

Understanding the Impact of Pythia Hadronization in Top Quark Mass Determinations at the LHC

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TOP 2017, Braga, Portugal, September 19, 2017

Gennaro Corcella, Roberto Franceschini and DK, 1709.xxxxx

Top Quark Mass Measurements

❑ Precision m_{top} measurement: extremely important in both SM and BSM (see Nathaniel's talk)

❑ From standard/conventional approaches to alternative ones

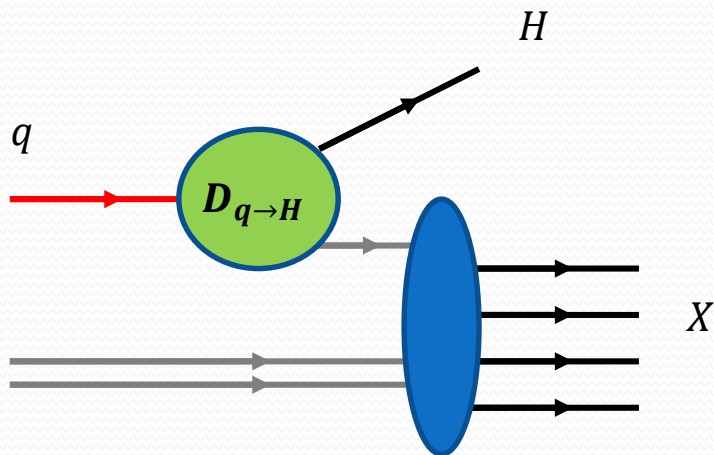
- ❖ Template method [ATLAS, Eur. Phys. J. C72 (2012)]
 - ❖ Ideogram method [CMS PAS TOP 14-001]
 - ❖ Matrix element method [DØ, arXiv:1501.07912]
 - ❖ Cross sections [ATLAS, Eur. Phys. J. C74 (2014), CONF 2014-053]
 - ❖ Endpoint method [CMS PAS TOP 11-027; CMS TOP 15-008]
 - ❖ b -jet energy-peak method [CMS PAS TOP 15-002]
 - ❖ Solvability method [DK, Matchev and Shyamsundar, to appear soon]
 - ❖ J/ψ method [CMS PAS TOP 15-014]
 - ❖ B -hadron 2D-decay length [CMS PAS TOP 12-030]
 - ❖ Leptonic final state [CMS PAS TOP 16-002]
 - ❖ B -hadron observables [Corcella, Franceschini and DK, to appear soon]
 - ❖ Many more which I can't exhaust
- Diagrammatic groupings:
- SM top assumed** (green bracket): Template method, Ideogram method, Matrix element method, Cross sections.
 - Kinematics-based** (yellow bracket): Endpoint method, b -jet energy-peak method, Solvability method.
 - No jetty objects in the final state → no JES, Th. uncertainty** (red bracket): J/ψ method, B -hadron 2D-decay length, Leptonic final state, B -hadron observables.
 - Jets in the final state → JES** (blue bracket): All methods listed above.

B-hadron Observables

- *B*-hadron observables: **crucial** to understand the transformation $b \rightarrow B$, but **challenging** because it is governed by non-perturbative QCD (similar conclusions hold for *B*-hadron decay length method [Hill, Incandela, Lamb (2005); CMS-PAS-TOP-12-030])

Filling the Gap: Theoretical

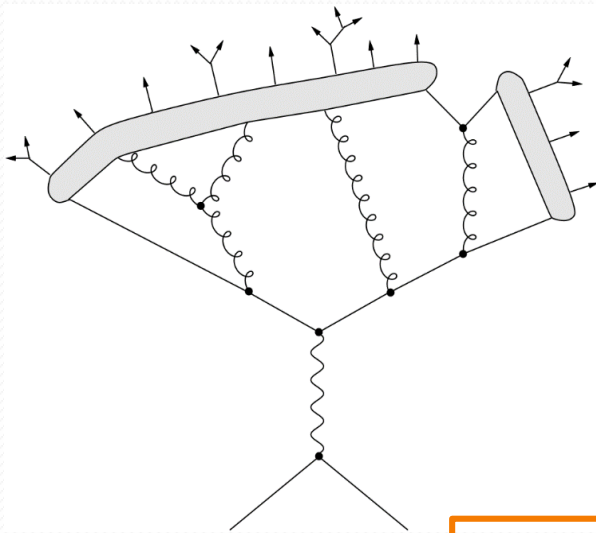
- ❑ B -hadron observables: **crucial** to understand the transformation $b \rightarrow B$, but **challenging** because it is governed by non-perturbative QCD (similar conclusions hold for B -hadron decay length method [Hill, Incandela, Lamb (2005); CMS-PAS-TOP-12-030])



