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# Precision Electroweak Constraints and New Physics

**Jens Erler (IF-UNAM)**

**Searching for Physics Beyond the Standard Model  
Using Charged Leptons**

**COFI, Old San Juan, Puerto Rico  
May 21–25, 2018**



# Outline

- Key Observables and Inputs
- Gauge Couplings at Lower Energies
- Electroweak Fits
- Conclusions

**Key observables and inputs**

# Z pole

- $M_Z = 91.1876 \pm 0.0021$  GeV (error no longer negligible)
- $\Gamma_Z$ ,  $\sigma_{\text{had}}$  and hadronic-to-leptonic BRs provide only  $\alpha_s$  constraints not limited by theory

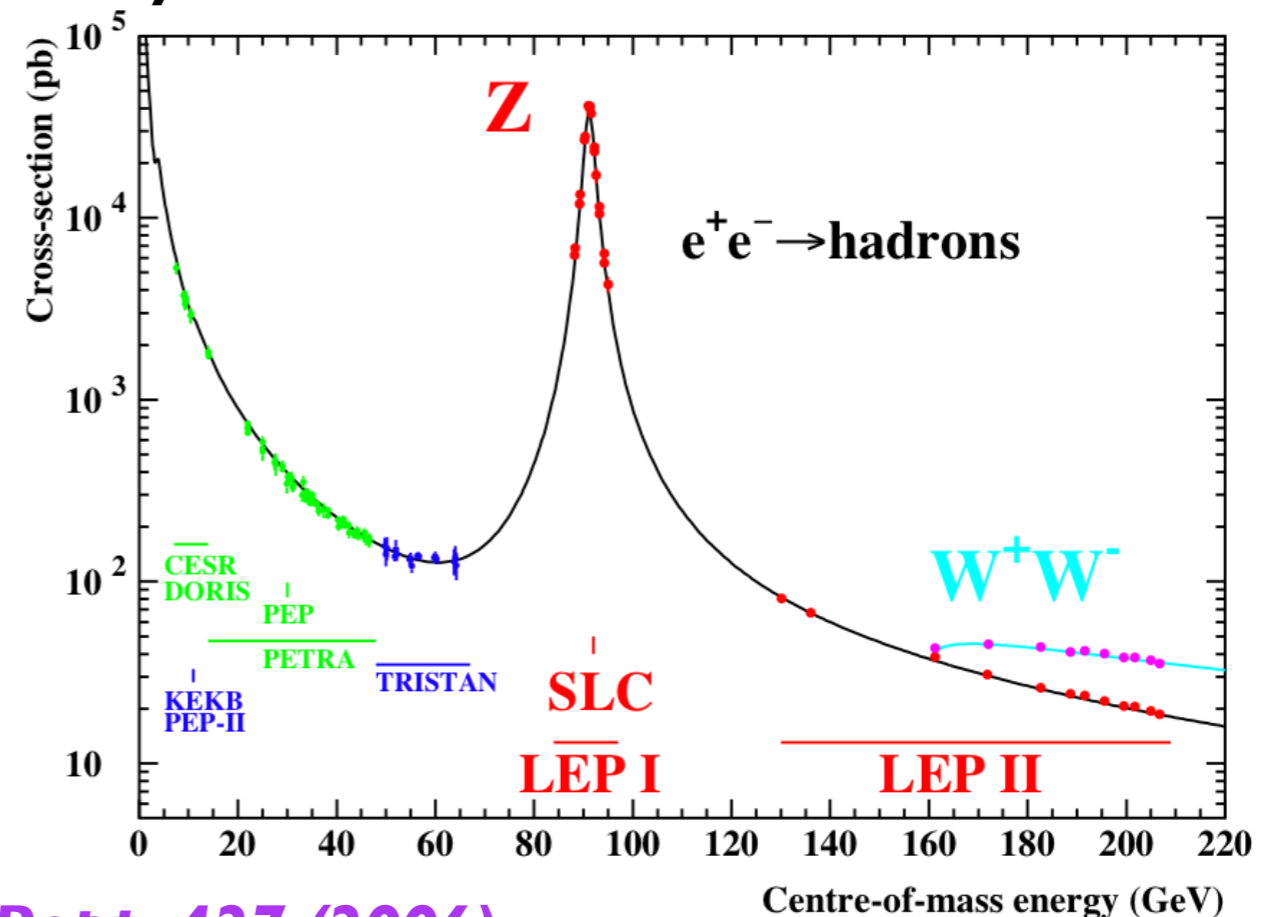
- forward-backward and left-right asymmetries

$$\propto A_e \sim 1 - 4 \sin^2\theta_W(M_Z)$$

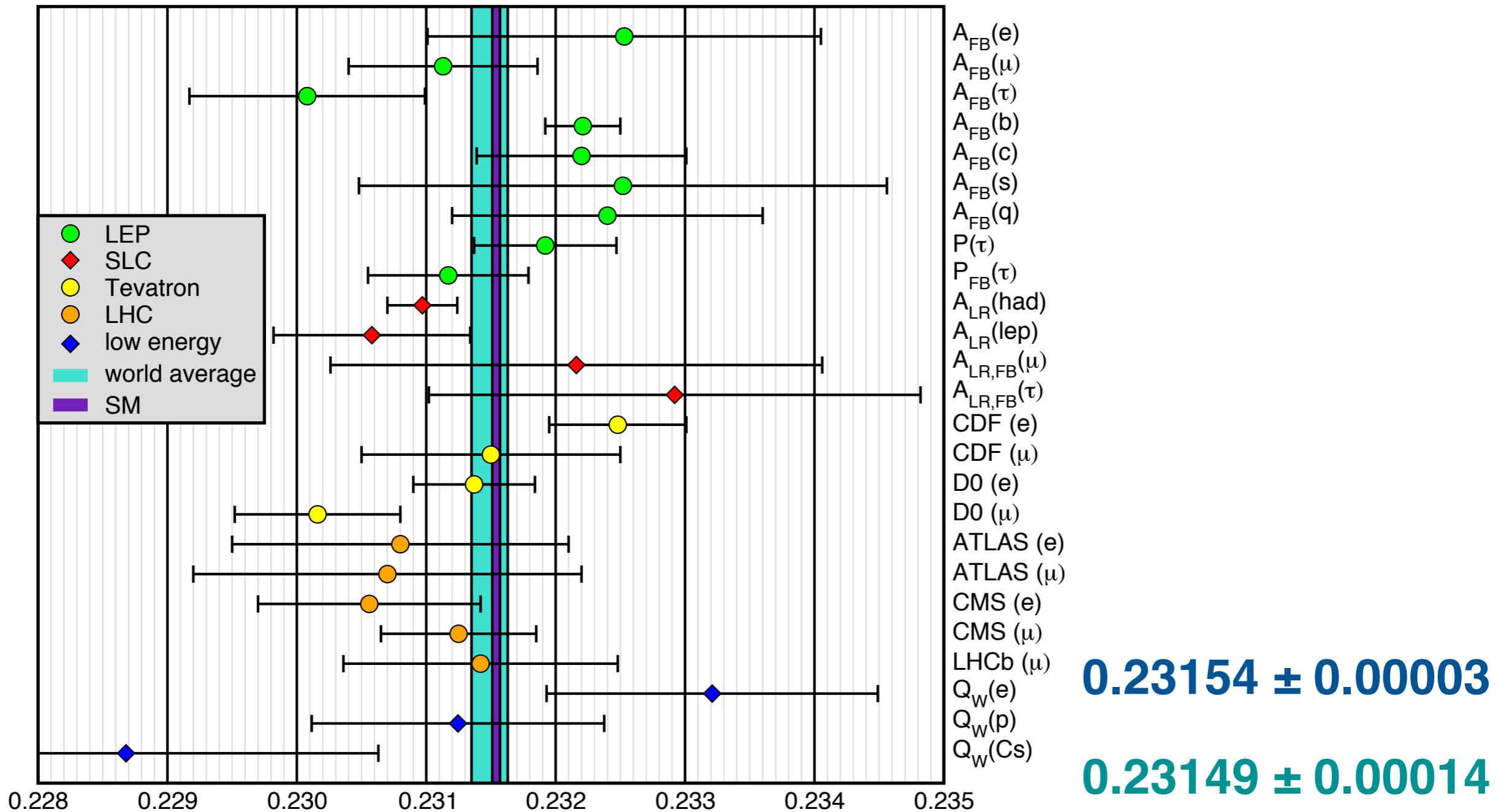
have strong sensitivity to

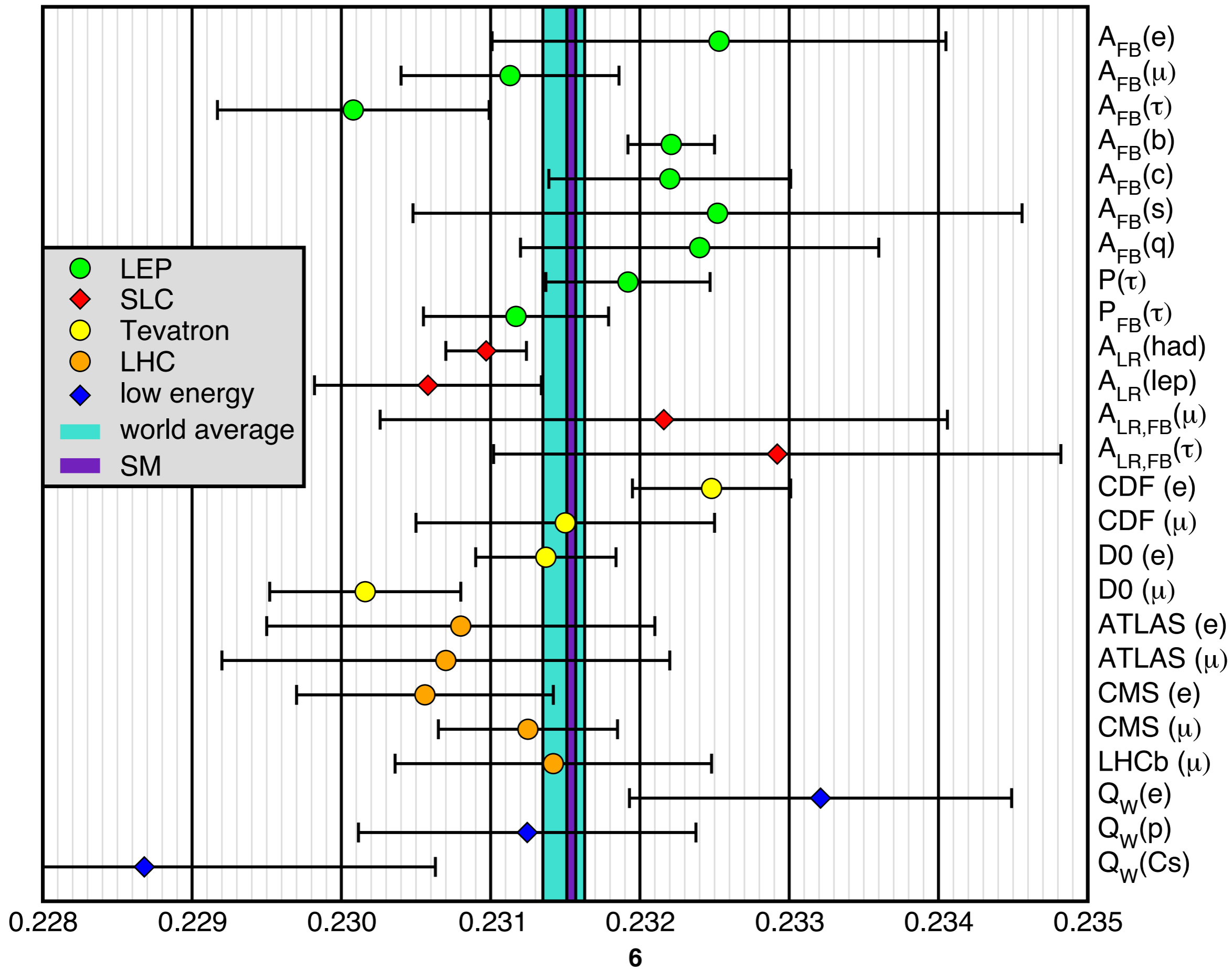
$$\sin^2\theta_W = g'^2 / (g^2 + g'^2)$$

