

**V_{xb} 2009
SLAC
29-31 October 2009**

Update of the Unitarity Triangle Analysis (UTA):
(on behalf of the  Collaboration)

www.utfit.org

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Status of the UTA within the Standard Model (SM)

→ high precision and global success (but few tensions: $\varepsilon_K \leftrightarrow \sin(2\beta)$, $\text{Br}(B \rightarrow \tau \nu)$)

Status of the UTA beyond the SM

→ hint of New Physics (NP) at $\sim 2.9 \sigma$ in the B_s system

→ bounds on the NP scale from the effective field theory analysis

The UTA within the Standard Model



The experimental constraints:

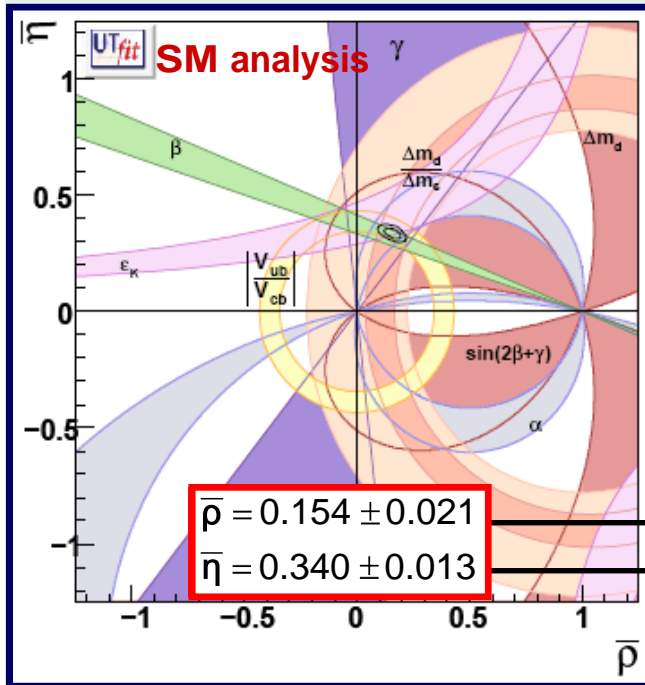
$$\epsilon_K, \Delta m_d, \left| \frac{\Delta m_s}{\Delta m_d} \right|, \left| \frac{V_{ub}}{V_{cb}} \right|$$

relying on theoretical calculations
of hadronic matrix elements

$$\sin 2\beta, \cos 2\beta, \alpha, \gamma(2\beta + \gamma)$$

independent from theoretical
calculations of hadronic parameters

overconstrain the CKM parameters consistently



~14%

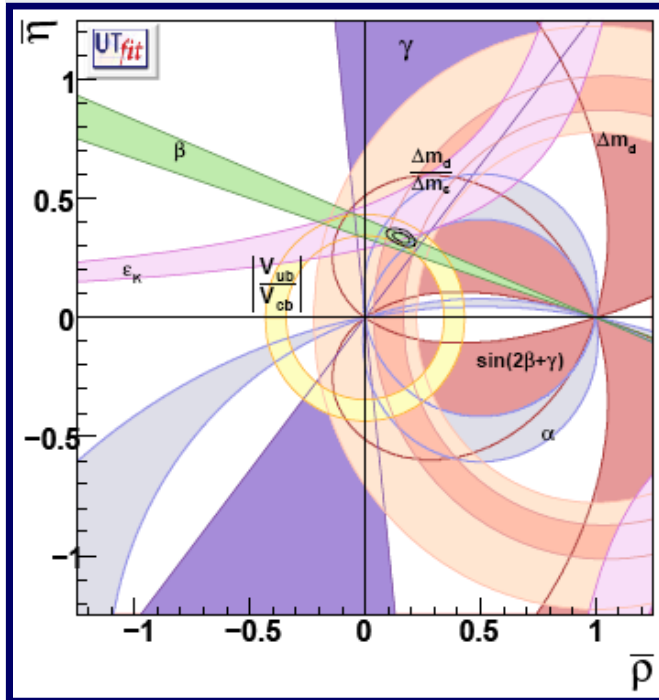
~4%

The UTA has established that the CKM matrix is the dominant source of flavour mixing and CP violation



From a closer look

UTfit, 0909.5065



Update of the UTA after summer (after the EPS):
Inclusion of the corrections to ϵ_K pointed out by
A.J.Buras and D.Guadagnoli (0805.3887)