- ❑ Computing *fragmentation function*, $D_{q \rightarrow H}(z)$
- ❑ Precision data available at LEP [arXiv: 1102.4748, hep-ex/01120282] and SLC [hep-ex/0202031]
- ❑ For b quark, the extraction of the fragmentation function at NNLO in α_s available [Fickinger, Fleming, Kim, Merechetti (2016)]
- ❑ Higher order corrections necessary (including resummation sometimes)
- ❑ Relying on factorization of the cross section to a very high accuracy
- ❑ Not clear to apply lepton collider data to hadron colliders

Filling the Gap: Phenomenological

- ❑ B -hadron observables: **crucial** to understand the transformation $b \rightarrow B$, but **challenging** because it is governed by non-perturbative QCD (similar conclusions hold for B -hadron decay length method [Hill, Incandela, Lamb (2005); CMS-PAS-TOP-12-030])



- ❑ Employing **hadronization model** with phenomenological parameters [Andersson, Gustafson, Ingelman, Sjostrand (1983)]
- ❑ “Tuning” of the parameters to reproduce the available data
- ❑ **Not obvious** that the **tuned model** (with $e^+e^- \rightarrow \text{hadrons}$) **describes the future data** [D. d’Enterria et al. (2013)]
- ❑ Should be tested at hadron collider environment (**incredible amount of statistics** available!!)

❑ Our approach/goal

- top quark mass *sensitivity to parameters*,
- **what/how to constrain** to achieve better precision

B-hadron-related Parameters

- ❑ $e\mu$ channel of LO $t\bar{t}$ ($gg/q\bar{q} \rightarrow t\bar{t}$) at the LHC 13 TeV with NNPDF2.3 QCD+QED LO, PartonLevel:MPI = off, HadronLevel:Decay = off, and anti- k_t jets of $R = 0.5$
- ❑ $p_{T,j(\ell)} > 30$ (20) GeV, $|\eta_{j(\ell)}| < 2.4$ (2.4)
- ❑ Input top quark mass varies from 170 GeV to 180 GeV by an interval of 0.5 GeV

Parameter	PYTHIA8 setting	Variation range	
r_B in the Bowler modification for heavy quarks	StringZ:rFactB	0.713 – 0.813	Heavy flavor-specific hadronization parameters
a parameter in the non-standard Lund ansatz for b quarks	StringZ:aNonstandardB	0.54 – 0.82	
b parameter in the non-standard Lund ansatz for b quarks	StringZ:bNonstandardB	0.78 – 1.18	
$P_{T,\min}^{\text{FSR}}$ [GeV]	TimeShower:pTmin	0.25 – 1.00	Showering parameters
recoiler switch	TimeShower:recoilToColoured	on or off	
$\alpha_s^{\text{FSR}}(m_Z)$	TimeShower:alphaSvalue	0.1092 – 0.1638	
b quark mass [GeV]	5:m0	3.84 – 5.76	

Summary of Results: Mellin Moments

$$\mathcal{M}_1 = \int_{FWHM} dx x f(x)$$

$$\Delta_{\theta}^{(m_t)} = \frac{\delta m_t / m_t}{\delta \theta / \theta}$$

Mellin-1	$\Delta_{\theta}^{(m_t)}$					
	$\alpha_{s,FSR}$	$p_{T,min}$	recoil	r_B	a	b
E_B	0.43	0.019	0.028	0.039	0.020	0.039
$E_B + E_B$	0.42	0.019	0.032	0.046	0.023	0.034
$p_{T,B}$	0.43	0.021	0.027	0.043	0.022	0.042
$p_{T,B} + p_{T,B}$	0.36	0.017	0.024	0.042	0.016	0.044
$m_{B\ell,min}$	0.26	0.011	0.016	0.041	0.011	0.031
$m_{B\ell,true}$	0.24	0.008	0.013	0.031	0.009	0.022
$m_{T2,B\ell,true}^{(mET)}$	0.24	0.012	0.012	0.034	0.013	0.032
$m_{T2,B\ell,min}^{(mET)}$	0.23	0.011	0.012	0.031	0.012	0.03
$m_{T2,B\ell,min,\perp}^{(mET)}$	0.24	0.01	0.011	0.038	0.011	0.03
$m_{T2,B\ell,true}^{(ISR)}$	0.21	0.007	0.01	0.022	0.008	0.022
$m_{T2,B\ell,min}^{(ISR)}$	0.2	0.008	0.013	0.024	0.01	0.027
$m_{T2,B\ell,min,\perp}^{(ISR)}$	0.22	0.01	0.01	0.03	0.013	

Preliminary

- Top quark mass measurements in these observables are **sensitive most to $\alpha_{s,FSR}$** , e.g., 10% uncertainty in $\alpha_{s,FSR}$ corresponds to 2 – 4% uncertainty in the top quark mass \Rightarrow affecting radiation in the final state, in turn, changing energy scale of B -hadrons!

