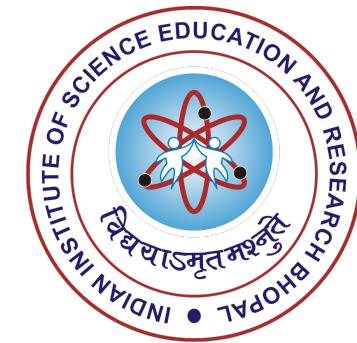
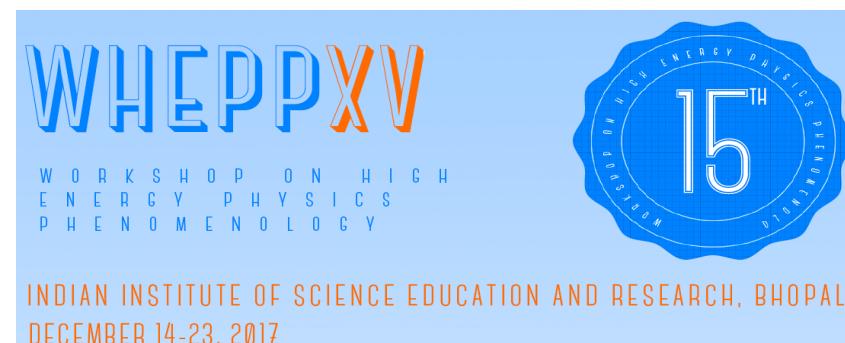


# A Fermionic Dark Matter Model

Ref: arXiv:1704.03417

Nirakar Sahoo  
IOP, Bhubaneswar

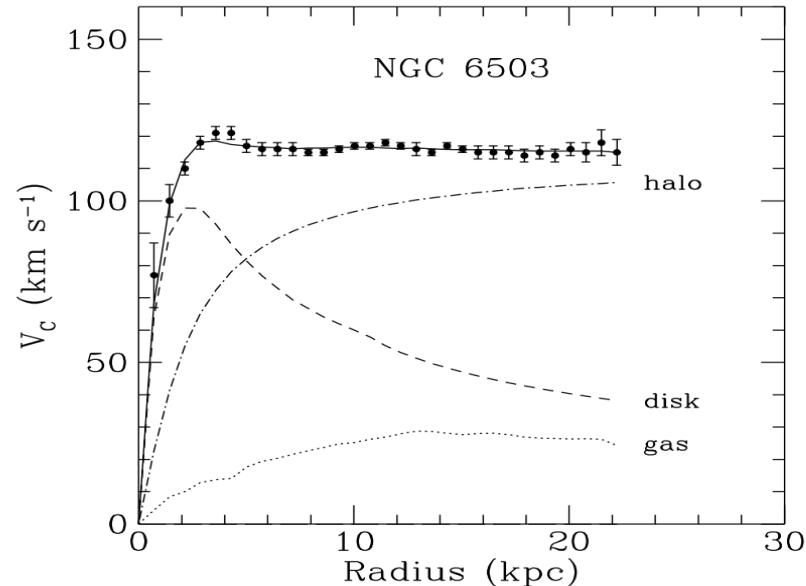


## ✧ ROTATIONAL VELOCITY OF SPIRAL GALAXY

✓ Virial theorem implies:

$$\langle K.E \rangle = -\frac{1}{2} \langle P.E \rangle$$

$$\langle v^2 \rangle = -\frac{GM}{R}$$



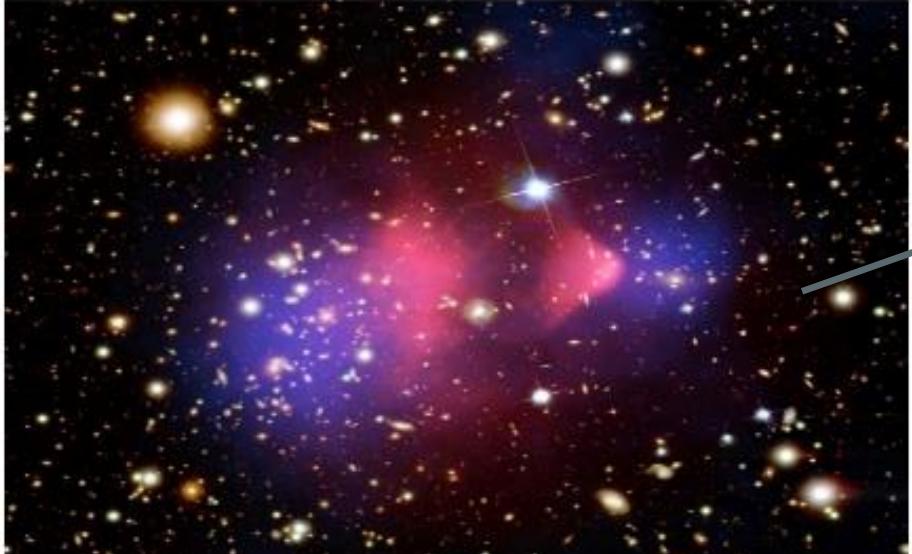
Mass inferred >> Baryonic mass

G. Jungman, M. Kamionkowski, and K. Griest  
1996

MISSING MASS IS NON-BARYONIC



# Other Astrophysical Evidences



Collision of galaxies  
in  
Bullet Cluster

[www.nasa.gov](http://www.nasa.gov)

