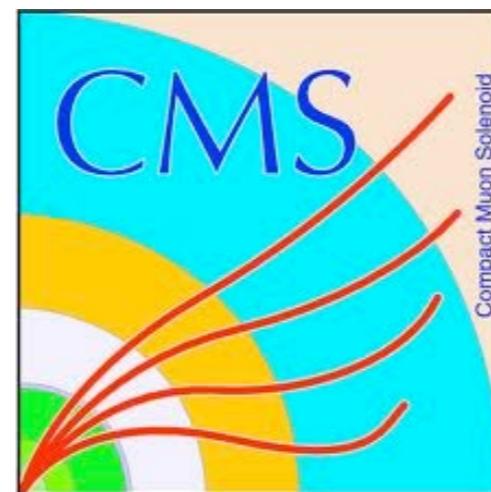


Andrea Castro
- University of Bologna and INFN -



Top Quark Mass



On behalf of the CMS collaboration

The top quark mass at LHC

22 years after its discovery at Tevatron we are still measuring the top quark mass:

- exploiting different channels (for completeness)
- recurring to different techniques (with different systematics)
- exploring its connection to the theory

Taking advantage of:

- **LHC is a top factory**
 - \approx 5 million pairs per experiment in 2012, \approx 30 million in 2016, each t decays \approx 100% to W+b
 - single top EWK production ($\sigma_t \sim \sigma_{tt}/3$)
- **Physics objects**
 - isolated energetic e or μ
 - energetic jets
 - b-tagged jets
 - momentum imbalance (MET)
 - *boosted top jets*

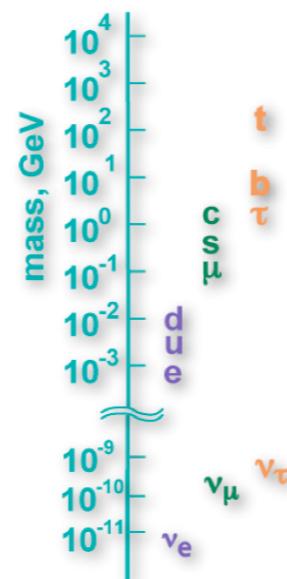
Why measure M_t ?

- 1) free parameter of SM
 - t decays well before hadronizing \Rightarrow measure M_t directly from decay products
 - we usually compare to Monte Carlo expectations, so what we really measure is $M_{t\text{MC}}$ parameter
 - there are ‘standard methods’ and ‘alternative methods’ (based on specific features)
 - this is complemented by a pole mass measurement

Why measure M_t ?

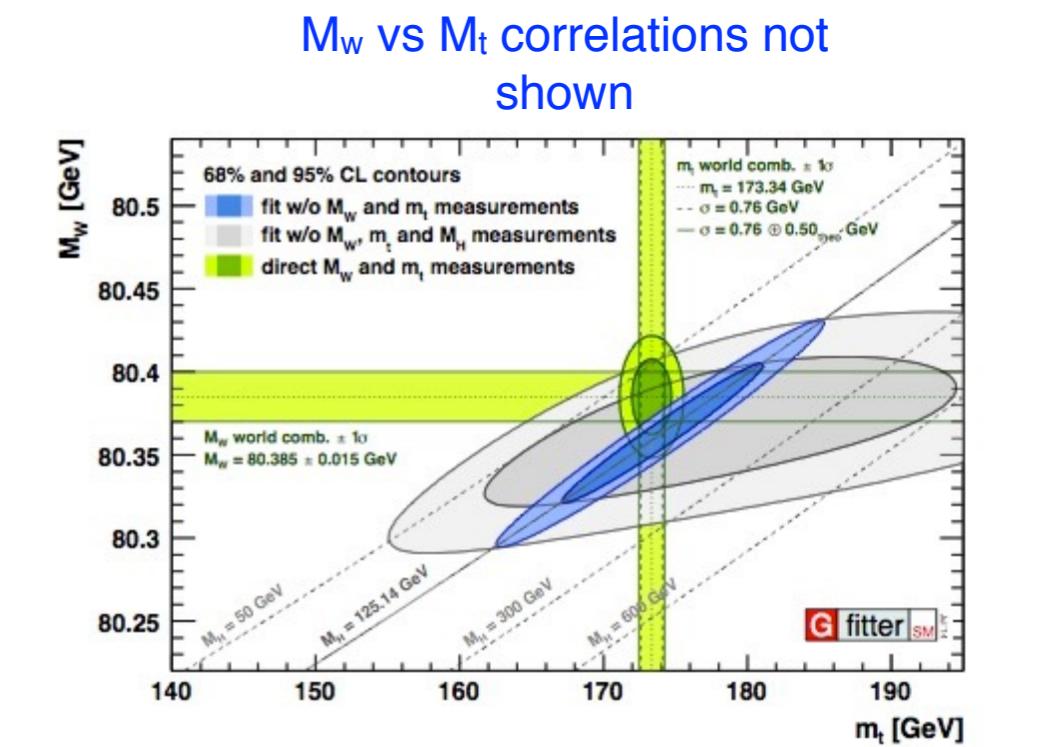
- 2) Participates in quantum loop radiative corrections to M_W constraining M_H
 ⇒ assessment of self-consistency within SM

- 3) M_t is close to scale of EWSB, so t might play a special role in it

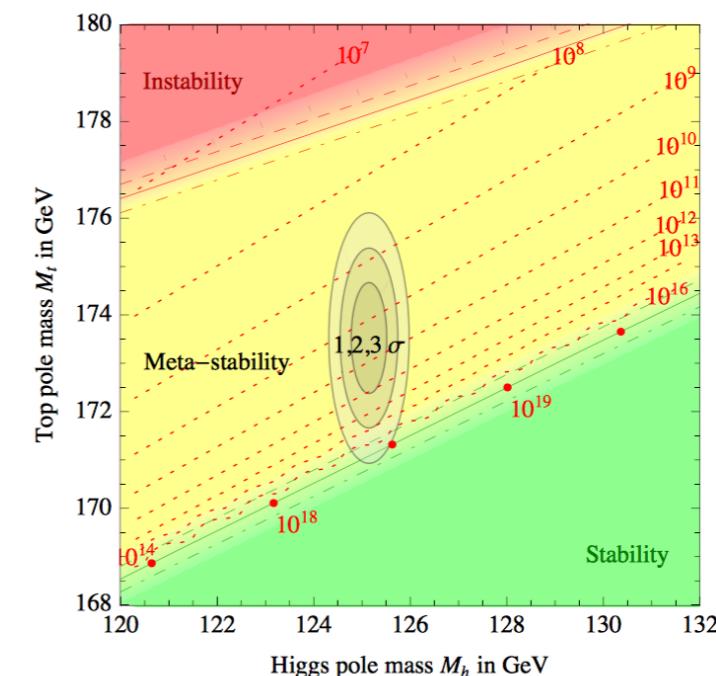


- 4) M_t related with M_H and vacuum stability of SM (and of Universe): near criticality of M_H

[arXiv:1307.3536](https://arxiv.org/abs/1307.3536)



[EPJC 74 \(2014\) 3046, arXiv:1407.3792](https://doi.org/10.1140/epjc/s10050-014-3046-4)



Latest results on M_t

There are several M_t measurements made so far by CMS at 7 and 8 TeV

We will discuss here only the most recent ones at 8 TeV (19.7 fb^{-1}):

- single top $\mu+\text{jets}$
- boosted top
- pole mass
- “alternative” methods

Systematic uncertainties

Statistical uncertainties becoming smaller and smaller
⇒ systematic uncertainties become dominant

Different sources of systematics, related to:

- Experimental effects
- Signal modeling
- Background modeling
- Features of the method

For every source, measurements performed (usually with pseudo-experiments) with modified parameters:
change of M_t ⇒ syst. uncertainty



Single top: $\mu + \text{jets}$

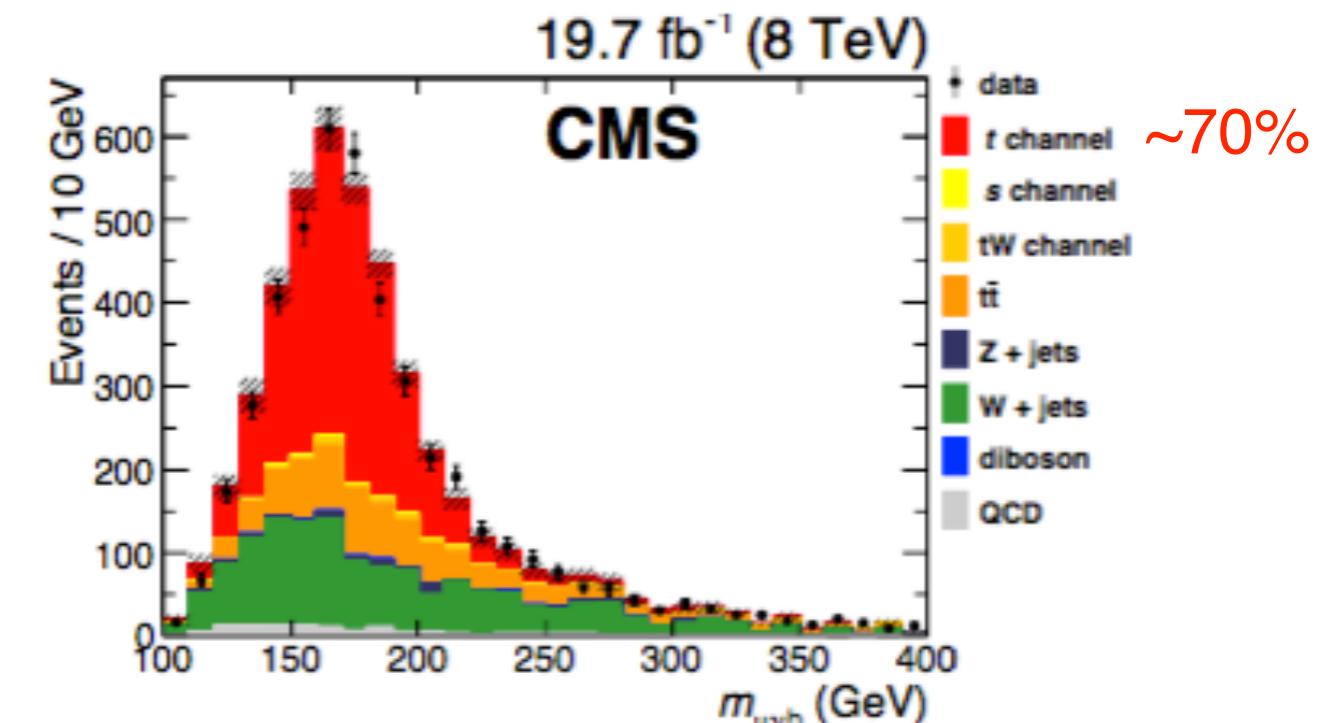
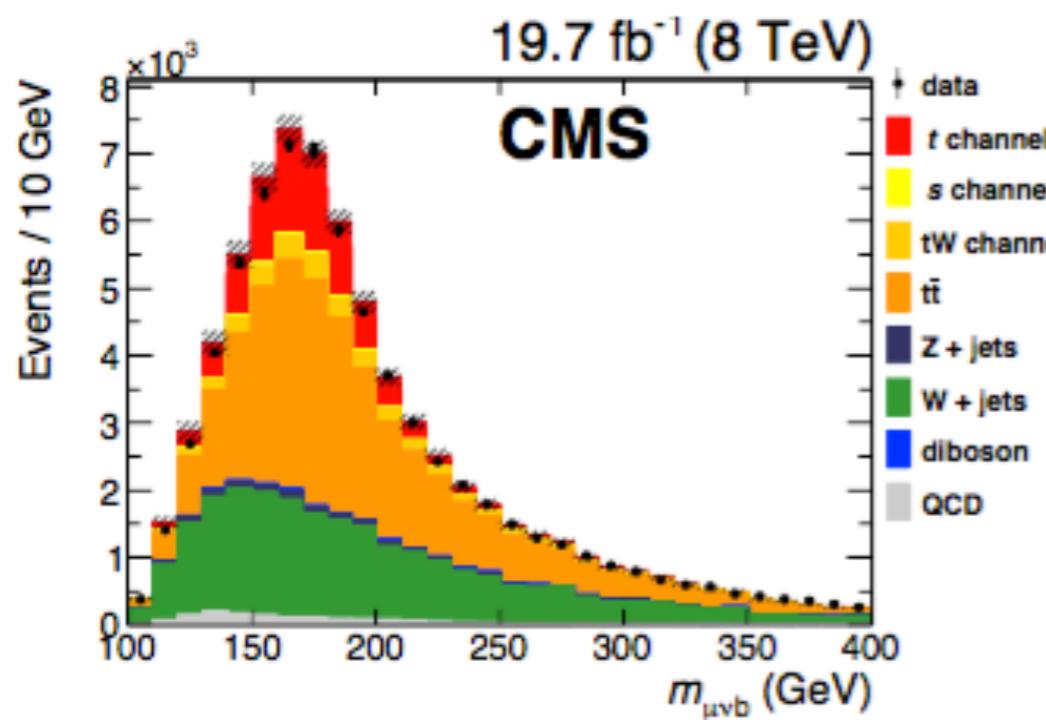
Template method ($m_{\mu\text{vb}}$)