Summary of Results: Shape

\mathcal{O}	$\Delta_{\theta}^{(m_t)}$					
	$\alpha_{s,FSR}$	PT_{\min}	recoil	r_B	a	b
$E_{B,\text{peak}}$	0.46	0.018	0.032	0.057	0.018	0.042
$\dot{m}_{B\ell,\min}$	0.0021	$< 10^{-3}$	$< 10^{-3}$	0.0038	0.0021	0.0017
$\dot{m}_{B\ell,\text{true}}$	0.005	$< 10^{-3}$	$< 10^{-3}$	0.0035	0.0014	$< 10^{-3}$
$\dot{m}_{T2,B\ell,\text{true}}^{(\text{mET})}$	$< 10^{-3}$	$< 10^{-3}$	$< 10^{-3}$	0.0029	0.0015	0.0014
$\dot{m}_{T2,B\ell,\min}^{(\text{mET})}$	0.0016	$< 10^{-3}$	$< 10^{-3}$	0.0035	0.0012	$< 10^{-3}$
$\dot{m}_{T2,B\ell,\min,\perp}^{(\text{mET})}$	0.016	$< 10^{-3}$	$< 10^{-3}$	0.0071	0.0019	0.0017
$\dot{m}_{T2,B\ell,\text{true}}^{(\text{ISR})}$	0.0067	0.0015	$< 10^{-3}$	0.0042	0.0021	0.0024
$\dot{m}_{T2,B\ell,\min}^{(\text{ISR})}$	0.0063	0.0015	$< 10^{-3}$	0.0031	0.0021	0.0023
$\dot{m}_{T2,B\ell,\min,\perp}^{(\text{ISR})}$	0.0056	$< 10^{-3}$	$< 10^{-3}$	0.0042	0.0016	0.0014

$$\Delta_{\theta}^{(m_t)} = \frac{\delta m_t / m_t}{\delta \theta / \theta}$$

Preliminary

- ❑ Top quark mass measurements in shape observables are **less sensitive to $\alpha_{s,FSR}$** (except energy-peak in B -hadron energy spectrum) \Rightarrow kinematic endpoints are less affected by process dynamics
- ❑ Sensitivities of top quark mass to the Lund-Bowler parameter becomes comparable!
- ❑ Statistics will be a major challenge in performing precision measurements of endpoints.

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How to achieve a small uncertainty? \Rightarrow Calibration variables!

Calibration Variables

- ❑ Calibration variables: **no/little sensitivity** to (input) **top quark mass**, but having **decent sensitivities** to **hadronization and showering parameters** in $t\bar{t}$ events [see for similar effort, e.g. ATL-PHYS-PUB-2015-007]

$$\diamond \frac{p_{T,B}}{p_{T,j_b}}, \rho(r) = \frac{1}{\Delta r} \frac{1}{E_j} \sum_{\text{track}} E_{\text{track}} \cdot \Theta(|r -$$



Don't call the Nobel Committee just yet:
We forgot to calibrate the instruments
before the experiment...

Calibration Variables

● In what sense

□ Calibration variables sensitive to hadronization and showering parameters

- ❖ Variables $\frac{p_{T,B}}{p_{T,j_b}}$ and $\rho(r)$ are sensitive to the importance of the heavy-quark hadron in the jet and to the energy distribution in the jet \Rightarrow suitable to **probe the dynamics on the conversion of a single parton into a hadron**
- ❖ χ_B variables are more sensitive to global nature (i.e., $b\bar{b}$ system) \Rightarrow probing **“cross-talk” between partons** in the process of forming color-singlet hadrons
- ❖ Various aspects probed by different χ_B options

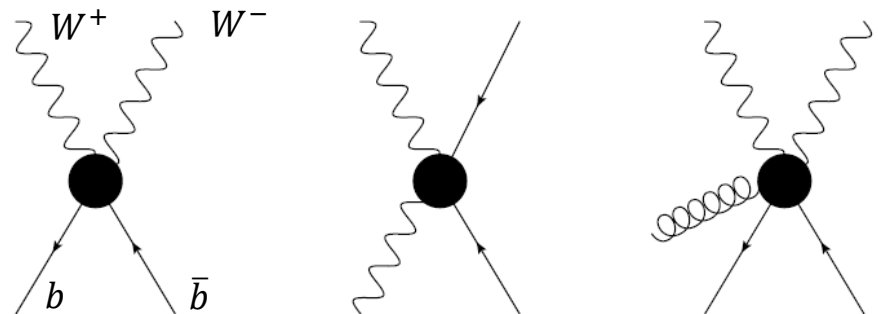


Figure 2: Three kinematical configurations distinguished by the X_B choices. The first two can have same $|p_{T,j_b}| + |p_{T,\bar{j}_b}|$ but differ for $m_{b\bar{b}}$, whereas the first and the third differs for $\sqrt{s_{min}}$, despite having same $m_{b\bar{b}}$ and same $|p_{T,j_b}| + |p_{T,\bar{j}_b}|$.