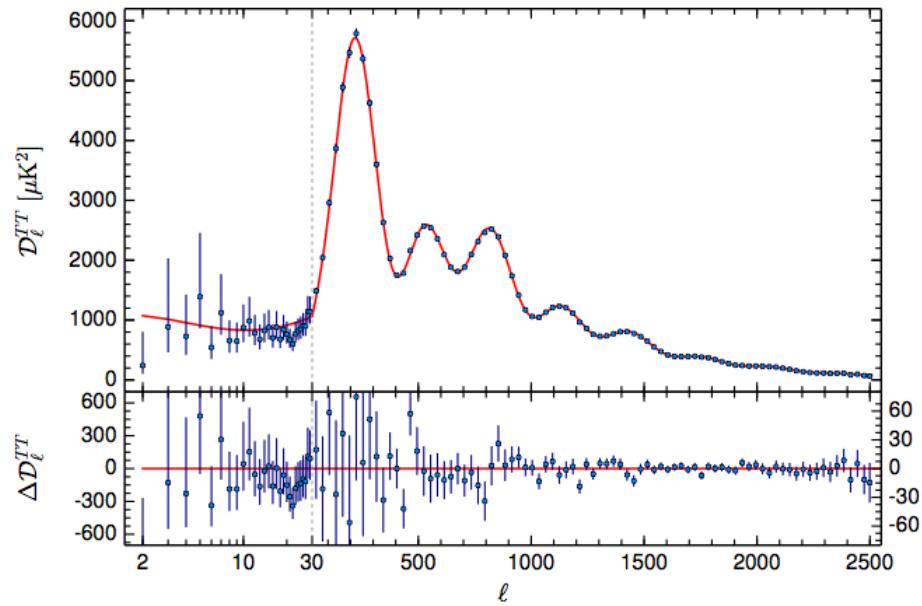
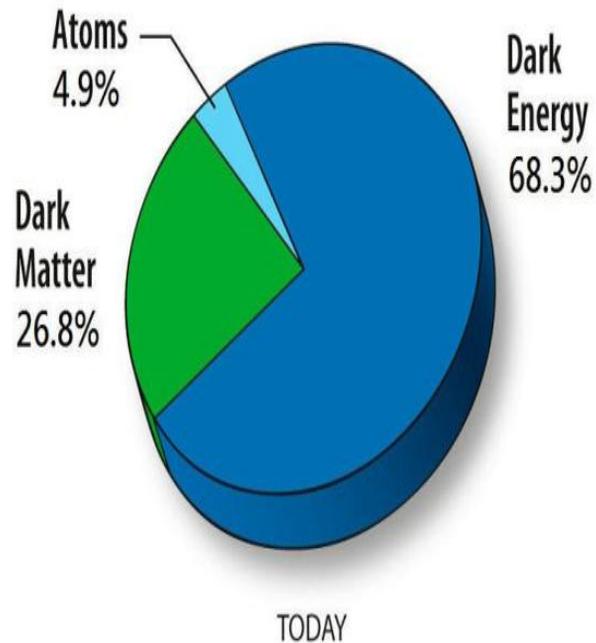
Gravitational lensing  
effect due to  
presence of DM



Gravitational lensing observed  
by Hubble in Abel 1689 cluster



# Evidence from Cosmic Microwave Data



PLANCK 2015

$$\Omega_{CDM} h^2 = 0.1199 \pm 0.0027$$



# Vector Like Mixed Singlet-doublet Fermionic DM In Presence Of A Scalar Triplet

Particle contents of the model

$$N = \begin{pmatrix} N^0 \\ N^- \end{pmatrix} = (1, 2, -1), \chi^0 = (1, 1, 0) \quad \Delta = \begin{pmatrix} \frac{\Delta^+}{\sqrt{2}} & \Delta^{++} \\ \Delta^0 & -\frac{\Delta^+}{\sqrt{2}} \end{pmatrix} (1, 3, 2)$$

The new Lagrangian apart from the SM is

$$-L_{NEW} = \bar{N} \not{D} N + \bar{\chi}^0 \not{\partial} \chi^0 + (D^\mu \Delta)^\dagger (D_\mu \Delta) + M_N \bar{N} N + M_\chi \bar{\chi}^0 \chi^0 + L_{yuk} - V(\Delta, H)$$

**where**  $L_{yuk} = \frac{1}{\sqrt{2}} \left[ (f_L)_{\alpha\beta} \bar{L}_\alpha^c i\tau_2 \Delta L_\beta + f_N \bar{N}_\alpha^c i\tau_2 \Delta N_\beta + h.c \right] + [Y \bar{N} \tilde{H} \chi^0 + h.c]$

$$V(\Delta, H) = -M_H^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + M_\Delta^2 \Delta^\dagger \Delta + \lambda_\Delta (\Delta^\dagger \Delta)^2$$

$$+ \lambda_{\Delta H} (H^\dagger H)(\Delta^\dagger \Delta) + \frac{1}{\sqrt{2}} \left[ \mu \Delta^\dagger H H + h.c \right]$$

Hep-th/0501082, hep-ph/0510064,  
 arXiv: 0705.4493, arXiv:0706.0918,  
 arXiv:0804.4080, arXiv:1109.2604,  
 arXiv:1311.5896, arXiv:1504.07892,  
 arXiv:1505.03867



# Neutrino Mass

$$(M_\nu)_{\alpha\beta} = \sqrt{2}(f_L)_{\alpha\beta} \langle u_\Delta \rangle \approx (f_L)_{\alpha\beta} \frac{-\mu v^2}{\sqrt{2}M_\Delta^2}$$

$$\mu \approx M_\Delta = O(10^{14}) \text{ GeV}$$

$$f_L \approx 1$$

$$M_\nu \approx 0.1 \text{ eV}$$

However mass of  $M_\Delta$  can be in TeV scale  
By choosing appropriate coupling



# Doublet Fermion as Dark Matter

$$L_{DM} = \bar{N} i\gamma^\mu D_\mu N + M_N \bar{N}N \text{ where } N = \begin{pmatrix} N^- \\ N^0 \end{pmatrix} \equiv (1, 2, -1)$$

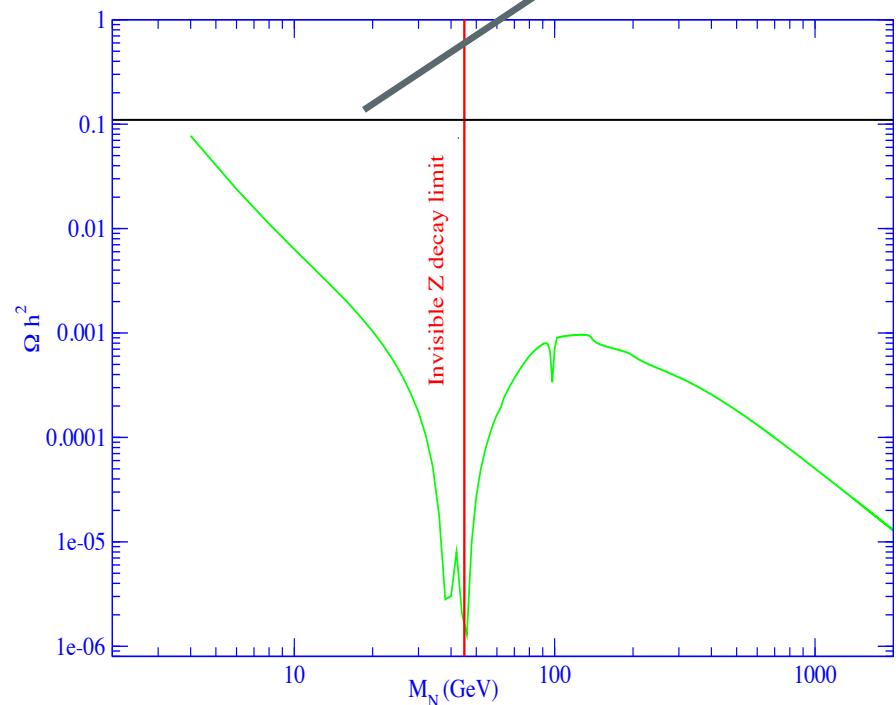
PLANCK

- ♣ The neutral field  $N^0$  can be stable by a  $Z_2$  symmetry.