[EPJC 77 \(2017\) 354](#)

Selection:

- 1 μ
- MET
- 2 jets (1 b-tagged)

- $m_{\mu\text{vb}}$ distributions as a proxy for m_t^{reco}
- μ charge >0 to improve S/B since $\sigma_t \sim 2 \sigma_{t\bar{t}}$
- light jet at large η expected from single top



$q_\mu > 0 \text{ & } |\eta_j| > 2.5 \Rightarrow$



Single top: $\mu + \text{jets}$

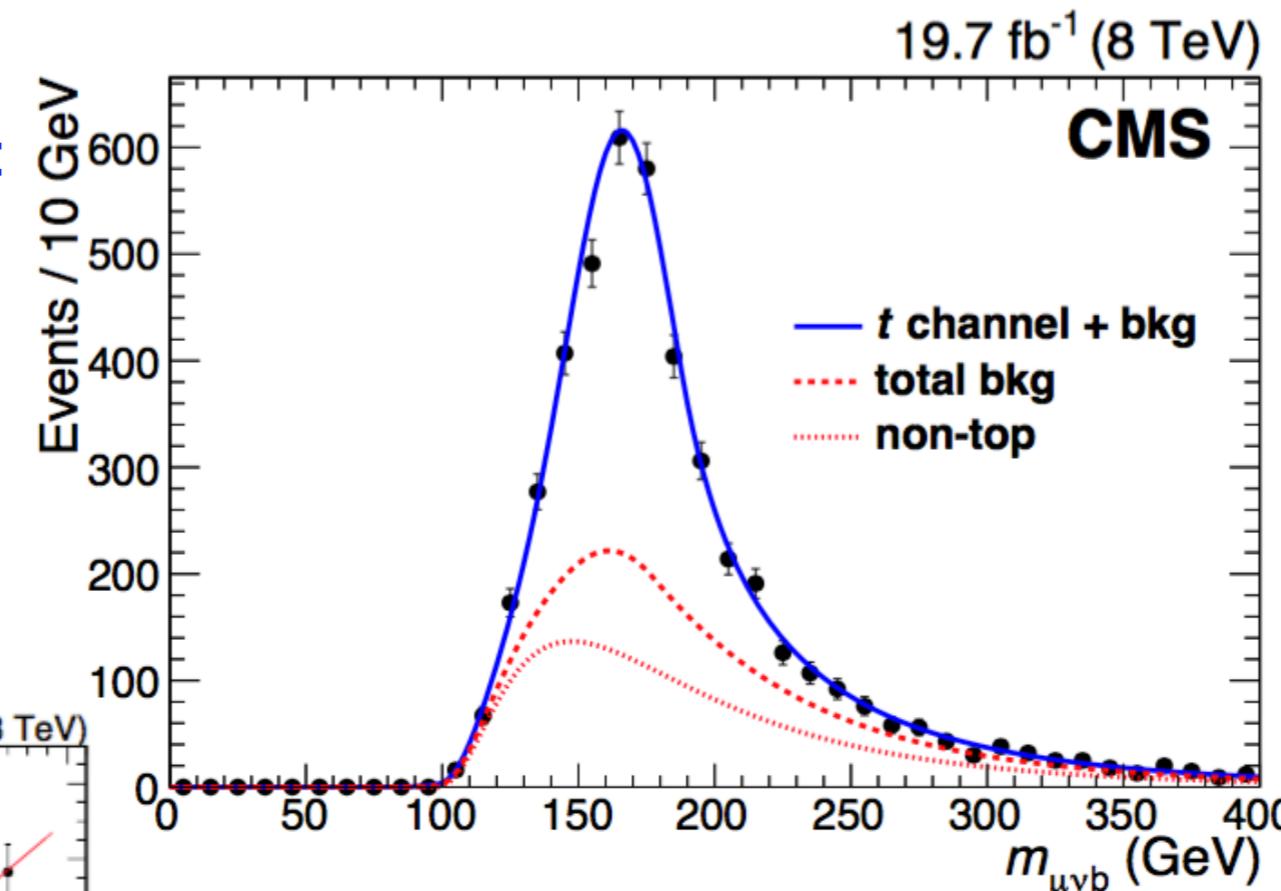
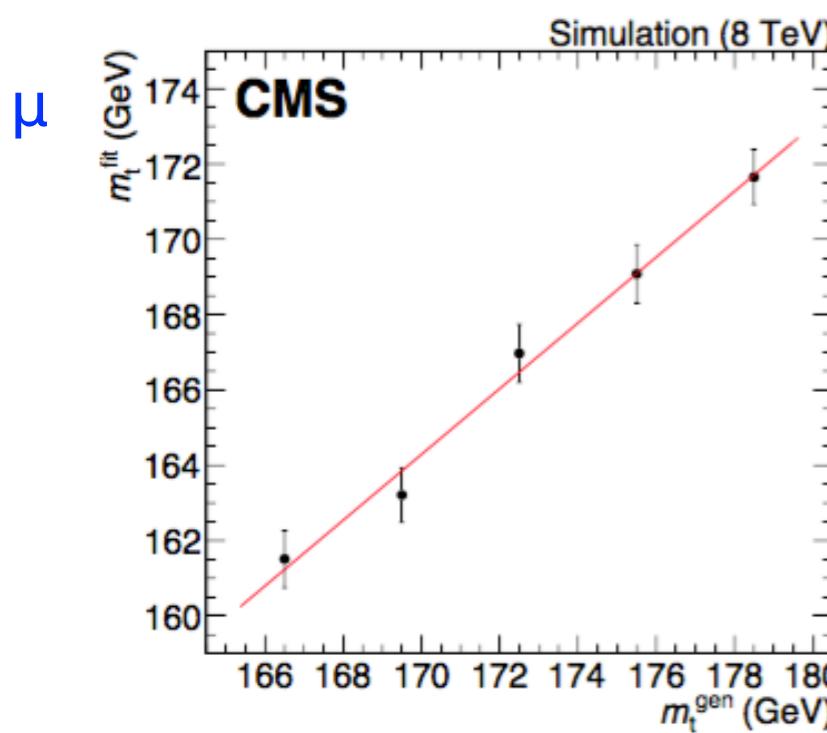
Template method ($m_{\mu\text{vb}}$)

$m_{\mu\text{vb}}$ parametrization

Gaussian core (μ, σ) + tails:

- t: Crystal Ball
- tt: Crystal Ball
- non-t: Novosibirsk

Calibration for the mass parameter μ



Main systematics

source	GeV
JES	0.68
bkgd	0.39
fit calibration	0.39

Fit

$M_t = 172.95 \pm 0.77(\text{stat}) + 0.97 - 0.93(\text{syst}) \text{ GeV}$

$M_t = 172.95 \pm 1.24 \text{ GeV} \quad (\pm 0.72\%)$



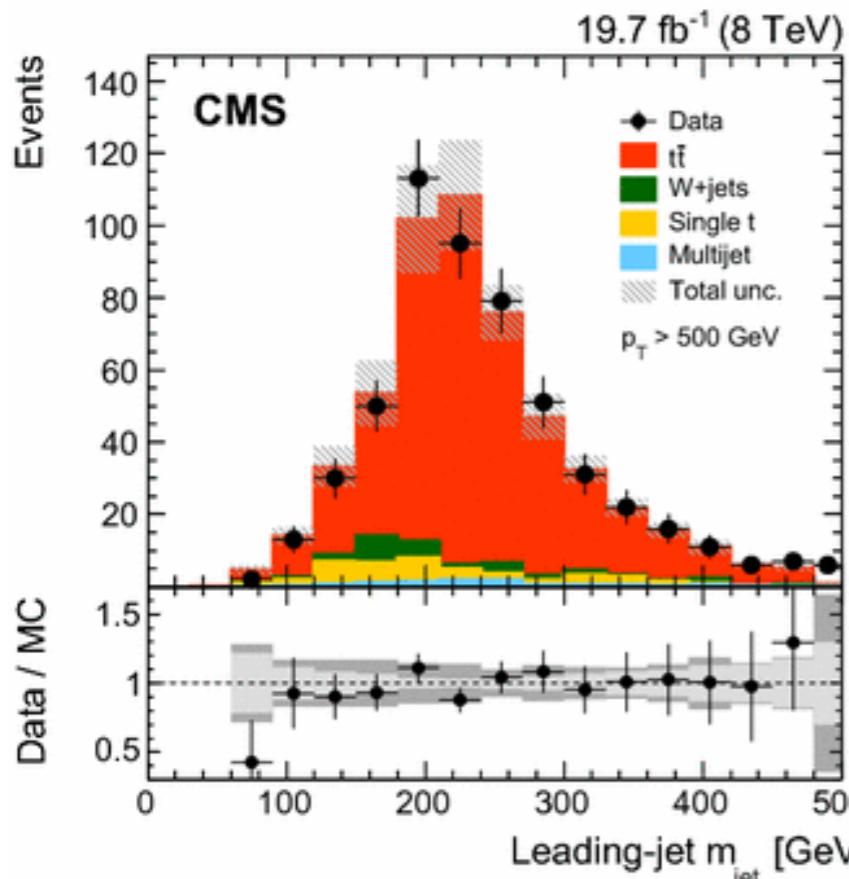
Boosted top: e/ μ + jets

Template method (m_{jet})

Selection:

