

# Global fits to EW and Higgs observables at the FCC

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# Introduction

## The Future Circular Collider

- In short:

FCC-ee: Experimental PRECISION

FCC-hh: ENERGY

FCC-eh: a good mix of both

- All 3 options can do a lot in terms of improving our understanding of the electroweak sector:
  - Properties of the Electroweak gauge bosons (EWPO)
  - Mass and couplings of the Higgs boson
- The point of this talk:
  - To get an overall idea of what one could get after combining the different collider options ...
  - ... emphasizing complementarities between them (to convince you that they are all important)
- **DISCLAIMER:** All results in this talk are VERY preliminary

# Outline

- **Theoretical framework: The dimension-6 SMEFT**
- **Electroweak precision observables at the FCC**
- **Higgs observables at the FCC**
- **The Global fit to EW and Higgs observables**

# The dimension-6 SMEFT

# The dimension-6 SMEFT

- The dimension 6 SMEFT: Assumes new physics is heavy + decoupling  
Particles and symmetries of the low-energy theory: SM  
Power counting: EFT expansion in canonical dimension of operators

$$\mathcal{L}_{\text{Eff}} = \sum_{d=4}^{\infty} \frac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

$$\mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \quad [\mathcal{O}_i] = d \longrightarrow \left(\frac{q}{\Lambda}\right)^{d-4}$$

Effects suppressed by  $q = v, E < \Lambda$

$\Lambda$ : Cut-off of the EFT

- LO new physics effects “start” at dimension 6
- With current precision, and assuming  $\Lambda \sim \text{TeV}$ , sensitivity to  $d > 6$  is small

$$\frac{M_Z^2}{(1\text{TeV})^2} \sim 0.8\% \quad \frac{M_Z^4}{(1\text{TeV})^4} \sim 0.007\%$$

Truncate at  $d=6$ : 59 types of operators (2499 counting flavor)

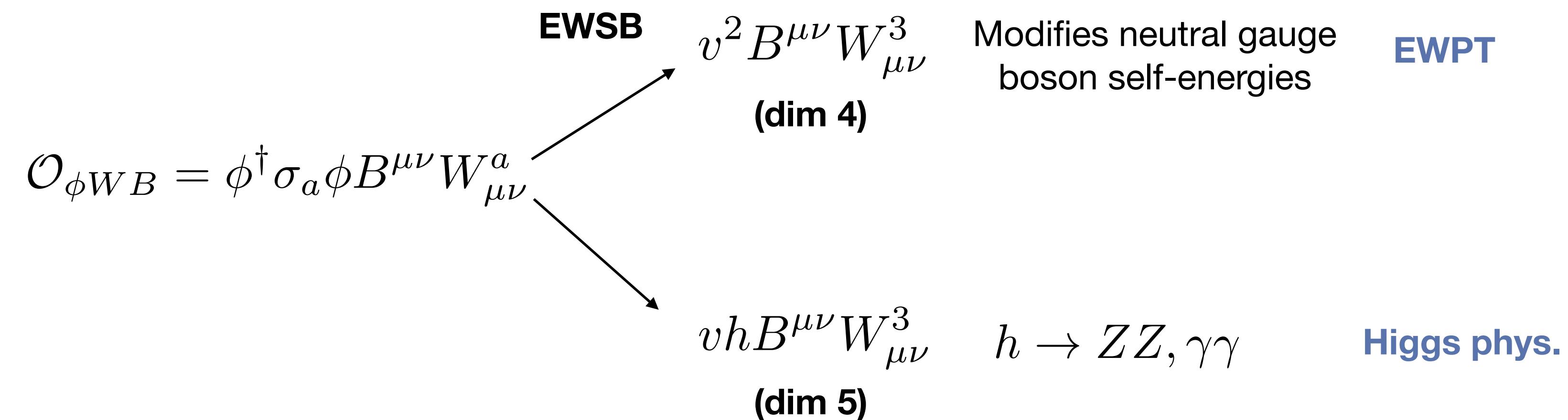
W. Buchmüller, D. Wyler, Nucl. Phys. B268 (1986) 621  
C. Arzt, M.B. Einhorn, J. Wudka, Nucl. Phys. B433 (1995) 41

► B.Grzadkowski, M.Iskrynski, M.Misiak, J.Rosiek, JHEP 1010 (2010) 085

First complete basis, aka Warsaw basis

# The dimension-6 SMEFT

- Advantages of EFTs:
  - Completely model-independent description of new physics  
(Consistent with assumptions of SM at low energies)
  - Well-defined perturbative expansion (can compute at  $N^n\text{LO}$ )
  - Well-defined way of connecting with explicit UV completions via matching/integrating out heavy degrees of freedom
  - Describes correlations of new physics effects in different types of observables, e.g.



# FCC sensitivity to New Physics

## General strategy for calculation of future sensitivities

- Fit to new physics effects parameterized by the dimension 6 SMEFT:
  - Bayesian fit using 
  - **FCC sensitivity:** from posterior info (NP parameter errors/limits)
- Assumptions:
  - **Likelihood:** SM predictions as central values for future “experimental” measurements. Errors given by projected experimental uncertainties. (For comparison, results using current data are also shown assuming SM central values, with current uncertainties.)
  - **SM theory uncertainties:** SM intrinsic and parametric uncertainties reduced according to future projections. Included in the analysis when available.
  - **New physics effects:** Working at the linear-level in the EFT effects (interference with SM amplitudes)

$$O = O_{\text{SM}} + \delta O_{\text{NP}} \frac{1}{\Lambda^2}$$

# Electroweak precision measurements at the FCC

# Electroweak precision observables at FCC ee

- Electroweak precision measurements at FCC-ee

## Precision measurements at the Z pole

See R. Tenchini's talk for details

Observable	Expected uncertainty	(Relative uncertainty)
$M_Z$ [GeV]	$10^{-4}$	$(10^{-6})$
$\Gamma_Z$ [GeV]	$10^{-4}$	$(4 \times 10^{-5})$
$\sigma_{\text{had}}^0$ [nb]	$5 \times 10^{-3}$	$(10^{-4})$
$R_e$	0.006	$(3 \times 10^{-4})$
$R_\mu$	0.001	$(5 \times 10^{-5})$
$R_\tau$	0.002	$(10^{-4})$
$R_b$	0.00006	$(3 \times 10^{-4})$
$R_c$	0.00026	$(15 \times 10^{-4})$

Observable	Expected uncertainty	(Relative uncertainty)
$A_e$	$10^{-4}$	$(7 \times 10^{-4})$
$A_\mu$	$1.5 \times 10^{-4}$	$(10^{-3})$
$A_\tau$	$3 \times 10^{-4}$	$(2 \times 10^{-3})$
$A_b$	$30 \times 10^{-4}$	$(32 \times 10^{-4})$
$A_c$	$80 \times 10^{-4}$	$(12 \times 10^{-3})$
$\sin^2 \theta_{\text{Eff}}^e (P_\tau)$	$6.6 \times 10^{-6}$	$(3 \times 10^{-5})$
<del><math>\sin^2 \theta_{\text{Eff}}^\ell (A_{FB}^\mu)</math></del>	<del><math>5 \times 10^{-6}</math></del>	<del><math>(2 \times 10^{-5})</math></del>

Z-lineshape parameters and normalized partial decay widths

Z-pole left-right asymmetry parameters

Latest Results from FCC-ee CDR

Thanks to Roberto Tenchini and Alain Blondel for info about the FCC-ee EW measurements

Not independent

# Electroweak precision observables at FCC ee

- Electroweak precision measurements at FCC-ee

## Precision measurements at the $WW$ threshold

See P. Azurri's talk for details

Observable	Expected uncertainty	(Relative uncertainty)
$M_W$ [GeV]	$6.5 \times 10^{-4}$	$(8 \times 10^{-6})$
$\Gamma_W$ [GeV]	$1.59 \times 10^{-3}$	$(8 \times 10^{-4})$
$R_{\text{inv}}$	0.002	$(3 \times 10^{-4})$

“First Look at the Physics Case of TLEP”  
JHEP 1401 (2014) 164

Latest Results from FCC-ee CDR

# Electroweak precision observables at FCC ee

- EWPO sensitive to modifications of NC couplings

$$\mathcal{L}_{\text{NC}} = -\frac{e}{sc} (1 + \delta^U g_{\text{NC}}) Z_\mu \sum_\psi \overline{\psi^i} \gamma^\mu \left[ \left( g_L^\psi \delta_{ij} + (\delta^D g_L^\psi)_{ij} \right) P_L + \left( g_R^\psi \delta_{ij} + (\delta^D g_R^\psi)_{ij} \right) P_R + \delta^Q g_{\text{NC}} \delta_{ij} \right] \psi^j$$

## Flavor non-universal contributions

$$\begin{aligned} \delta^D g_L^e &= -\frac{1}{2} \left( C_{\phi l}^{(1)} \mp C_{\phi l}^{(3)} \right) \frac{v^2}{\Lambda^2}, & \delta^D g_R^e &= -\frac{1}{2} C_{\phi e}^{(1)} \frac{v^2}{\Lambda^2} \\ \delta^D g_L^d &= -\frac{1}{2} \left( C_{\phi q}^{(1)} \mp C_{\phi q}^{(3)} \right) \frac{v^2}{\Lambda^2}, & \delta^D g_R^d &= -\frac{1}{4} C_{\phi_d^u}^{(1)} \frac{v^2}{\Lambda^2} \end{aligned}$$

## 10 Operators

$$\begin{aligned} \mathcal{O}_{\phi f}^{(1)} &= (\phi^\dagger i \overleftrightarrow{D}_\mu \phi) (\bar{f} \gamma^\mu f) \\ \mathcal{O}_{\phi f}^{(3)} &= (\phi^\dagger i \overleftrightarrow{D}_\mu^a \phi) (\bar{f} \gamma^\mu \sigma_a f) \end{aligned}$$

## Flavor-universal contributions

$$\delta^U g_{\text{NC}} = -\frac{1}{2} \left[ \Delta_{G_F} + \frac{C_{\phi D}}{2} \right] \frac{v^2}{\Lambda^2}$$

$$\mathcal{O}_{\phi D} = |\phi^\dagger i D_\mu \phi|^2$$

$$\delta^Q g_{\text{NC}} = -Q \left( \frac{sc}{c^2 - s^2} C_{\phi WB} + \frac{s^2 c^2}{c^2 - s^2} \left[ \Delta_{G_F} + \frac{C_{\phi D}}{2} \right] \right) \frac{v^2}{\Lambda^2}$$

$$\mathcal{O}_{\phi WB} = (\phi^\dagger \sigma_a \phi) W_{\mu\nu}^a B^{\mu\nu}$$

## Indirect effect associated to modifications in $\mu$ decay ( $G_F$ )

$$\Delta_{G_F} = \left( C_{\phi l}^{(3)} \right)_{22} + \left( C_{\phi l}^{(3)} \right)_{11} - (C_{ll})_{1221}$$

$$\mathcal{O}_{ll} = (\bar{l} \gamma_\mu l)(\bar{l} \gamma^\mu l)$$

# Electroweak precision observables at FCC ee

- EWPO sensitive to modifications of CC couplings (ignoring CKM effects)

$$\mathcal{L}_{\text{CC}} = -\frac{e}{\sqrt{2}s} (1 + \delta^U g_{\text{CC}}) W_\mu^+ \left[ \left( \delta_{ij} + (\delta^D U_L)_{ij} \right) \bar{\nu}_L^i \gamma^\mu e_L^j + (\delta^D V_R)_{ij} \bar{u}_R^i \gamma^\mu d_R^j + \left( \delta_{ij} + (\delta^D V_L)_{ij} \right) \bar{u}_L^i \gamma^\mu d_L^j \right] + \text{h.c.}$$

## Flavor non-universal contributions

$$\delta^D U_L = C_{\phi l}^{(3)} \frac{v^2}{\Lambda^2},$$

Does not interfere with SM

$$\delta^D V_L = C_{\phi q}^{(3)} \frac{v^2}{\Lambda^2}, \quad \delta^D V_R = \frac{1}{2} C_{\phi ud} \frac{v^2}{\Lambda^2}$$

## Operators

$$\mathcal{O}_{\phi ud} = (\tilde{\phi}^\dagger i D_\mu \phi)(\bar{u}_R \gamma^\mu d_R)$$

## Flavor-universal contributions

$$\delta^U g_{\text{CC}} = \left[ \frac{sc}{s^2 - c^2} C_{\phi WB} - \frac{c^2}{2(c^2 - s^2)} \left( \Delta_{G_F} + \frac{C_{\phi D}}{2} \right) \right] \frac{v^2}{\Lambda^2}$$

No more operators but constraints  
1 more direction

- W mass:

$$M_W^2 = M_Z^2 c^2 \left( 1 - \frac{c^2}{c^2 - s^2} \left( \frac{C_{\phi D}}{2} + \frac{2s}{c} C_{\phi WB} + \frac{s^2}{c^2} \Delta_{G_F} \right) \frac{v^2}{\Lambda^2} \right)$$

## EWPO: Z-pole + W properties

Constrain 8 independent  
combinations (in the FU case)

# Electroweak precision observables at FCC ee

- Electroweak precision measurements at FCC-ee of SM inputs

## Precision measurements of Top and Higgs masses

Observable	Expected uncertainty	(Relative uncertainty)
$m_t$ [GeV]	$50 \times 10^{-3}$	$(3 \times 10^{-4})$
$M_H$ [GeV]	$7 \times 10^{-3}$	$(6 \times 10^{-5})$

“First Look at the Physics Case of TLEP”  
JHEP 1401 (2014) 164

Only ~10 MeV experimental

## Precision measurements of the EM constant from $A_{FB}^{\mu\mu}$ above and below the Z pole

See note in P. Janot’s talk for details

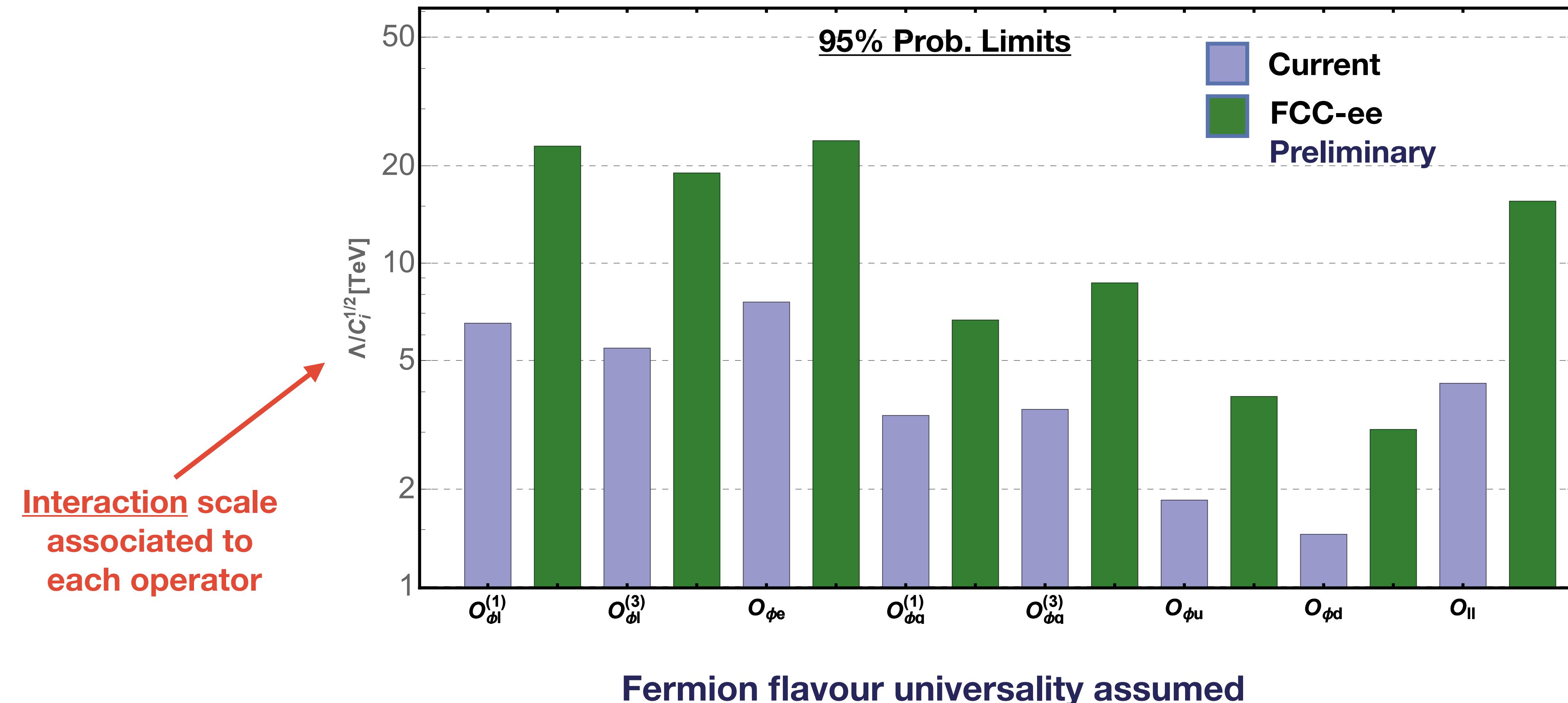
$$\frac{\delta \alpha_{\text{QED}}^{-1}(M_Z^2)}{\alpha_{\text{QED}}^{-1}(M_Z^2)} \sim 3 \times 10^{-5}$$

With a 1-year running period at 87.9/94.3 GeV (85 ab<sup>-1</sup>):

Latest Results from FCC-ee CDR

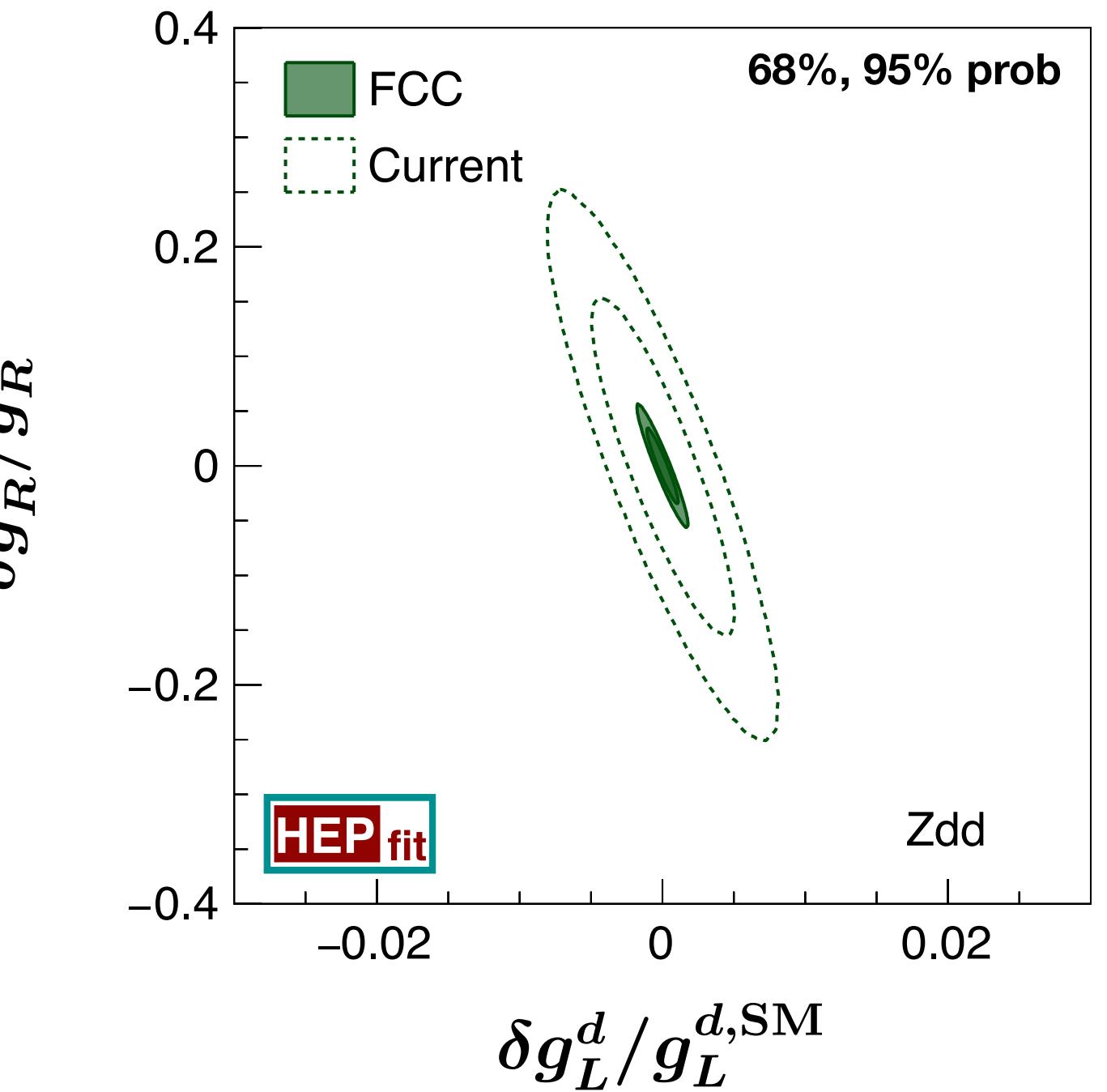
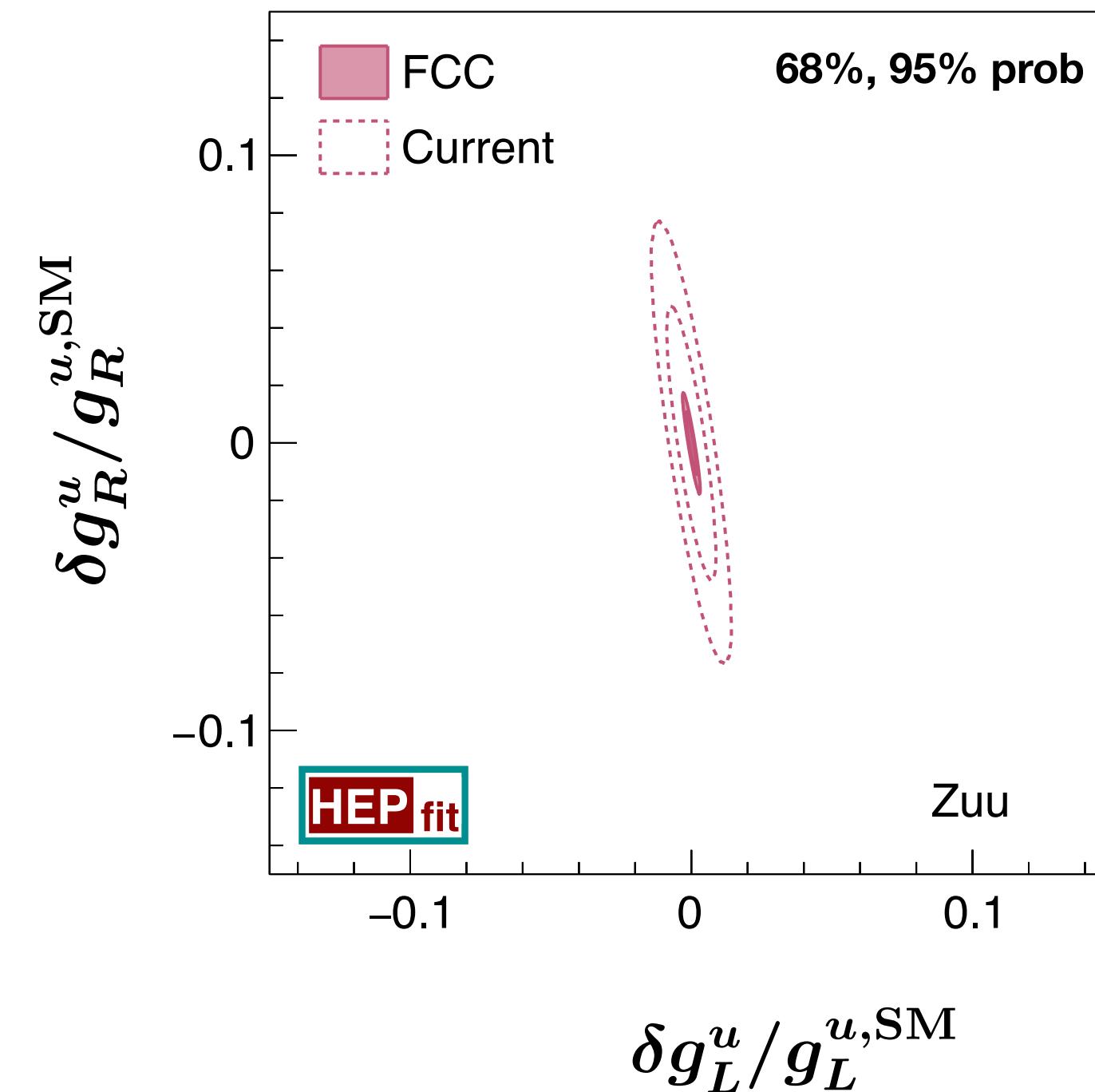
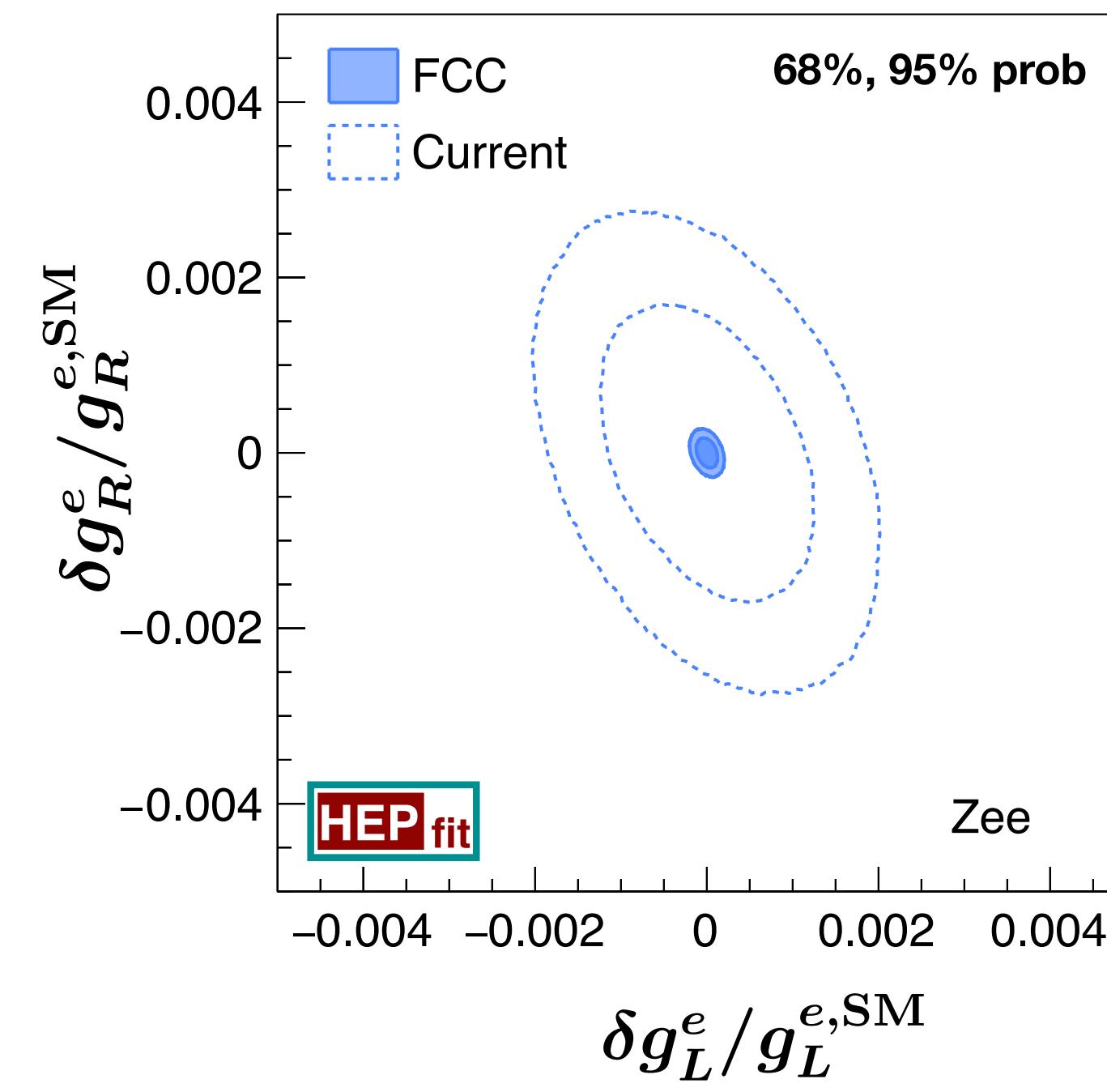
# Electroweak precision observables at FCC ee