- ♣ Large annihilation cross-section

Of process  $\bar{N}^0 N^0 \rightarrow W^+ W^-$

kills the relic density.



Hence  $N^0$  can not alone be a DM candidate



# Mixed Singlet-Doublet Dark Matter

After Electro-Weak Symmetry Breaking, the mass for vector like neutral fermions

$$\begin{pmatrix} \overline{N^o}, \overline{\chi^0} \end{pmatrix} \begin{pmatrix} M_N & m_D \\ m_D & M_\chi \end{pmatrix} \begin{pmatrix} N^o \\ \chi^0 \end{pmatrix} \quad m_D = Yv$$

☞ The new mass eigenstates are  $N_1$ ,  $N_2$  and  $N^\pm$  :

$$N_1 = \cos\theta\chi^0 + \sin\theta N^0$$

$$N_2 = \cos\theta N^0 - \sin\theta\chi^0$$

$$N^\pm$$

☞ And mixing angle

$$\tan 2\theta = \frac{m_D}{M_N - M_\chi}$$

$$M_1 = M_\chi - \frac{m_D^2}{M_N - M_\chi}$$

$$M_2 = M_N + \frac{m_D^2}{M_N - M_\chi}$$

$$M^\pm = M_1 \sin^2\theta + M_2 \cos^2\theta = M_N$$

$$Y = \frac{\Delta M \sin 2\theta}{2v}$$

The lightest particle  $N_1$  is the Dark Matter



# Pseudo Dirac Nature of Dark Matter

After EW symmetry breaking ,when  $\Delta$  gets an induced vev, a Majorana mass term is created for  $N_1$ .

$$m_1 = \sqrt{2} f_N \sin^2 \theta \langle \Delta \rangle \approx f_N \sin^2 \theta \frac{-\mu v^2}{\sqrt{2} M_\Delta^2}$$

Majorana mass splits Dirac state into two Majorana states ( $\Psi_1$ )<sup>a,b</sup>

Mass Eigenvalues  $M_1 \pm m_1$

$$\text{Mass splitting } \delta m = 2m_1 = 2\sqrt{2} f_N \sin^2 \theta \langle \Delta \rangle$$

Small mass-splitting does not play any role  
In relic abundance of DM, but plays important  
role in direct detection of DM



# Relic Abundance of Dark Matter

$$\Omega_{DM} h^2 = \frac{1.09 \times 10^9}{g_*^{1/2} (M_{Pl} / GeV)} \frac{1}{J(x_f)}$$

$$J(x_f) = \int_{x_f}^{\infty} \frac{\langle \sigma | v | \rangle_{eff}}{x^2} dx$$

$$\langle \sigma | v | \rangle_{eff} = \sum_{i,j}^3 (\sigma v)_{ij} \frac{g_i g_j}{g_{eff}^2} (1 + \omega_i)^{3/2} (1 + \omega_j)^{3/2} \exp[-x(\omega_i + \omega_j)]$$

with

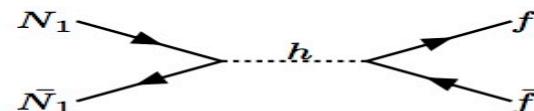
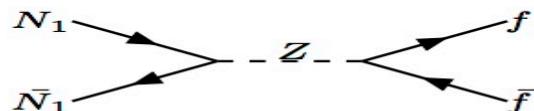
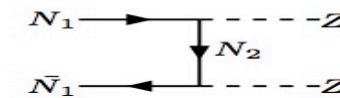
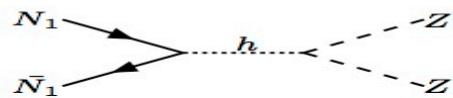
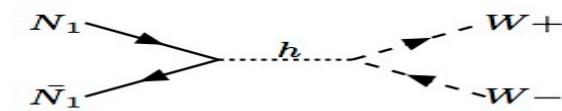
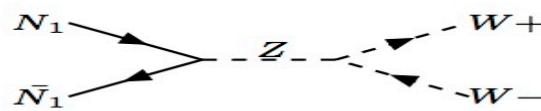
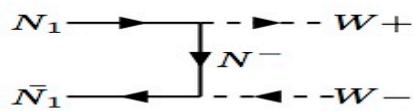
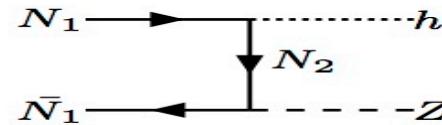
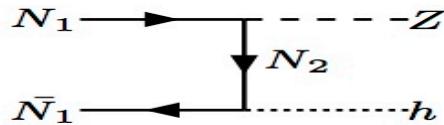
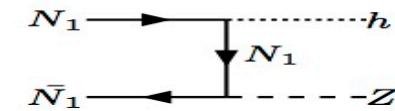
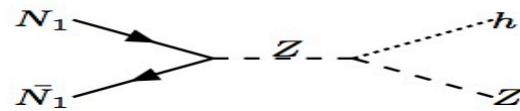
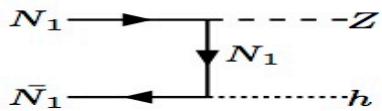
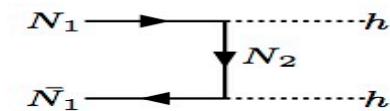
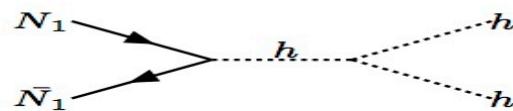
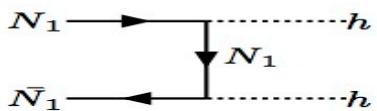
$$g_{eff} = \sum_{i=1}^3 g_i (1 + \omega_i)^{3/2} \exp(-x\omega_i)$$