- 1 e or μ
- ≥ 2 narrow jets (≥ 1 b-tagged)
- ≥ 2 wide jets
- MET

The M_t -sensitive quantity is the leading-jet mass

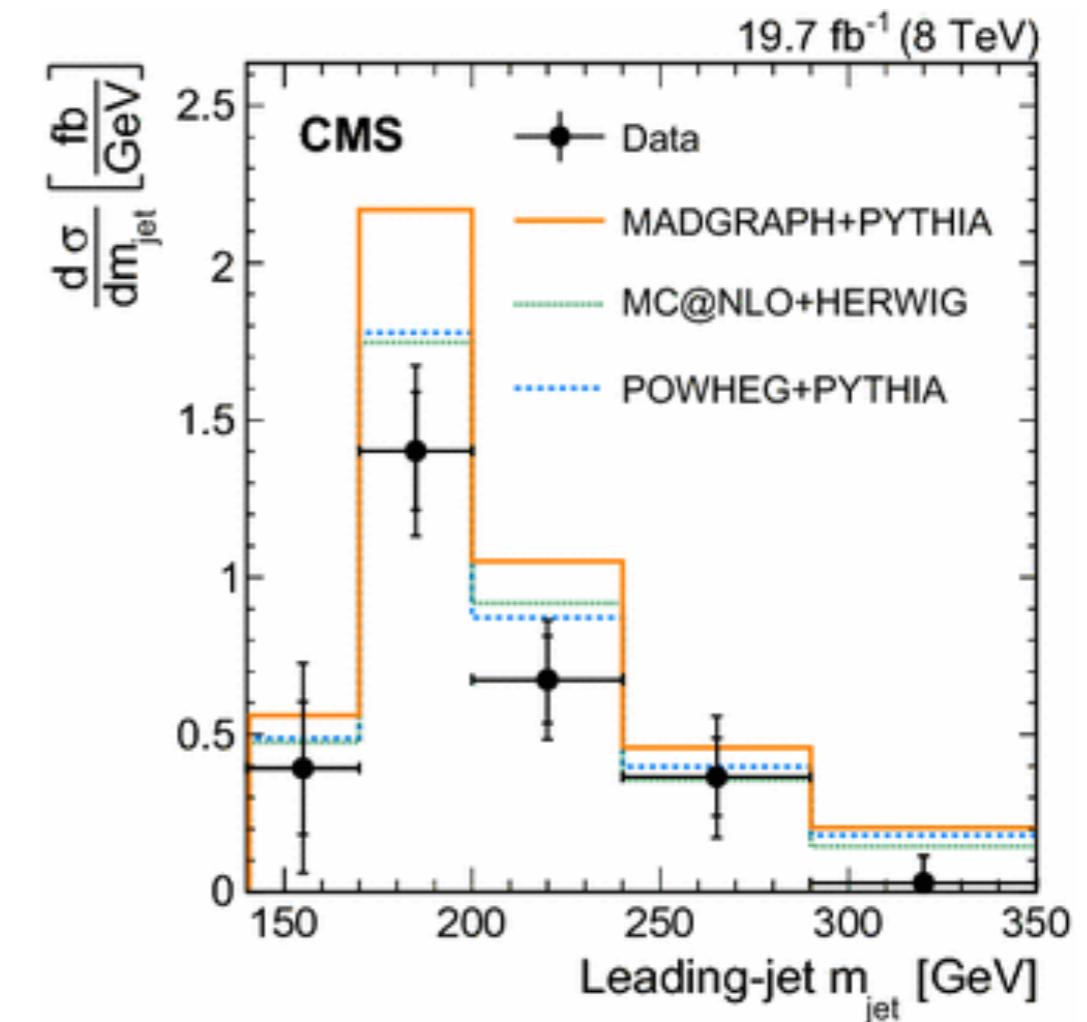


$p_T > 500$ GeV

[EPJC 77 \(2017\) 467](#)

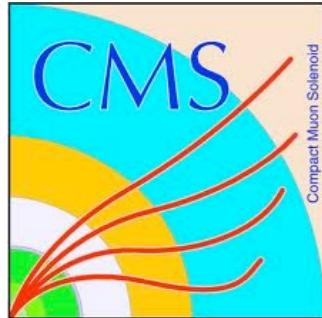
Distributions are translated (unfolded) from reconstruction to particle-level (fiducial)

$d\sigma/dM$



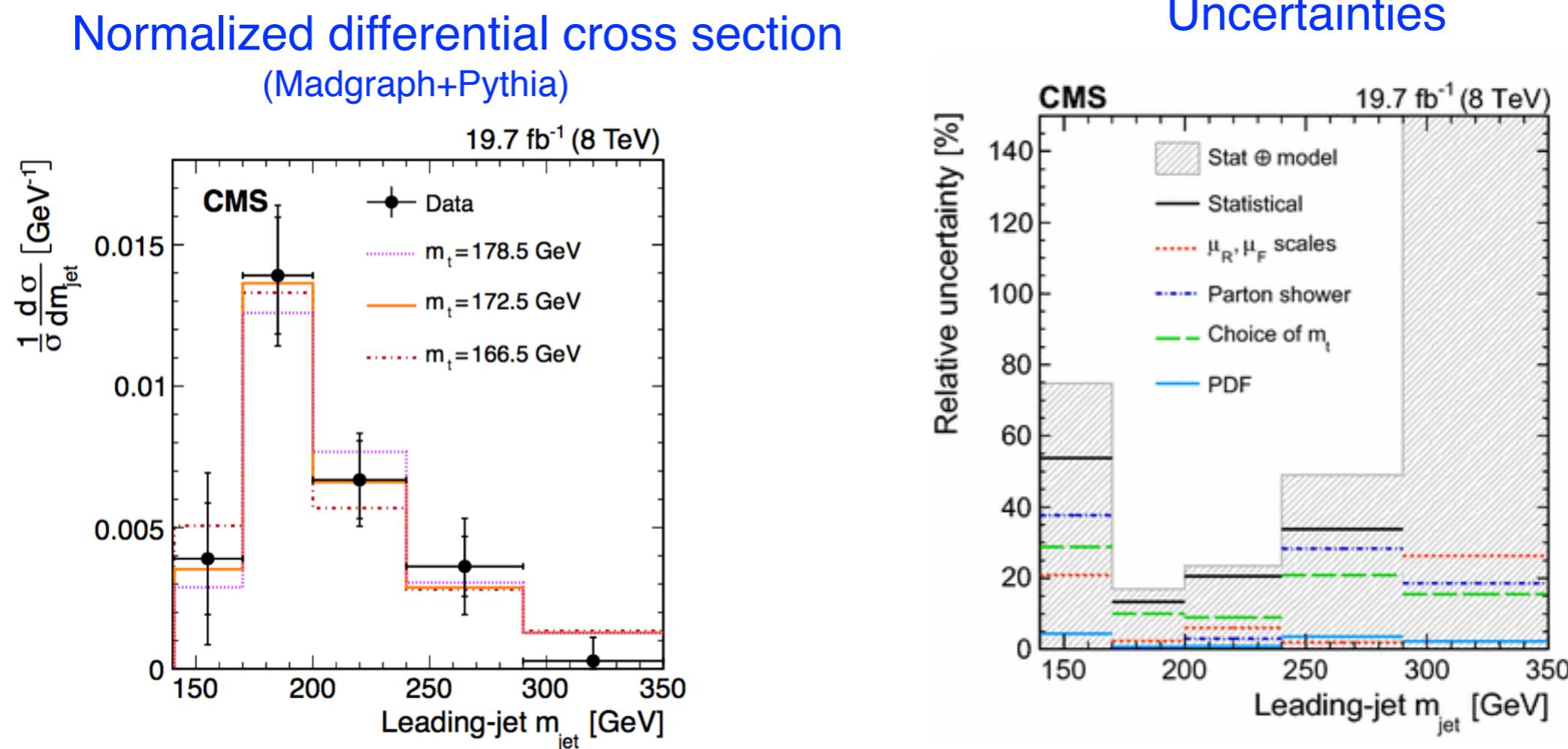
$\sigma_{\text{fid}}^{\text{obs}} \approx (0.7-0.8)\sigma_{\text{fid}}^{\text{th}}$

TOP2017



Boosted top: e/ μ + jets

Template method (m_{jet})



Fit

$M_t = 170.8 \pm 6.0 (\text{stat}) \pm 2.8 (\text{syst}) \pm 4.6 (\text{model}) \pm 4.0 (\text{theory}) \text{ GeV}$

$M_t = 170.8 \pm 9.0 \text{ GeV} \quad (\pm 5.3\%)$

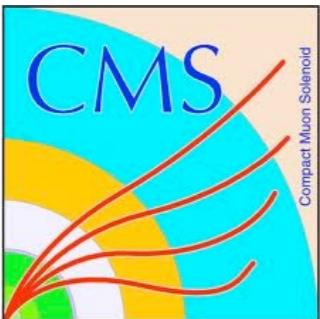
Method works but need more data, better modeling and higher-order calculations

What mass are we measuring?

The mass measured so far is the M_t^{MC} (typically LO or NLO) and is affected by perturbative/non-perturbative sub-1% uncertainties

The increasing level of accuracy requires to relate this to theory-based quantities like:

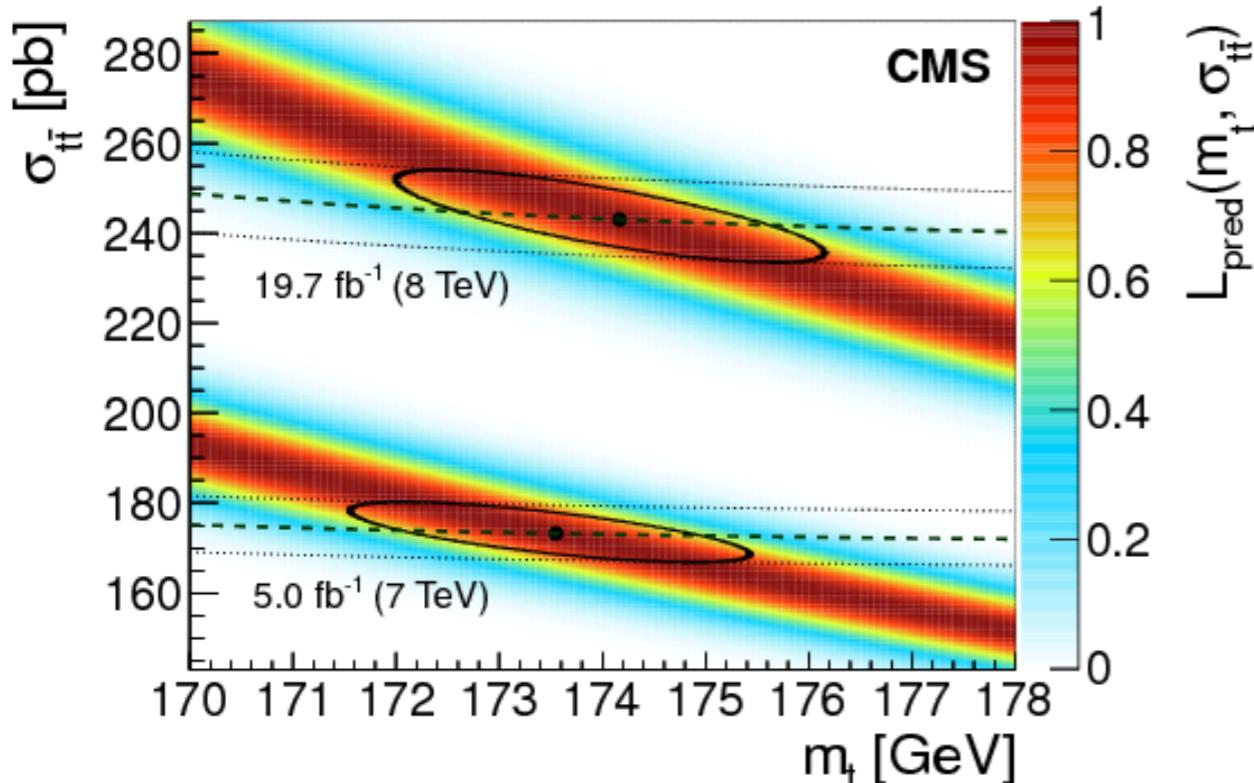
- the *pole mass*, universal but theoretically ambiguous by amounts $\mathcal{O}(\Lambda_{\text{QCD}})$ due to soft gluon radiation (*infrared renormalon problem*)
- *Lagrangian masses*, theoretically unambiguous but not universal, like the *MS* mass which is defined only in perturbation theory