*ALEPH, DELPHI, L3 & OPAL, Phys. Rept. 427 (2006)*

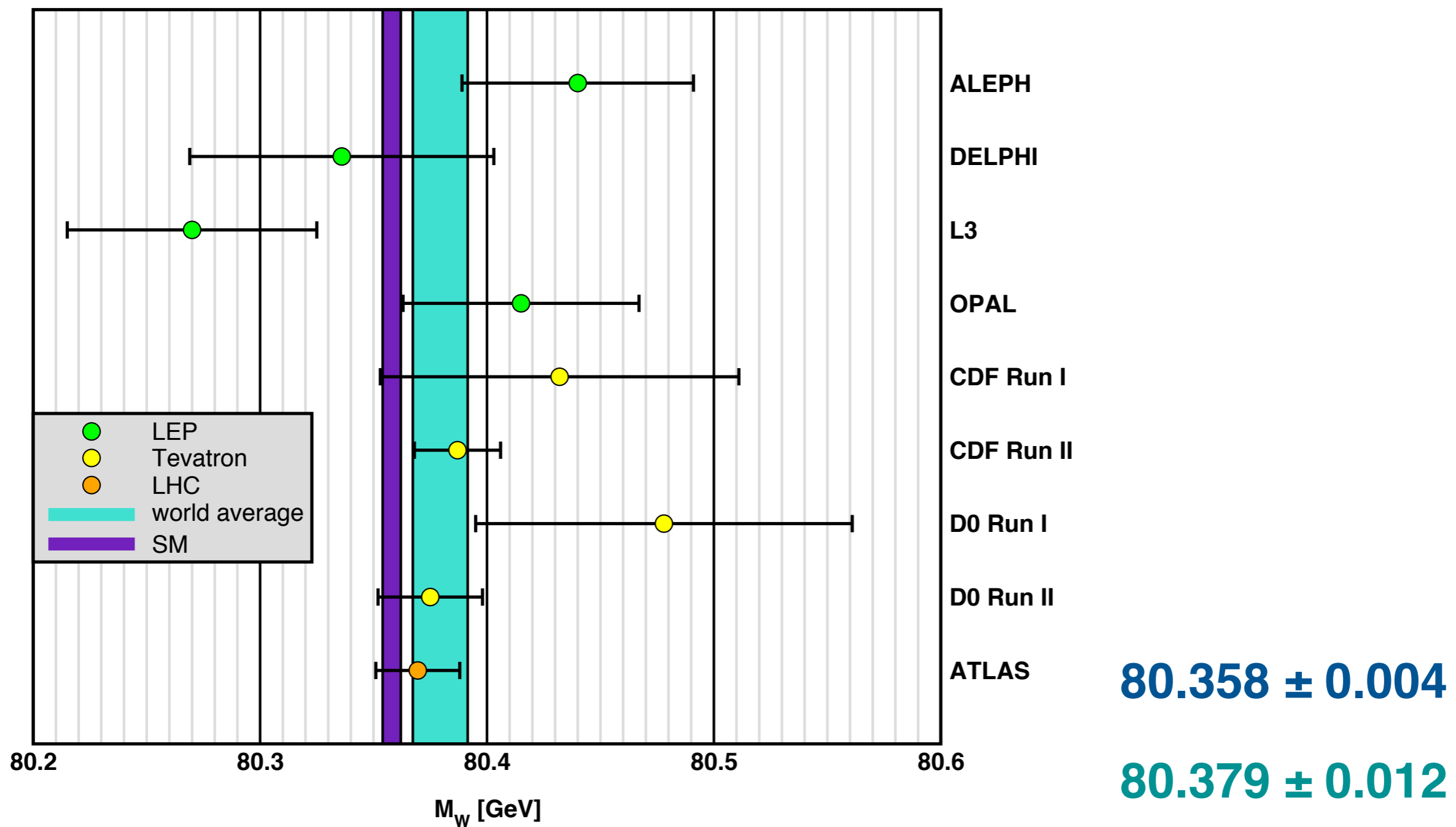


# $\sin^2\theta_W$ measurements





# $M_W$ measurements



# $m_t$ measurements

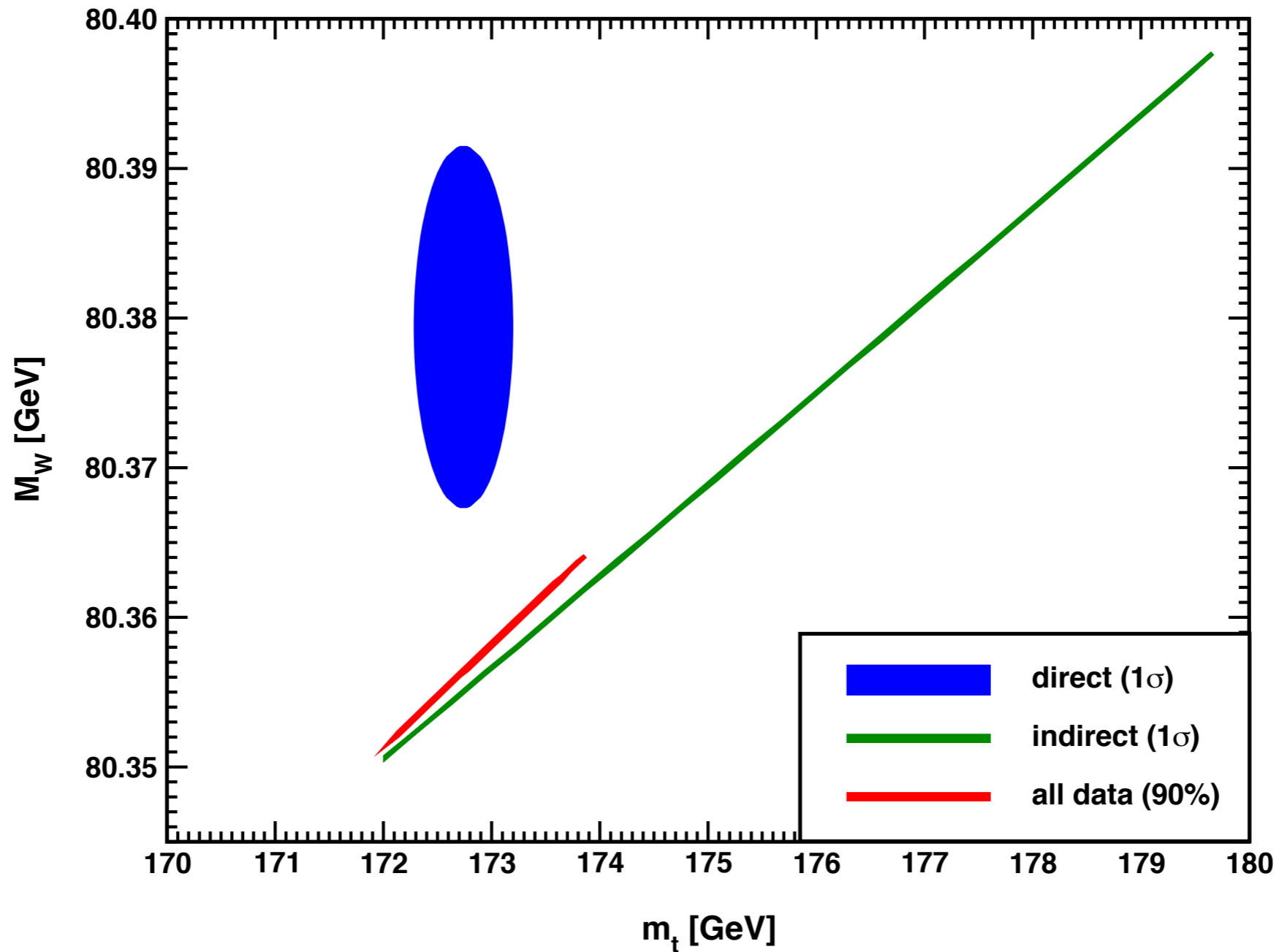
	central	statistical	systematic	total
<b>Tevatron</b>	<b>174.30</b>	<b>0.35</b>	<b>0.54</b>	<b>0.64</b>
<b>ATLAS</b>	<b>172.51</b>	<b>0.27</b>	<b>0.42</b>	<b>0.50</b>
<b>CMS</b>	<b>172.43</b>	<b>0.13</b>	<b>0.46</b>	<b>0.48</b>
<b>CMS Run 2</b>	<b>172.25</b>	<b>0.08</b>	<b>0.62</b>	<b>0.63</b>
<b>grand</b>	<b>172.74</b>	<b>0.11</b>	<b>0.31</b>	<b>0.33</b>

*JE, EPJC 75 (2015)*

- $m_t = 172.74 \pm 0.25_{\text{uncorr.}} \pm 0.21_{\text{corr.}} \pm 0.32_{\text{QCD}} \text{ GeV} = 172.74 \pm 0.46 \text{ GeV}$
- somewhat larger shifts and smaller errors conceivable in the future  
*Butenschoen et al., PRL 117 (2016); Andreassen & Schwartz, JHEP 10 (2017)*
- 2.8  $\sigma$  discrepancy between lepton + jet channels from DØ and CMS Run 2
- indirectly from EW fit:  $m_t = 176.4 \pm 1.8 \text{ GeV}$  *Freitas & JE (PDG 2018)*