→ **decrease of the SM prediction of ϵ_K by ~8%**

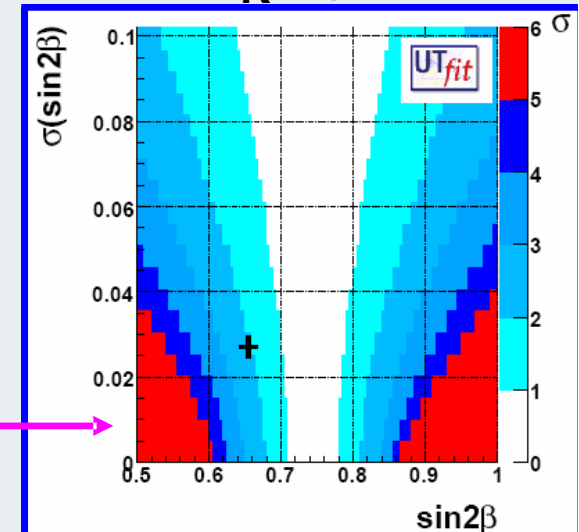
$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12}^K)}{\Delta M_K} + \xi \right)$$

$$\xi = \frac{\text{Im}A_0}{\text{Re}A_0} \rightarrow \text{decreases } \epsilon_K \text{ by } \sim 5\%$$

$$\phi_\epsilon = (43.51 \pm 0.05)^\circ \rightarrow \text{decreases } \epsilon_K \text{ by } \sim 3\%$$

The inclusion of the **corrections** to ϵ_K generates some **tension** between the experimental constraints: ϵ_K and $\sin(2\beta)$

The **indirect determination** of $\sin(2\beta)$ turns out to be at **~2.0 σ** from the **experimental measurement**



Due to many experimental constraints
the UT turns out to be overconstrained

The UTA can be used to improve the knowledge
of some inputs: e.g. $|V_{ub}|$

Input value in the UTA: $(36.7 \pm 2.1) \cdot 10^{-4}$
Output value from the UTA: $(35.2 \pm 1.1) \cdot 10^{-4}$

HFAG
Inclusive-exclusive
average by V.Lubicz
and C.T. (0807.4605)

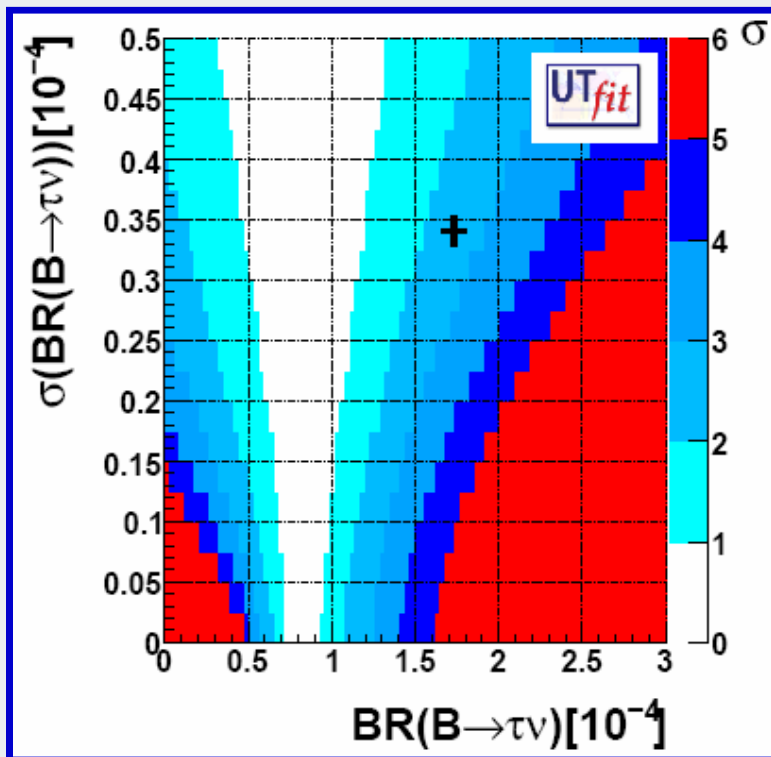
[unquenched lattice results from
FNAL/MILC 04 and HPQCD 06
+ QCD sum rules (Duplancic et al. 08)]

The improved (by other constraints) determination of $|V_{ub}|$
(and f_B) reduces the uncertainty on the SM prediction of
 $\text{Br}(B \rightarrow \tau \nu)$ (from 24% to 13%) [UTfit, 0908.3470]

$$BR(B \rightarrow \tau \nu) = \frac{G_F^2 m_B m_\tau^2}{8\pi} \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$

$B \rightarrow \tau \nu$

$BR(B \rightarrow \tau \nu)_{SM} = (0.84 \pm 0.11) \cdot 10^{-4}$ turns out to be
larger by $\sim 2.5 \sigma$ than the experimental value
(BaBar+Belle, 2006-2008) $BR(B \rightarrow \tau \nu)_{exp} = (1.73 \pm 0.34) \cdot 10^{-4}$



By generalizing the analysis,
allowing for NP effects in $\Delta F=2$ processes,
the deviation remains ($\sim 2.2 \sigma$)

**N.B. a charged Higgs could not easily
explain the enhancement**
(due to other constraints, mainly $b \rightarrow s \gamma$)

$$\frac{BR(B \rightarrow \tau \nu)_{2HDM}}{BR(B \rightarrow \tau \nu)_{SM}} = \left(1 - \tan^2 \beta \frac{m_B^2}{m_{H^+}^2} \right)^2$$

Some hadronic quantities can be extracted from the (overconstraint) UTA and compared to Lattice calculations *



* assuming the SM validity!!!