Pole mass

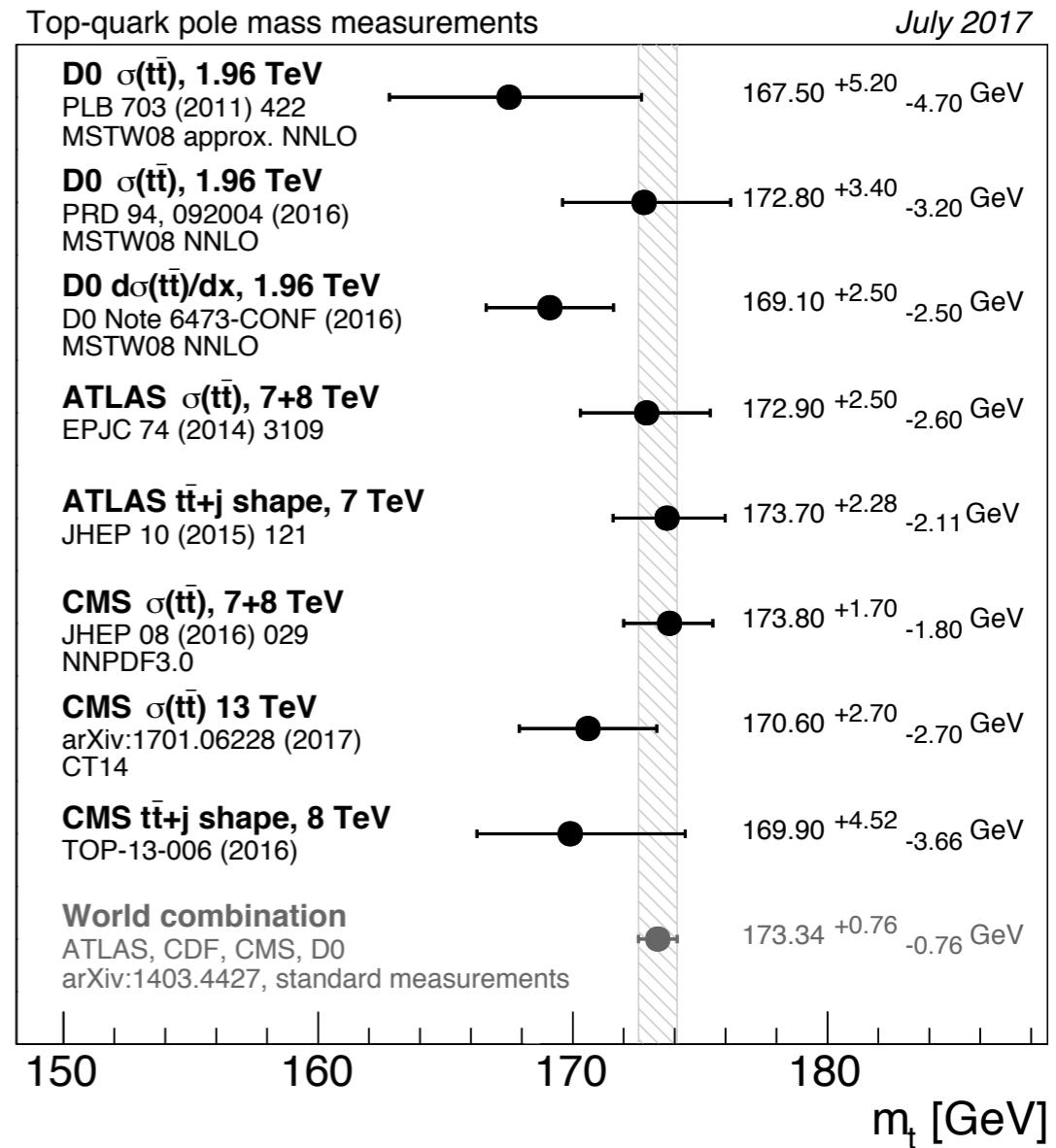
[JHEP 08 \(2016\) 029](#)

M_t^{pole} can be derived from $\sigma_{t\bar{t}}$ cross section measurements compared to NNLO
(but need to assume M_t^{MC})



Syst. dominated by PDF and lumi

1% precision, but could reach 0.5%
(CMS-PAS-FTR-16-006)



$M_t^{\text{pole}} = 173.8 + 1.7 - 1.8 \text{ GeV}$



Alternative methods

The precision of the standard measurements depends strongly on the hadronization modeling

To improve it we can use cleaner observables, i.e. avoid jets

M_t can be derived from observables other than the reconstructed mass, like for instance:

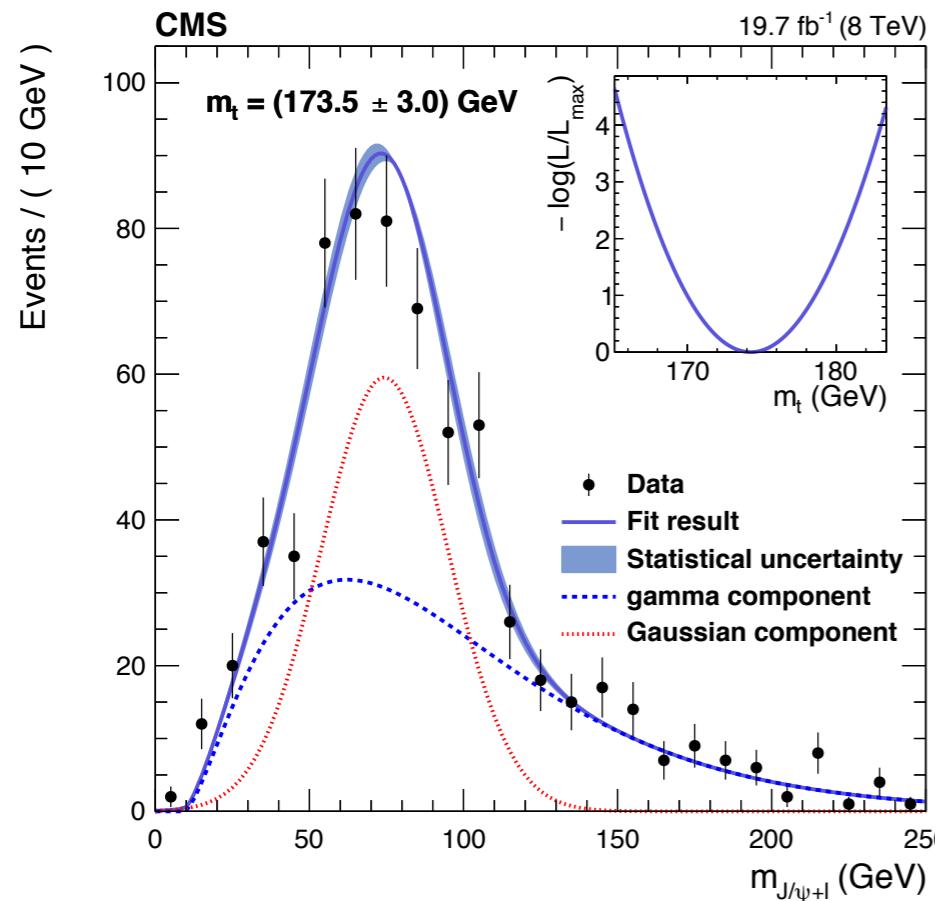
- lepton + J/ Ψ mass
- lepton + secondary vertex mass
- mass observables: $M_{b\ell}$, $M_{bb}^{bb}\text{T}_2$, $M_{b\ell\nu}$



Lepton+J/ Ψ

[JHEP 12 \(2016\) 123](#)

The M_t -sensitive quantity is the lepton+J/ Ψ mass



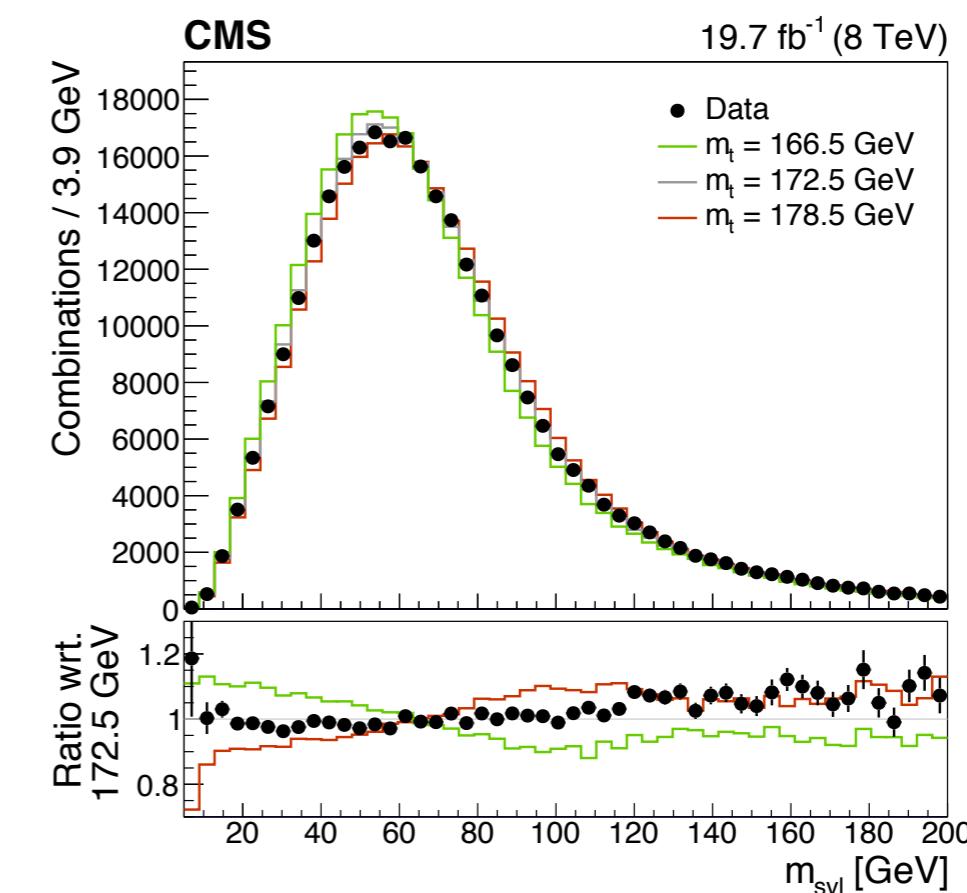
Limited by statistics, top p_T modeling and QCD scales

$$M_t = 175.3 \pm 3.0 \text{ (stat)} \pm 0.9 \text{ (syst)} \text{ GeV}$$

Lepton+sec.vertex

[PRD 93 \(2016\) 092006](#)

The M_t -sensitive quantity is the lepton+S.V. mass in bins of track multiplicity



Limited by b-fragmentation and top p_T modeling

$$M_t = 173.68 \pm 0.20 \text{ (stat)}^{+1.58}_{-0.97} \text{ (syst)} \text{ GeV}$$