Sensitivities investigated from different angles!!

Sensitivity Measure

□ Sensitivity measure: $\Delta_{\theta}^{(\mathcal{O})} = \frac{\delta \mathcal{O} / \mathcal{O}}{\delta \theta / \theta}$

❖ \mathcal{O} : Observable

❖ θ : hadronization and showering parameters

□ Observables with larger Δ : **best diagnostics** of the accuracy of the tunes

Summary of Results

\mathcal{O}	Range	$\Delta_m^{(\mathcal{M}_O)}$	$\Delta_\theta^{(\mathcal{M}_O)}$						
			$\alpha_{s,FSR}$	m_b	PT_{min}	a	b	r_B	recoil
$\rho(r)$	0-0.04	-0.007(7)	0.78(1)	0.204(4)	-0.1286(8)	0.029(3)	-0.043(4)	0.056(7)	0.020(1)
PT_B/PT_{j_b}	0.6-0.998	-0.053(1)	-0.220(3)	-0.1397(8)	0.0353(5)	-0.0187(4)	0.0451(6)	-0.0518(9)	-0.0108(3)
E_B/E_{j_b}	0.6-0.998	-0.049(1)	-0.220(3)	-0.1381(8)	0.0360(5)	-0.0186(4)	0.0447(6)	-0.052(1)	-0.0107(3)
E_B/E_ℓ	0.05-1.5	-0.155(7)	-0.156(3)	-0.053(3)	0.0149(7)	-0.007(2)	0.016(2)	-0.016(10)	-0.0087(7)
$E_B/(E_\ell + E_{\bar{\ell}})$	0.05-1.0	0.021(5)	-0.231(2)	-0.082(4)	0.0228(4)	-0.011(2)	0.026(2)	-0.028(6)	-0.0113(3)
$m(j_b)/\text{GeV}$	8-20	0.229(3)	0.218(1)	0.022(1)	-0.0219(7)	0.000(1)	-0.001(1)	0.001(3)	0.0050(3)
$\chi_B(\sqrt{s_{min,b\bar{b}}})$	0.075-0.875	-0.177(4)	-0.262(4)	-0.086(1)	0.0255(3)	-0.0105(10)	0.027(1)	-0.031(3)	-0.0137(2)
$\chi_B(E_{j_b} + E_{\bar{j}_b})$	0.175-1.375	-0.109(2)	-0.357(4)	-0.134(1)	0.0373(3)	-0.016(1)	0.040(1)	-0.045(4)	-0.0175(3)
$\chi_B(m_{j_b j_{\bar{b}}})$	0.175-1.375	-0.089(3)	-0.252(3)	-0.080(1)	0.0248(3)	-0.010(1)	0.024(1)	-0.028(5)	-0.0126(2)
$\chi_B(PT_{j_b} + PT_{\bar{j}_b})$	0.46-1.38	-0.15(2)	-0.47(1)	-0.189(10)	0.054(3)	-0.023(10)	0.06(1)	-0.07(4)	-0.022(2)
$m_{BB}/m_{j_b j_{\bar{b}}}$	0.8-0.95	-0.0191(8)	-0.0623(7)	-0.0464(5)	0.0146(2)	-0.0093(3)	0.0180(4)	-0.0212(9)	-0.00296(10)
$\Delta\phi(j_b j_{\bar{b}})$	0.28-3.	-0.210(7)	0.027(3)	0.001(2)	-0.0014(5)	-0.000(3)	-0.000(1)	-0.003(9)	0.0003(5)
$\Delta R(j_b j_{\bar{b}})$	1.4-3.3	-0.071(3)	0.010(1)	0.0005(10)	-0.0004(2)	-0.000(1)	0.0004(9)	0.001(3)	0.0001(2)
$\Delta\phi(BB)$	0.28-3.	-0.207(7)	0.026(2)	0.001(1)	-0.0008(4)	0.000(4)	0.000(2)	-0.000(8)	0.0002(5)
$\Delta R(BB)$	1.4-3.3	-0.070(3)	0.009(1)	0.000(1)	-0.0003(2)	-0.0003(10)	0.0002(9)	-0.000(4)	0.0001(2)
$ \Delta\phi(BB) - \Delta\phi(j_b j_{\bar{b}}) $	0-0.0488	0.06(1)	0.734(6)	0.099(5)	-0.088(2)	0.006(5)	-0.004(5)	0.01(2)	0.0001(2)
$ \Delta R(BB) - \Delta R(j_b j_{\bar{b}}) $	0-0.0992	0.10(1)	0.920(3)	0.079(5)	-0.075(1)	-0.000(4)	0.005(4)	0.0001(2)	0.0418(8)

preliminary

- ❑ $\rho(r)$: (typically) **most sensitive variable** to both hadronization and shower parameters
- ❑ Nevertheless, **other variables contain useful/orthogonal information** to constrain parameters (unless they are perfectly correlated)!!

