

Revisiting fine-tuning in the MSSM.

...or thoughts about the SUSY WIMP...

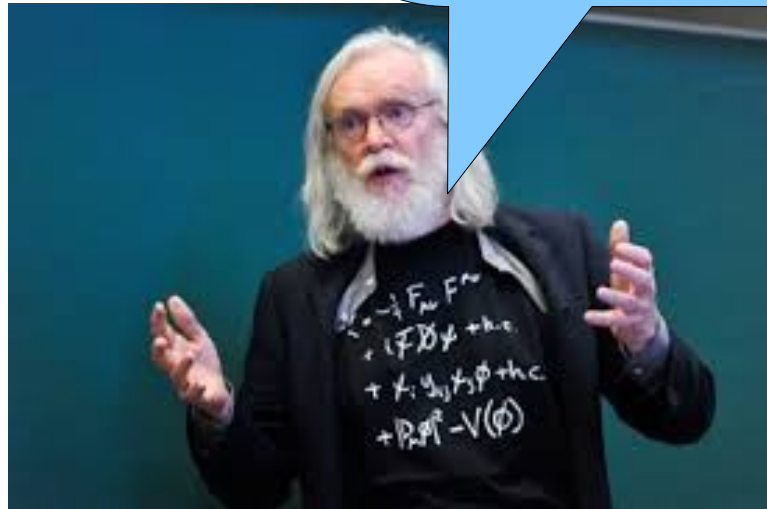
Kai Schmidt-Hoberg

“SUSY anywhere is better than SUSY nowhere!”

Largely based on

1603.09347
1701.03480

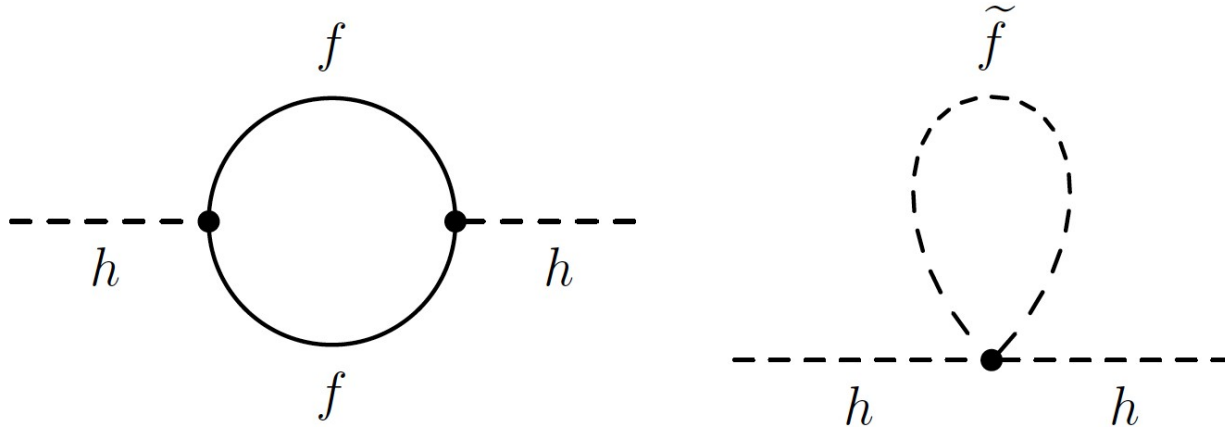
with G Ross and F Staub



Corfu Summer Institute

Why we like(d) SUSY

- > **Hierarchy problem:** stabilizes the weak against the Planck scale



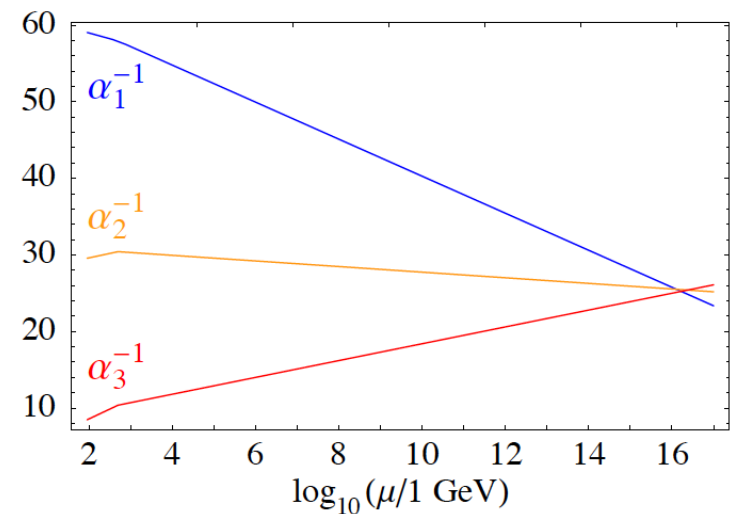
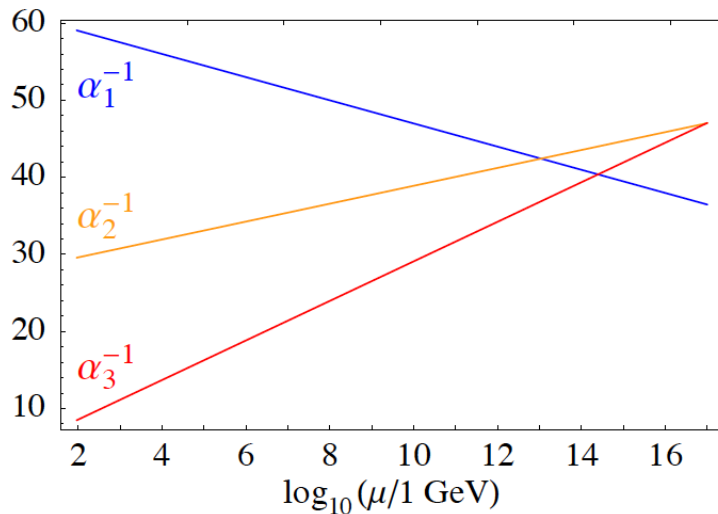
Why we like(d) SUSY

- > **Hierarchy problem:** stabilizes the weak against the Planck scale
- > **Dark matter:** If lightest SUSY particle stable \rightarrow dark matter candidate



Why we like(d) SUSY

- > Hierarchy problem: stabilizes the weak against the Planck scale
- > Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- > Gauge coupling unification:



Why we like(d) SUSY

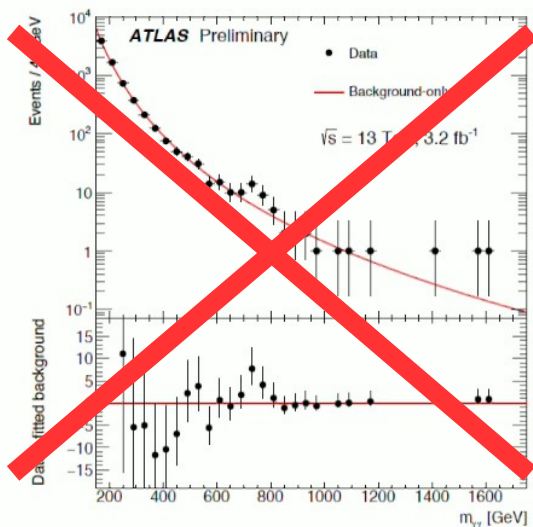
- > Hierarchy problem: stabilizes the weak against the Planck scale
- > Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- > Gauge coupling unification:
- > A 125 GeV Higgs boson: Additional hint for SUSY?