Extracting them as free parameters from the UTA: Averaging accurate Lattice results:

(UTfit, update of hep-ph/0606167) (V.Lubicz, C.T., 0807.4605)

$$\hat{B}_K^{UT} = 0.81 \pm 0.08$$

$$f_{B_s} \sqrt{\hat{B}_{B_s}^{UT}} = 265 \pm 4 \text{ MeV}$$

$$\xi^{UT} = 1.25 \pm 0.06$$

$$\hat{B}_K^{LAT} = 0.75 \pm 0.07$$

$$f_{B_s} \sqrt{\hat{B}_{B_s}^{LAT}} = 270 \pm 30 \text{ MeV}$$

$$\xi^{LAT} = 1.21 \pm 0.04$$

• Recent unquenched Lattice calculations point towards (~3-4%) smaller values

(see new averages by J.Lahio, E.Lunghi, R.Van de Water, 0910.2928) $\hat{B}_K = 0.73 \pm 0.03$

Remarkable agreement:

- Additional evidence of the SM success in describing flavour physics
- Reliability of Lattice QCD

N.B. a smaller value of B_K increases the tension $\varepsilon_K - \sin(2\beta)$

An update of the Lattice averages is in program (~once per year)

Update after summer (after the EPS): decrease of the SM prediction of ε_K by $\sim 3\%$ [due to $\phi_\varepsilon = 43.51^\circ$, while $\xi = 0$ (difficult to estimate from ε'/ε beyond MFV)]

Model-independent UTA: bounds on deviations from the SM (+CKM)

- Parametrize generic NP in $\Delta F=2$ processes, in all sectors
- Use all available experimental info
- Fit simultaneously the CKM and NP parameters

NP contributions in the mixing amplitudes:

$$H^{\Delta F=2} = m + \frac{i}{2} \Gamma \quad A = m_{12} = \langle M | m | \bar{M} \rangle \quad \Gamma_{12} = \langle M | \Gamma | \bar{M} \rangle$$

K mixing amplitude (2 real parameters):

$$\text{Re} A^K = C_{\Delta m_K} \text{Re} A_K^{SM} \quad \text{Im} A_K = C_{\xi_K} \text{Im} A_K^{SM}$$

B_d and B_s mixing amplitudes (2+2 real parameters):

$$A_q e^{2i\phi_q} = C_{B_q} e^{2i\phi_{B_q}} A_q^{SM} e^{2i\phi_q^{SM}} = \left(1 + \frac{A_q^{NP}}{A_q^{SM}} e^{2i(\phi_q^{NP} - \phi_q^{SM})} \right) A_q^{SM} e^{2i\phi_q^{SM}}$$

SM	→	SM+NP
$(V_{ub}/V_{cb})^{SM}$ γ^{SM}	tree level	$(V_{ub}/V_{cb})^{SM}$ γ^{SM}
β^{SM} α^{SM} Δm_d	Bd Mixing	$\beta^{SM} + \phi_{Bd}$ $\alpha^{SM} - \phi_{Bd}$ $C_{Bd} \Delta m_d$
Δm_s^{SM} $-\beta_s^{SM}$	Bs Mixing	$C_{Bs} \Delta m_s^{SM}$ $-\beta_s^{SM} + \phi_{Bs}$
ε_K^{SM} Δm_K^{SM}	K Mixing	$C_{\varepsilon_K} \varepsilon_K^{SM}$ $C_{\Delta m_K} \Delta m_K^{SM}$

For K - \bar{K} mixing, the NP parameters are found
in agreement with the SM expectations
(some effect can be seen once B_K is updated to a smaller value)

For B_d - \bar{B}_d mixing, the mixing phase ϕ_{B_d} is found
 1.5σ away from the SM expectation
(reflecting a slight tension between $\sin(2\beta)$ and $|V_{ub}/V_{cb}|$)

In 2008 both CDF and DØ published the tagged time-dependent angular analysis of $B_s \rightarrow J/\psi \phi$



2D likelihood ratio for $\Delta\Gamma$ and ϕ_s
2-fold ambiguity present, no assumption on the strong phases

arXiv:0712.2397



7-parameter fit + correlation matrix or 1D likelihood profiles of $\Delta\Gamma$ and ϕ_s
2-fold ambiguity removed using strong phases from $B \rightarrow J/\psi K^* + SU(3) + ?$

arXiv:0802.2255

At ICHEP '08:



1. DØ released the 2D likelihood scan w/o assumptions on the strong phases

2. New measurement of A_{SL}^s , now $A_{SL}^s = (-0.20 \pm 1.19) \%$

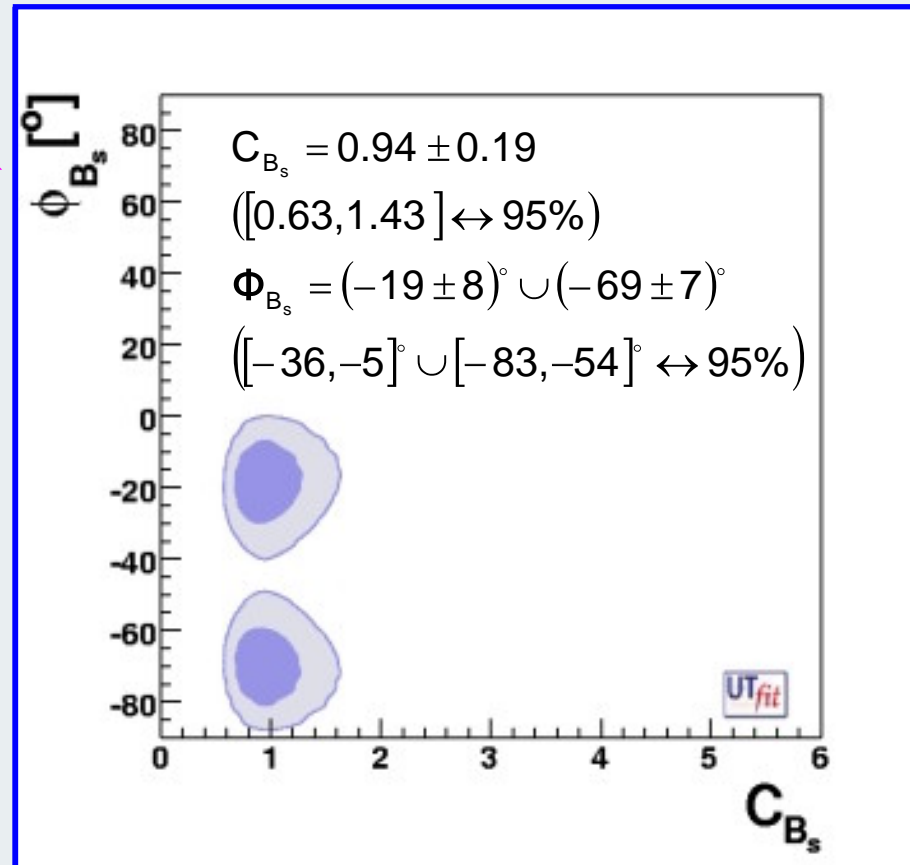
All the **exp. info** have been **combined**
(UPDATE OF UTfit Coll. 0803.0659)

**UPDATED UTfit analysis:
SM @ 2.9 σ**



HFAG: 2.2 σ (0808.1297)
CKMfitter: 2.5 σ (0810.3139)

**More than 2 σ deviation for
every statistical approach!**



Enlarged data sample: 1.35 fb⁻¹ -> 2.8 fb⁻¹
opposite-side tagging only (equivalent to ~2 fb⁻¹)

CDF analysis: SM compatibility 15%(1.5 σ) -> 7%(1.8 σ)



(New CDF likelihood not available yet)

If this deviation is confirmed 
NP with new sources of flavour violation is required

- Minimal Flavour Violation (MFV) models are ruled out (including the simplest MSSM)
- A clear pattern of flavour violation in NP emerges:
 - 1 ↔ 2: suppressed
 - 1 ↔ 3: $\leq O(10\%)$
 - 2 ↔ 3: $O(1)$
- This pattern can be explained by nonabelian flavour symmetries and in some SUSY-GUTs

**Flavour Physics is highly sensitive to NP:
The Effective Field Theory (EFT) analysis**



The mixing amplitudes $A_q e^{2i\phi_q} = \langle \bar{M}_q | H_{eff}^{\Delta F=2} | M_q \rangle$

$$H_{eff}^{\Delta B=2} = \sum_{i=1}^5 C_i(\mu) Q_i(\mu) + \sum_{i=1}^3 \tilde{C}_i(\mu) \tilde{Q}_i(\mu)$$

