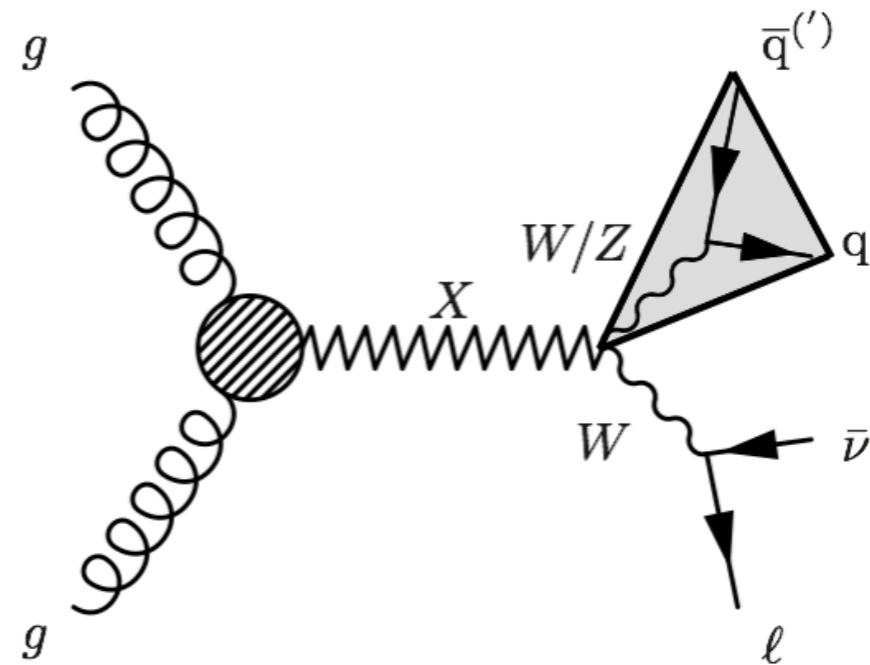




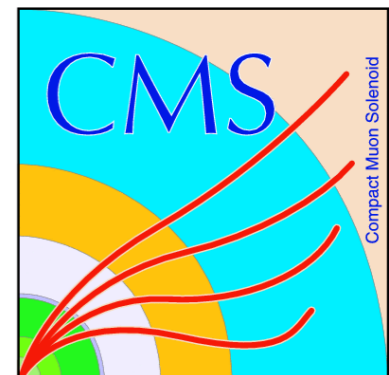
Search for heavy resonances in diboson final states at CMS

Highlights of the CMS diboson resonances search programme



Clemens Lange (CERN)
on behalf of the CMS Collaboration

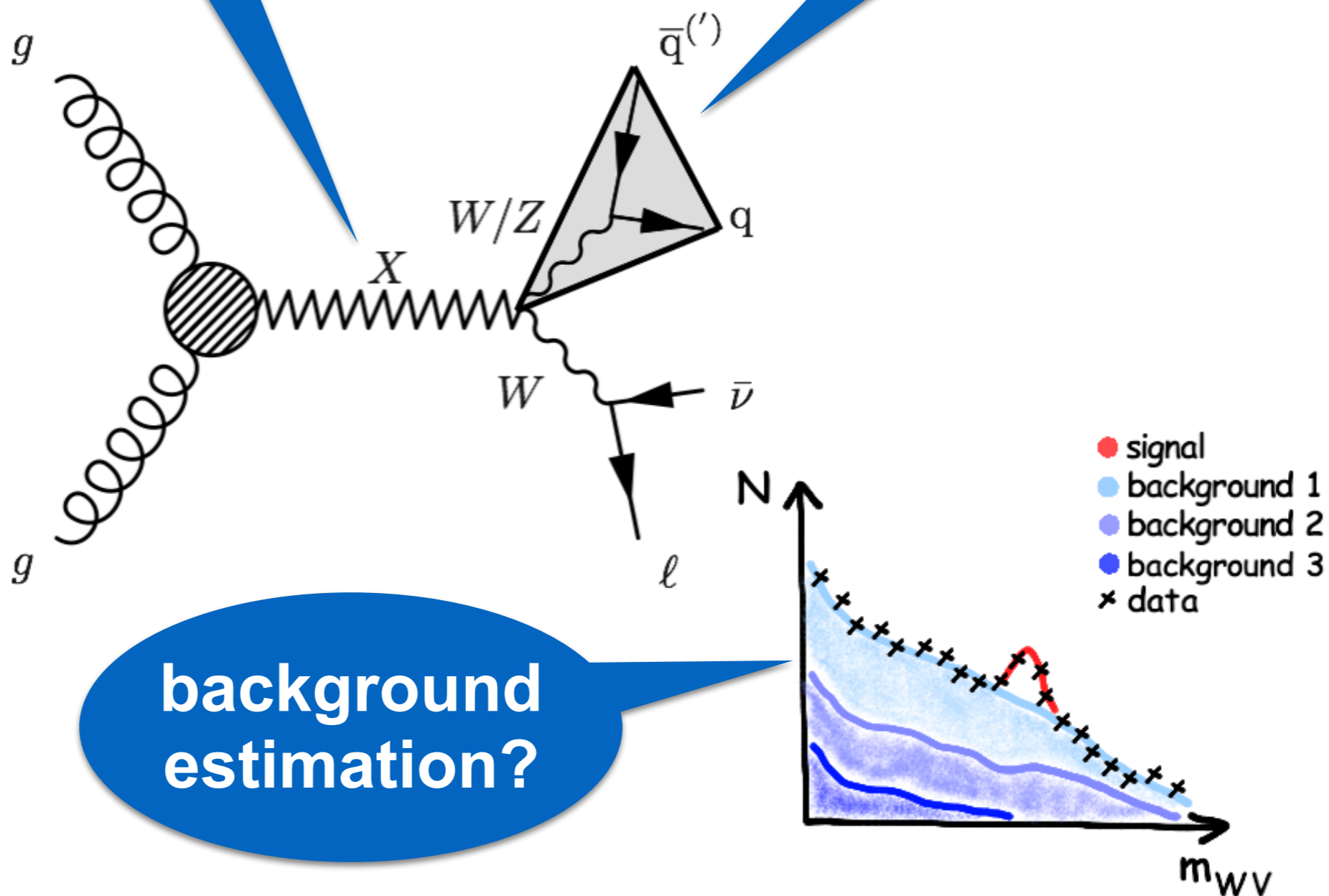
CERN-LHC Seminar
5th December 2017



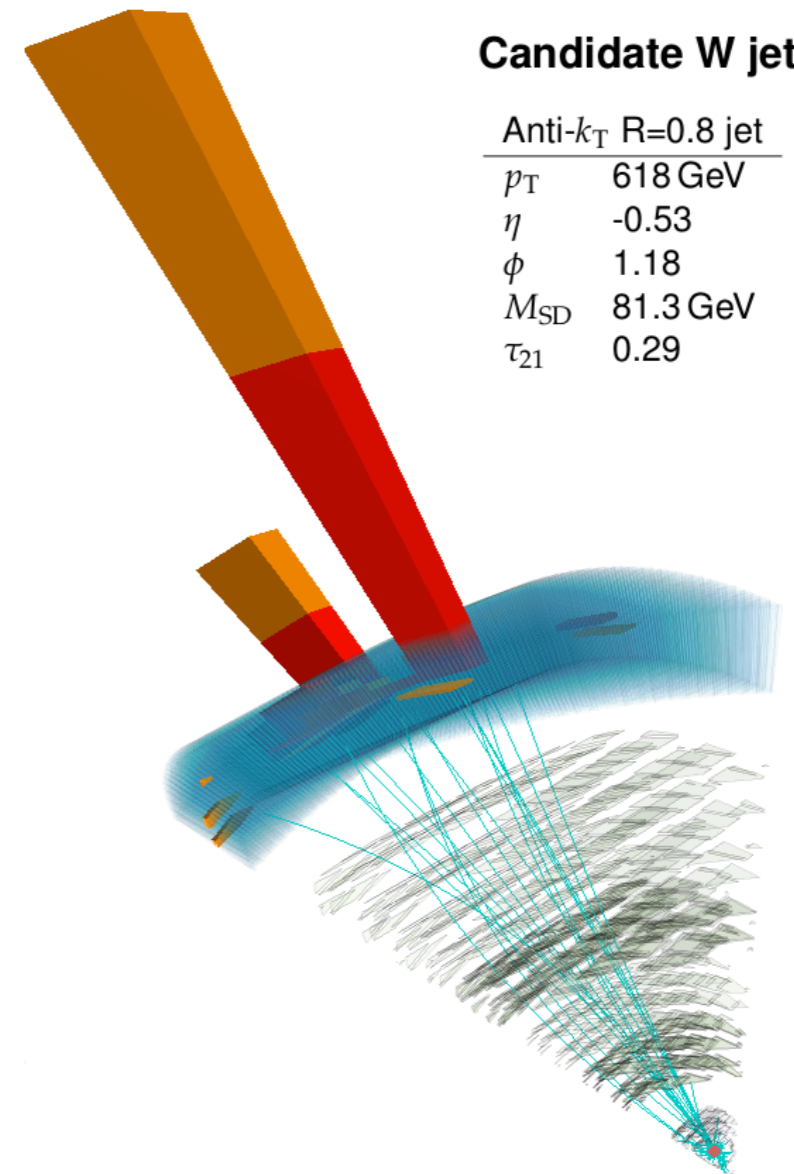
heavy resonances?

reconstruction of bosons?

background estimation?



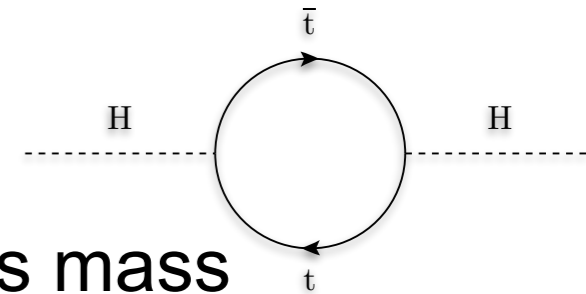
- > Why search for diboson resonances?
- > Boson reconstruction
 - jet substructure
- > Diboson resonance searches
 - all-hadronic VV/Vh final states
 - di-Higgs (hh) final states
 - (semi-)leptonic VV final states



Results of 4 analyses shown publicly for the first time today

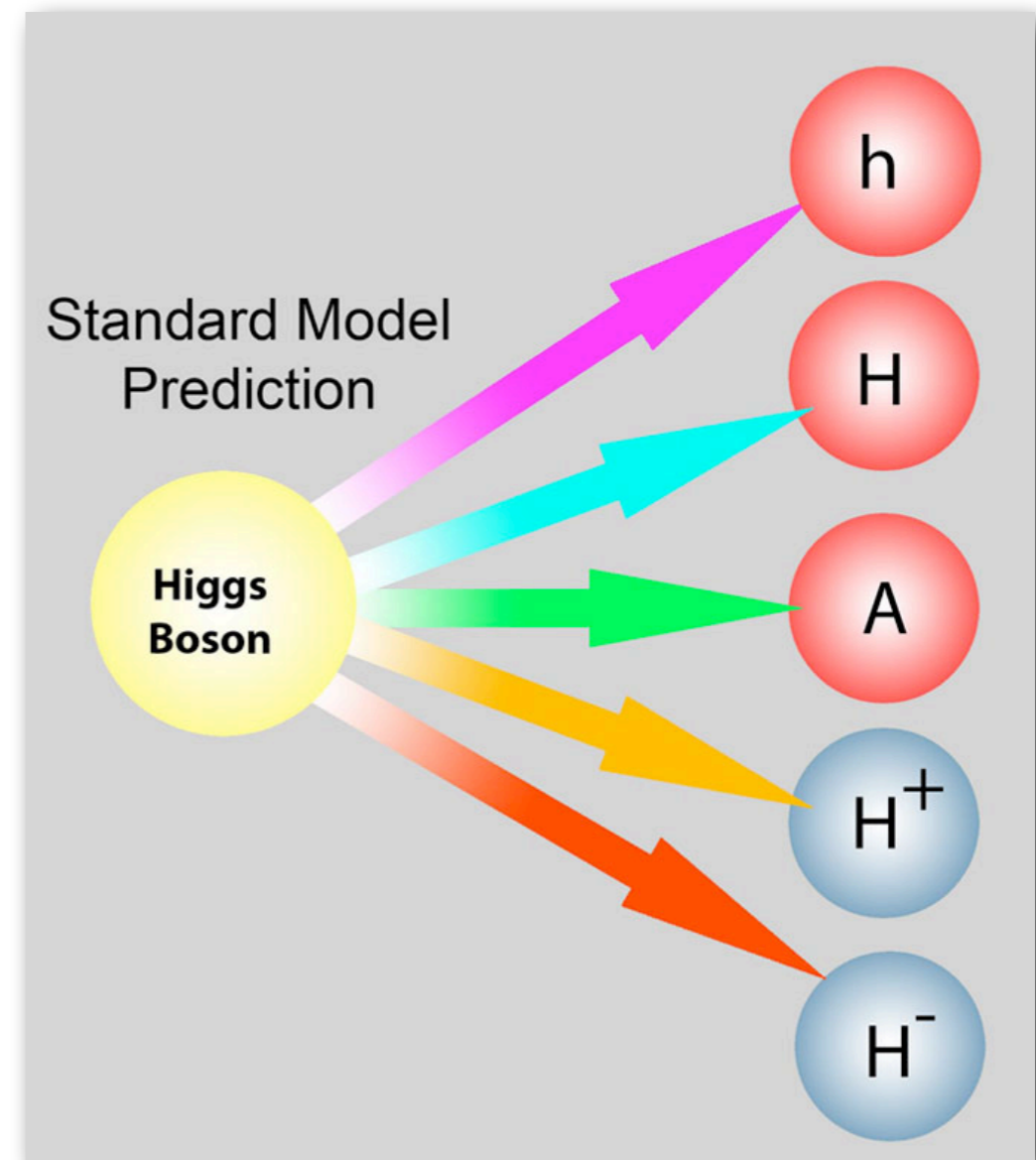
- > how can we explain the big difference between the **weak force** and **gravitation**?

$$\mu^2 = \lambda v^2 = \frac{\lambda}{g^2} 4M_W^2 \sim 10^4 \text{ GeV}^2 \ll M_{Pl} \sim 10^{38} \text{ GeV}^2$$



- > no **symmetry** in the standard model (SM) protects the Higgs mass
- > „natural“ explanation would be that SM is replaced/extended by another theory at the TeV scale: $\mu^2 \sim (\text{heavier scale})^2 \rightarrow$ **new particles**
- > these theories could be (among others):
- > **SUSY**: protecting the Higgs mass by a symmetry
- > **Composite Higgs**: the Higgs is not elementary
- > **Large/warped extra dimensions**: gravity is strong at electroweak scale

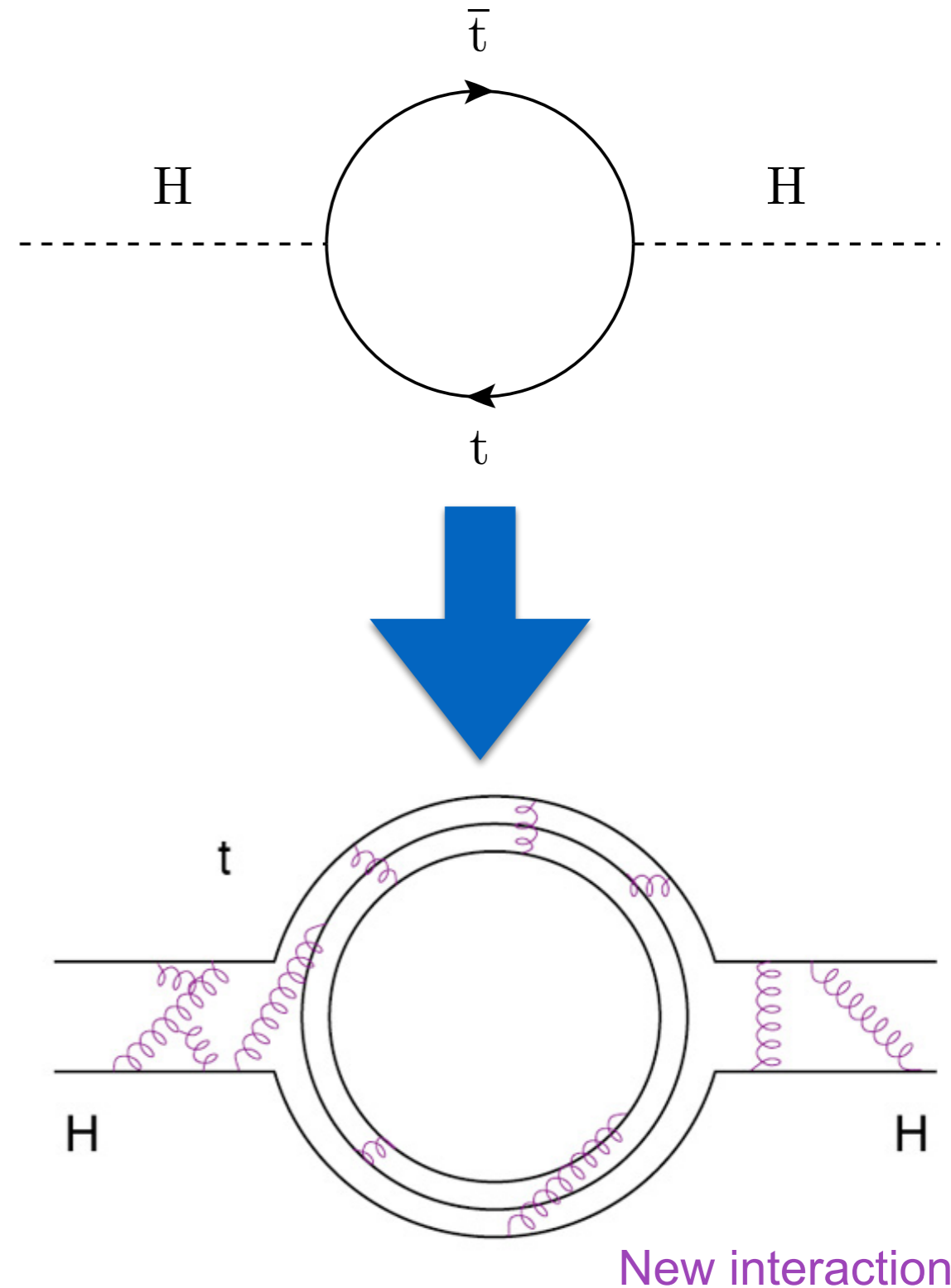
- > 2-Higgs-Doublet-Models (2HDM)
 - one of the simplest extensions of the SM
 - motivation from SUSY, axion models, baryon asymmetry, ...
- > usually two complex scalar SU(2) doublets leading to five scalar Higgs fields:
 - charged scalar (H^\pm), two neutral scalars (H and h), one pseudoscalar (A)
- > Here: light scalar (h) 125 GeV boson
 - heavier particles' mass $> 2 \times m_W / m_Z / m_H$



Here:
 mostly scalar (spin-0) resonances
 decaying to pairs of SM h bosons
 (decays into W/Z pairs also possible)

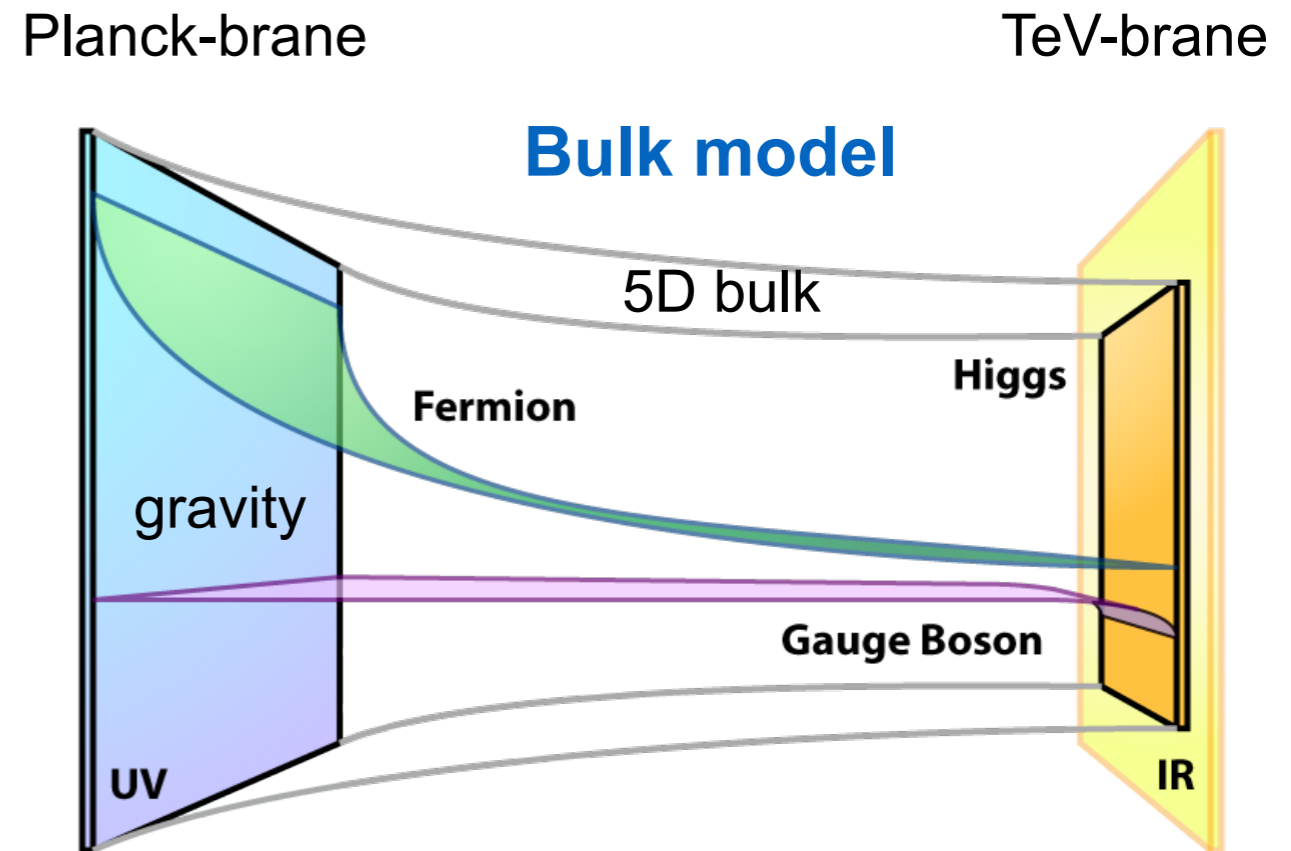
- > The **Higgs** could be **non-fundamental**
- > Instead: **bound state** of a new strong interaction
- > Brings along **new heavy particles/** states
- > Heavy partners of SM particles decay to lighter ones (W, Z, h, top, ...)

Here:
spin-1 (W' , Z') resonances
decaying to pairs/combinations of W, Z, h
interpretation in “heavy vector triplet”
framework (W'/Z' , V')



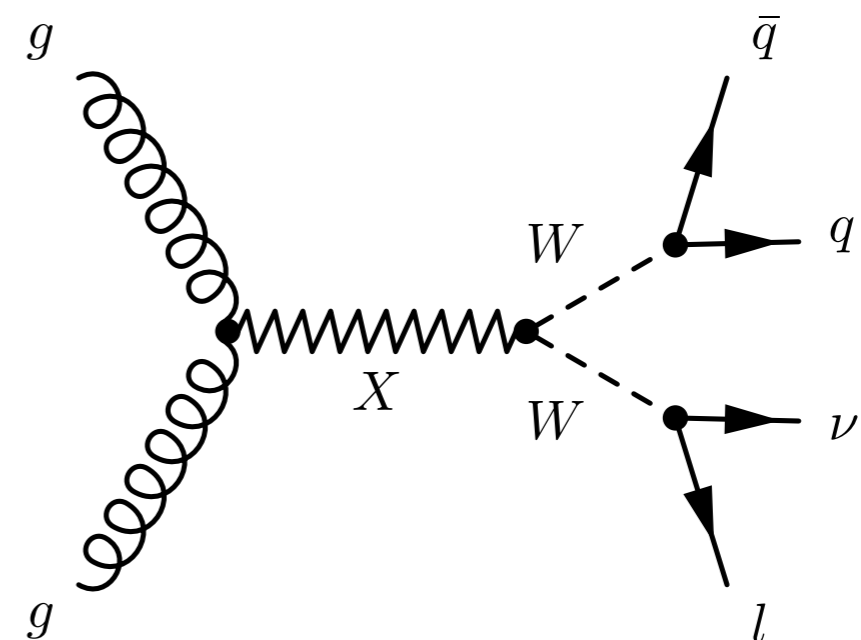
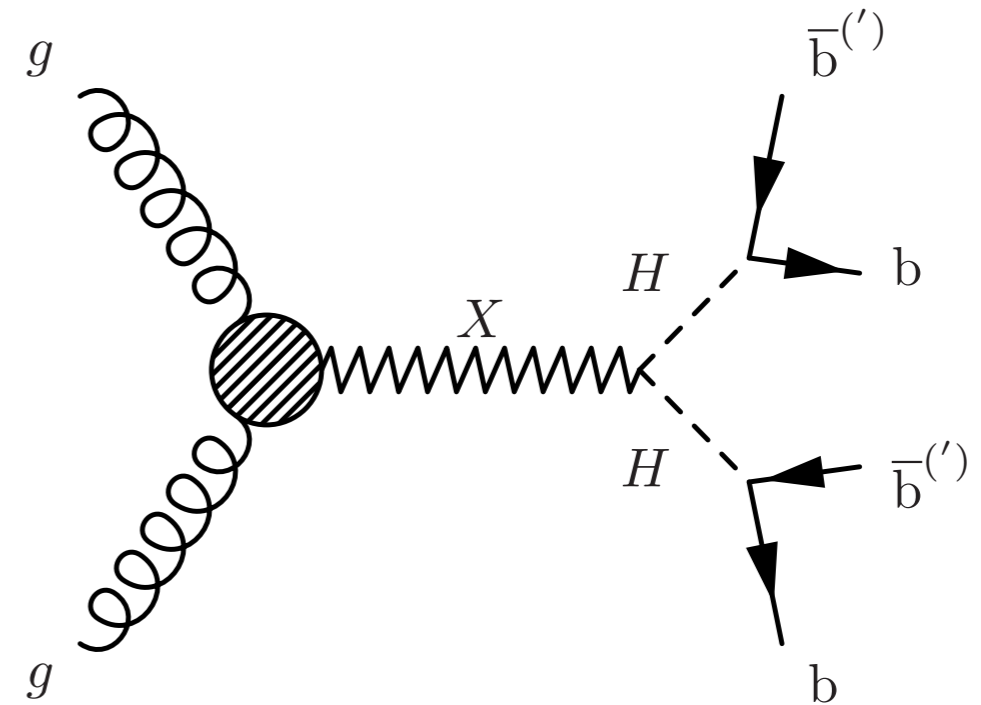
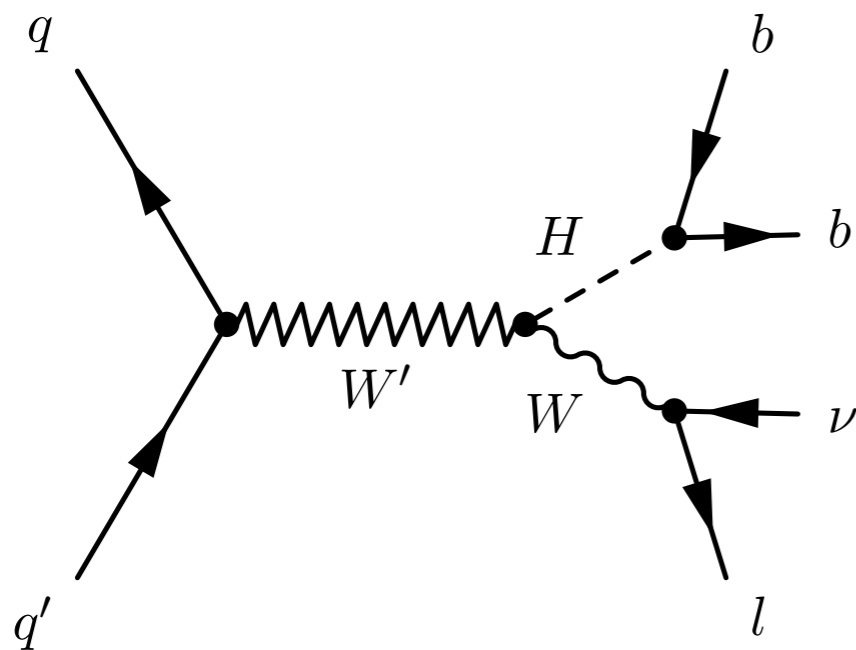
(Warped) extra dimensions?