- Global fit to electroweak precision measurements at FCC-ee



# Electroweak precision observables at FCC ee

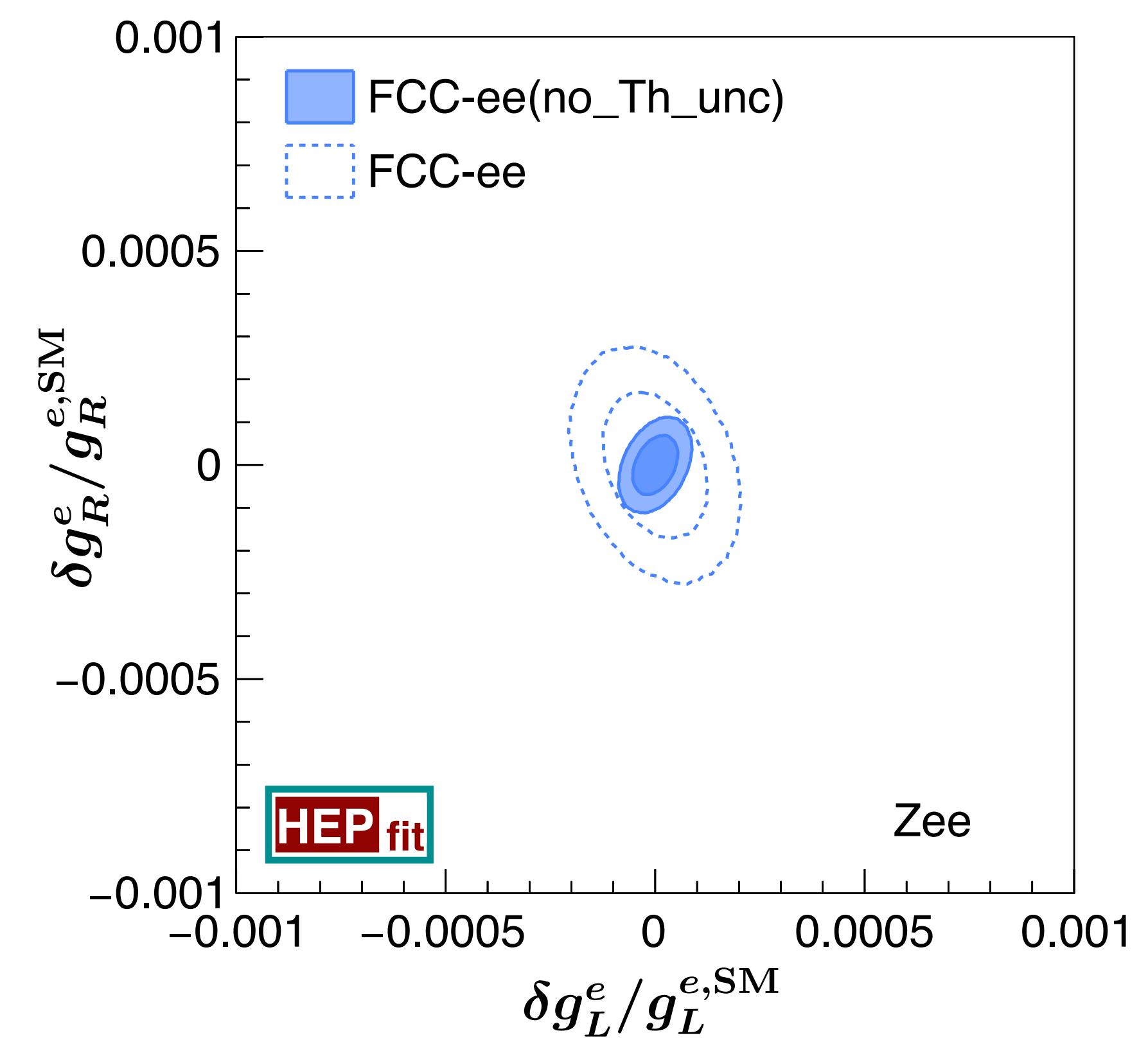
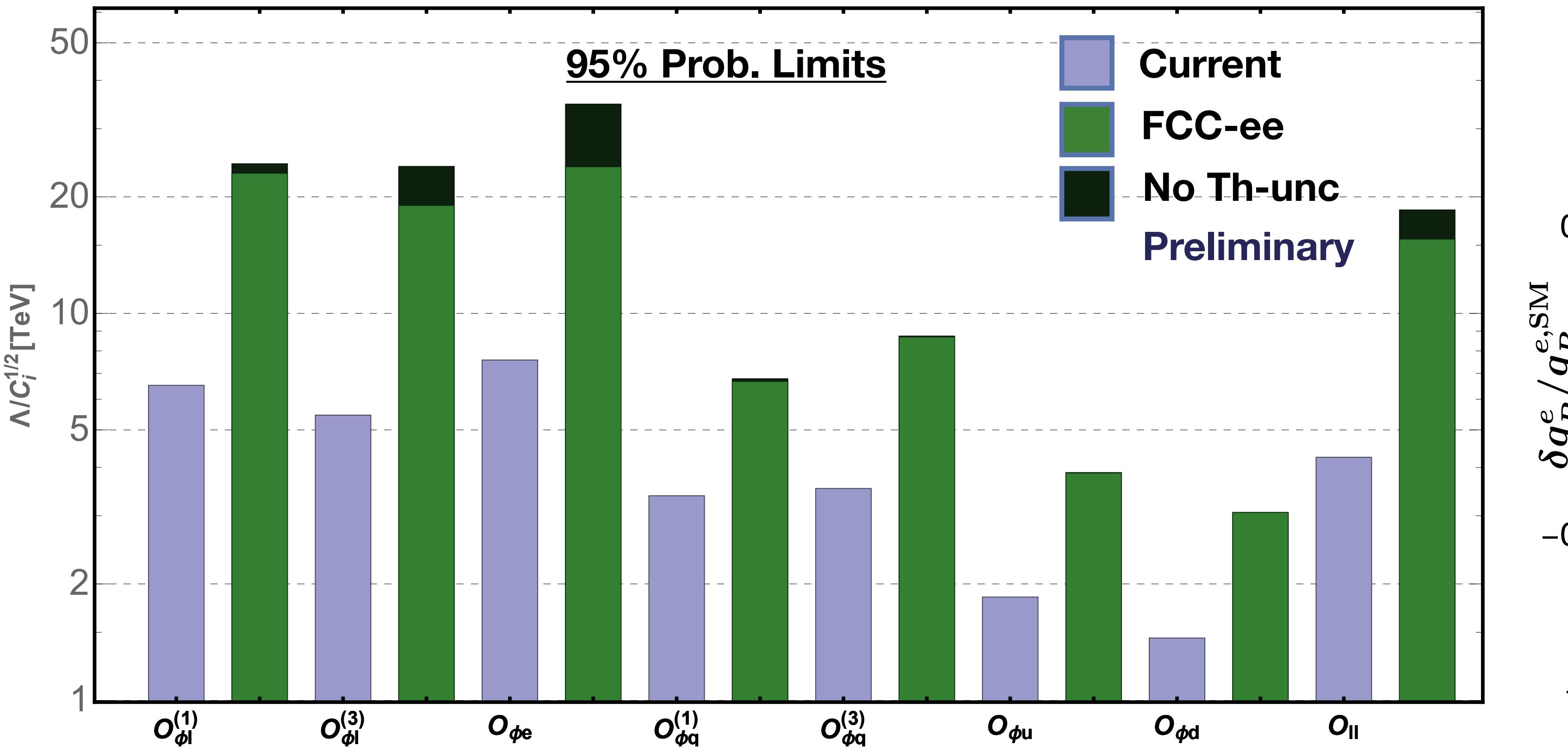
- Global fit to electroweak precision measurements at FCC-ee



Fermion flavour universality assumed

# Electroweak precision observables at FCC ee

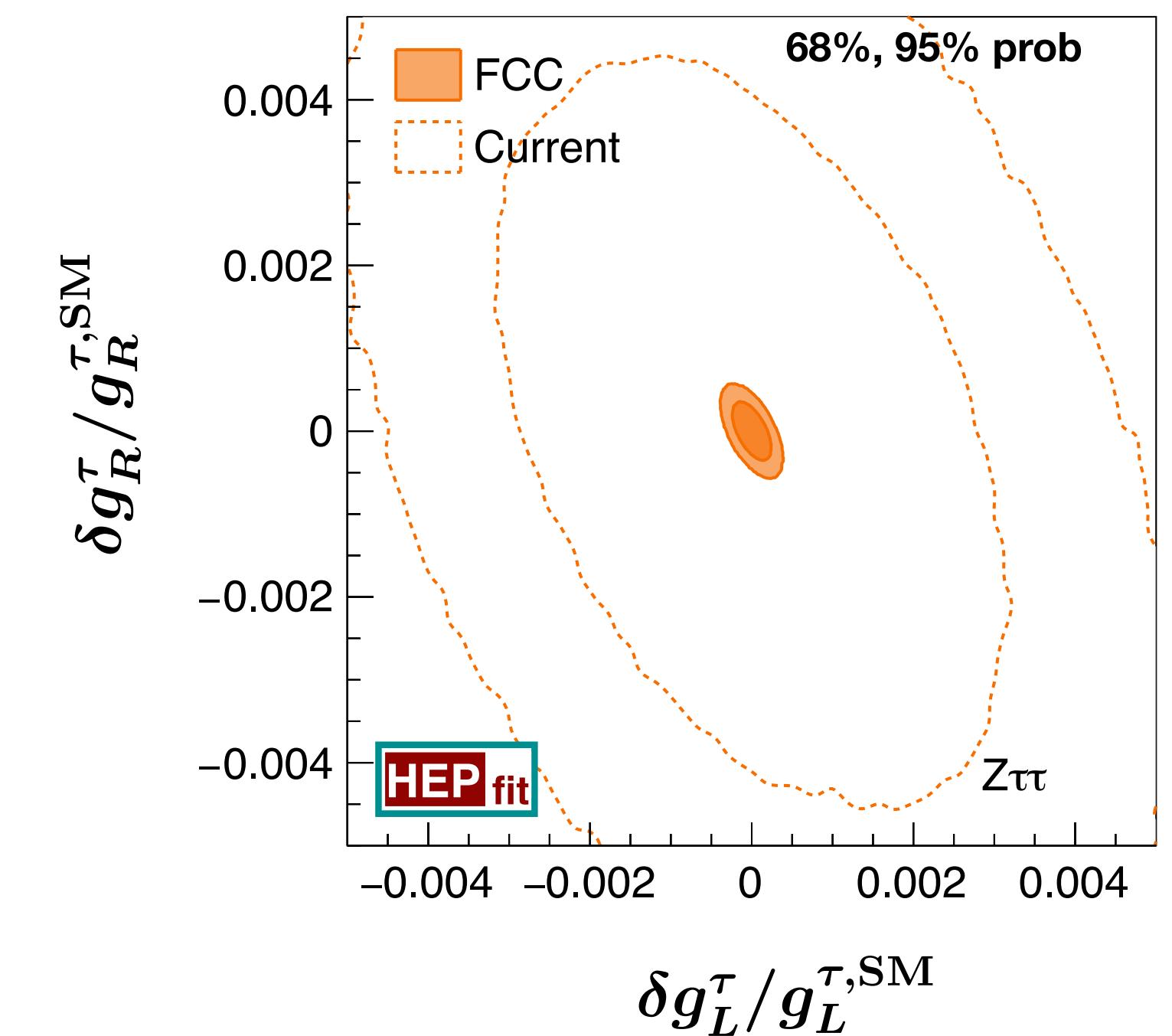
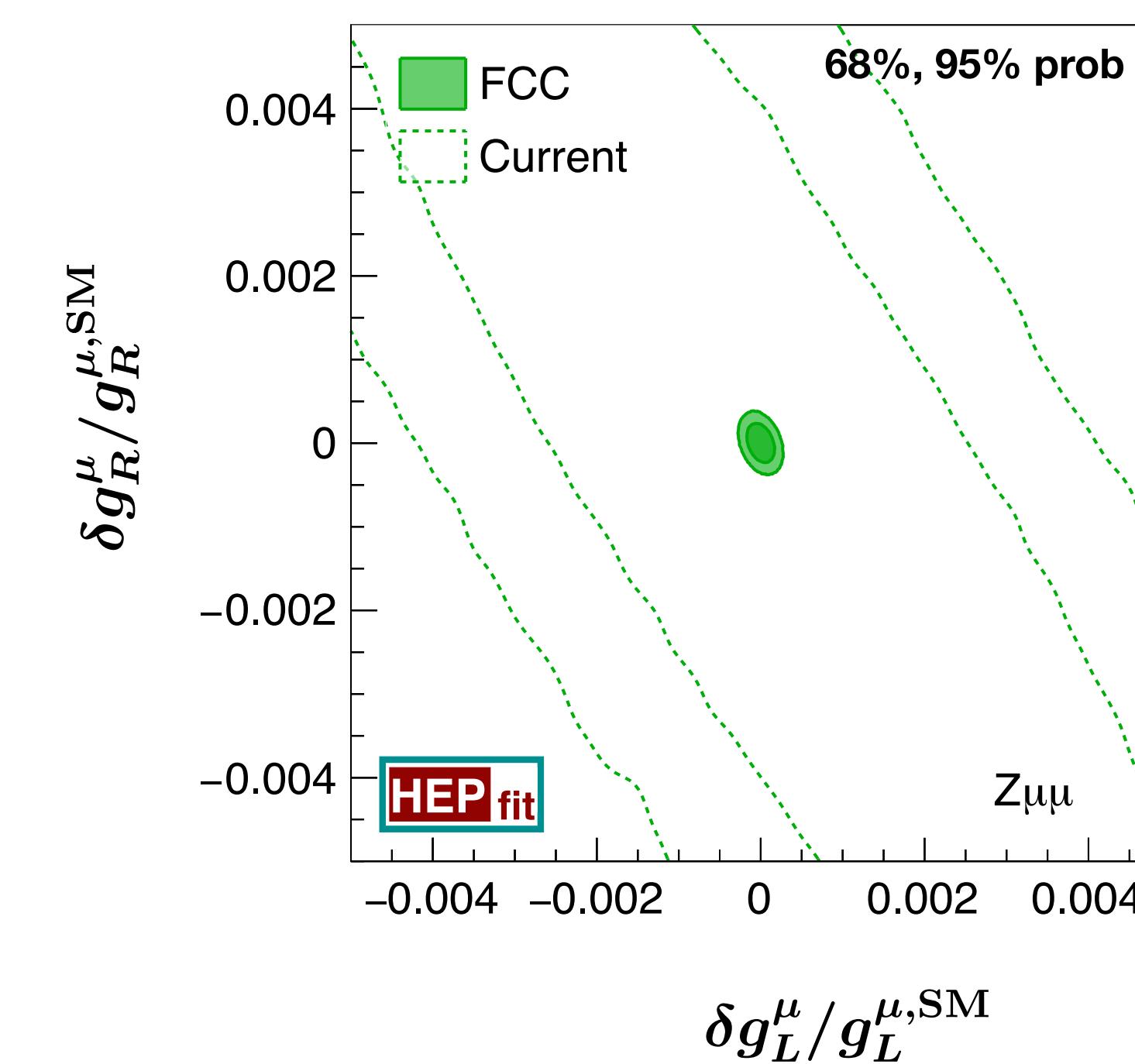
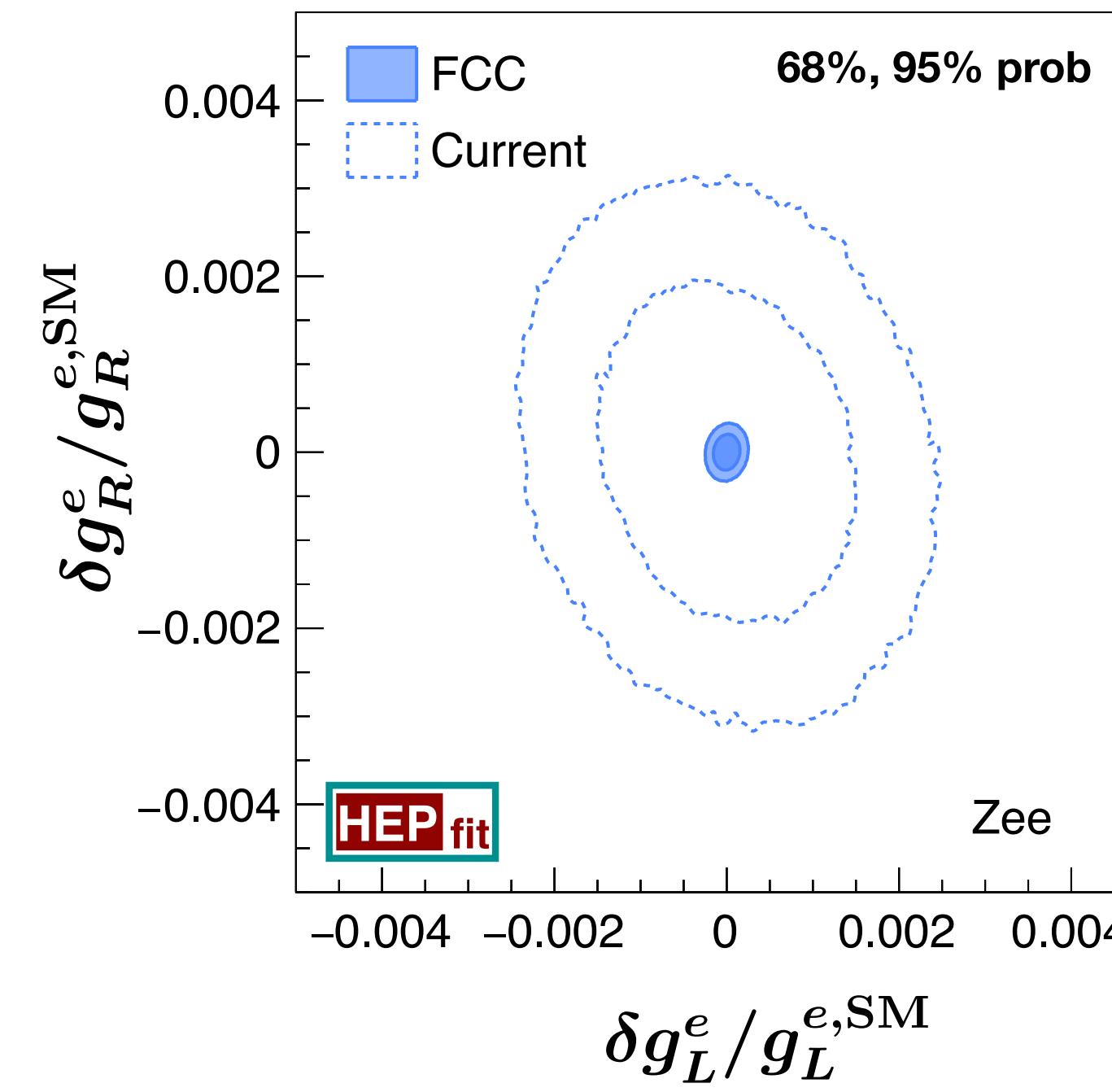
- Global fit to electroweak precision measurements at FCC-ee: Impact of theory uncertainties



Fermion flavour universality assumed

# Electroweak precision observables at FCC ee

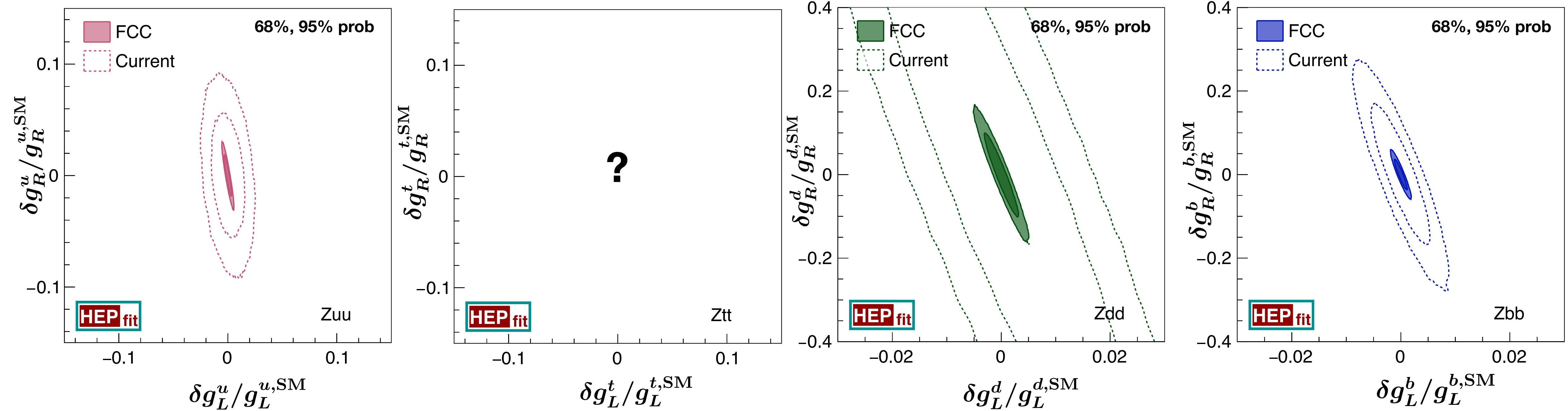
- Global fit to electroweak precision measurements at FCC-ee