$$\omega_i = \frac{M_i - M_1}{M_1}$$

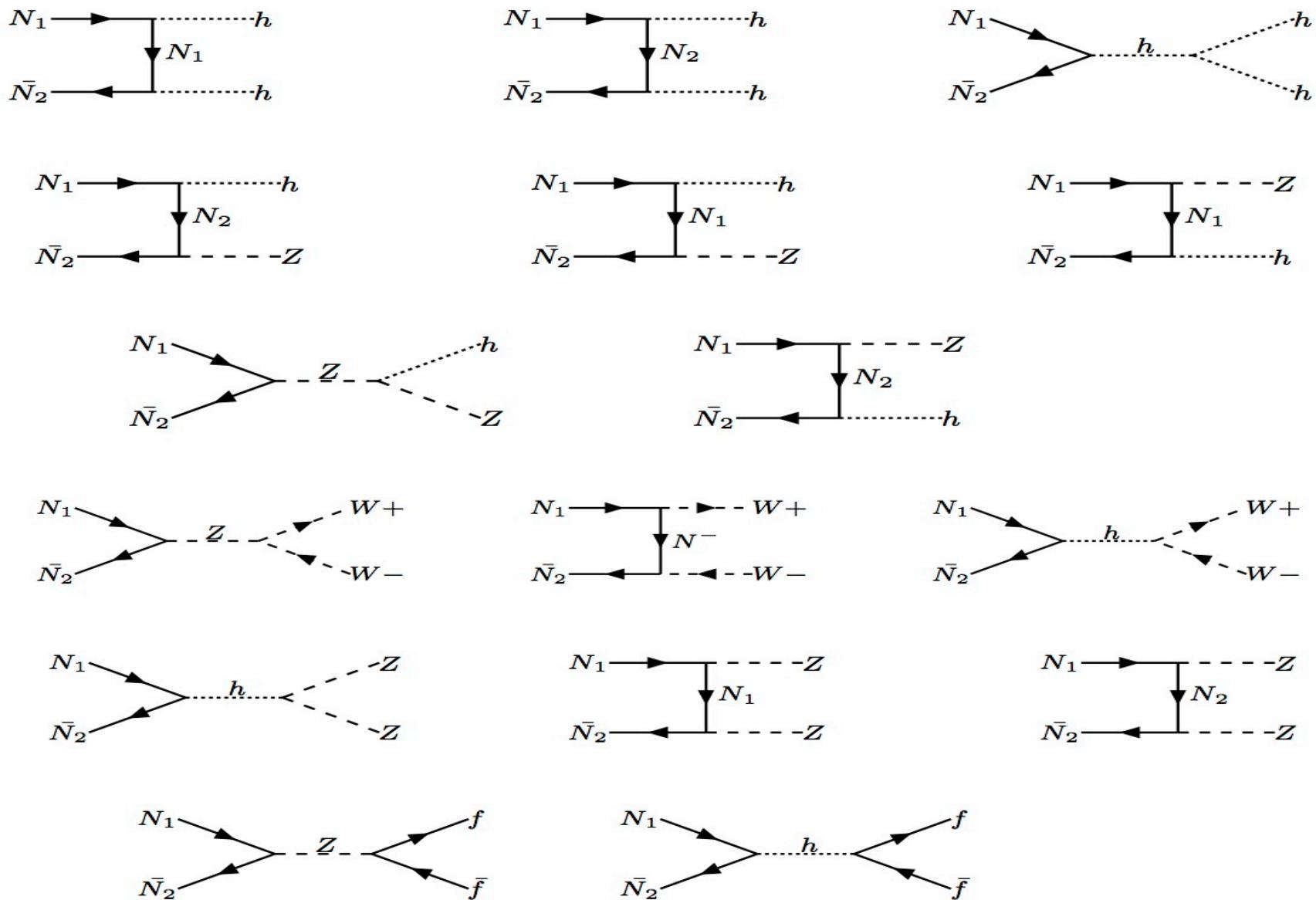
- [Griest and Seckel: PRD 1991](#)

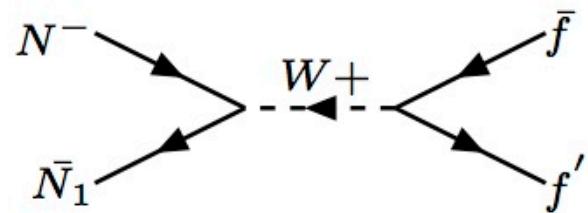
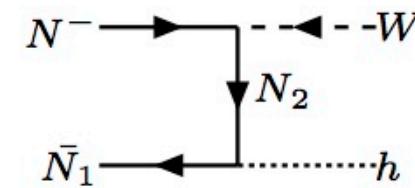
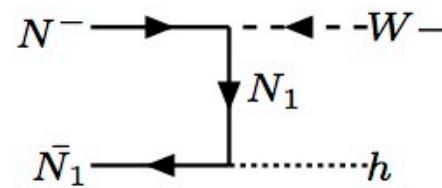
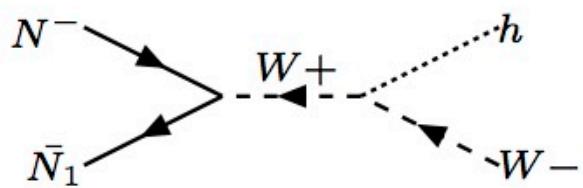
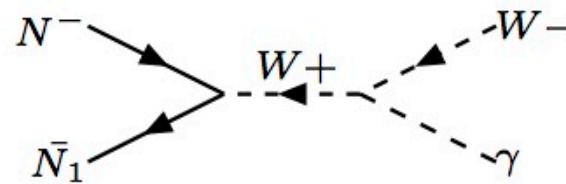
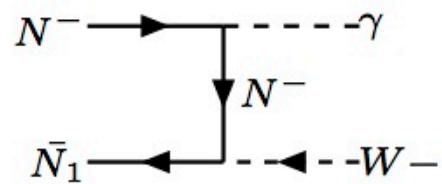
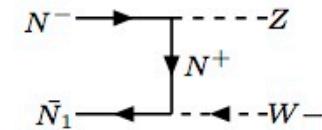
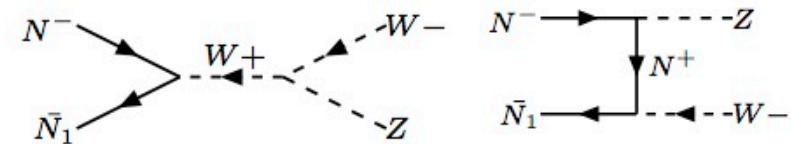
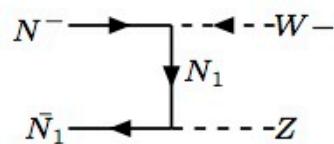
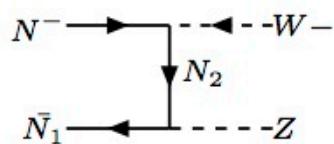