...somebody still owes me...



Why we like(d) SUSY

- Hierarchy problem: stabilizes the weak against the Planck scale
- Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- Gauge coupling unification:
- A 125 GeV Higgs boson: Additional hint for SUSY?
- Also hard to get 750 GeV diphoton excess ;-)



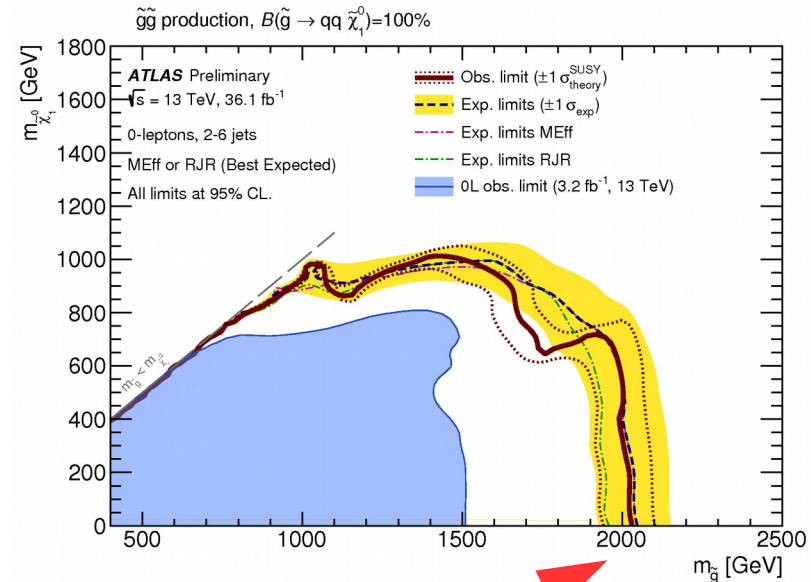
Congrats Jamie!



Why we like(d) SUSY

- > Hierarchy problem: stabilizes the weak against the Planck scale
- > Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- > Gauge coupling unification:
- > A 125 GeV Higgs boson: Additional hint for SUSY?

> So why do people get worried?



$\sim 2 \text{ TeV}$

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NATURALNESS!



DM naturalness in the MSSM

> How naturally can the dark matter relic abundance be achieved?

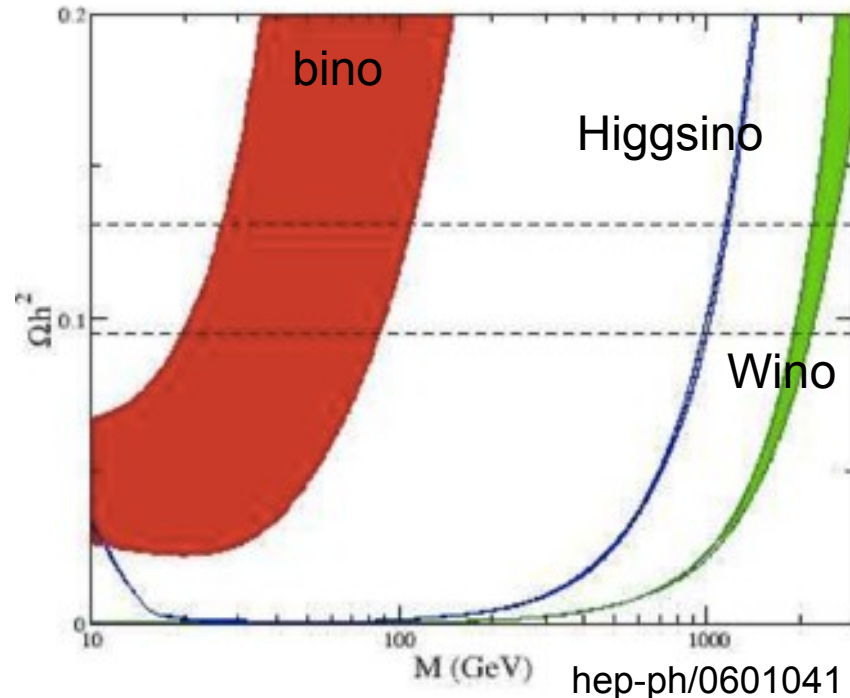
Often universal gaugino masses assumed at high scale, at low scale $M_3:M_2:M_1 \sim 6:2:1 \rightarrow$ bino LSP

Bino: Typically need to finely tune relic density via coannihilations or resonances :-)

Crucially depends on assumption of SUSY breaking terms!
Other patterns possible...

2-3 TeV Wino challenged by ID
Mariangela Lisanti et al [1307.4082](#)

1 TeV Higgsino looking good :-)



EW naturalness in the MSSM

- > How naturally can we achieve the correct Higgs vev?
- > Electroweak vev (or m_Z) determined by SUSY parameters (from minimization condition for scalar potential)

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 - m_{H_u}^2 \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \mu^2$$

- > Cancellation (tuning) needed for large SUSY masses
- > How to quantify this?

$$\Delta_p \equiv \frac{\partial \ln v^2}{\partial \ln p} = \frac{p}{v^2} \frac{\partial v^2}{\partial p} \quad \text{'sensitivity measure'}$$

- > Large Δ implies large tuning

The usual story

- > What does this tell us about a natural SUSY spectrum?
- > μ is a superpotential parameter and hardly runs: $\mu_{\text{EW}} \sim \mu_{\text{GUT}}$