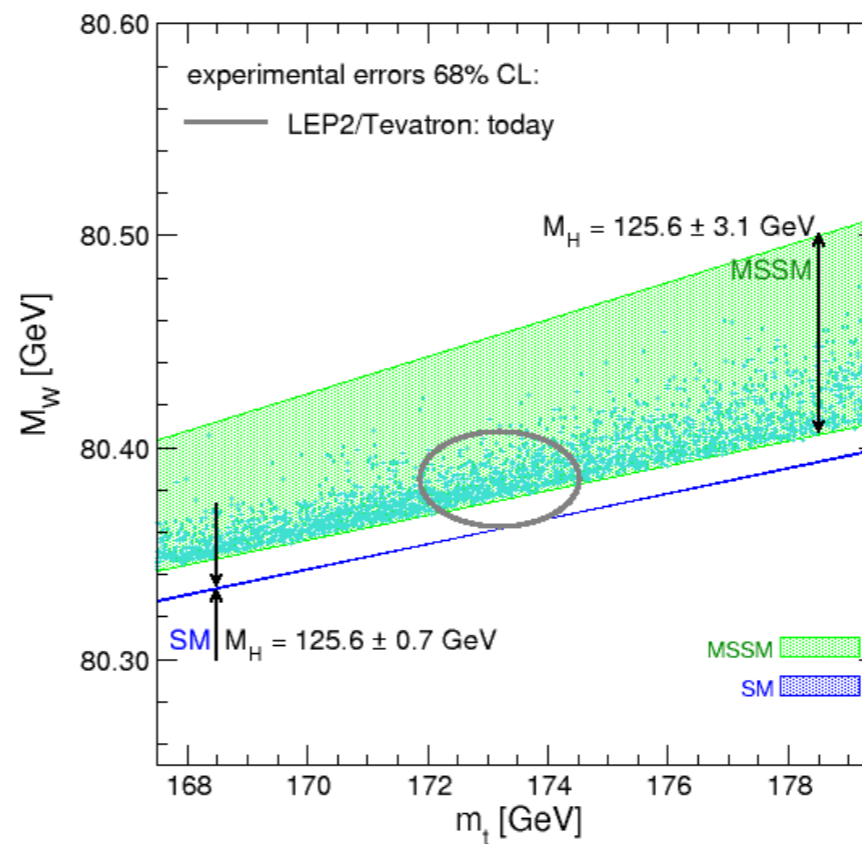
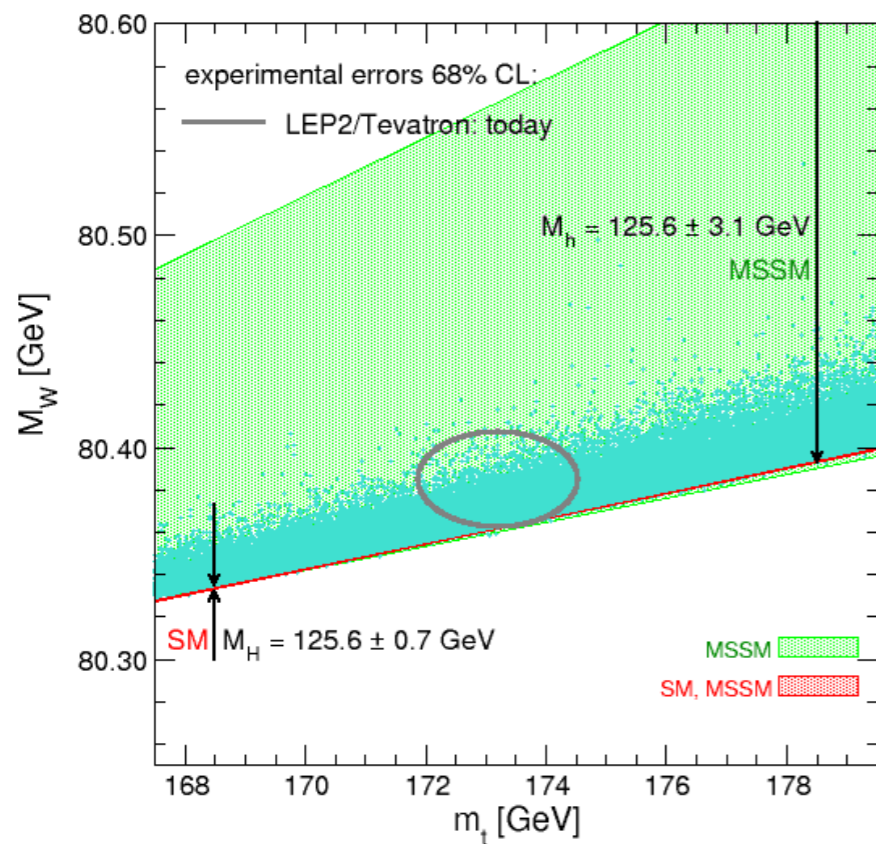
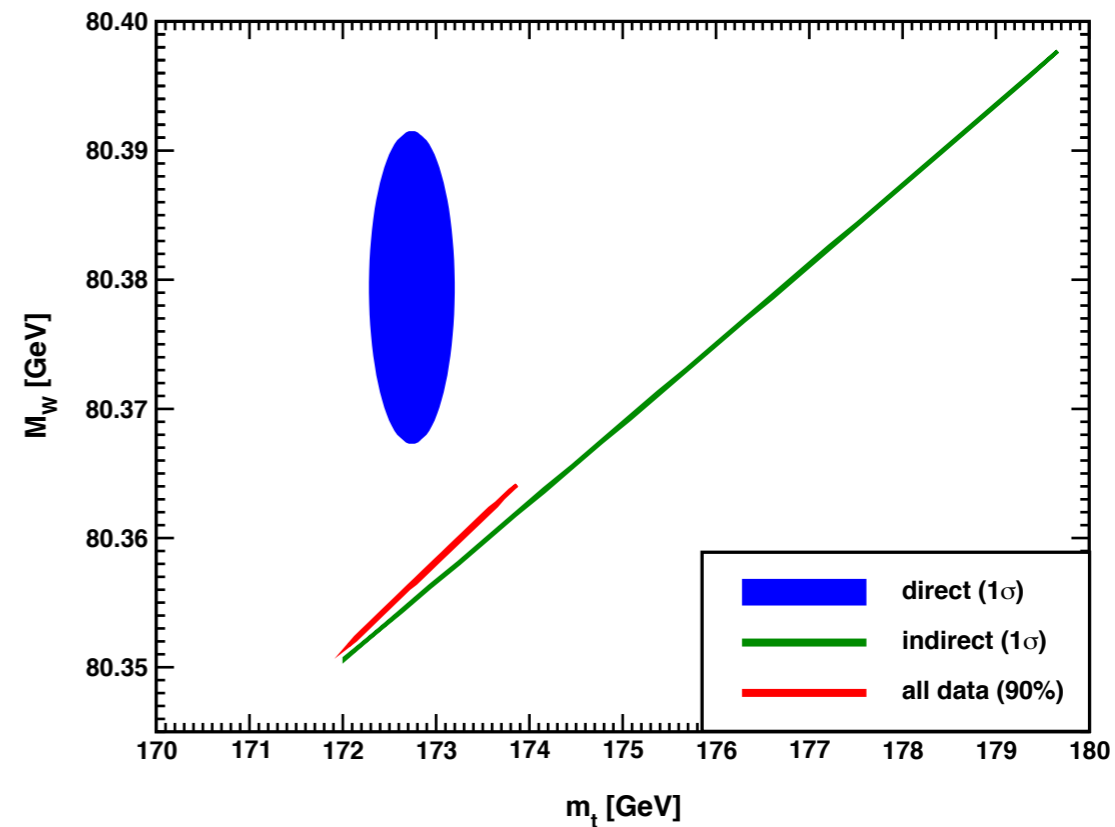


# $M_W - m_t$



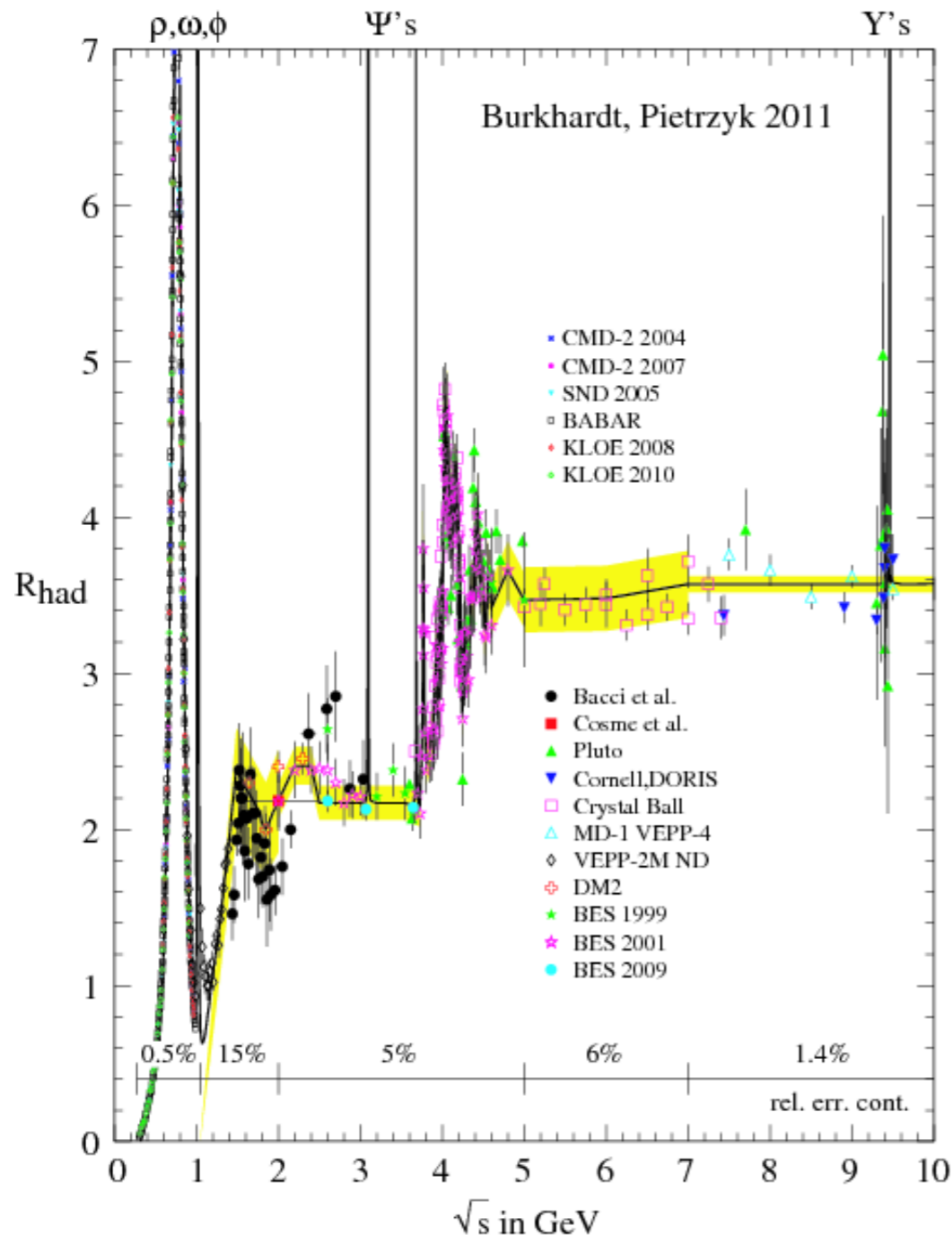
*Freitas & JE (PDG 2018)*

# $M_W$ in the MSSM



*Heinemeyer,  
Hollik,  
Weiglein,  
Zeune 2013*

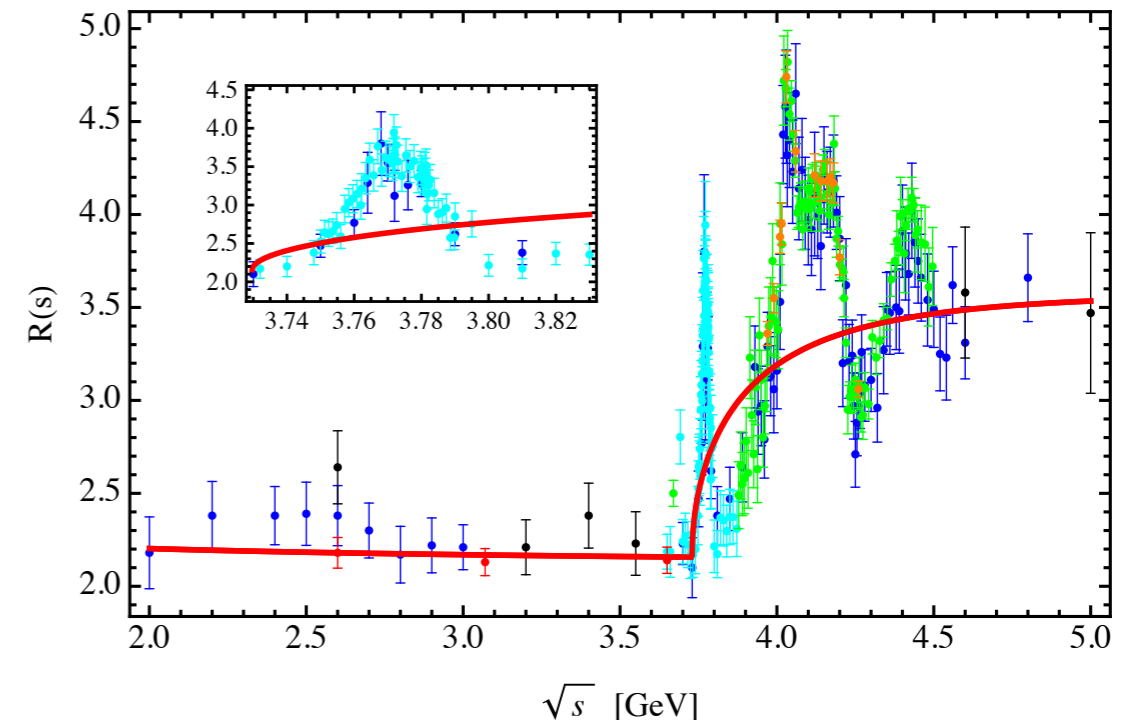
# $m_c$



- $\alpha(M_Z)$  and  $\sin^2\theta_W(0)$ : can use PQCD for heavy quark contribution if masses are known.
- $g-2$ : c quark contribution to muon  $g-2$  similar to  $\gamma^*\gamma$ ;  $\pm 70$  MeV uncertainty in  $m_c$  induces an error of  $\pm 1.6 \times 10^{-10}$  comparable to the projected errors for the FNAL and J-PARC experiments.
- Yukawa coupling – mass relation (in single Higgs doublet SM):  $\Delta m_b = \pm 9$  MeV and  $\Delta m_c = \pm 8$  MeV to match precision from HiggsBRs @ FCC-ee
- QCD sum rule:  $m_c = 1272 \pm 8$  MeV  
*Masjuan, Spiesberger & JE, EPJC 77 (2017)*  
(expect about twice the error for  $m_b$ )

