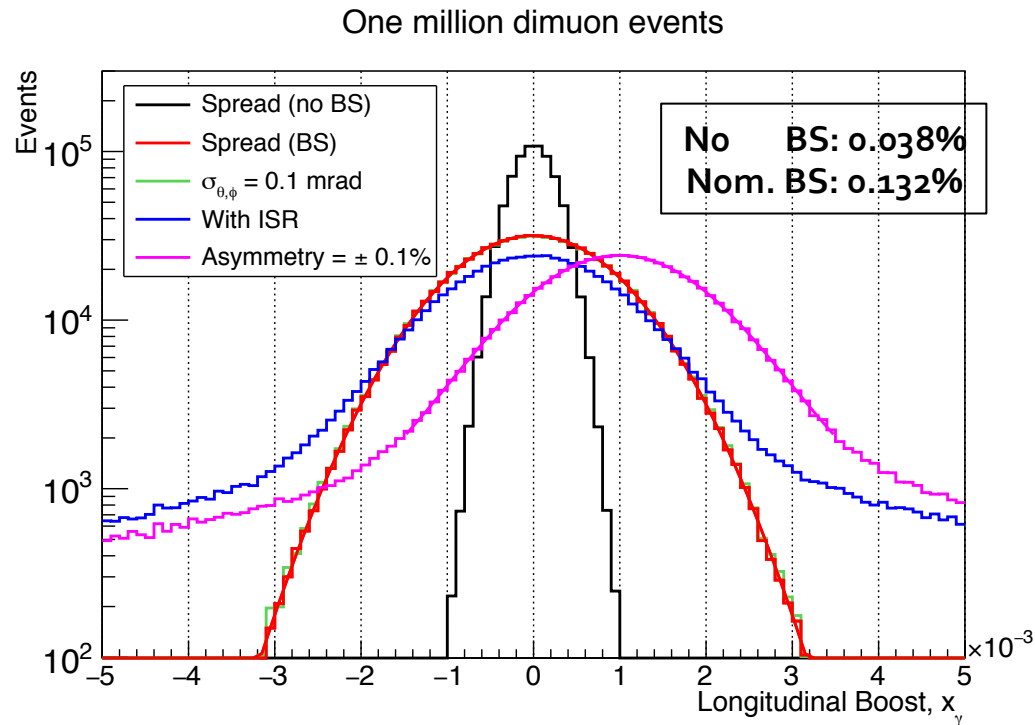
$$x_\gamma = -\frac{x_+ \cos \theta^+ + x_- \cos \theta^-}{\cos(\alpha/2) + |x_+ \cos \theta^+ + x_- \cos \theta^-|}$$

$$\text{With } x_\pm = \frac{\mp \sin \theta^\mp \sin \varphi^\mp}{\sin \theta^+ \sin \varphi^+ - \sin \theta^- \sin \varphi^-}$$

- ◆ As a function of the muon angles only (usually better determined than momenta)
 - Assume initially that α and the “natural” (x,y,z) axes are known perfectly
 - ➔ z axis = bisector of the two beam axes
 - ➔ (x,z) plane = plane that contains the two beam axes
 - ➔ y = axis going upwards perpendicularly to that plane
 - Assume that these axes match the local (detector) axes

Distribution of x_γ

- With 10^6 dimuon events (every 5 minutes at the Z pole)



- Returns the \sqrt{s} spectrum (with asymmetry) for perfect angular resolution and no ISR
 - An angular resolution of 0.1 mrad seems adequately small (typical of CLD / IDEA)
 - ISR slightly degrades the Gaussian core and add tails, needs to be unfolded
- Returns beam energy relative asymmetry with 10^{-6} precision