Constraining Power

□ Sensitivity matrix $\left(\Delta_{\theta}^{(\mathcal{O})}\right)_{ij} : \left(\frac{\delta \mathcal{O}}{\delta \theta}\right)_i = \left(\Delta_{\theta}^{(\mathcal{O})}\right)_{ij} \left(\frac{\delta \theta}{\delta}\right)_j$

❖ Singular value decomposition of the sensitivity matrix:

1.5, 0.17, 0.05, 0.027, 0.01, 0.0071, 0.0033 \Rightarrow (Roughly saying) a determination of the calibration variables with some accuracy leads to a constrain of **one** Monte Carlo parameter with a similar precision.

□ Relative uncertainties

$\alpha_{s,FSR}$	m_b	$p_{T,min}$	a	b	r_B	recoil
0.021	0.49	0.46	0.92	3.6	4.8	4.1

Preliminary

Table 5: Relative uncertainty on the physical parameters that could be achieved for a 1% measurement of the calibration observables.

□ (Note that precision level of the calibration observables is limited to JES.)

Differential Constraining Power

- ❑ Coarse-grained bin counts as calibration observables

preliminary

$p_{T,B}/p_{T,j_b}$	$\alpha_{s,FSR}$	m_b	$p_{T,min}$	a	b	r_{recoil}	recoil
[0,0.08]	0.82(5)	-0.30(3)	0.046(8)	-0.05(5)	0.11(4)	-0.2(1)	0.051(6)
[0.08,0.16]	1.21(5)	-0.19(4)	0.016(7)	-0.04(3)	0.06(4)	-0.1(1)	0.067(8)
[0.16,0.24]	1.55(5)	-0.00(2)	-0.029(8)	-0.03(2)	0.00(1)	0.01(8)	0.091(8)
[0.24,0.32]	1.83(3)	0.14(2)	-0.071(6)	-0.00(2)	-0.03(2)	0.02(6)	0.112(6)
[0.32,0.4]	2.04(3)	0.31(3)	-0.113(5)	0.02(2)	-0.08(1)	0.11(4)	0.133(5)
[0.4,0.48]	2.17(2)	0.50(2)	-0.160(4)	0.03(1)	-0.13(1)	0.13(5)	0.146(3)
[0.48,0.56]	2.21(1)	0.76(1)	-0.213(3)	0.05(1)	-0.20(1)	0.18(4)	0.153(3)
[0.56,0.64]	2.091(9)	1.081(9)	-0.277(2)	0.087(8)	-0.30(1)	0.33(3)	0.143(2)
[0.64,0.72]	1.48(2)	1.20(2)	-0.307(2)	0.129(8)	-0.407(8)	0.43(2)	0.096(2)
[0.72,0.80]	0.16(3)	0.57(3)	-0.207(2)	0.129(4)	-0.27(1)	0.33(2)	0.014(1)
[0.80,0.88]	-1.41(5)	-0.71(3)	0.061(5)	0.009(7)	0.140(9)	-0.12(2)	-0.067(3)
[0.88,0.96]	-2.32(6)	-1.74(2)	0.41(1)	-0.296(8)	0.631(9)	-0.74(1)	-0.119(5)

Table 6: Sensitivity of the bin counting of the $p_{T,B}/p_{T,j_b}$ distribution to variations of the Pythia8 parameters.

Decent sensitivity to all parameters

Differential Constraining Power

❑ Relative uncertainties

$\alpha_{s,FSR}$	m_b	$p_{T,min}$	a	b	r_B	recoil
0.0065	0.038	0.093	0.23	0.62	1	2.8