The high scale coefficients $C_i(\Lambda)$ can be extracted from the data

(switching on one operator per time)

$$Q_1 = \bar{q}_L^\alpha \gamma_\mu b_L^\alpha \bar{q}_L^\beta \gamma^\mu b_L^\beta \quad (\text{SM/MFV})$$

$$Q_2 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_R^\beta b_L^\beta$$

$$Q_3 = \bar{q}_R^\alpha b_L^\beta \bar{q}_R^\beta b_L^\alpha$$

$$Q_4 = \bar{q}_R^\alpha b_L^\alpha \bar{q}_L^\beta b_R^\beta$$

$$Q_5 = \bar{q}_R^\alpha b_L^\beta \bar{q}_L^\beta b_R^\alpha$$

$$\tilde{Q}_1 = \bar{q}_R^\alpha \gamma_\mu b_R^\alpha \bar{q}_R^\beta \gamma^\mu b_R^\beta$$

$$\tilde{Q}_2 = \bar{q}_L^\alpha b_R^\alpha \bar{q}_L^\beta b_R^\beta$$

$$\tilde{Q}_3 = \bar{q}_L^\alpha b_R^\beta \bar{q}_L^\beta b_R^\alpha$$

$$C_i(\Lambda) = \frac{LF_i}{\Lambda^2} \Rightarrow \Lambda = \sqrt{\frac{LF_i}{C_i(\Lambda)}}$$

Tree/strong inter. NP: $L \sim 1$
Perturbative NP: $L \sim \alpha_s^2, \alpha_w^2$

MFV	next-to-MFV	generic
- $F_1 = F_{SM} \sim (V_{tq} V_{tb}^*)^2$	- $ F_i \sim F_{SM}$	- $ F_i \sim 1$
- $F_{i \neq 1} = 0$	- arbitrary phases	- arbitrary phases

Present lower bound on the NP scale

From B and K data (TeV@95%)

Scenario	strong/tree	α_s loop	α_W loop
MFV	5.5	0.5	0.2
NMFV	62	6.2	2
General	24000	2400	800

* $\Delta F=2$ **chirality-flipping operators** are **RG enhanced** and thus probe **larger NP scale** (that can be pushed beyond the LHC reach)

In the presence of a **NP evidence**,
also an **upper bound** is provided

From the B_s system (TeV@95%)

Scenario	strong/tree	α_s loop	α_W loop
NMFV	35	4	2
General	800	80	30

upper bound \ll lower bound!!



The pattern of **NP flavour couplings** cannot be **SM-like** nor **general**

Data suggest some **hierarchy in NP**, **stronger** than in the **SM** (e.g. some SUSY-GUTs)



We are **looking for increased exp. accuracy** in order to **confirm or exclude a NP effect** in the **B_s system!**