Effect of $\sigma_{\theta,\phi}$ and ISR

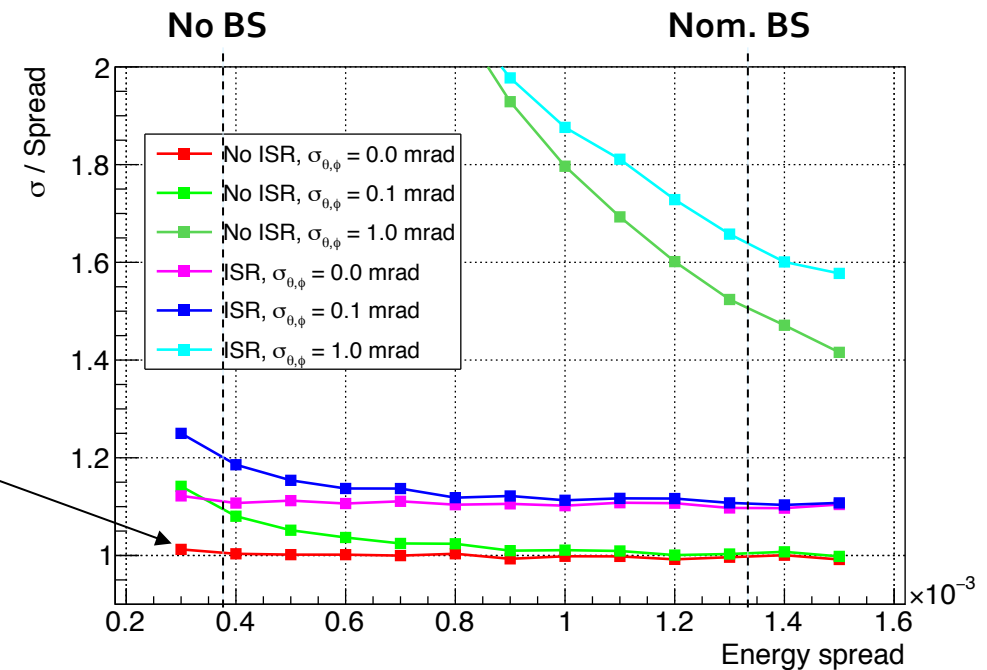
- Fit the core of the x_γ distribution with a Gaussian and compare to E_{spread}

- With 10^6 dimuons
- With $\sigma_{\theta,\phi} = 0$ and no ISR

$$\sqrt{2}\sigma = (0.03804 \pm 0.00004)\% \quad \text{No BS}$$

$$\sqrt{2}\sigma = (0.13185 \pm 0.00011)\% \quad \text{Nom. BS}$$

- ~0.1% relative precision
- ➔ BINGO !



- Effect of ISR is to increase the x_γ width by 10%
 - Need to be known with a precision better than 1% for a per-mil precision on E_{spread}
- Effect of $\sigma_{\theta,\phi} = 0.1$ mrad is to increase the x_γ width by 0.5% for nominal BS
 - Need to be known with a precision of ~10% over the whole tracker acceptance
- Effect of $\sigma_{\theta,\phi} = 1$ mrad is devastating (similar to energy spread, difficult to control)

Control the angular resolution to 0.01 mrad ?

- **Q: How to measure the angular resolution to 10% or better**
 - ◆ For any value of θ and ϕ ?

- **A: Take a muon track in dimuon events**
 - ◆ Refit it with the odd hits, on the one hand, and with the even hits, on the other
 - And compare the angles
 - ◆ Need only 100 tracks in each (θ, ϕ) bin for a 10% precision
 - 10^6 dimuon events = 5 minutes at the Z pole = bins of 3×3 (mrad)²
 - ◆ Expected to be stable in time
 - Precision (or bin size) improves with dimuon statistics

Summary (extract from the CDR)

- **A three-point scan of the Z resonance (87.8, 91.2, 93.9 GeV)**
 - ◆ Is adequate for the measurement of m_Z , Γ_Z , $\sin^2\theta_W^{\text{eff}}$ and α_{QED} with target precisions
 - If the beam energy spread δE is known to a few per mil (for Γ_Z and α_{QED})
- **Dimuon events offer a permanent monitoring of δE**
 - ◆ With adequate precision, every few minutes
 - And with the same as those used to measure the EW observables