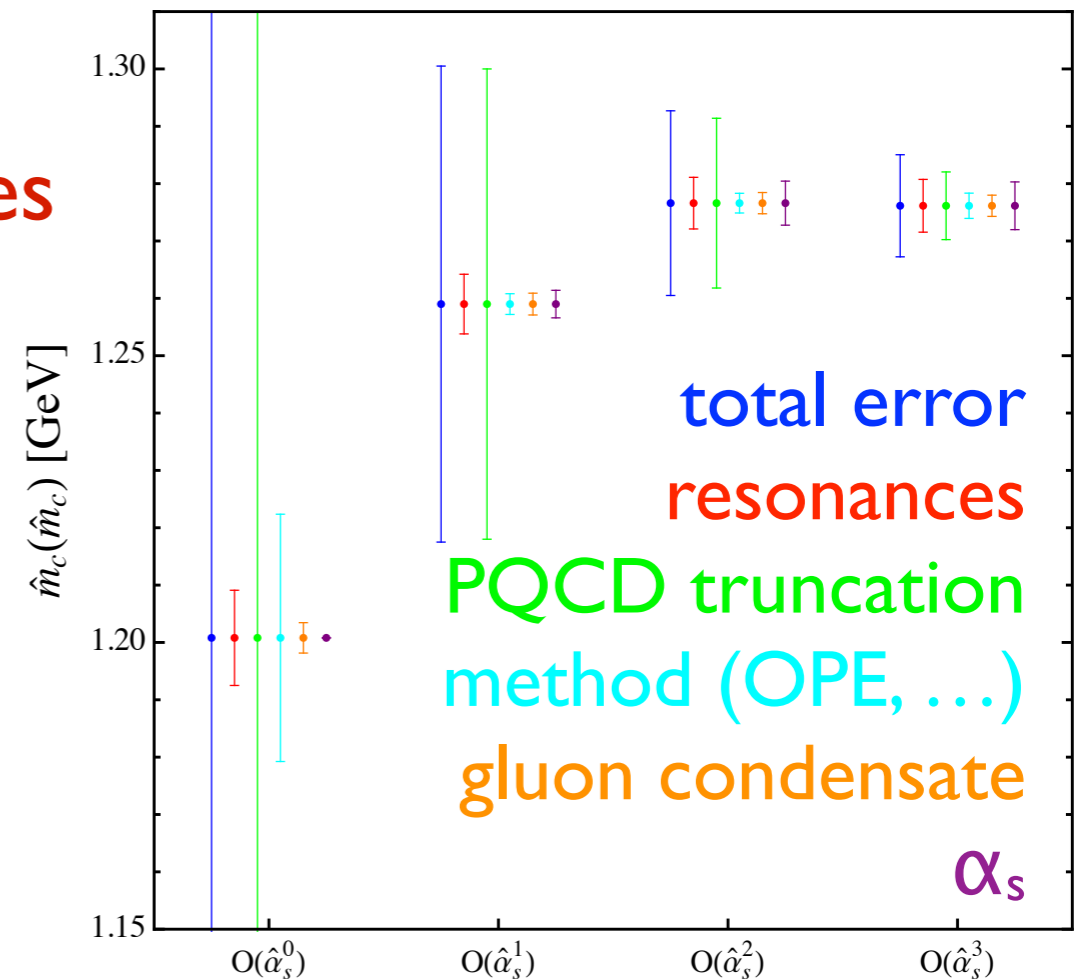
# Features of our approach

- only experimental input: **electronic widths** of  $J/\psi$  and  $\psi(2S)$
- continuum contribution from **self-consistency between sum rules**
- include  $M_0 \rightarrow$   
stronger (milder) sensitivity  
to continuum ( $m_c$ )
- quark-hadron duality needed  
only in finite region (**not locally**)
- $\bar{m}_c(\bar{m}_c) = 1272 \pm 8 + 2616 [\bar{\alpha}_s(M_Z) - 0.1182] \text{ MeV}$   
*Masjuan, Spiesberger & JE, EPJC 77 (2017)*



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# **Gauge Couplings at Lower Energies**

# $\alpha_s$ from $\tau$ decays

$$\tau_\tau = \hbar \frac{1 - \mathcal{B}_\tau^s}{\Gamma_\tau^e + \Gamma_\tau^\mu + \Gamma_\tau^{ud}} = 290.75 \pm 0.36 \text{ fs},$$

$$\Gamma_\tau^{ud} = \frac{G_F^2 m_\tau^5 |V_{ud}|^2}{64\pi^3} S(m_\tau, M_Z) \left( 1 + \frac{3}{5} \frac{m_\tau^2 - m_\mu^2}{M_W^2} \right) \times$$

$$\left[ 1 + \frac{\alpha_s(m_\tau)}{\pi} + 5.202 \frac{\alpha_s^2}{\pi^2} + 26.37 \frac{\alpha_s^3}{\pi^3} + 127.1 \frac{\alpha_s^4}{\pi^4} + \frac{\hat{\alpha}}{\pi} \left( \frac{85}{24} - \frac{\pi^2}{2} \right) + \delta_{\text{NP}} \right]$$

- $\tau_\tau$  result includes leptonic branching ratios
- $\mathcal{B}_{\tau^s} = 0.0292 \pm 0.0004$  ( $\Delta S = -1$ ) *PDG 2018*
- $S(m_\tau, M_Z) = 1.01907 \pm 0.0003$  *JE, Rev. Mex. Fis. 50 (2004)*
- $\delta_{\text{NP}} = 0.003 \pm 0.009$  (within OPE & OPE breaking) based on (controversial) *Boito et al., PRD 85 (2012) & PRD 91 (2015); Davier et al., EPJC 74 (2014); Pich & Rodríguez-Sánchez, PRD 94 (2016)*

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- dominant uncertainty from PQCD truncation (FOPT vs. CIPT vs. geometric continuation)
- $\alpha_s^{(4)}(m_\tau) = 0.323^{+0.018}_{-0.014}$
- $\alpha_s^{(5)}(M_Z) = 0.1184^{+0.0020}_{-0.0018}$
- updated from [Luo & JE, PLB 558 \(2003\)](#) in [Freitas & JE \(PDG 2018\)](#)



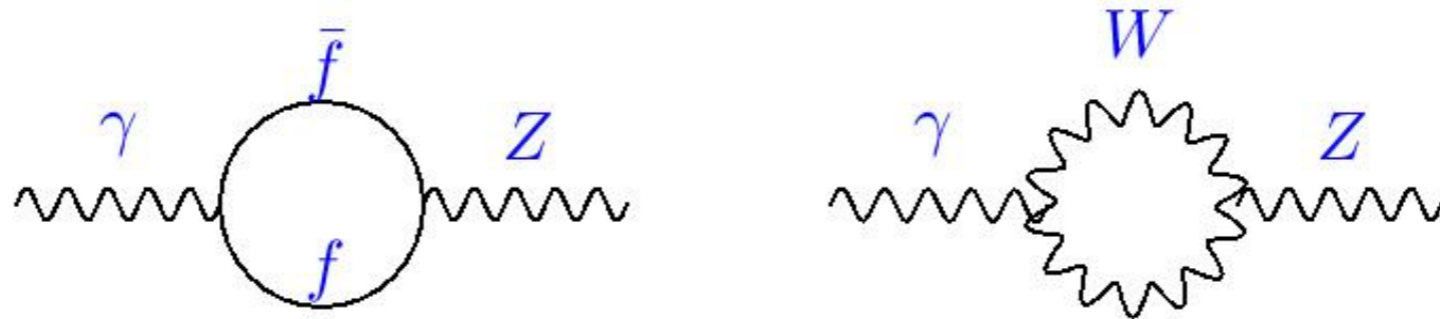
# $\alpha(M_Z)$

- Dispersive approach:
  - $\alpha^{-1}(M_Z) = 128.947 \pm 0.012$  *Davier et al., EPJC 77 (2017)*
  - $\alpha^{-1}(M_Z) = 128.958 \pm 0.016$  *Jegerlehner, arXiv:1711.06089*
  - $\alpha^{-1}(M_Z) = 128.946 \pm 0.015$  *Keshavarzi et al., arXiv:1802.02995*
- $\alpha^{-1}(M_Z) = 128.949 \pm 0.010$  *Ferro-Hernández & JE, JHEP 03 (2018)*
  - **This value** is converted from the  $\overline{MS}$  scheme and uses both  $e^+e^-$  annihilation and  $\tau$  decay spectral functions *Davier et al., EPJC 77 (2017)*
  - $\tau$  data corrected for  $\gamma$ - $\rho$  mixing *Jegerlehner & Szafron, EPJC 71 (2011)*
  - PQCD for  $\sqrt{s} > 2$  GeV (using  $\overline{m}_c$  &  $\overline{m}_b$ ) *Ferro-Hernández & JE, in preparation*

# $g_{\mu-2}$ hadronic effects

- $a_{\mu}^{\text{had,LO}} = (69.31 \pm 0.34) \times 10^{-9}$  *Davier et al., EPJC 77 (2017)*
- $a_{\mu}^{\text{had,LO}} = (68.81 \pm 0.41) \times 10^{-9}$  *Jegerlehner, EPJ Web Conf. 166 (2018)*
- $a_{\mu}^{\text{had,LO}} = (68.88 \pm 0.34) \times 10^{-9}$  (incl.  $\tau$  data) *Jegerlehner, EPJ Web Conf. 166 (2018)*
- $a_{\mu}^{\text{had,LO}} = (69.33 \pm 0.25) \times 10^{-9}$  *Keshavarzi et al., arXiv:1802.02995*
- $a_{\mu}^{\text{had,NLO}} = (-1.01 \pm 0.01) \times 10^{-9}$  (anti-correlated with  $a_{\mu}^{\text{had,LO}}$ ) *Krause, PLB 390 (1997)*
- $a_{\mu}^{\text{had,NNLO}} = (0.124 \pm 0.001) \times 10^{-9}$  *Kurz et al., EPJ Web Conf. 118 (2016)*
- $a_{\mu}^{\text{had,LBLS}} (\alpha^3) = (1.05 \pm 0.33) \times 10^{-9}$  ( $\bar{m}_c$  treatment!) *Toledo-Sánchez & JE, PRL 97 (2006)*
- $a_{\mu}^{\text{had,LBLS}} (\alpha^4) = (0.03 \pm 0.02) \times 10^{-9}$  *Colangelo et al., PLB 735 (2014)*
- $a_{\mu} (\text{exp.}) - a_{\mu} (\text{SM}) = (2.55 \pm 0.77) \times 10^{-9}$  ( $3.3 \sigma$ ) *Freitas & JE, PDG 2018*

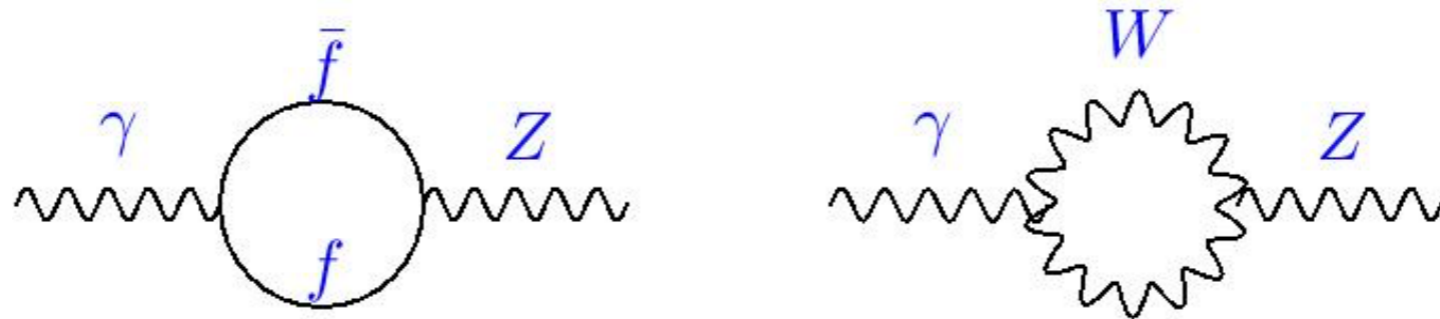
# $\sin^2\theta_W(0)$ : RGE



$$\mu^2 \frac{d\hat{v}_f}{d\mu^2} = \frac{\hat{\alpha} Q_f}{24\pi} \left[ \sum_i K_i \gamma_i \hat{v}_i Q_i + 12\sigma \left( \sum_q Q_q \right) \left( \sum_q \hat{v}_q \right) \right]$$