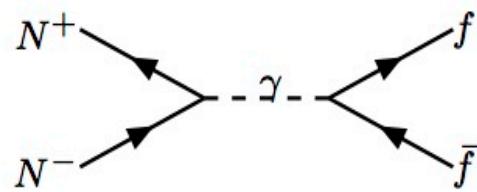
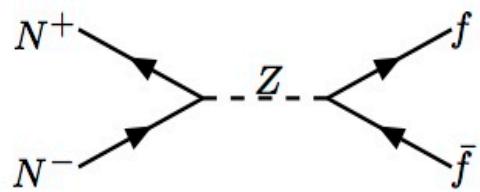
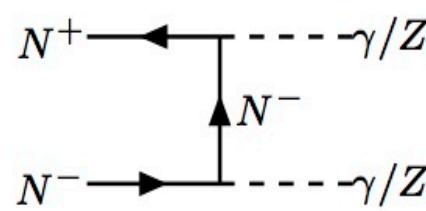
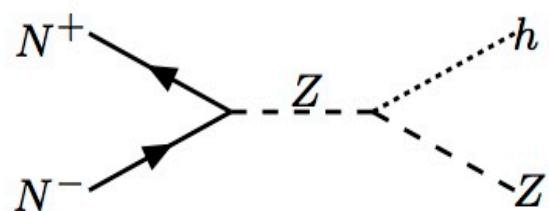
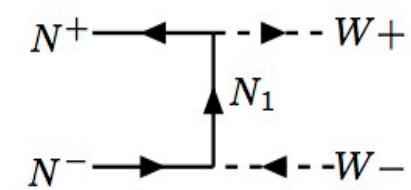
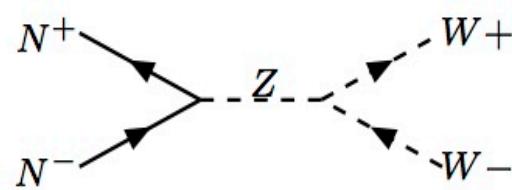
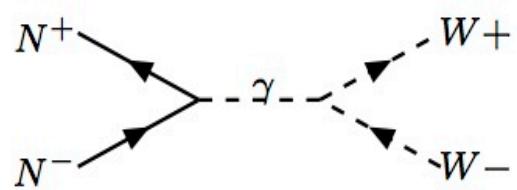
# Annihilation processes