Pseudo Observable	Γ_Z			$\alpha_{\text{QED}}(m_Z^2)$		Γ_W	Γ_{top}
Acceptable error	35 keV			10^{-5}		0.5 MeV	9 MeV
\sqrt{s} (GeV)	87.9	91.2	93.8	87.9	93.8	161	350
$\sigma(\delta E)/\delta E$	0.8%	0.2%	0.8%	0.7%		7%	15%
$N_{e^+e^- \rightarrow \mu^+\mu^-}$	$5 \cdot 10^4$	$8 \cdot 10^5$	$5 \cdot 10^4$	$6.5 \cdot 10^4$		650	150
L ($10^{34} \text{cm}^{-2} \text{s}^{-1}$)	230					32	1.8
$\sigma_{\mu\mu}$ (pb)	185	1450	460	185	460	4.0	0.8
Dimuon rate (Hz)	425	3325	1050	425	1050	1.3	0.015
Time needed	2 min	4 min	< 1 min	3 min	1 min	8 min	2 h 30

Notes added

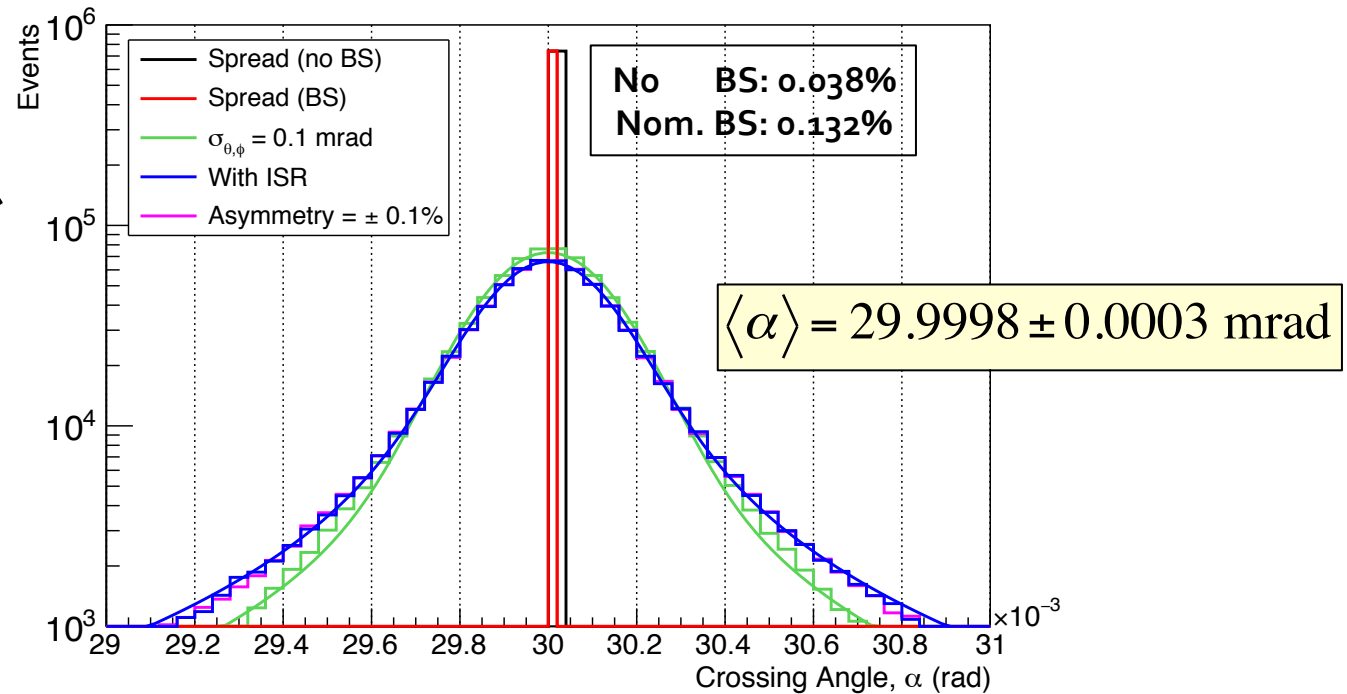
- **All these numbers are valid for nominal luminosities / energy spread**
 - ◆ For smaller energy spread, the precision requirements are quadratically less stringent
- **Example: 0.066% instead of 0.132% at the Z pole (still BS dominated)**
 - ◆ Precision required increases by a factor 4 to 0.8% instead of 0.2%
 - Number of events needed decrease by a factor of 16 to 5×10^4 instead of 8×10^5 .
 - ◆ In the same time, the luminosity is reduced much less, by a factor $2^{2/3} \sim 1.58$
 - If the E spread decrease is achieved by increasing the vertical β^*
 - ➔ Keeping beam currents and vertical emittance unchanged
 - ◆ Therefore the time needed decreases to 25 seconds instead of 4 minutes
- **All times displayed in the table are absolute maxima**
- **All these numbers were derived from the muon angles only**
 - ◆ Angular resolution must be better than 0.1 mrad, + measured to better than 0.01 mrad
 - Angular biases must be kept below 0.01 mrad too (absolute alignment)
 - ◆ To be studied: (probably marginal) influence of muon momenta