- $v_f$ : Z vector coupling to fermion f
- $K_i$ : QCD factor known to  $O(\alpha_s^4)$  *Baikov et al., JHEP 07 (2012)*
- $\sigma$ : singlet piece at  $O(\alpha_s^3)$  and  $O(\alpha_s^4)$  *Baikov et al., JHEP 07 (2012)*
- $\gamma_i$ : field type dependent constants *Ramsey-Musolf & JE, PRD 72 (2005)*

# $\sin^2\theta_W(0)$ and $\Delta\alpha(M_Z)$



$$\mu^2 \frac{d\hat{v}_f}{d\mu^2} = \frac{\hat{\alpha} Q_f}{24\pi} \left[ \sum_i K_i \gamma_i \hat{v}_i Q_i + 12\sigma \left( \sum_q Q_q \right) \left( \sum_q \hat{v}_q \right) \right]$$

compare with

$$\mu^2 \frac{d\hat{\alpha}}{d\mu^2} = \frac{\hat{\alpha}^2}{\pi} \left[ \frac{1}{24} \sum_i K_i \gamma_i Q_i^2 + \sigma \left( \sum_q Q_q \right)^2 \right]$$

➔ coupled system of differential equations

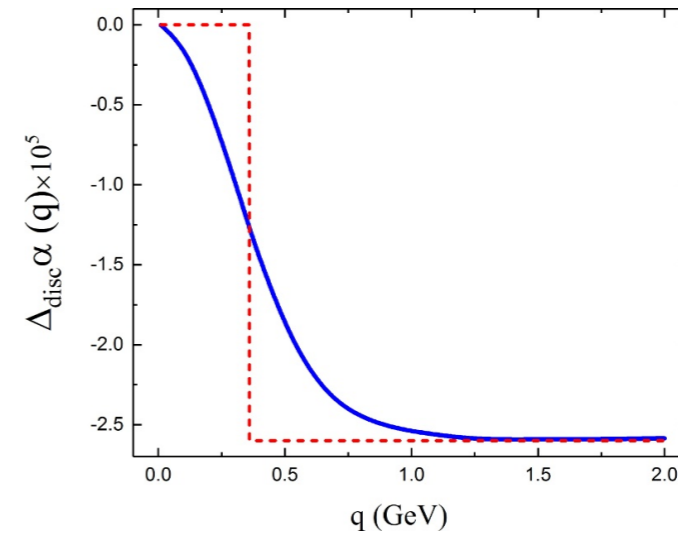
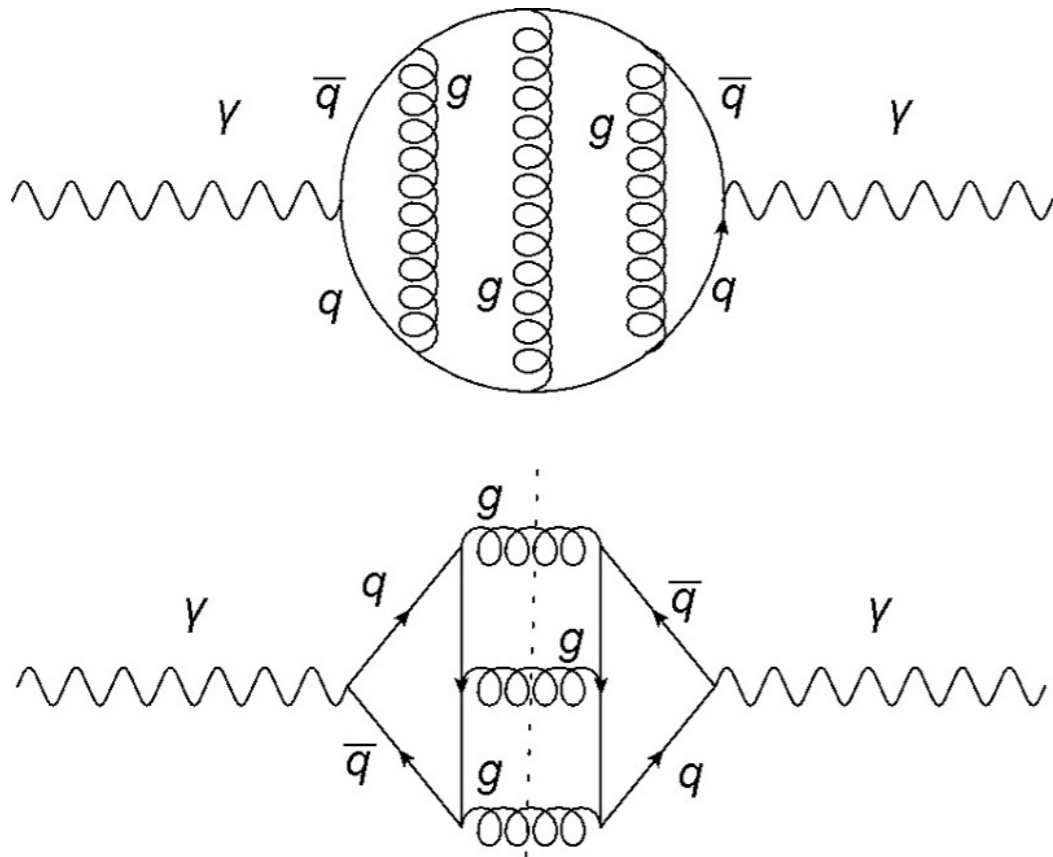
*Ramsey-Musolf & JE, PRD 72 (2005)*

# $\sin^2\theta_W(0)$ : RGE solution

$$\hat{s}^2(\mu) = \hat{s}^2(\mu_0) \frac{\hat{\alpha}(\mu)}{\hat{\alpha}(\mu_0)} + \lambda_1 \left[ 1 - \frac{\hat{\alpha}(\mu)}{\hat{\alpha}(\mu_0)} \right] + \frac{\hat{\alpha}(\mu)}{\pi} \left[ \frac{\lambda_2}{3} \ln \frac{\mu^2}{\mu_0^2} + \frac{3\lambda_3}{4} \ln \frac{\hat{\alpha}(\mu)}{\hat{\alpha}(\mu_0)} + \tilde{\sigma}(\mu_0) - \tilde{\sigma}(\mu) \right]$$

- $\lambda_i$ : rational numbers depending on active particle content of the EFT
- theory uncertainty from RGE running  $\sim 1.6 \times 10^{-6}$  (negligible)
- theory error from b and c matching  $\sim 3 \times 10^{-6}$  (again using  $\bar{m}_c$  &  $\bar{m}_b$ )
- we recycle the on-shell result for  $\alpha(2 \text{ GeV})$  *Davier et al., EPJC 77 (2017)*  
→ scheme conversion introducing  $4.8 \times 10^{-6}$  uncertainty
- **total uncertainty from PQCD  $\sim 6 \times 10^{-6}$  in  $\sin^2\theta_W(0) \equiv \overline{S}^2$**

# $\sin^2\theta_W(0)$ : singlet separation



*Ferro-Hernández & JE, JHEP 03 (2018)*  
 adapted from lattice  $g_{\mu-2}$  calculation  
*RBC/UKQCD, PRL 116 (2016)*

- use of result for  $\alpha(2 \text{ GeV})$  needs singlet piece isolation  $\Delta_{\text{disc}} \alpha(2 \text{ GeV})$
- then  $\Delta_{\text{disc}} \bar{s}^2 = (\bar{s}^2 \pm 1/20) \Delta_{\text{disc}} \alpha(2 \text{ GeV}) = (-6 \pm 3) \times 10^{-6}$
- **step function**  $\Rightarrow$  singlet threshold mass  $\bar{m}_s^{\text{disc}} \approx 350 \text{ MeV}$

# $\sin^2\theta_W(0)$ : flavor separation

strange quark external current	ambiguous external current
$\Phi$	$K\bar{K}$ (non- $\Phi$ )
$K\bar{K}\pi$ [almost saturated by $\Phi(1680)$ ]	$K\bar{K}2\pi$ , $K\bar{K}3\pi$
$\eta\Phi$	$K\bar{K}\eta$ , $K\bar{K}\omega$

- use of result for  $\alpha(2 \text{ GeV})$  also needs isolation of strange contribution  $\Delta_s\alpha$
- left column assignment assumes OZI rule
- expect right column to originate mostly from strange current ( $m_s > m_{u,d}$ )
- quantify expectation using averaged  $\Delta_s(g_\mu-2)$  from lattices as Bayesian prior  
*RBC/UKQCD, JHEP 04 (2016); HPQCD, PRD 89 (2014)*
- $\Delta_s\alpha(1.8 \text{ GeV}) = (7.09 \pm 0.32) \times 10^{-4}$  (threshold mass  $\bar{m}_s = 342 \text{ MeV} \approx \bar{m}_s^{\text{disc}}$ )

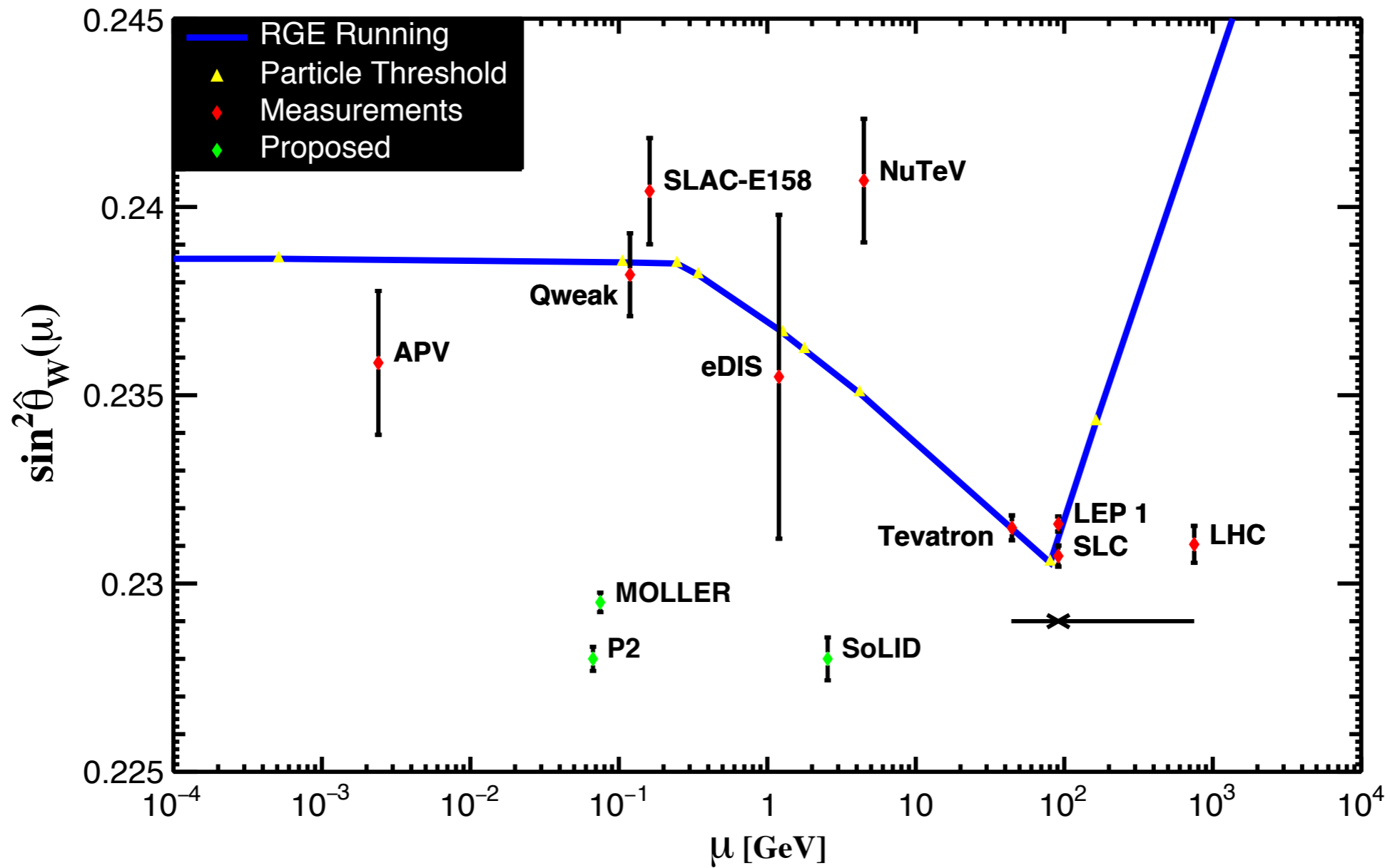
# $\sin^2\theta_W(0)$ : result

source	uncertainty in $\sin^2\theta_W(0)$
$\Delta\alpha^{(3)}(2 \text{ GeV})$	$1.2 \times 10^{-5}$
flavor separation	$1.0 \times 10^{-5}$
isospin breaking	$0.7 \times 10^{-5}$
singlet contribution	$0.3 \times 10^{-5}$
PQCD	$0.6 \times 10^{-5}$
Total	$1.8 \times 10^{-5}$