- > SM fields are confined to **four-dimensional „membrane“**, gravity propagates in additional dimensions
- > Change effective Planck constant, **reducing Planck scale** to close **to electroweak scale**
- > Overlap of 5D profiles at TeV-brane (and Higgs) determine **particle masses**
- > Additionally, if distance between two branes is not fixed, **additional fluctuations** can occur



Here:
spin-2 gravitons
 decaying to pairs of W, Z, (γ)
 —
spin-0 radion fluctuations
 decaying to pairs of SM h bosons

- > In **diboson final states**, should be able to observe excitations/resonances/fluctuations
- > Spin 0, 1, or 2
- > Depending on **model parameters**, resonances can be narrow or wide
 - model-independent interpretation important
- > Majority of analyses presented here focus on **narrow resonances** (width < detector resolution), **mass ≥ 130 GeV**



➤ Performed **statistical combination** of all CMS $\sqrt{s} = 8$ TeV (2012) and 2015 13 TeV **diboson high mass analyses**

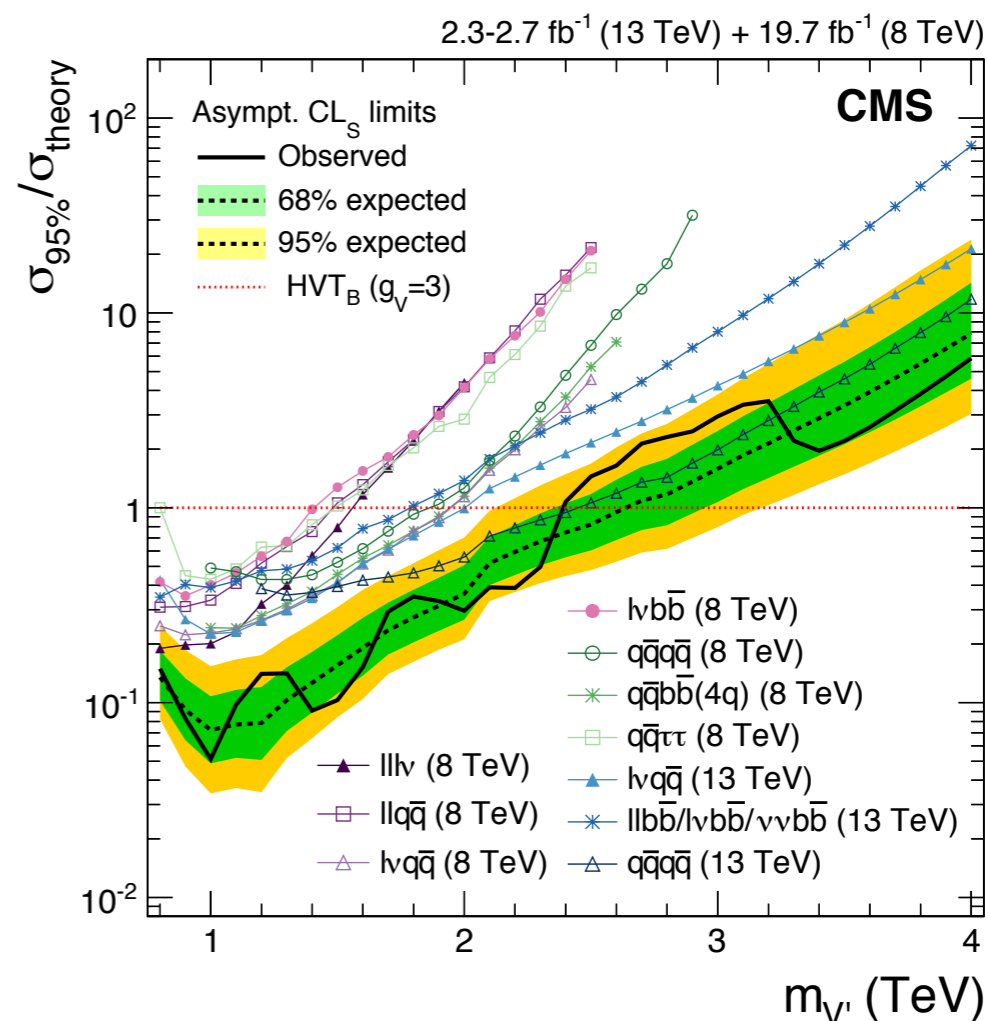
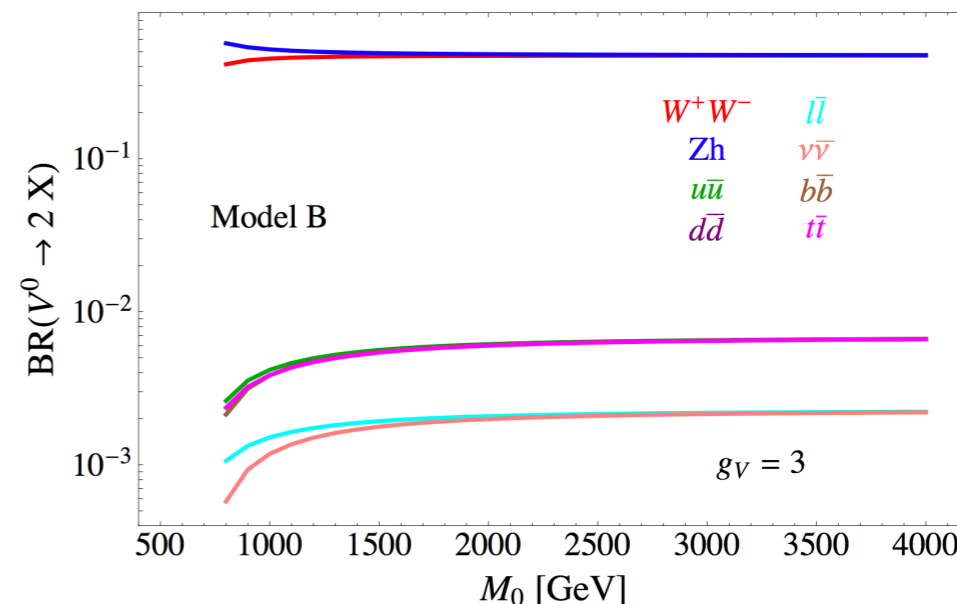
➤ Benchmark models:

- Heavy vector triplet (HVT) models A and B
- Bulk graviton model

➤ While individual analyses showed deviations from SM background expectation $> 2\sigma$, combination showed none

➤ Full **2016 data (35.9 fb⁻¹) analyses have significantly higher sensitivity**

- signal cross sections 13 TeV to 8 TeV a factor 2–5 higher



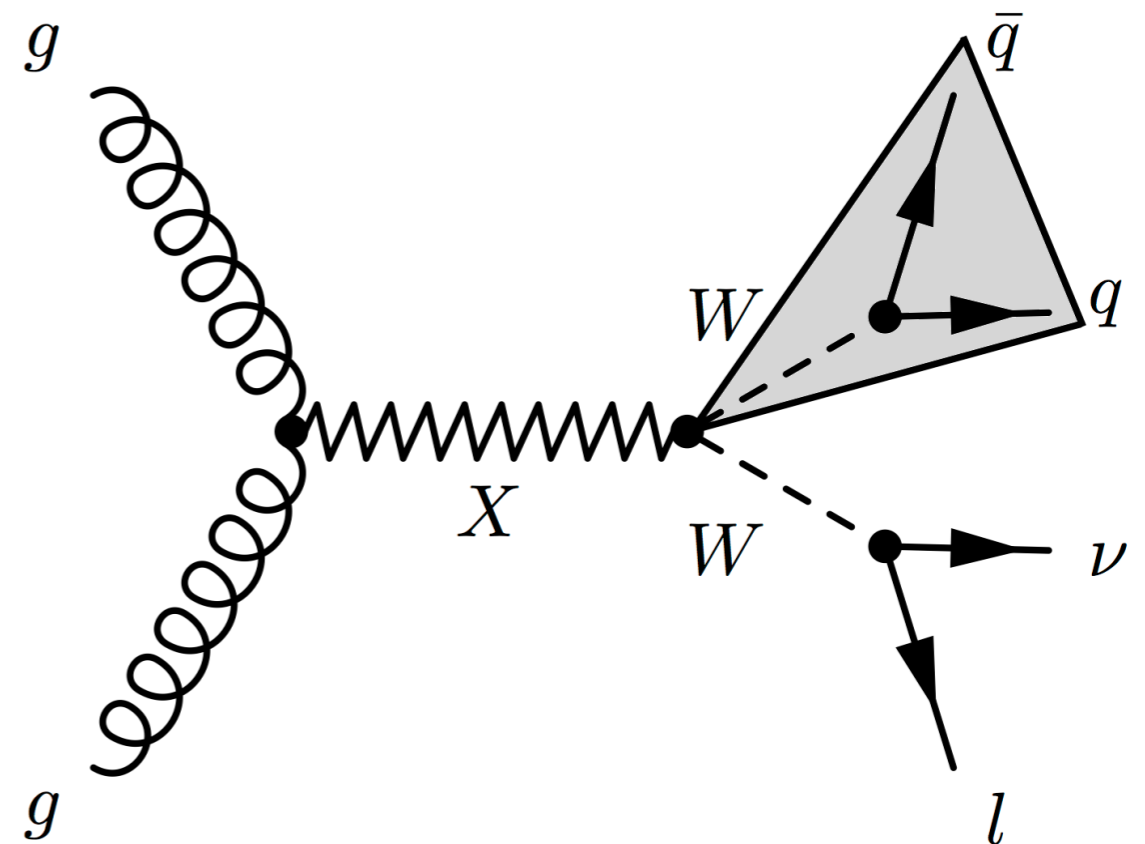
all-hadronic 13 TeV analyses most sensitive



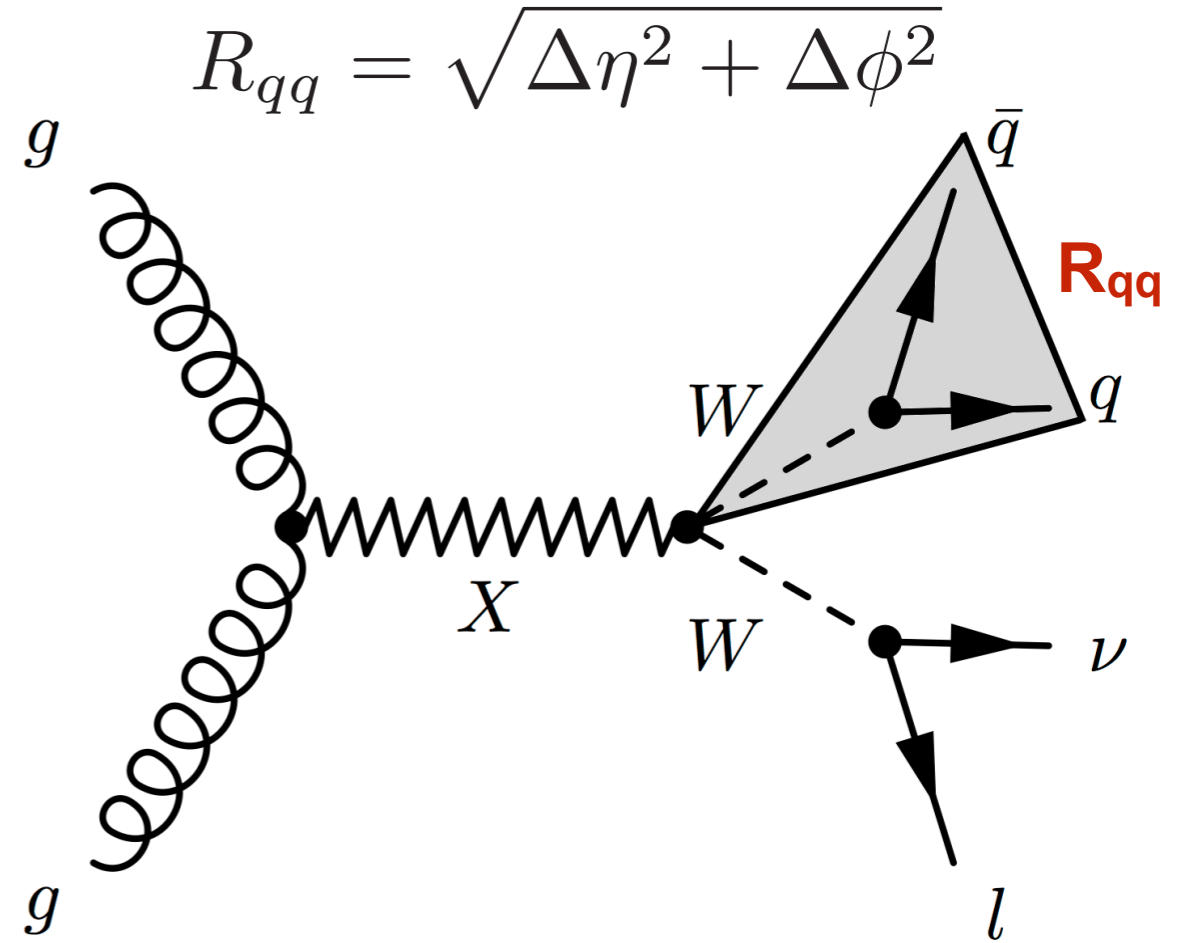
Boson reconstruction



- > For high mass resonances, bosons will be very energetic
 - **collimated decay products**
- > Need to develop dedicated reconstruction methods
- > Hadronic decays of bosons:
 - „**boson-tagging**“
 - exploiting **substructure** of jets
- > Leptonic decays:
 - special **isolation** for dileptonic decays
 - dedicated reconstruction algorithms for high- p_T leptons
 - new τ -identification algorithms



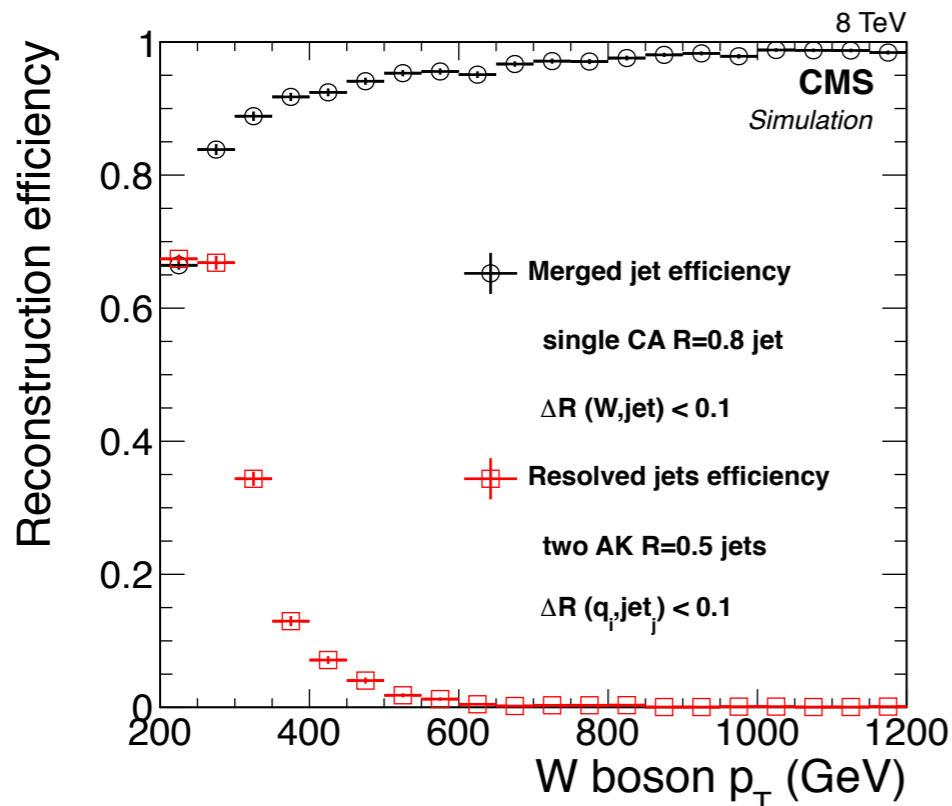
- > At CMS use anti- k_T jet algorithm with $R = 0.4$
- > Already for resonances of **1 TeV** a significant fraction of cases where the boson **decay** is contained **in a single jet**
- > Increase jet size to $R = 0.8$ to contain full decay within „**fat**“ jet



$$R_{qq} \approx 2 \frac{M_V}{p_T^V}$$

back-of-the-envelope calculation:

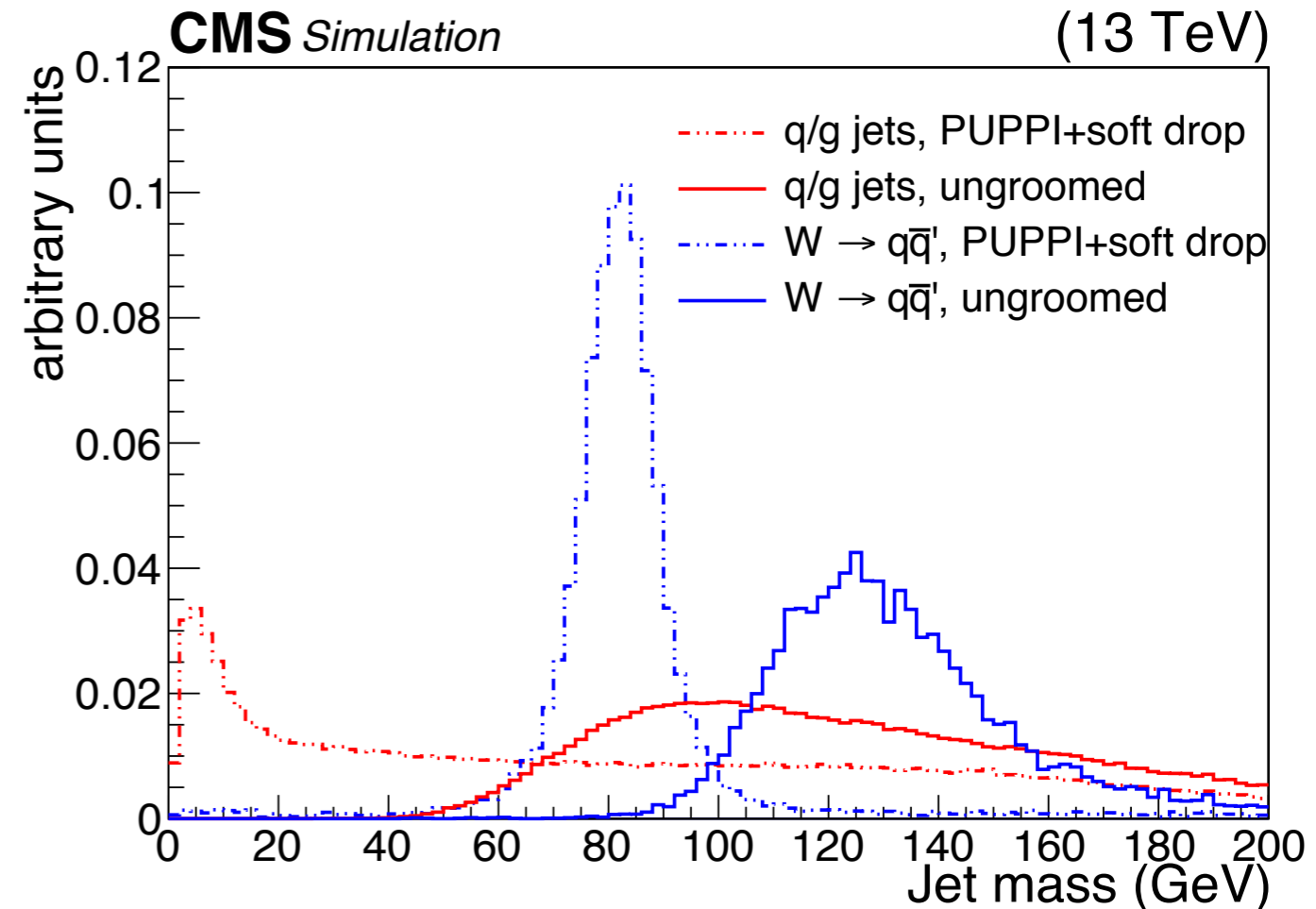
for a resonance of mass 1 TeV the bosons from the decay will have $p_T \sim 0.4$ TeV $\rightarrow \Delta R \approx 0.4$



Jet grooming

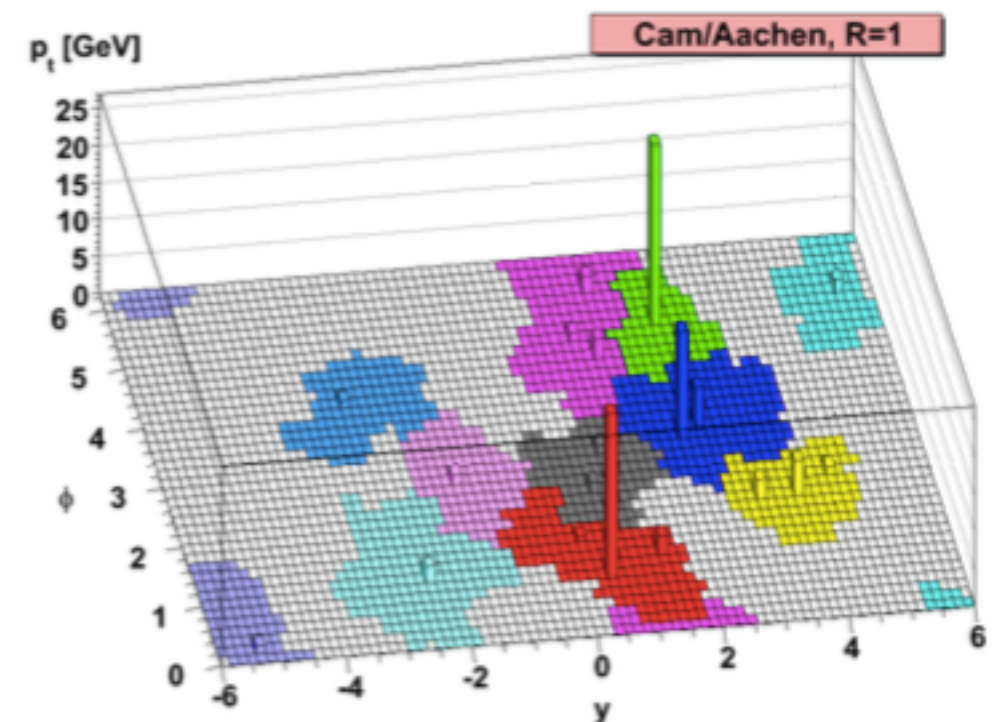
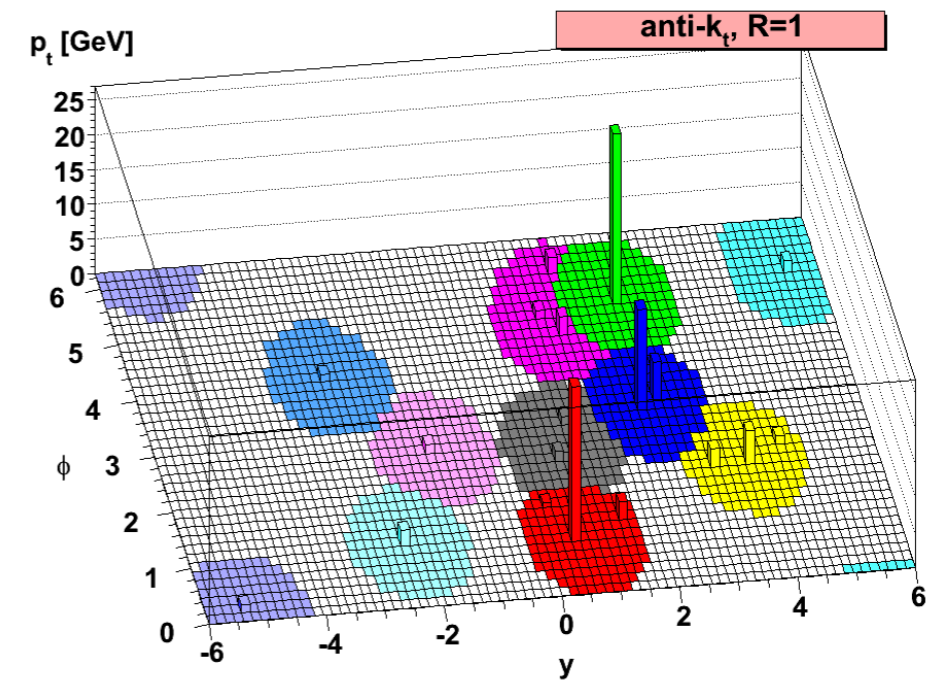
- > We know the **masses** of W, Z and Higgs very well → can use them as **constraints**
- > However, large number of particles in jet → rather bad resolution
- > Jet grooming removes **soft and large angle radiation**
- > Strategy:
 - recluster jet using Cambridge-Aachen (CA) jet algorithm
 - iteratively break into two subjets

?



Reminder: jet clustering algorithms

- > k_T -algorithms: **sequential clustering**
- > Examine four-vector inputs pairwise and construct jets hierarchically
- > **anti- k_T** : preferentially merge constituents with **high p_T** with respect to their nearest neighbours first
 - undoing merging yields one high- and one low- p_T subjet
- > **Cambridge-Aachen**: no p_T -weighting, merge based on **spatial separation** only → undoing clustering yields more **p_T -symmetric subjets**



> We know the **masses** of W, Z and Higgs very well → can use them as **constraints**

> However, large number of particles in jet → rather bad resolution

> Jet grooming removes **soft and large angle radiation**

> Strategy:

- recluster jet using Cambridge-Aachen (CA) jet algorithm
- iteratively break into two subjets
- remove softer contribution (and continue with harder one) if:

soft-drop condition

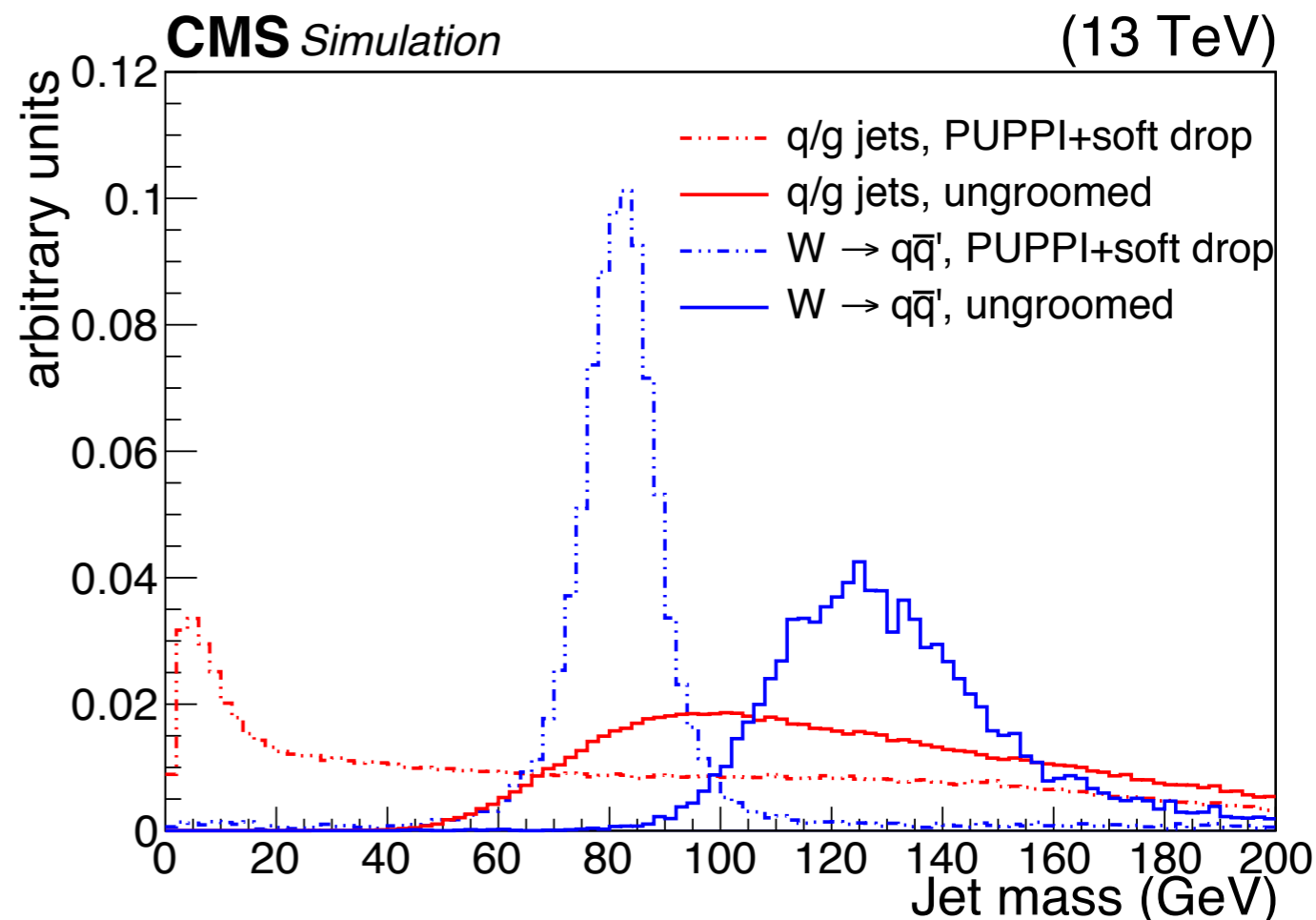
$$\frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}} < 0.1$$

modified mass-drop algorithm (mMDT)

- stop otherwise

> Cut on mass window ($\sim \pm 10$ GeV)

before grooming, pileup removal is performed (PUPPI)



use of soft-drop algorithm new w.r.t. 2015 — previously used pruning → perturbatively robust

N-subjettiness

> For **boson-tagging**: want to quantify how **2-subjetty** a jet is

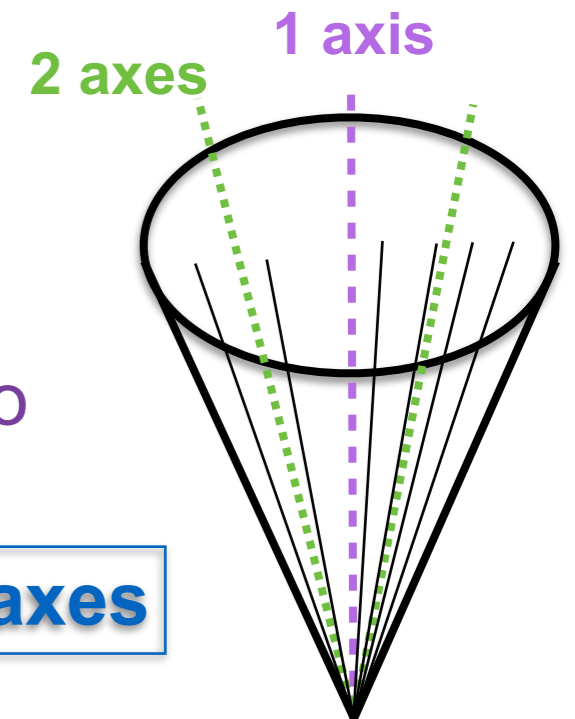
> → To what extent is energy flow aligned along 2 momentum directions (N=2)?

$$\tau_N = \frac{1}{d_0} \sum_k p_{T,k} \min(\Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k})$$

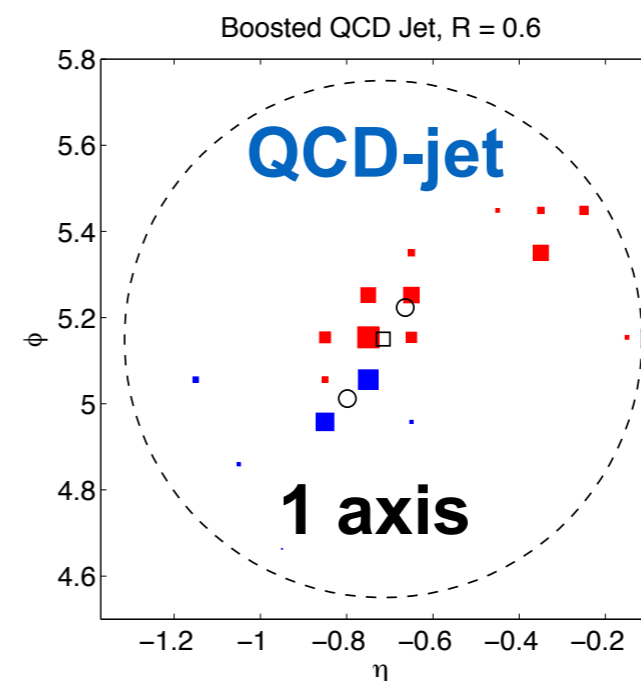
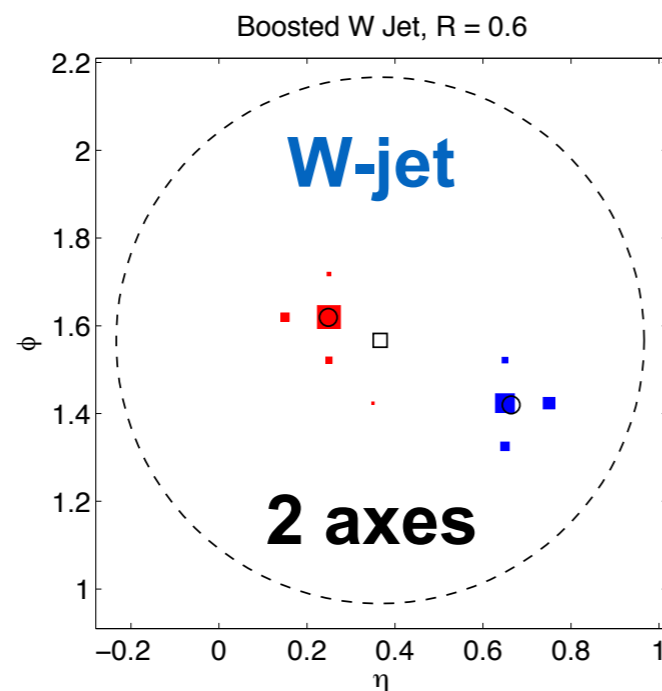
normalisation