THANKS





BACKUP

The effect of quenching the strange quark has never been seen

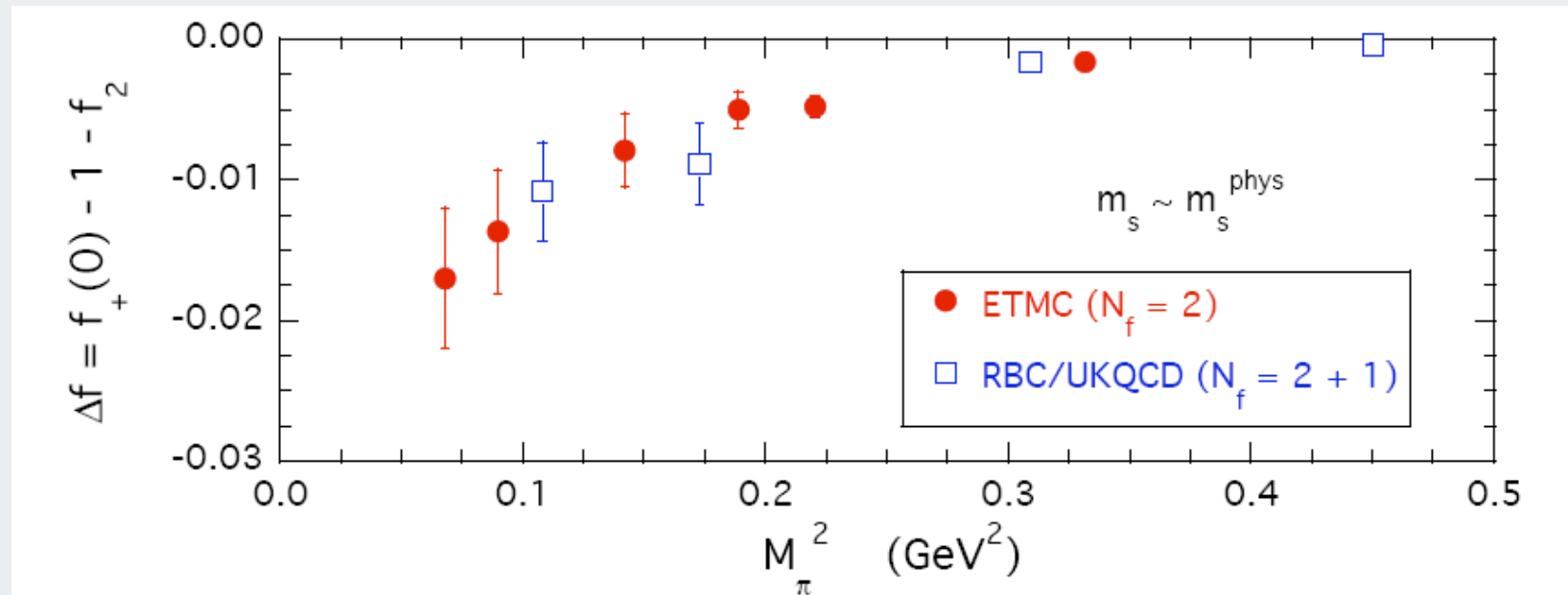
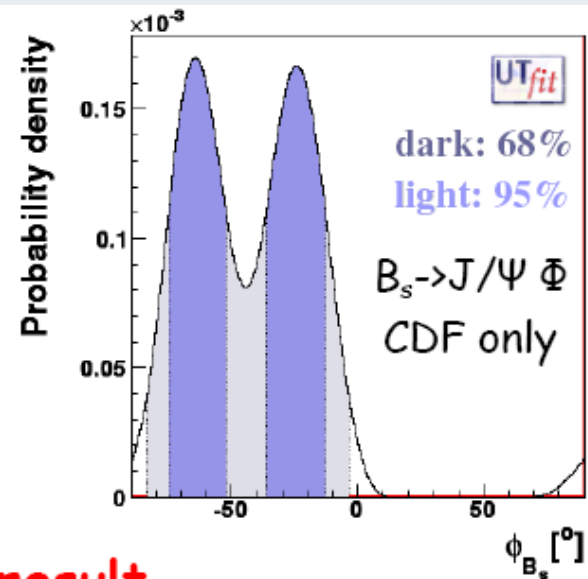
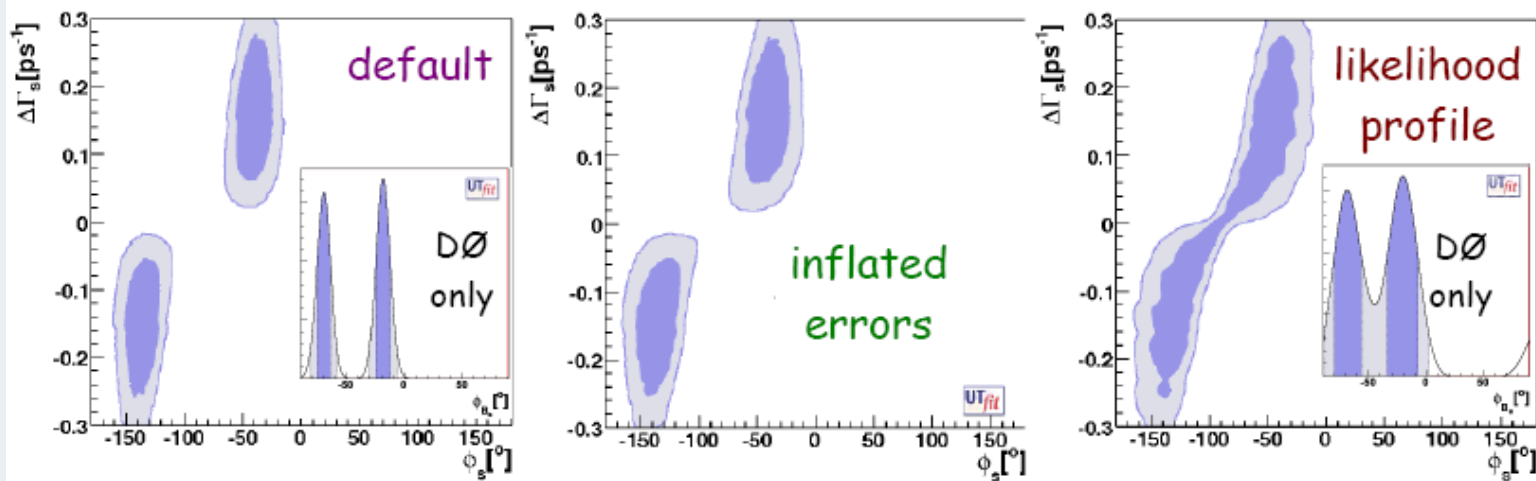


Figure 2: Values of the $\mathcal{O}(p^6)$ term Δf [Eq. (2.6)] obtained by the ETM [8] and RBC/UKQCD [7] Collaborations taking into account the values of the NLO term f_2 appropriate for $N_f = 2$ and $N_f = 2 + 1$.