Mass observables: $2e/\mu + \text{jets}$

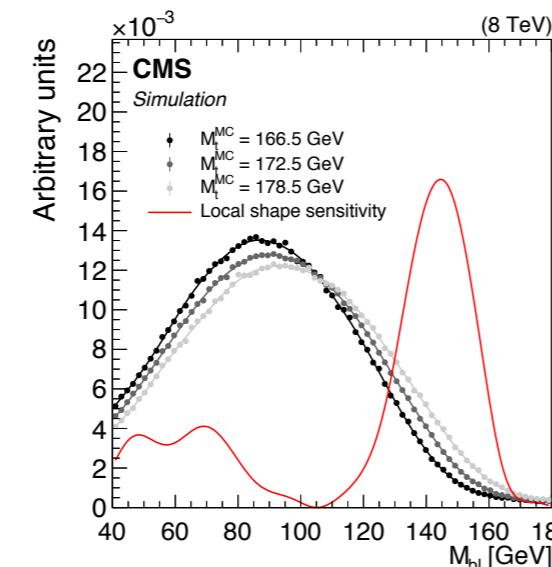
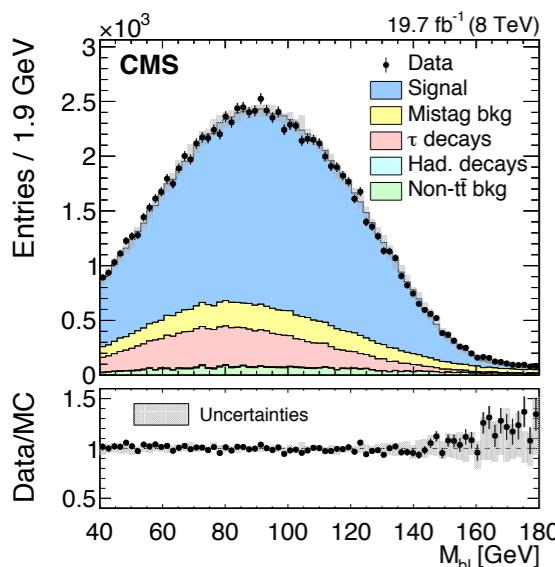
[PRD 96 \(2017\) 032002](#)

The M_t -sensitive quantities M_x are:

$$M_{b\ell}$$

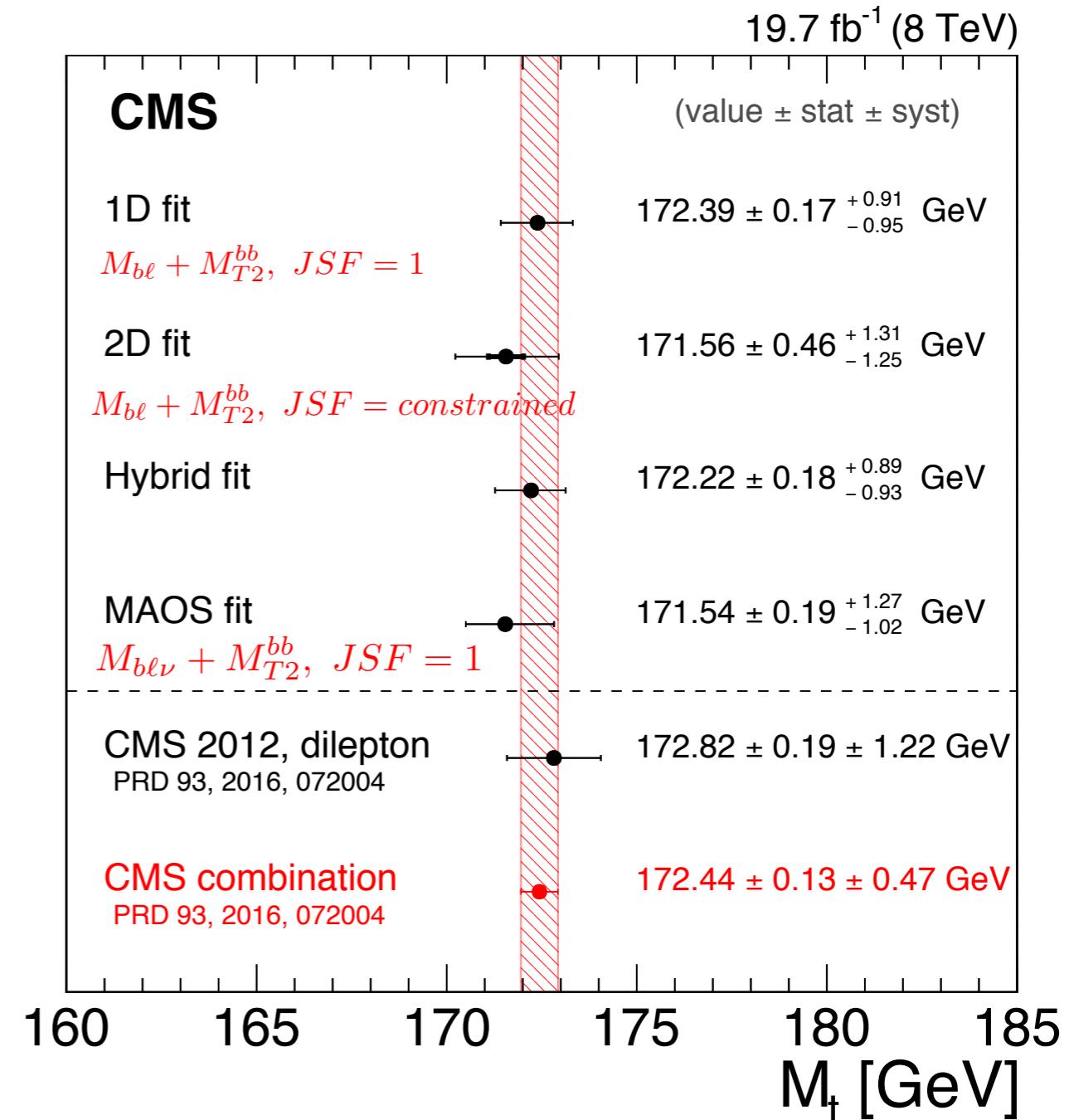
$$M_{T2}^{bb}$$

$$M_{b\ell\nu}$$



3D (M_x, M_t, JSF) non-parametrical modeling of these distributions

Systematics dominated by JES, b fragmentation and top p_T reweighting



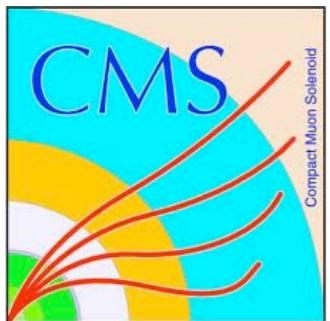
Hybrid fit = linear comb. of 1D and 2D

$M_t = 172.22 \pm 0.18 \text{ (stat)}^{+0.89}_{-0.93} \text{ (syst) GeV}$

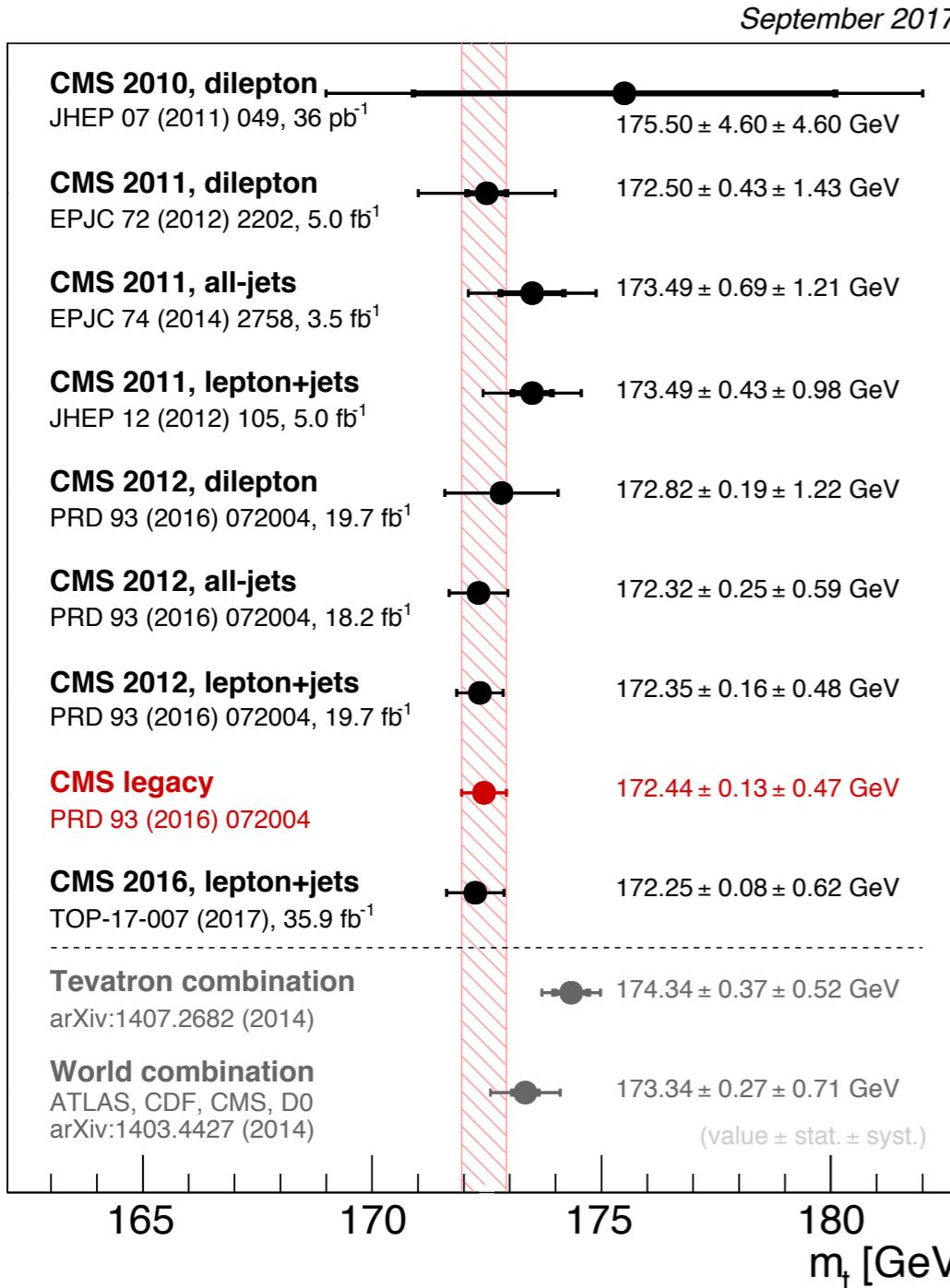
M_t combinations

The individual measurements are then combined into experiment/world averages to gain in precision

Results computed with the Best Linear Unbiased Estimator, accounting for correlations ρ in the systematics (ρ signs are relevant for large systematics)