No Lepton flavour universality assumed

# Electroweak precision observables at FCC ee

- Global fit to electroweak precision measurements at FCC-ee



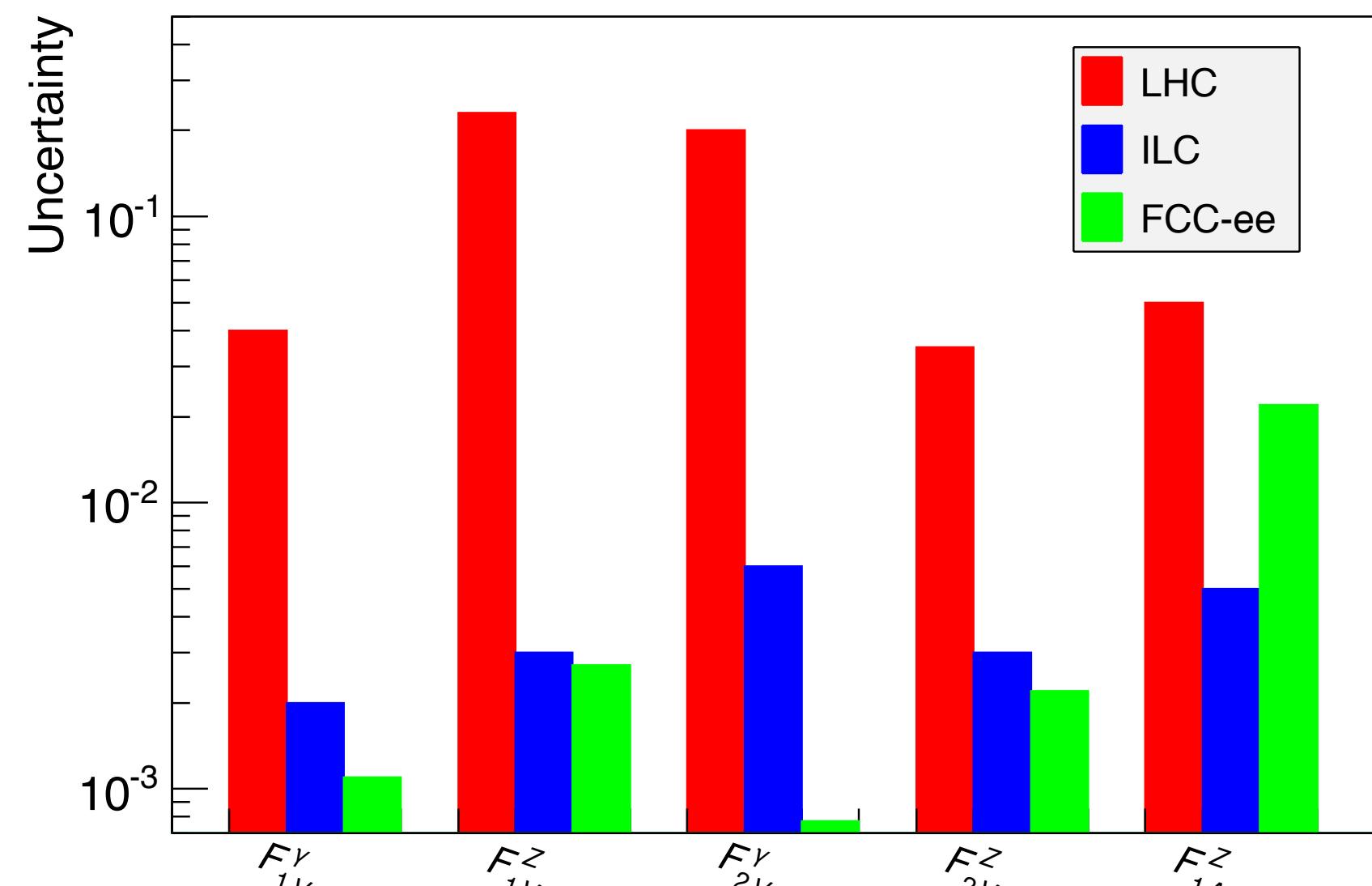
Partial Quark flavour universality assumed (1st 2 families with equal couplings)

# Electroweak precision observables at FCC ee

- Electroweak precision measurements at FCC-ee

## Precision measurements of Top couplings at 365 GeV

$$\Gamma_{\mu}^{ttX} = -ie \left\{ \gamma_{\mu} (F_{1V}^X + \gamma_5 F_{1A}^X) + \frac{\sigma_{\mu\nu}}{2m_t} (p_t + p_{\bar{t}})^{\nu} (iF_{2V}^X + \gamma_5 F_{2A}^X) \right\}$$



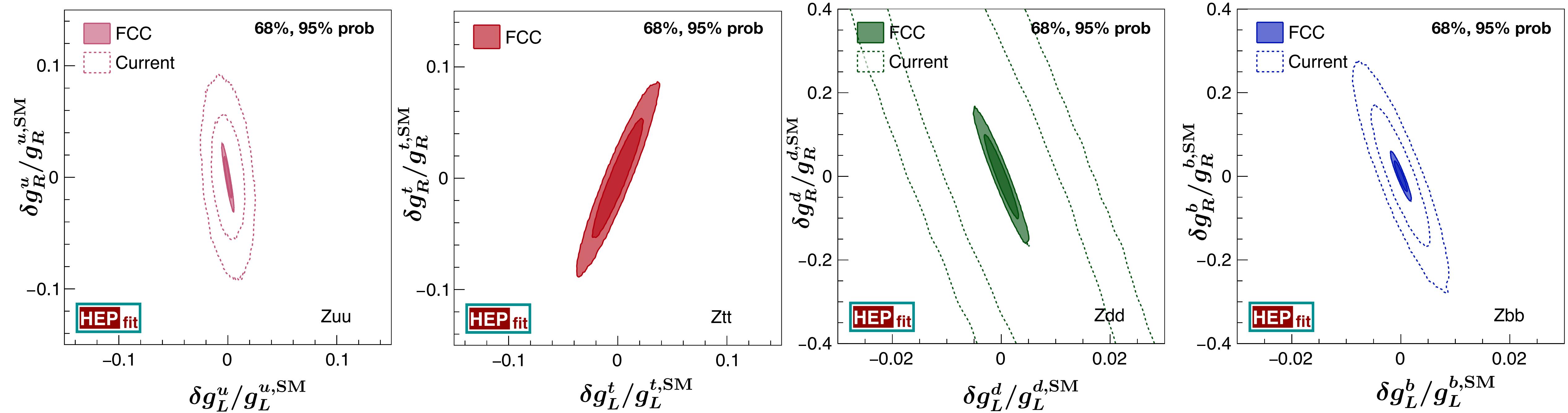
Precision on	$F_{1V}^{\gamma}$	$F_{1V}^Z$	$F_{1A}^{\gamma}$	$F_{1A}^Z$
Only three $F_{1V,A}^X$	$1.2 \cdot 10^{-3}$	$2.9 \cdot 10^{-3}$	$0.01 \cdot 10^{-2}$	$2.2 \cdot 10^{-2}$

FCC-ee 365 GeV			
Coupling	Uncertainty	Correlation Matrix	
$\delta g_V^t / g_V^t$ SM	0.009	1	0.25
$\delta g_A^t / g_A^t$ SM	0.021	1	
$\delta g_L^t / g_L^t$ SM	0.008	1	-0.96
$\delta g_R^t / g_R^t$ SM	0.016	1	

Assuming new physics only in  $Ztt$  couplings

# Electroweak precision observables at FCC ee

- Global fit to electroweak precision measurements at FCC-ee

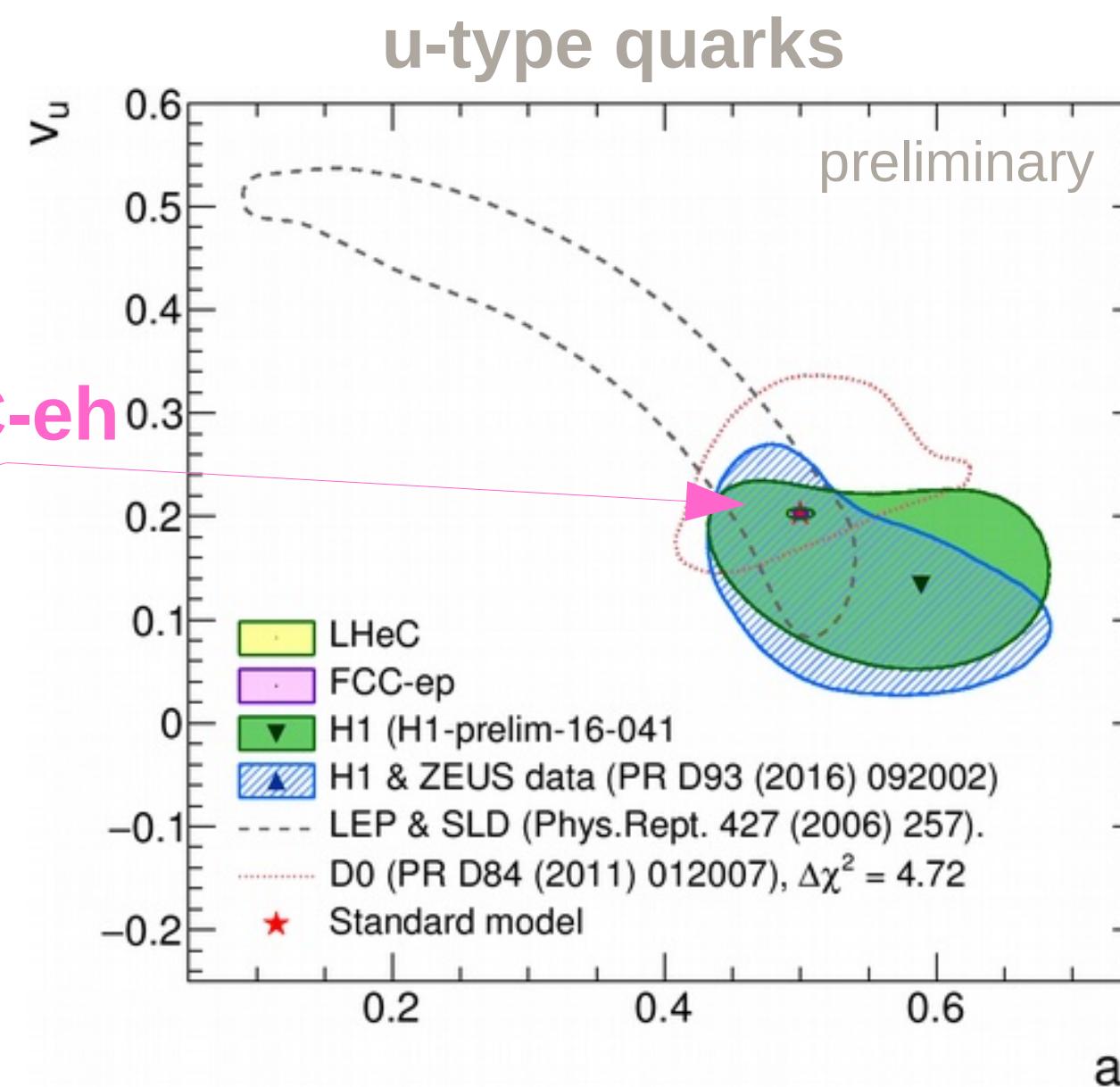
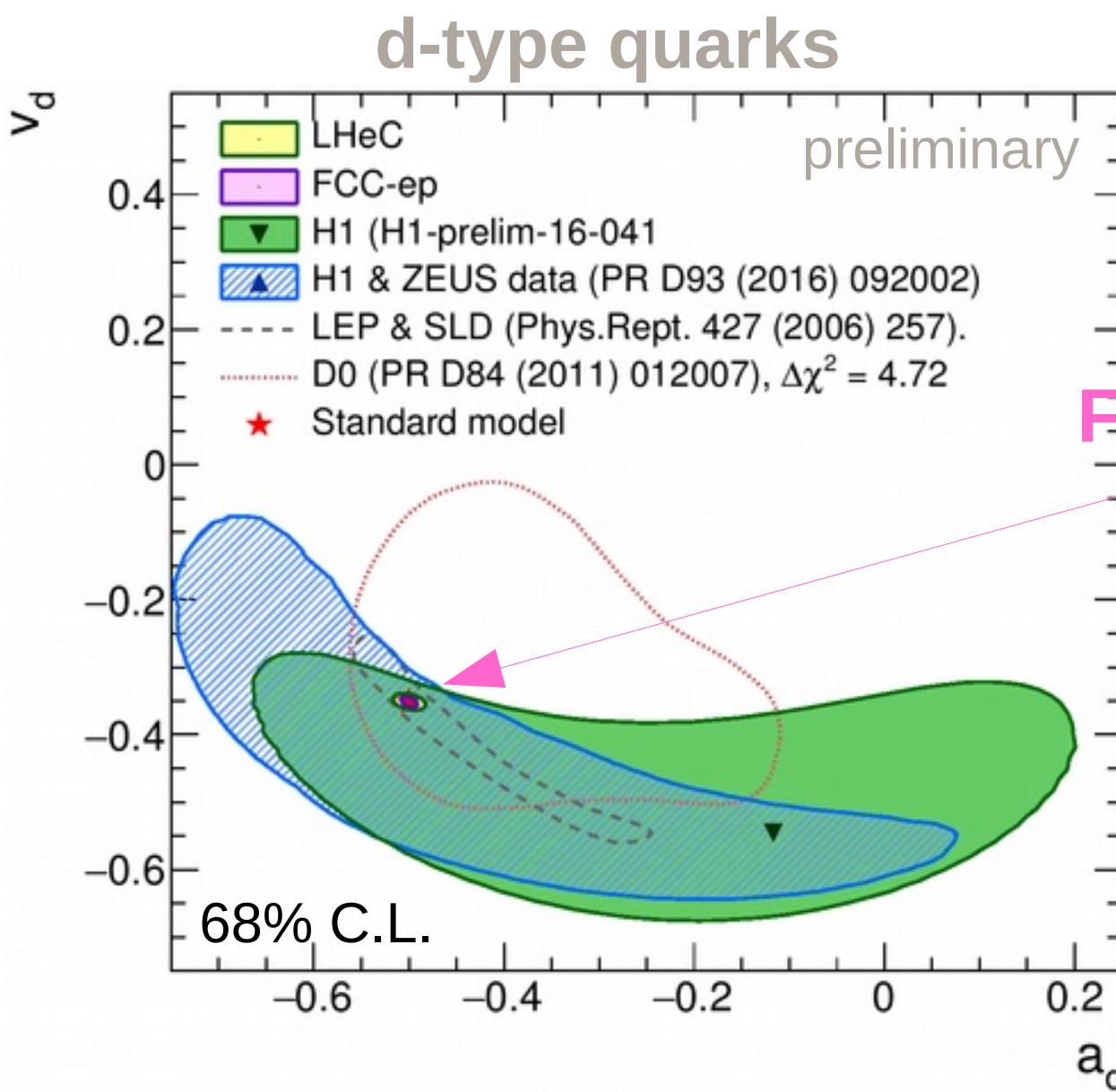


Partial Quark flavour universality assumed (1st 2 families with equal couplings)

# Electroweak precision observables at FCC eh

- Electroweak precision measurements at FCC-eh

## Precision measurements of couplings to light quark families



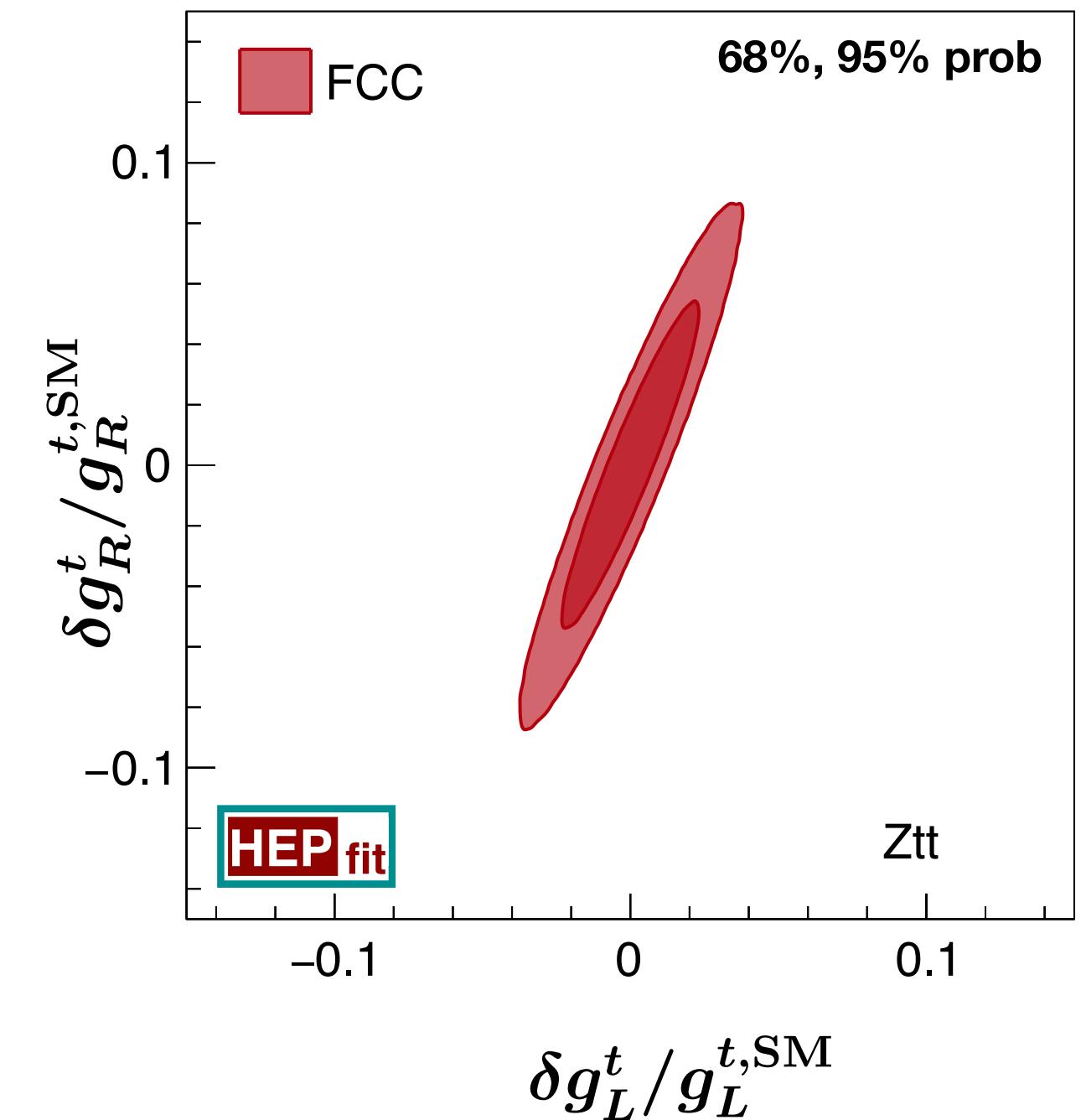
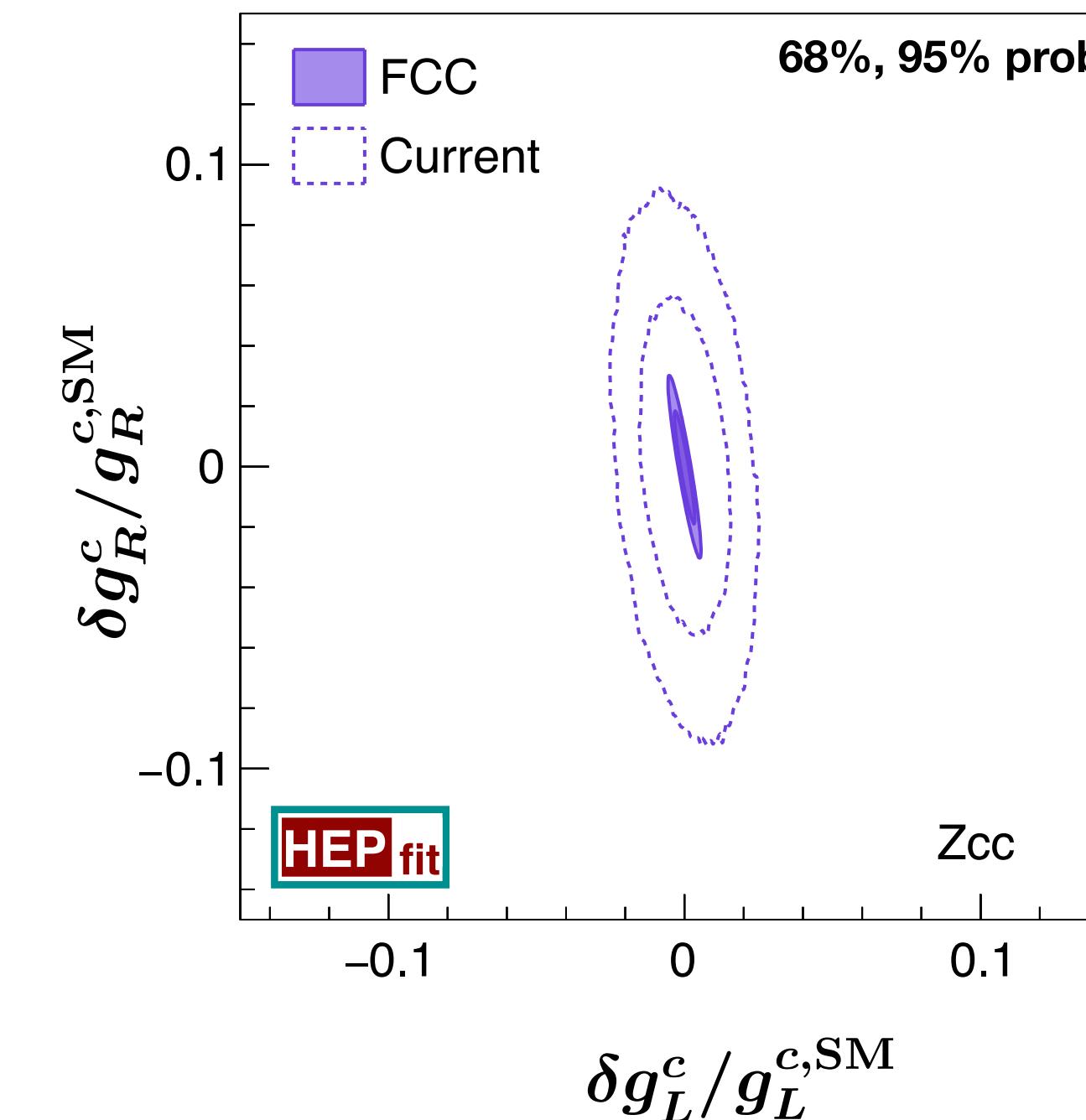
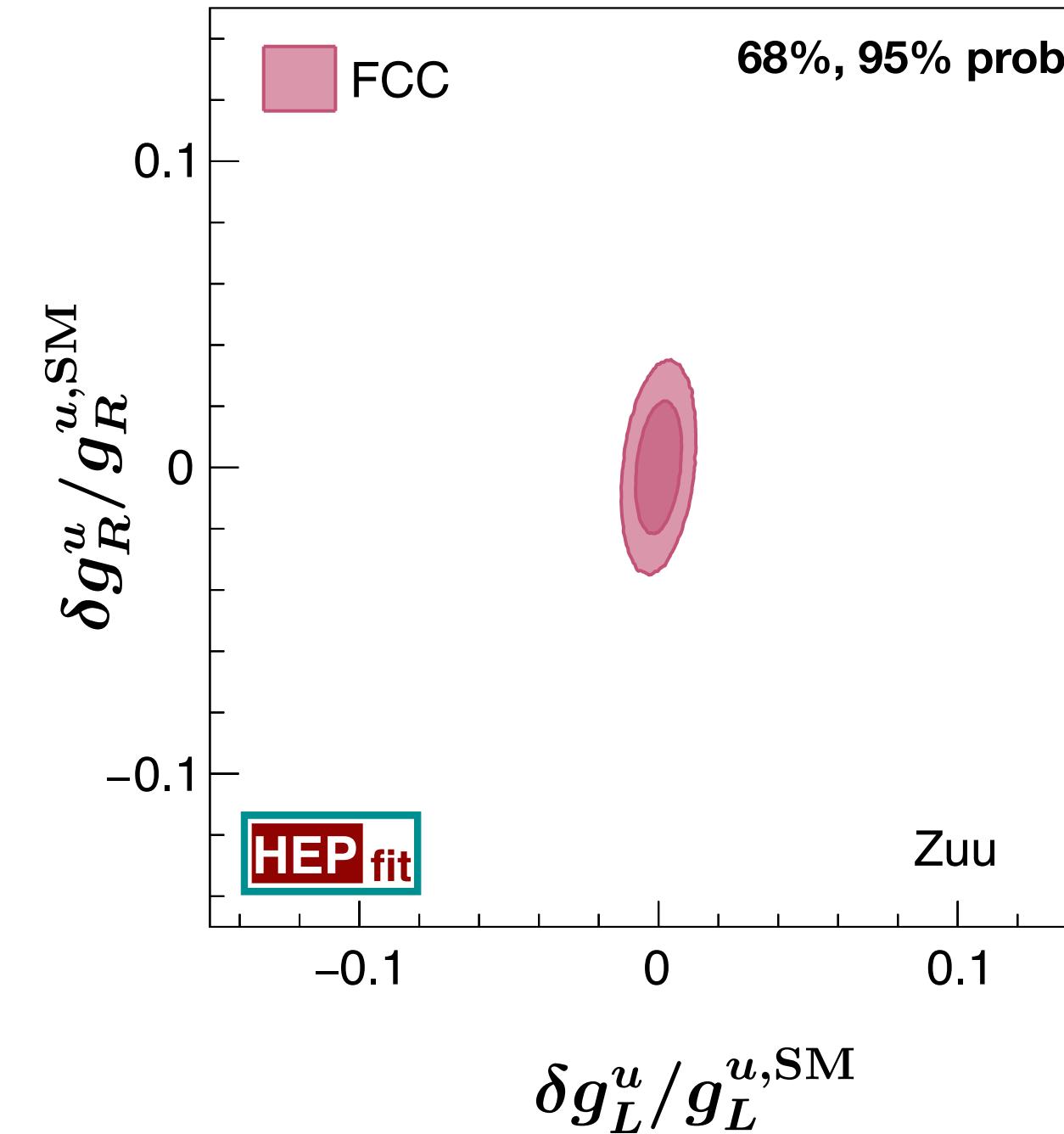
Observable	Uncertainty	(Relative uncertainty)
$g_V^u$	0.0022	(1.1%)
$g_A^u$	0.0031	(0.6%)
$g_V^d$	0.0049	(1.4%)
$g_A^d$	0.0049	(0.97%)