sum over particles

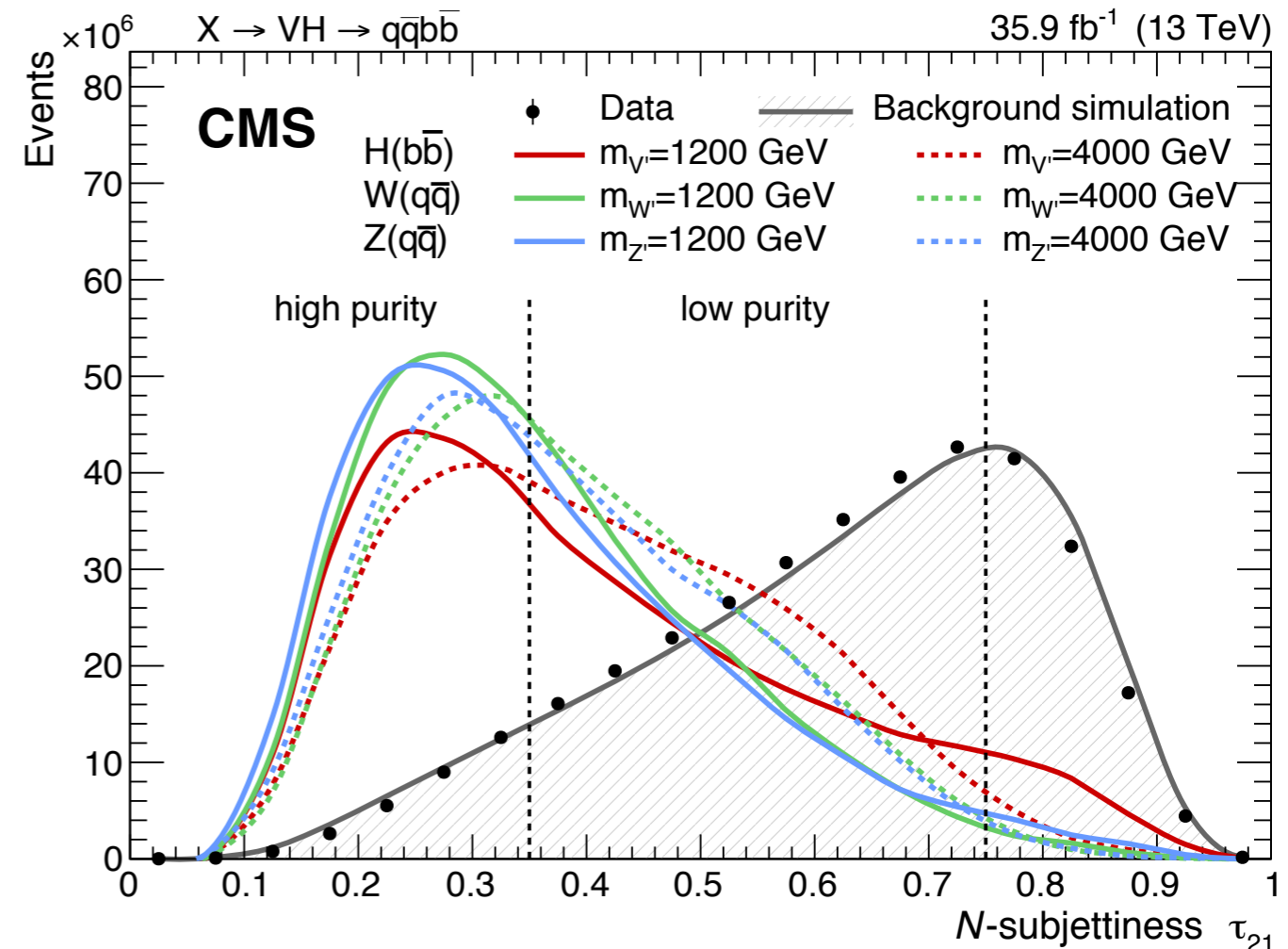
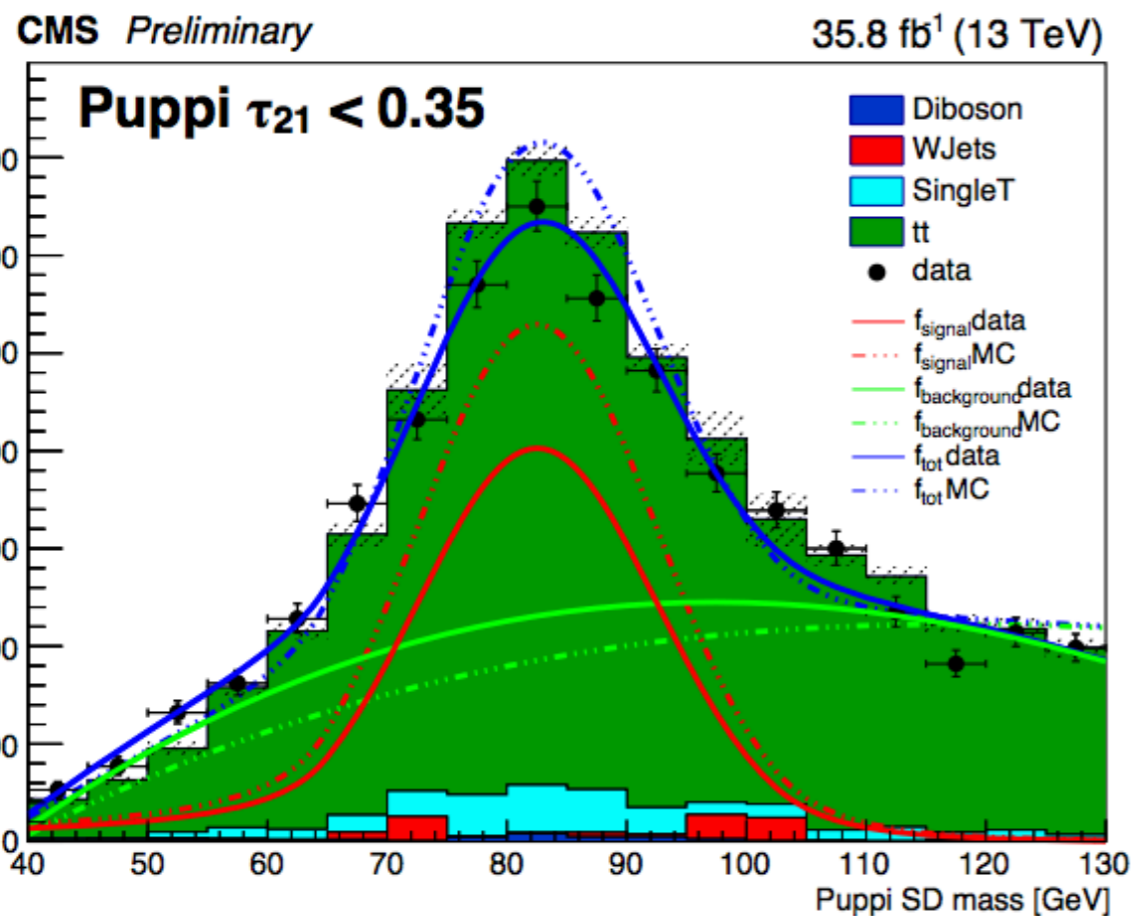
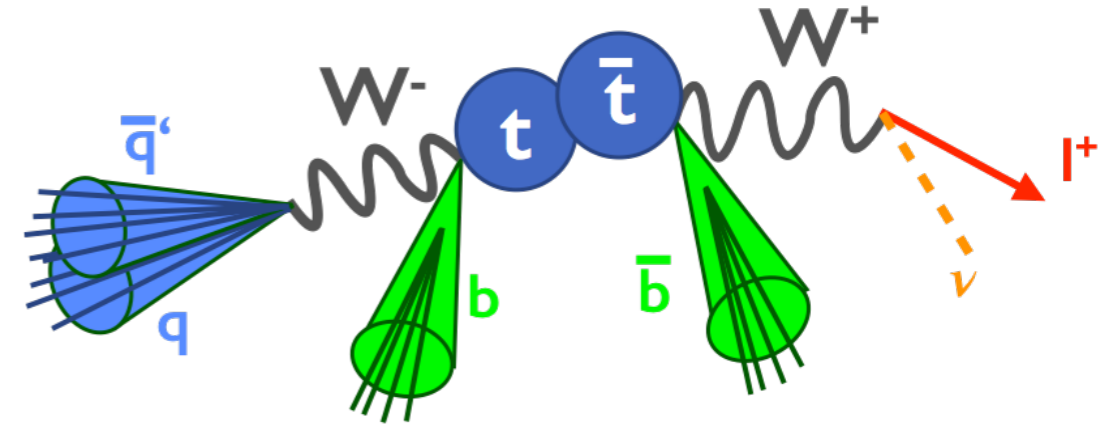
minimise distance to candidate subsets



low values of τ_N → compatibility with the hypothesis of N axes

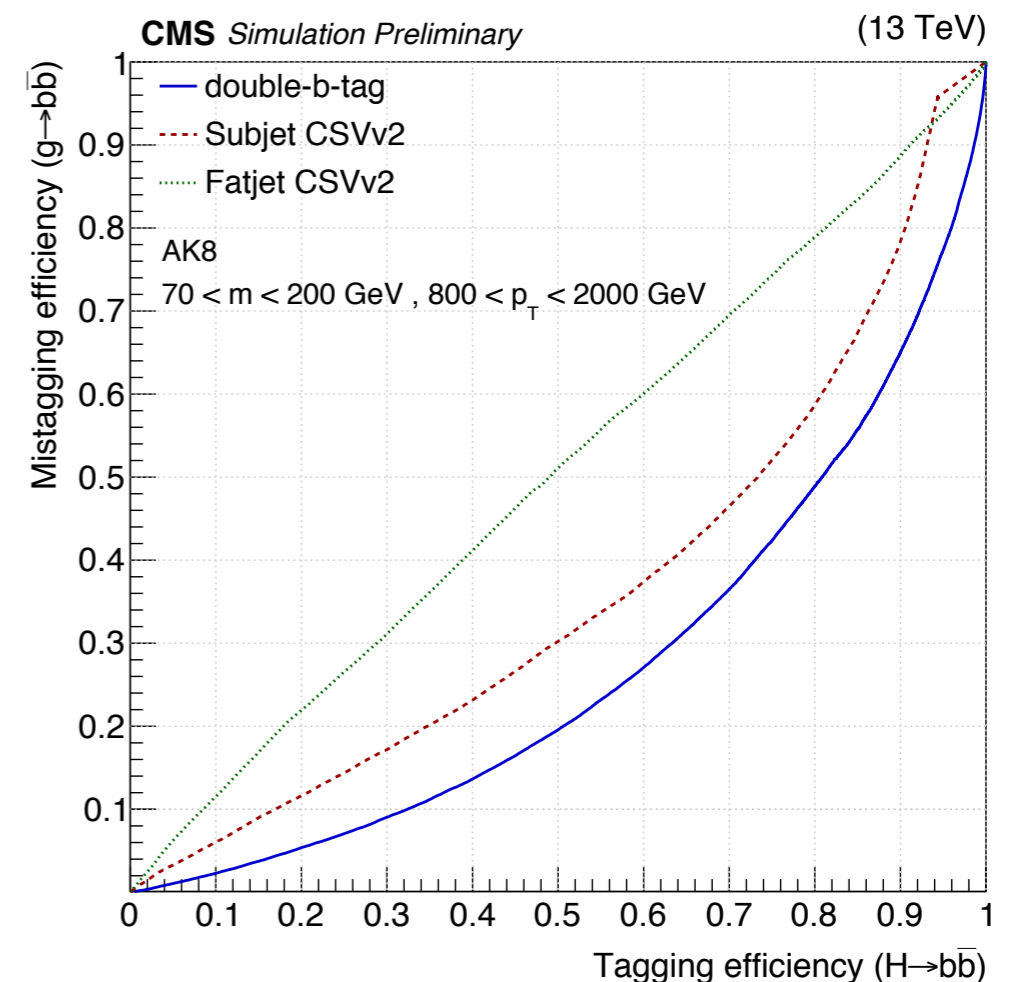
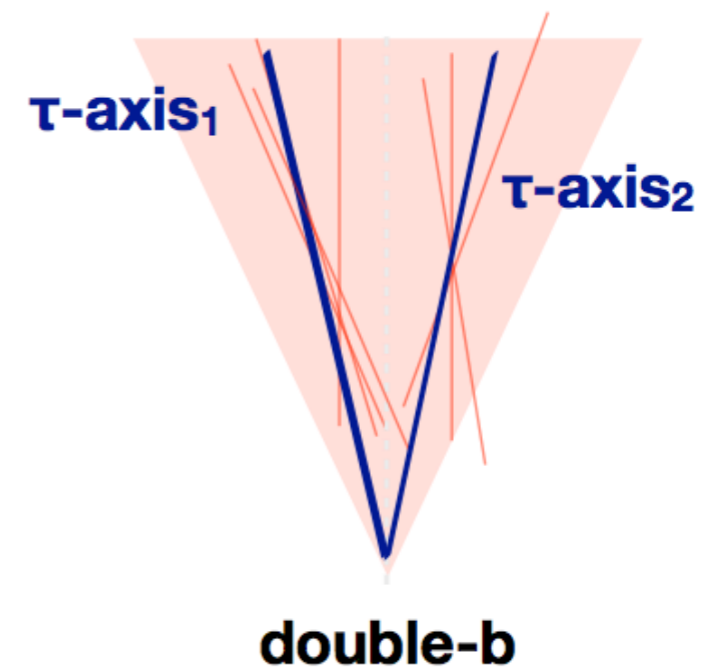


- > Bare τ_N has very little discrimination power
- > Take **ratio** τ_2/τ_1 instead
- > Note: rather **complicated variable**, difficult to model \rightarrow need to validate in data
- > Clean sample of W-jets: **top-antitop quark pairs** used for calibration (W-jet $p_T \sim 200$ GeV)



Higgs \rightarrow bb tagging

- > Higgs has higher mass than W/Z bosons \rightarrow τ_2/τ_1 less important, **exploit b-jet content** instead
- > Previously, two different strategies:
 - identify b-subjets
 - tag fat jet
- > Already **50% lower mis-tagging rate** than W-/Z-tagging
- > Run-2: dedicated Higgs-tagger — **double-b tagger** (MVA-based):
 - inputs based on observables from secondary vertex and tracks associated to each τ -axis (27 total)
 - significantly better background rejection
- > For resolved Higgs decays: new DeepCSV discriminator



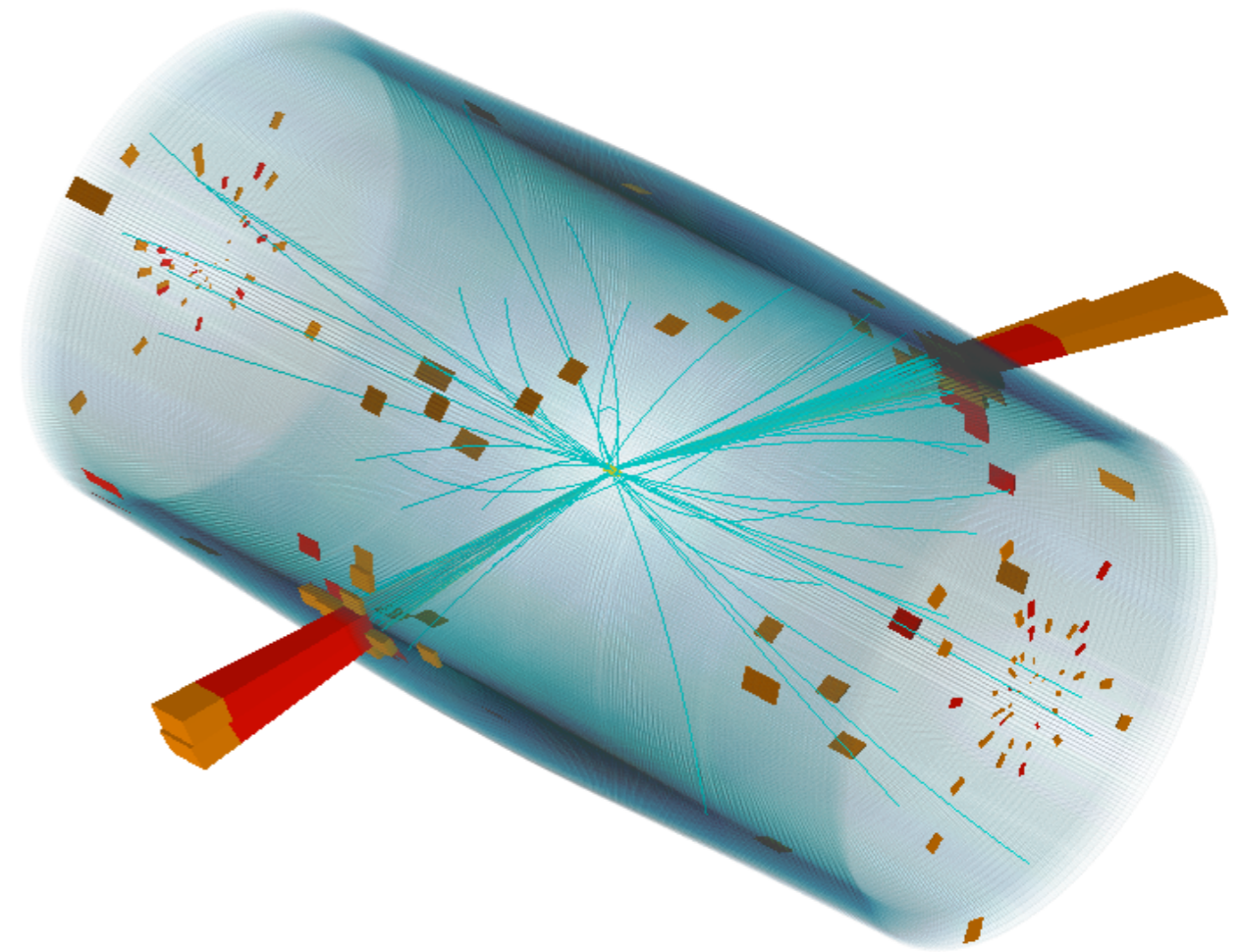


All-hadronic VV/Vh final states

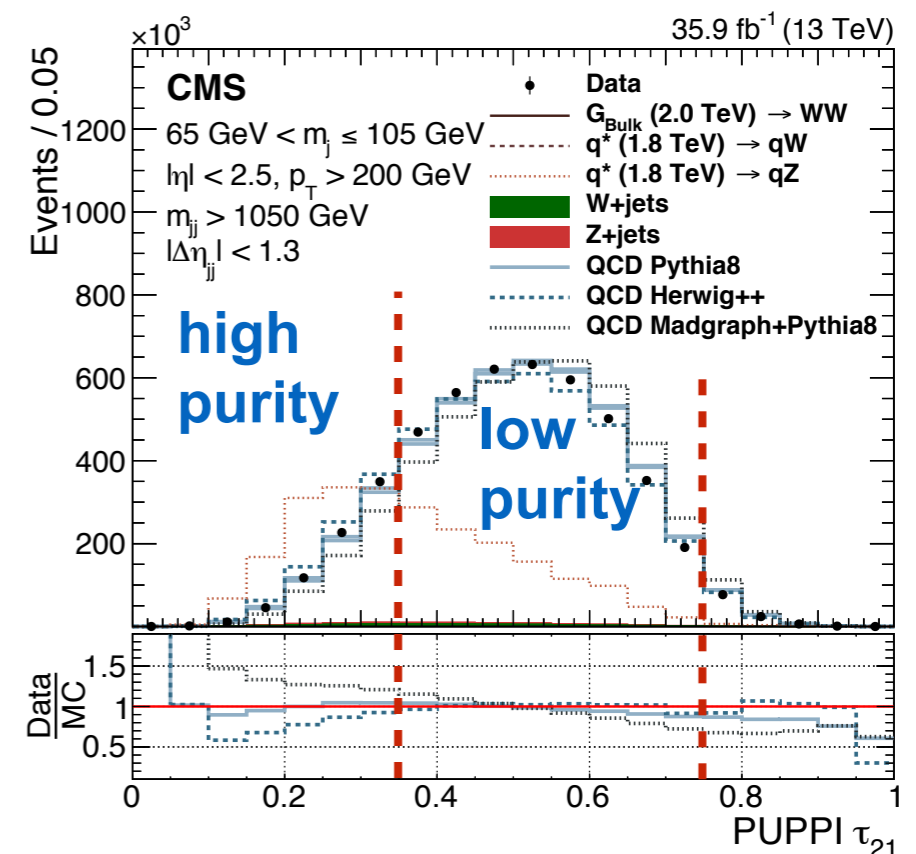
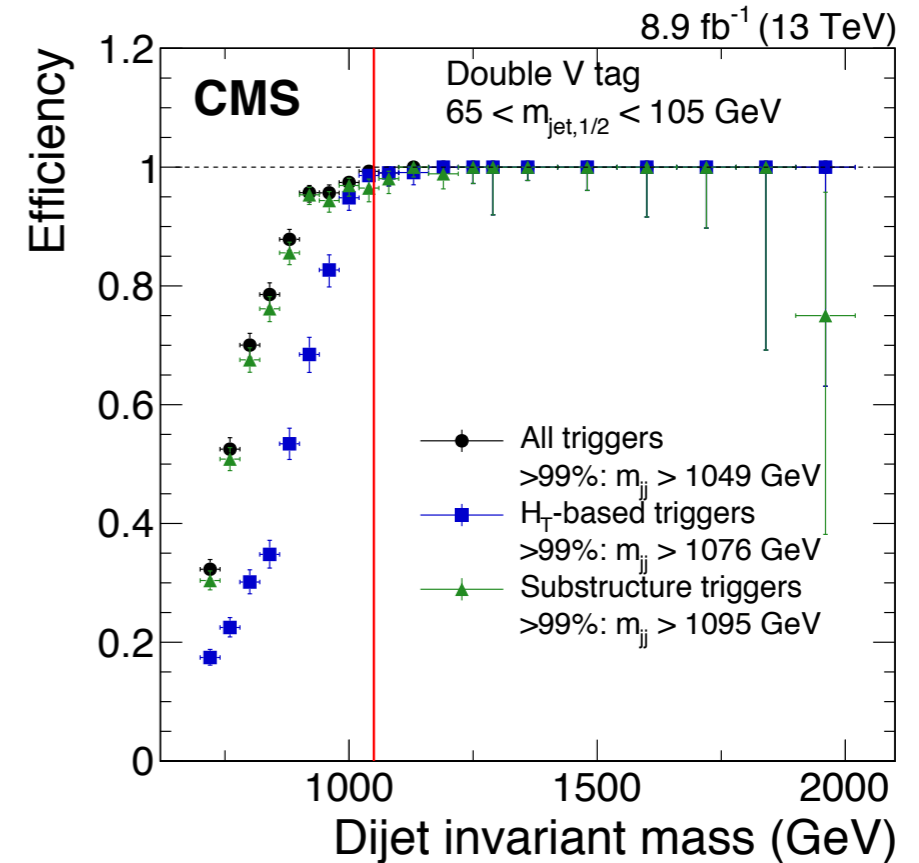


- Boson branching fractions to quarks very large:
 - $W/Z \rightarrow$ quarks $\sim 69\%$
 - Higgs $\rightarrow bb$ $\sim 57\%$
- High acceptance
 - large reach in resonance mass
- For heavy resonances: dijet final state
- Challenge: background estimation

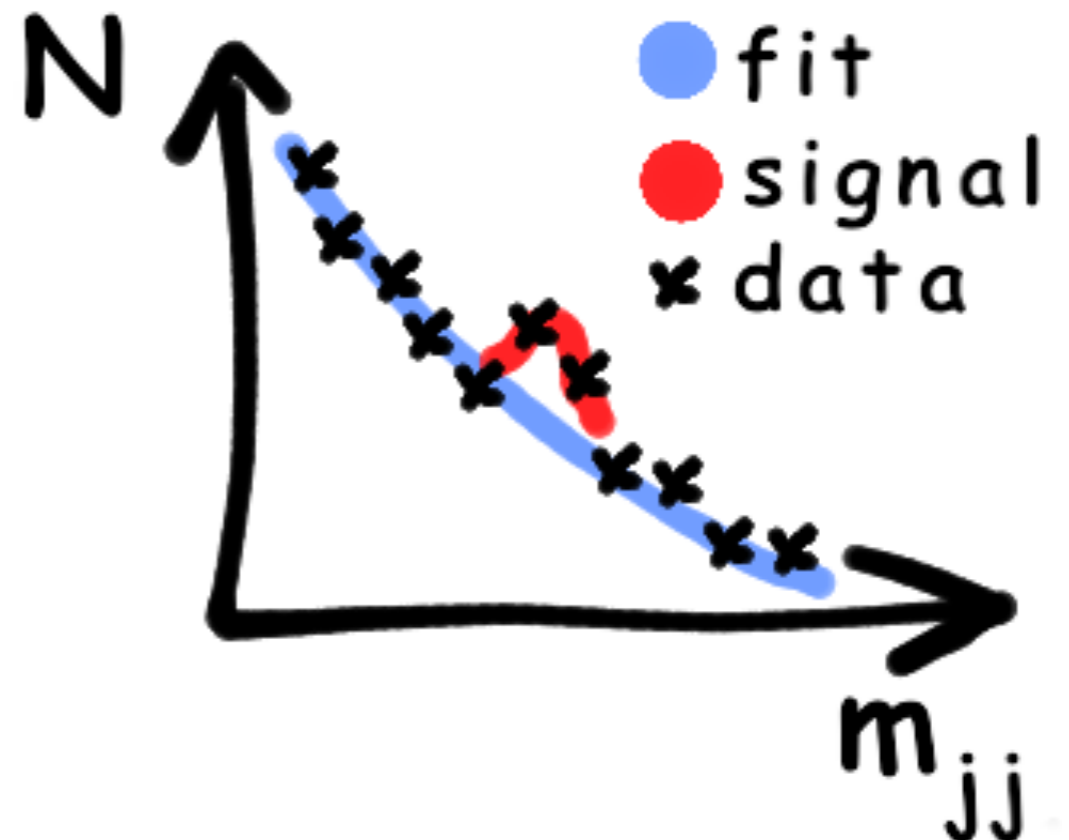
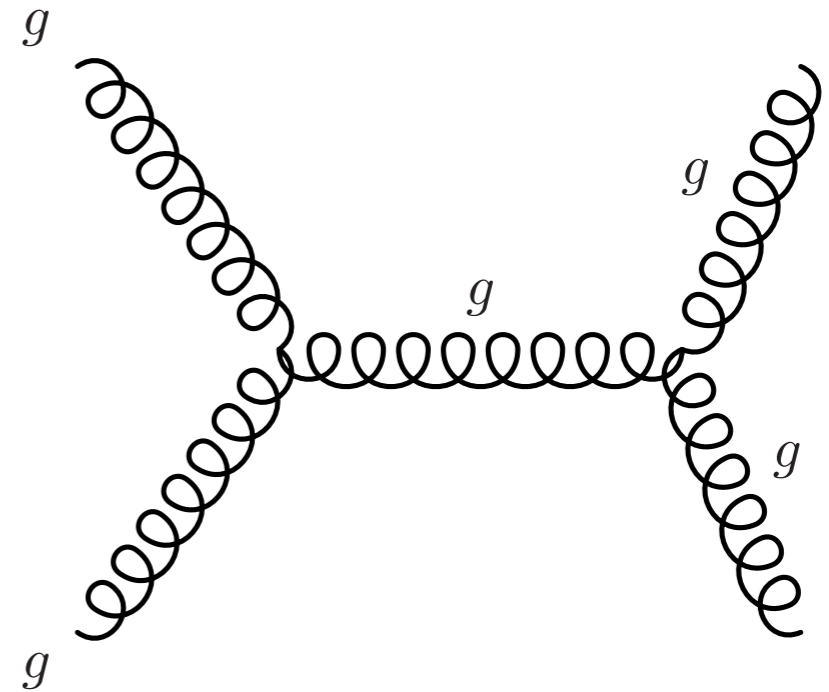
Candidate ZZ event
Dijet mass: 3.2 TeV



- > Dijet final state: WW, WZ, ZZ (=VV)
- > High trigger thresholds
 - H_T -based (scalar p_T sum of all objects in the event)
 - high- p_T jet with substructure
- > **Trigger** at $\sim 100\%$ efficiency for $m_{JJ} > 1.05 \text{ TeV}$
 - apply cut on reconstructed dijet system
- > Define different τ_2/τ_1 regions:
 - high purity to suppress background
 - low purity to recover signal efficiency at high masses
- > **Split W and Z** samples based on soft-drop jet mass (65-85, 85-105 GeV)



- > After selection, still dominated by **QCD multijet** events
- > Difficult to obtain sufficient number of simulated events
- > Need a data-driven approach
- > Exponentially falling spectrum
 - since we are in the trigger efficiency plateau
- > Can use **fit function** and perform a so-called **bump hunt**



- > Naively, fitting the m_{JJ} spectrum could swallow signal
- > Also, need to **avoid claiming false discovery** (in particular in tail)

> Fit function:

$$\frac{dN}{dm_{jj}} = \frac{P_0}{(m_{jj}/\sqrt{s})^{P_2}} \text{ or } \frac{dN}{dm_{jj}} = \frac{P_0(1 - m_{jj}/\sqrt{s})^{P_1}}{(m_{jj}/\sqrt{s})^{P_2}}$$

2 parameters

3 parameters

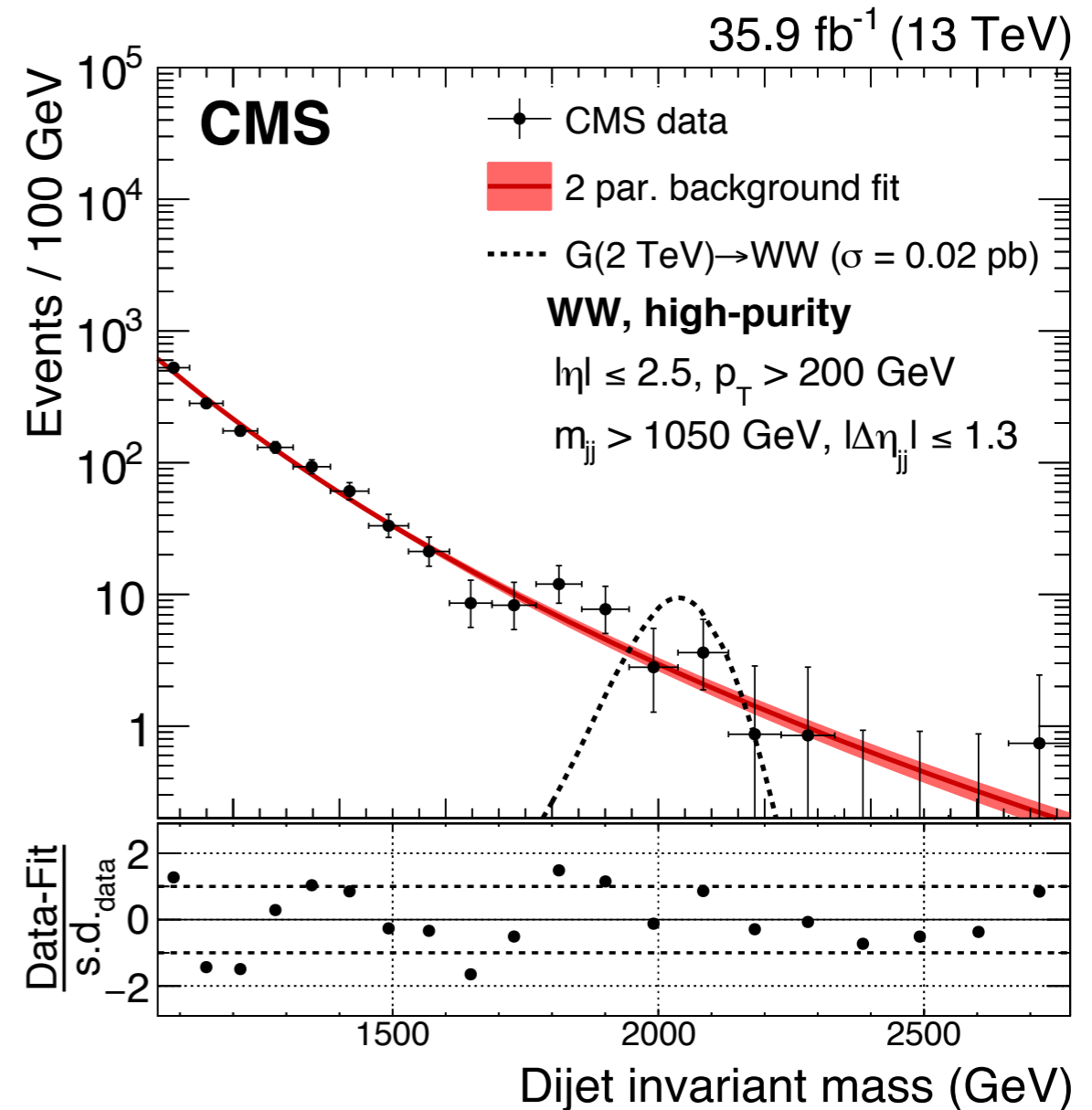
... N parameters

> Number of free parameters determined by **F-test**:

- check if quality of fit improves by > 10% confidence level
- if not, stick with current fit function

> **Extensive bias tests** conducted

> Combined **signal+background fit** performed



spin-1 and -2 interpretation

see also Vh analysis (B2G-17-002)