→  $\sin^2\theta_W(0) = 0.23861 \pm 0.00005_{\text{Z-pole}} \pm 0.00002_{\text{theory}} \pm 0.00001_{\alpha_s}$   
*Ferro-Hernández & JE, JHEP 03 (2018); Freitas & JE, PDG 2018*




# $\sin^2\theta_w(\mu)$



Ferro-Hernández & JE, JHEP 03 (2018)

# **Electroweak Fits**

# Inputs

- 5 inputs needed to fix the **bosonic sector** of the SM:  
 $SU(3) \times SU(2) \times U(1)$  gauge couplings and 2 Higgs parameters
- fine structure constant:  $\alpha$  e.g. from the Rydberg constant  
(leaves  $g_e^{-2}$  as derived quantity and extra SM test)
- Fermi constant:  $G_F$  from PSI (muon lifetime)
- Z mass:  $M_Z$  from LEP
- Higgs mass:  $M_H$  from the LHC
- strong coupling constant:  $\alpha_s(M_Z)$  is fit output 

# Standard global fit

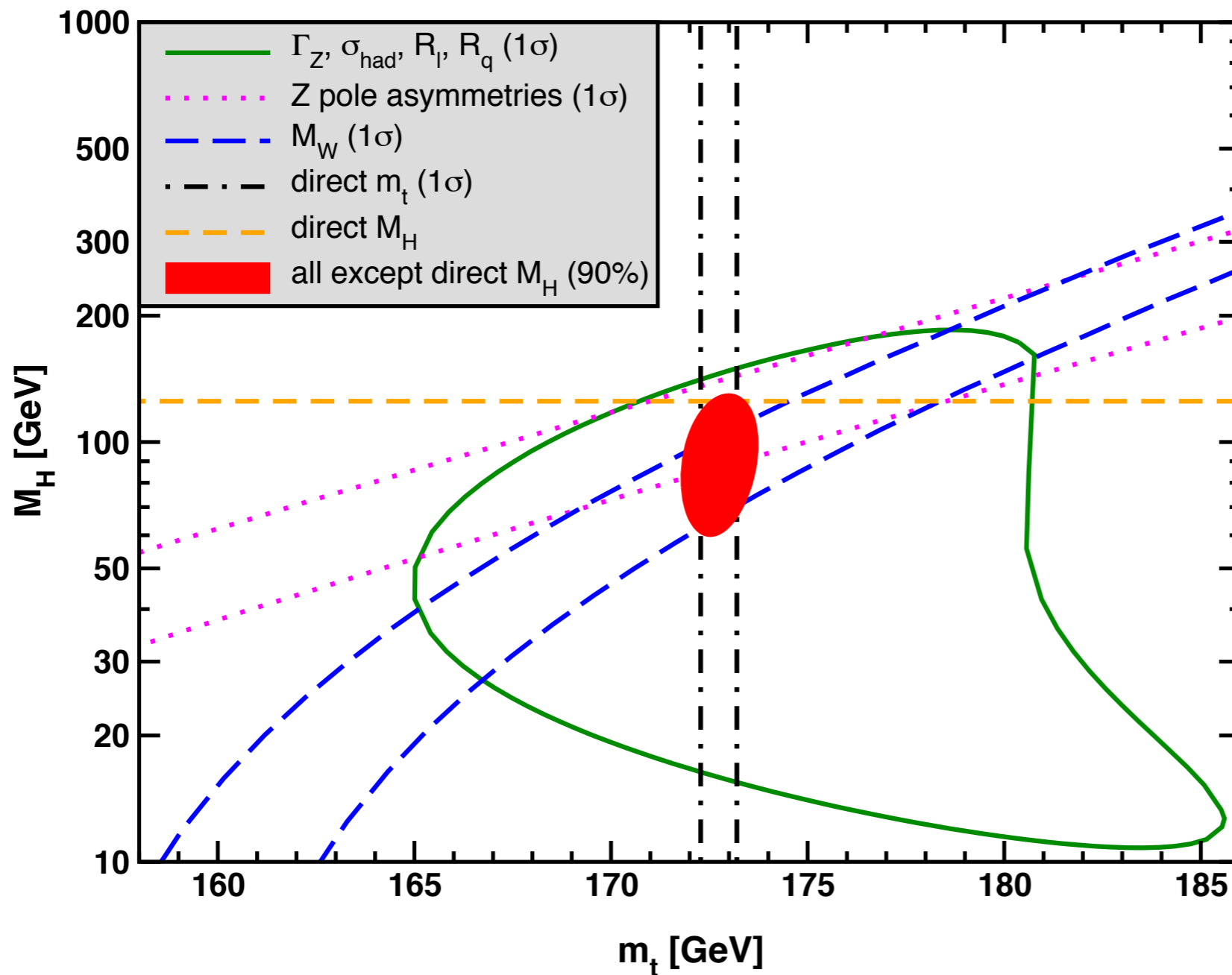
$M_H$	$125.14 \pm 0.15 \text{ GeV}$
$M_Z$	$91.1884 \pm 0.0020 \text{ GeV}$
$\bar{m}_b(\bar{m}_b)$	$4.180 \pm 0.021 \text{ GeV}$
$\Delta\alpha_{\text{had}}^{(3)}(2 \text{ GeV})$	$(59.0 \pm 0.5) \times 10^{-4}$

$\bar{m}_t(\bar{m}_t)$	$163.28 \pm 0.44 \text{ GeV}$	1.00	-0.13	-0.28
$\bar{m}_c(\bar{m}_c)$	$1.275 \pm 0.009 \text{ GeV}$	-0.13	1.00	0.45
$\alpha_s(M_Z)$	$0.1187 \pm 0.0016$	-0.28	0.45	1.00

other correlations small

*Freitas & JE, PDG 2018*

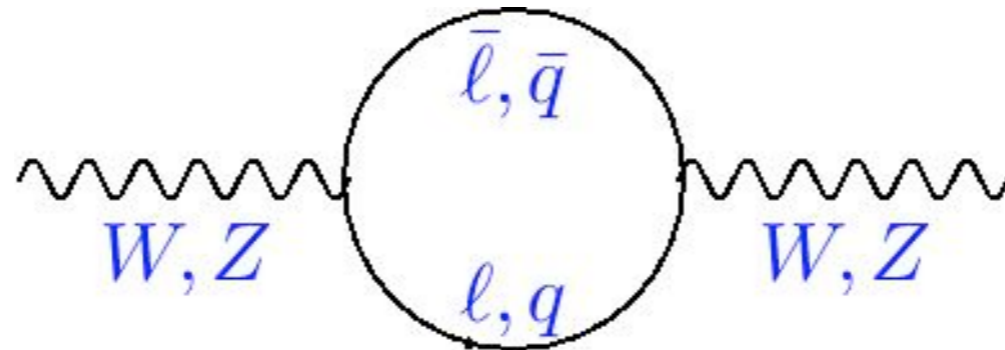
# $M_H - m_t$



indirect  $M_H$ :  
 $90^{+17}_{-16}$  GeV  
 (1.9  $\sigma$  low)

indirect  $m_t$ :  
 $176.4 \pm 1.8$  GeV  
 (2.0  $\sigma$  high)

# Oblique physics beyond the SM



- STU describe corrections to gauge-boson self-energies
- T breaks custodial  $SO(4)$
- a multiplet of heavy **degenerate** chiral fermions contributes  $\Delta S = N_C/3\pi \sum_i [t_{3L}^i - t_{3R}^i]^2$
- extra **degenerate** fermion family yields  $\Delta S = 2/3\pi \approx 0.21$
- S and T (U) correspond to dimension 6 (8) operators

# $\rho_0$ fit

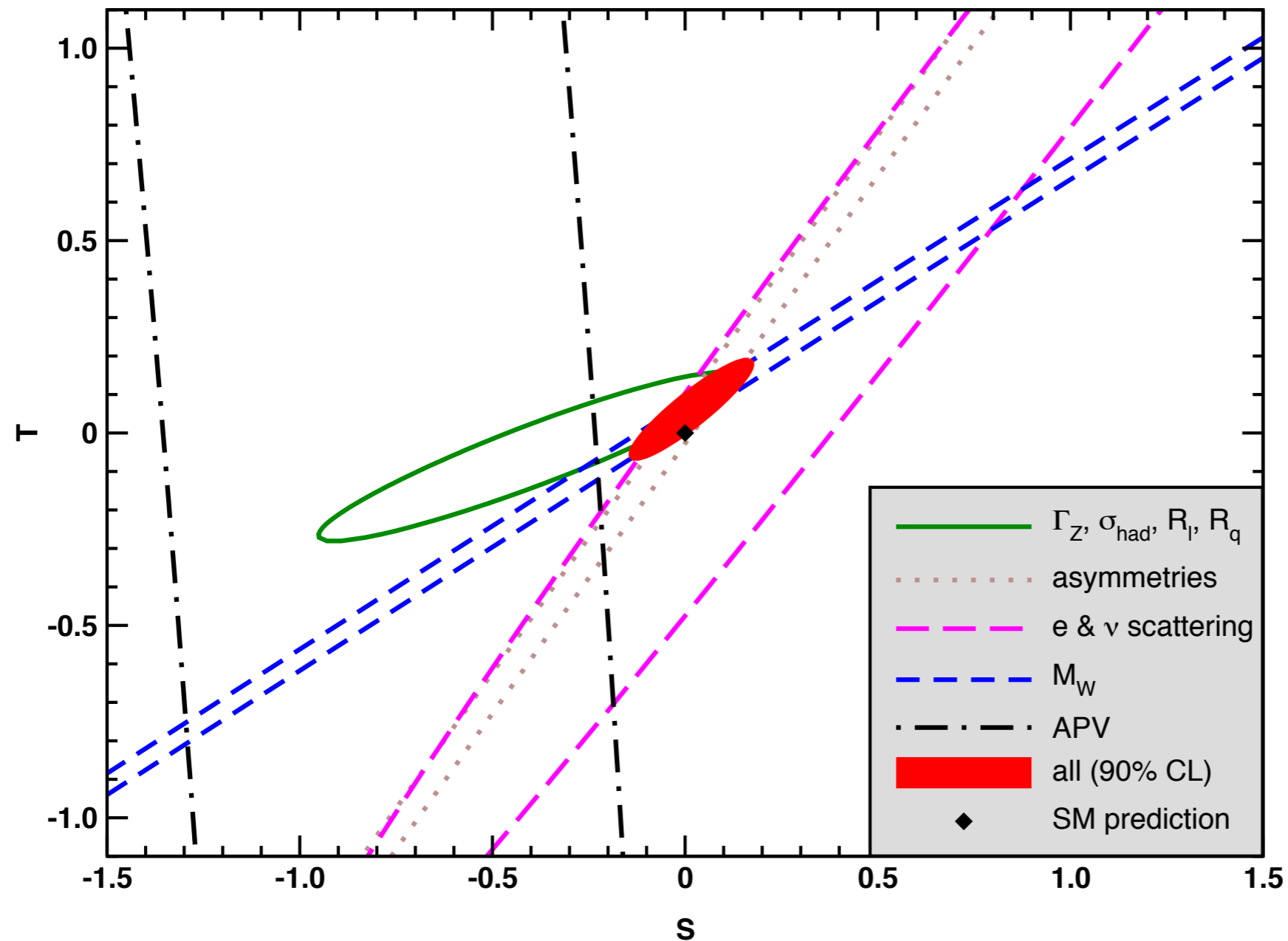
- $\Delta\rho_0 = G_F \sum_i C_i / (8\sqrt{2}\pi^2) \Delta m_i^2$ 
  - where  $\Delta m_i^2 \geq (m_1 - m_2)^2$
  - despite appearance there is decoupling (see-saw type suppression of  $\Delta m_i^2$ )
- $\rho_0 = 1.00039 \pm 0.00019$  (2.0  $\sigma$ )
  - $(16 \text{ GeV})^2 \leq \sum_i C_i / 3 \Delta m_i^2 \leq (48 \text{ GeV})^2$  @ 90% CL
  - $Y = 0$  Higgs triplet VEVs  $v_3$  strongly disfavored ( $\rho_0 < 1$ )
  - consistent with  $|Y| = 1$  Higgs triplets if  $v_3 \sim 0.01 v_2$