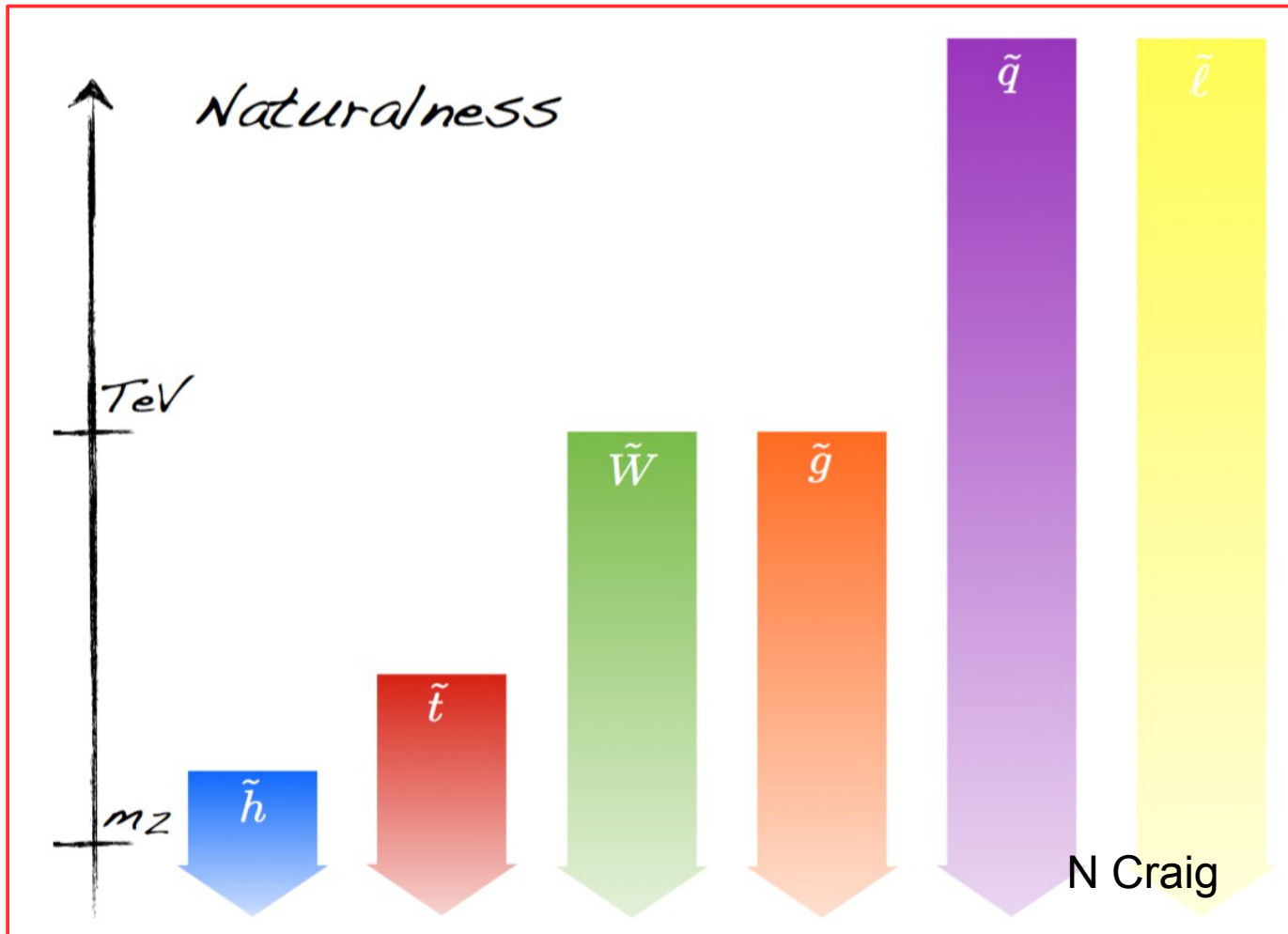
$$\Delta_\mu \sim \frac{2\mu^2}{M_Z^2} \quad \text{Higgsino mass} \sim \mu \sim 1 \text{ TeV} \rightarrow \quad \Delta \sim 250$$

- > “Natural SUSY requires light Higgsino”
- > What about the m_{H_u} part?
- > Loop effects introduce a large sensitivity to stop and gluino masses

$$\delta m_{H_u}^2 = -\frac{3y_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln(\Lambda/m_{\tilde{t}})$$

$$\delta m_{\tilde{t}}^2 = \frac{2g_s^2}{3\pi^2} m_{\tilde{g}}^2 \ln(\Lambda/m_{\tilde{g}})$$

The 'natural SUSY' spectrum



The 'natural SUSY' spectrum



EW naturalness in the MSSM – the GUT picture

- Starting from the high scale, all soft terms contribute to m_{H_u} and m_Z

$$\begin{aligned} m_Z^2 \simeq & -2.18\mu^2 + 3.84M_3^2 + 0.32M_3M_2 + 0.047M_1M_3 - 0.42M_2^2 \\ & + 0.011M_2M_1 - 0.012M_1^2 - 0.65M_3A_t - 0.15M_2A_t \\ & - 0.025M_1A_t + 0.22A_t^2 + 0.004M_3A_b \\ & - 1.27m_{H_u}^2 - 0.053m_{H_d}^2 \\ & + 0.73m_{Q_3}^2 + 0.57m_{U_3}^2 + 0.049m_{D_3}^2 - 0.052m_{L_3}^2 + 0.053m_{E_3}^2 \\ & + 0.051m_{Q_2}^2 - 0.11m_{U_2}^2 + 0.051m_{D_2}^2 - 0.052m_{L_2}^2 + 0.053m_{E_2}^2 \\ & + 0.051m_{Q_1}^2 - 0.11m_{U_1}^2 + 0.051m_{D_1}^2 - 0.052m_{L_1}^2 + 0.053m_{E_1}^2, \end{aligned}$$

- We don't just want m_{H_u} to be small, but every contribution to it. Assuming no correlations among the terms, need rather light stops and gluinos

EW naturalness in the MSSM – the GUT picture

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$$\begin{aligned} & -1.27m_{H_u}^2 - 0.053m_{H_d}^2 \\ & + 0.73m_{Q_3}^2 + 0.57m_{U_3}^2 + 0.049m_{D_3}^2 - 0.052m_{L_3}^2 + 0.053m_{E_3}^2 \\ & + 0.051m_{Q_2}^2 - 0.11m_{U_2}^2 + 0.051m_{D_2}^2 - 0.052m_{L_2}^2 + 0.053m_{E_2}^2 \\ & + 0.051m_{Q_1}^2 - 0.11m_{U_1}^2 + 0.051m_{D_1}^2 - 0.052m_{L_1}^2 + 0.053m_{E_1}^2 \end{aligned}$$

$$m_i^2 = m_0^2$$

$$\sim 0.01 m_0^2$$

- We don't just want m_{H_u} to be small, but every contribution to it. Assuming no correlations among the terms, need rather light stops and gluinos
- But we know correlations should be present...
- Example: the scalar focus point.