7+8 TeV combination



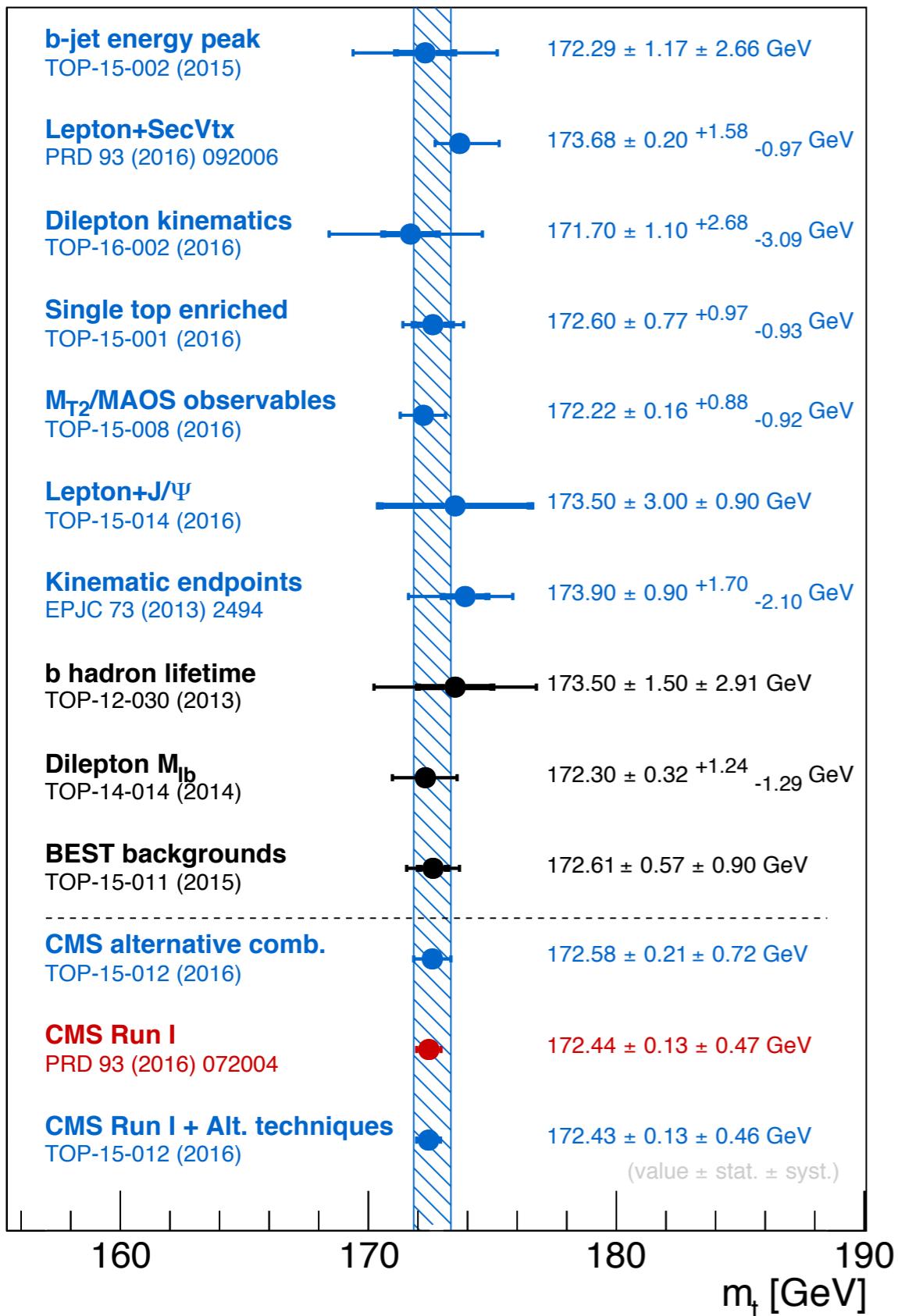
[**PRD 93 \(2016\) 072004**](#)

- Several decay channels pursued (some with in situ JSF)
- Different main systematics
- Combination to check consistency and increase precision

September 2015 value:

$M_t = 172.44 \pm 0.13(\text{stat}) \pm 0.47(\text{syst}) \text{ GeV}$

$M_t = 172.44 \pm 0.49 \text{ GeV} (\pm 0.28\%)$



7+8 TeV alternative methods

[CMS-PAS-TOP-15-012](#)

Slightly less precise (0.4%) than the standard measurements

Independent verification with different systematics

$M_t = 172.58 \pm 0.21 \text{ (stat)} \pm 0.72 \text{ (syst)} \text{ GeV}$

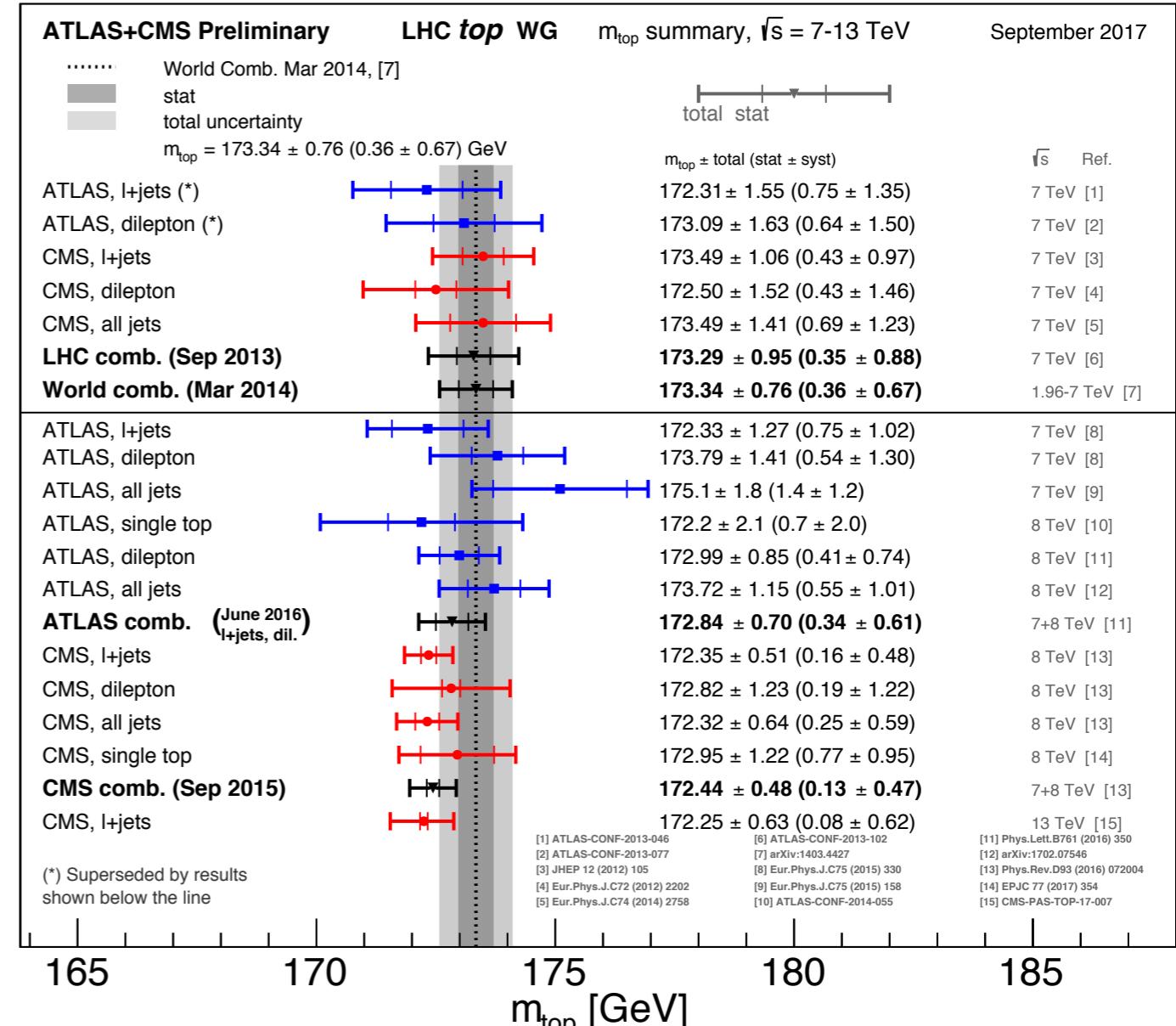


World combination

Combination of all
ATLAS+CMS (7+8 TeV) +
Tevatron results computed
accounting for correlations
in the systematics

March 2014 value:
[arXiv:1403.4427](https://arxiv.org/abs/1403.4427)

$M_t = 173.34 \pm 0.36(\text{stat}) \pm 0.67(\text{syst}) \text{ GeV}$
 $M_t = 173.34 \pm 0.76 \text{ GeV} \quad (\pm 0.44\%)$



Tevatron 2016 combination
 $174.30 \pm 0.65 \text{ GeV}$
[arXiv:1608.01881](https://arxiv.org/abs/1608.01881)



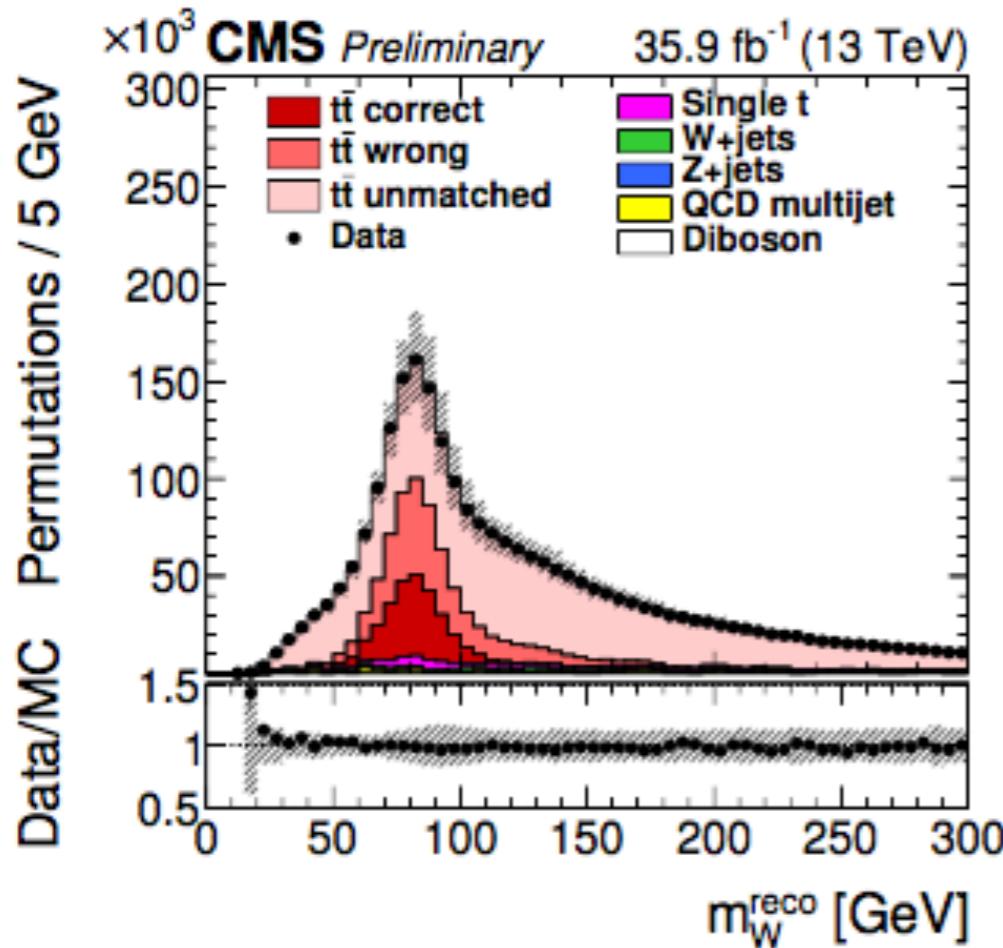
@13 TeV: $e/\mu + 4$ jets

Ideogram method
 2D fit (m_t^{reco} , m_w^{reco})
 w. *in situ JSF* (35.9 fb^{-1})

[CMS-PAS-TOP-17-007](#)