Assuming new physics only in  $Zqq$  couplings

Thanks to Daniel Britzger for info about the FCC-eh EW measurements

# Electroweak precision observables at FCC ee/eh

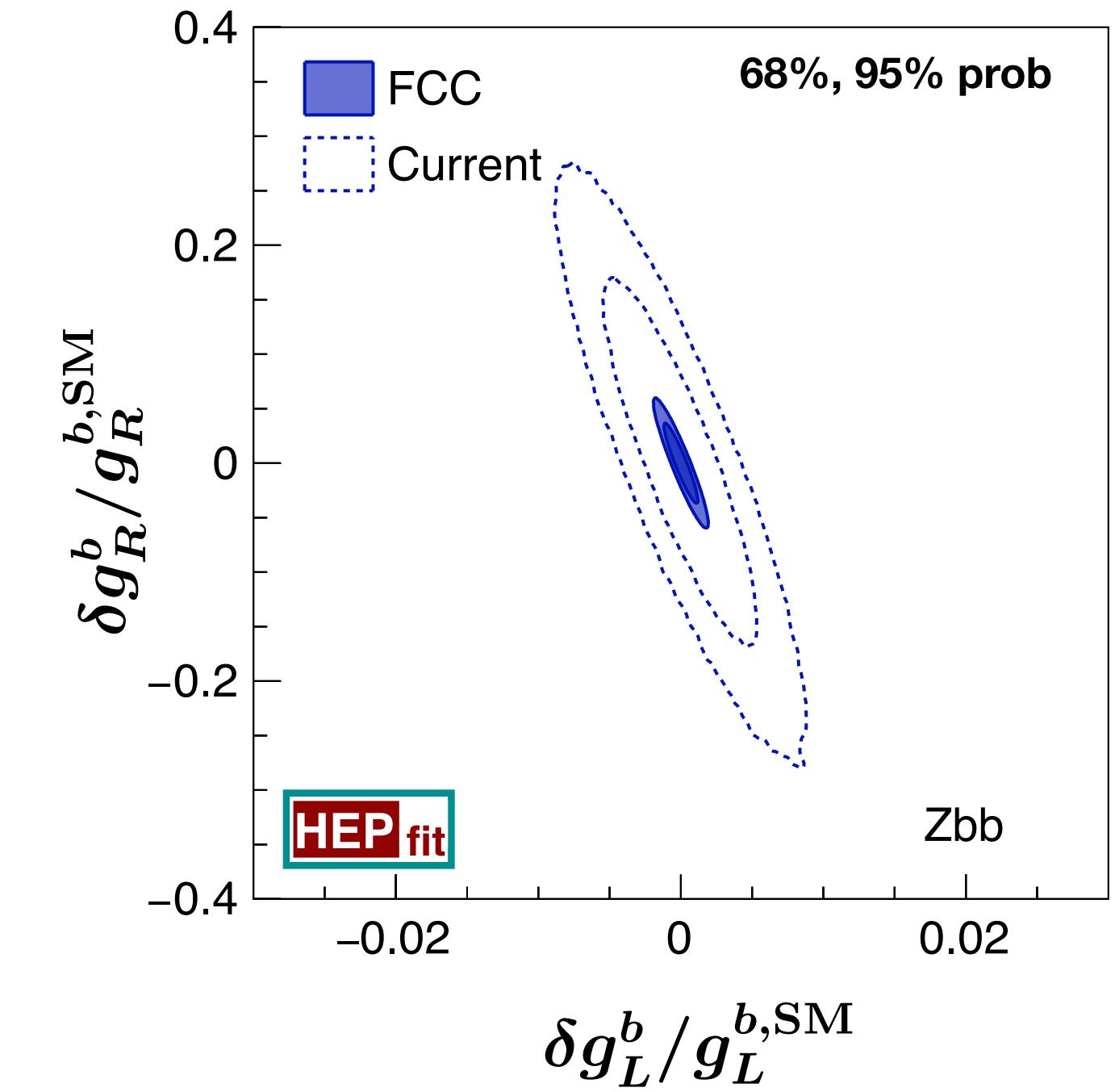
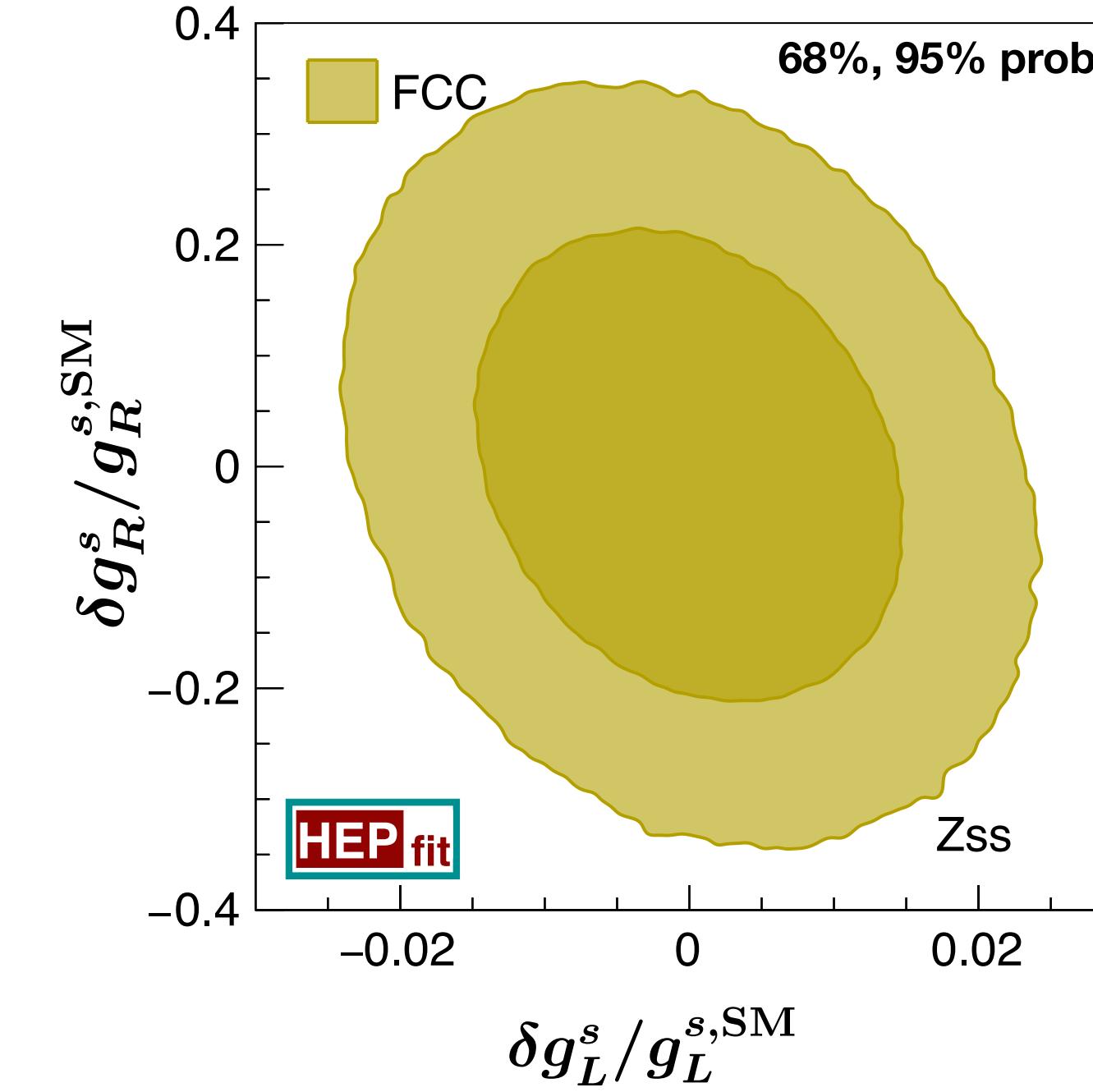
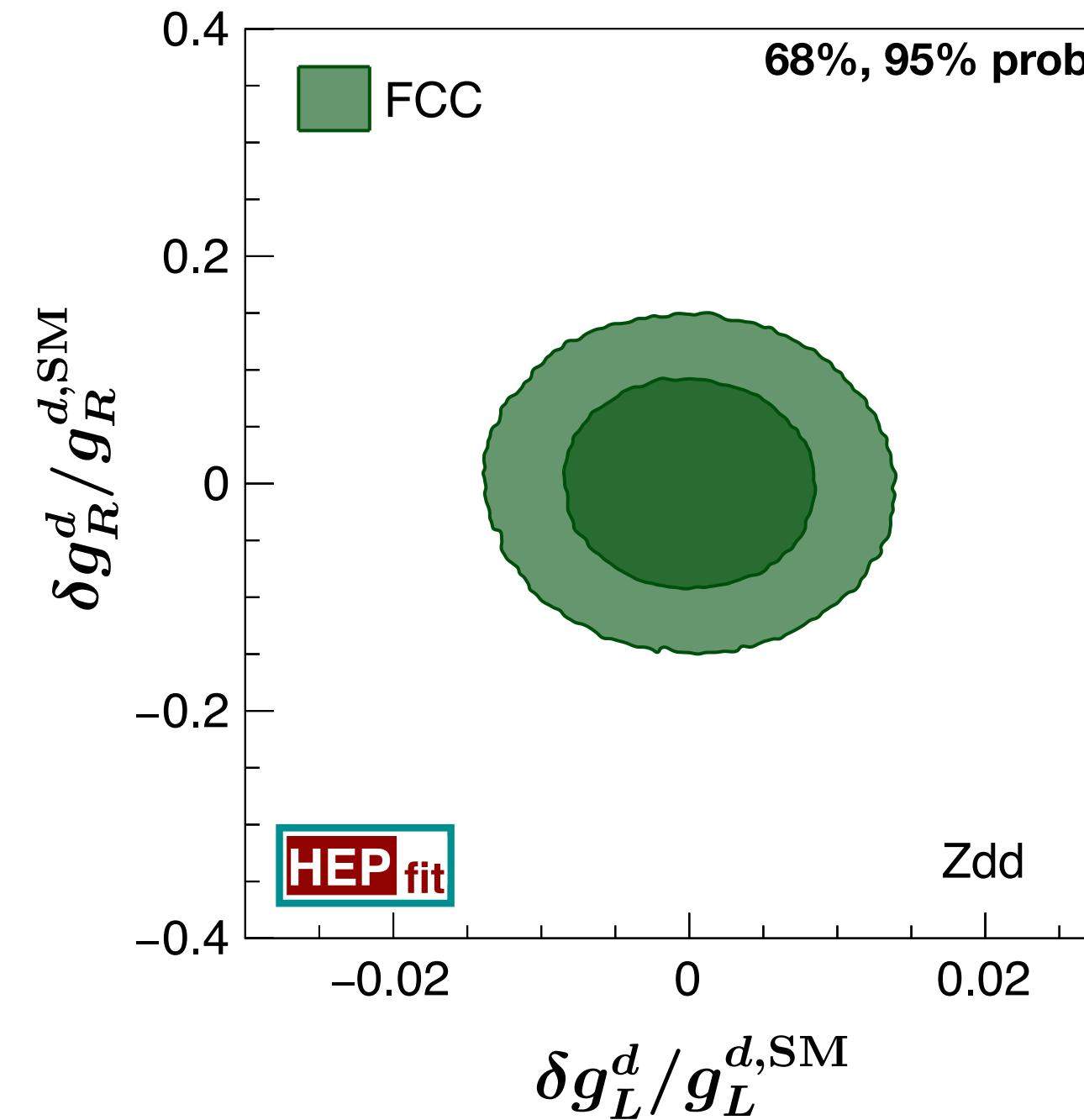
- Global fit to electroweak precision measurements at FCC-ee + FCC-eh



No Fermion flavour universality assumed

# Electroweak precision observables at FCC ee/eh

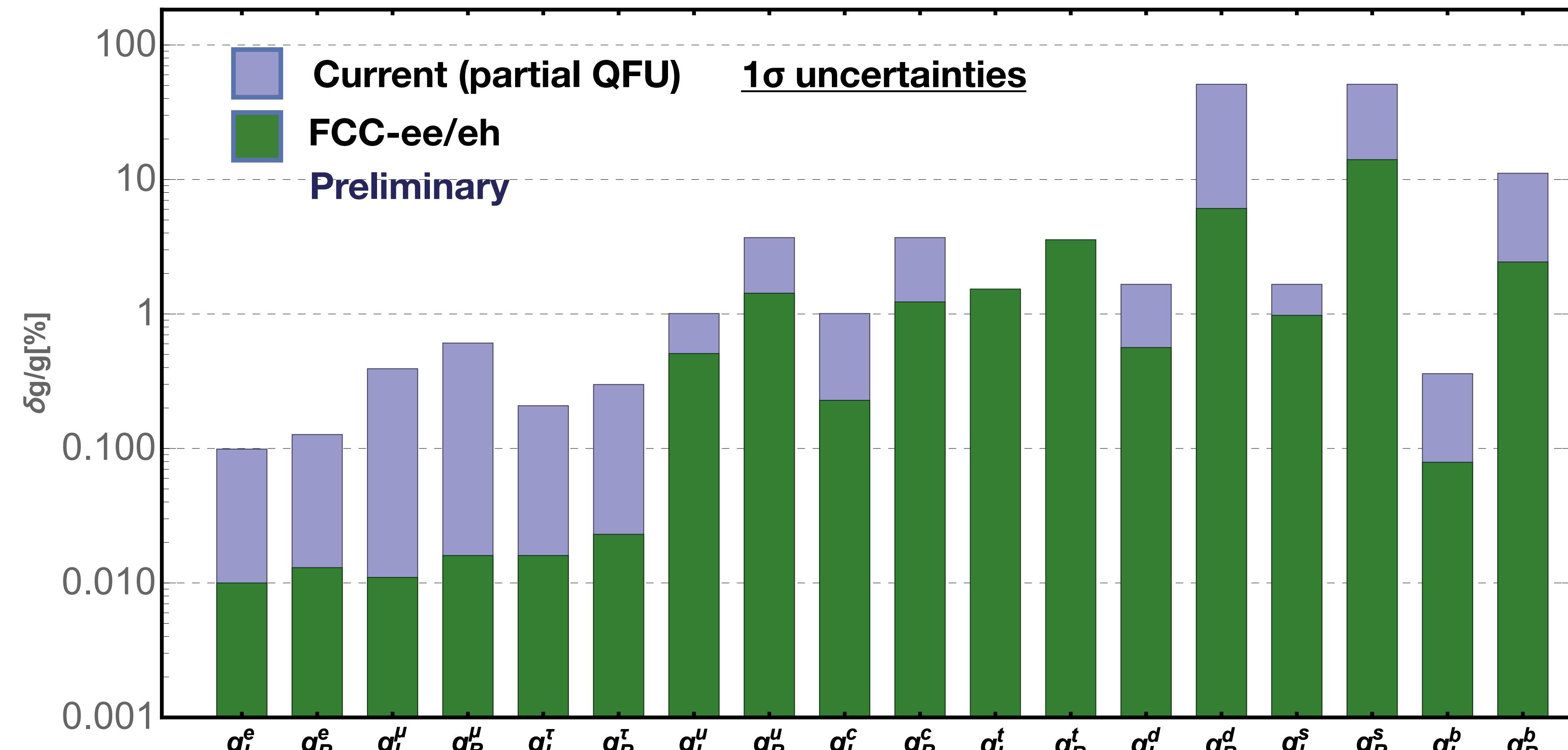
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**No Fermion flavour universality assumed**

# Electroweak precision observables at FCC ee/eh

- Global fit to electroweak precision measurements at FCC-ee + FCC-eh

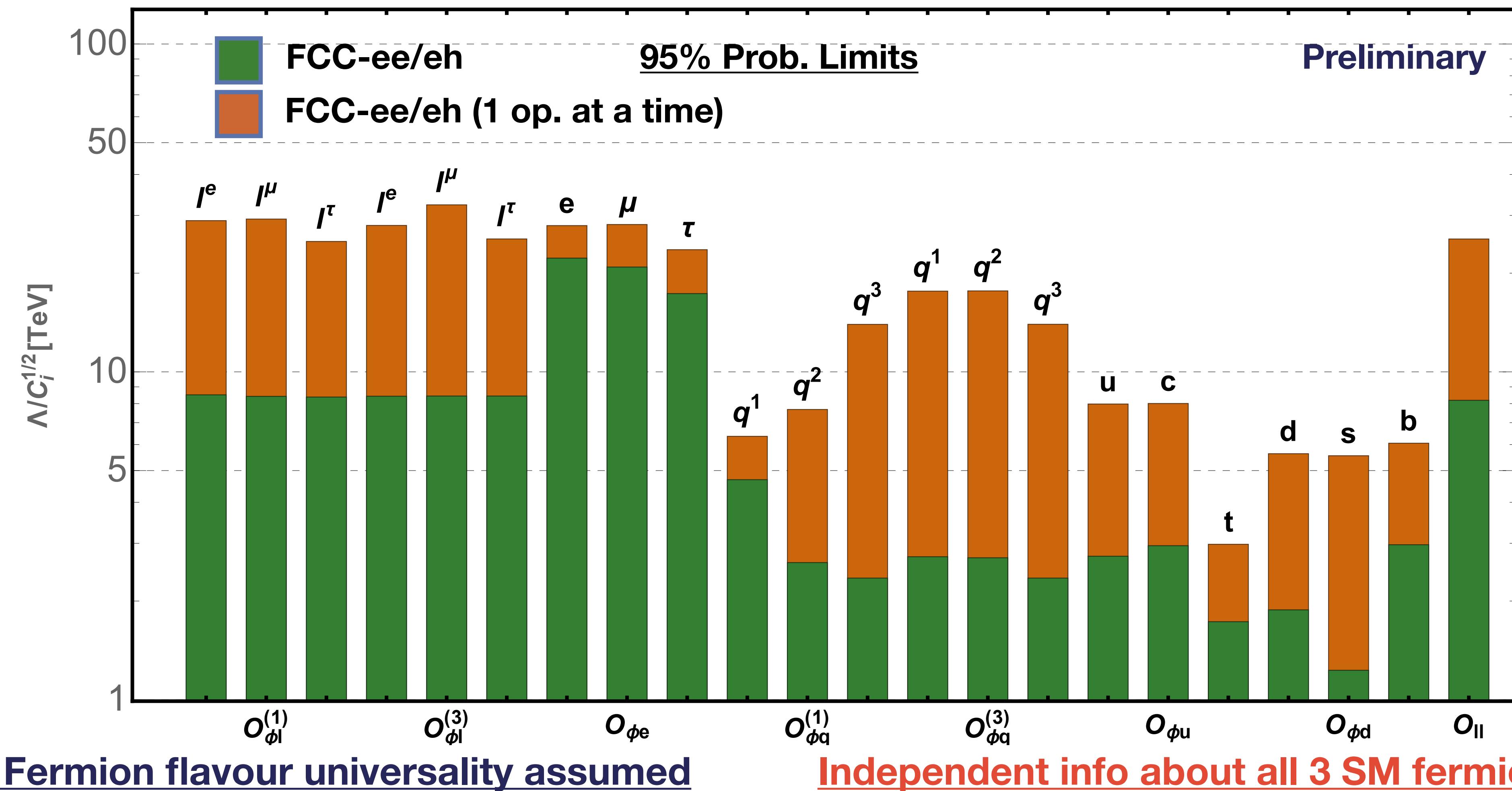


No Fermion flavour universality assumed

Independent info about all 3 SM fermion families

# Electroweak precision observables at FCC ee/eh

- Global fit to electroweak precision measurements at FCC-ee + FCC-eh



# Higgs precision measurements at the FCC

# Higgs precision observables at FCC ee

- Uncertainties rescaled to 2 IP wrt. “First Look at the Physics Case of TLEP” JHEP 1401 (2014) 164

## Associated ZH production

Observable	Expected uncertainty
$\sigma_{HZ}$	0.57%
$\sigma_{HZ} \text{ Br}(H \rightarrow b\bar{b})$	0.28%
$\sigma_{HZ} \text{ Br}(H \rightarrow c\bar{c})$	1.7%
$\sigma_{HZ} \text{ Br}(H \rightarrow gg)$	2.0%
$\sigma_{HZ} \text{ Br}(H \rightarrow W^\pm W^{\mp*})$	1.3%
$\sigma_{HZ} \text{ Br}(H \rightarrow \tau^+ \tau^-)$	1.0%
$\sigma_{HZ} \text{ Br}(H \rightarrow ZZ^*)$	4.4%
$\sigma_{HZ} \text{ Br}(H \rightarrow \gamma\gamma)$	4.2%
$\sigma_{HZ} \text{ Br}(H \rightarrow \mu^+ \mu^-)$	18.4%

## W boson fusion (WBF) production

Observable	Expected uncertainty
$\sigma_{WBF}^{(240\text{GeV})} \text{ Br}(H \rightarrow b\bar{b})$	3.1%
$\sigma_{WBF}^{(350\text{GeV})} \text{ Br}(H \rightarrow b\bar{b})$	0.79%