Beam crossing angle determination

- With 10^6 dimuon events (every 5 minutes at the Z pole)

$$\alpha = 2 \arcsin \left[\frac{\sin(\varphi^- - \varphi^+) \sin \theta^+ \sin \theta^-}{\sin \varphi^- \sin \theta^- - \sin \varphi^+ \sin \theta^+} \right]$$

One million dimuon events

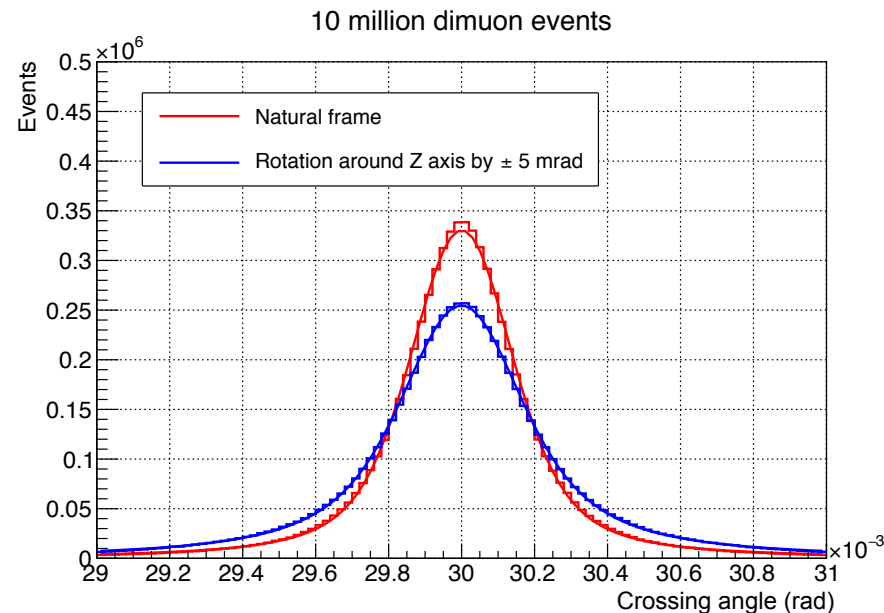


“1000 recipes to use up dimuons”
See also poster

- Spread sensitive to anything happening in the transverse plane
 - ϕ resolution, p_T of emitted photons, and (x, y, z) axes knowledge