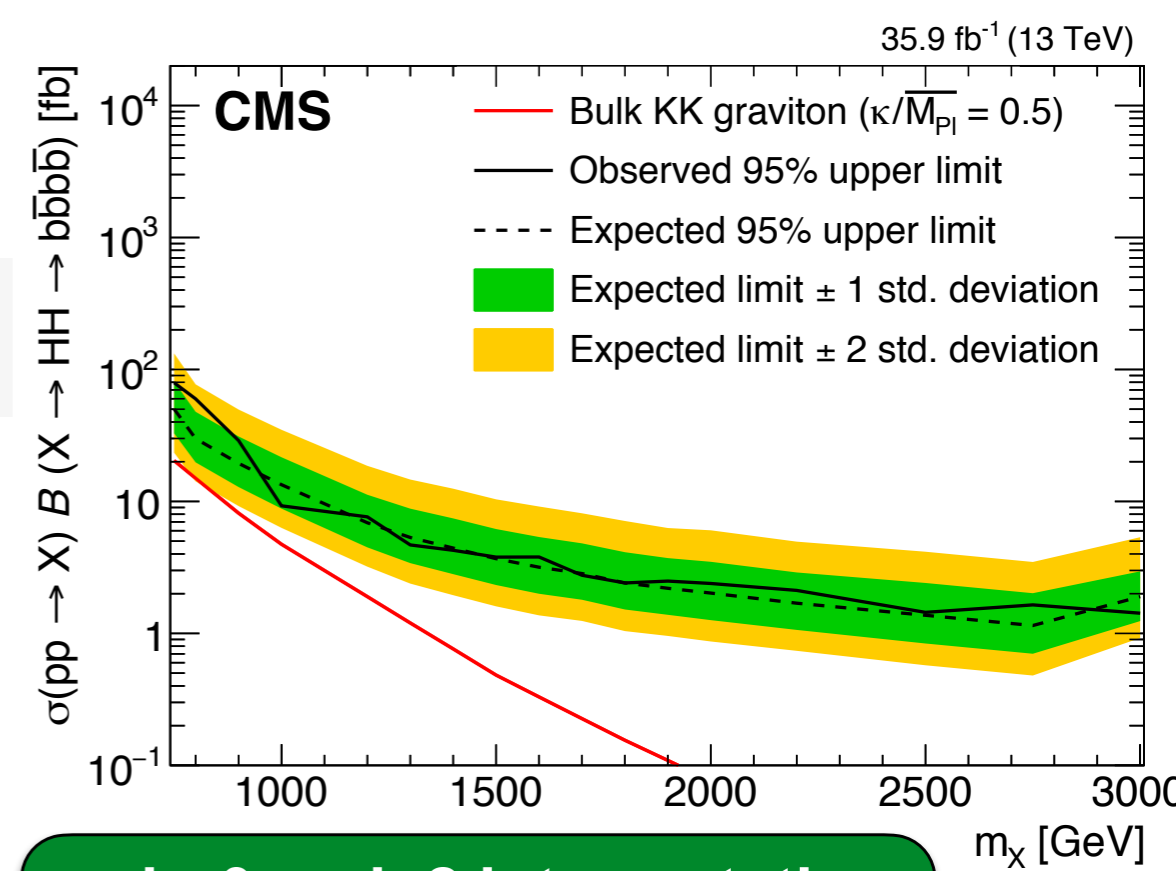
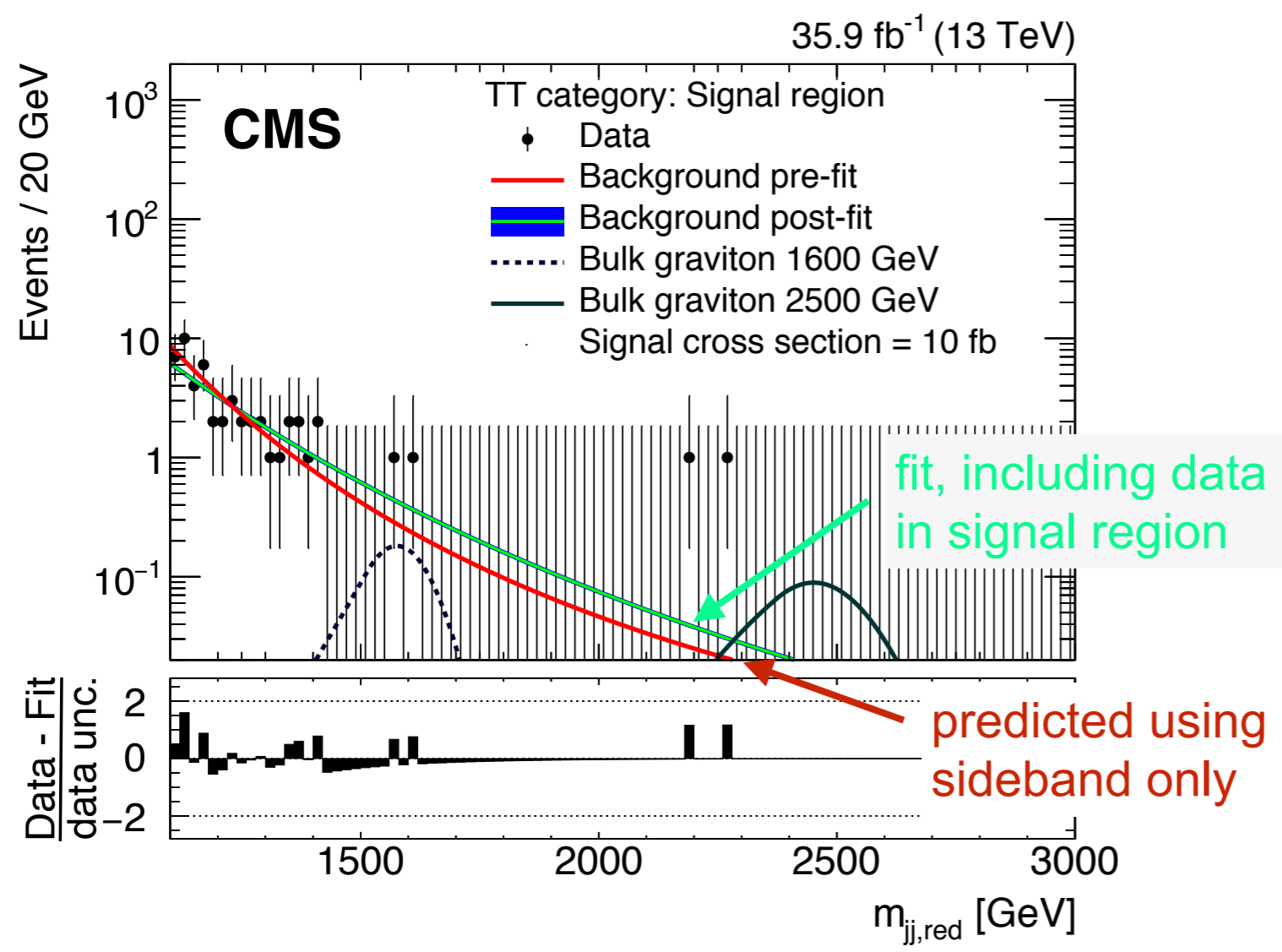
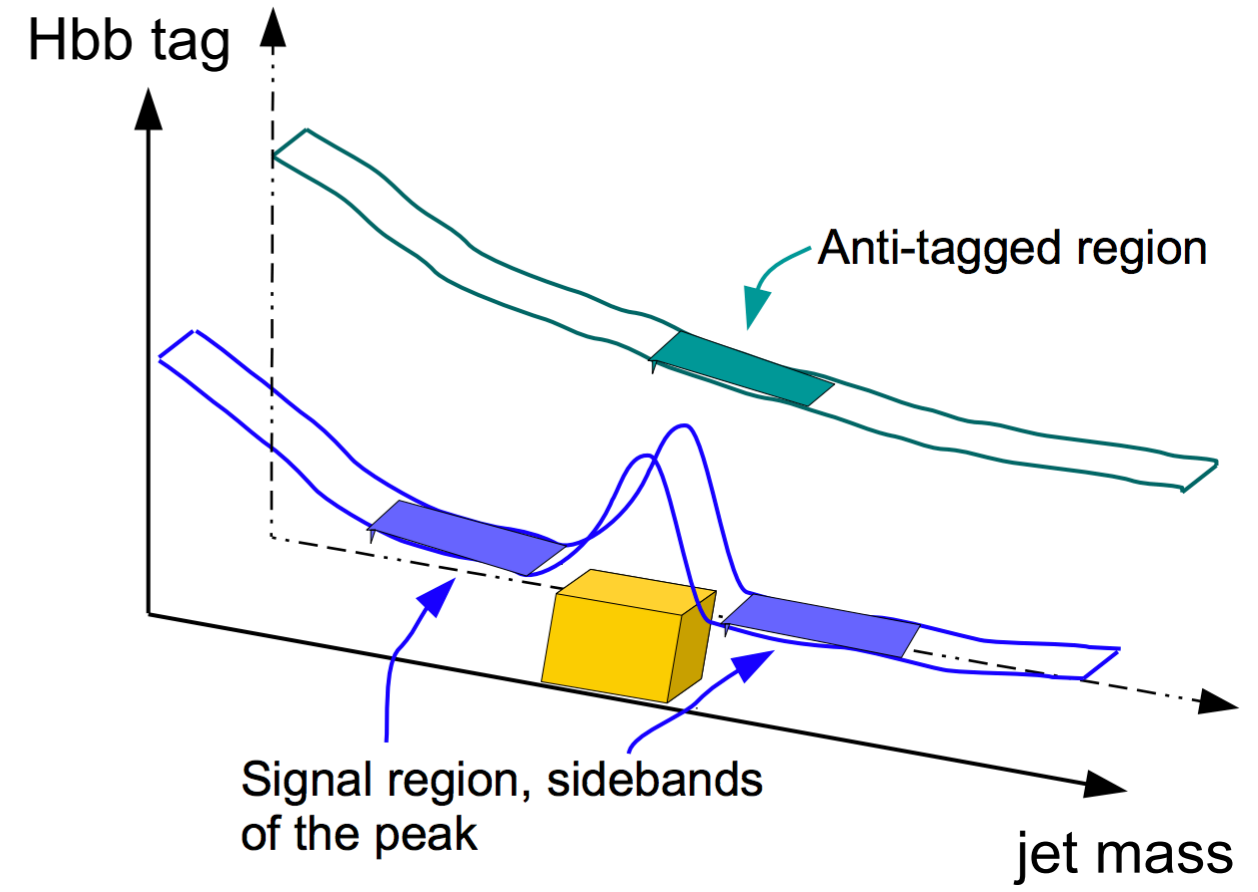


Di-Higgs (hh) analyses



hh analysis — high mass

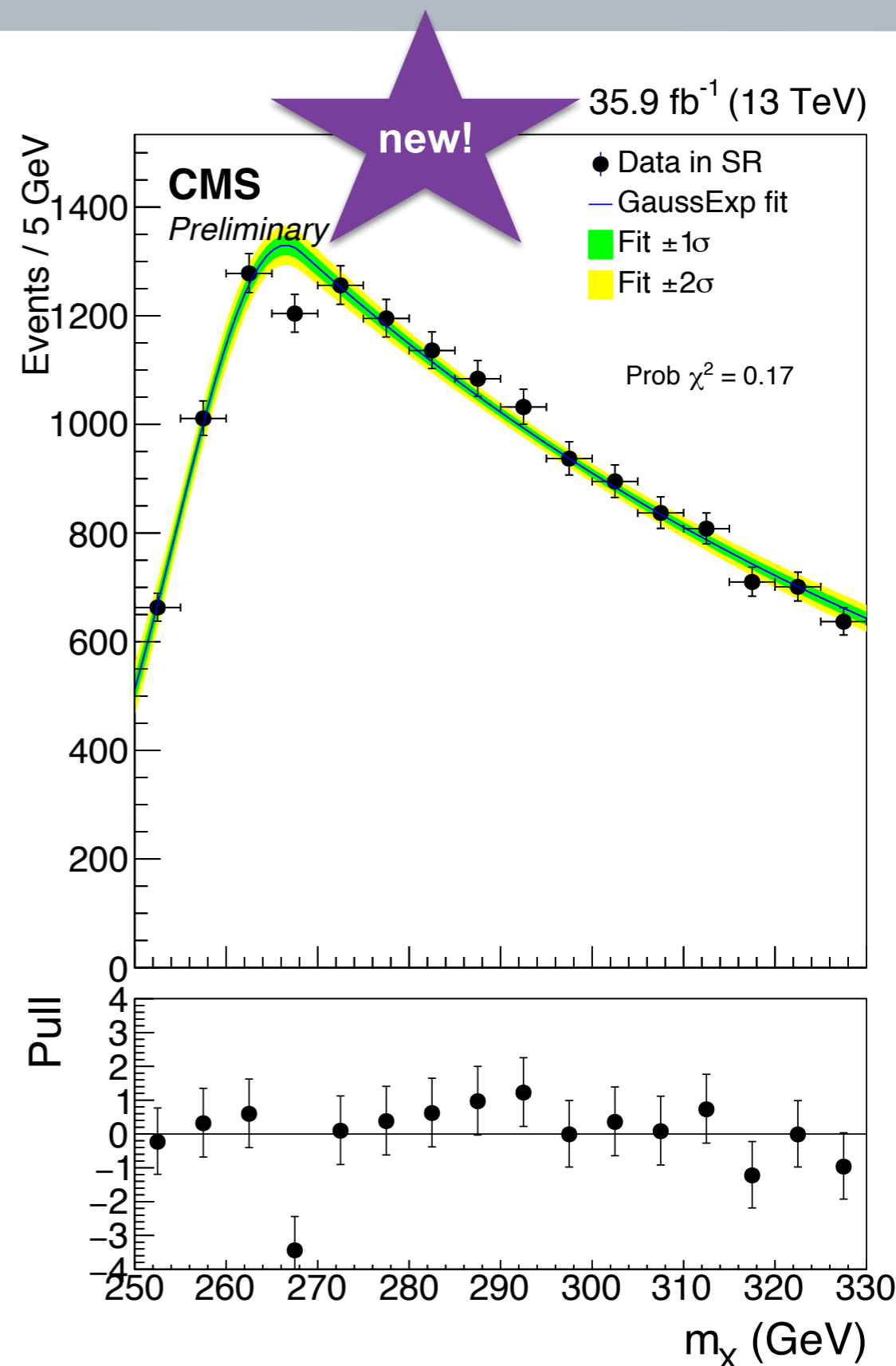
- > Dijet final state: $hh \rightarrow 2 \text{ bb-jets}$
- > New background estimation technique: Alphabet-assisted bump hunt
 - define anti-Hbb tagged region
 - parametrise anti-tagged to tagged ratio to constrain background from sidebands
 - perform dijet fit



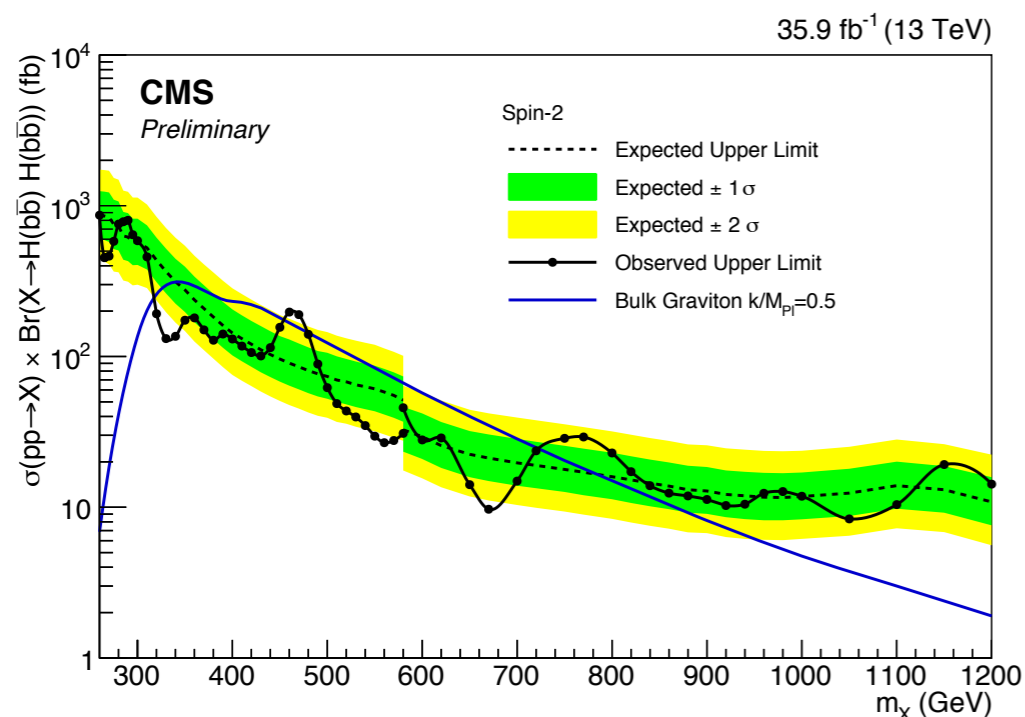
spin-0 and -2 interpretation

hh→4b analysis — low/medium mass

- Can go lower in mass
 - go from boosted dijet final state to (semi-) resolved final state
- Select 4 anti- k_T ($R=0.4$) jets
 - low mass region: 260-620 GeV
 - medium mass region: 550-1200 GeV
 - profits from DeepCSV b-tagging algorithm
- Also uses parametric background estimation



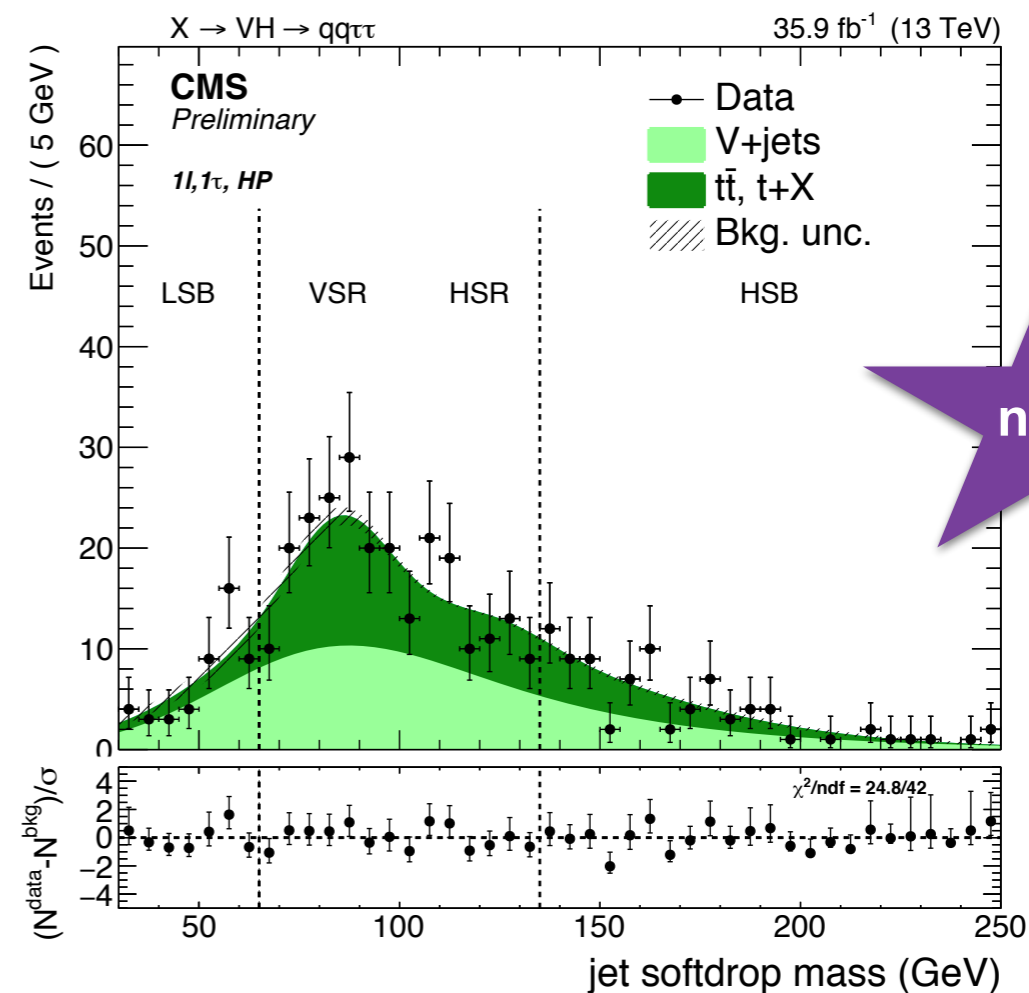
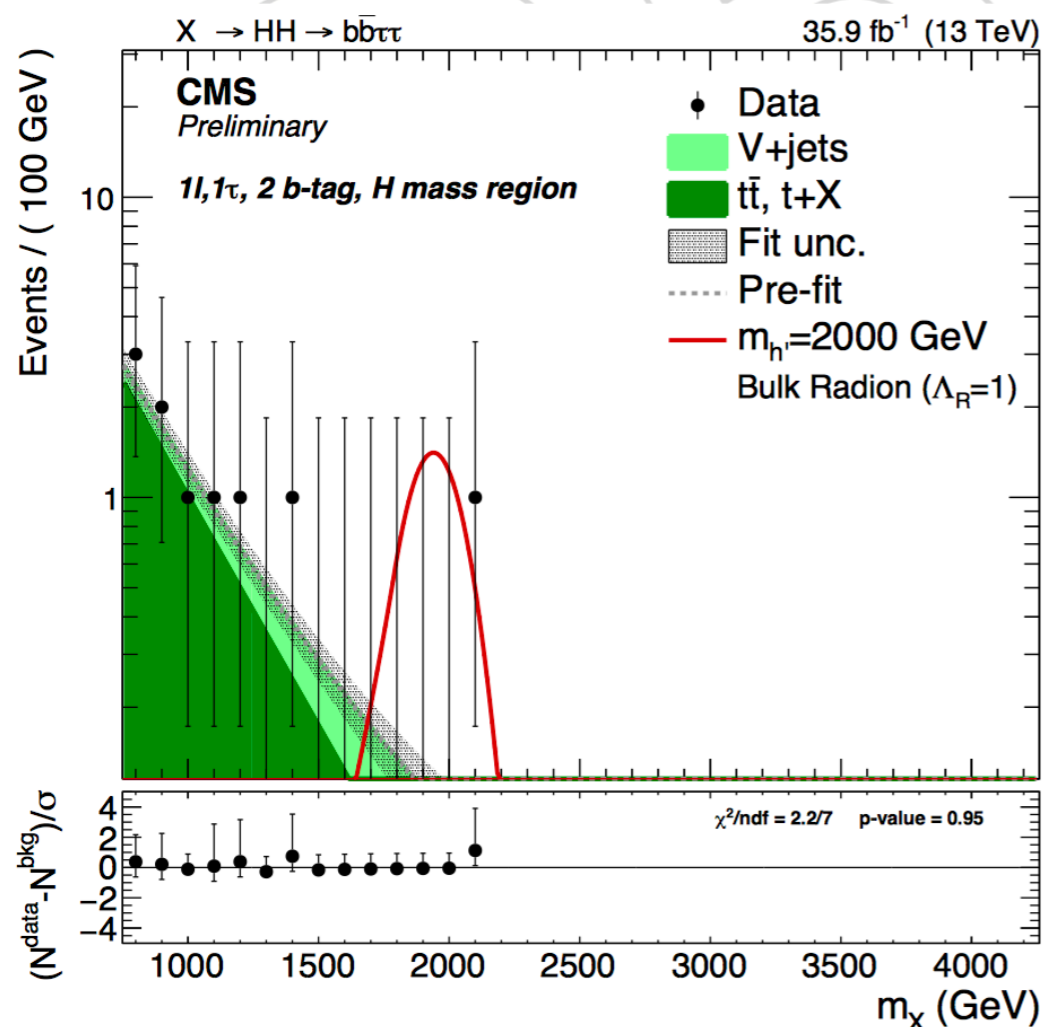
spin-0 and -2 interpretation



More on di-Higgs production

> Large number of di-Higgs analyses in CMS

- resonant and non-resonant (self-coupling, probing the electroweak symmetry breaking potential) interpretation
- $hh \rightarrow 4b$ analyses most sensitive at high invariant hh mass



> Another example: $hh \rightarrow b\bar{b}\tau\tau$ high mass analysis

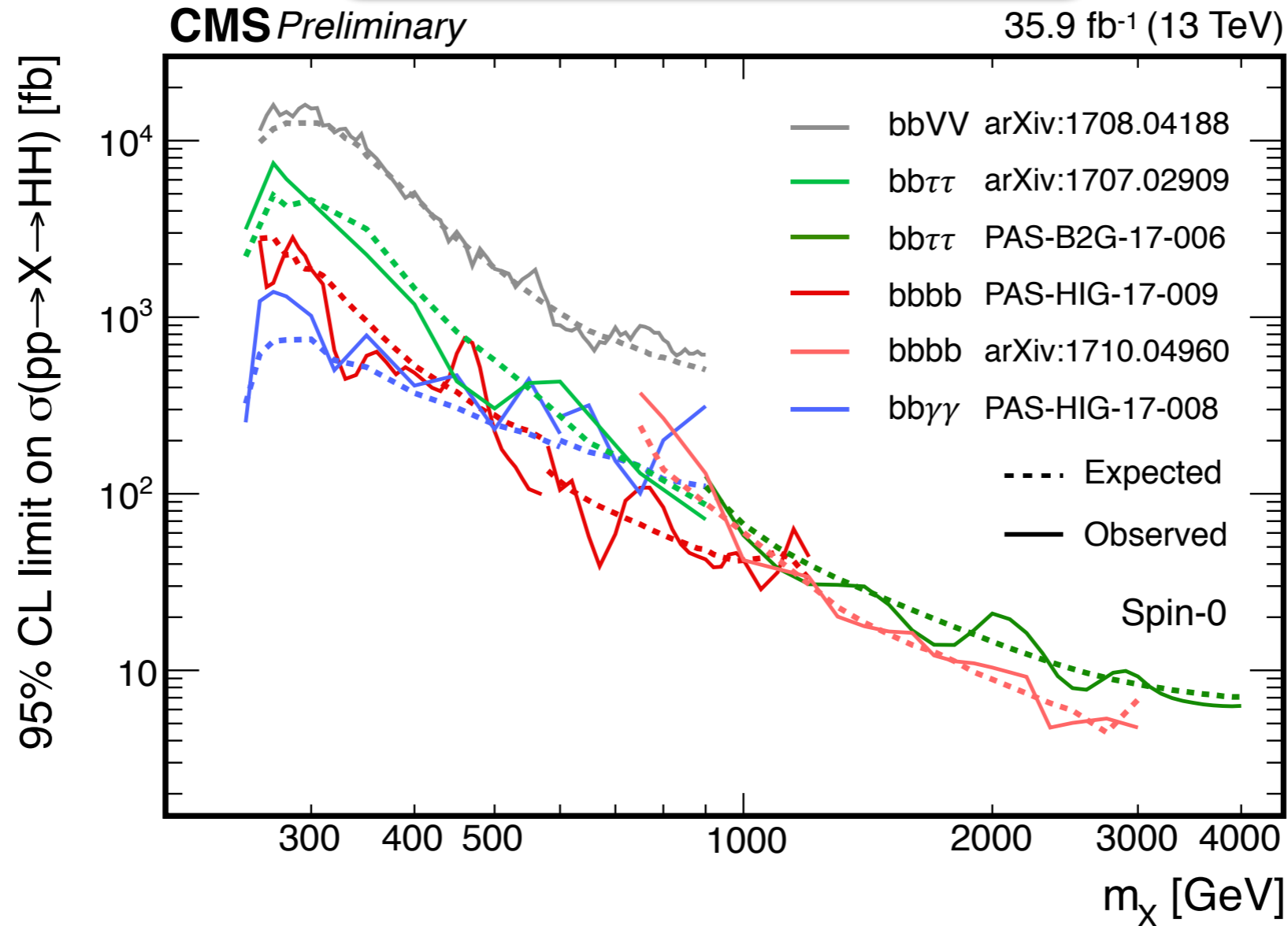
- dedicated high- p_T $h \rightarrow \tau\tau$ reconstruction
- spin-0, -1 ($Vh \rightarrow qq\tau\tau$), and -2 interpretation

> Also have analyses covering $hh \rightarrow b\bar{b}\gamma\gamma$ (HIG-17-008), $b\bar{b}l\nu l\nu$ (HIG-17-006) and low mass $b\bar{b}\tau\tau$ (HIG-16-002)

spin-0 and -2 interpretation

More on di-Higgs production

spin-0 interpretation



bbllvv (HIG-17-006)

bbbb (HIG-17-009)

bb $\tau\tau$ (HIG-17-002, submitted to Phys. Lett. B)

bbbb (B2G-16-026, submitted to Phys. Lett. B)

bb $\tau\tau$ (B2G-17-006, submitted to J. High Energy Phys.)

bby γ (HIG-17-008)

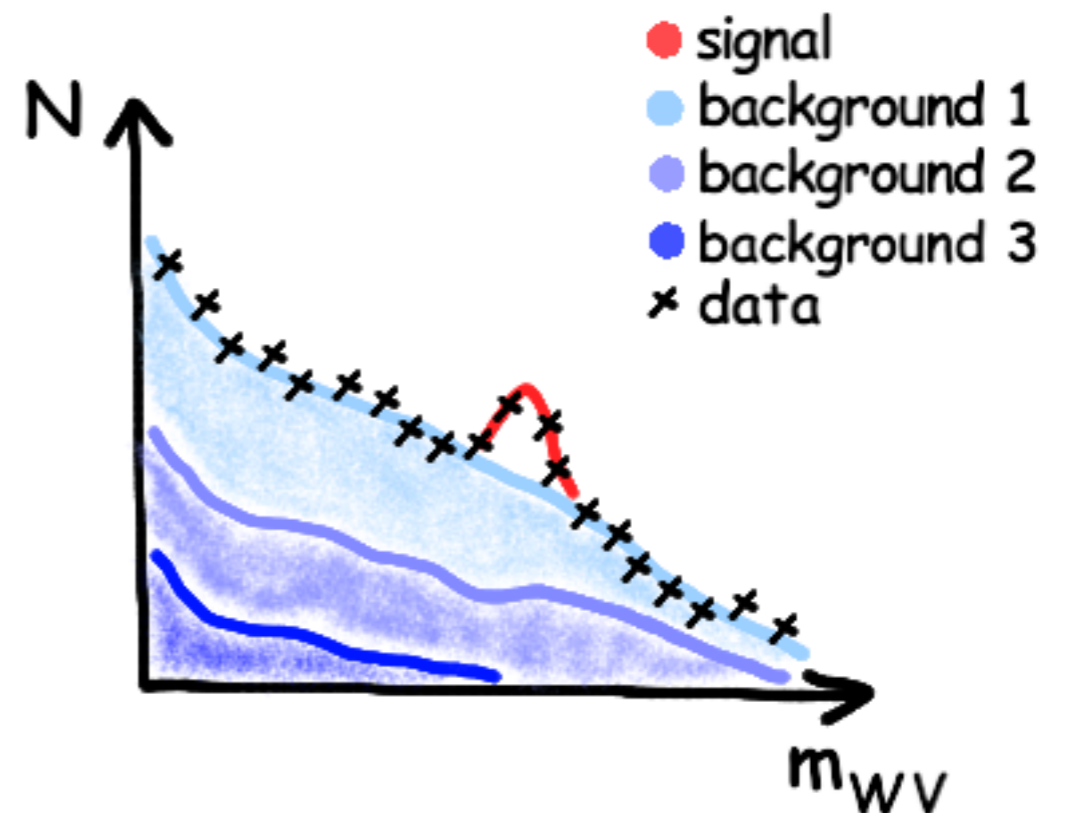
spin-2 interpretation in backup



(Semi-)leptonic VV final states



- Requiring one or more lepton(s) in the final state reduces background significantly
 - covering large number of final states ($l\nu 2q$, $4l$, $2l2\nu$, $2q2l$, $2q2\nu$) and large mass range
- Lower trigger thresholds
- Simulation-assisted background estimation performed
 - distinguish different background processes
 - ratios for sideband to signal region extrapolation
 - or directly using simulation
- All analyses combined cover mass range from 130 to 4500 GeV