See M. Klute's talk for details

# Higgs precision observables at FCC eh

- FCC-eh (60 GeV e - 50 TeV p): Precisions for 2 ab<sup>-1</sup> of data

## CC DIS: $W$ boson fusion (WBF)

Observable	Expected uncertainty
$\sigma_{WBF} \text{ Br}(H \rightarrow b\bar{b})$	0.27%
$\sigma_{WBF} \text{ Br}(H \rightarrow c\bar{c})$	2.36%
$\sigma_{WBF} \text{ Br}(H \rightarrow gg)$	1.78%
$\sigma_{WBF} \text{ Br}(H \rightarrow W^\pm W^{\mp*})$	2.45%
$\sigma_{WBF} \text{ Br}(H \rightarrow \tau^+ \tau^-)$	1.65%
$\sigma_{WBF} \text{ Br}(H \rightarrow ZZ^*)$	3.94%
$\sigma_{WBF} \text{ Br}(H \rightarrow \gamma\gamma)$	4.7%

## NC DIS: $Z$ boson fusion (ZBF)

Observable	Expected uncertainty
$\sigma_{ZBF} \text{ Br}(H \rightarrow b\bar{b})$	0.83%
$\sigma_{ZBF} \text{ Br}(H \rightarrow c\bar{c})$	7.08%
$\sigma_{ZBF} \text{ Br}(H \rightarrow gg)$	5.62%
$\sigma_{ZBF} \text{ Br}(H \rightarrow W^\pm W^{\mp*})$	4.29%
$\sigma_{ZBF} \text{ Br}(H \rightarrow \tau^+ \tau^-)$	5.25%
$\sigma_{ZBF} \text{ Br}(H \rightarrow ZZ^*)$	11.8%
$\sigma_{ZBF} \text{ Br}(H \rightarrow \gamma\gamma)$	14.1%

Thanks to Max and Uta Klein for info about the FCC-eh Higgs measurements

# Higgs precision observables at FCC ee and eh

- Fit to modified Higgs couplings (assuming no extra invisible decays)

FCC-ee	
Coupling	Relative precision
$\kappa_b$	0.58%
$\kappa_t$	—
$\kappa_\tau$	0.78%
$\kappa_c$	1.05%
$\kappa_\mu$	9.6%
$\kappa_Z$	0.16%
$\kappa_W$	0.41%
$\kappa_g$	1.23%
$\kappa_\gamma$	2.18%
$\kappa_{Z\gamma}$	—

FCC-eh	
Coupling	Relative precision
$\kappa_b$	0.74%
$\kappa_t$	—
$\kappa_\tau$	1.10%
$\kappa_c$	1.35%
$\kappa_\mu$	—
$\kappa_Z$	0.43%
$\kappa_W$	0.26%
$\kappa_g$	1.17%
$\kappa_\gamma$	2.35%
$\kappa_{Z\gamma}$	—

See M. Klute's talk for MI fit

$$\kappa_i \equiv g_{hi}/g_{hi}^{SM}$$

# Higgs precision observables at FCC ee and eh

- Fit to modified Higgs couplings (assuming no extra invisible decays)

FCC-ee	
Coupling	Relative precision
$\kappa_b$	0.58%
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$\kappa_Z$	0.16%
$\kappa_W$	0.41%
$\kappa_g$	1.23%
$\kappa_\gamma$	2.18%
$\kappa_{Z\gamma}$	—

FCC-eh	
Coupling	Relative precision
$\kappa_b$	0.74%
$\kappa_t$	—
$\kappa_\tau$	1.10%
$\kappa_c$	1.35%
$\kappa_\mu$	—
$\kappa_Z$	0.43%
$\kappa_W$	0.26%
$\kappa_g$	1.17%
$\kappa_\gamma$	2.35%
$\kappa_{Z\gamma}$	—

$$\kappa_i \equiv g_{hi}/g_{hi}^{SM}$$

# Higgs precision observables at FCC ee and eh

- Fit to modified Higgs couplings (assuming no extra invisible decays)

FCC-ee	
Coupling	Relative precision
$\kappa_b$	0.58%
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$\kappa_c$	1.05%
$\kappa_\mu$	9.6%
$\kappa_Z$	0.16%
$\kappa_W$	0.41%
$\kappa_g$	1.23%
$\kappa_\gamma$	2.18%
$\kappa_{Z\gamma}$	—

FCC-eh	
Coupling	Relative precision
$\kappa_b$	0.74%
$\kappa_t$	—
$\kappa_\tau$	1.10%
$\kappa_c$	1.35%
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$\kappa_Z$	0.43%
$\kappa_W$	0.26%
$\kappa_g$	1.17%
$\kappa_\gamma$	2.35%
$\kappa_{Z\gamma}$	—

$$\kappa_i \equiv g_{hi}/g_{hi}^{SM}$$

See U. Klein's talk for tH at FCC-eh

# Higgs precision observables at FCC ee and eh

- Fit to modified Higgs couplings (assuming no extra invisible decays)

Using HLLHC Results (Only 1 Experiment) to fill the gaps

HLLHC (ATLAS)	
Coupling	Relative precision
$\kappa_b$	10.4%
$\kappa_t$	7.6%
$\kappa_\tau$	9.43%
$\kappa_c$	—
$\kappa_\mu$	7.4%
$\kappa_Z$	3.7%
$\kappa_W$	4.2%
$\kappa_g$	5.2%
$\kappa_\gamma$	4.3%
$\kappa_{Z\gamma}$	15%

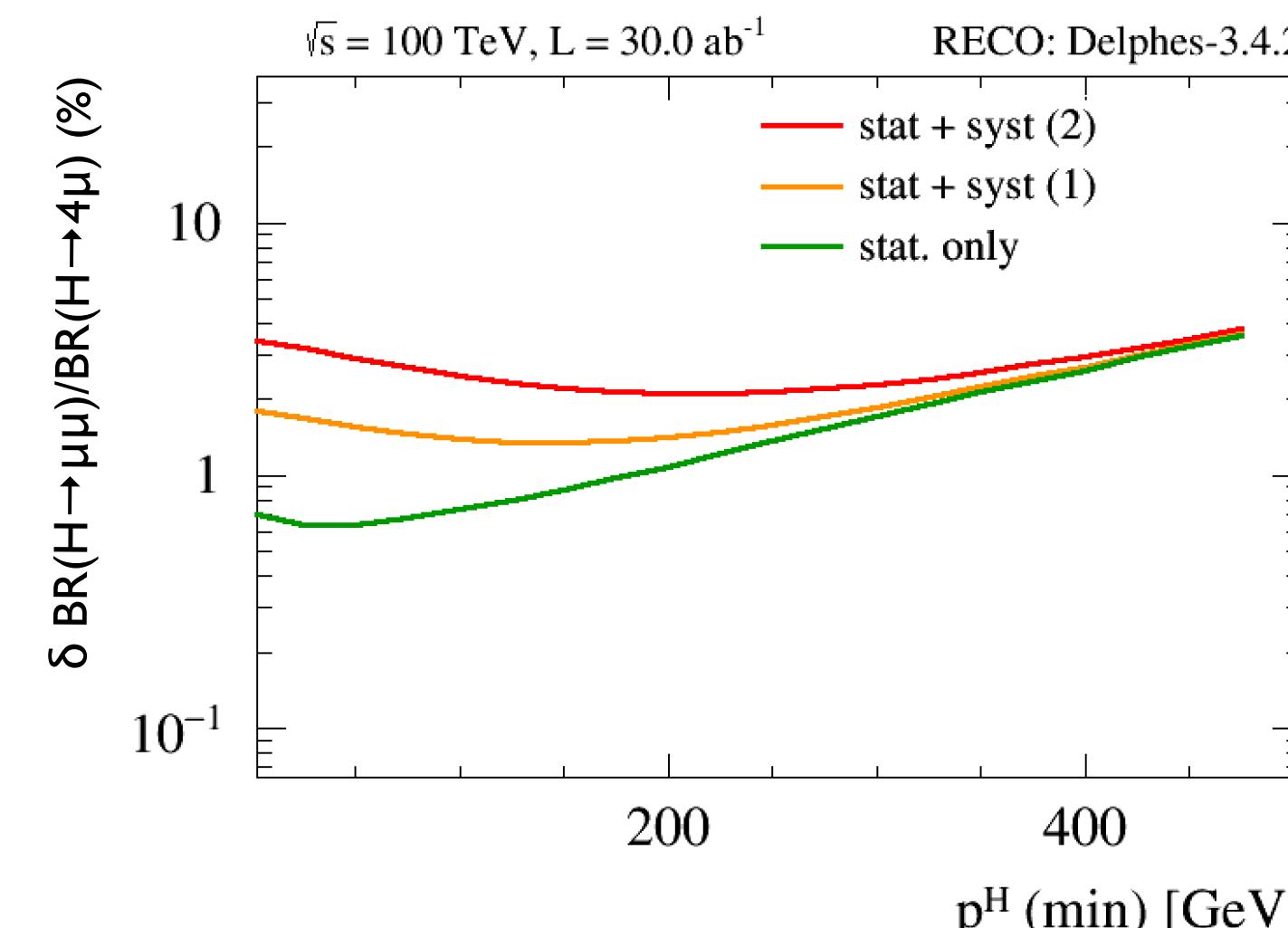
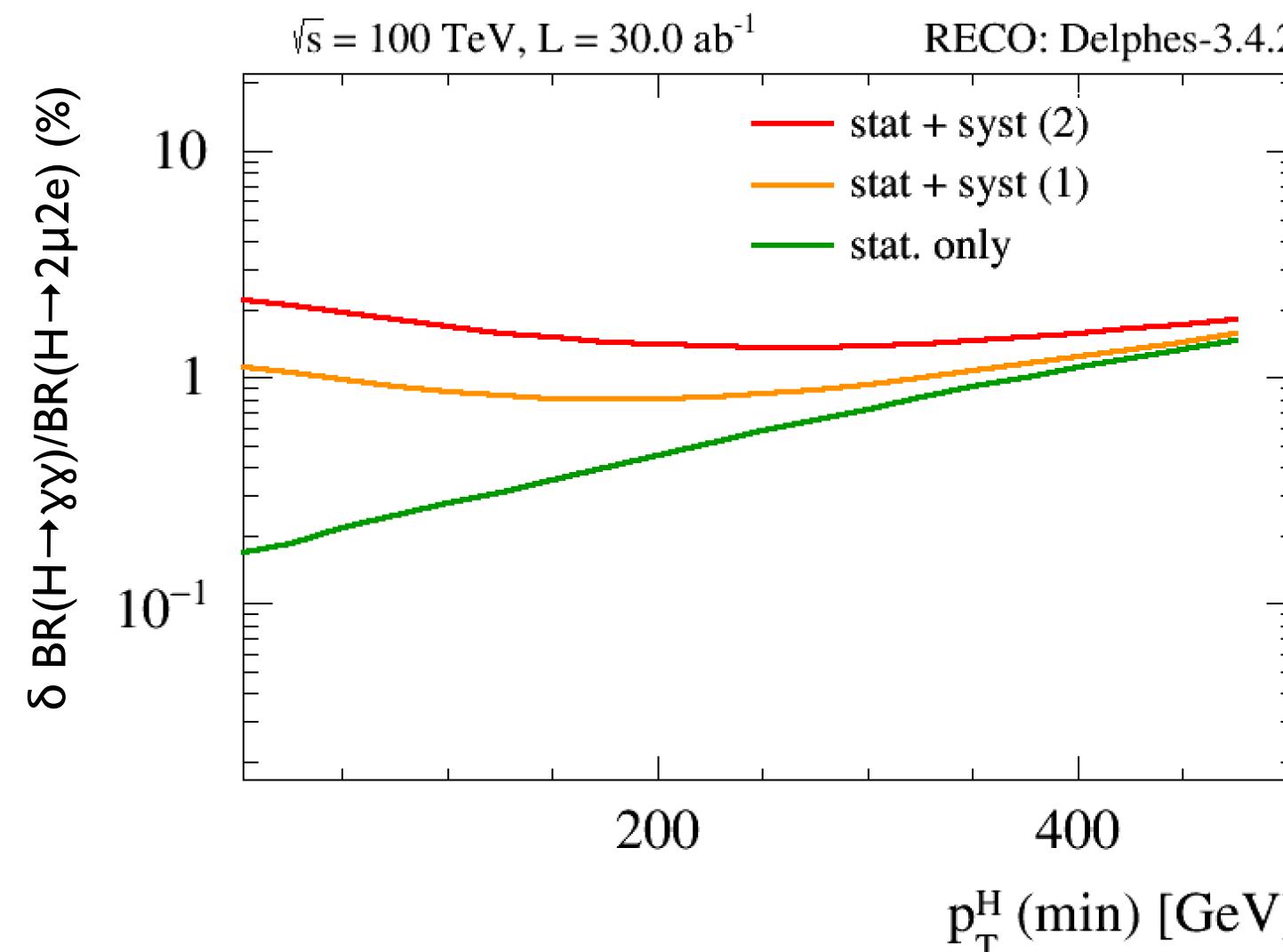
HLLHC + FCC-ee	
Coupling	Relative precision
$\kappa_b$	0.55%
$\kappa_t$	5.19%
$\kappa_\tau$	0.76%
$\kappa_c$	1.03%
$\kappa_\mu$	5.20%
$\kappa_Z$	0.16%
$\kappa_W$	0.40%
$\kappa_g$	1.03%
$\kappa_\gamma$	1.41%
$\kappa_{Z\gamma}$	14.2%

HLLHC + FCC-eh	
Coupling	Relative precision
$\kappa_b$	0.69%
$\kappa_t$	5.27%
$\kappa_\tau$	1.06%
$\kappa_c$	1.33%
$\kappa_\mu$	6.25%
$\kappa_Z$	0.41%
$\kappa_W$	0.25%
$\kappa_g$	1.02%
$\kappa_\gamma$	1.44%
$\kappa_{Z\gamma}$	14.26%

$$\kappa_i \equiv g_{hi}/g_{hi}^{SM}$$

# Higgs complementarities

- Rare Higgs decays statistically limited at FCC-ee/eh
  - Can be measured at FCC-hh with 1% stat. precision (in  $\delta\mu/\mu$ )
  - Systematics can be further cancelled by measuring ratios of BR ( $\gamma\gamma/4l$ ,  $\mu\mu/4l$ ,  $Z\gamma/4l$ ,  $\gamma\gamma/\mu\mu$ )



M. Selvaggi, Talk at 2nd FCC Physics Workshop

1% accuracy (stat + sys)  
within reach

Provided  $\text{BR}(\text{H} \rightarrow 4l)$  know to <<1%  
( $pp \rightarrow H \rightarrow 4l$  measurable at FCC-hh to ~1%)

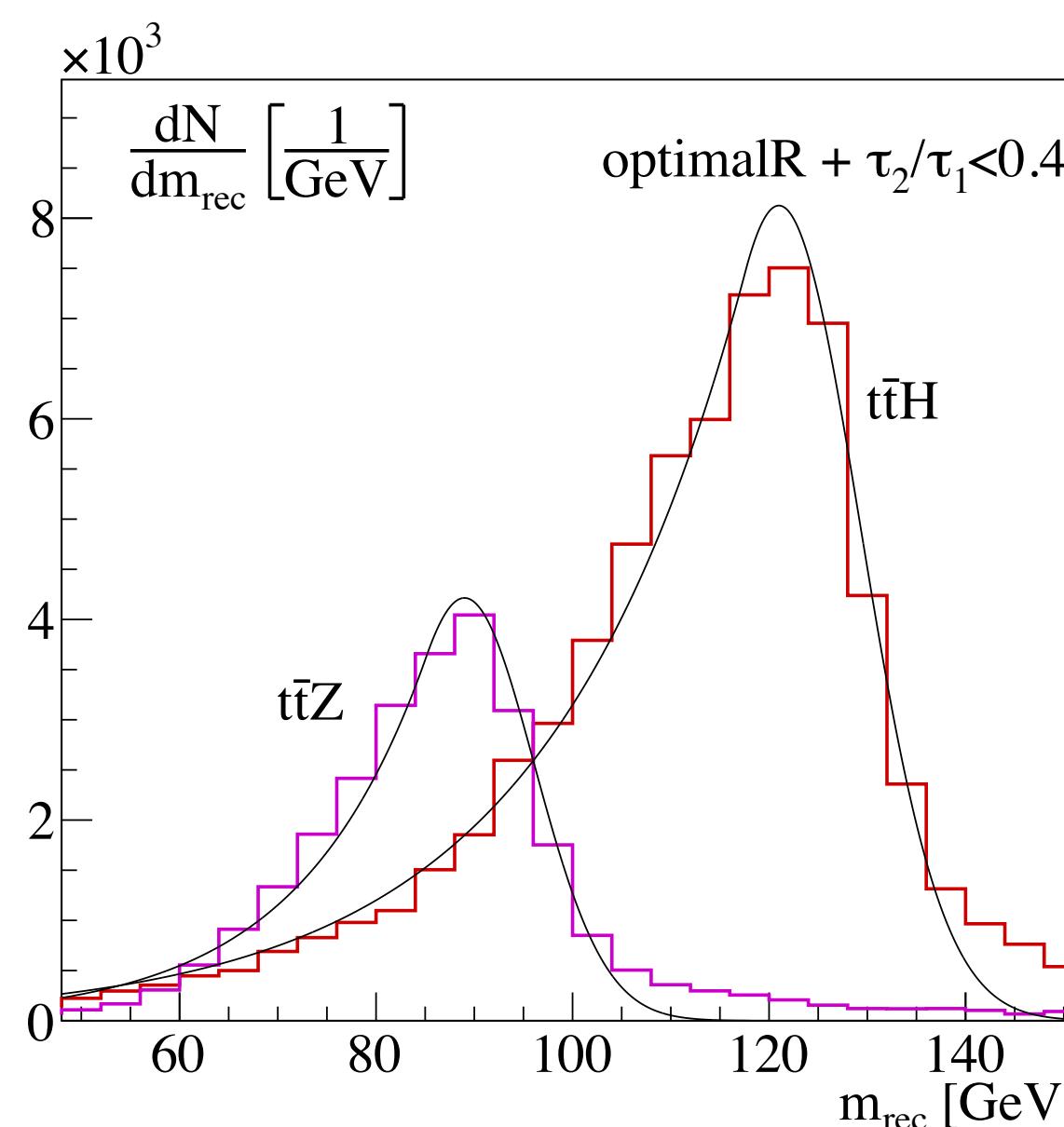
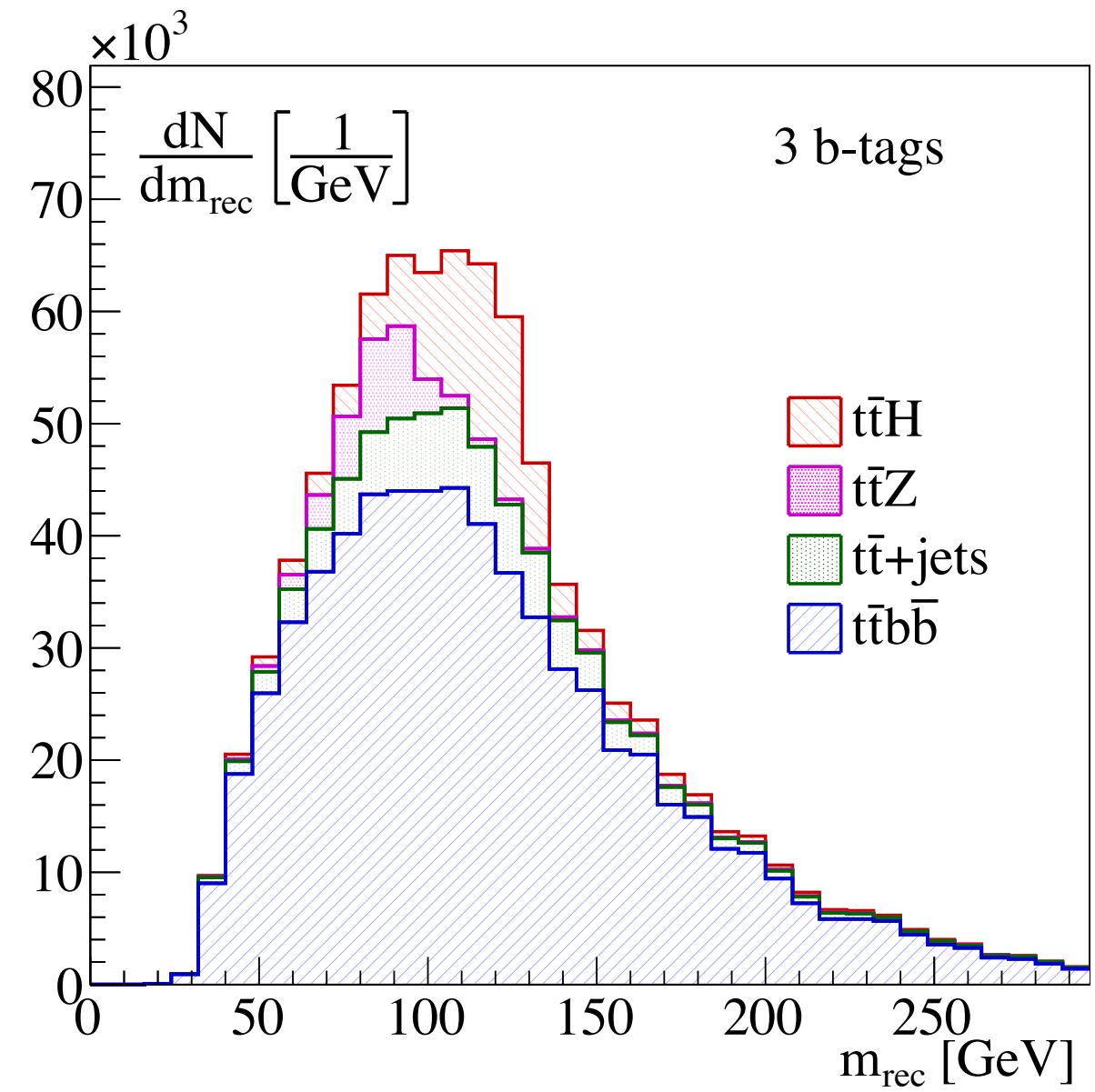
Measurable at FCC-ee/eh with  
required precision

- Robust determination by this method requires both FCC-hh and FCC-ee/eh

Thanks to Michele Selvaggi for info about the FCC-hh Higgs measurements

# Higgs complementarities

- Top Yukawa coupling not directly accessible at FCC-ee. Could be measured in single tH at FCC-eh
  - Can be measured at FCC-hh from  $\sigma(t\bar{t}H)/\sigma(t\bar{t}Z)$  boosted



e.g. Fit and extract  $N_H/N_Z$  to 1% accuracy  
⇒  $\delta_{stat+th} y_t/y_t \sim 1\%$

More on this later

M.L. Mangano et al., arXiv: 1507.08169 [hep-ph]

# Higgs complementarities: Global fit to Higgs couplings at FCC

- All single Higgs couplings can be determined below the 1%

FCC-ee/FCC-eh

Precise determinations for the leading couplings

**HZZ** Crucial for normalization of FCC-hh results

FCC-hh

Completes the picture with precise determinations of Top and coupling associated to rare decays

NOT MODEL-INDEPENDENT:

Results assume that, if there is New physics, it can only be in the Higgs couplings

Coupling	HLLHC + FCC Relative precision
$\kappa_b$	0.38%
$\kappa_t$	0.51%
$\kappa_\tau$	0.58%
$\kappa_c$	0.79%
$\kappa_\mu$	0.42%
$\kappa_Z$	0.14%
$\kappa_W$	0.17%
$\kappa_g$	0.74%
$\kappa_\gamma$	0.40%
$\kappa_{Z\gamma}$	0.52%

$$\kappa_i \equiv g_{hi}/g_{hi}^{SM}$$

# Higgs precision observables at the FCC

- Operators considered in this analysis: Higgs couplings

## Vector couplings

$$\begin{aligned}\mathcal{L}_{hVV} = & g_{hgg} G_{\mu\nu}^A G^{A\mu\nu} h + g_{hWW}^{(1)} W^{\mu\nu} W_{\mu\nu}^\dagger h + \left( g_{hWW}^{(2)} W^{+\nu} \partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.} \right) + g_{hWW}^{(3)} W_\mu^+ W^{-\mu} h \\ & + g_{hZZ}^{(1)} Z_{\mu\nu} Z^{\mu\nu} h + g_{hZZ}^{(2)} Z_\nu \partial_\mu Z^{\mu\nu} h + g_{hZZ}^{(3)} Z_\mu Z^\mu h \\ & + g_{hZA}^{(1)} Z_{\mu\nu} F^{\mu\nu} h + g_{hZA}^{(2)} Z_\nu \partial_\mu F^{\mu\nu} h + g_{hAA} F_{\mu\nu} F^{\mu\nu} h\end{aligned}$$