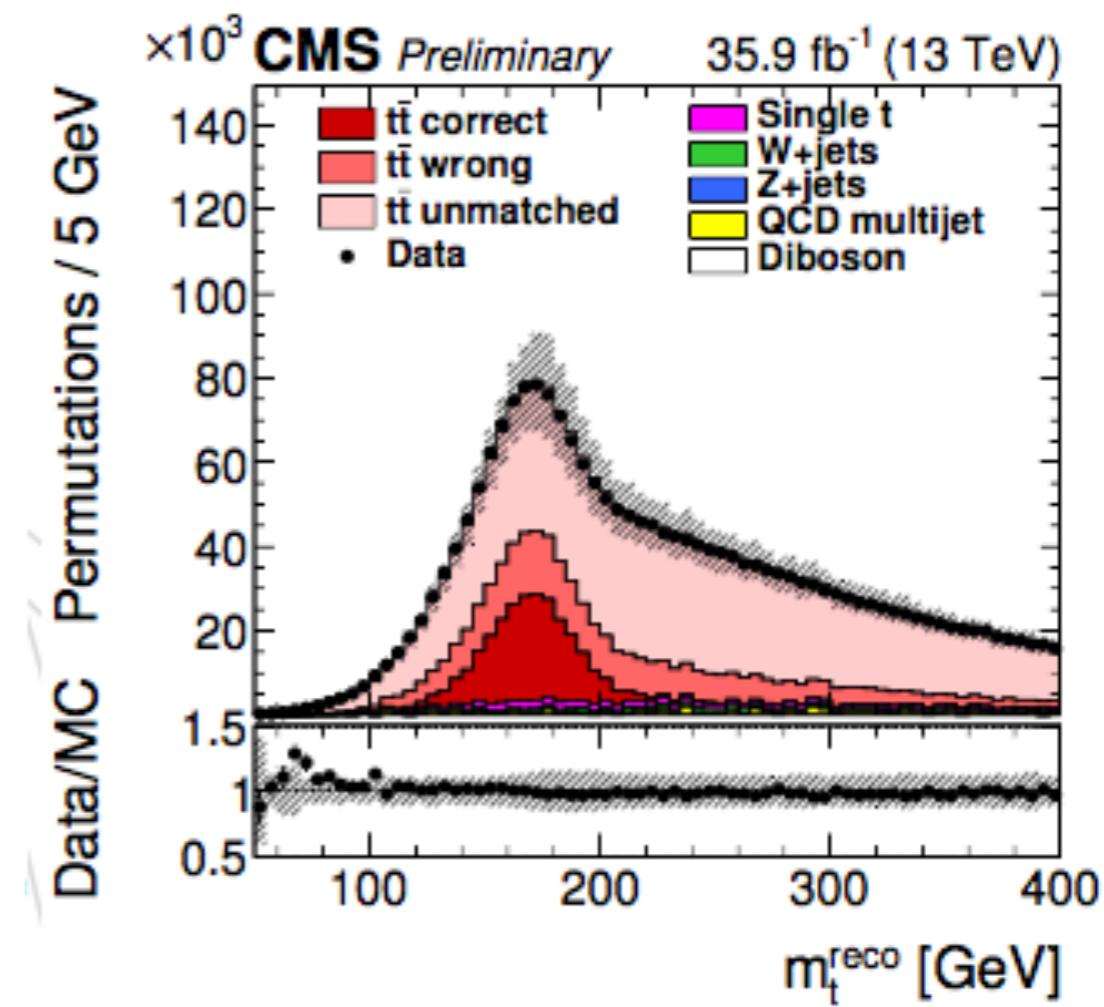
Selection:

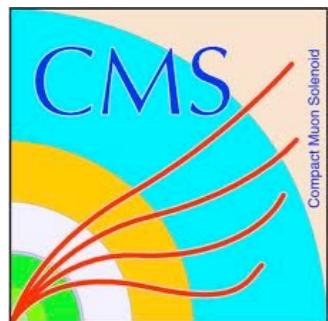
- 1 e or μ
- ≥ 4 jets (2 b-tagged)



Kinematic fit to the $t\bar{t} \rightarrow WbWb$ hypothesis
 Possible combinations treated separately:

- correct:
- wrong:
- unmatched

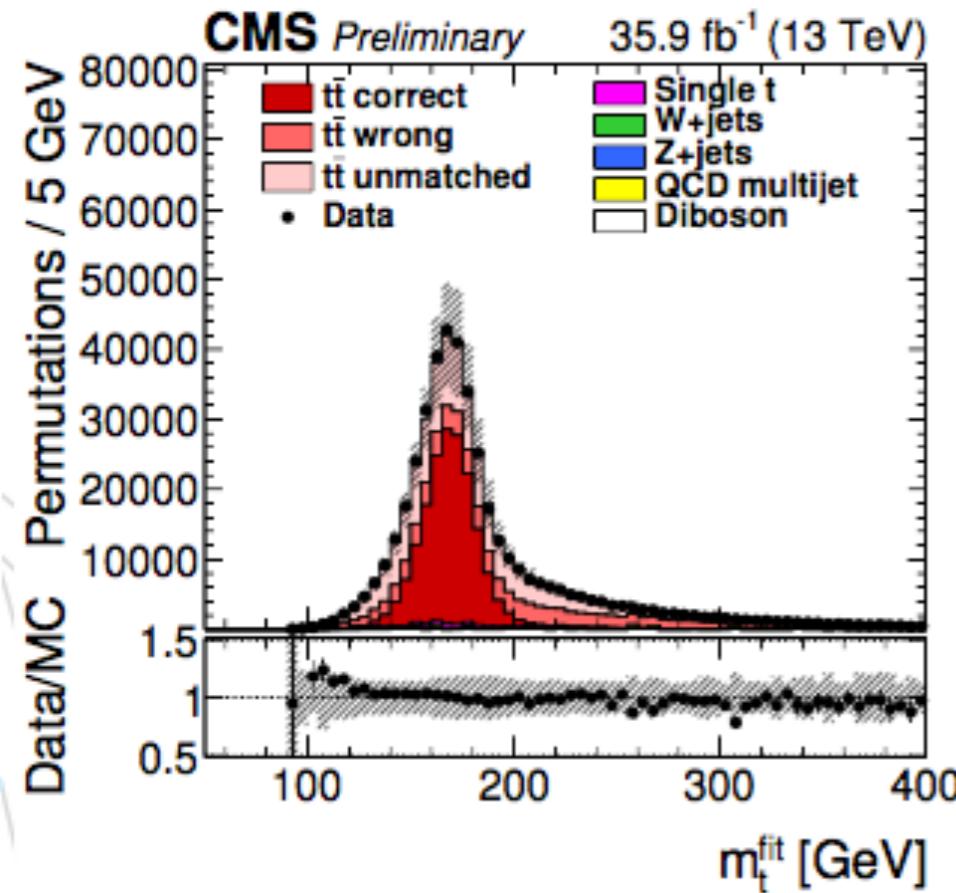




@13 TeV: $\mu + 4$ jets

*Ideogram method (m_t^{reco} , m_w^{reco})
w. in situ JSF*

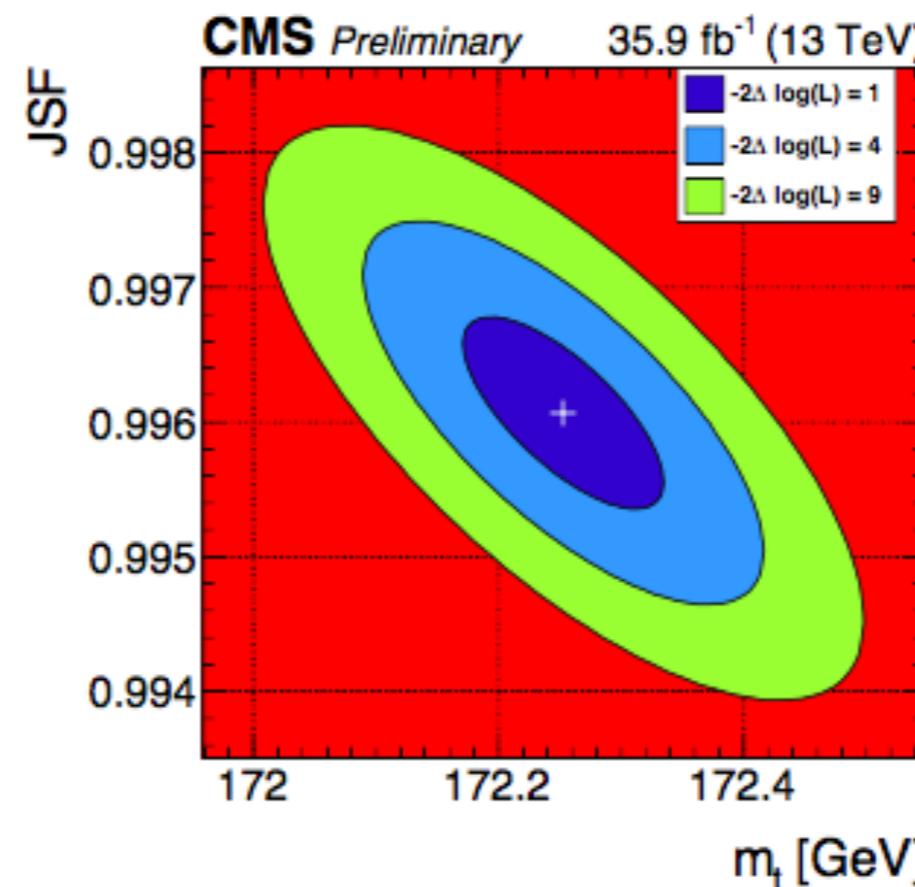
See E. Yazgan's talk



Multiple permutations weighted
by $P_{\text{gof}} = \exp(-\chi^2/2)$

$$P_{\text{gof}} = \exp(-\chi^2/2) > 0.2$$

to favor the correct combin.