Detector alignment

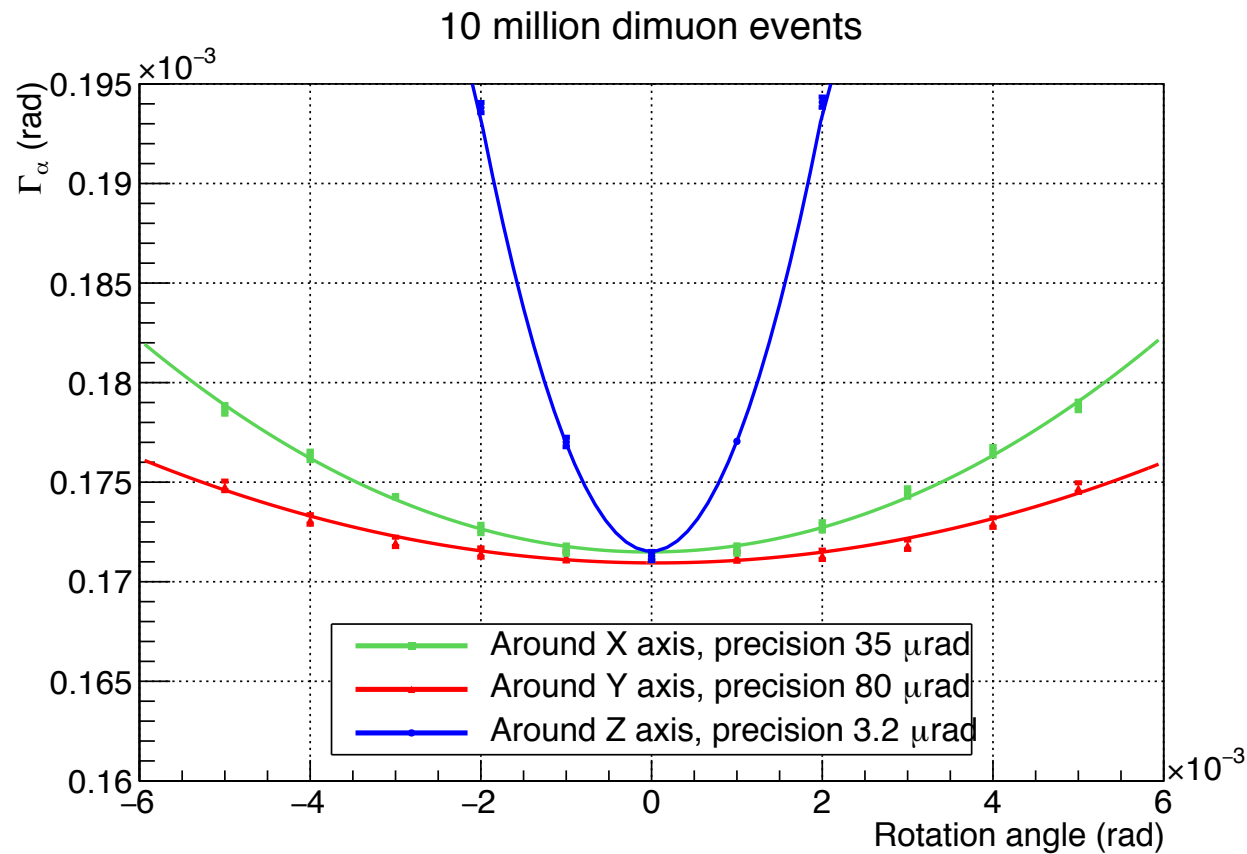
- ❑ **Absolute angle determination is (usually) not an easy task**
 - ◆ Requires alignment of the local (detector) frame with the natural (FCC-ee) frame
 - Z axis = solenoid axis vs bisector of the two beam axes
 - (X,Z) plane = horizontal plane vs plane containing the two beam axes
- ❑ **Spread of α increases with anything happening in the transverse plane**
 - ◆ E.g., rotation around the Z axis changes both X and Y directions