## Several New Operators

$$\mathcal{O}_{\phi G} = (\phi^\dagger \phi) G_{\mu\nu}^A G^{A\mu\nu}$$

$$\mathcal{O}_{\phi B} = (\phi^\dagger \phi) B_{\mu\nu} B^{\mu\nu}$$

$$\mathcal{O}_{\phi W} = (\phi^\dagger \phi) W_{\mu\nu}^a W^{a\mu\nu}$$

$$\mathcal{O}_{\phi\square} = (\phi^\dagger \phi) \square (\phi^\dagger \phi)$$

Modifies Higgs kinetic term.  
Enters in all Higgs observables

## Already present in the EWPO analysis

$$\mathcal{O}_{\phi WB} = (\phi^\dagger \sigma_a \phi) W_{\mu\nu}^a B^{\mu\nu}$$

$$\mathcal{O}_{\phi D} = |\phi^\dagger i D_\mu \phi|^2$$

$$\mathcal{O}_{D\phi B} = i D^\mu \phi^\dagger D^\nu \phi B_{\mu\nu}$$

$$\mathcal{O}_{D\phi W} = i D^\mu \phi^\dagger \sigma_a D^\nu \phi W_{\mu\nu}^a$$

Field redefinition: trade by this 2 operators (do not enter in EWPO)

# Higgs precision observables at the FCC

- Operators considered in this analysis: Higgs couplings

## Fermionic couplings

$$\mathcal{L}_{hff} = g_{hee}^{ii} \bar{e}_L^i e_R^i h + g_{huu}^{ii} \bar{u}_L^i u_R^i h + g_{hdd}^{ii} \bar{d}_L^i d_R^i h + \text{h.c.}$$

$$g_{hff} = -\frac{m_f}{v} \left( 1 + \left[ (C_{\phi\square} - \frac{1}{4}C_{\phi D}) - \frac{v}{\sqrt{2}m_f} C_{f\phi} \right] \frac{v^2}{\Lambda^2} \right)$$

## Higgs self-coupling

$$\mathcal{L}_{h^3} = g_{hhh} h^3$$

$$g_{hhh} = -\frac{M_h^2}{2v} \left( 1 + \left[ 3(C_{\phi\square} - \frac{1}{4}C_{\phi D}) - 2\frac{v^2}{M_h^2} C_\phi \right] \frac{v^2}{\Lambda^2} \right)$$

## Operators

$$\mathcal{O}_{e\phi} = (\phi^\dagger \phi) (\bar{l}_L \phi e_R)$$

$$\mathcal{O}_{u\phi} = (\phi^\dagger \phi) (\bar{q}_L \tilde{\phi} u_R)$$

$$\mathcal{O}_{d\phi} = (\phi^\dagger \phi) (\bar{q}_L \phi d_R)$$

$$\mathcal{O}_\phi = (\phi^\dagger \phi)^3$$

Only enters in Higgs self-interactions

# Diboson observables at FCC ee

- $W$  pair production at FCC-ee

See P. Azurri's talk and J. Gu's poster for details

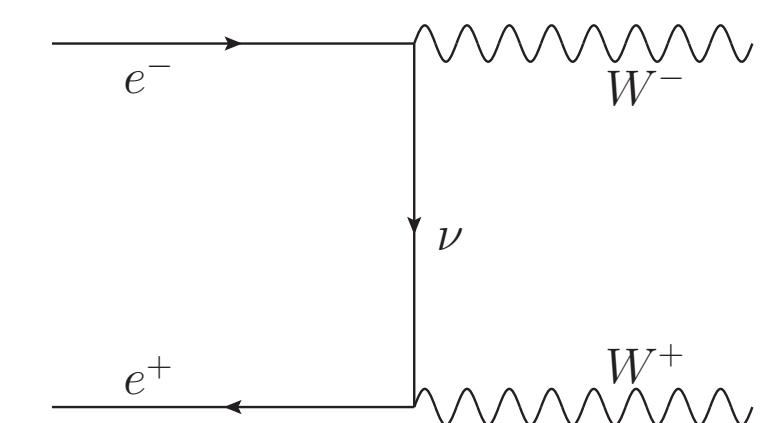
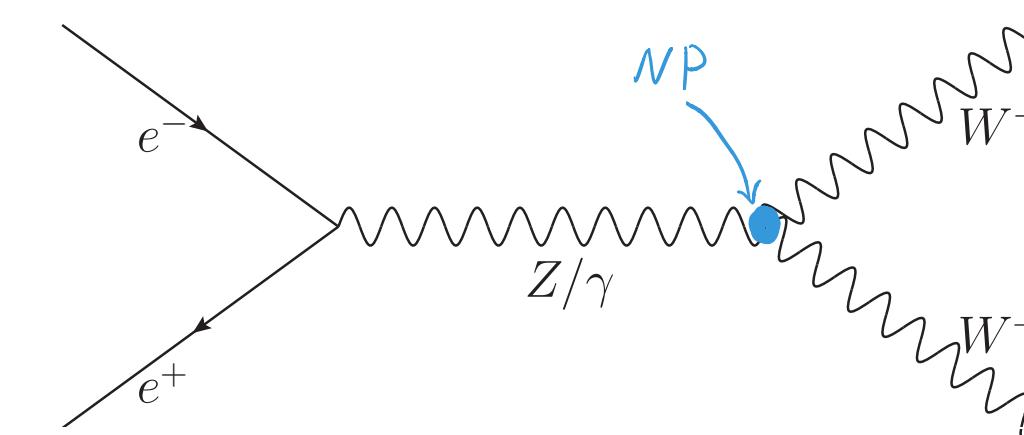
FCC-ee 161/240/350/365 GeV

aTGC	Uncertainty (stat)	Correlation Matrix
$\delta g_{1Z}$	$8.1 \times 10^{-4}$	1 <b>-0.28</b> <b>-0.87</b>
$\delta \kappa_\gamma$	$5.2 \times 10^{-4}$	1 <b>-0.12</b>
$\lambda_Z$	$7.9 \times 10^{-4}$	1

Using full angular information of the W's and decays

Statistical Analysis only

Assumes aTGC dominance



(Full EFT study in progress)

From J. Gu's Talk at FCC-ee Physics Meeting, March 19 2018

# Diboson observables at FCC ee

- Operators entering in anomalous Triple gauge couplings

$$\begin{aligned}
 \mathcal{L}_{\text{TGC}} = & ie \left[ \left( W_{\mu\nu}^+ W_\mu^- - W_{\mu\nu}^- W_\mu^+ \right) A_\nu + (1 + \delta\kappa_\gamma) A_{\mu\nu} W_\mu^+ W_\nu^- \right] \\
 & + ig \cos \theta_W \left[ (1 + \delta g_{1,Z}) \left( W_{\mu\nu}^+ W_\mu^- - W_{\mu\nu}^- W_\mu^+ \right) Z_\nu + (1 + \delta\kappa_Z) Z_{\mu\nu} W_\mu^+ W_\nu^- \right] \\
 & + ie \frac{\lambda_\gamma}{m_W^2} W_{\mu\nu}^+ W_{\nu\rho}^- A_{\rho\mu} + ig \cos \theta_W \frac{\lambda_Z}{m_W^2} W_{\mu\nu}^+ W_{\nu\rho}^- Z_{\rho\mu},
 \end{aligned}$$

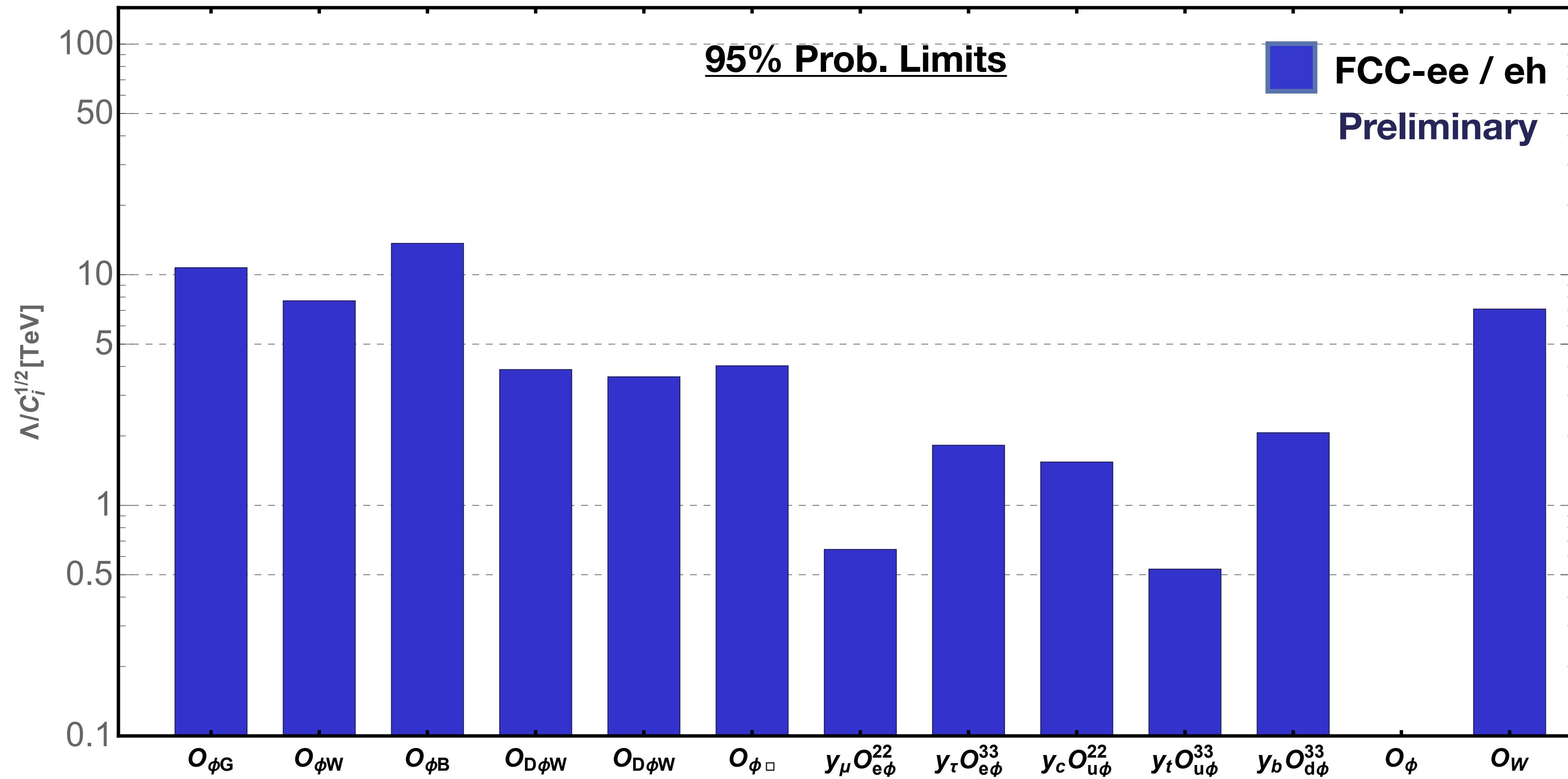
## 2 aTGC related to Higgs couplings

$$\begin{aligned}
 \delta\kappa_\gamma = & -2vc_W^2 \left( g_{hAA} - g_{hZZ}^{(1)} + \frac{1}{2s_W c_W} g_{hZA}^{(1)} (c_W^2 - s_W^2) \right) \\
 \delta g_{1,Z} = & \frac{v}{2(c_W^2 - s_W^2)} \left( c_W^2 g_{hZZ}^{(2)} + 4s_W^2 (g_{hAA} - g_{hZZ}^{(1)} + \frac{1}{4} g_{hZZ}^{(2)}) + \frac{2s_W}{c_W} g_{hZA}^{(1)} (c_W^2 - s_W^2) \right) \\
 \delta\kappa_Z = & \delta g_{1,Z} - \frac{g'^2}{g^2} \delta\kappa_\gamma \\
 \lambda_\gamma = \lambda_Z = & -\frac{3}{2} \frac{e}{s_W} C_{3W} \frac{v^2}{\Lambda^2} \quad \xrightarrow{\hspace{10em}} \quad \mathcal{O}_W = i\epsilon_{abc} W_\mu^{a\nu} W_\nu^{b\rho} W_\rho^{c\mu}
 \end{aligned}$$

Only one more operator  
(enters only in aNGC)

# Higgs precision observables at the FCC

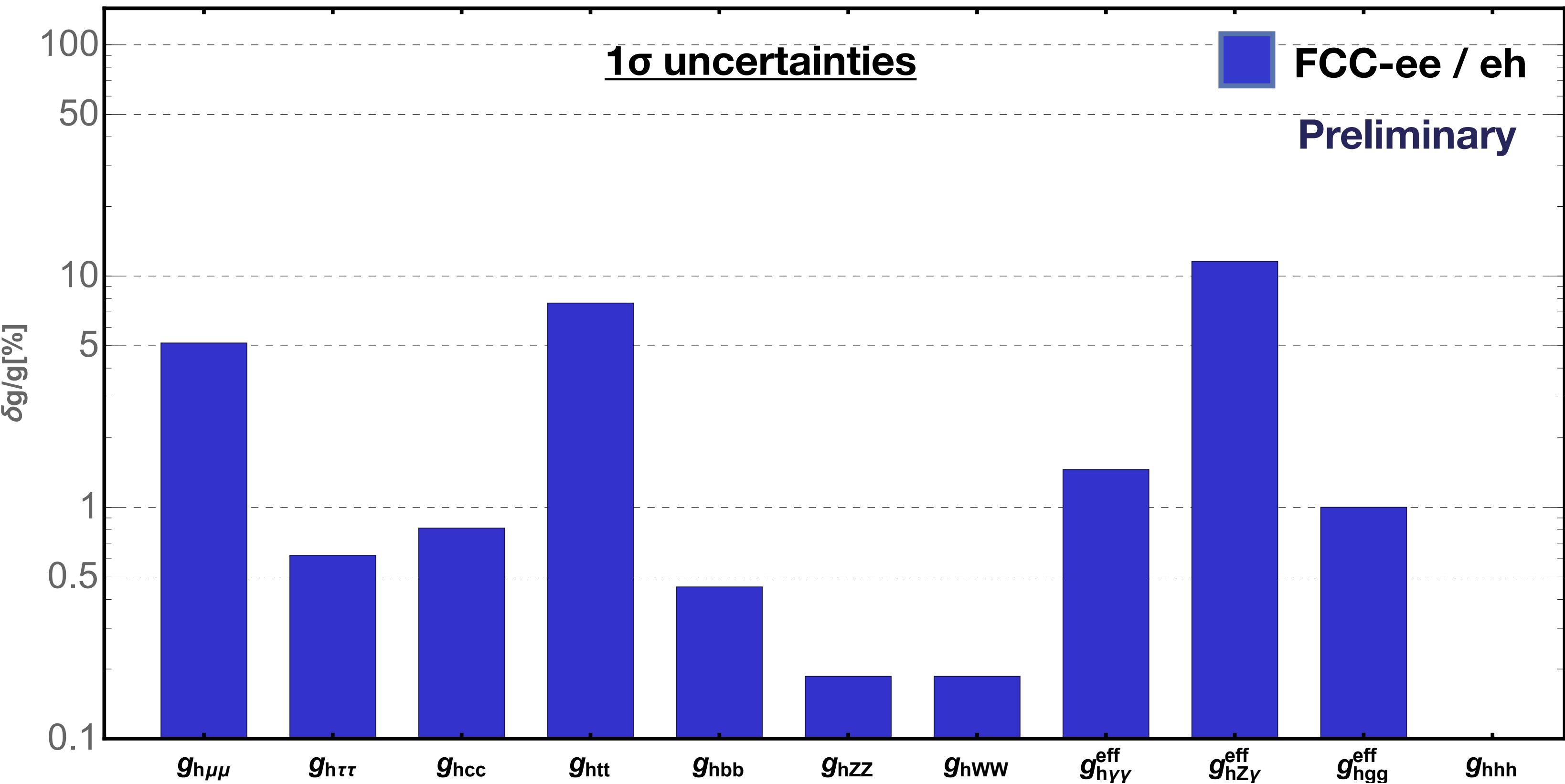
- EFT fit to Higgs precision observables at the FCC-ee and FCC-eh:



**Results show only for operators non-entering in EWPO (stay unchanged)**

# Higgs precision observables at the FCC

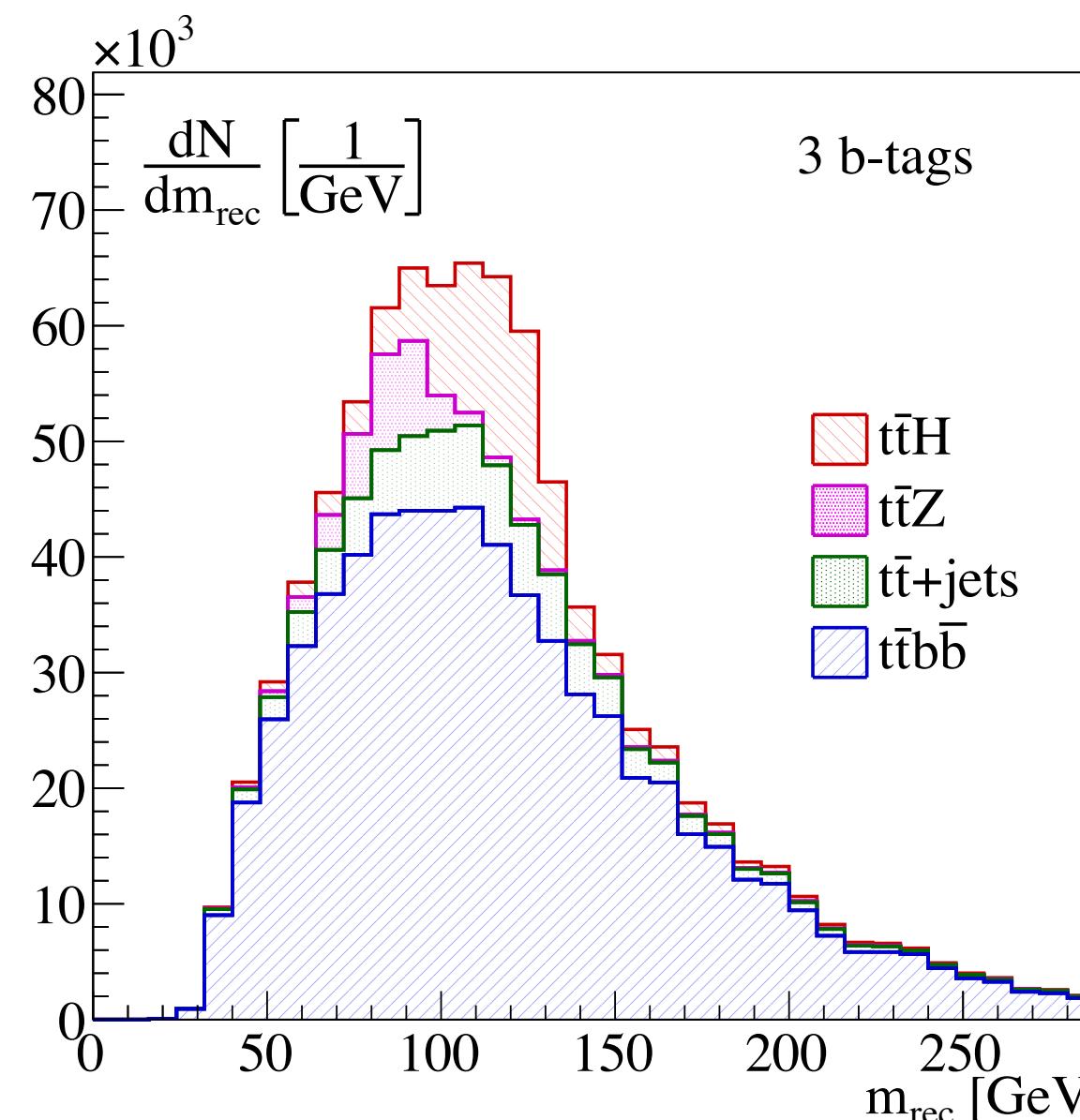
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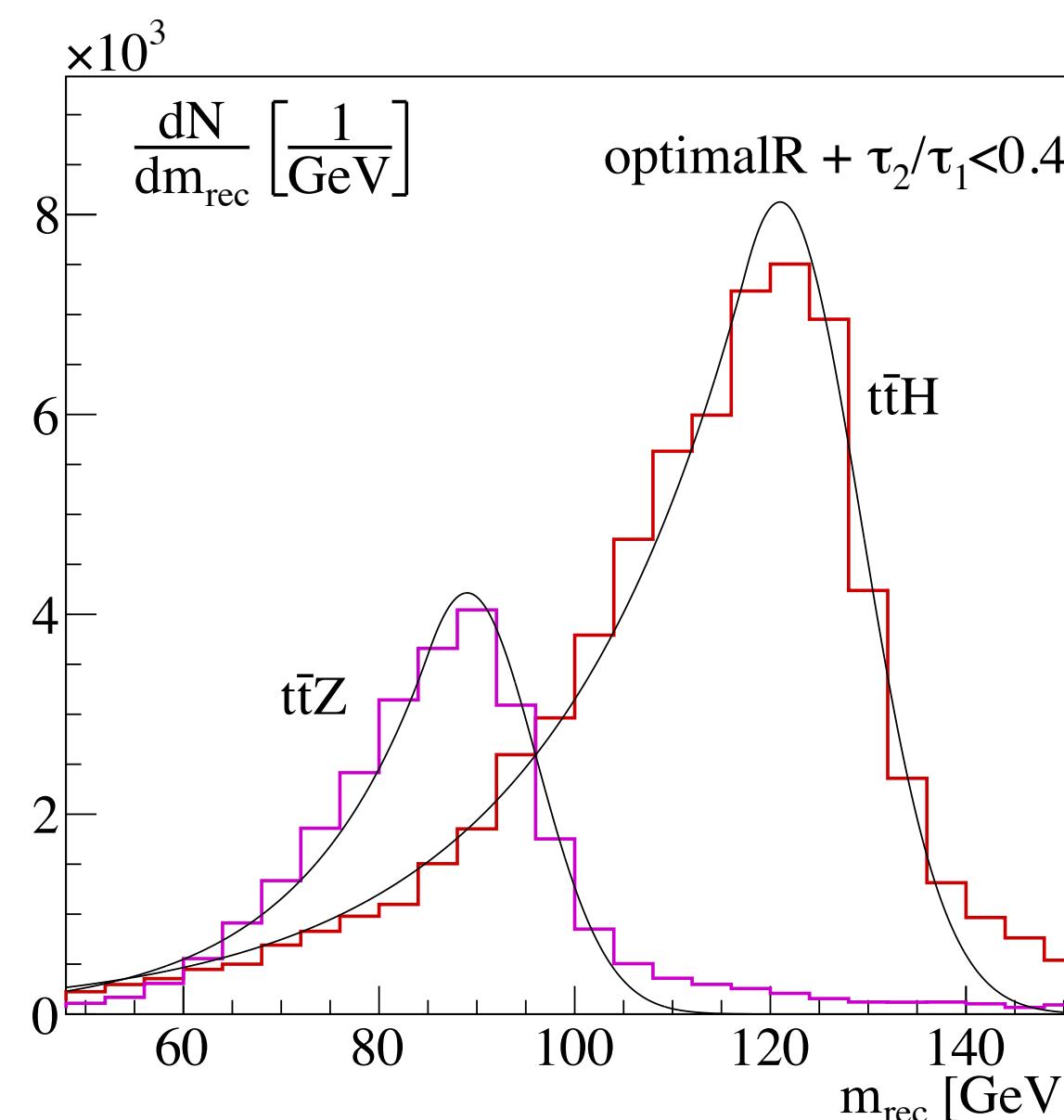
At this point, it compares relatively well with the kappa analysis for the SM-like couplings

# Going back to Higgs complementarities...

- Top Yukawa coupling not directly accessible at FCC-ee. Could be measured in single tH at FCC-eh
  - Can be measured at FCC-hh from  $\sigma(t\bar{t}H)/\sigma(t\bar{t}Z)$  boosted



M.L. Mangano et al., arXiv: 1507.08169 [hep-ph]



e.g. Fit and extract  $N_H/N_Z$  to 1% accuracy  
⇒  $\delta_{\text{stat+th}} y_t/y_t \sim 1\%$