$\alpha_{s,FSR}$	m_b	$p_{T,min}$	a	b	r_B	recoil
0.008	0.24	1.2	0.69	1	2.4	3.1

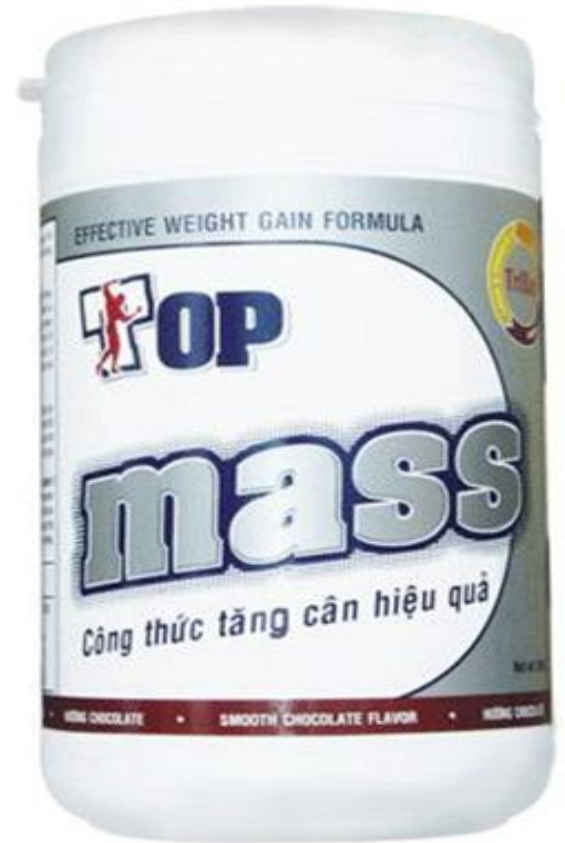
Preliminary

Table 7: Relative uncertainty on the physical parameters that could be achieved for a 1% measurement of the whole spectrum of $p_{T,B}/p_{T,j_b}$ (top) or $\rho(r)$ (bottom).

- ❖ More parameters are constrained.
- ❖ $\alpha_{s,FSR}$ can be best constrained.
- ❖ More dedicated studies are needed.

Conclusions

- ❑ We **first** perform a **systematic study** on *B*-hadron observable methods and potential impact of Pythia parameters on them
 - ❖ Non-jetty nature \Rightarrow **free from JES**
 - ❖ Most sensitive to α_s^{FSR} , so a better “tune” reduces the theoretical uncertainty of top mass in *B*-hadron observables
 - ❖ α_s^{FSR} should be **constrained at 1-2% level**, while the others at 10-20% to achieve $\sim 0.5\%$ precision in m_t ($\alpha_s^{\text{FSR}} \rightarrow r_b \rightarrow \dots$)
 - ❖ Parameters should/can be **constrained/tuned by calibration observables probing various aspects**
- ❑ The same exercises with HERWIG is available.





Thank you!



Back-up

“Tuning” of PYTHIA8 Parameters

A study of the sensitivity to the PYTHIA8 parton shower parameters of $t\bar{t}$ production measurements in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS experiment at the LHC

The ATLAS Collaboration

Abstract

Various measurements of $t\bar{t}$ observables, performed by the ATLAS experiment in pp collisions at $\sqrt{s} = 7$ TeV, are used to constrain the initial- and final-state radiation parameters of the PYTHIA8 Monte Carlo generator. The resulting tunes are compared to previous tunes to the Z boson transverse momentum at the LHC, and to the LEP event shapes in Z boson hadronic decays. Such a comparison provides a test of the universality of the parton shower model. The tune of PYTHIA8 to the $t\bar{t}$ measurements is applied to the next-to-leading-order generators MadGraph5_aMC@NLO and PowHEG, and additional parameters of these generators are tuned to the $t\bar{t}$ data. For the first time in the context of parton shower tuning in Monte Carlo simulations, the correlation of the experimental uncertainties has been used to constrain the parameters of the Monte Carlo models.



B-hadron Decay

□ Fully reconstructible with tracks

J/ψ modes

$$b \xrightarrow{\text{few } 10^{-3}} J/\psi + X \xrightarrow{10^{-1}} \ell^+ \ell^- + X$$

- $B_s^0 \rightarrow J/\psi \phi \rightarrow \mu^- \mu^+ K^- K^+ \quad (1106.4048) \quad B^0 \rightarrow J/\psi K_s^0 \rightarrow \mu^- \mu^+ \pi^- \pi^+ \quad (1104.2892)$
- $B^+ \rightarrow J/\psi K^+ \rightarrow \mu^- \mu^+ K^+ \quad (1101.0131, 1309.6920) \quad \Lambda_b \rightarrow J/\psi \Lambda \rightarrow \mu^- \mu^+ p \pi^- \quad (1205.0594)$