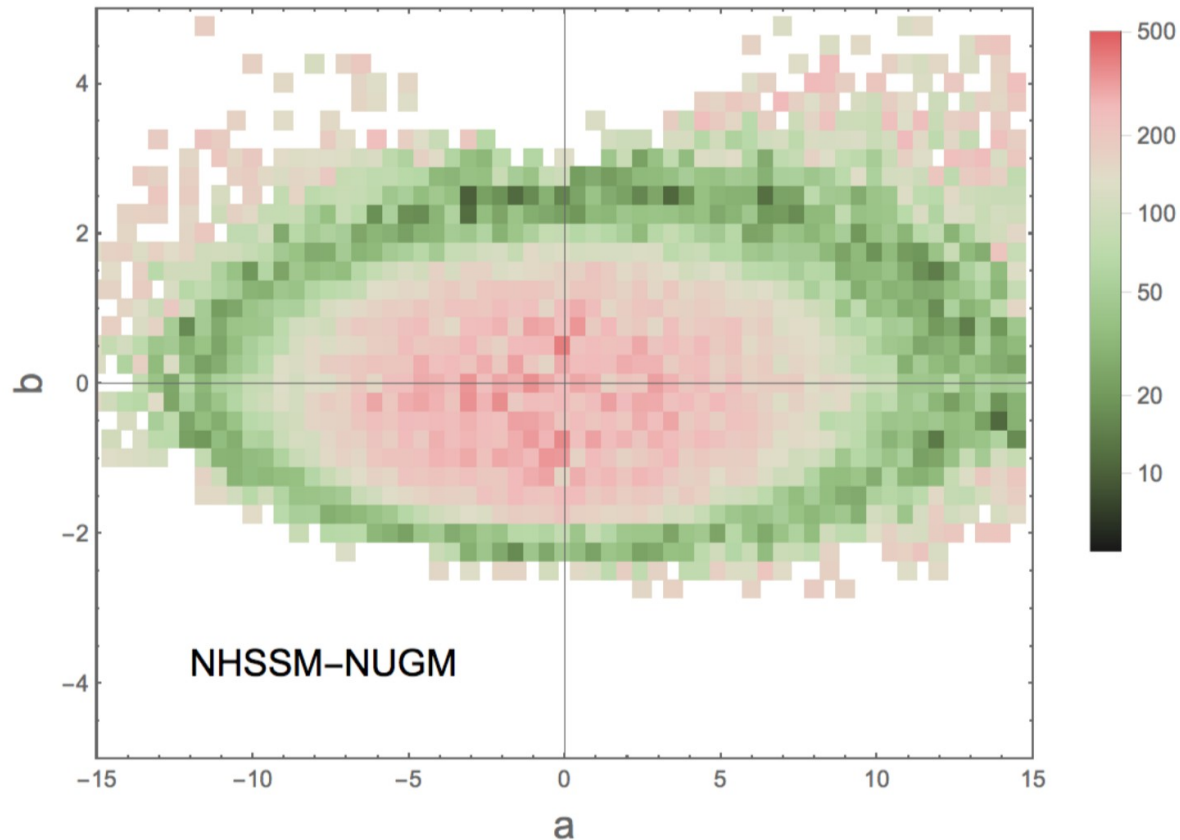
The gaugino focus point

- > Assume fixed ratios of gaugino masses
- > Possible also in GUTs Horton, Ross, 0908.0857

$$M_1 = a \cdot m_{1/2}$$

$$M_2 = b \cdot m_{1/2}$$

$$M_3 = m_{1/2}$$



Comment on loop corrections

> So far assumed tree-level relation for EWSB condition

Slide from F Staub

$$\left. \frac{\partial V^{(L)}}{\partial v_u} \right|_{\tan\beta \rightarrow \infty} \equiv 0 = (m_{h_u}^2 + \mu^2 + \frac{1}{8}(g_1^2 + g_2^2)v^2)v + \Sigma_u$$

How to parametrise Σ_u ?

① $\Sigma_u \equiv v\Sigma_{uu}$

$$\rightarrow \frac{1}{2}M_Z^2 = -|\mu|^2 - m_{H_u}^2 + \Sigma_{uu}$$

no change in FT; only valid if Σ_{uu} is independent of v !

② $\Sigma_u \equiv \Sigma_1 v + \Sigma_2 v^2 + \Sigma_3 v^3$

$$\rightarrow \Delta\mu = \frac{8\mu^2}{(g_1^2 + g_2^2 + 8\Sigma_3)v^2 + 4\Sigma_2 v}$$

Comment on loop corrections

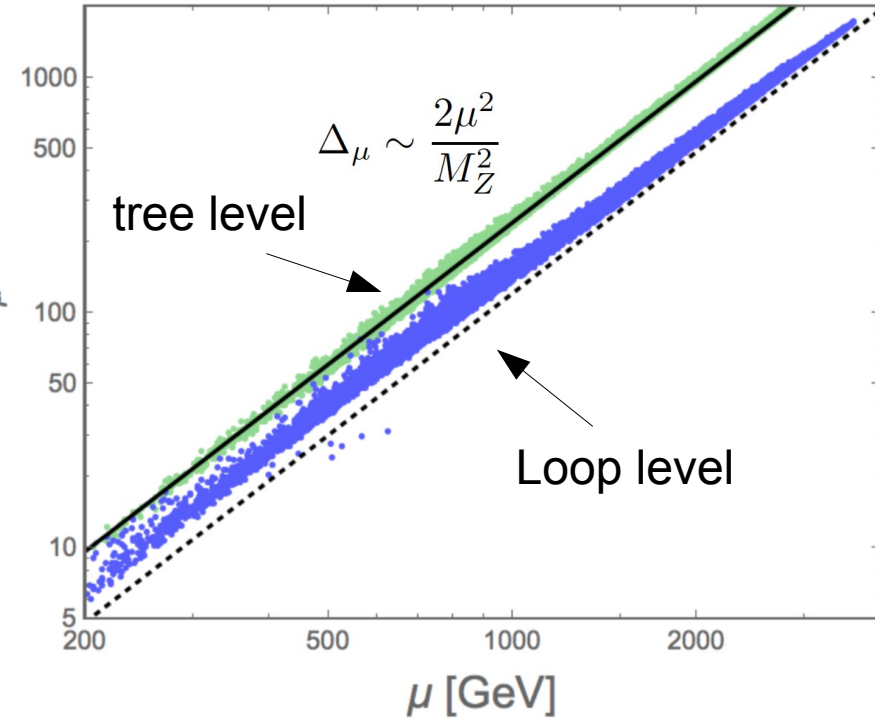
- > Result for leading stop correction

$$\Sigma_1 = -\frac{3m_t^2 Y_t^2 \left(-2\log\left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2}\right) + 2 + \log(4) \right)}{32\pi^2} \Delta_\mu$$

$$\Sigma_2 = 0$$

$$\Sigma_3 = \frac{3Y_t^4 \left(\log\left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2}\right) - \log\left(\frac{v^2 Y_t^2}{Q^2}\right) \right)}{32\pi^2} \neq 0$$

$$r^{FT} \equiv \frac{\Delta_\mu^{\text{Loop}}}{\Delta_\mu^{\text{Tree}}} \simeq \left(1 + \frac{3}{4\pi^2} \frac{Y_t^4}{g_1^2 + g_2^2} \log\left(\frac{m_t^2}{m_t^2}\right) \right)^{-1} \simeq \frac{1}{2}$$