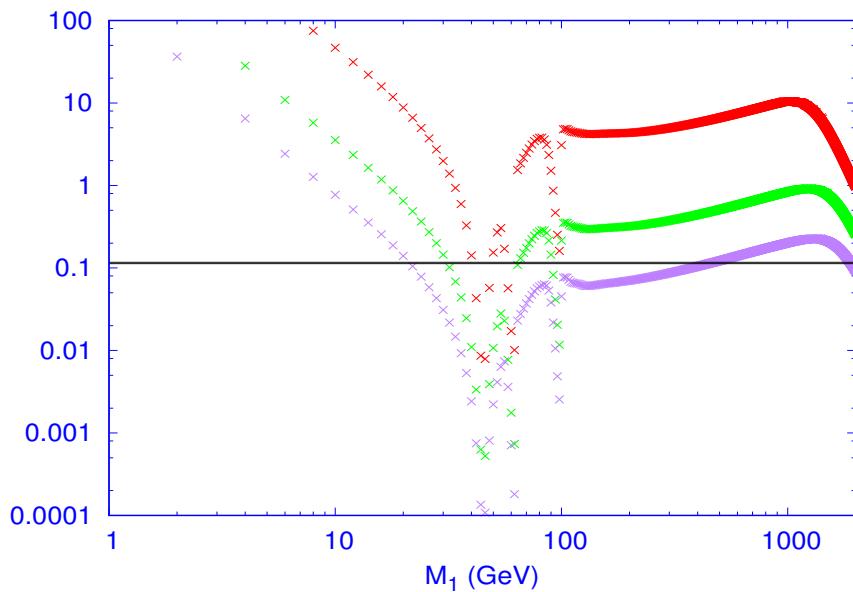


# Dominant Co-annihilation processes

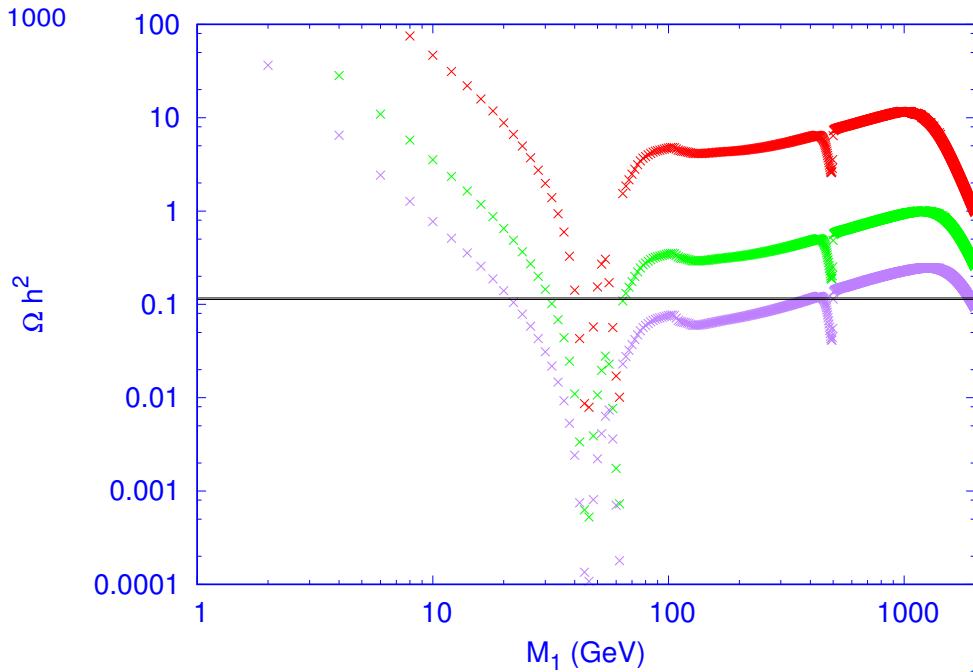




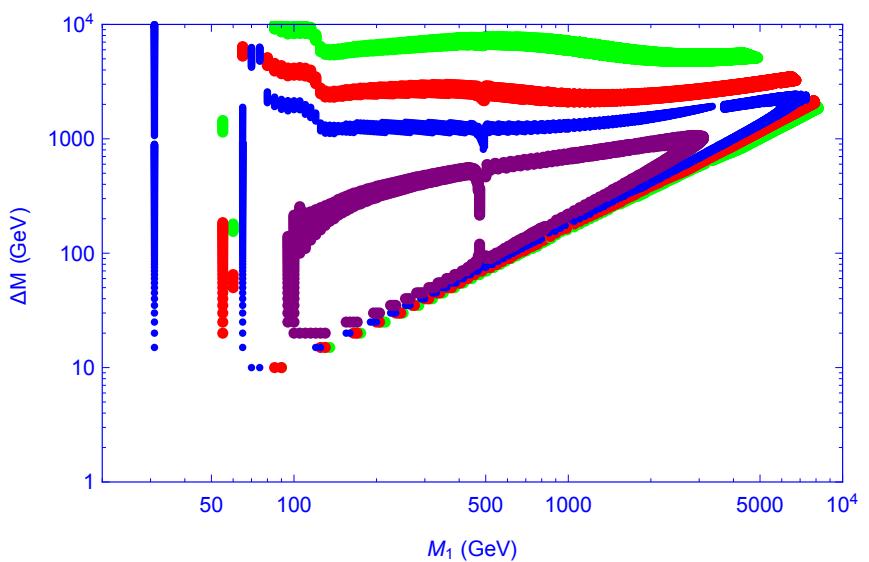




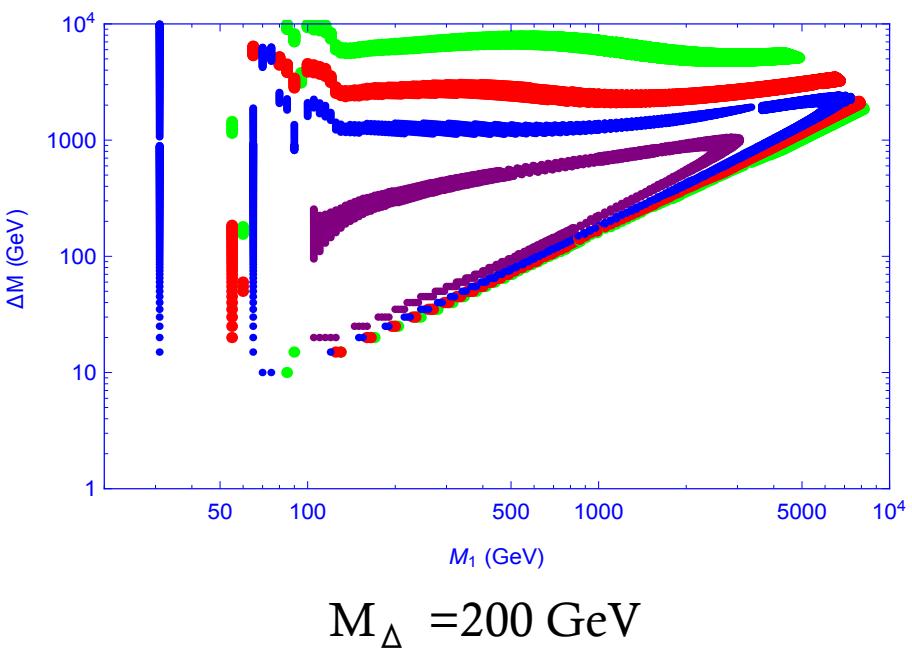
$\sin \theta = 0.1$  (Red)  
 $\sin \theta = 0.2$  (Green)  
 $\sin \theta = 0.3$  (Purple)  
 $f_L/f_N = 0.001$

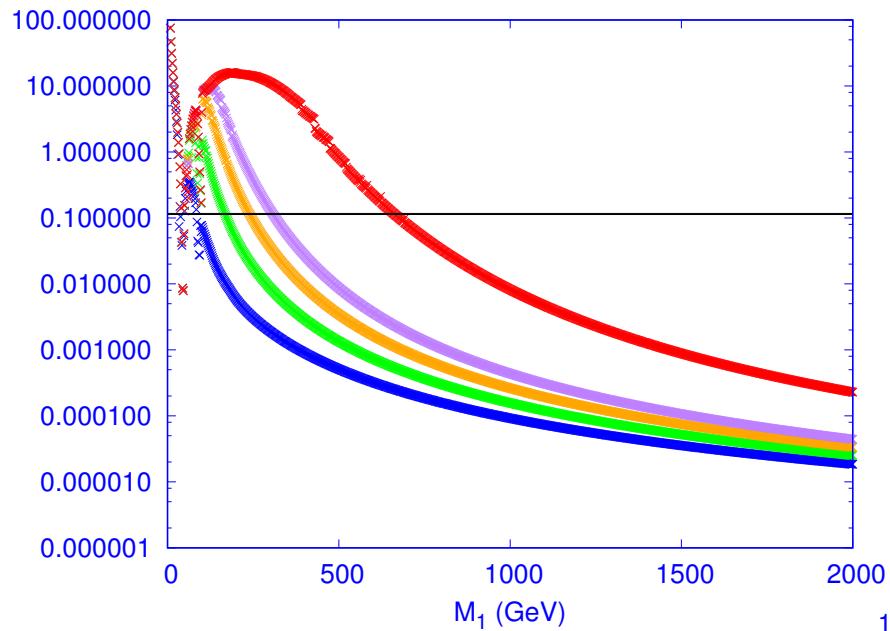


Sin  $\theta = 0.1$  (Green)  
 Sin  $\theta = 0.15$  (Red)  
 Sin  $\theta = 0.2$  (Blue)  
 Sin  $\theta = 0.3$  (Purple)