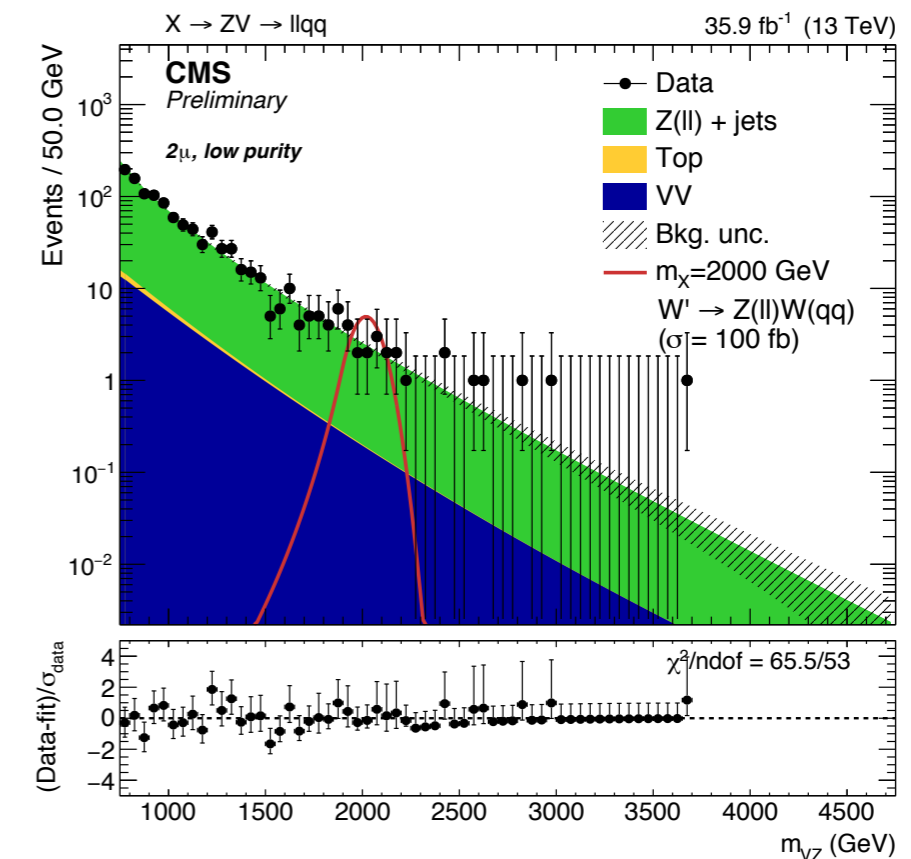
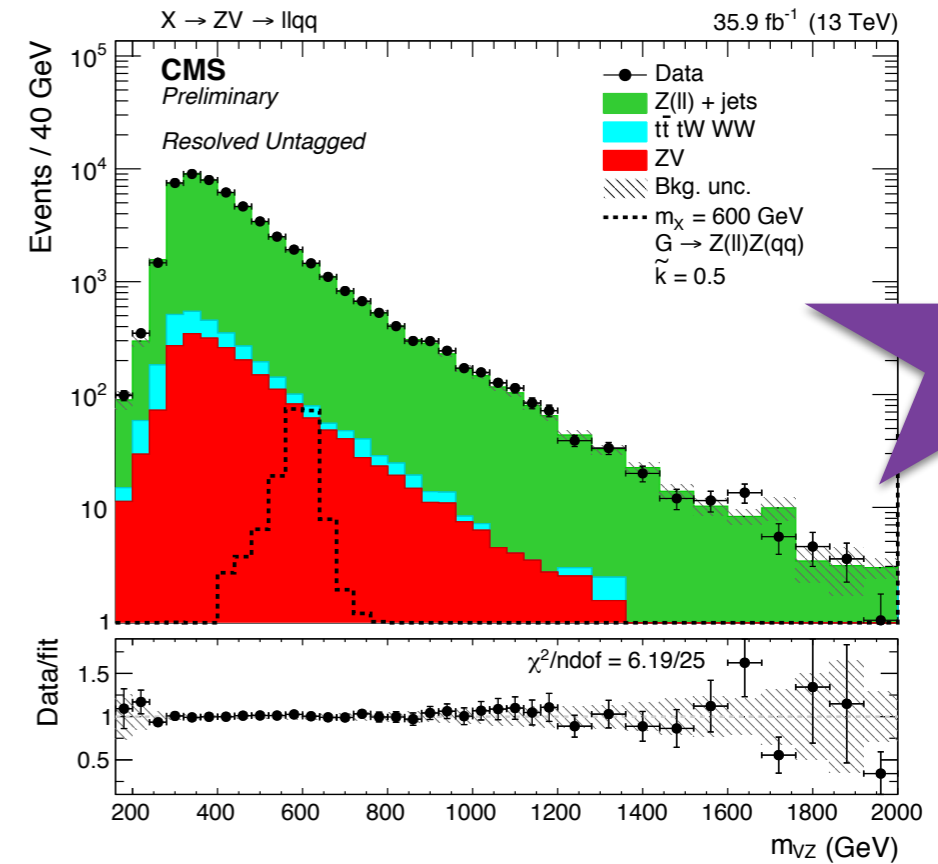


> Several different analyses:

- high mass spin-2: $2\ell 2\nu$ (B2G-16-023)
- high mass spin-1/-2: $2q 2\nu$ (B2G-17-005)
- intermediate to high mass spin-1/-2: $2\ell 2q$ (B2G-17-013)
- low to high mass spin-0: 4ℓ , $2\ell 2\nu$, $2\ell 2q$ (HIG-17-012)

> Recently released: B2G-17-013

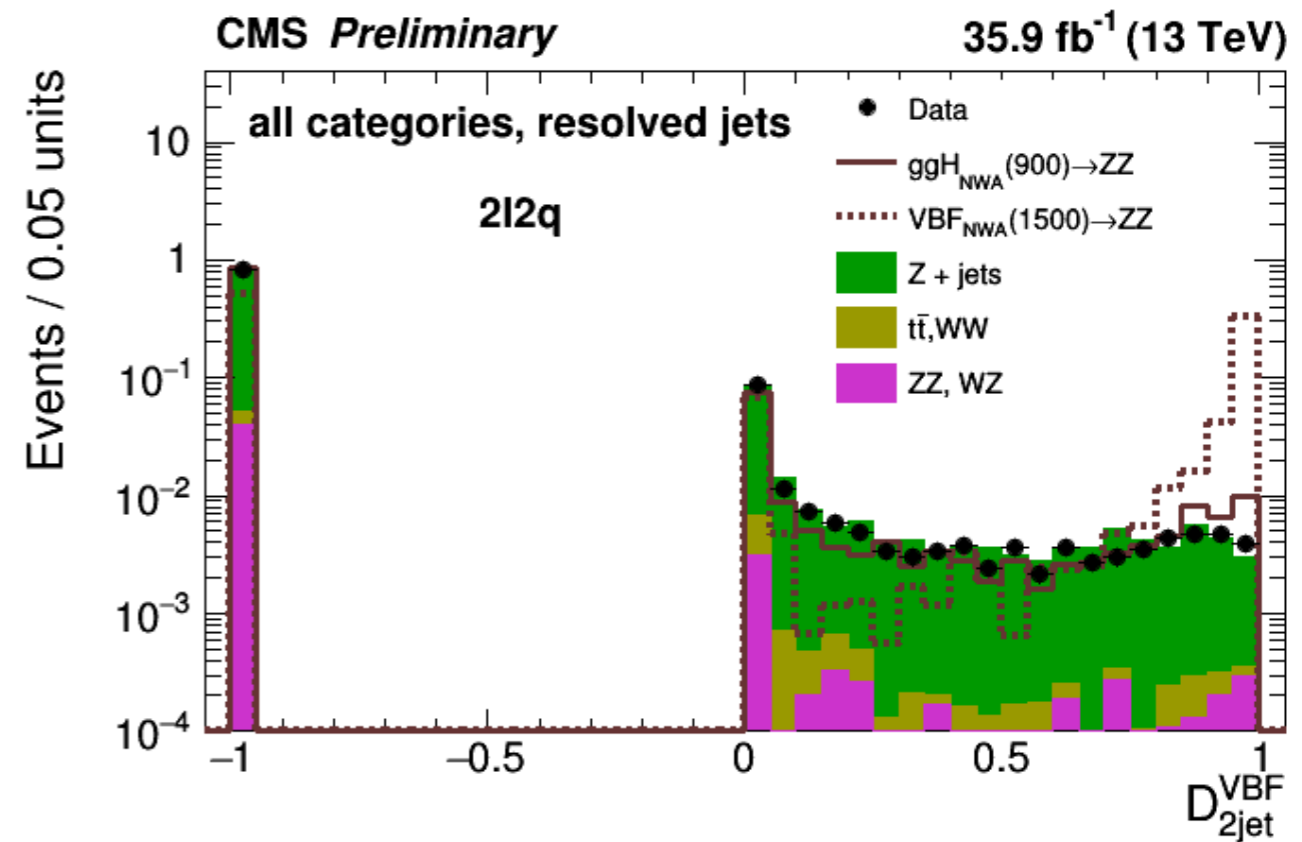
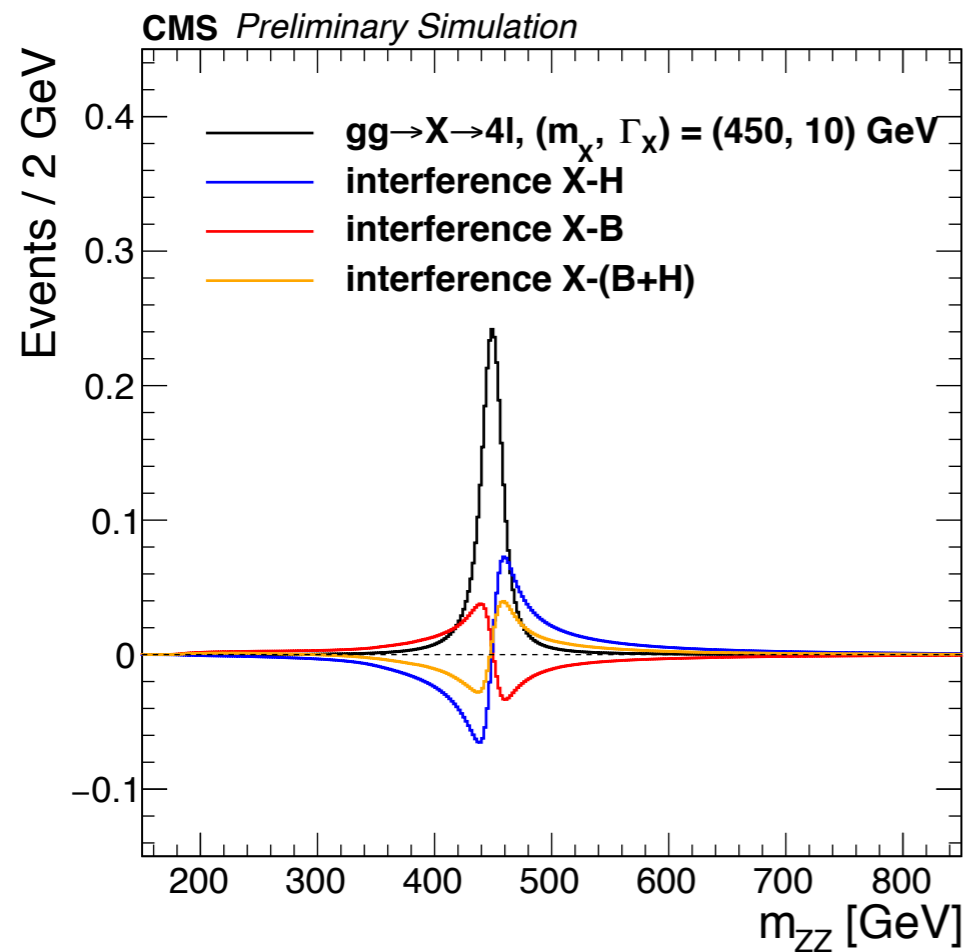
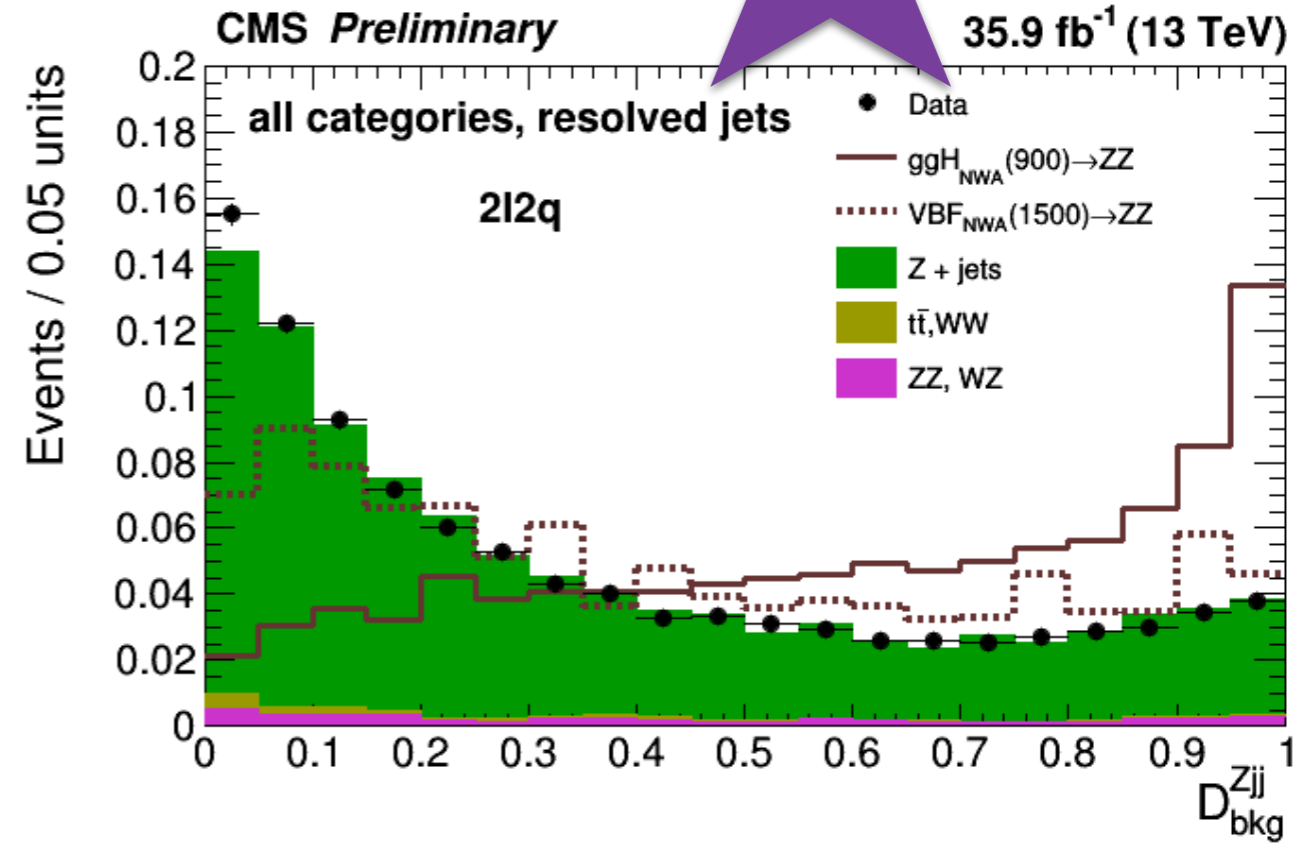
- $2\ell 2q$ final state
- dedicated low and high mass analyses
- analysis extends from 400 GeV up to 4.5 TeV



ZV/ZZ ($Z \rightarrow \nu\nu/\ell\ell$) analyses

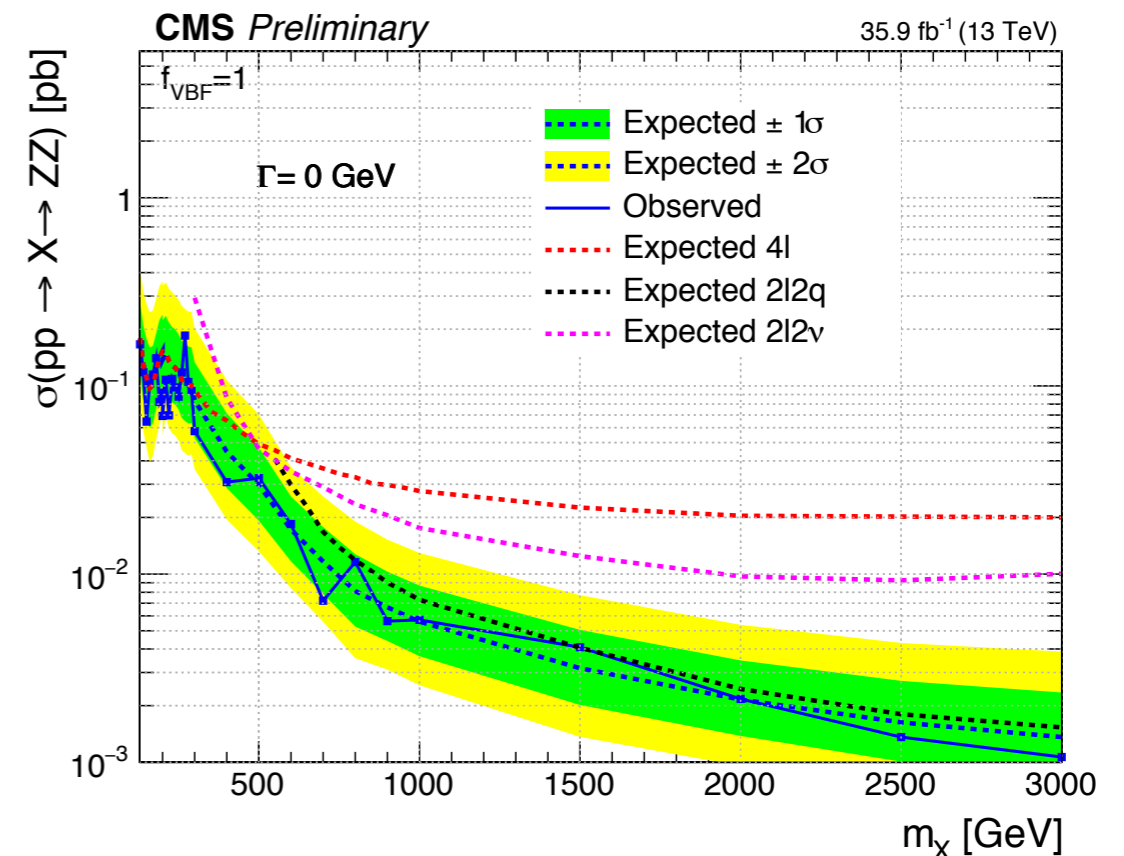
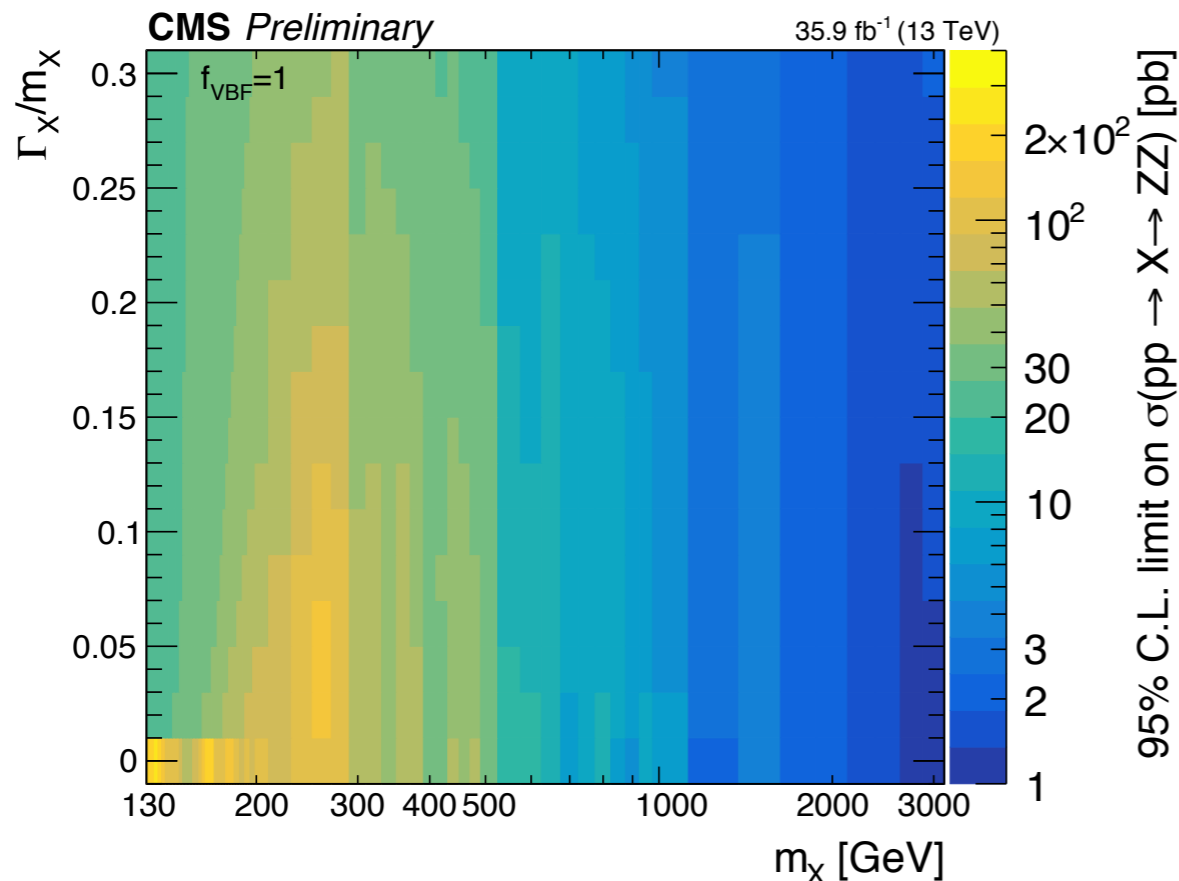
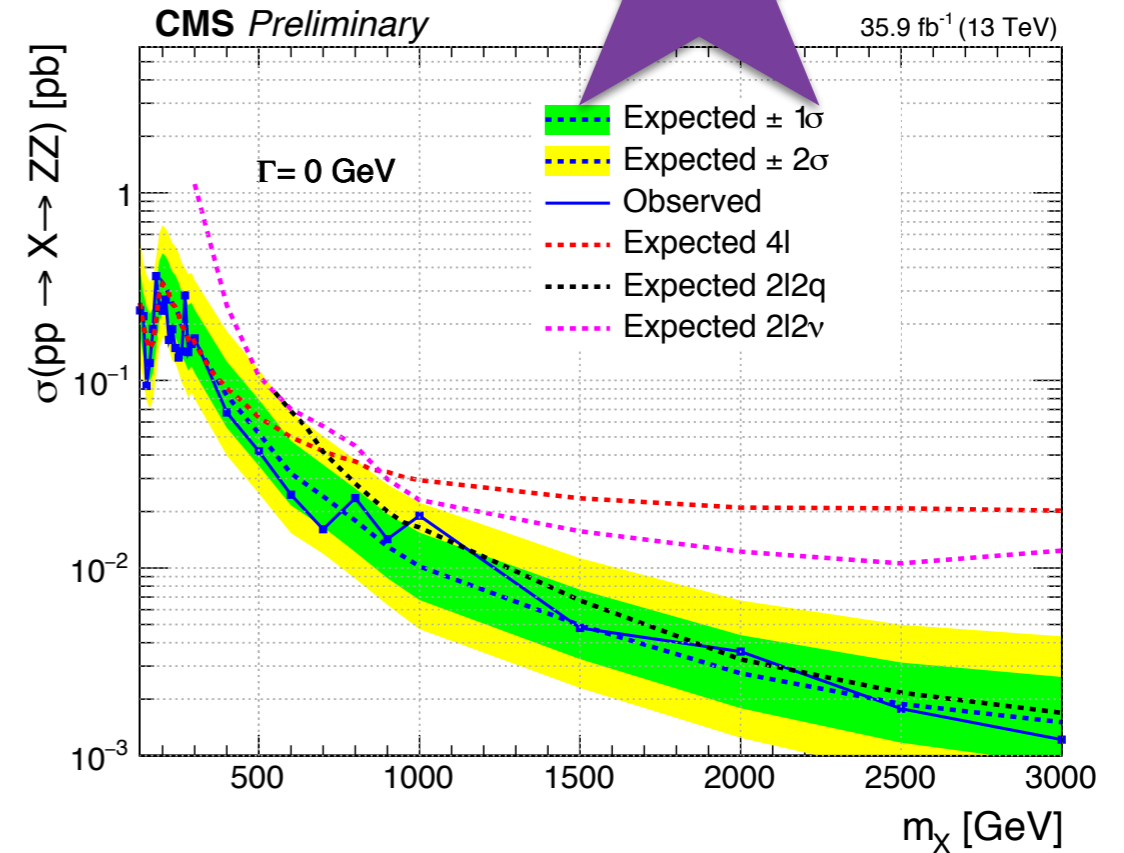


- > Recently released: HIG-17-012
- > Combination of 4l, 2l2v, and 2l2q
- > MELA (Matrix Element Likelihood Analysis) discriminators for 4l & 2l2q (for spin 0 & spin 2)
- > Considers S+B interference



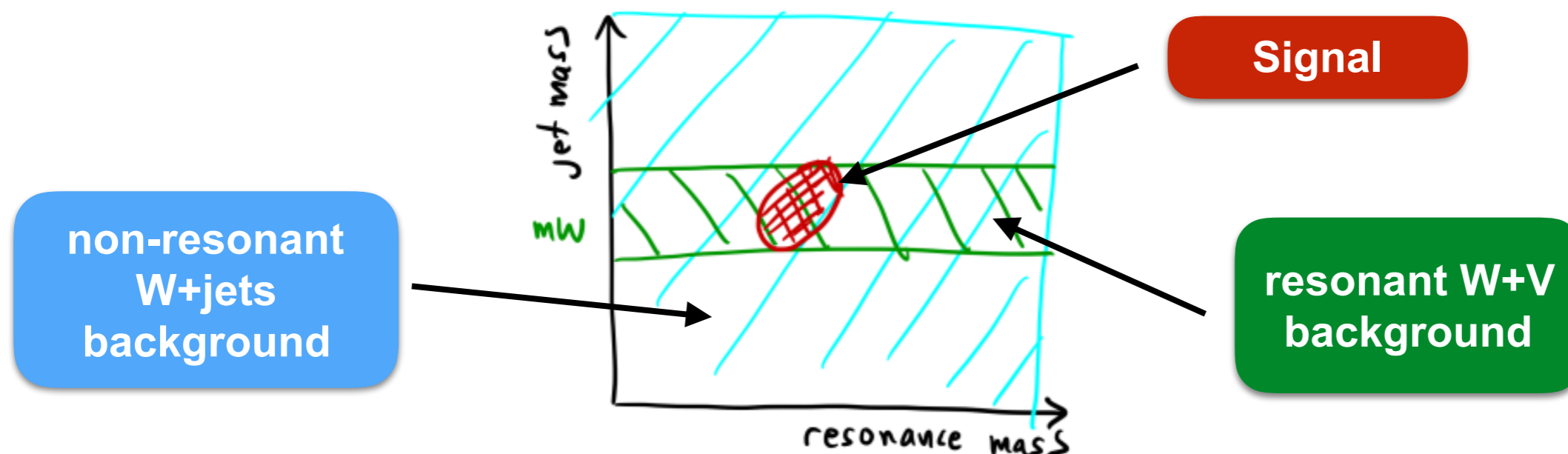
ZV/ZZ ($Z \rightarrow \nu\nu/\ell\ell$) analyses

- > Considers ggH and VBF production
- > Model-independent limits from 130 to 3000 GeV
 - up to 30% resonance width





- > New analysis approach for $WV \rightarrow l\nu qq$ analysis
 - do not use jet mass windows anymore
- > Instead: 2D fit in (m_{WV}, m_{jet}) plane - use full V jet mass range: $30 < m_{\text{jet}} < 210$ GeV
 - make better use of correlations between m_{WV} and m_{jet}
 - much more sideband statistics — use full line-shape of jet mass
 - become less dependent on simulation — learn from data
- > 2D fit: distinguish between
 - non-resonant W+jets ($W(l\nu)+\text{jets}$, $t\bar{t}$ bar with non-W V jet)
 - resonant W+V ($t\bar{t}$ bar, diboson) background processes





> Using conditional probabilities for signal and each background component starting from simulation

- taking into account scale and resolution

> Signal peaks in both m_{WV} and m_{jet} :

$$P_{sig}(m_{WV}, m_{jet} | \theta(M_X)) = P_{WV}(m_{WV} | \theta_1(M_X)) \times P_j(m_{jet} | \theta_2(M_X))$$

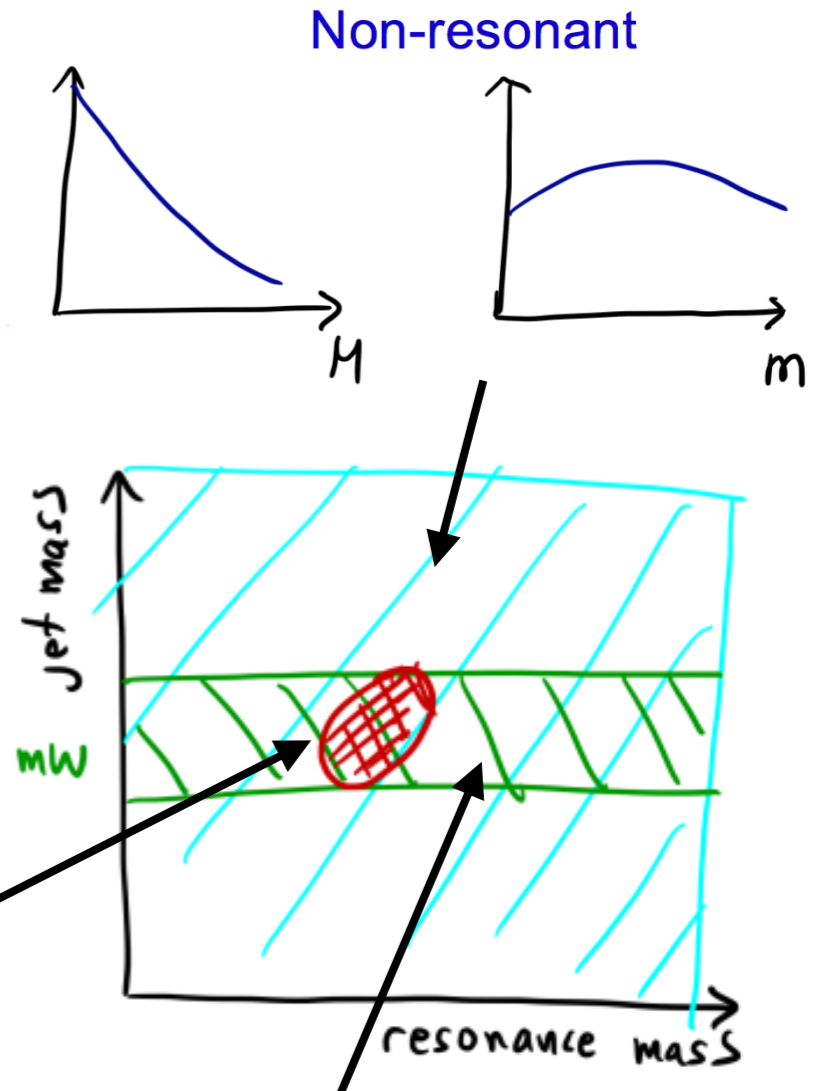
> W+jets background:

- Conditional probability of m_{WV} as function of m_{jet} :

$$P_{W+jets}(m_{WV}, m_{jet}) = P_{WV}(m_{WV} | m_{jet}, \theta_1) \times P_j(m_{jet} | \theta_2)$$

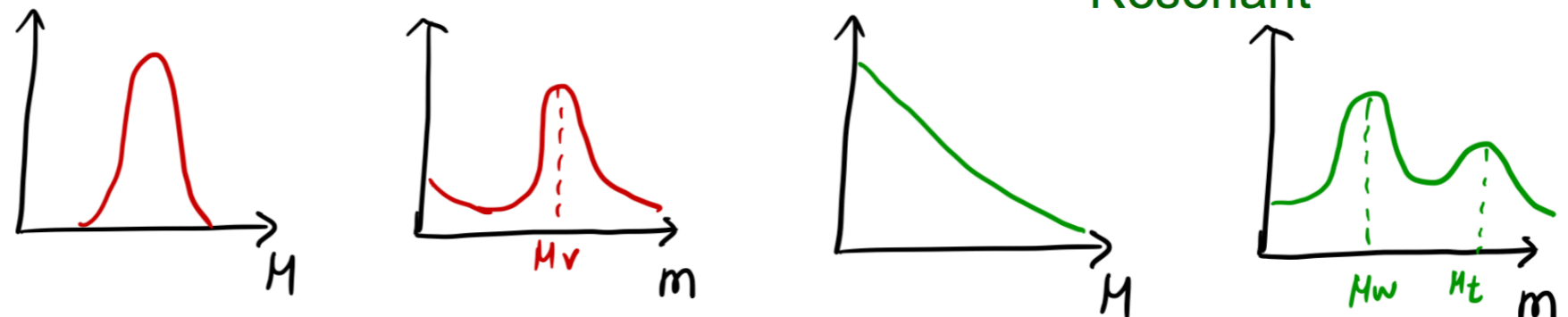
> W+V background:

$$P_{W+V}(m_{WV}, m_{jet} | \theta) = P_{WV}(m_{WV} | \theta_1) \times P_j(m_{jet} | \theta_2(m_{WV}))$$