See also Ghilencea, Ross +
0903.1115, 1001.3884

- > Reduction of about 1/2 when including loop corrections

Comment on loop corrections

- > Result for leading stop correction

$$3m_t^2 Y_t^2 \left(-2 \log \left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2} \right) + 2 + \log(4) \right)$$

$$\Sigma_1 = -$$

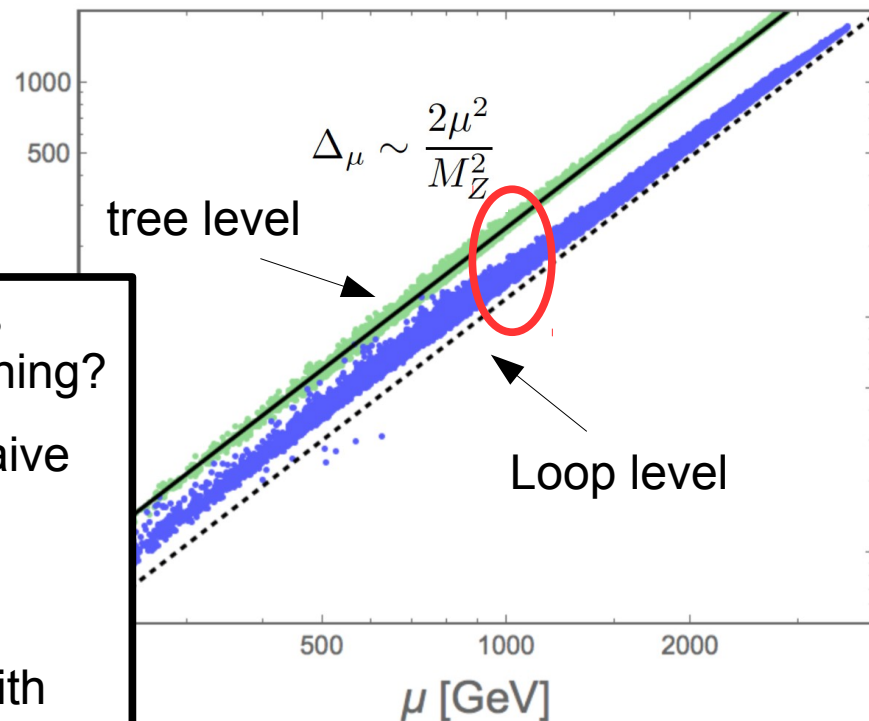
$$\Sigma_2 = 0$$

$$3Y$$

$$\Sigma_3 = -$$

$$r^{FT} \equiv \frac{\Delta I}{\Delta A}$$

- > So does a 1 TeV Higgsino always require significant electroweak tuning?
- > Two possible ways beyond the naive FT arguments
- > The μ -term could be correlated with the soft terms ($\rightarrow \mu$ problem)
- > The Higgsino mass could arise from something beyond the μ -term



See also Ghilencea, Ross +
0903.1115, 1001.3884

- > Reduction of about $\frac{1}{2}$ when including loop corrections

A new contribution to the Higgsino mass

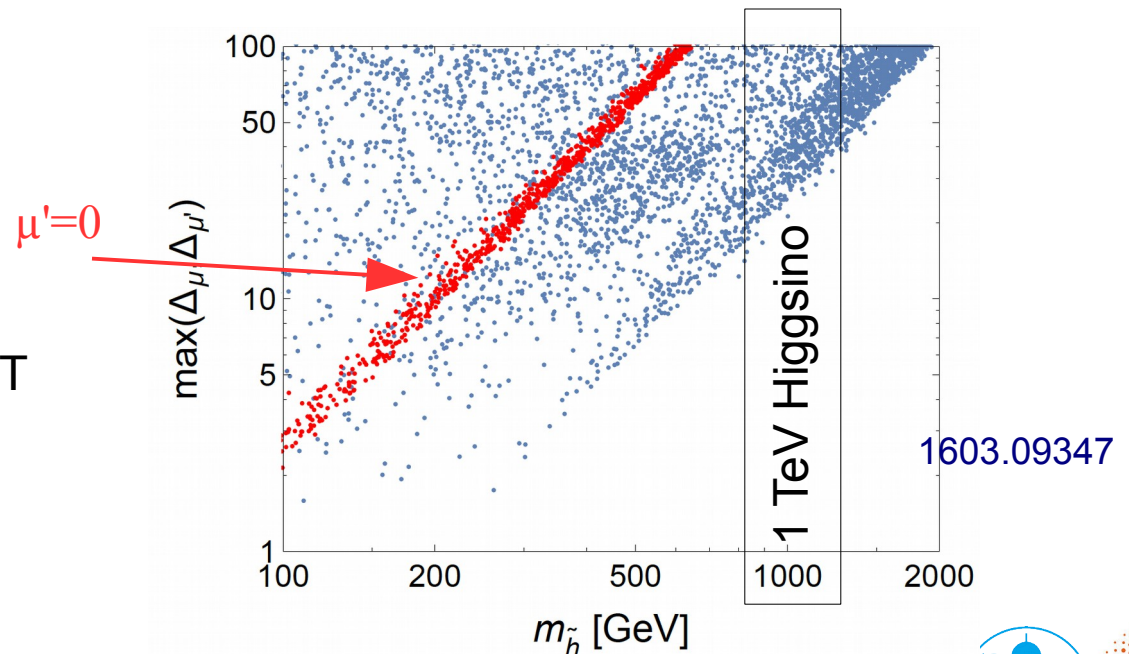
- > Non-standard SUSY breaking terms (in the classification of S Martin: 'maybe-soft')

$$\mathcal{L}_{NH} = \mu' \tilde{h}_d \tilde{h}_u + T'_{u,ij} h_d^* \tilde{u}_{R,i}^* \tilde{q}_j + T'_{d,ij} h_u^* \tilde{d}_{R,i}^* \tilde{q}_j + T'_{e,ij} h_u^* \tilde{e}_{R,i}^* \tilde{l}_j + \text{h.c.}$$

Girardello, Grisaru (1982)

- > μ' contributes to the Higgsino mass ($m_{\tilde{h}} \sim \mu + \mu'$) but does not enter the scalar potential