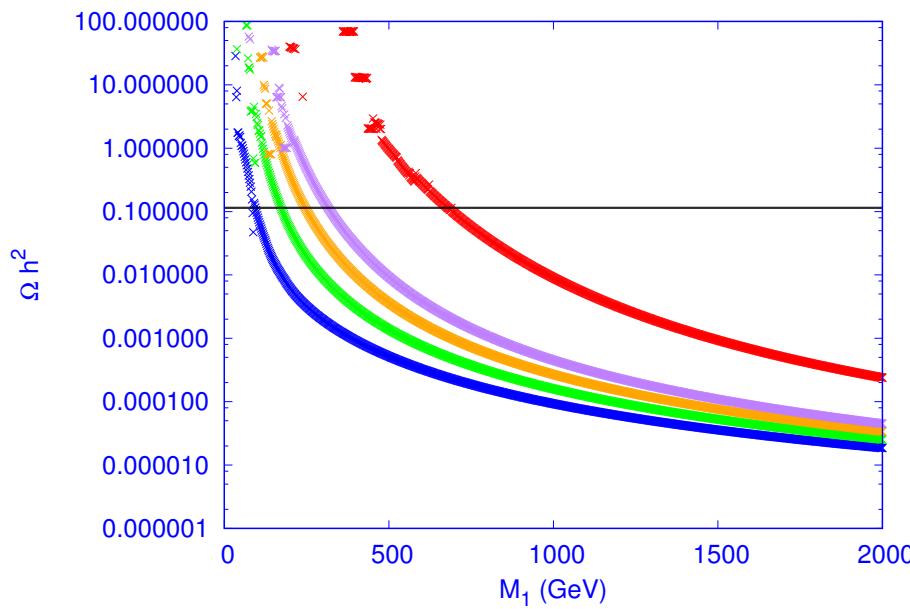
$M_\Delta = 1000$  GeV



$\Omega h^2$ 

$\Delta M = 10$  GeV(Blue)  
 $\Delta M = 20$  GeV(Green)  
 $\Delta M = 30$  GeV(Orange)  
 $\Delta M = 40$  GeV(Purple)  
 $\Delta M = 100$  GeV(Red)

$M_\Delta = 200$  GeV



# Inelastic DM and its Direct Detection

The DM interaction Lagrangian with the nuclei through  $Z$  boson

$$L_{Z-DM} = \bar{N}_1 i(\gamma^\mu \partial_\mu + g_z \gamma^\mu Z_\mu) N_1$$

Where

$$g_Z = \frac{g}{2 \cos \theta_W} \sin^2 \theta$$

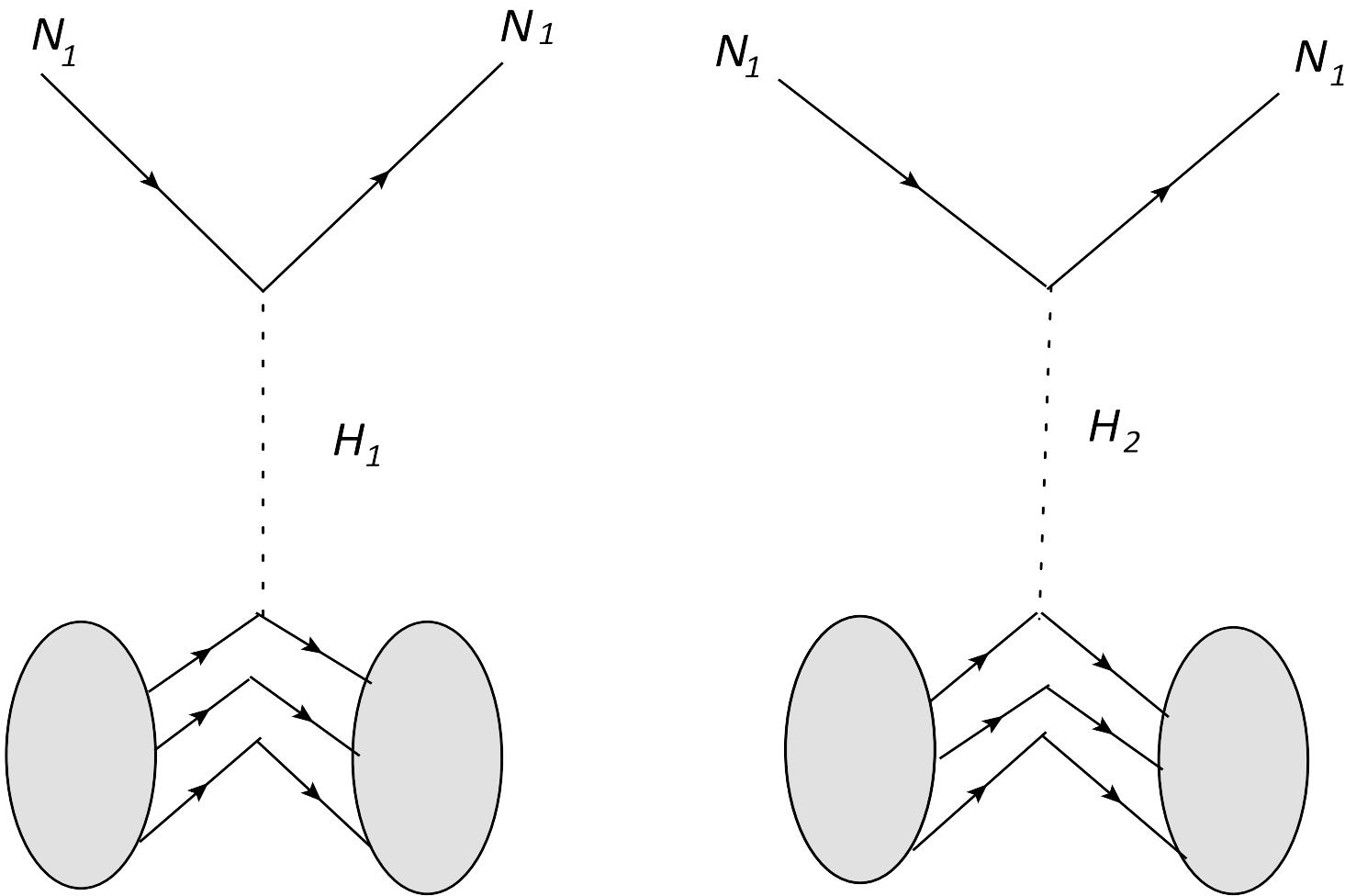
- ☞ Presence of Majorana mass splits the DM into two new states ( $\psi_1$ )<sup>a,b.</sup>
- ☞ Mass splitting  $\delta m = 2\sqrt{2}f_N \sin^2 \theta \langle \Delta \rangle$
- ☞ Dominant  $Z$  interaction now off-diagonal  $\bar{\psi}_1^a i g_z \gamma^\mu Z_\mu \psi_1^b$
- ☞ Minimum velocity required

$$v = c \sqrt{\frac{1}{2m_n E_R}} \left( \frac{m_n E_R}{\mu_R} + \delta M \right)$$