Signal

Resonant



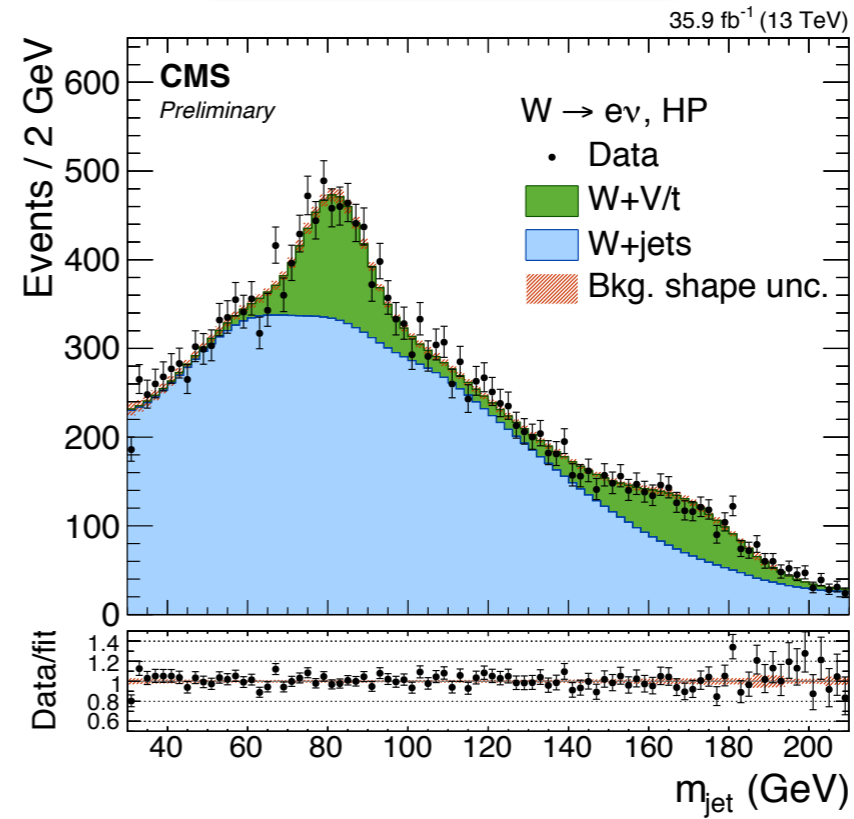
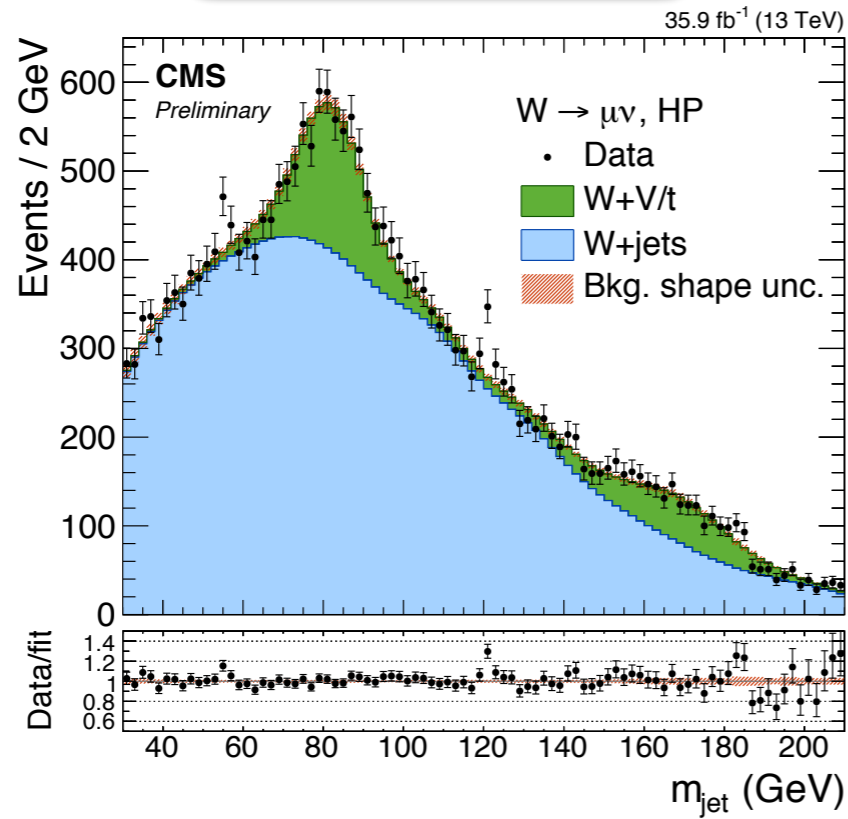
Post-fit projections



muon

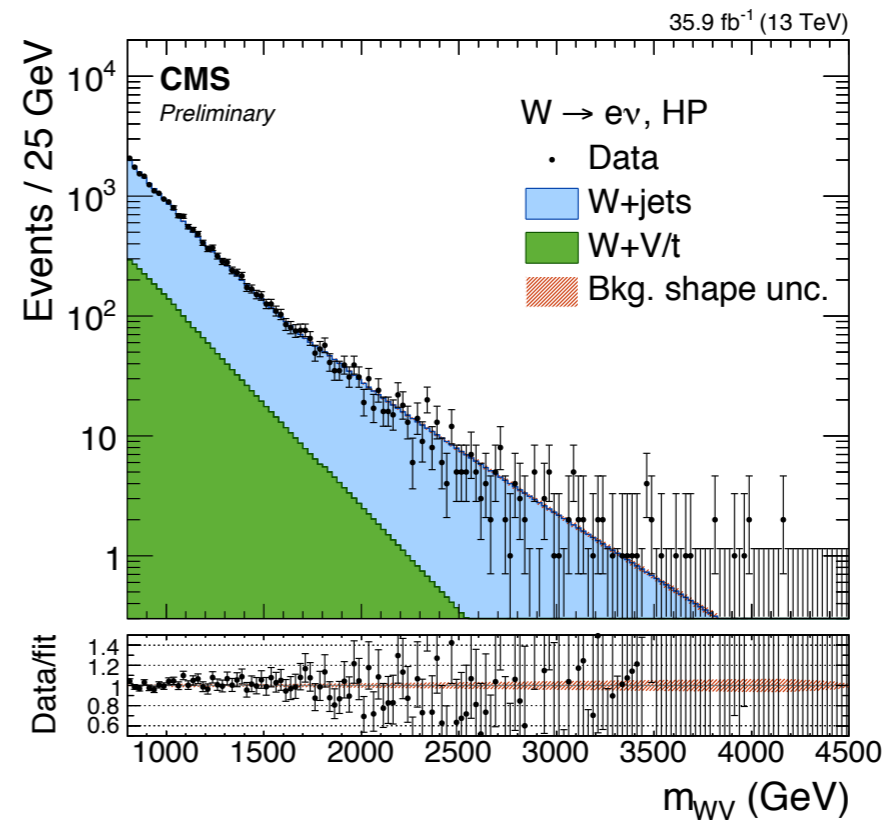
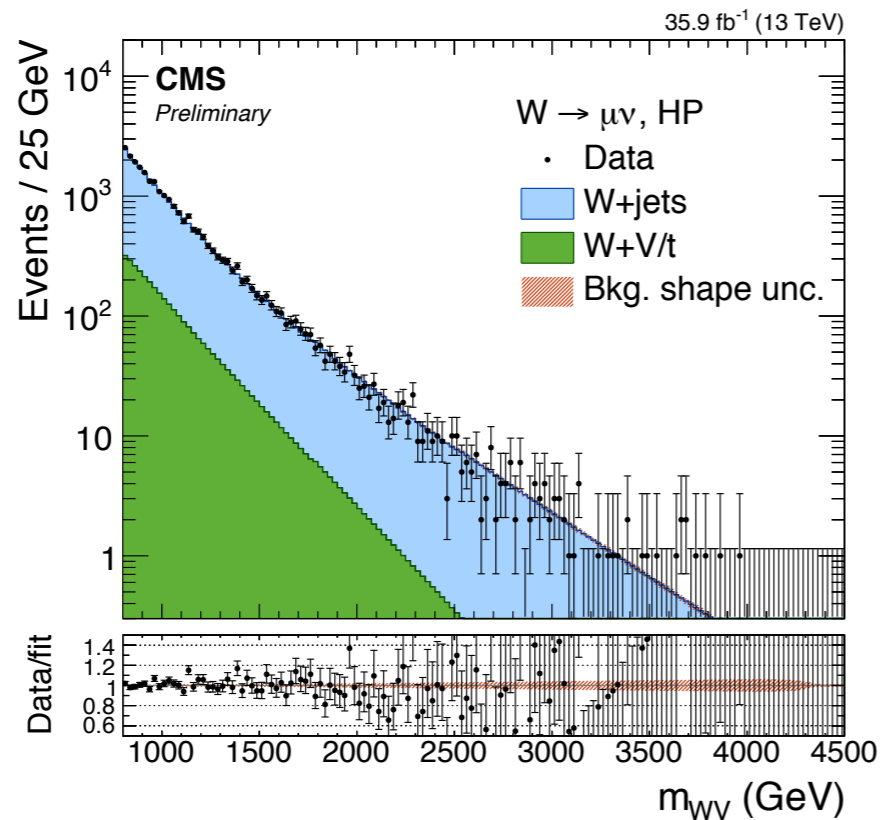
electron

m_{jet}



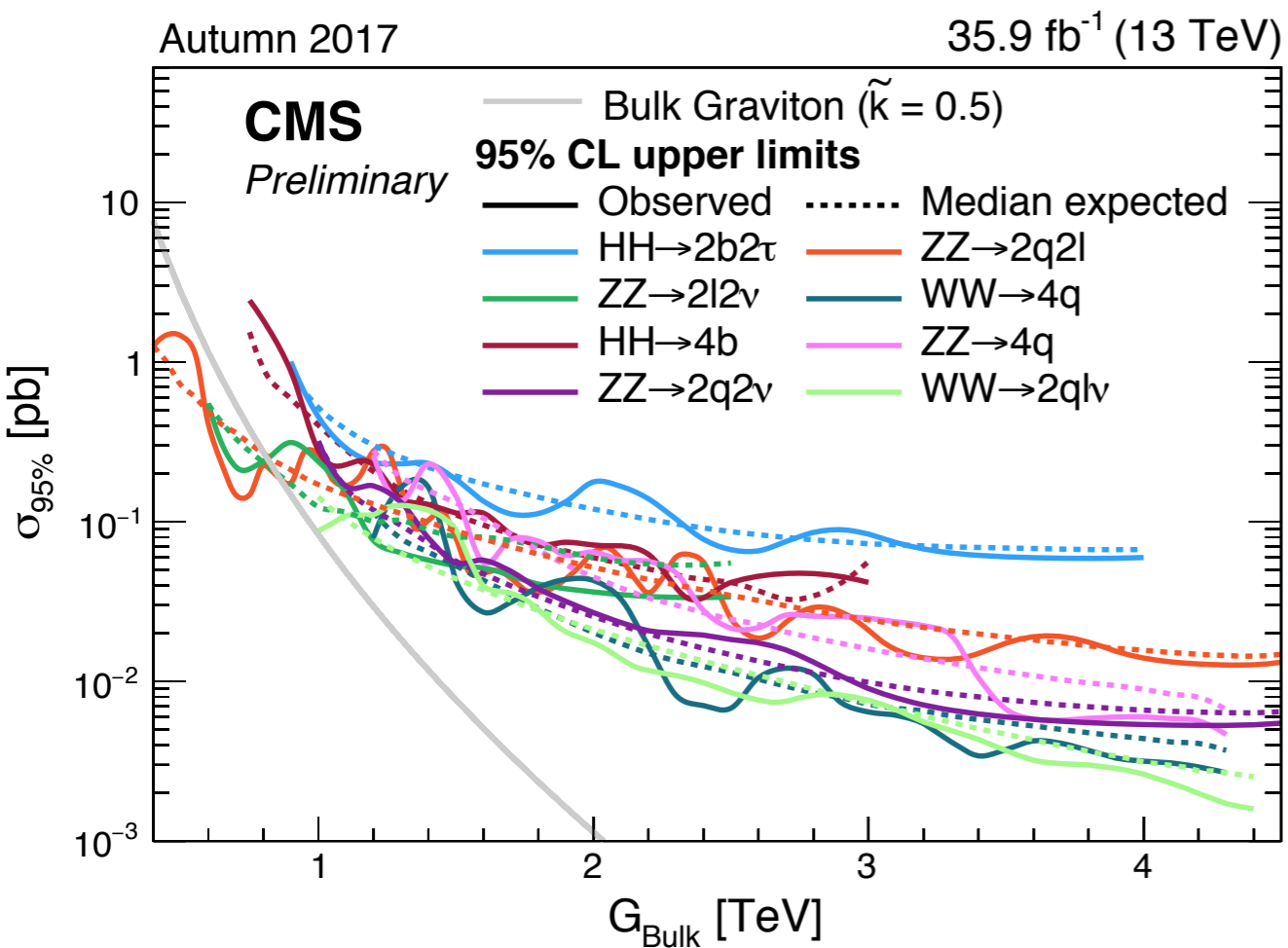
**Nicely
modelling
both W and
top peaks!**

m_{WV}

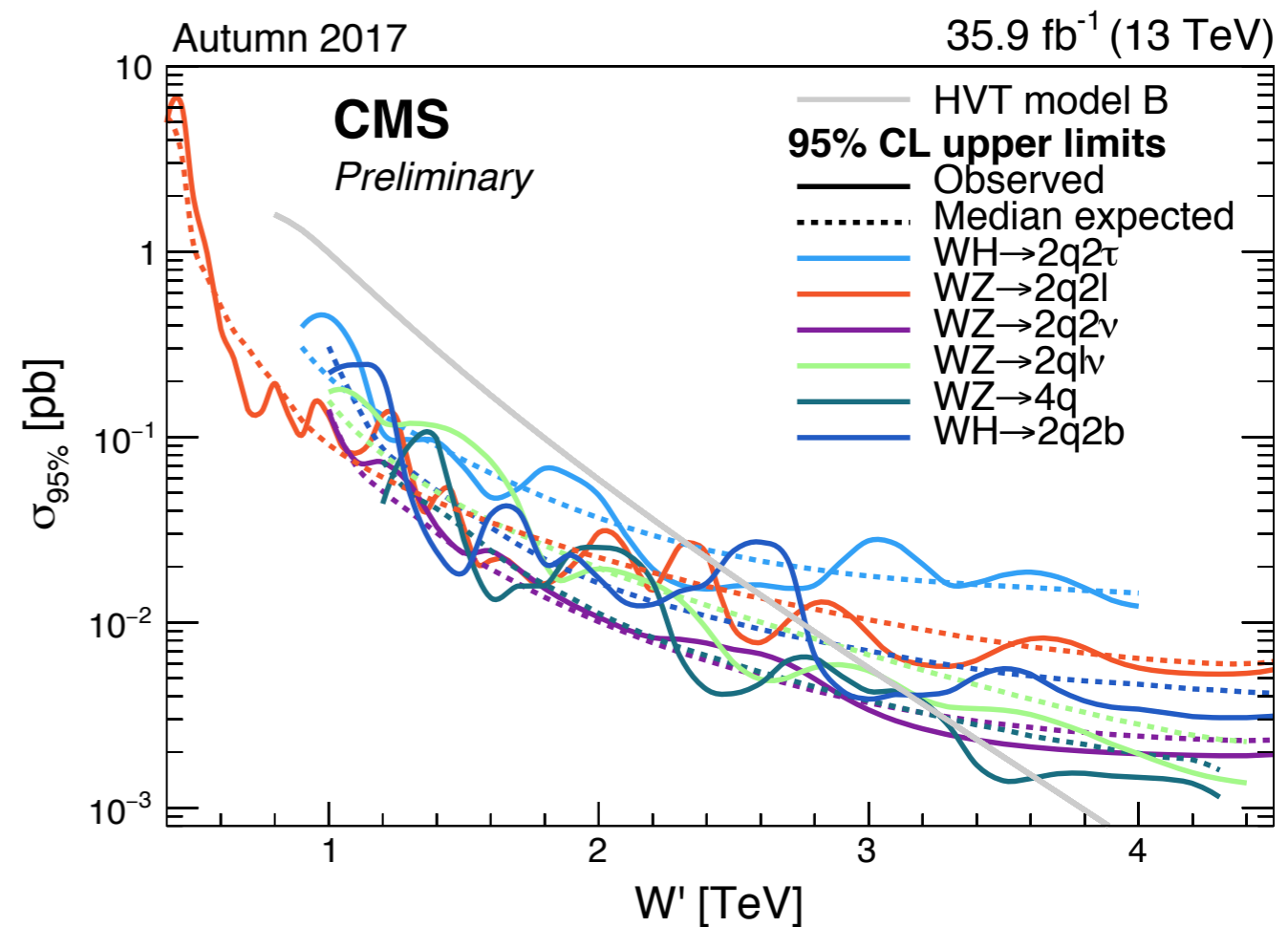


High mass diboson resonances summary

bulk graviton (spin-2)



heavy vector triplet (W', spin-1)



HH \rightarrow 2b2 τ /WH \rightarrow 2q2 τ (B2G-17-006)

ZZ/WZ \rightarrow 2q2l (B2G-17-013)

ZZ \rightarrow 2l2 ν (B2G-16-023, submitted to J. High Energy Phys.)

WW/WZ/ZZ \rightarrow 4q (B2G-17-001, submitted to Phys. Rev. D)

WW/WZ \rightarrow 2ql ν (B2G-16-029)

ZZ \rightarrow 2q2 ν (B2G-17-005)

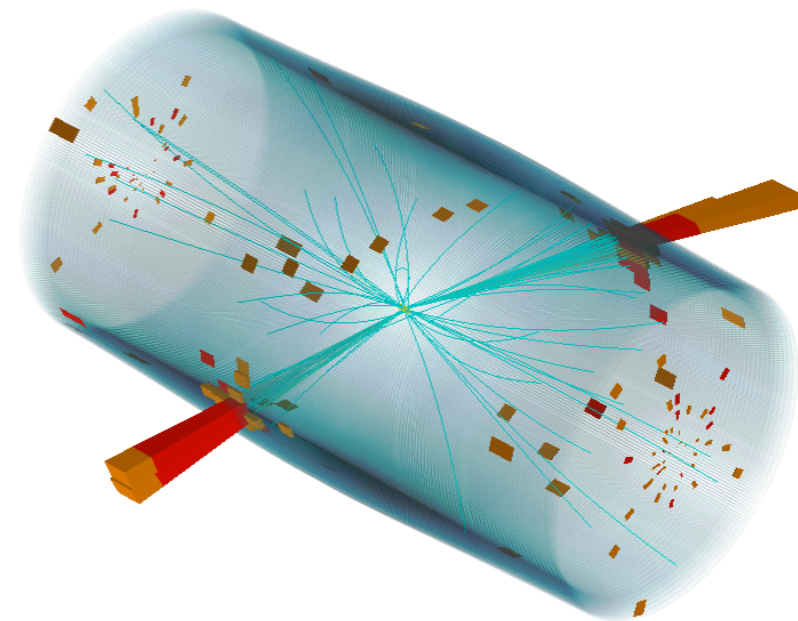
HH \rightarrow 4b (B2G-16-026, submitted to Phys. Lett. B)

WH \rightarrow 2q2b (B2G-17-002, Eur. Phys. J. C 77 (2017) 636)

Challenges:

- > Run-2 data set will yield higher statistics, but no increase in centre-of-mass energy
- > Imperative to further improve methodology
 - multi-dimensional fits make best use of statistics
- > Further work needed on understanding jet substructure
 - new boson tagging algorithms on the market
 - can also profit from machine learning

Candidate ZZ event
Dijet mass: 3.2 TeV

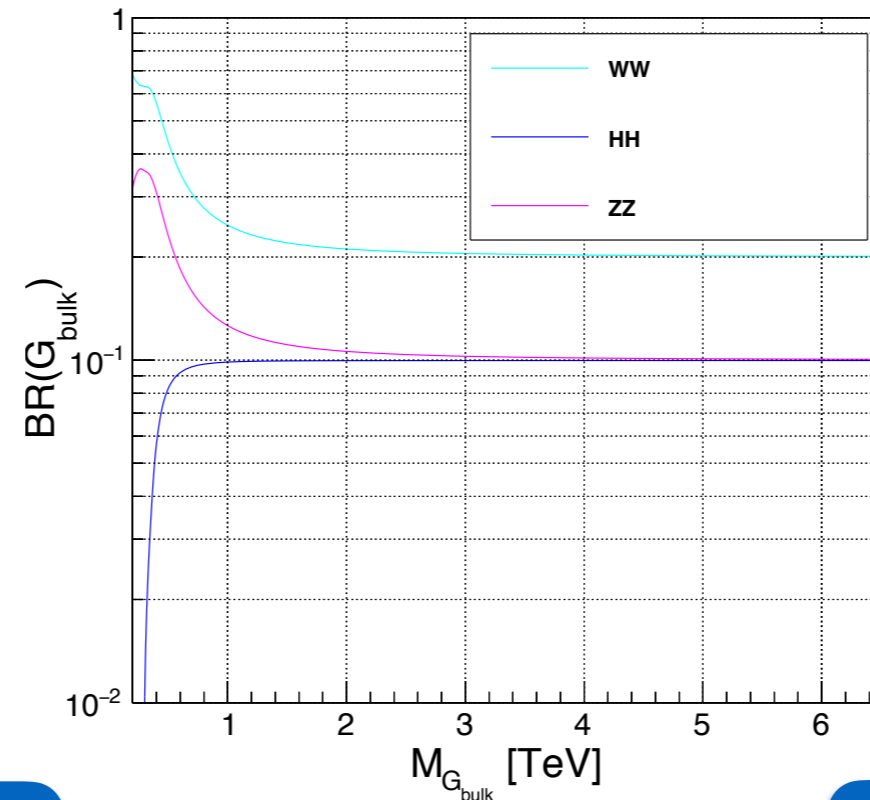


- > CMS has an extensive diboson resonance search programme
- > Several additions and improvements w.r.t. previous analyses
 - significantly higher sensitivity and extended mass range
- > Well set up to make best use of full Run-2 data set



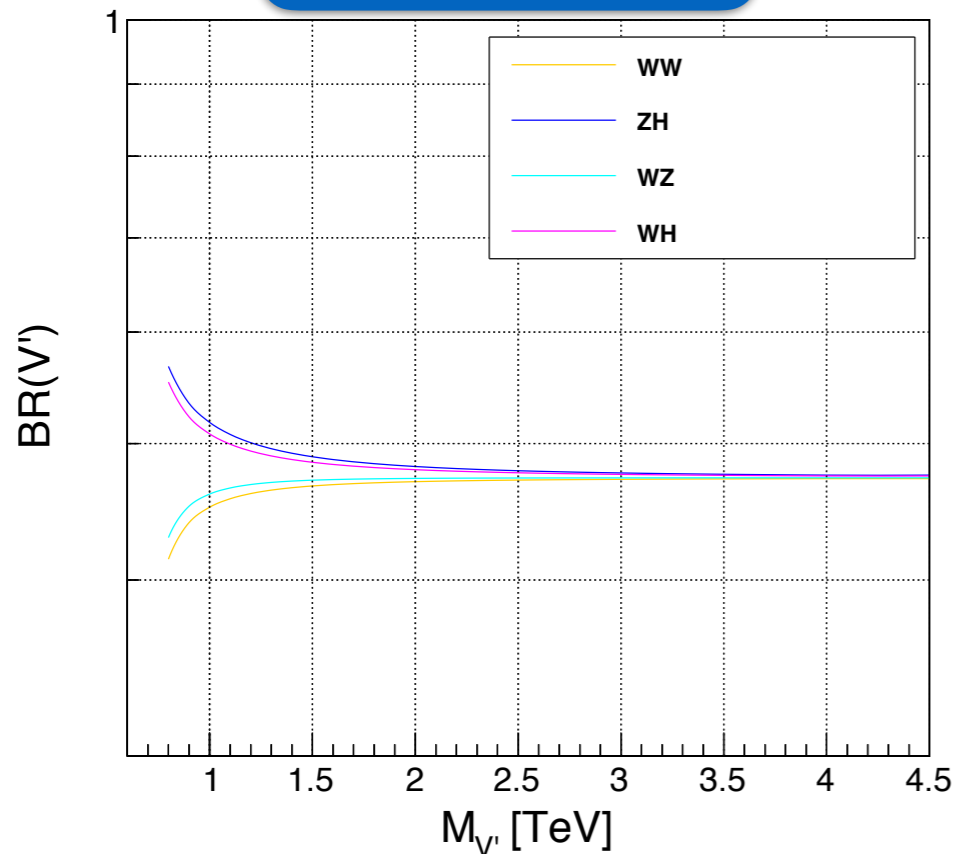
Branching fractions

bulk graviton

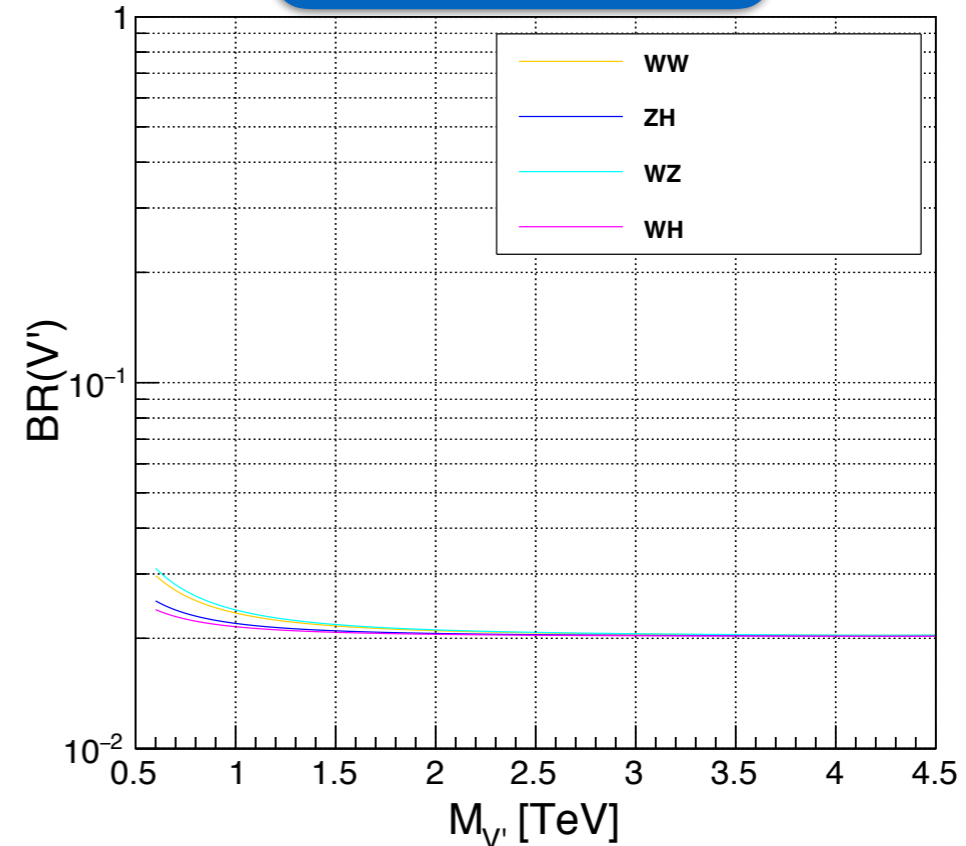


vector bosons are produced with a **longitudinal polarization** in more than **99%** of the cases, resulting in a **~24% higher acceptance per boson** than for models producing **transversally polarized** vector bosons

HVT model B

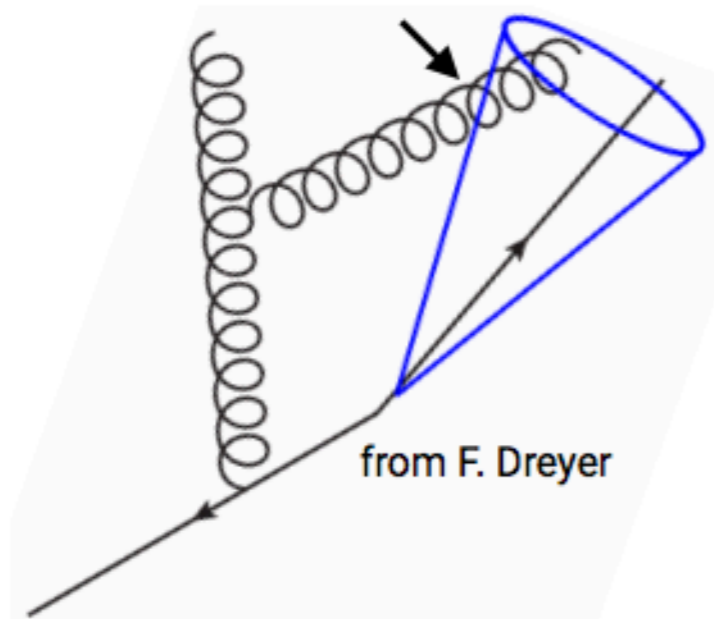


HVT model A

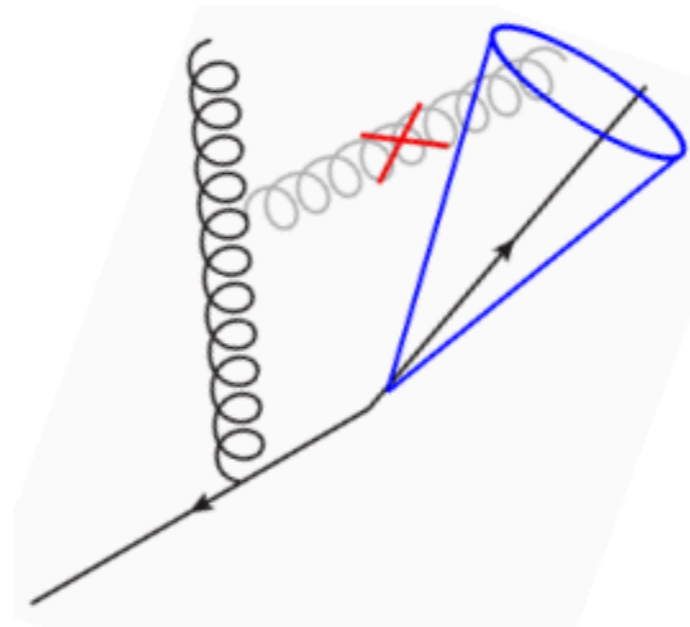


Soft drop vs. pruning

Pruning: Soft gluon radiating into jet not removed



Softdrop: All soft stuff removed. NGL free

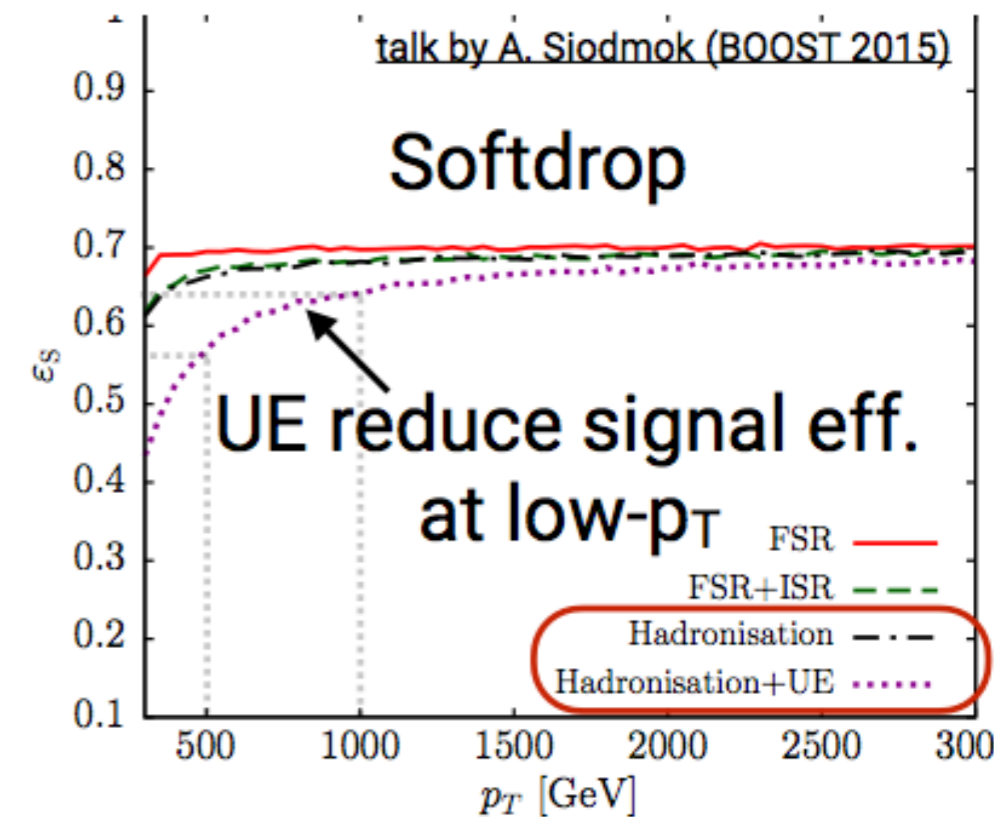


> Pruning not “perturbatively robust”