- Similarly, rotation around the X (Y) axis changes Y (X) direction

Detector alignment

- Minimize the spread of the α distribution to find the three Euler angles

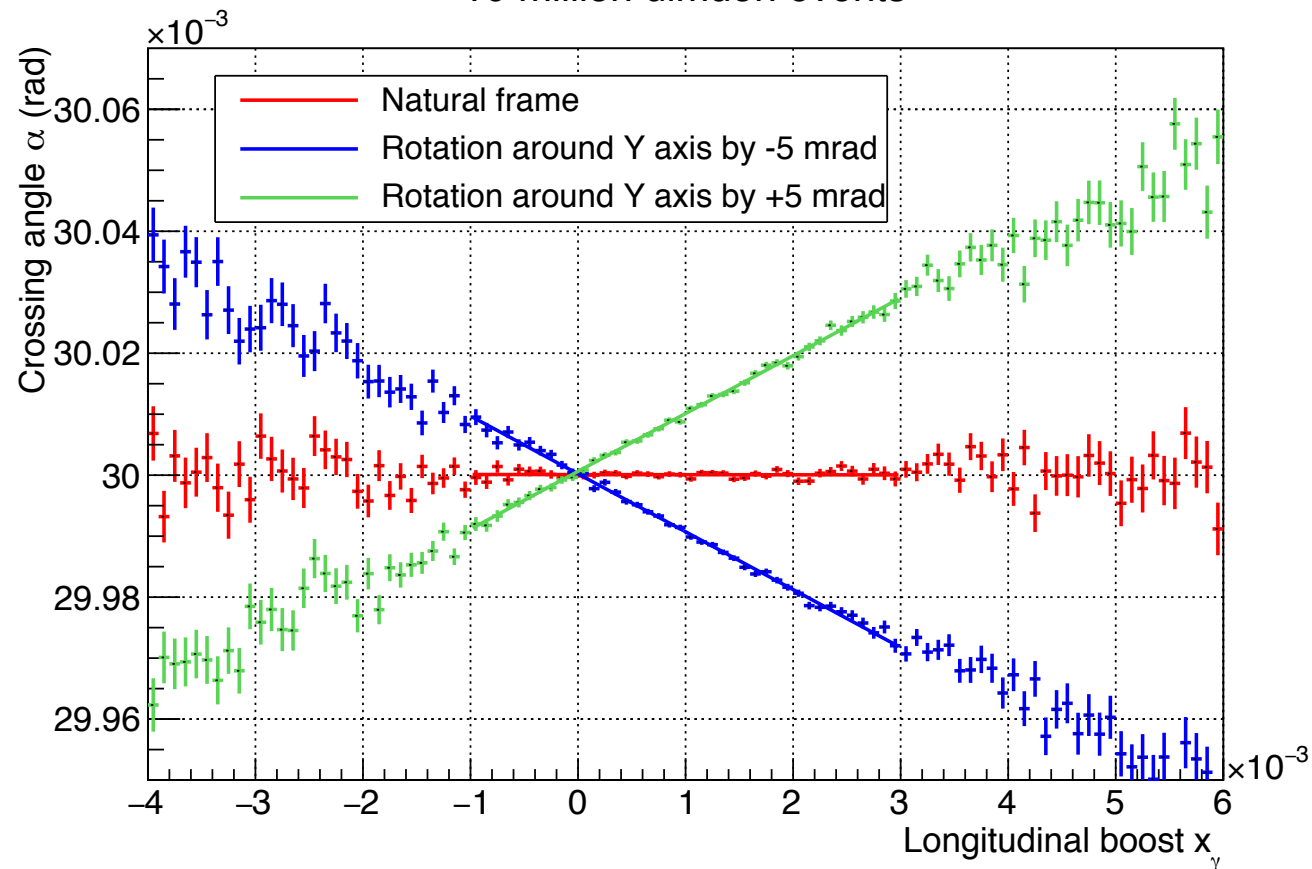


- Note: α spread dominated by the ϕ resolution (here 0.1 mrad)
 - Precisions quadratically improves with the resolution in ϕ (here 0.1 mrad)