- > Can significantly reduce FT



Embedding this into a model

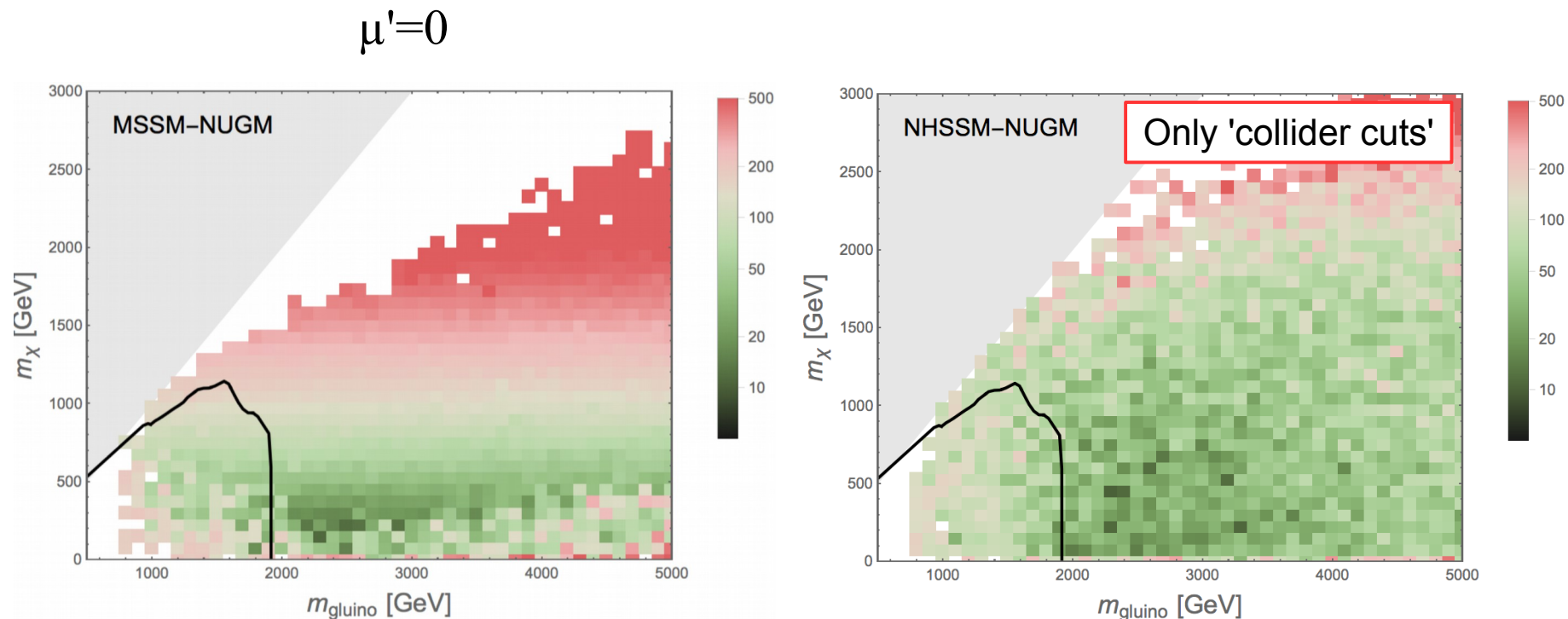
- Studied different MSSM variants with GUT boundary conditions

	$m_{h_u}^2$	$m_{h_d}^2$	M_1	M_2	M_3	μ'	A'_0
CMSSM	m_0^2	m_0^2	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	-	-
MSSM-NUHM	$m_{h_u}^2$	$m_{h_d}^2$	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	-	-
MSSM-NUGM	m_0^2	m_0^2	$a \cdot m_{1/2}$	$b \cdot m_{1/2}$	$m_{1/2}$	-	-
CNHSSM	m_0^2	m_0^2	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	μ'	A'_0
NHSSM-NUHM	$m_{h_u}^2$	$m_{h_d}^2$	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	μ'	A'_0
NHSSM-NUGM	m_0^2	m_0^2	$a \cdot m_{1/2}$	$b \cdot m_{1/2}$	$m_{1/2}$	μ'	A'_0

1701.03480

Results non-universal gaugino masses

- Region of small FT can be well beyond LHC reach



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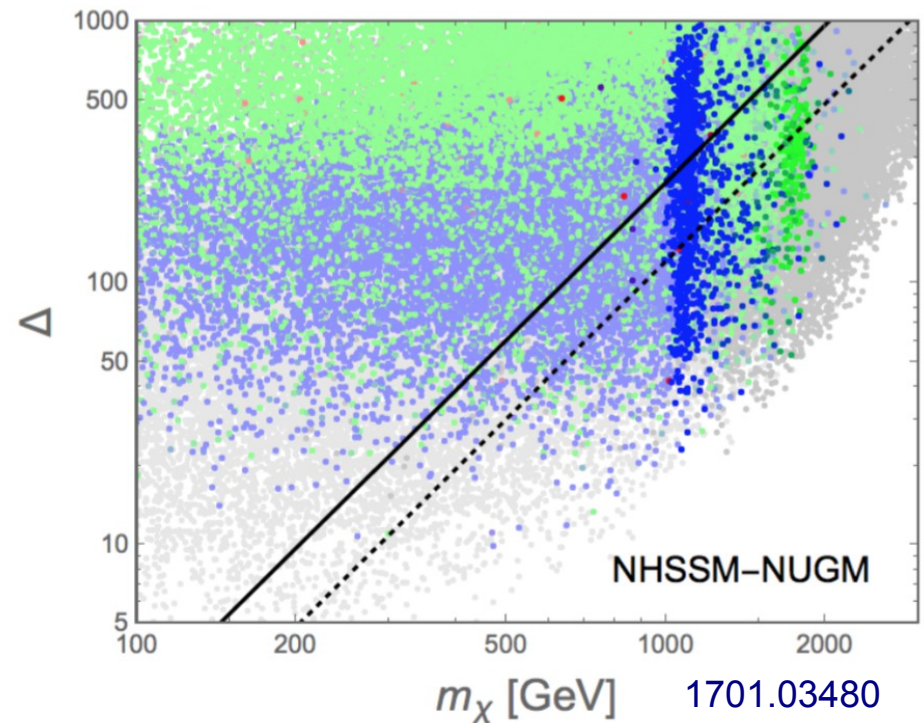
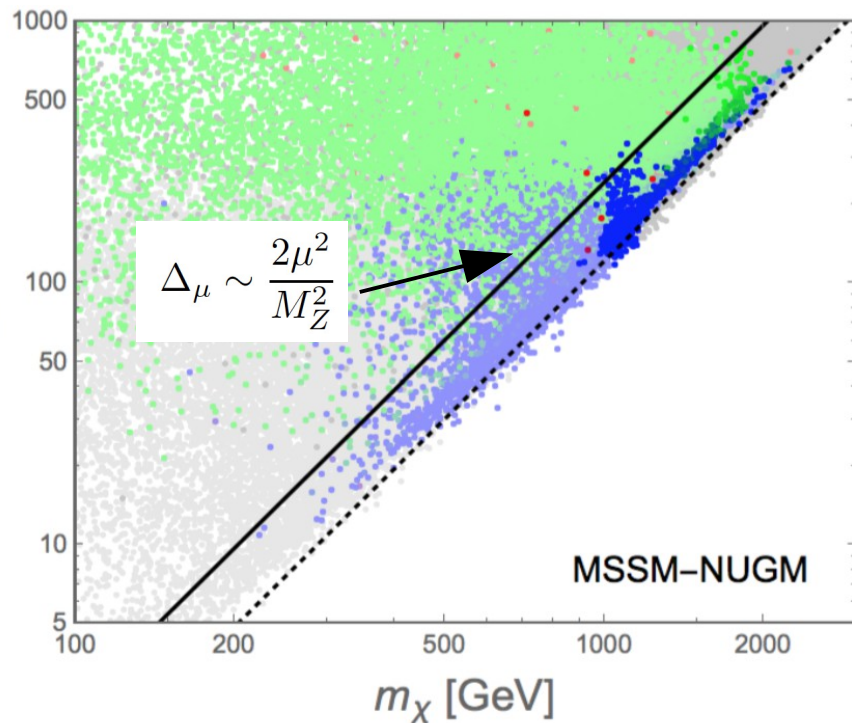
- Allowing for DM underabundance FT can be as small as 10.