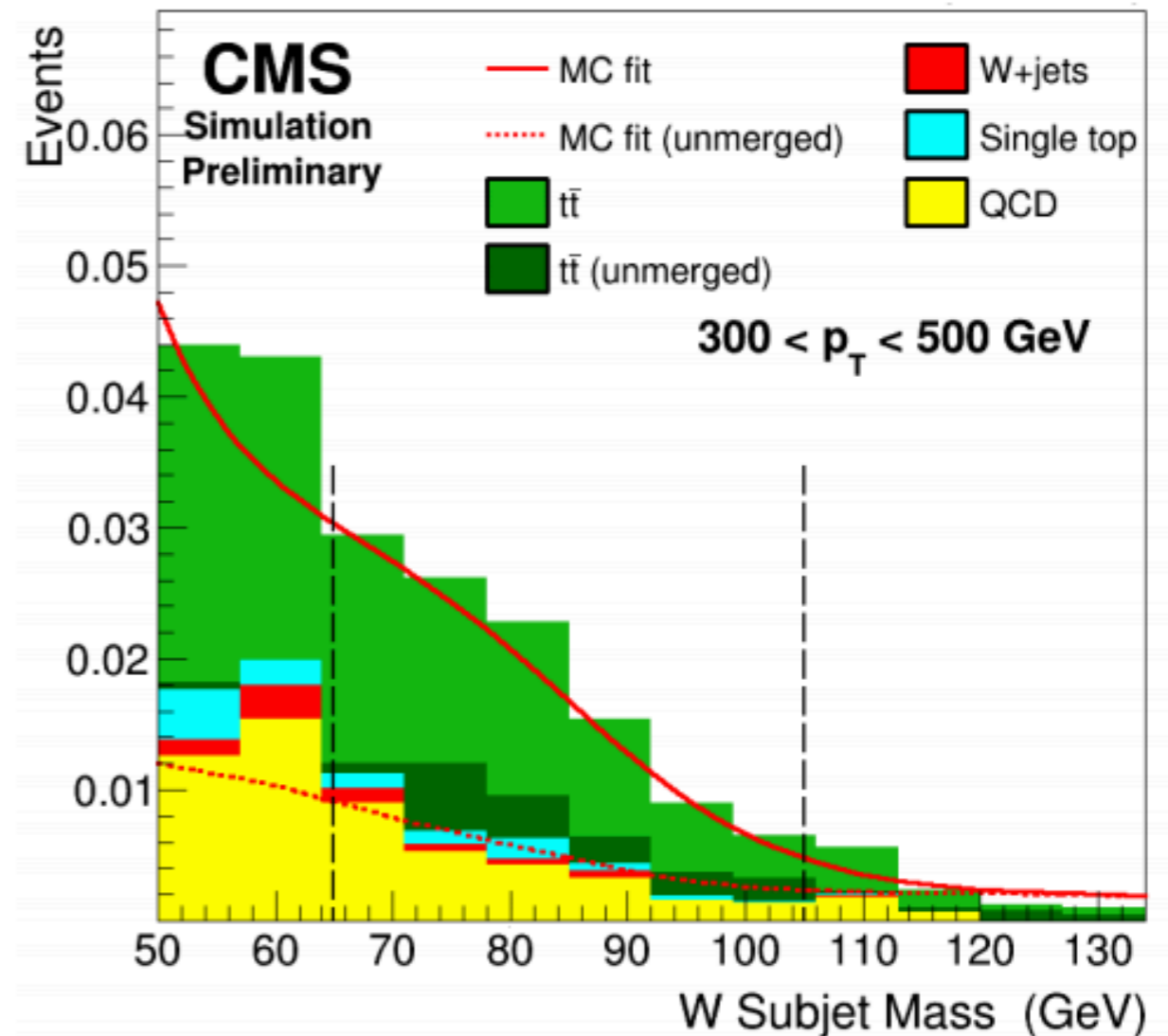
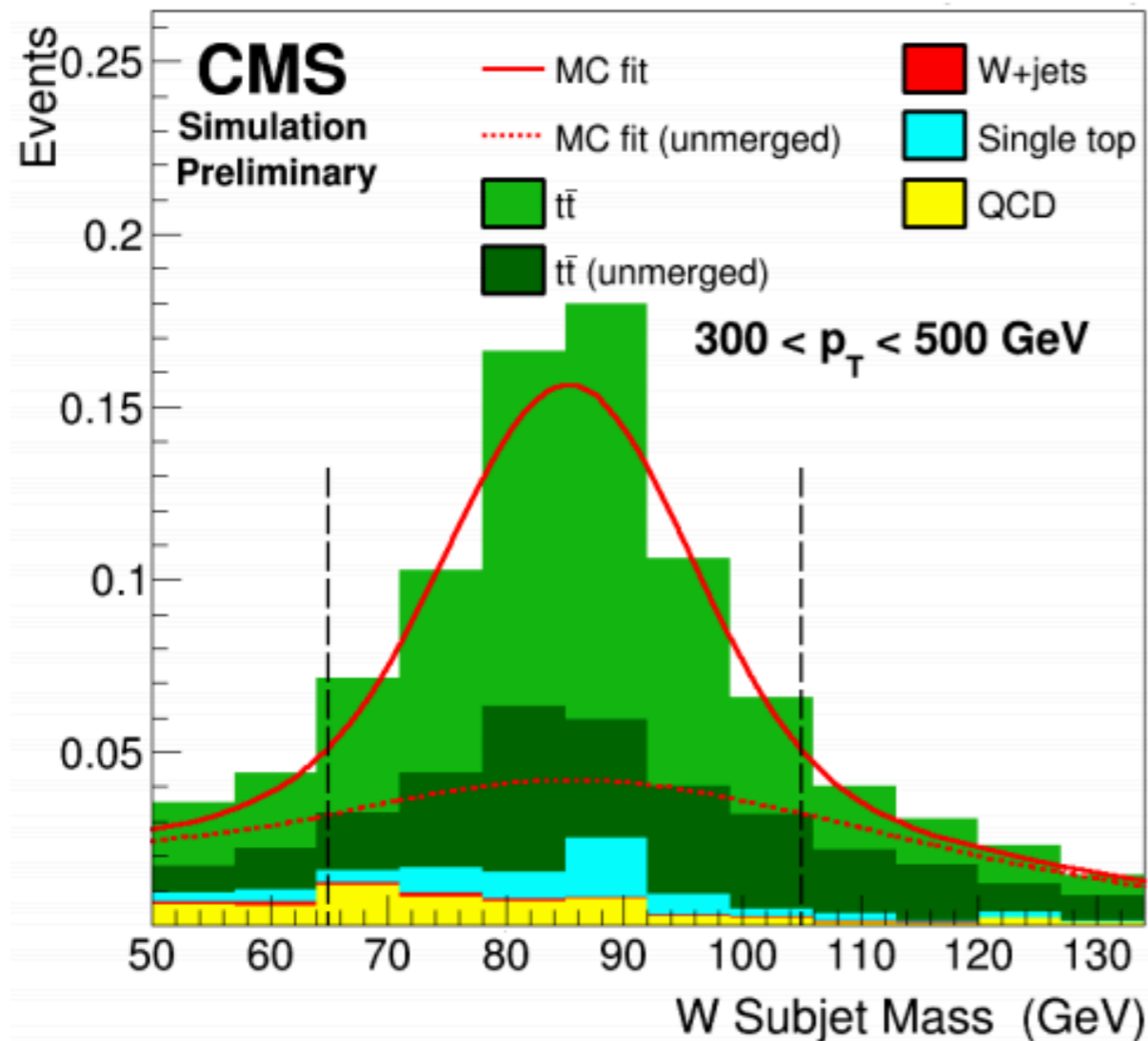
- not soft radiation free
- nonglobal logarithmic terms (NGLs) in mass

> Soft drop:

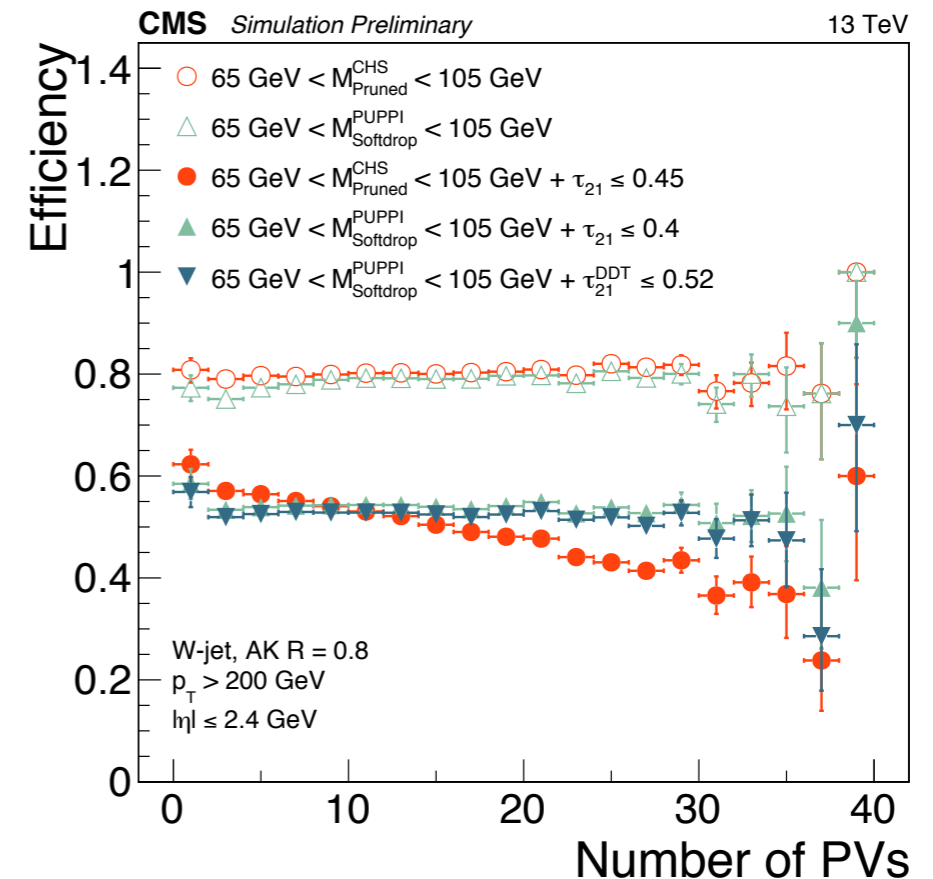
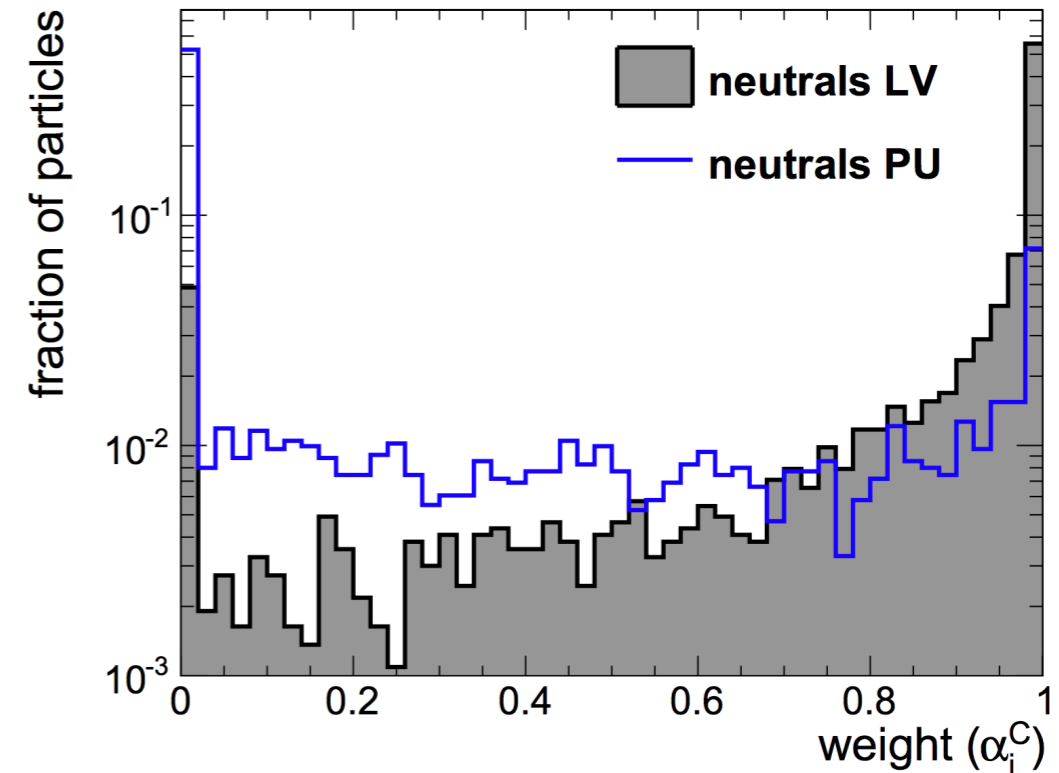
- removes all sensitivity to soft divergences
- however, signal sensitive to underlying event (for multijets, pruning more sensitive to UE)
- correct with PUPPI and mass corrections



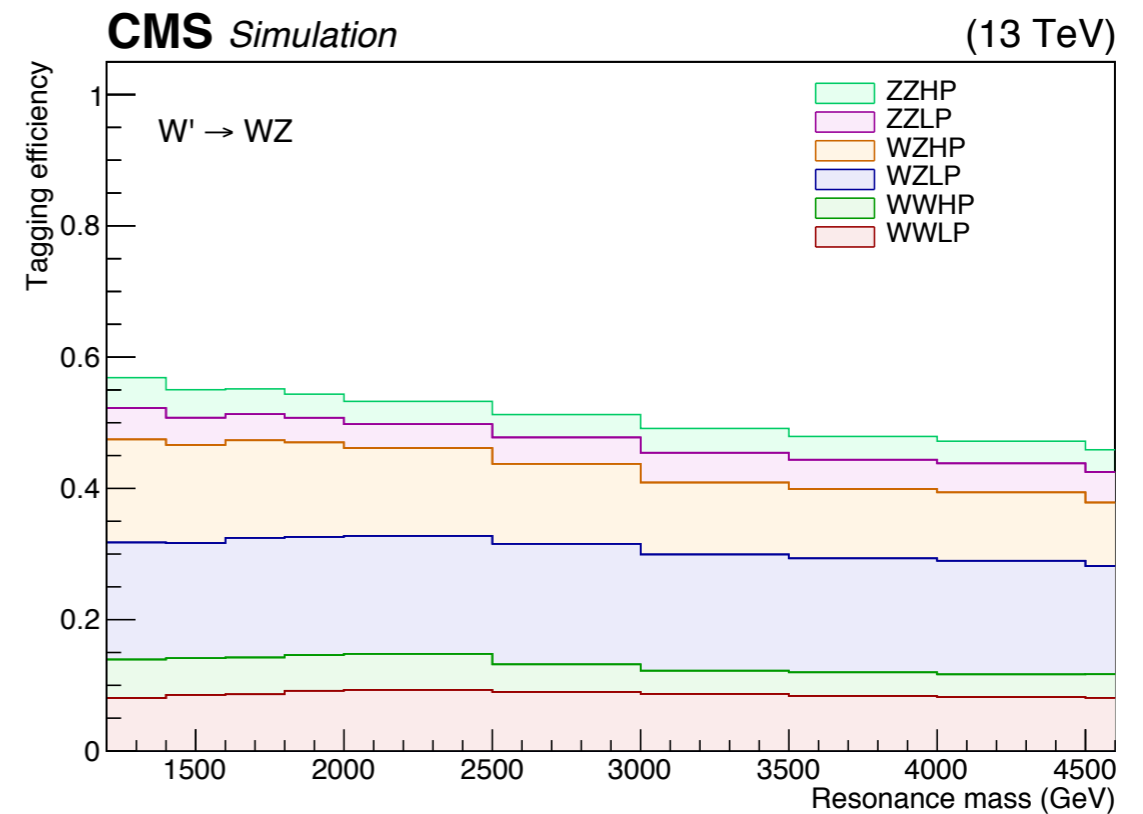
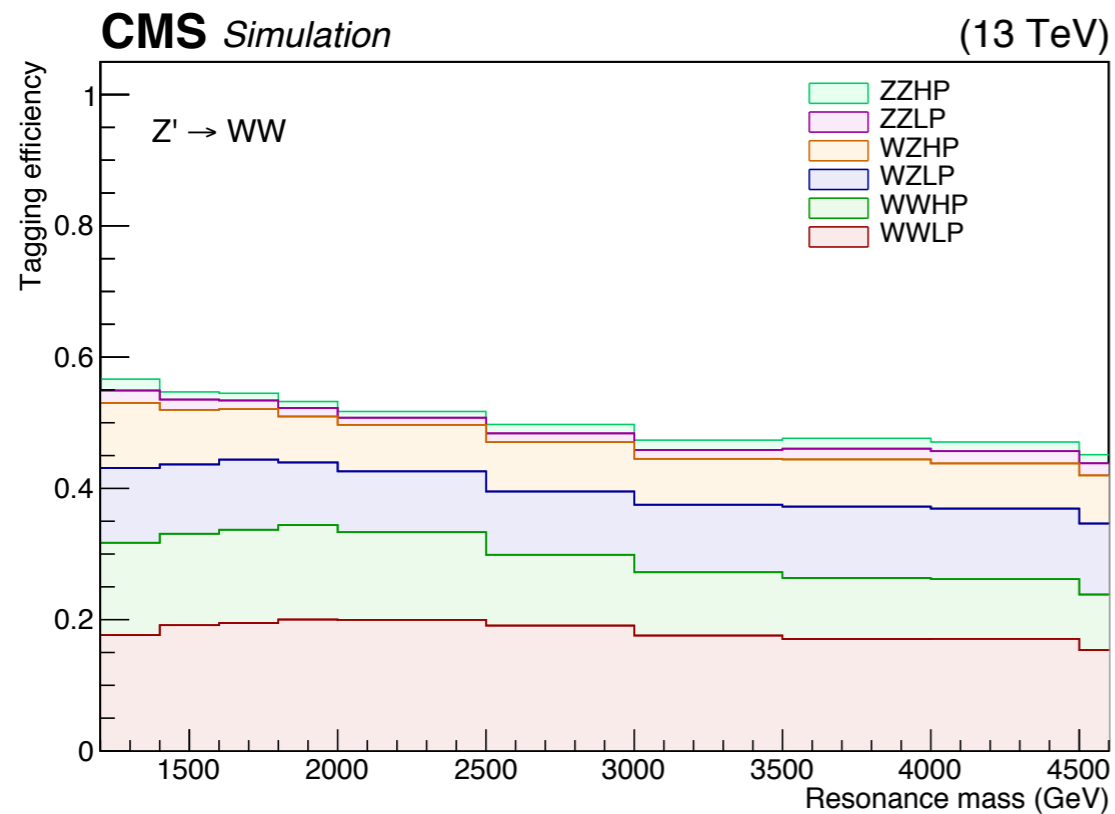
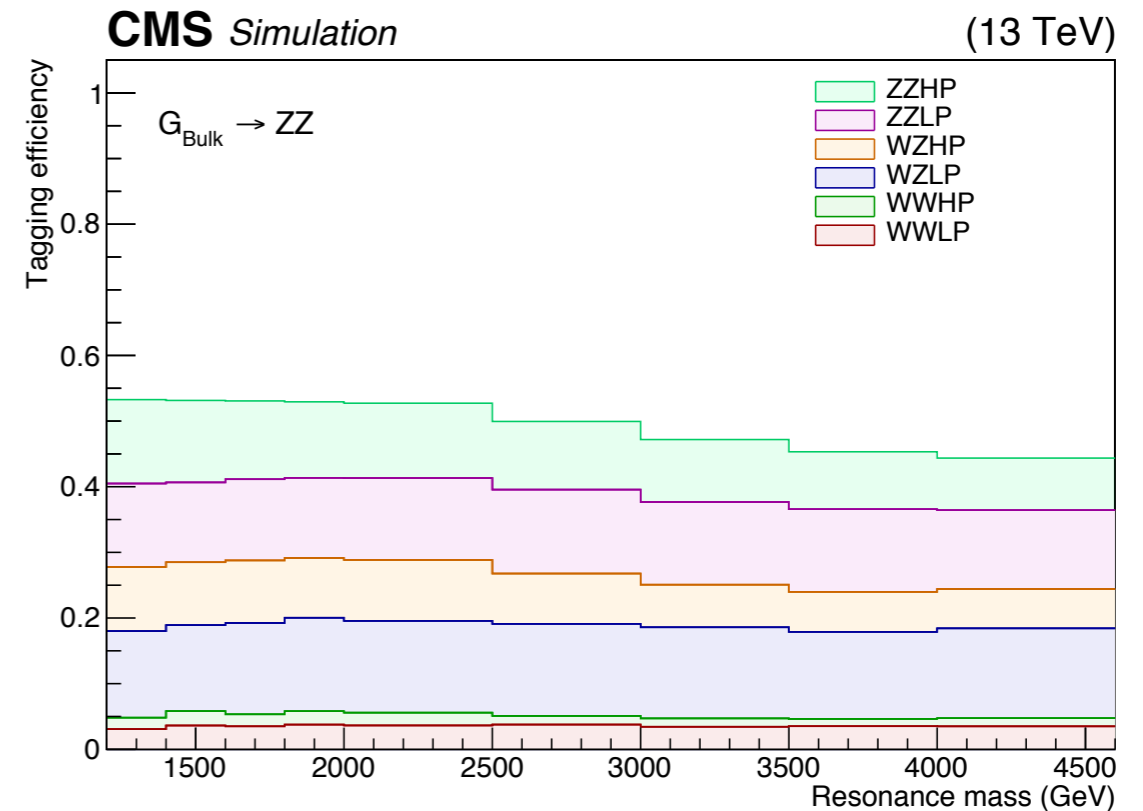
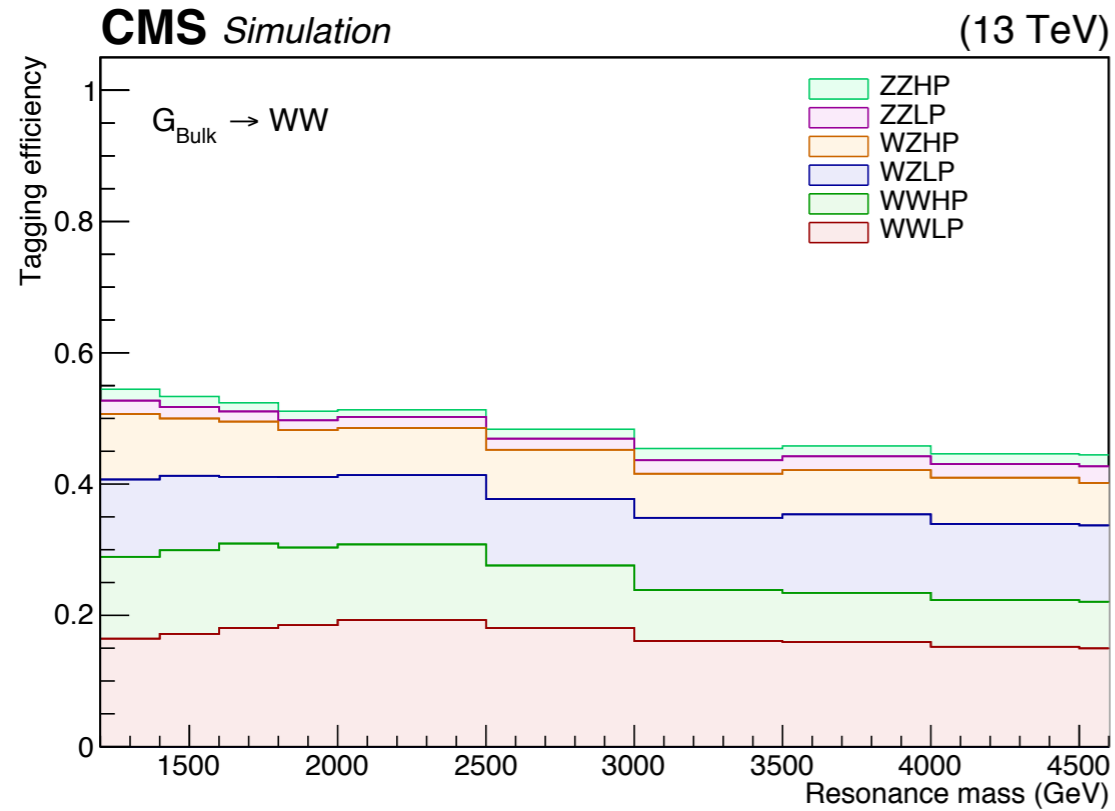
- Previously used simulation-based extrapolation of scale factors
- Now have data-based calibration available



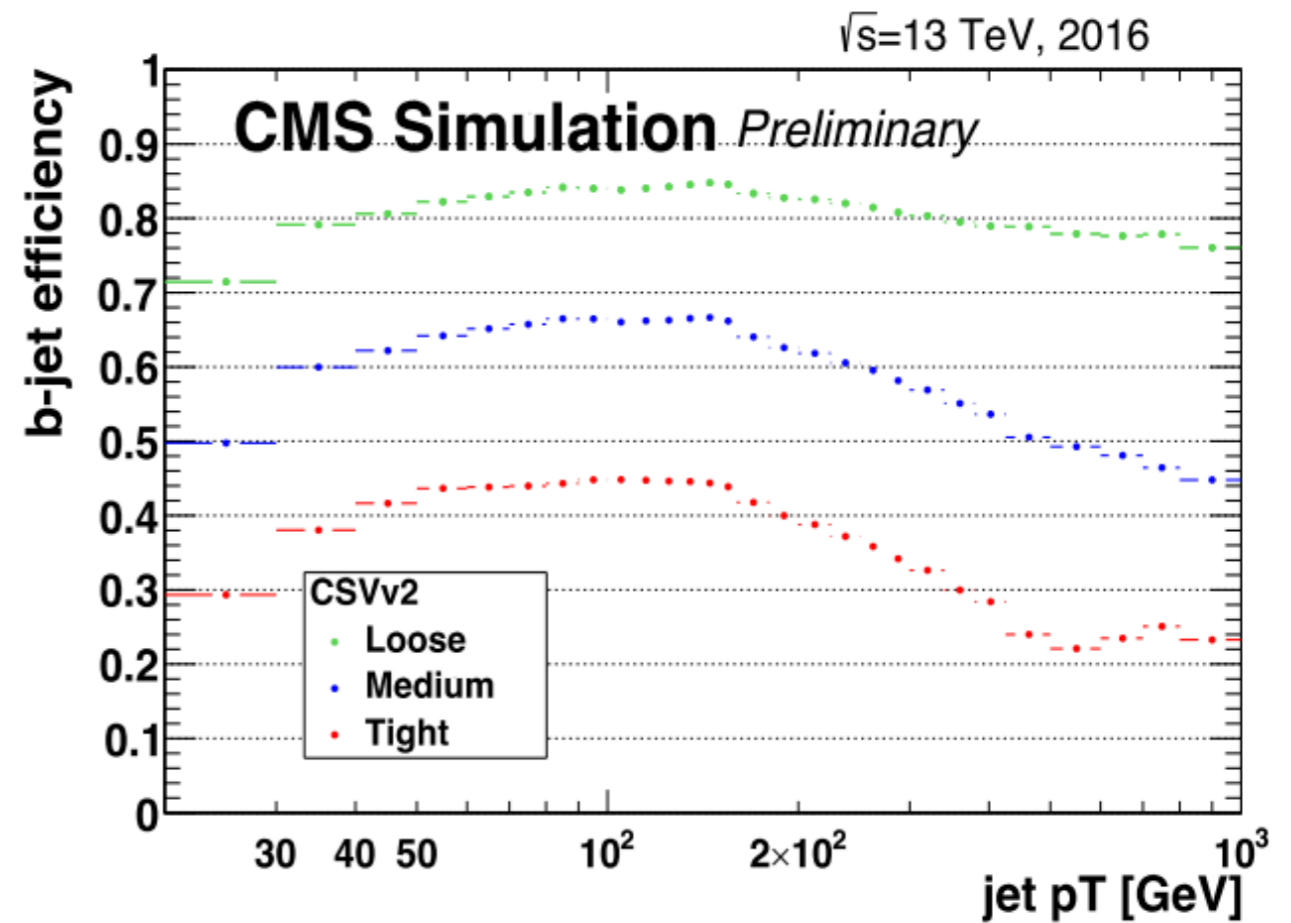
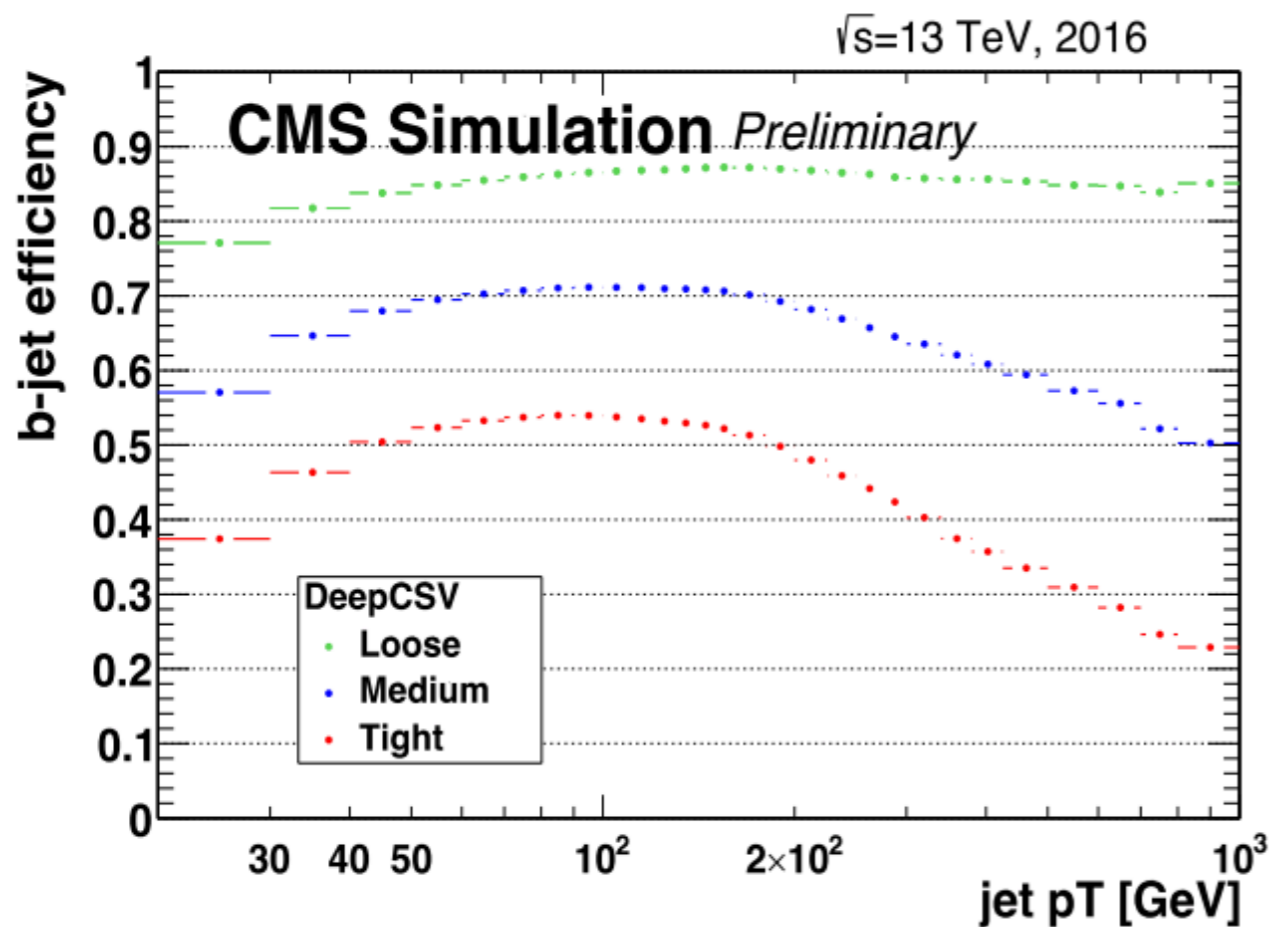
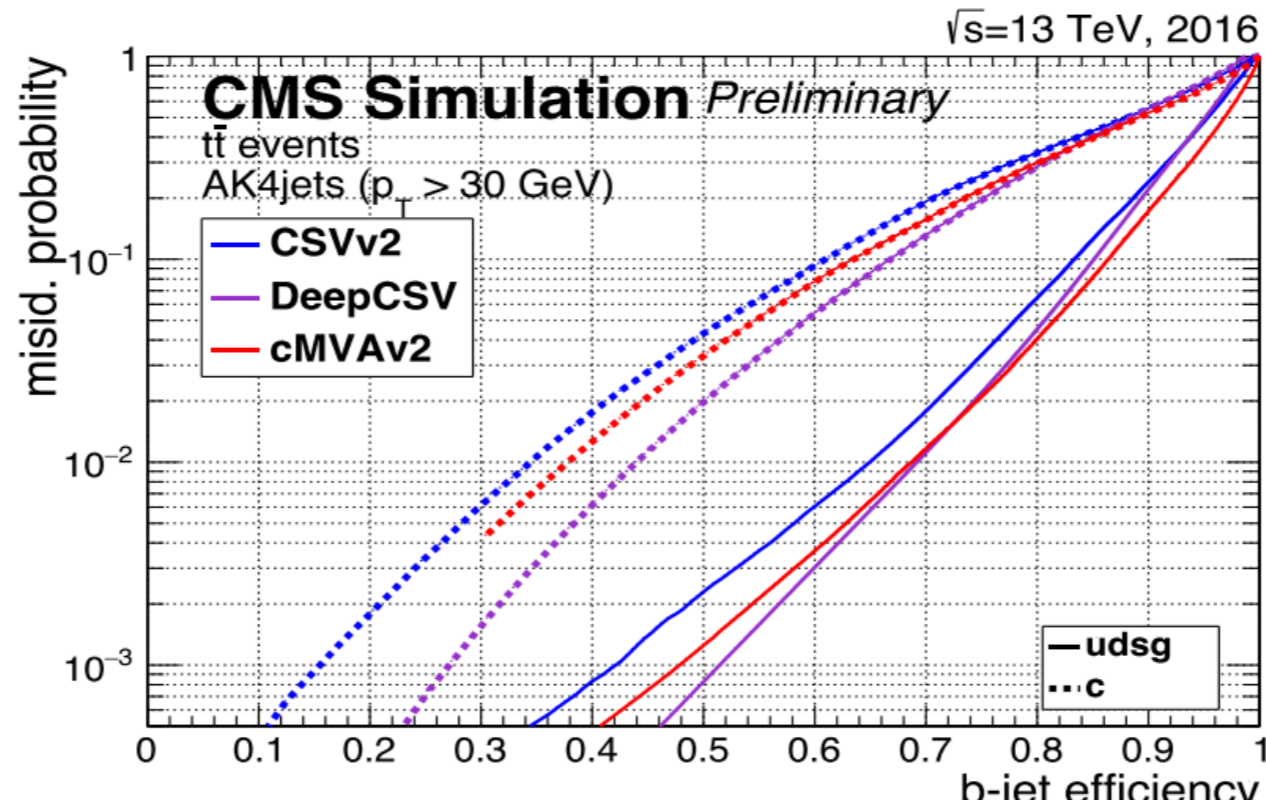
- Tracks can point to PU vertices, only keep charged tracks that come from the primary vertex
- Draw a cone around each neutral Particle Flow candidate
- Define a local metric, α , that differs between pileup (PU) and leading vertex (LV)
- For the neutrals, ask “how PU-like is α for this particle?”, compute a weight for how LV-like it is
- Reweight the four-vector of the particle by this weight