Detector alignment

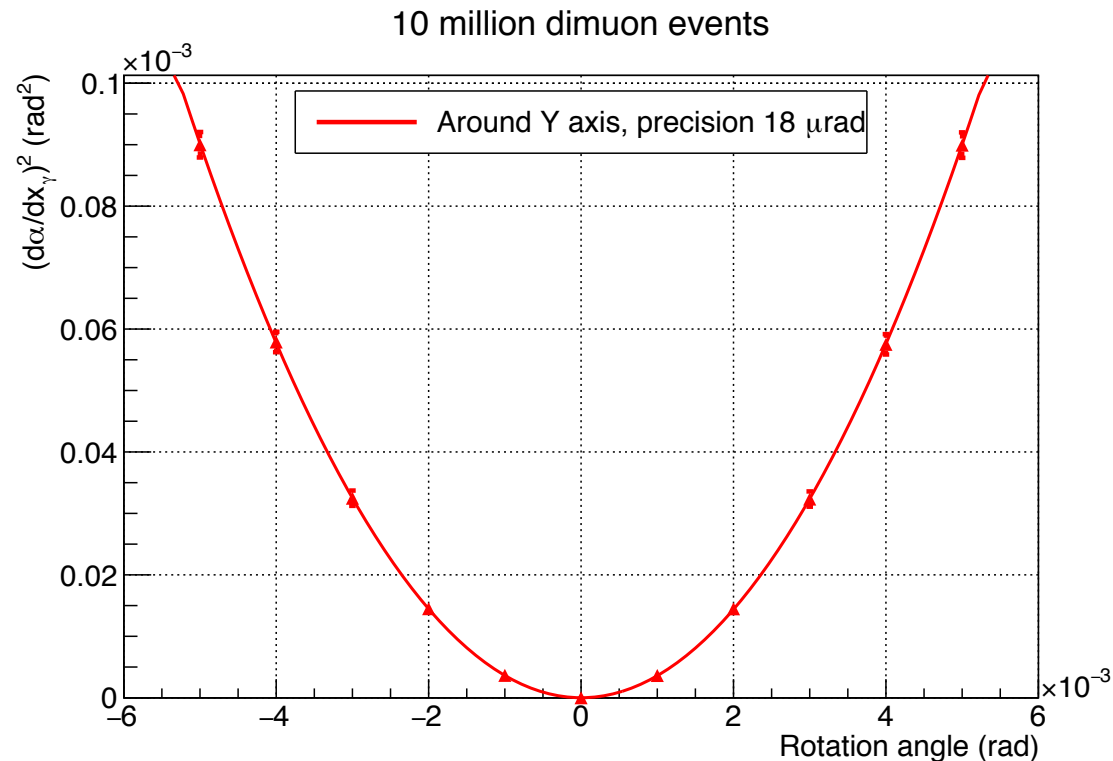
- Improve the angle corresponding to a rotation around the Y axis
 - ◆ X and Z information get mixed by such a rotation
 - Resulting in a strong (linear) correlation between x_y and α :

10 million dimuon events



Detector alignment

- Minimize the correlation between x_γ and α :



- ◆ Improves the precision on that crossing angle by a factor of five.
 - Reach a precision of $0.1 \mu\text{rad}$ on α and of 10^{-7} on x_γ
 - Variation of the x_γ spread already insignificant with 100 times less events

Detector alignment

□ A finite crossing angle is the key

- ◆ A zero crossing angle cannot be determined with precision ($\sin\phi^+ = \sin\phi^-$, $\sin\theta^+ = \sin\theta^-$)

$$\alpha = 2 \arcsin \left[\frac{\sin(\varphi^- - \varphi^+) \sin\theta^+ \sin\theta^-}{\sin\varphi^- \sin\theta^- - \sin\varphi^+ \sin\theta^+} \right]$$

- Spread is infinitely large even with a perfect detector alignment
 - Leads to inaccurate absolute alignment
 - ➔ With corresponding biases on the muon angles
 - ◆ With a finite crossing angle, only events with $\sin\phi^+ \sim \sin\phi^- \sim 0$ need to be rejected
 - Cut $|\sin\phi^\pm| > 0.2$ was applied in all previous plots
- ## □ Possible bonus (to be studied)
- ◆ Built-in energy asymmetry between electrons and positrons