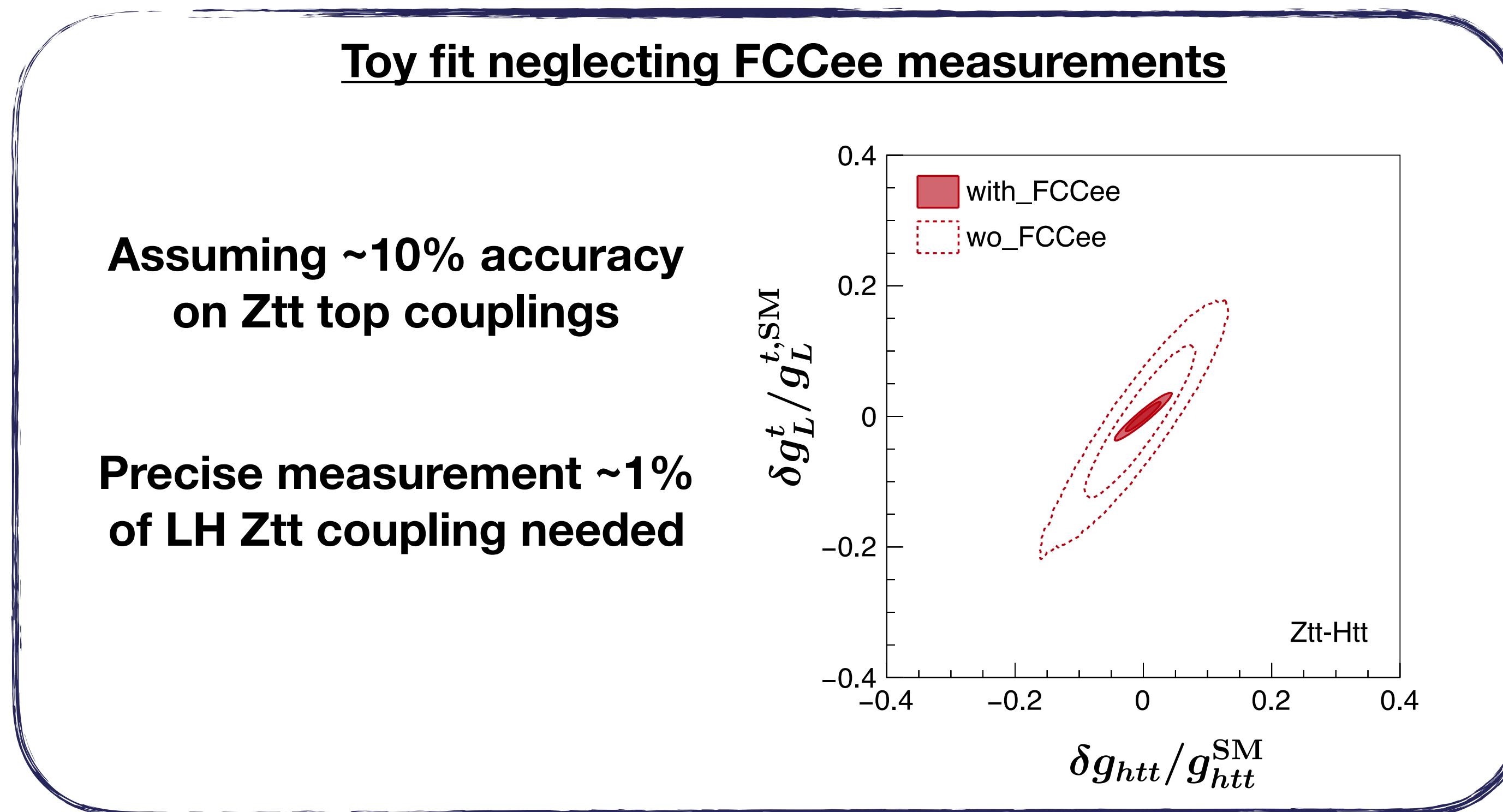
Assumes no NP in  $Ztt$  and  $BR(H \rightarrow bb)$  known to 1%

Both measurable at FCC-ee with required precision

- Robust determination by this method requires both FCC-hh and FCC-ee

# Going back to Higgs complementarities...

- Top Yukawa coupling not directly accessible at FCC-ee. Could be measured in single tH at FCC-eh
  - Can be measured at FCC-hh from  $\sigma(ttH)/\sigma(ttZ)$  boosted



- Robust determination by this method requires both FCC-hh and FCC-ee

# Going back to Higgs complementarities...

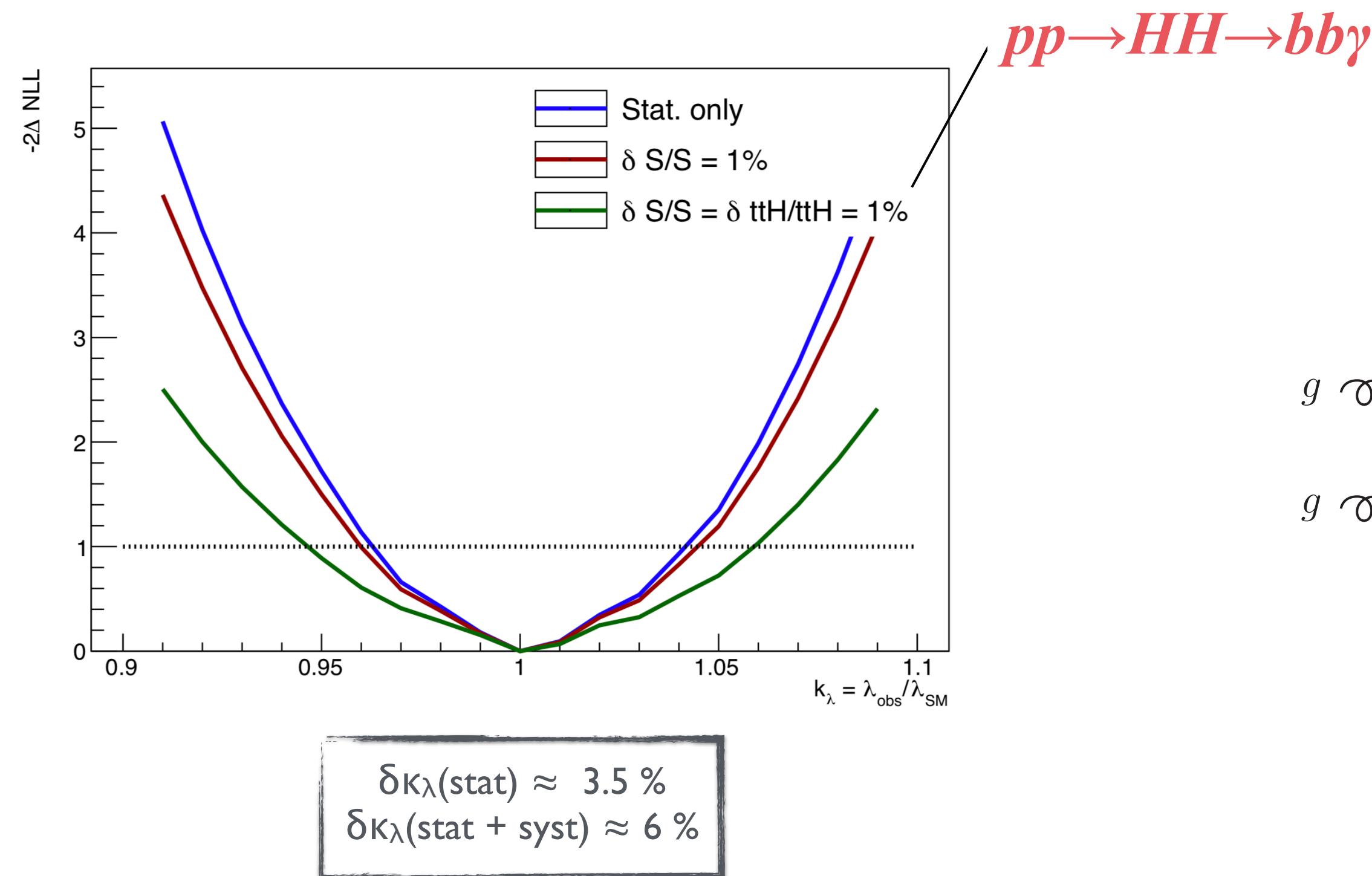
- **Higgs self-interaction:**

- **Can be tested at FCC-ee via NLO effects: Limited (~50%) precision**

S. Di Vita et al., JHEP 1802 (2018) 178

- **Direct HH production at FCC-hh:**

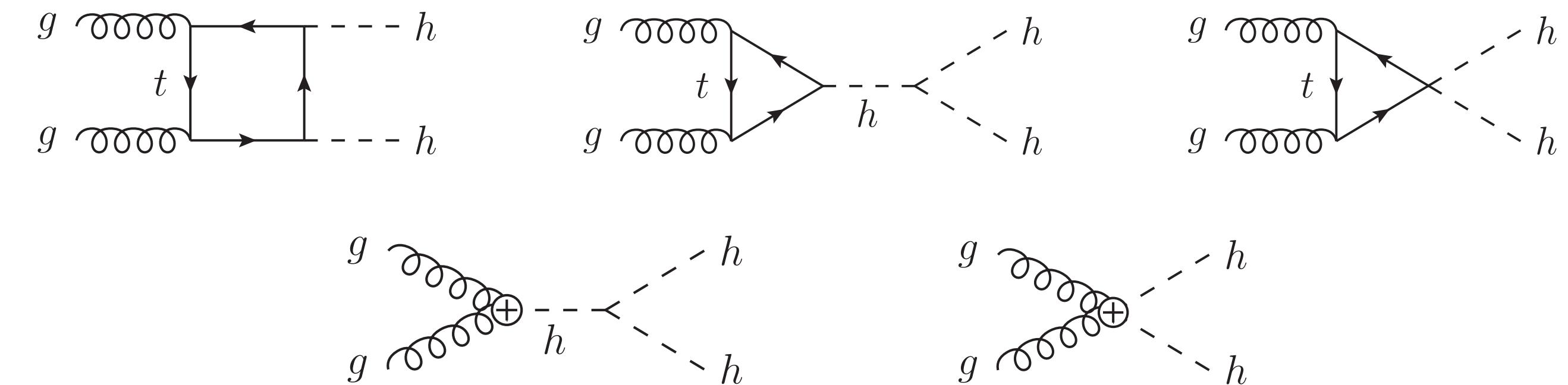
**Assumes all uncertainty goes into  $\kappa_\lambda$**



M. Selvaggi, Talk at 2nd FCC Physics Workshop

$$\delta\kappa_\lambda = \left[ 3(C_{\phi\square} - \frac{1}{4}C_{\phi D}) - 2\frac{v^2}{M_h^2}C_\phi - \frac{1}{2}\Delta_{GF} \right] \frac{v^2}{\Lambda^2}$$

**But other NP parameters modify HH production and decays**



A. Azatov et al. PRD92 (2015) no.3, 035001

**They can be measured at FCC-ee/eh/hh with ~1% precision → Global FCC fit for robust estimates**

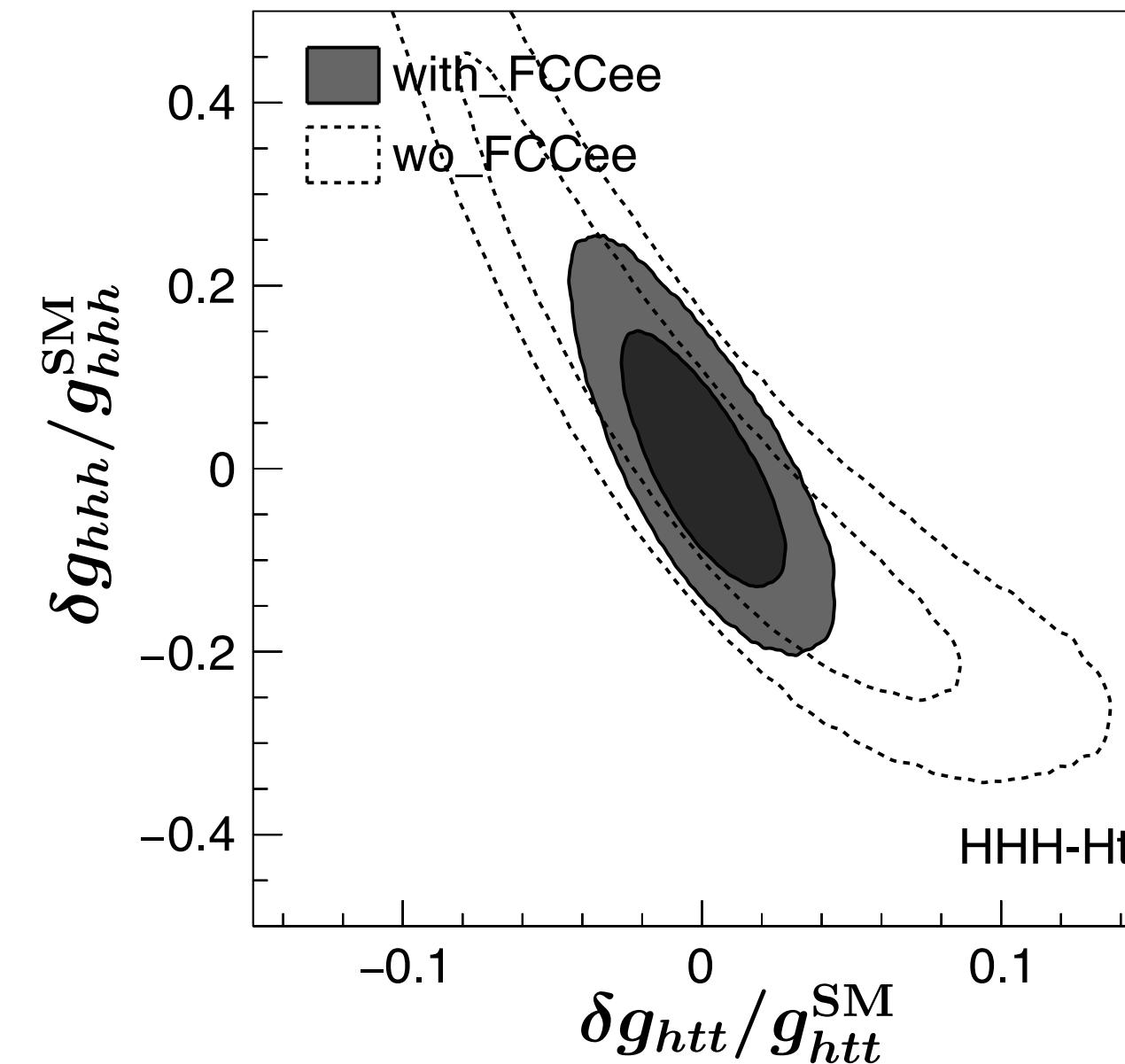
# Going back to Higgs complementarities...

- Higgs self-interaction:
  - Can be tested at FCC-ee via NLO effects: Limited (~50%) precision

S. Di Vita et al., JHEP 1802 (2018) 178

## Toy fit neglecting FCCee measurements

Precise determination  
of operators modifying  
 $Htt$  interactions needed

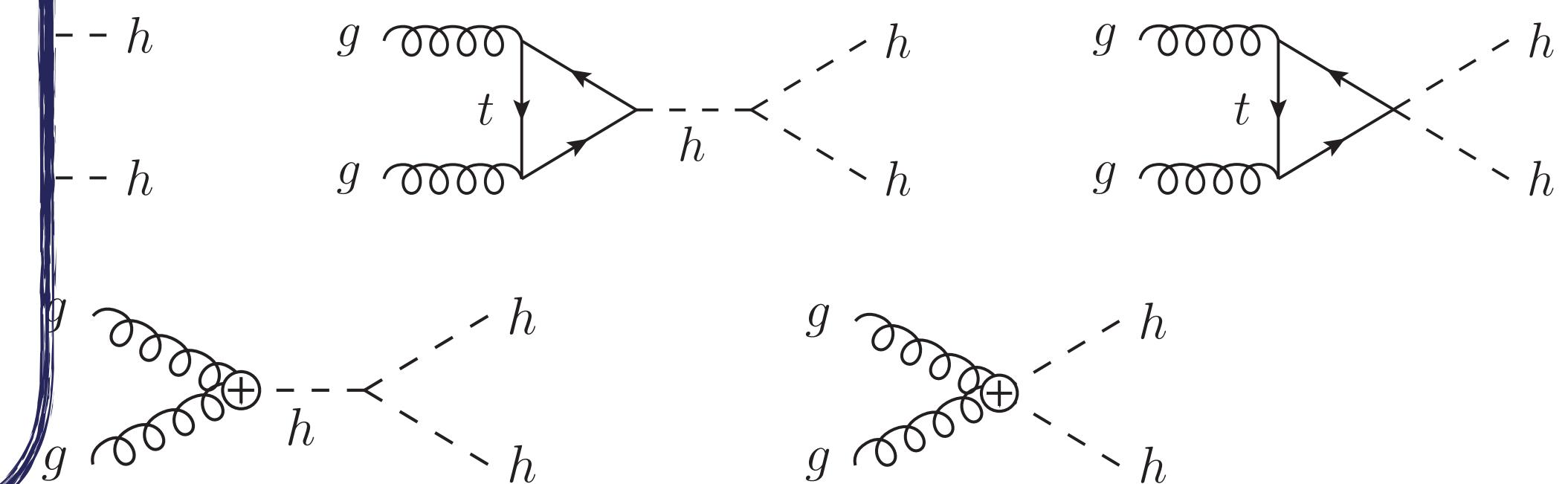


$\delta\kappa_\lambda(\text{stat + syst}) \approx 6\%$

Assumes all uncertainty goes into  $\kappa_\lambda$

$$\lambda = \left[ 3(C_{\phi\square} - \frac{1}{4}C_{\phi D}) - 2\frac{v^2}{M_h^2}C_\phi - \frac{1}{2}\Delta_{GF} \right] \frac{v^2}{\Lambda^2}$$

Other NP parameters modify HH production  
and decays



A. Azatov et al. PRD92 (2015) no.3, 035001

They can be measured at FCC-ee/eh/hh with ~1%  
precision → Global FCC fit for robust estimates

M. Selvaggi, Talk at 2nd FCC Physics Workshop

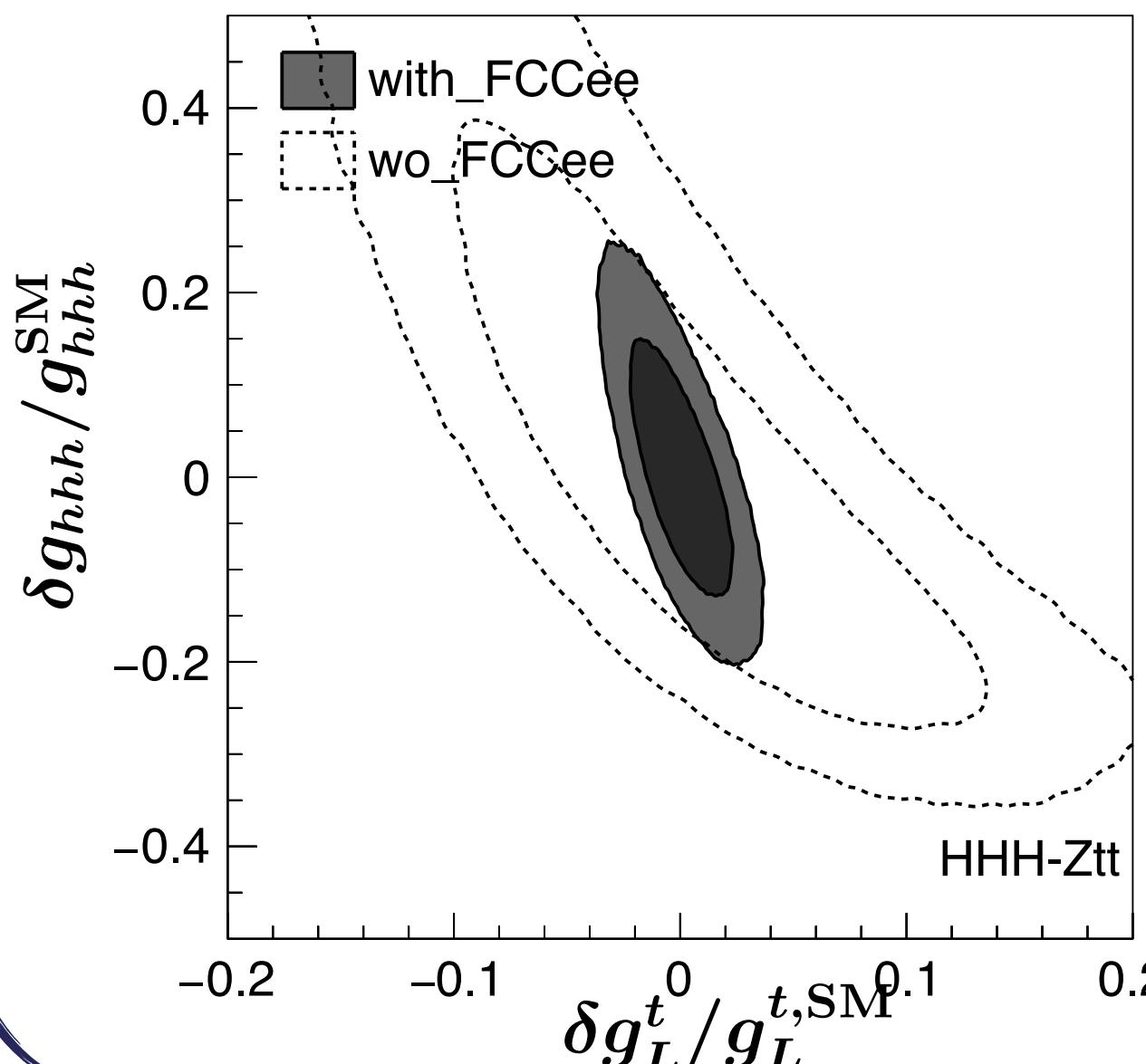
# Going back to Higgs complementarities...