V-tagging efficiencies

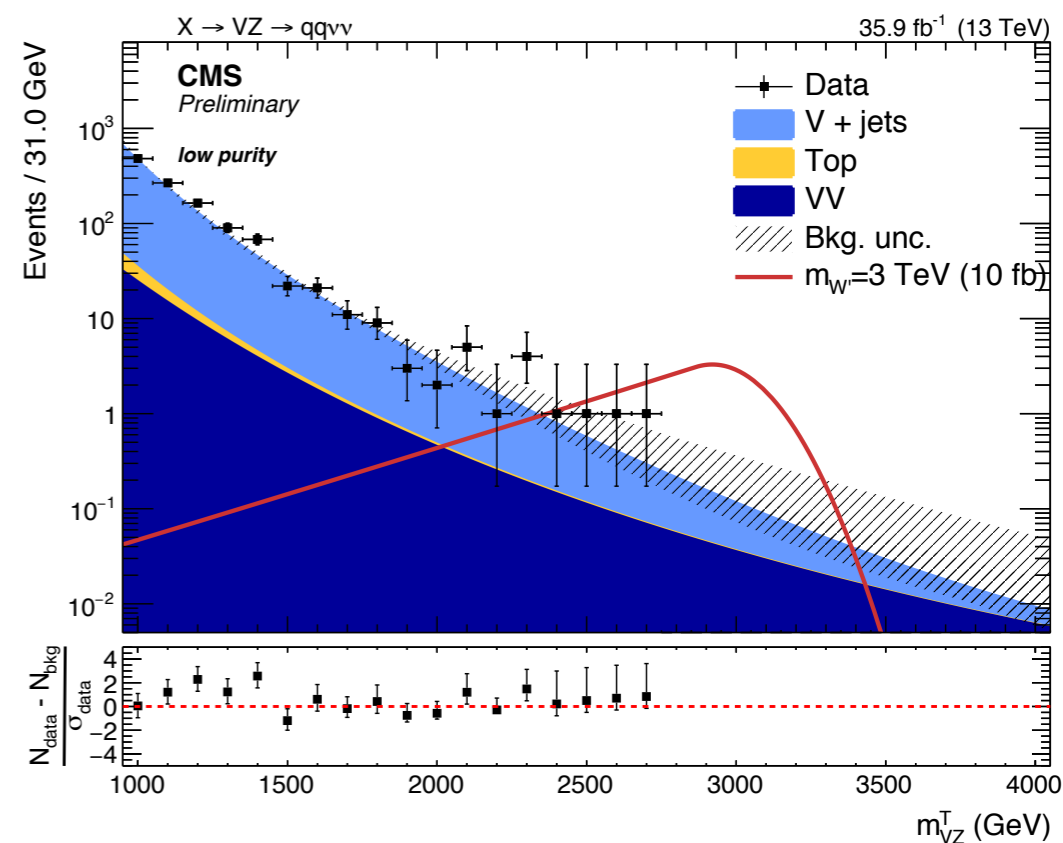
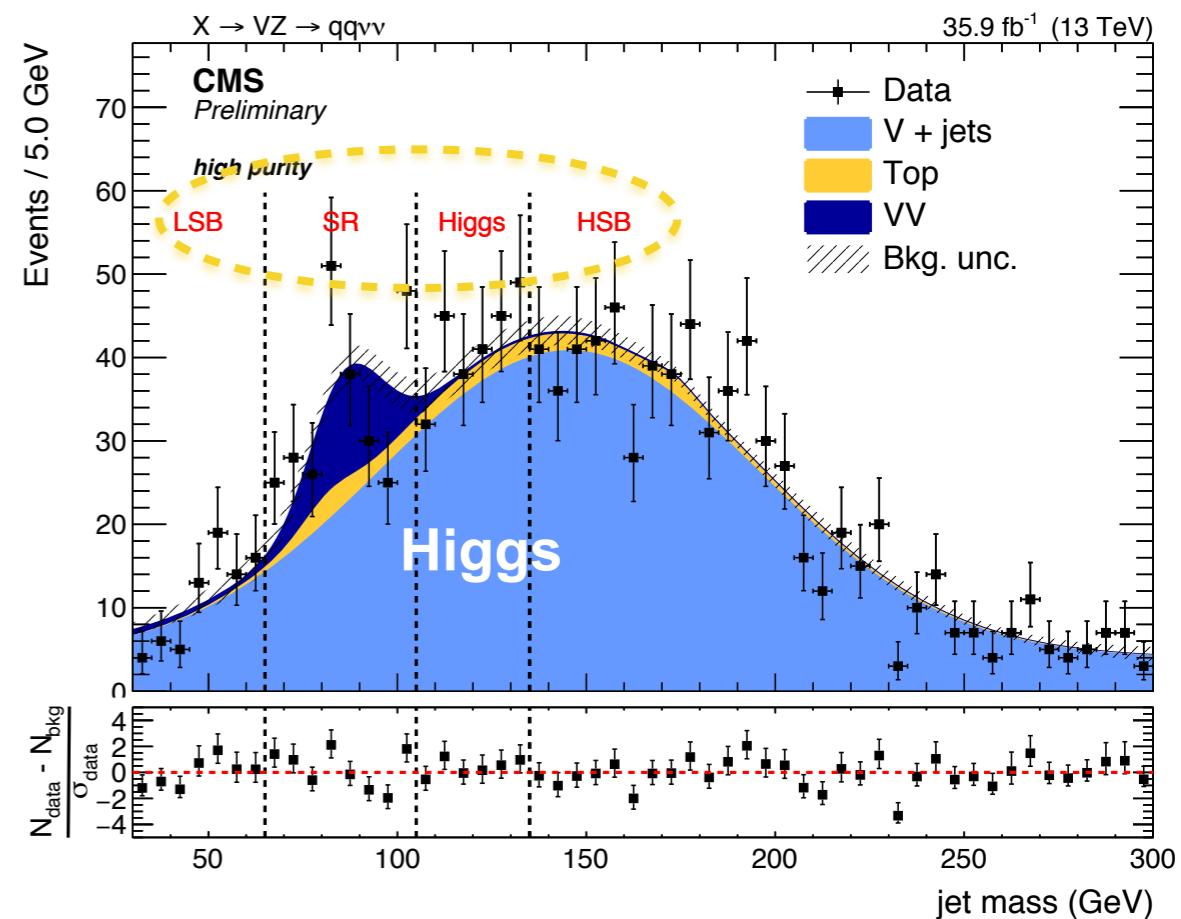


DeepCSV performance



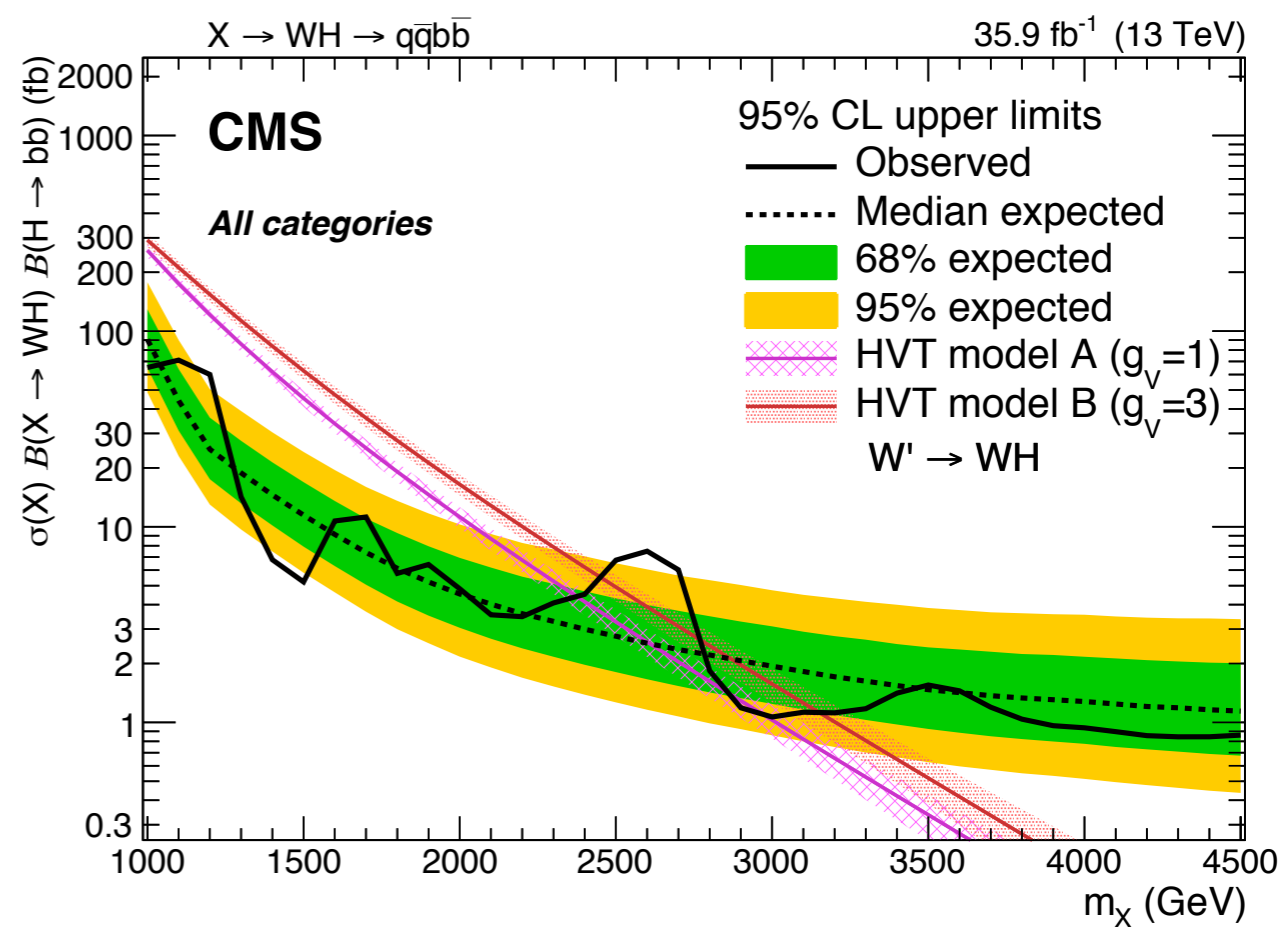
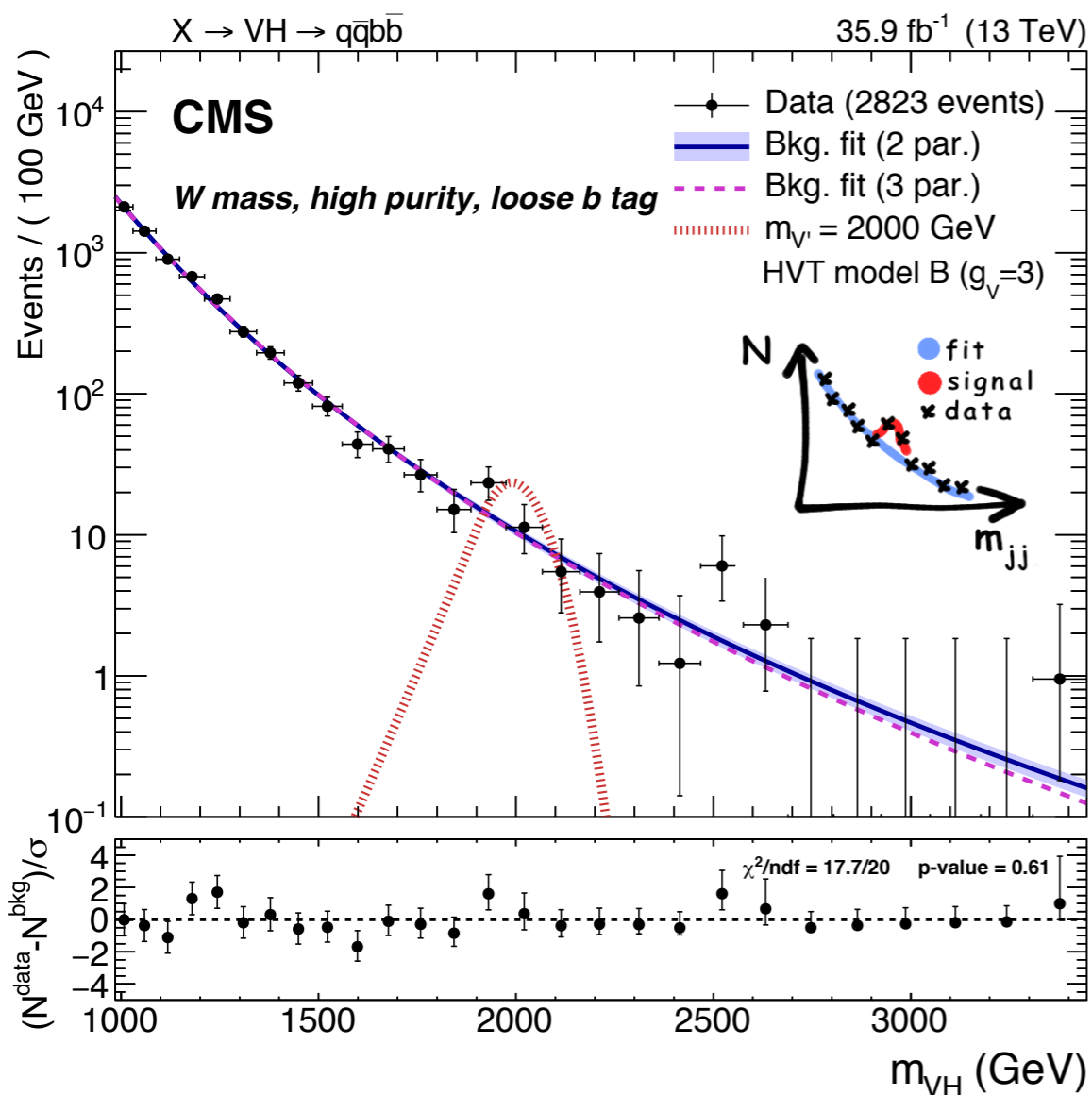
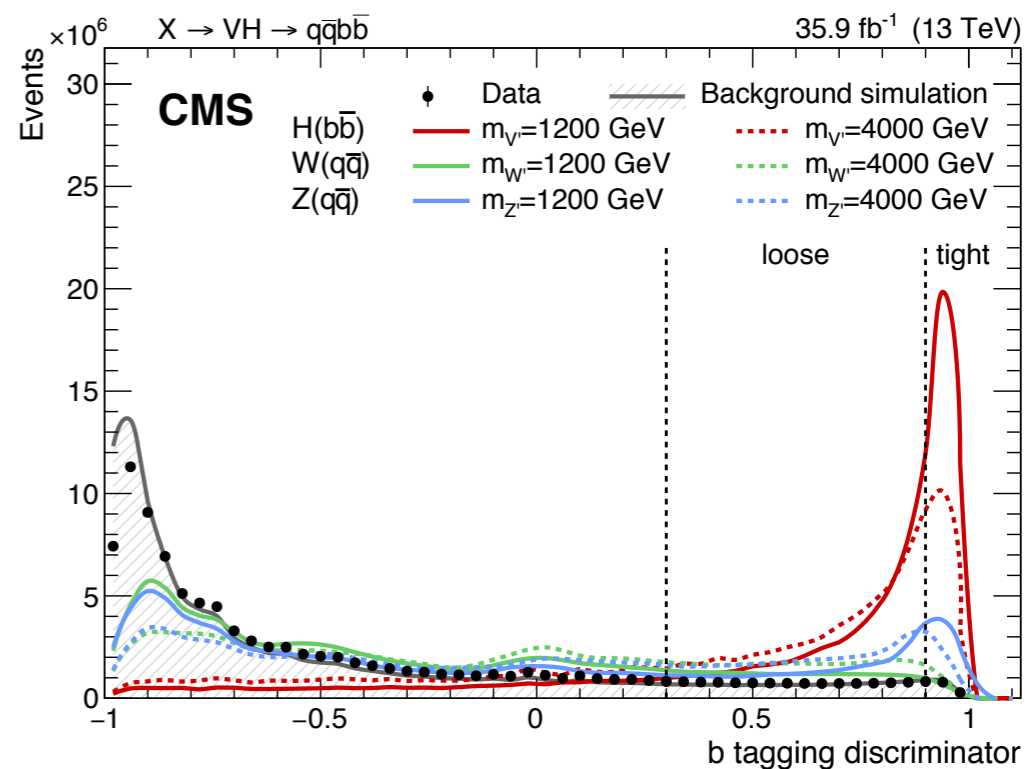
alpha-ratio background estimation

- Statistics in MC simulated samples still limited
- Furthermore, analysis performed in **extreme phase space**
- Use **jet mass sidebands** (40-65 GeV, 135-160 GeV) to exploit correlation between soft-drop jet mass and resonance mass
 - Higgs mass region kept blind
- Determine **ratio of simulated to data** distributions in sideband
- **Extrapolate** to signal region using **transfer function** (based on simulation)
- Method accounts for data-MC differences in shape and normalisation



Vh analysis

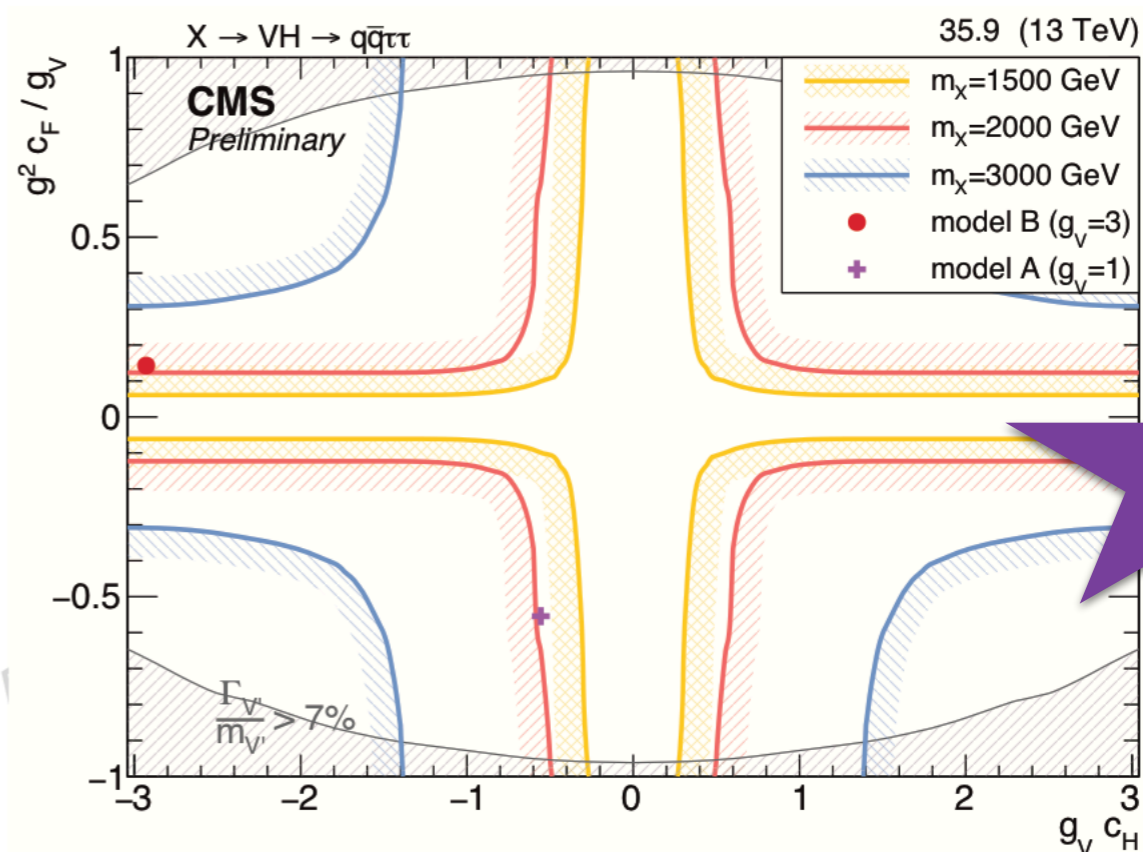
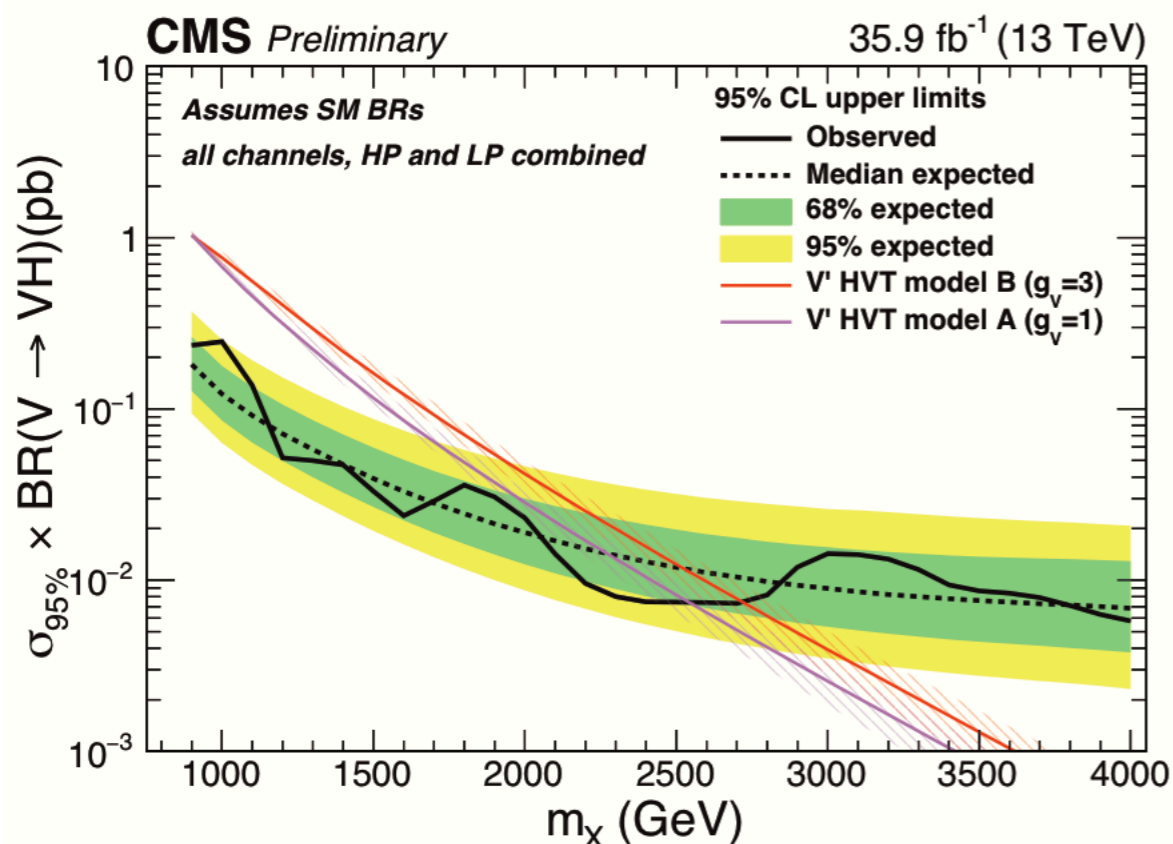
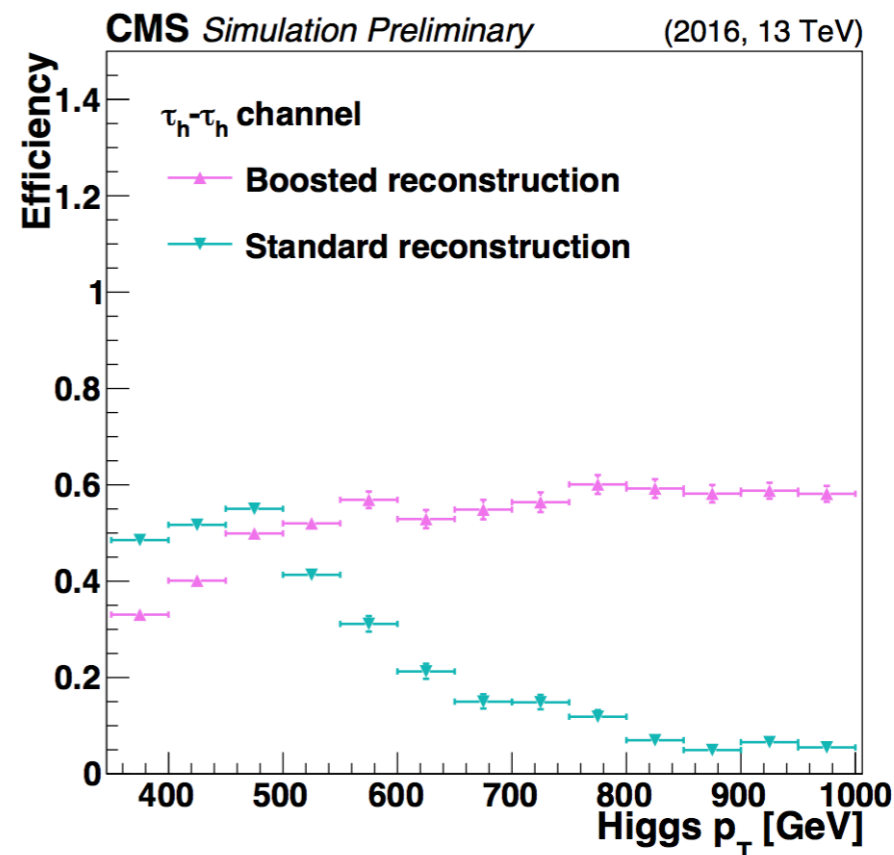
- > Dijet final state: Wh, Zh, h → bb
- > Using dedicated Higgs tagger
 - Loose and tight regions
- > Same background estimation strategy as for VV analysis



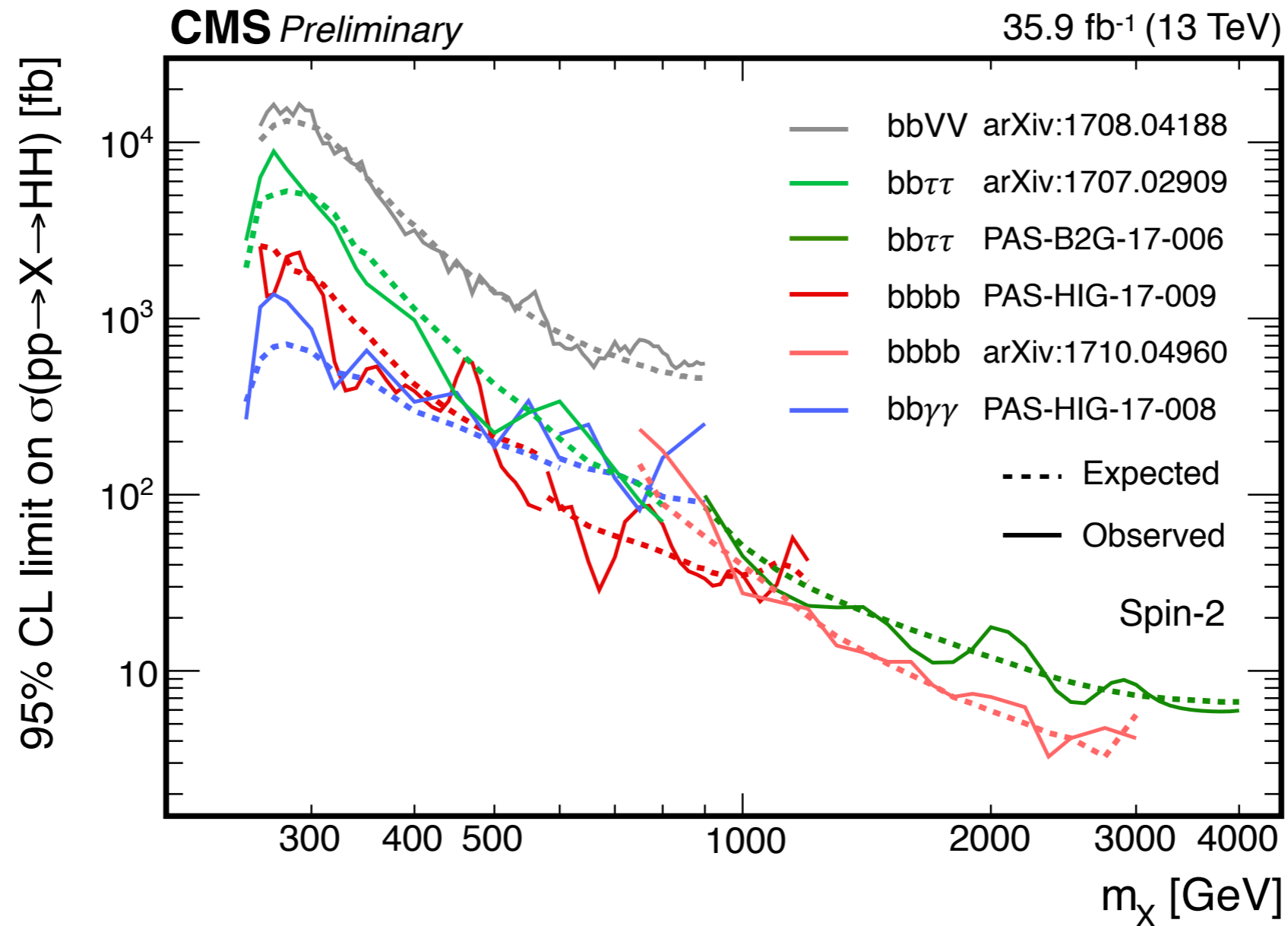
> Uses dedicated high- p_T $H \rightarrow \tau\tau$ reconstruction

- starting from fat jets, subjects reconstructed and then subject to τ reconstruction and identification algorithms

> Spin 0, 1, and 2 interpretation



hh spin-2 interpretation



bb $\nu\nu$ (HIG-17-006)

bbbb (HIG-17-009)

bb $\tau\tau$ (HIG-17-002, submitted to Phys. Lett. B)

bbbb (B2G-16-026, submitted to Phys. Lett. B)

bb $\tau\tau$ (B2G-17-006, submitted to J. High Energy Phys.)

bb $\gamma\gamma$ (HIG-17-008)