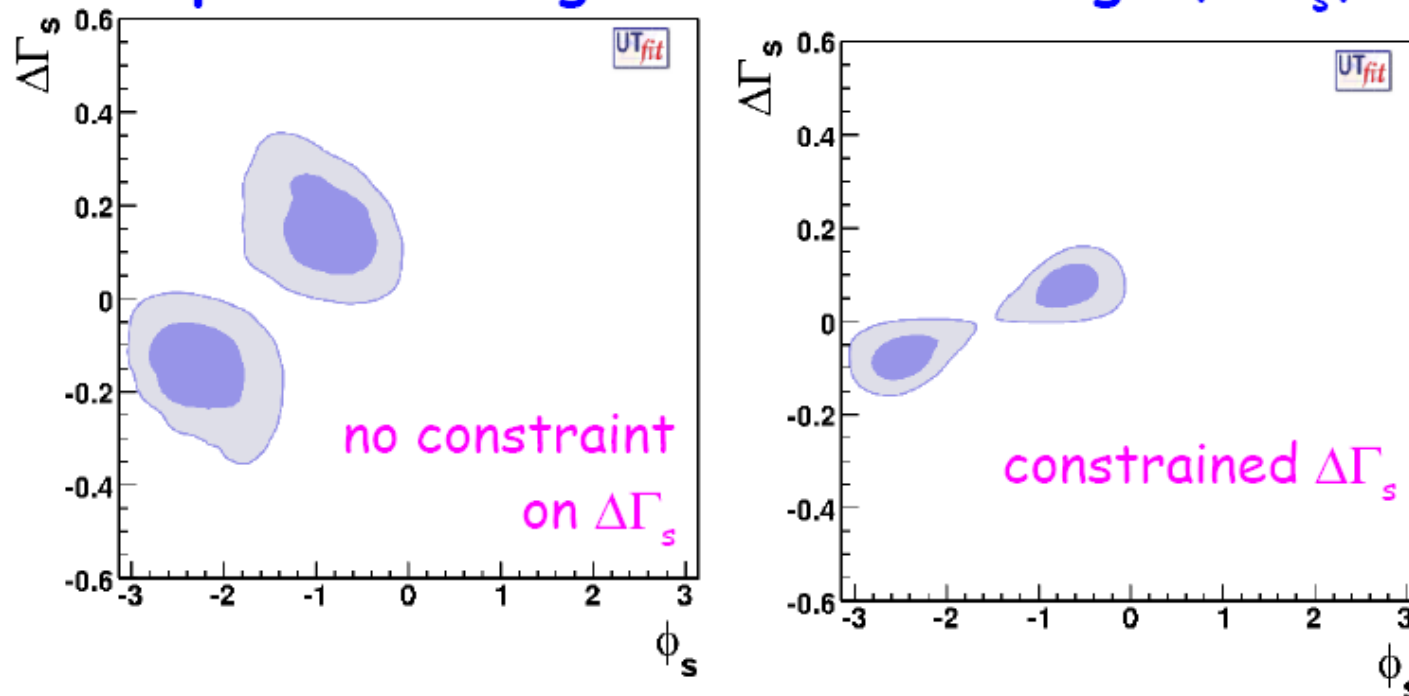
- * gaussian: CDF likelihood+Gaussian $D\emptyset$ result with 2x2 corr. matrix
- * inflated error: as above, but with error inflated to reproduce the 2σ range computed by $D\emptyset$
- * likelihood profile: using the 1D likelihood profiles for ϕ_s and $\Delta\Gamma_s$



ambiguity reintroduced in the $D\emptyset$ result



The θ input for $\Delta\Gamma_s$ is crucial: most of the exp allowed region has a too large $|\Delta\Gamma_s|$

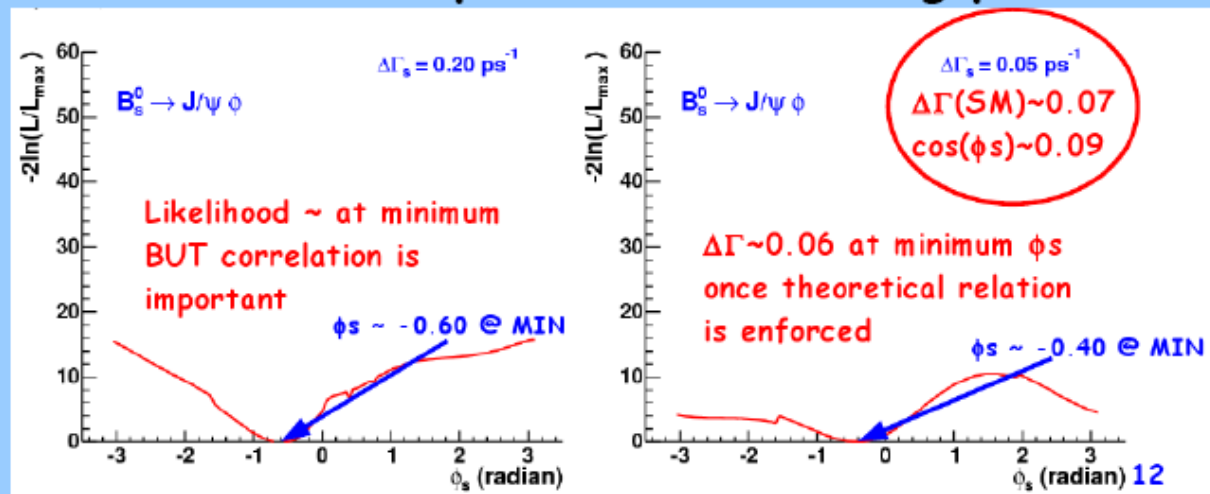


We use a conservative estimate of the SM error and allow NP to enter $\Delta\Gamma_s$ through NP penguins