Results non-universal gaugino masses

> A 1 TeV Higgsino can be quite natural

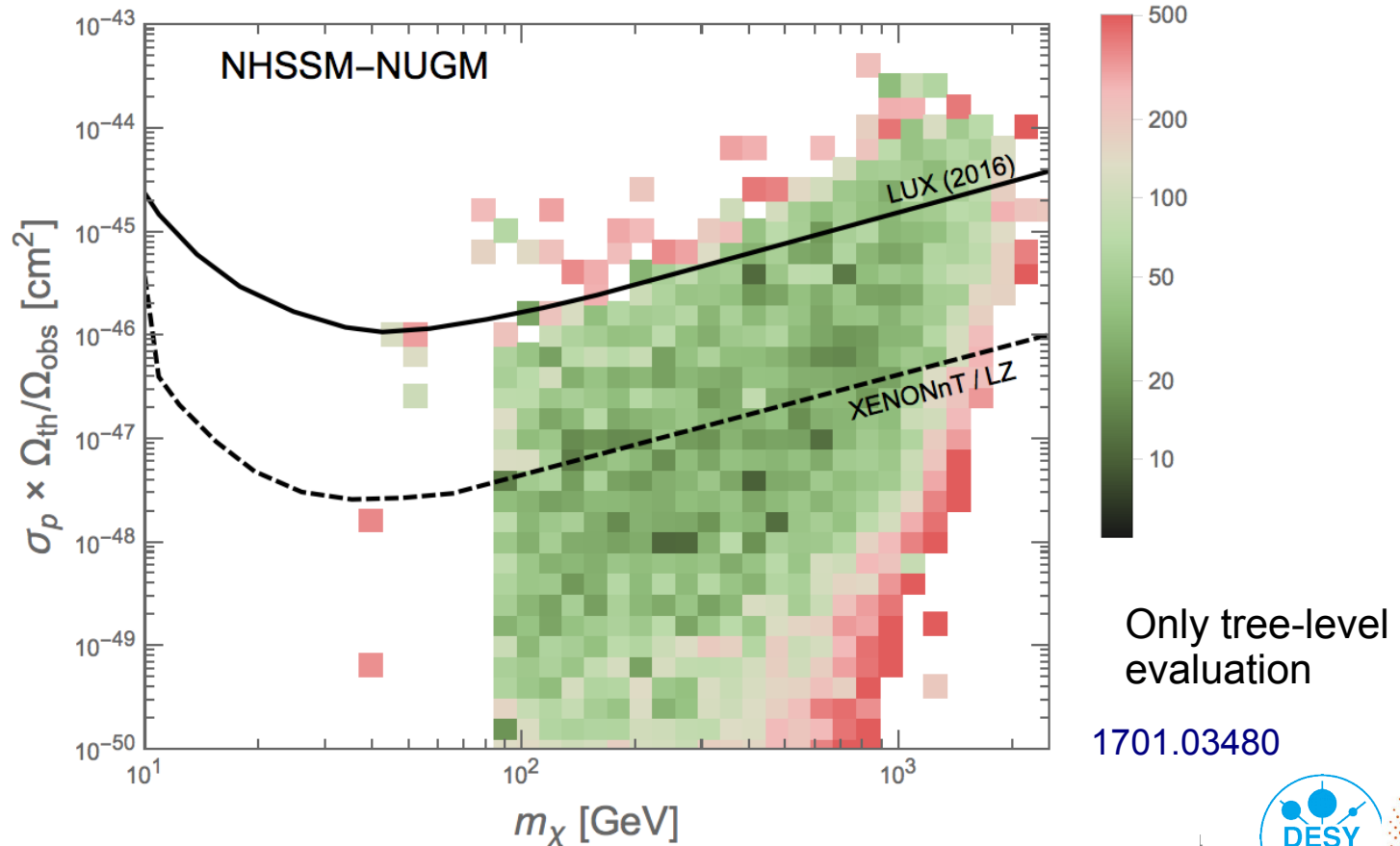
$$\mu' = 0$$

- bino
- Higgsino
- Wino



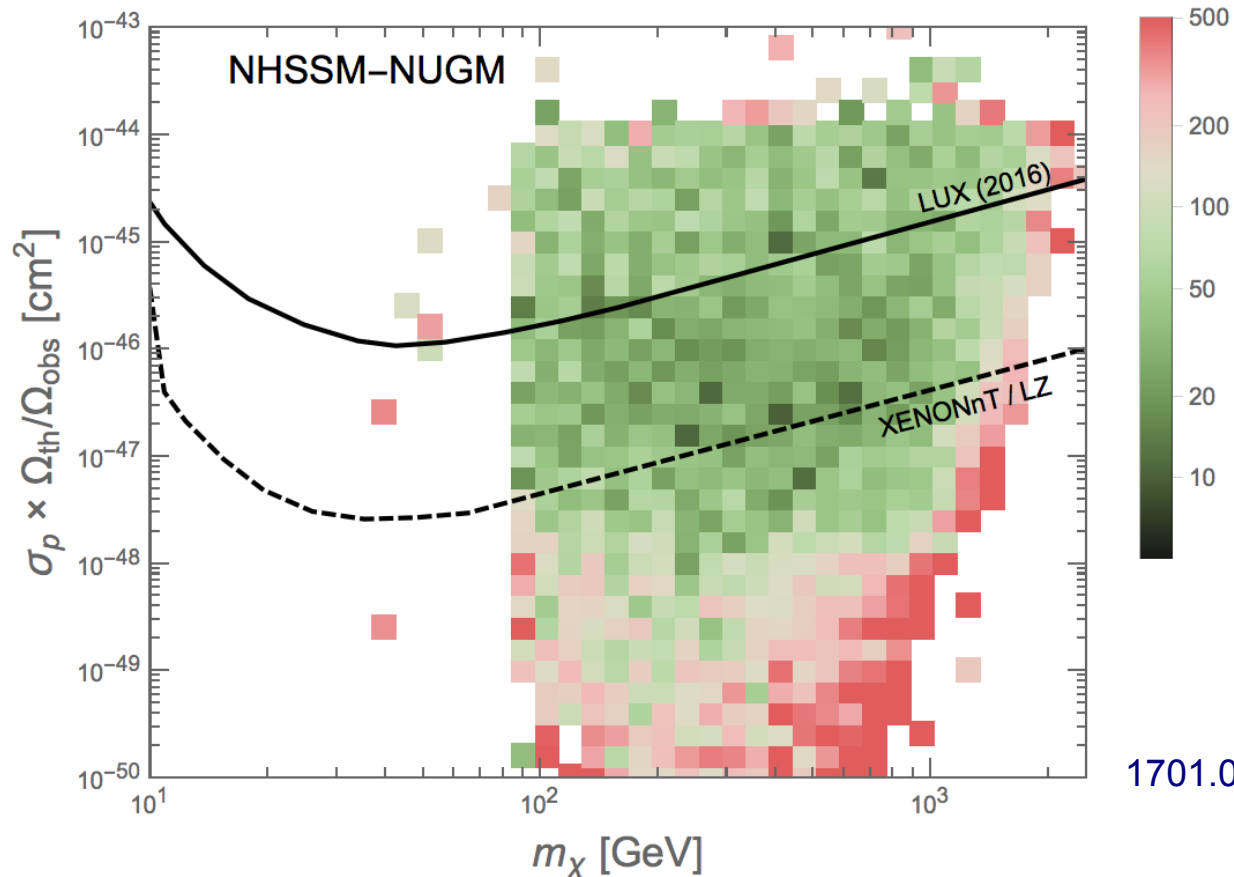
Prospects for direct detection

- Prospects for direct detection
- No lower bound on relic abundance (and rescaled) – other DM component



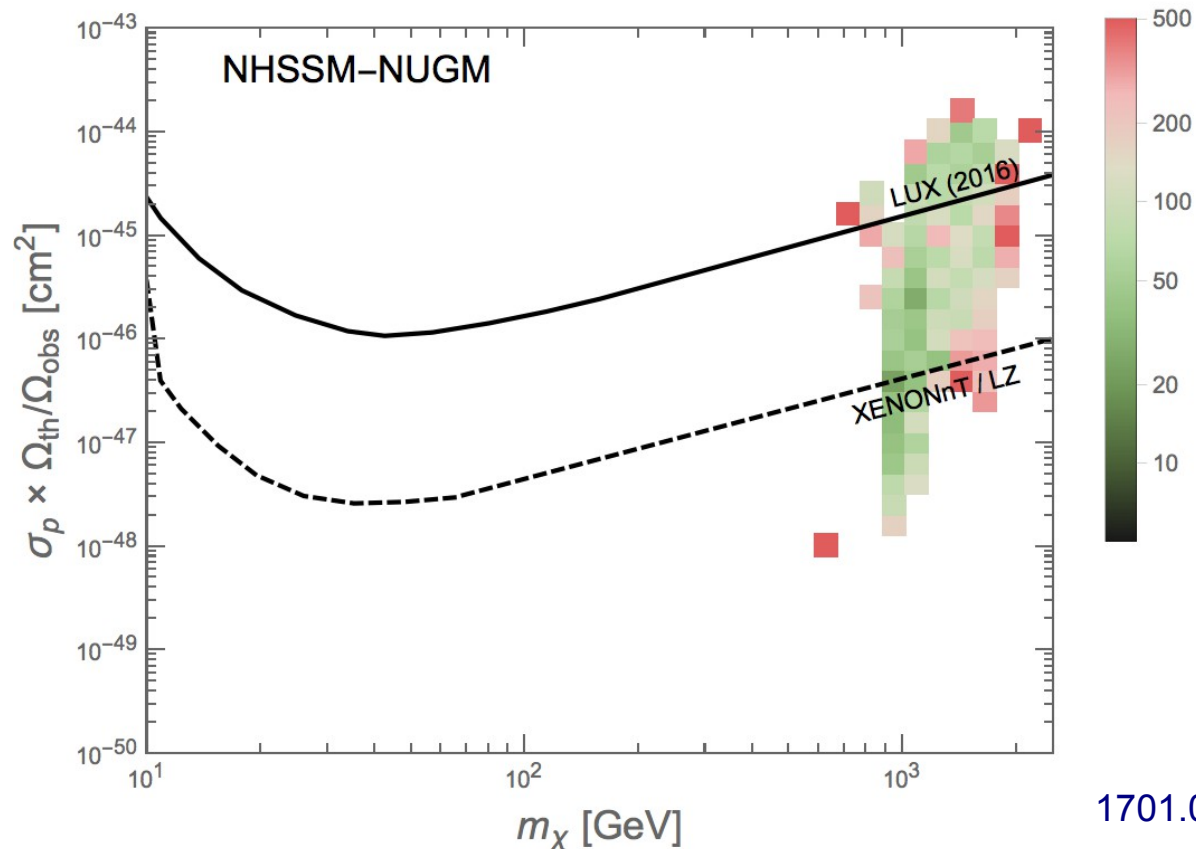
Prospects for direct detection

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- No lower bound on relic abundance (not rescaled) – non-thermal production (gravitino decay)



Prospects for direct detection

- Prospects for direct detection
- Correct (thermal) relic abundance



1701.03480

Summary

- What looks unnatural from an IR perspective **might** still look natural from the UV
- Extra Higgsino mass contribution μ' could help
- To do: build a UV model



- SUSY could well be beyond the LHC reach
- Good chances at direct detection experiments to find it

Summary

- > What looks unnatural from an IR perspective **might** still look natural from the UV
- > Extra Higgsino mass contribution μ' could help
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Thank you!



- > SUSY could well be beyond the LHC reach
- > Good chances at direct detection experiments to find it