Main systematics

source	GeV
flavor JEC	0,39
color reconn.	0,31
ME generator	0,22

Hybrid Fit
 $M_t = 172.25 \pm 0.08(\text{stat+JSF}) \pm 0.62(\text{syst}) \text{ GeV}$

$$M_t = 172.25 \pm 0.63 \text{ GeV} \quad (\pm 0.36\%)$$

$$\text{JSF} = 0.996 \pm 0.008$$

Conclusions

- Level of precision reached ($<0.3\%$) in measuring M_t is impressive but comes from 22 years of continuous improvements
- Even better precision expected from Run2; single top events and boosted top can contribute
- Inclusion of alternative methods will improve precision
- Test of the consistency of the SM, vacuum stability and new physics

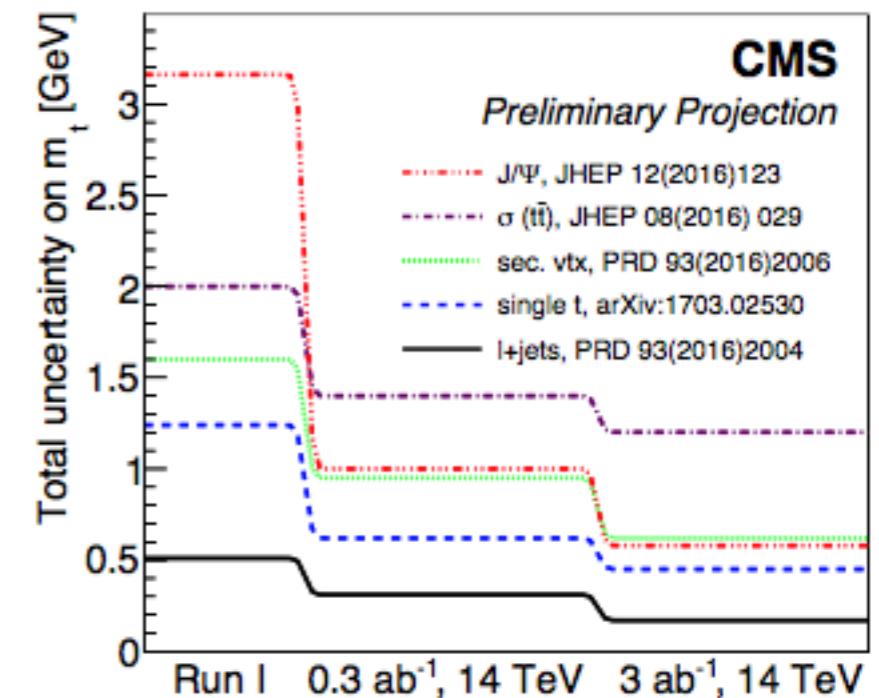
Important to work on reducing systematics
e.g. those related to theory and signal modeling

Outlook

Run1 legacy measurements of M_t being completed
⇒ published soon

[**CMS-PAS-FTR-16-006**](#)

Ultimate precision of few hundreds MeV expected merging measurements/experiments, accounting for correlations and taking advantage of improvements in MC modelling



Differences between M_t^{MC} and theoretical definitions (pole mass, Lagrangian mass): important issue to deal with

Backup

Methods for measuring M_t

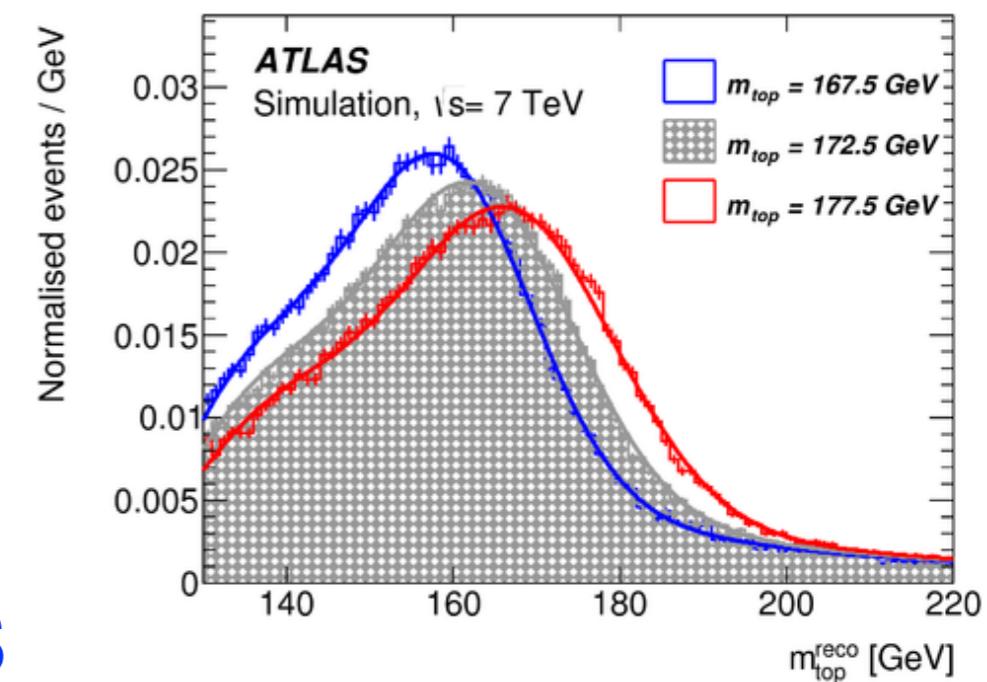
1) *Template method*: distributions of variables sensitive to M_t , e.g., reconstructed M_t^{reco} from χ^2 fit to WbWb

Pdf's derived for MC events assuming different M_t^{MC} ; parametrized vs M_t

Likelihood from pdf's; outcome calibrated for biases (pull-mean and pull-width of pseudo-experiments)

M_W templates for in-situ calibration of JES

Possible to add constraints on b-jet JES



Relatively simple, fast, but non optimal statistical uncertainty

Methods for measuring M_t

2) *Ideogram method*: modification of template method using multiple permutations with different weights

Starts from kinematical reconstruction, then computes event likelihood as a function of M_t

Different pdf's used for different jet-quark assignments

Event likelihoods (ideograms) are given by

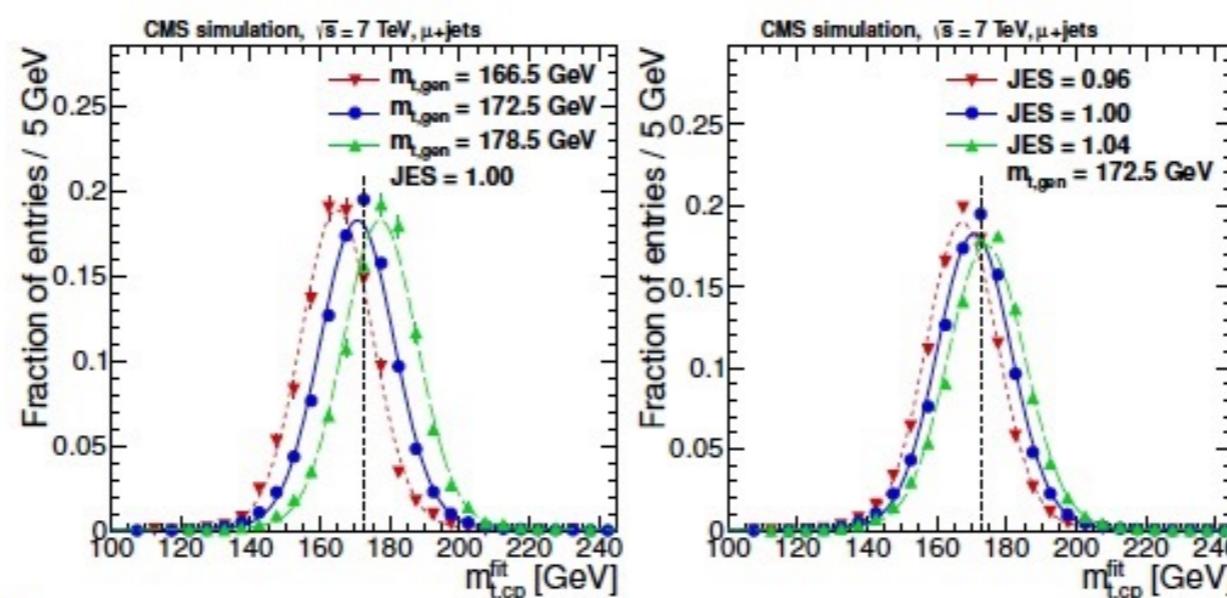
$$\mathcal{L}(\text{event}|m_t, \text{JSF}) = \sum_{i=1}^n P_{\text{gof}}(i) \left\{ f_{\text{sig}} P_{\text{sig}}(m_{t,i}^{\text{fit}}, m_{W,i}^{\text{reco}} | m_t, \text{JSF}) + (1 - f_{\text{sig}}) P_{\text{bkg}}(m_{t,i}^{\text{fit}}, m_{W,i}^{\text{reco}}) \right\}$$

$$P_{\text{gof}} = \exp(-\chi^2/2)$$

$$\mathcal{L}(\text{sample}|m_t, \text{JSF}) = \prod_{\text{events}} \mathcal{L}(\text{event}|m_t, \text{JSF})^{w_{\text{event}}}$$

$$w_{\text{event}} = \sum P_{\text{gof}}(i)$$

$$P_{\text{sig}}^{\text{cp}}(m_{t,\text{cp}}^{\text{fit}} | M_t, \text{JSF})$$



Gaussian process regression technique

GP shape determined by:

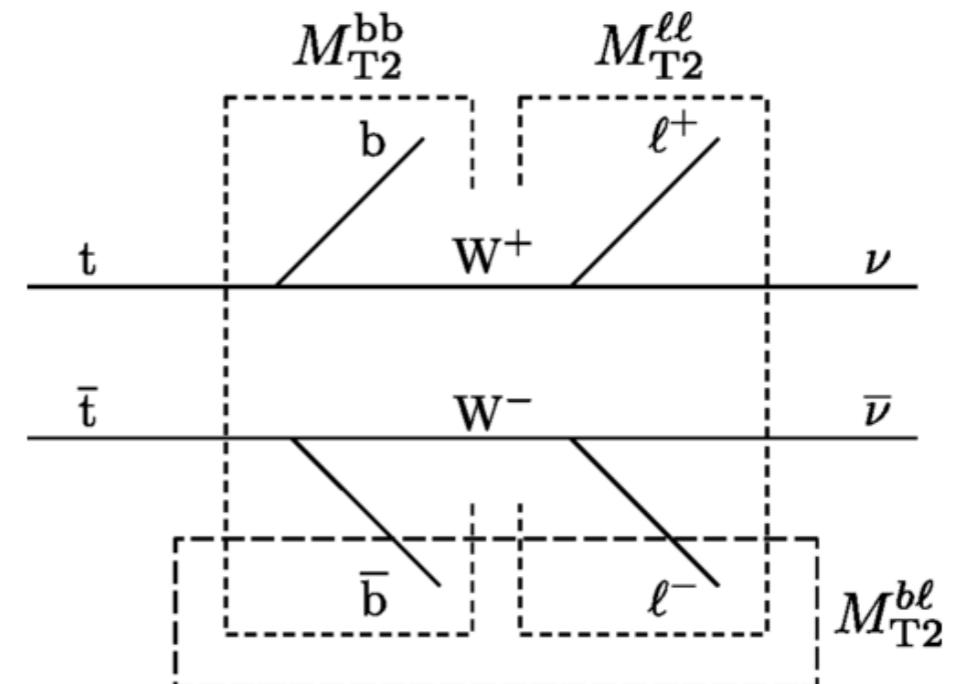
- a set of training points
- smoothing parameters
- 35 binned distributions (7 bins for M_t , 5 for JSF, 75 for M_x)
- Gaussian = values in each point are Gaussian distributions

Mass observables: M_{T2}^{bb}

$$M_T = \sqrt{m_\ell^2 + m_\nu^2 + 2(E_{T\ell}E_{T\nu} - \vec{p}_{T\ell} \cdot \vec{p}_{T\nu})},$$

$$M_{T2} = \min_{\vec{p}_T^a + \vec{p}_T^b = \vec{p}_T^{\text{miss}}} [\max\{M_T^a, M_T^b\}],$$

Impose the constraint on the invisible particle momenta + preserve the kinematic endpoint



Systematic uncertainties

Experimental (*i.e. imperfect knowledge of*):

- Jet Energy Scale (JES)
- b-Jet Energy Scale (bJES)
- jet energy resolution and reconstruction
- MET scale
- b-tagging scale factor
- lepton energy scale and reconstruction
- pileup
- trigger

Signal modeling (*i.e. imperfect knowledge of theory regarding*):

- MC generator
- hadronization
- amount of ISR/FSR
- flavor-dependent hadronization
- b-quark fragmentation and BRs
- renormaliz./factoriz. scales
- PDF's
- Color reconnection
- Underlying event

Background modeling (*i.e. uncertainty on*):

- MC normalization and shape
- normalization and shapes of data-driven backgrounds

Features of the method (*i.e. dependence on*):

- parametrization of pdf's
- calibration
- MC statistics

Agreement between ATLAS and CMS is essential