D modes

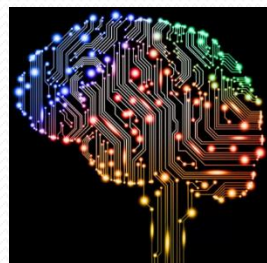
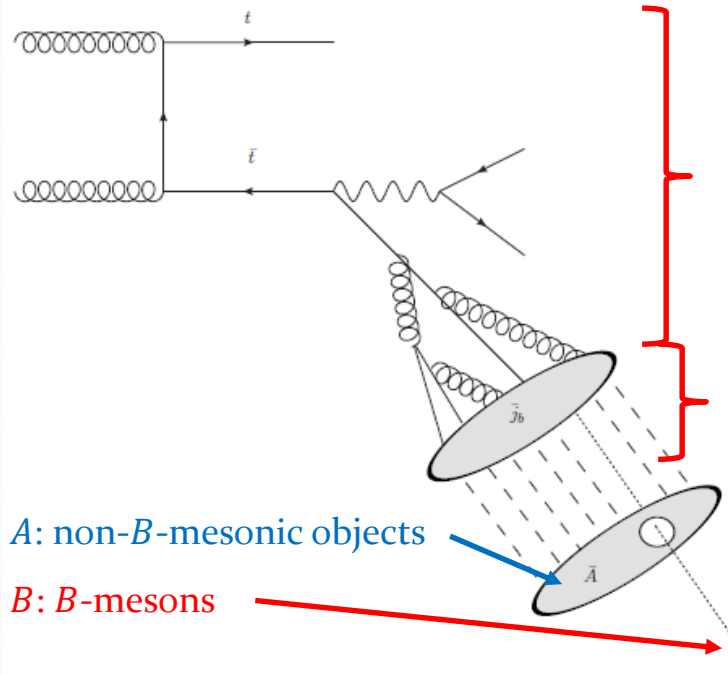
$$\text{➤ } B^0 \xrightarrow{3 \times 10^{-3}} D^- \pi^+ \xrightarrow{10^{-2}} K_s^0 \pi^- \pi^+, B^0 \xrightarrow{3 \times 10^{-3}} D^- \pi^+ \xrightarrow{10^{-2}} K^- \pi^+ \pi^- \pi^+,$$

$$B^0 \xrightarrow{3 \times 10^{-3}} D^- \pi^+ \xrightarrow{3 \times 10^{-2}} K_s^0 \pi^+ \pi^- \pi^+$$

$$\text{➤ } B^- \xrightarrow{5 \times 10^{-3}} D^0 \pi^- \xrightarrow{4 \times 10^{-2}} K^- \pi^+ \pi^-, B^- \xrightarrow{5 \times 10^{-3}} D^0 \pi^- \xrightarrow{2 \times 10^{-2}} K^{*-}(892) \pi^+ \pi^- \rightarrow K_s^0 \pi^+ \pi^- \pi^+,$$

$$B^- \xrightarrow{5 \times 10^{-3}} D^0 \pi^- \xrightarrow{6 \times 10^{-3}} K_s^0 \rho^0 \pi^-, B^- \xrightarrow{5 \times 10^{-3}} D^0 \pi^- \xrightarrow{5 \times 10^{-3}} K^- \pi^+ \rho^0 \pi^-$$

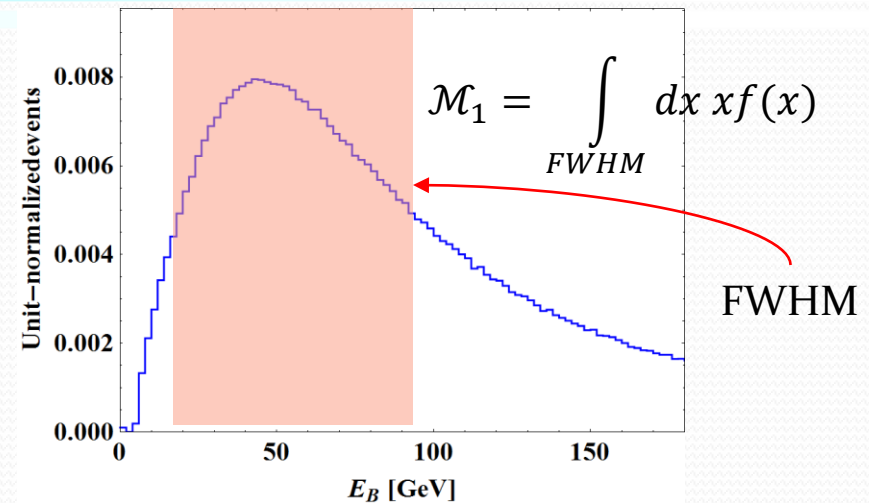
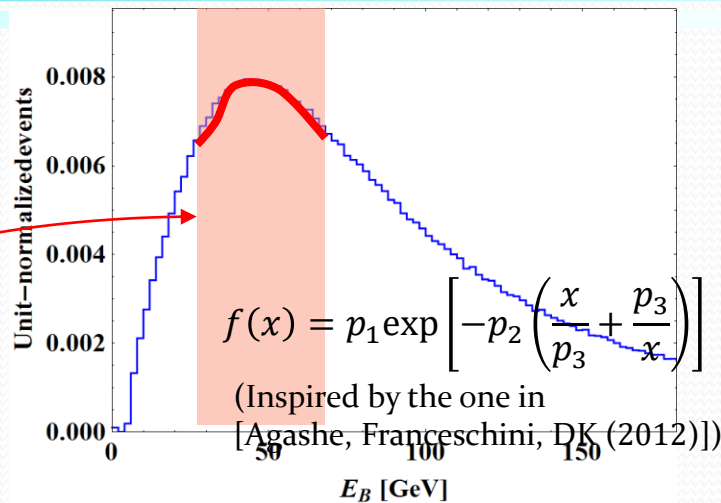
Strategy in a Nut-shell