ICHEP '08 update (i)



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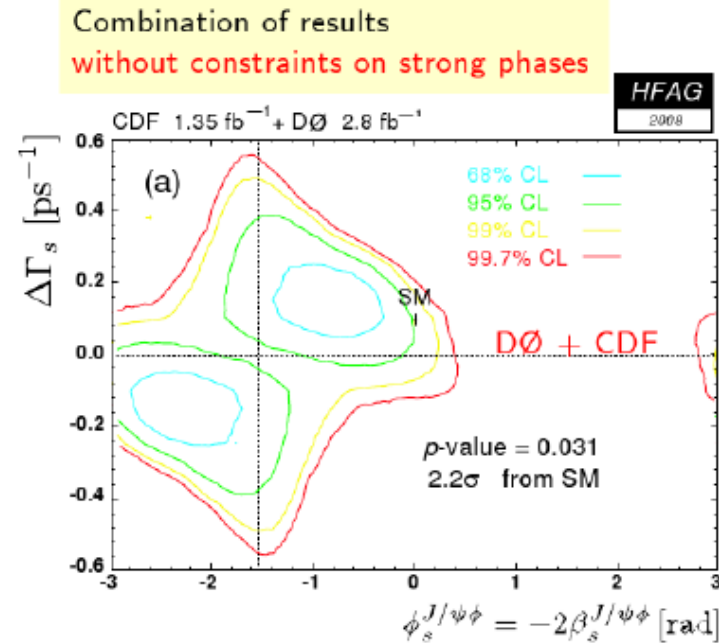
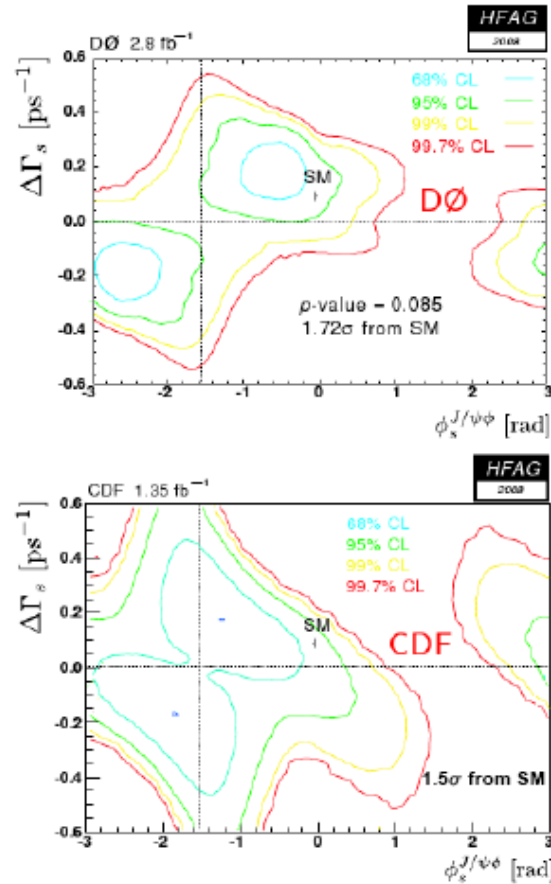
2. New measurement of A_{SL}^s , now $A_{SL}^s = (-0.20 \pm 1.19) \%$



Enlarged data sample: $1.35 \text{ fb}^{-1} \rightarrow 2.8 \text{ fb}^{-1}$
opposite-side tagging only (equivalent to $\sim 2 \text{ fb}^{-1}$)

CDF analysis: SM compatibility 15%(1.5 σ) \rightarrow 7%(1.8 σ)

New HFAG combination



$\phi_s = -2.37^{+0.38}_{-0.27}$ rad, $-0.75^{+0.27}_{-0.38}$ rad

$\Delta\Gamma_s = -0.150^{+0.66}_{-0.59}$ ps⁻¹, $0.150^{+0.59}_{-0.66}$ ps⁻¹

90% C.L. intervals (1-d regions):

$\phi_s \in [-2.85, -1.65], [-1.47, -0.29]$

$\Delta\Gamma_s \in [-0.265, -0.036], [0.036, 0.265]$

New CDF results not included yet

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