# S fit

- S parameter rules out QCD-like technicolor models
- S also constrains extra degenerate fermion families:
  - ➔  $N_F = 2.75 \pm 0.14$  (assuming  $T = U = 0$ )
- compare with  $N_\nu = 2.991 \pm 0.007$  from  $\Gamma_Z$



# T - S



S	$0.02 \pm 0.07$
T	$0.06 \pm 0.06$
$\Delta\chi^2$	- 4.0

*Freitas & JE, PDG 2018*

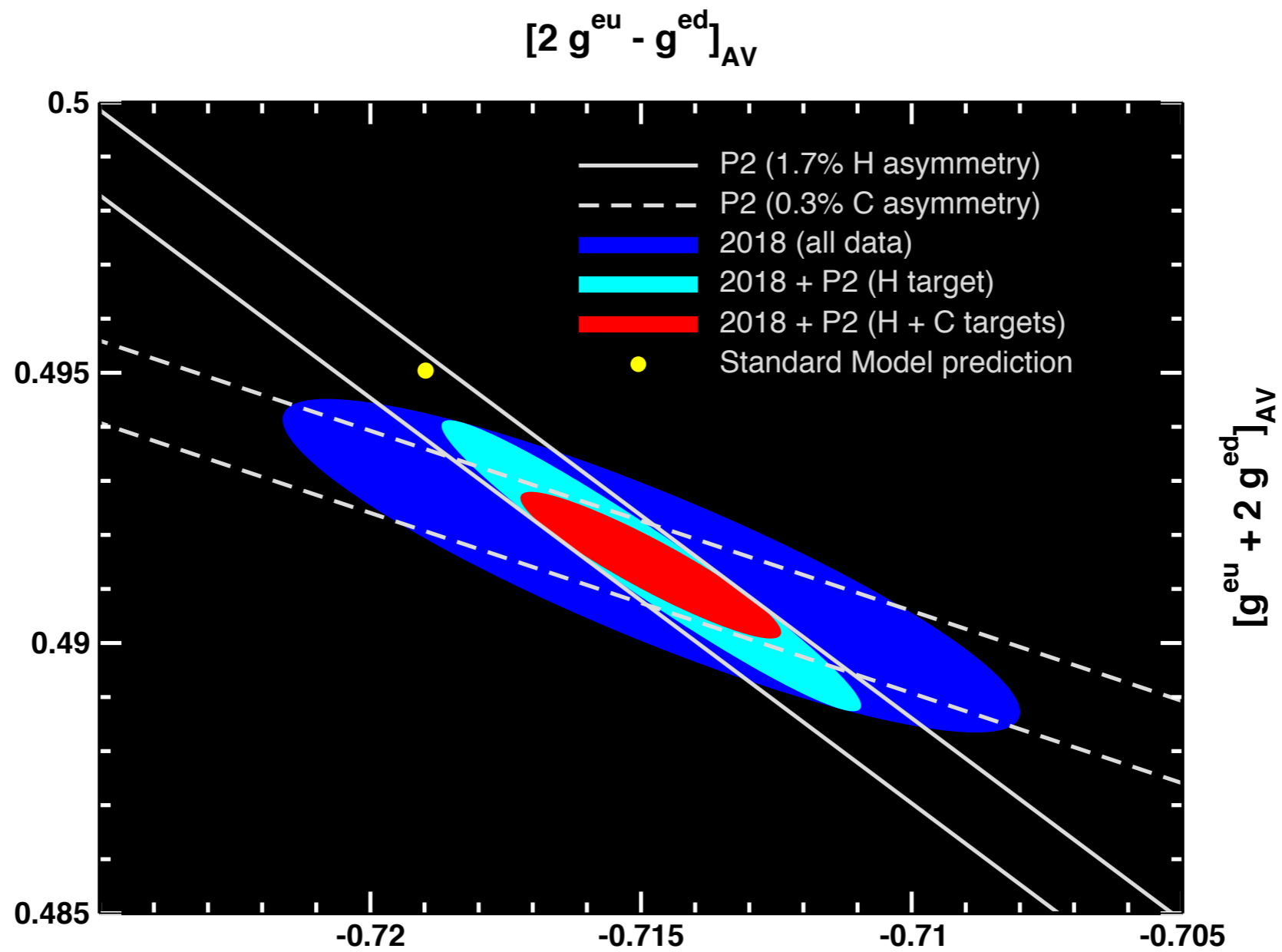
# STU fit

$\sin^2\theta_W(M_Z)$	$0.23113 \pm 0.00014$
$\alpha_s(M_Z)$	$0.1189 \pm 0.0016$

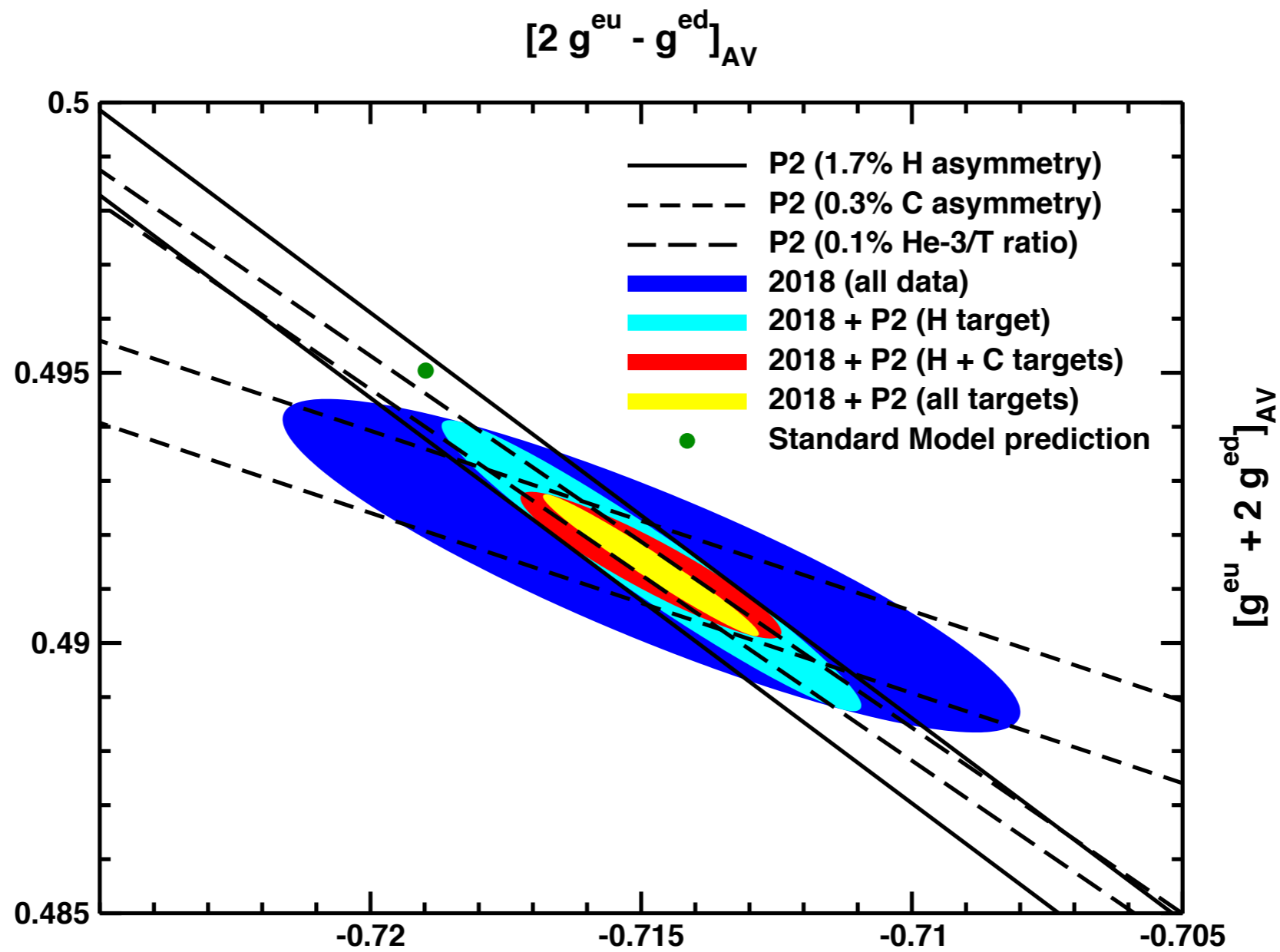
S	$0.02 \pm 0.10$	1.00	0.92	-0.66
T	$0.07 \pm 0.12$	0.92	1.00	-0.86
U	$0.00 \pm 0.09$	-0.66	-0.86	1.00

- $M_{KK} \approx 3.2 \text{ TeV}$  in warped extra dimension models
- $M_V \approx 4 \text{ TeV}$  in minimal composite Higgs models *Freitas & JE (PDG 2018)*