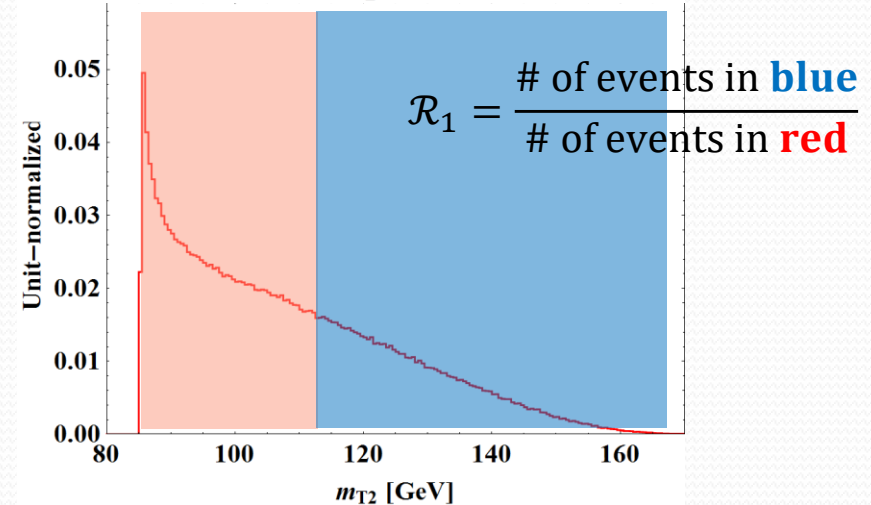
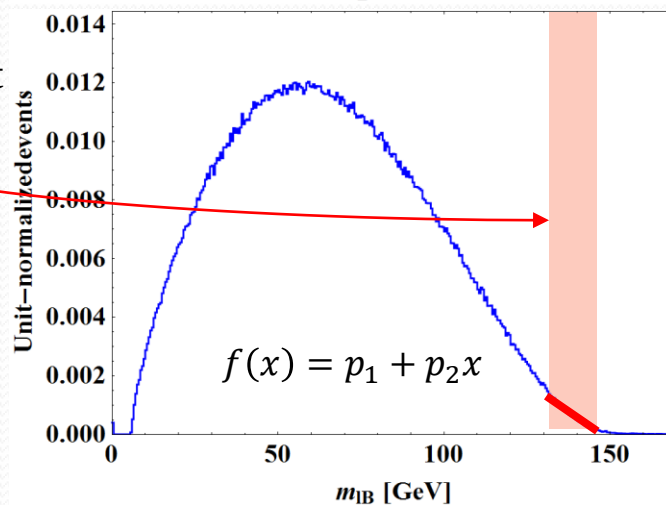
- For a given input top mass,
- 1) set relevant parameters (next slide),
 - 2) generate, shower, and hadronize leptonic $t\bar{t}$ events using PYTHIA 8.2.19,
 - 3) find anti- k_t jets using FastJet,
 - 4) find jets containing a B -hadron as a constituent, and extract its information,
 - 5) evaluate various B -hadron observables/ calibration variables along with (sometimes) leptons: Mellin moments, peak/endpoint,
 - 6) Correlate them with input top masses and find sensitivity measures (defined later),
 - 7) Repeat 1) through 6) for other parameter sets

Measurements: Peak/Endpoint, Mellin Moment

Full width
at $\frac{3}{4}$ height
($\chi^2/\text{dof} \sim 1$)



$\sim 5\%$ from
the endpoint
($\chi^2/\text{dof} \sim 1$)



B-hadron Observables

Observable	Mellin moment (\mathcal{M}_1)	Peak/Endpoint	Features
E_B	V	V (i.e., peak)	<ul style="list-style-type: none"> Expecting inheritance of “invariance” property of the energy-peak in the b-jet energy spectrum
$E_{B_1} + E_{B_2}$	V		<ul style="list-style-type: none"> Two B-meson tagging required
$P_{T,B}$	V		
$P_{T,B_1} + P_{T,B_2}$	V		<ul style="list-style-type: none"> Two B-meson tagging required
$m_{B\ell}$	V	V	<ul style="list-style-type: none"> True pairing (theory-level) Experimental observable paring: the smaller in each combination
m_{T2}	V (\mathcal{R}_1 for (B) subsystem)	V	<ul style="list-style-type: none"> (B) and ($B\ell$) subsystems True assignment (theory-level) for the ($B\ell$) subsystems Experimental observable paring for the ($B\ell$) subsystems: the smaller of the two possible assignments Different ISR and MET definitions
$m_{T2,\perp}[1]$	V (\mathcal{R}_1 for (B) subsystem)	V	<ul style="list-style-type: none"> ISR-free observables (B) and ($B\ell$) subsystems Different ISR and MET definitions

[1]: K. Matchev and M. Park (2009)

Differential Constraining