- Higgs self-interaction:

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S. Di Vita et al., JHEP 1802 (2018) 178

## Toy fit neglecting FCCee measurements



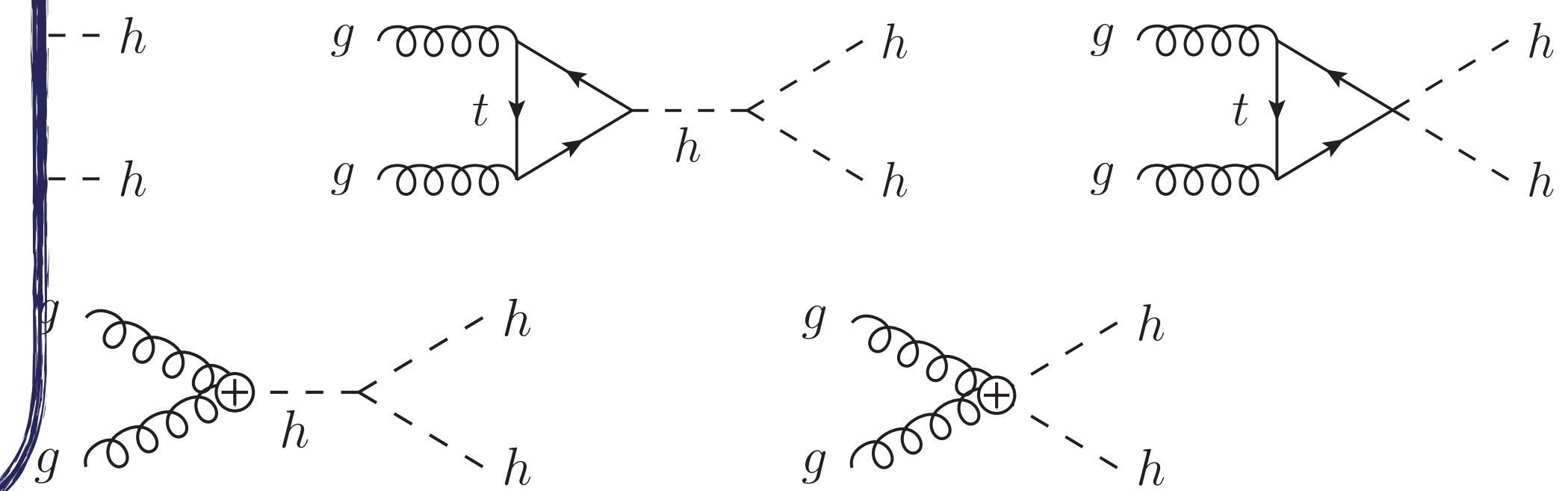
$\delta\kappa_\lambda(\text{stat + syst}) \approx 6\%$

Indirect dependence on  
FCC-ee  $Ztt$  via  
 $Htt/Ztt$  FCC-hh measurement

Assumes all uncertainty goes into  $\kappa_\lambda$

$$\kappa_\lambda = \left[ 3(C_{\phi\square} - \frac{1}{4}C_{\phi D}) - 2\frac{v^2}{M_h^2}C_\phi - \frac{1}{2}\Delta_{G_F} \right] \frac{v^2}{\Lambda^2}$$

Other NP parameters modify HH production  
and decays



A. Azatov et al. PRD92 (2015) no.3, 035001

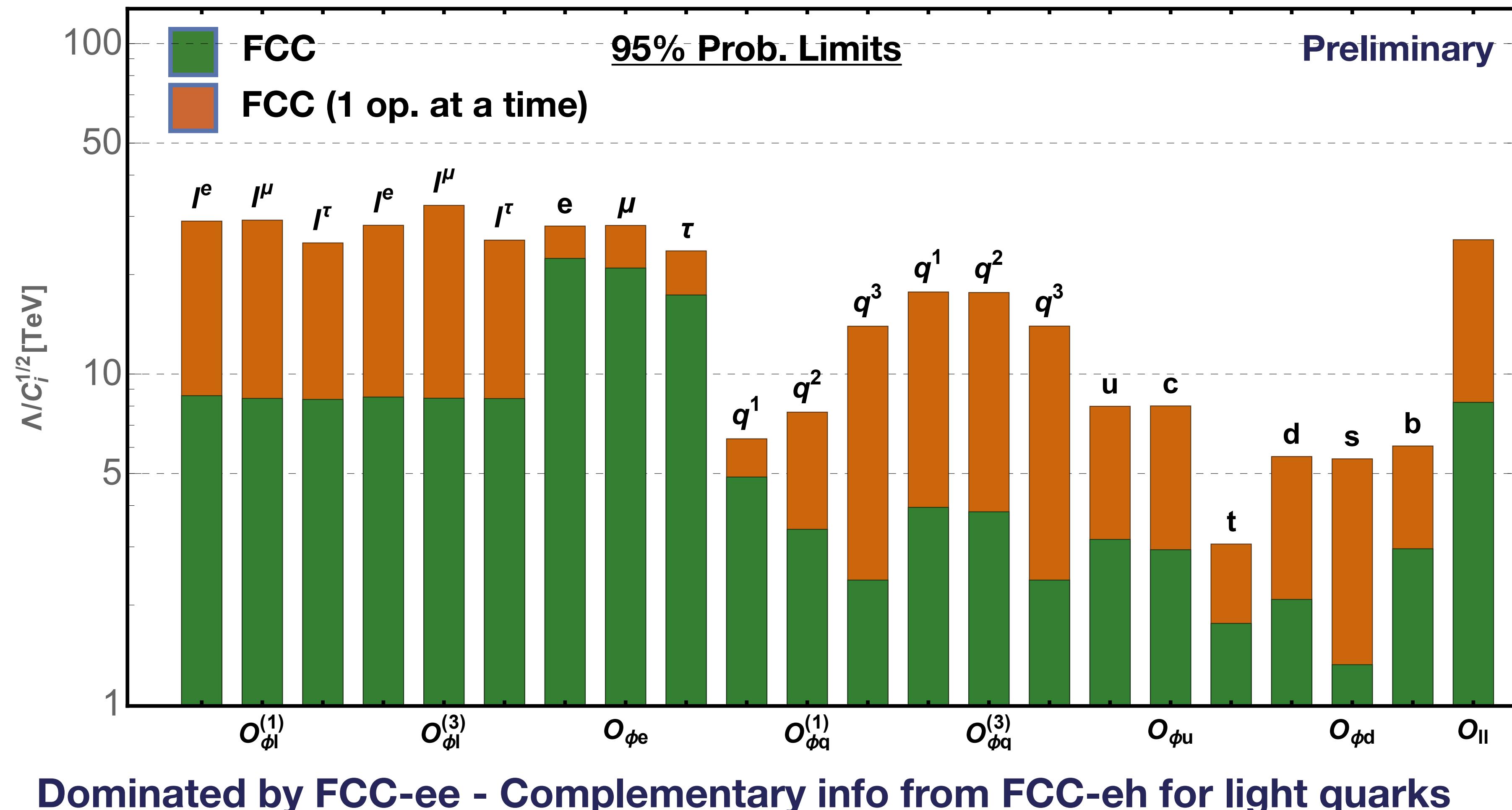
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M. Selvaggi, Talk at 2nd FCC Physics Workshop

# The Global Electroweak and Higgs fit

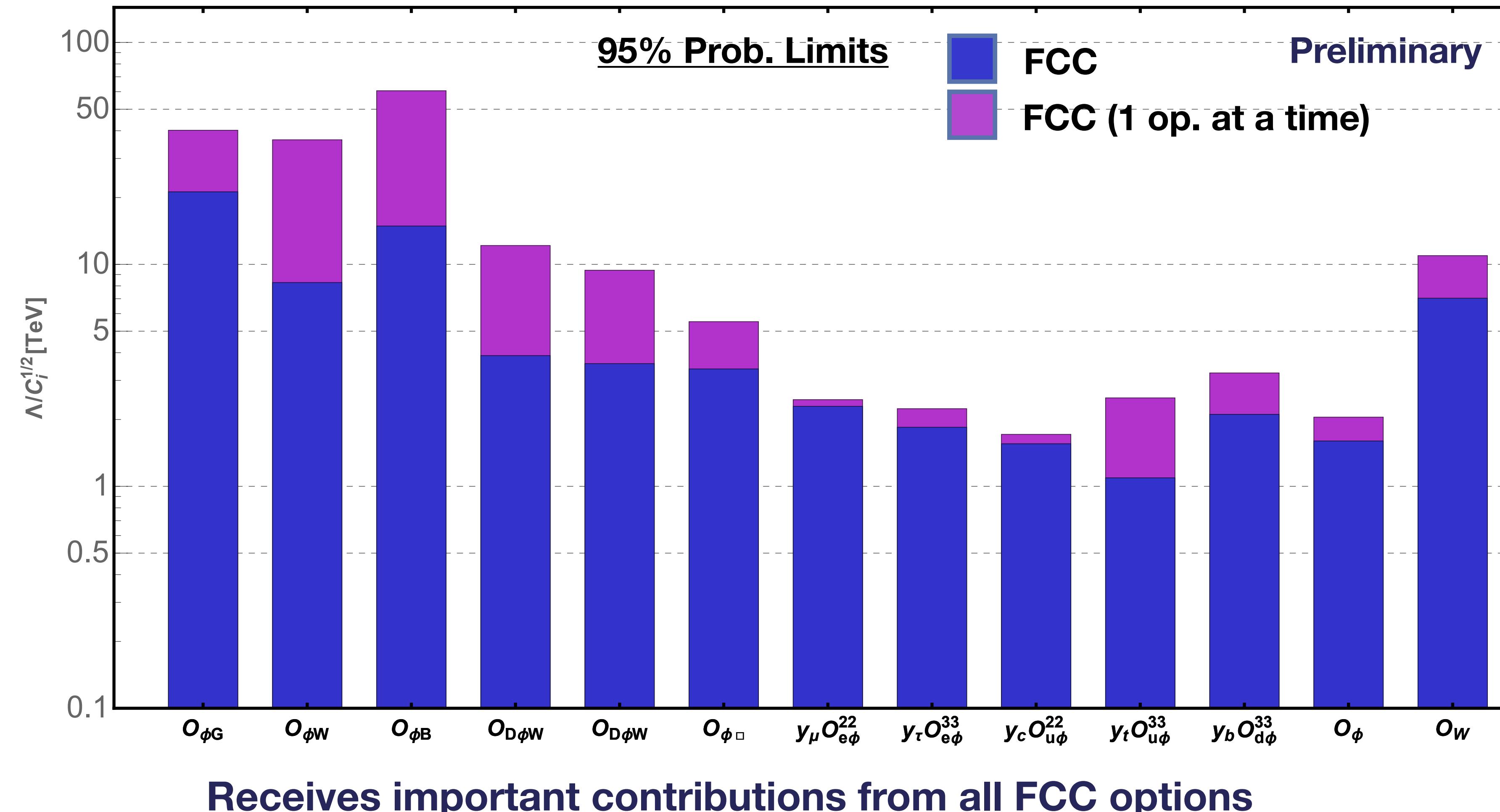
# The Global Electroweak and Higgs fit

- Putting all together... Fit to all 22 (EW) + 13 (Higgs+WW) operators



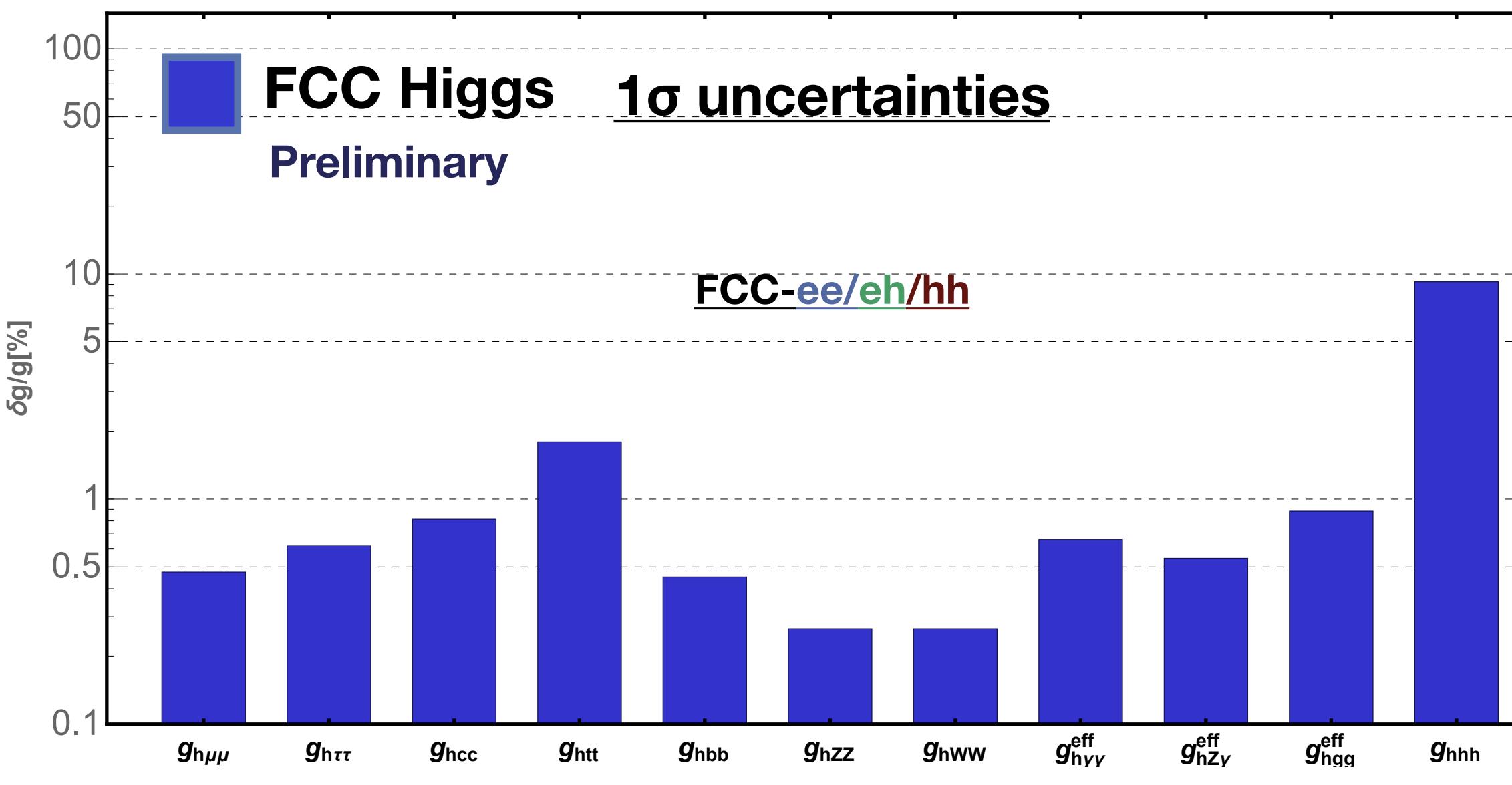
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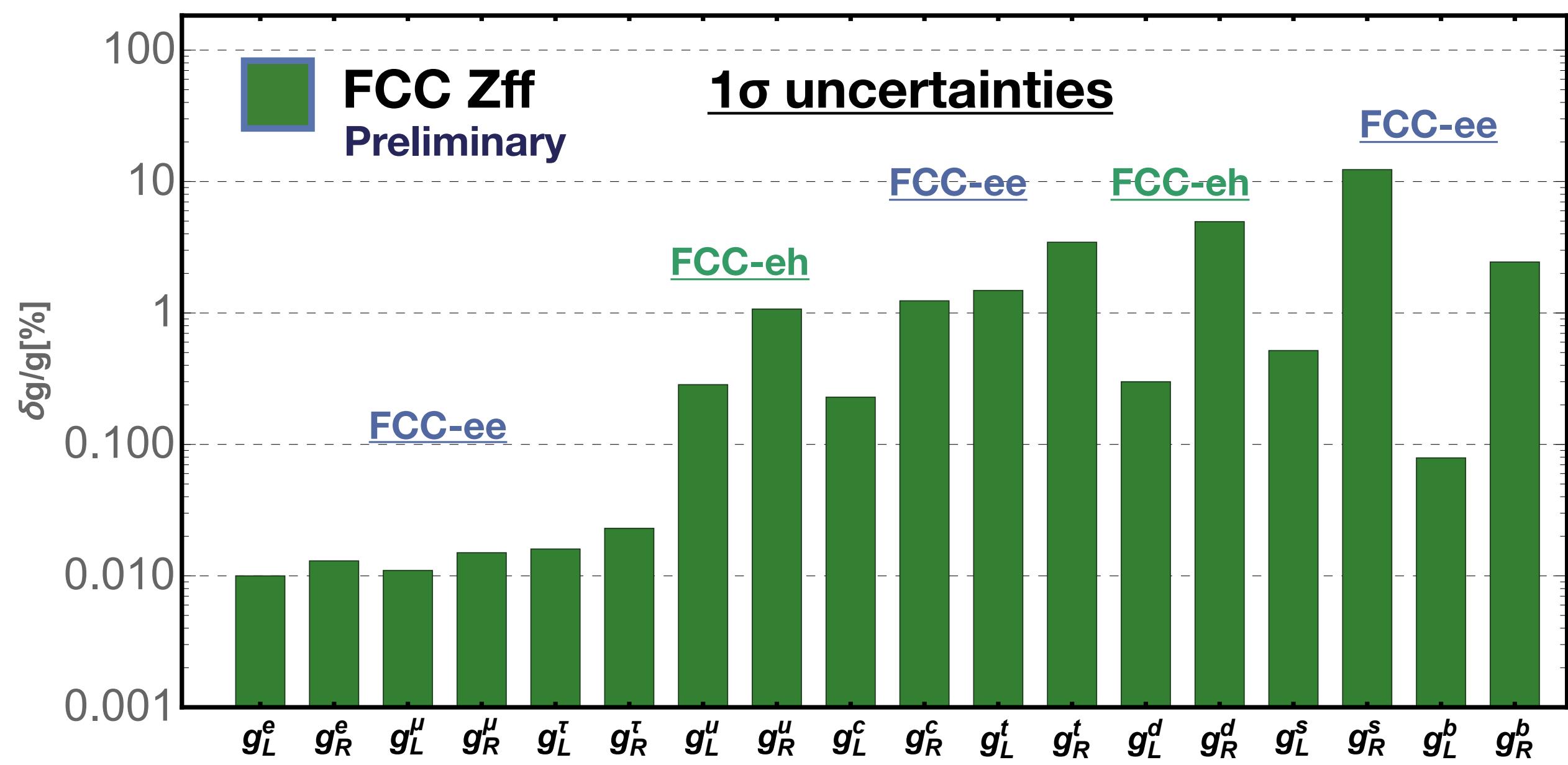


Single Higgs couplings could be known to ~1% or better

Higgs Self-coupling could be known to less than 10%

Independent information for the couplings to all fermion families

Lepton sector: precision down to  $\sim 10^{-4}$   
Quark sector: sub-percent level in most cases

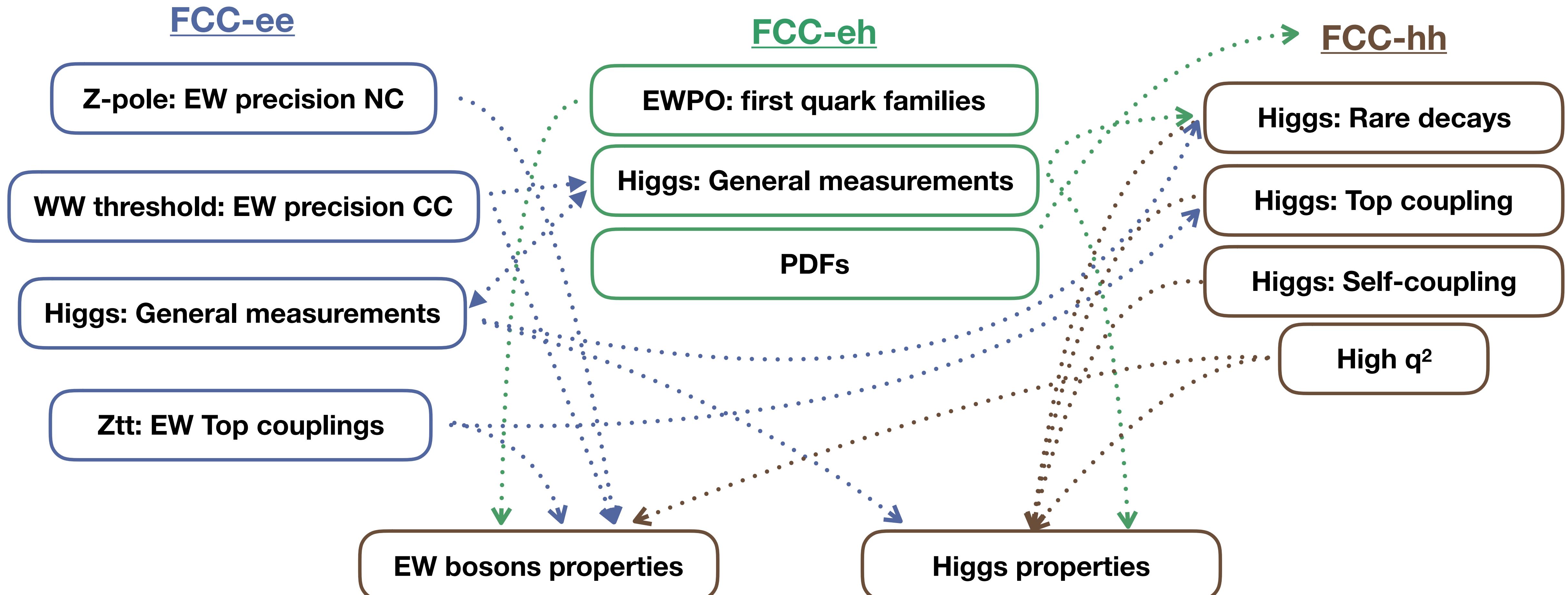


# Conclusions

- A global fit in a model-independent theoretical framework is required to accurately assess the sensitivity to new physics in the different EW and Higgs observables
- The FCC can test (indirectly) new physics interaction scales at the level of several tenths to TeV and increase the precision of properties of the SM particles by, roughly, one order of magnitude
- All three FCC options complement each other very well:
  - FCC-ee allows not only very precise measurements of the Higgs and EWPO but also provides the normalization for more precise measurements at the FCC-eh and FCC-hh
  - FCC-eh complements FCC-ee providing information about light quark EW couplings. Similar precision in the Higgs sector
  - FCC-hh fills gaps in precision Higgs measurements for rare decays, top and the Higgs self-coupling

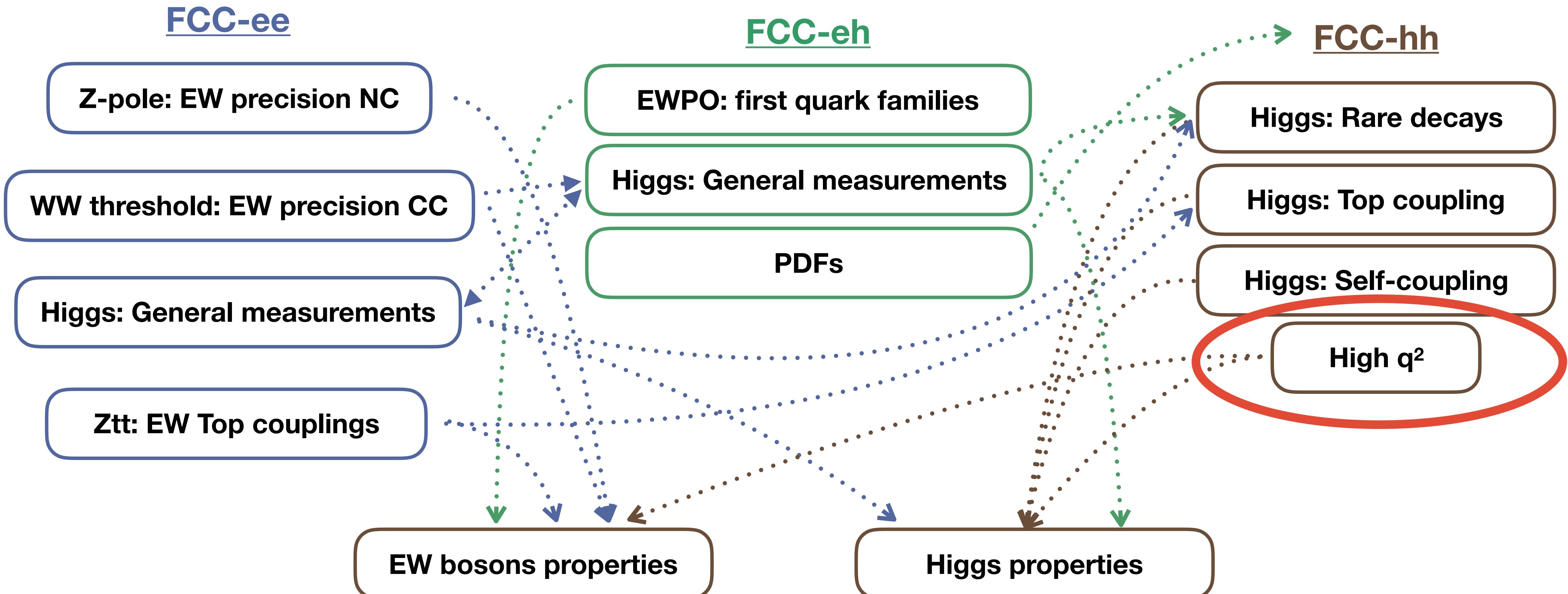
# Conclusions

- All three FCC options complement each other very well and are useful to complete the whole picture:



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# Conclusions

- The dimension 6 SMEFT:

$$\mathcal{L}_{\text{Eff}} = \sum_{d=4}^{\infty} \frac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

$$\mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \quad [\mathcal{O}_i] = d \longrightarrow \left(\frac{q}{\Lambda}\right)^{d-4}$$

$\Lambda$ : Cut-off of the EFT

Effects suppressed by  $q = v, E < \Lambda$

Most of the effects discussed so far

For  $E \gg v$  these effects can provide precise constraints on EFT interactions even if experimental precision is lower

Large Energies  $\Rightarrow$  FCC-hh

Look for  $E$ -enhanced effects in differential distributions

See A. Wulzer's talk this morning

# Conclusions

- The dimension 6 SMEFT:

$$\mathcal{L}_{\text{Eff}} = \sum_{d=4}^{\infty} \frac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

$$\mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \quad [\mathcal{O}_i] = d \xrightarrow{\text{Effects}} \left(\frac{q}{\Lambda}\right)^{d-4}$$

Example: Individual 95% prob bound on  $\frac{C_{\phi q}^{(3)}}{\Lambda^2} \sim \pm 3 \times 10^{-3} \text{ TeV}^{-2}$

For  $E \gg v$  these  
Contributes to WZ production via High-E primary  $a_q^{(3)} = 4 \frac{C_{\phi q}^{(3)}}{\Lambda^2} \sim \pm 12 \times 10^{-3} \text{ TeV}^{-2}$

Look  
Study of  $p_{\text{TV}}$  distrib. at FCC-hh  $20 \text{ ab}^{-1}$  ( $\delta_{\text{sys}}=5\%$ )  $\Rightarrow a_q^{(3)} \sim \pm 6 \times 10^{-3} \text{ TeV}^{-2}$

Energy helps to improve precision constraints on new physics

See A. Wulzer's talk this morning

To be included in updated global fits

# Thank you