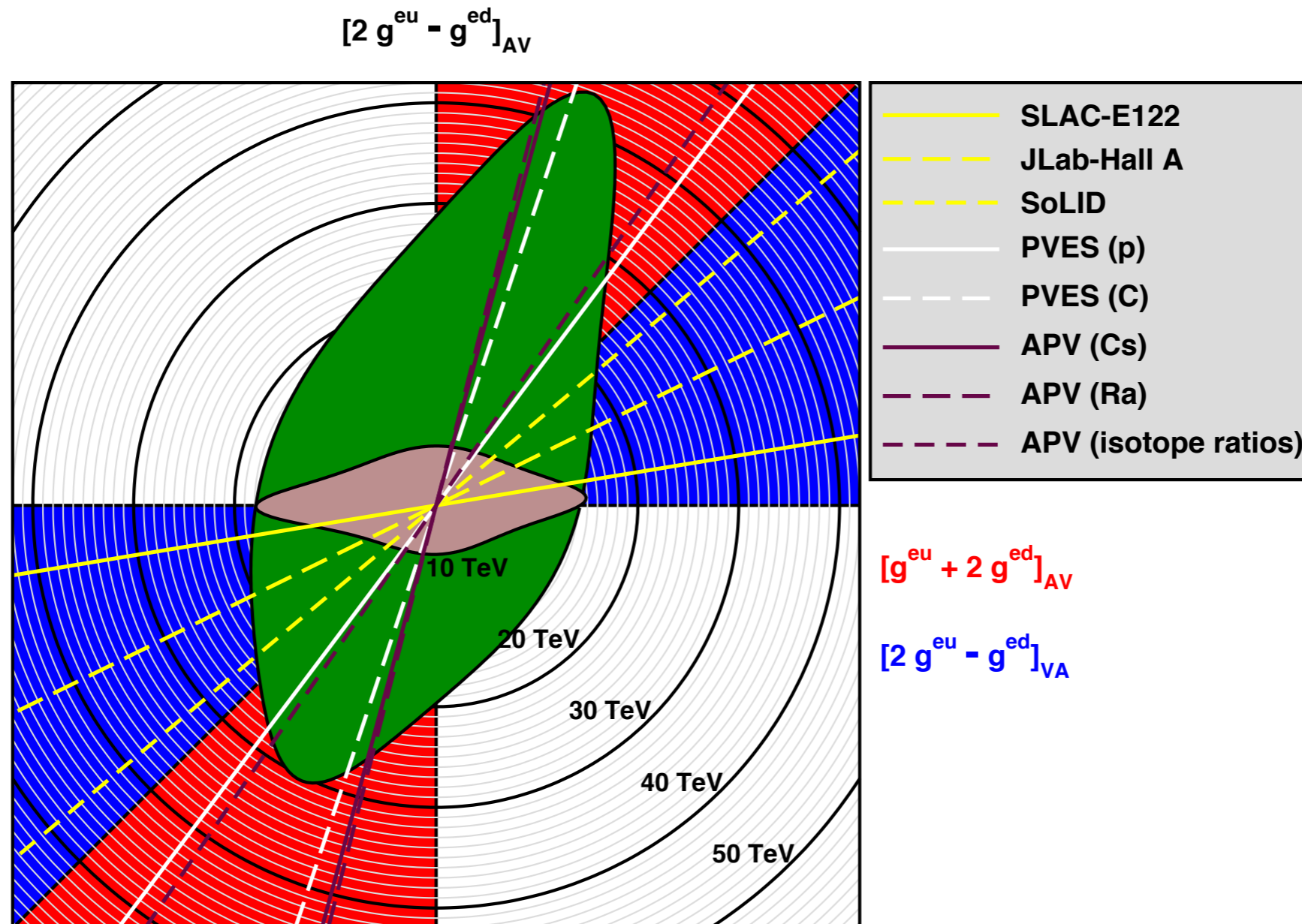
# Effective couplings



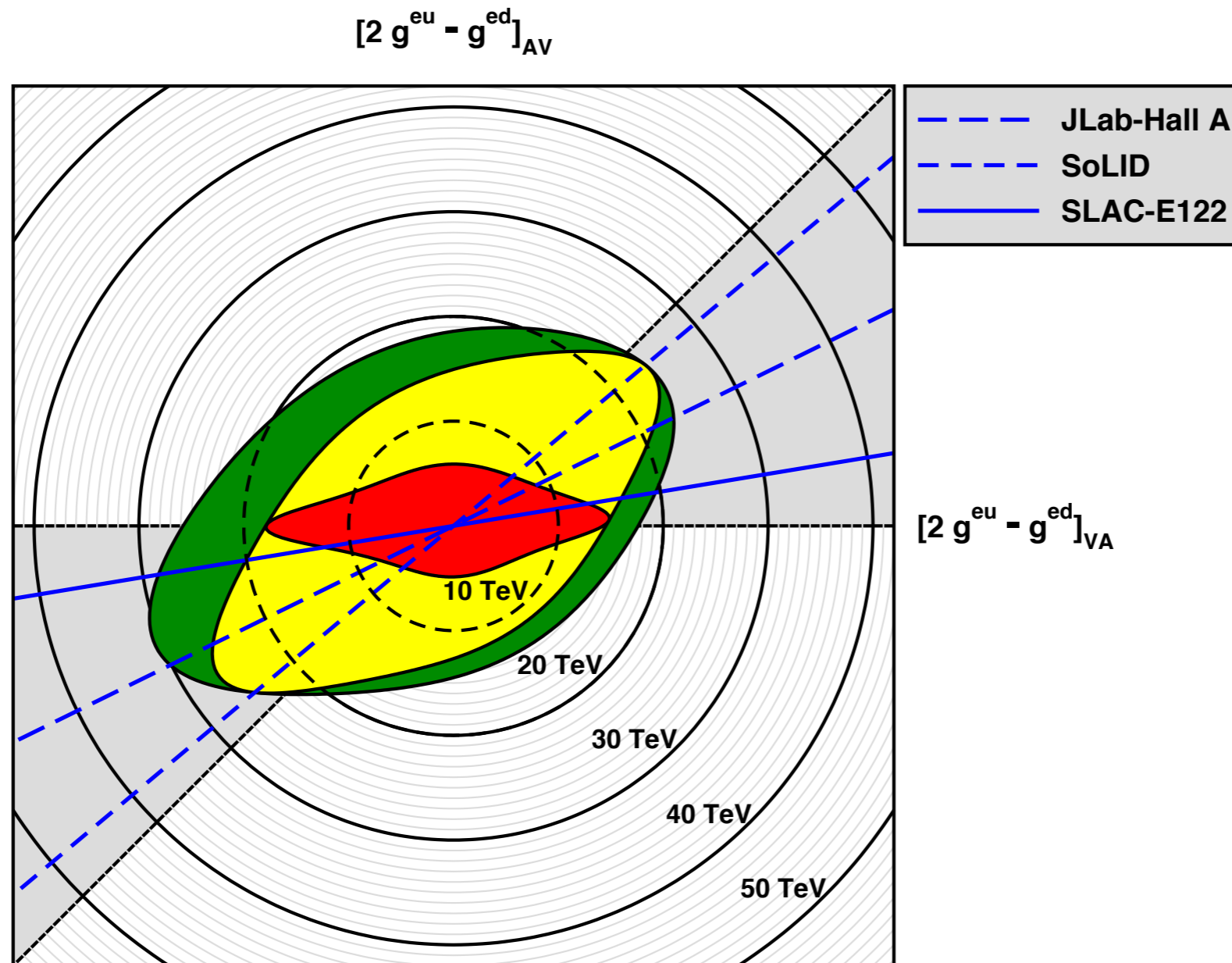
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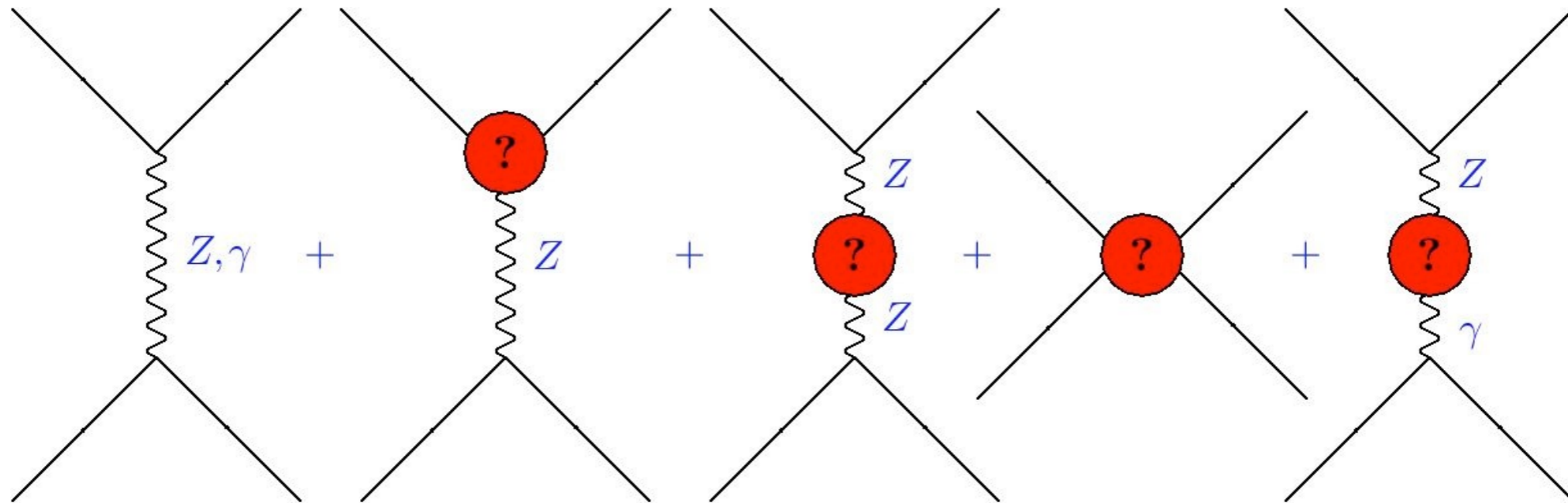
# Scale exclusions post Qweak



# Scale exclusions pre-SoLID / P2



# $\sin^2\theta_W$ beyond the SM

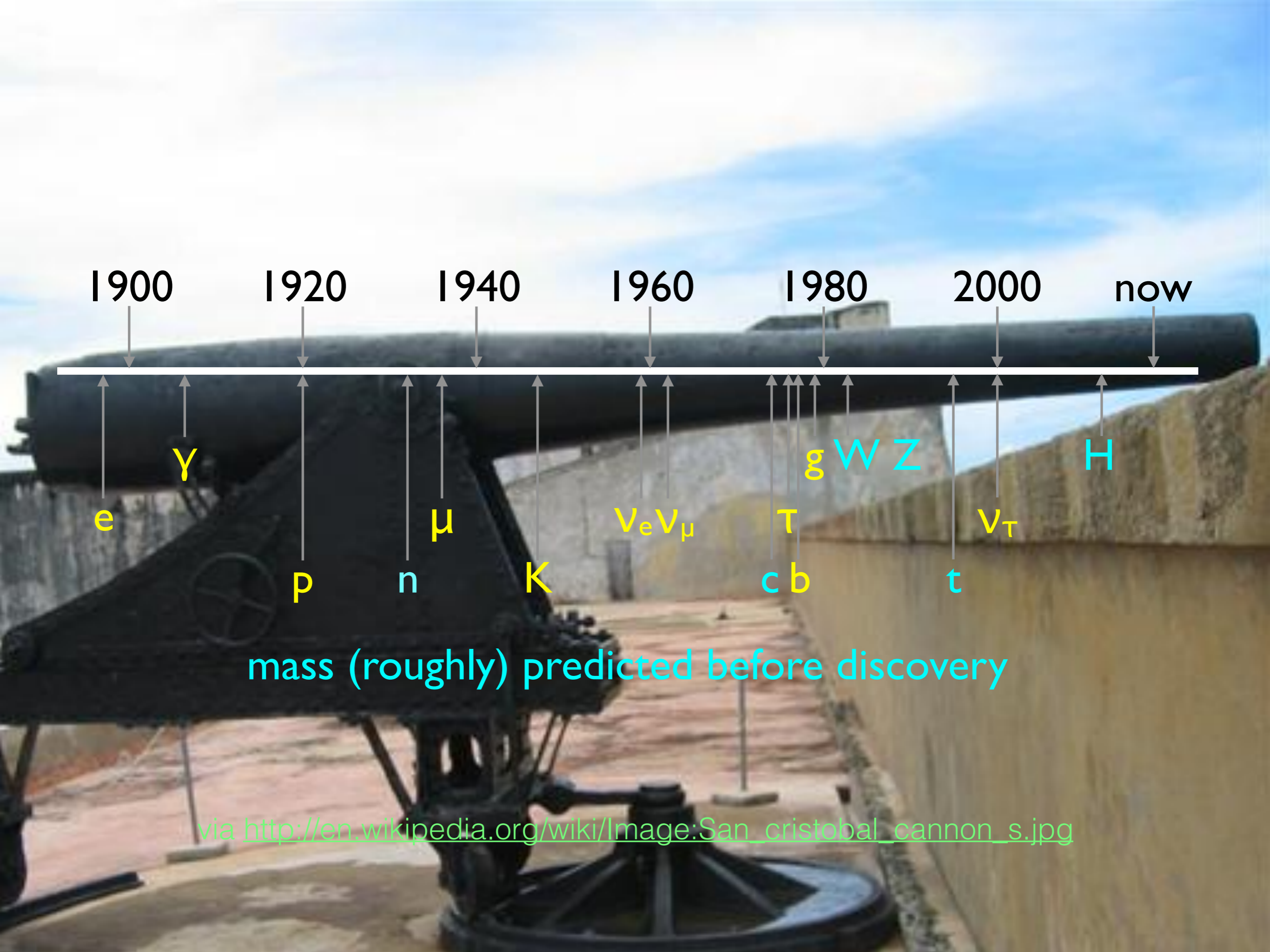


- **Z-Z' mixing:** modification of Z vector coupling
- **oblique parameters:** STU (also need  $M_W$  and  $\Gamma_Z$ )
- **new amplitudes:** off- versus on-Z pole measurements (e.g.  $Z'$ )
- **dark Z:** renormalization group evolution (running)

# Conclusions

- The SM is 50 years old and in great health — immortal?
- BSM desperados:
  - $g_{\mu-2}$  (Tuesday)
  - $M_W$  (surely only  $2\sigma$  ... but in a very special observable)
    - simplest possibility:  $\rho_0 > 1$
  - B-sector anomalies (mañana en la mañana 🤪)
- Precision in  $\sin^2\theta_W$  ( $A_{FB}$ ) &  $M_W$  and future  $Q_W(e)$  &  $Q_W(p)$  measurements challenge theory → needs major global effort





1900

1920

1940

1960

1980

2000

now

$e$

$\gamma$

$p$

$n$

$\mu$

$K$

$V_e V_\mu$

$c$

$b$

$\tau$

$g$

$W$   $Z$

$t$

$V_\tau$

$H$

mass (roughly) predicted before discovery

via [http://en.wikipedia.org/wiki/Image:San\\_cristobal\\_cannon\\_s.jpg](http://en.wikipedia.org/wiki/Image:San_cristobal_cannon_s.jpg)