

$|V_{ub}|$ **from** $B \rightarrow X \pi l \bar{\nu}$

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with I. Stewart and F. Tackmann

$B \rightarrow X \pi \ell \bar{\nu}$

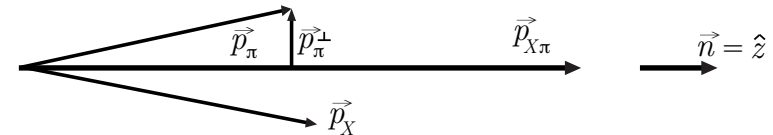
- Inclusive $B \rightarrow X_u \ell \bar{\nu}$: 3 independent kinematic variables
- Semi-inclusive decay, $|X_u\rangle \rightarrow |X\pi\rangle$, pion is identified: 3 more

$$\{p_{X\pi}^+, p_{X\pi}^-, p_{\pi}^+, p_{\pi}^-, E_{\ell}, \phi_{\ell}\}$$

$$p^{\pm} \equiv p^0 \mp p^z$$

$$p_{X\pi}^{\mu} = p_X^{\mu} + p_{\pi}^{\mu}$$

$$m_{X\pi}^2 = p_{X\pi}^{\mu} p_{X\pi\mu} = p_{X\pi}^+ p_{X\pi}^-$$



$$z \equiv \frac{p_{\pi}^-}{p_{X\pi}^-}, \quad 0 \leq z \leq 1$$

- $b \rightarrow c \ell \bar{\nu}$ background: the identified pion helps, we can use events with $m_{X\pi} > m_D$, if z sufficiently large !

Factorization theorem: $|V_{ub}|$ from $B \rightarrow X \pi \ell \bar{\nu}$

Integrate the fully differential decay rate over $E_\ell, \phi_\ell, p_\pi^\pm$ and a large range in $p_{X\pi}^+$, local OPE gives, $\{p_{X\pi}^-, p_\pi^-\} \rightarrow \{m_{X\pi}^2, z\}$:

$$\frac{d^2\Gamma}{dm_{X\pi}^2 dz} = \sum_{i=u, \bar{u}, d, g\dots} |V_{ub}|^2 \int_z^1 \frac{dx}{x} H_{u,i} \left(m_{X\pi}, \frac{z}{x}, \mu \right) D_i^\pi(x, \mu) + \dots$$

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- Hard function H is calculable in perturbation theory:

$$H \propto \delta(m_{X\pi}^2) \delta\left(1 - \frac{z}{x}\right) + H^{1\text{-loop}} + \dots$$

MP, I. Stewart and F. Tackmann, *forthcoming*

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● Fragmentation function $D(z)$:

- non-perturbative but universal, due to factorization theorems for single inclusive hadron production: $e^+ e^- \rightarrow h X, e N \rightarrow e' h X, H_A H_B \rightarrow h X$ at high $p_T \dots$

Collins, Soper and Sterman ...

- $D^\pi(x)$ from fits to (mainly) $e^+ e^- \rightarrow \pi X$ cross sections
- Unprecedented accuracy expected from analysis of B -factory data: few % uncertainty for medium-to-large z

M. Grosse Perdekamp, *private communication*

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Semi-inclusive analysis:

a new independent determination of $|V_{ub}|$ with competitive uncertainties

Subtracting the $b \rightarrow c$ background

● No $B \rightarrow X_c \pi \ell \bar{\nu}$:

$$m_X^2 < m_D^2 \quad \Longrightarrow \quad z > \frac{1}{2} \left(1 - \frac{m_D^2 - m_\pi^2}{m_{X\pi}^2} \right) \quad \text{if } p_\pi^+ \leq p_\pi^-$$

Includes events with $m_{X\pi}^2 > m_D^2$. Always $z \geq m_\pi^2 / m_{X\pi}^2$

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- No $B \rightarrow (X_c \rightarrow X\pi) \ell \bar{\nu}$:

$$m_{X\pi}^2 < m_D^2 \text{ always OK}$$

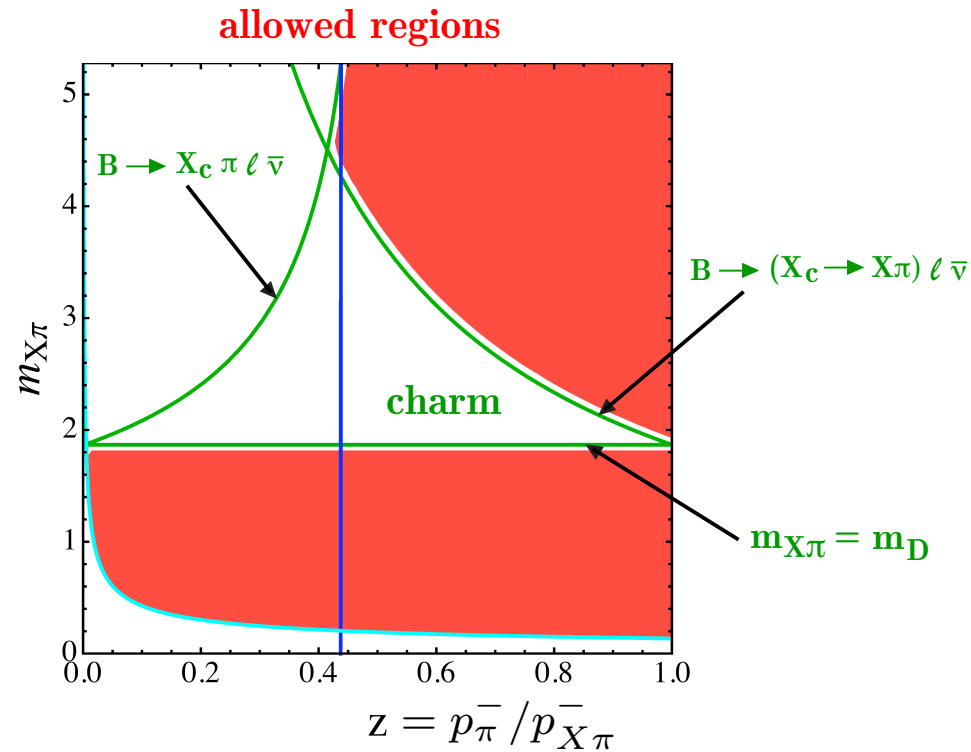
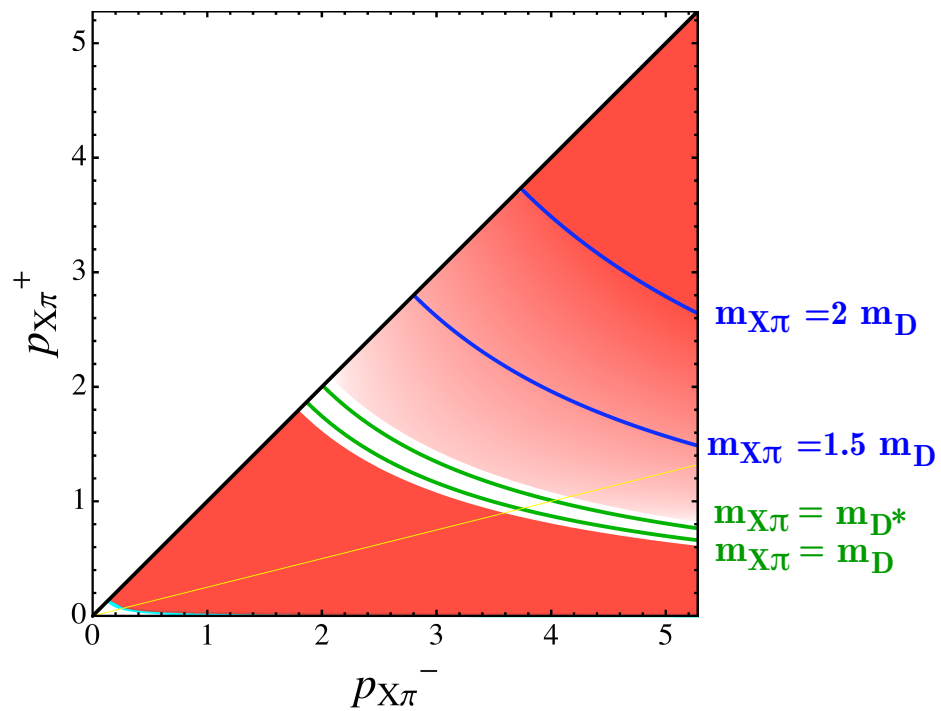
Otherwise, $m_{X\pi}^2 > m_D^2$. In the rest frame of the decaying X_c , it requires

$$p_\pi^{-\text{rest}} > m_D$$

Boosting to the frame where the B -meson decays at rest:

$$\frac{m_D}{\sqrt{p_{X\pi}^- p_{X\pi}^+}} = \frac{m_D}{m_{X\pi}} < z \leq 1$$

Subtracting the $b \rightarrow c$ background



Conclusion:

$$\frac{d\Gamma}{dm_{X_u}^2} \implies \frac{d^2\Gamma}{dm_{X\pi}^2 dz} \implies |V_{ub}|!$$