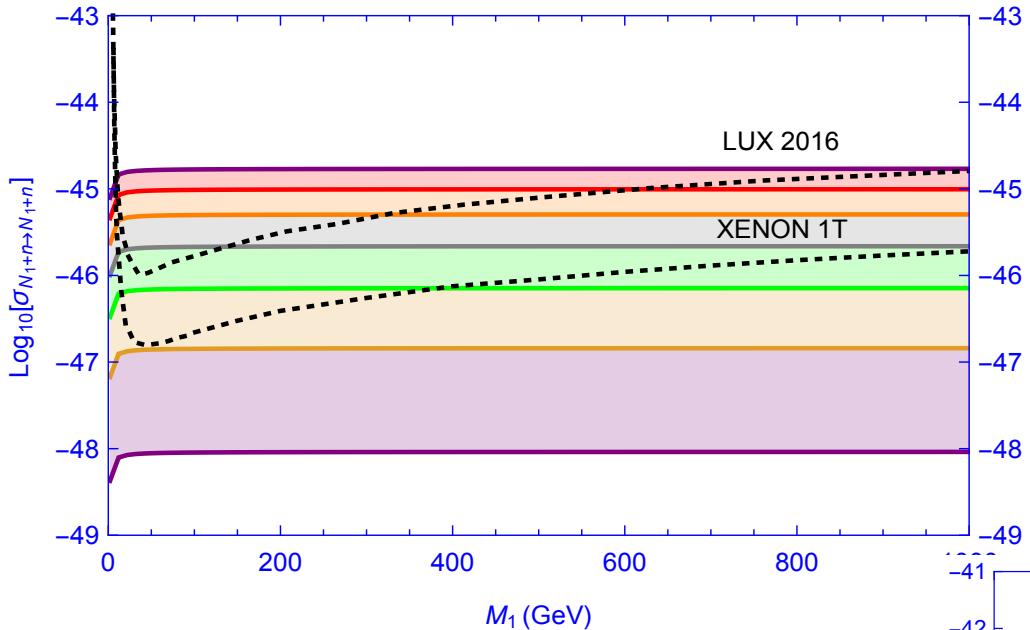
- ✓  $m_n$  = nucleon mass
- ✓  $E_R$  = Recoil Energy
- ✓  $\mu_R$  = Reduced mass



# Elastic Scattering of DM with the nucleus

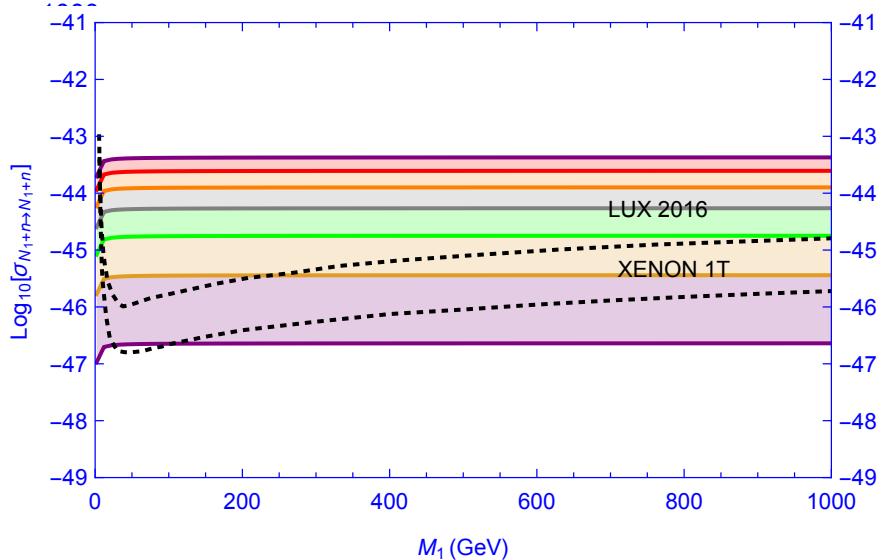


$\Delta M = 100 \text{ GeV}$

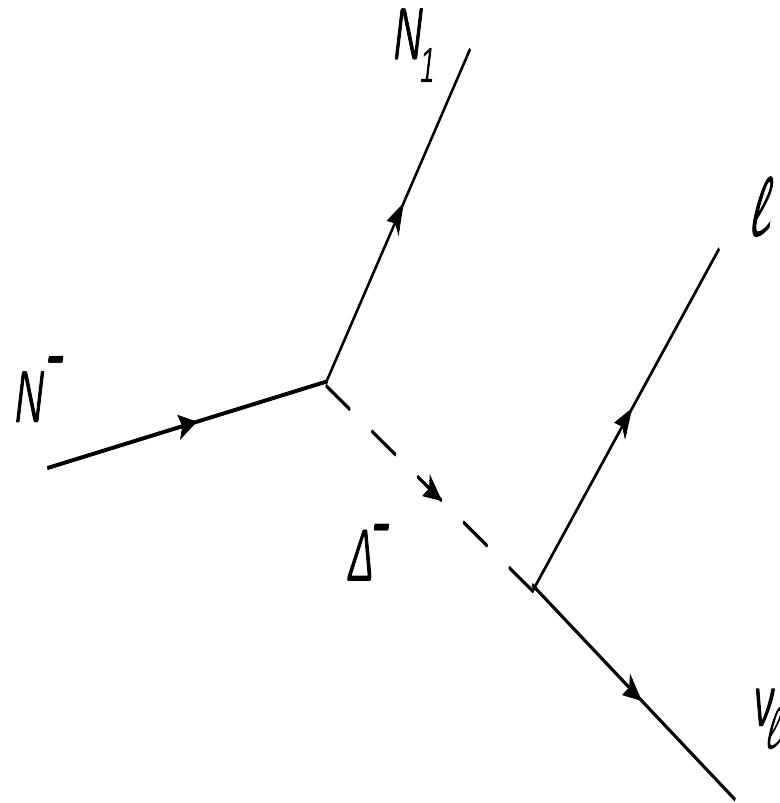
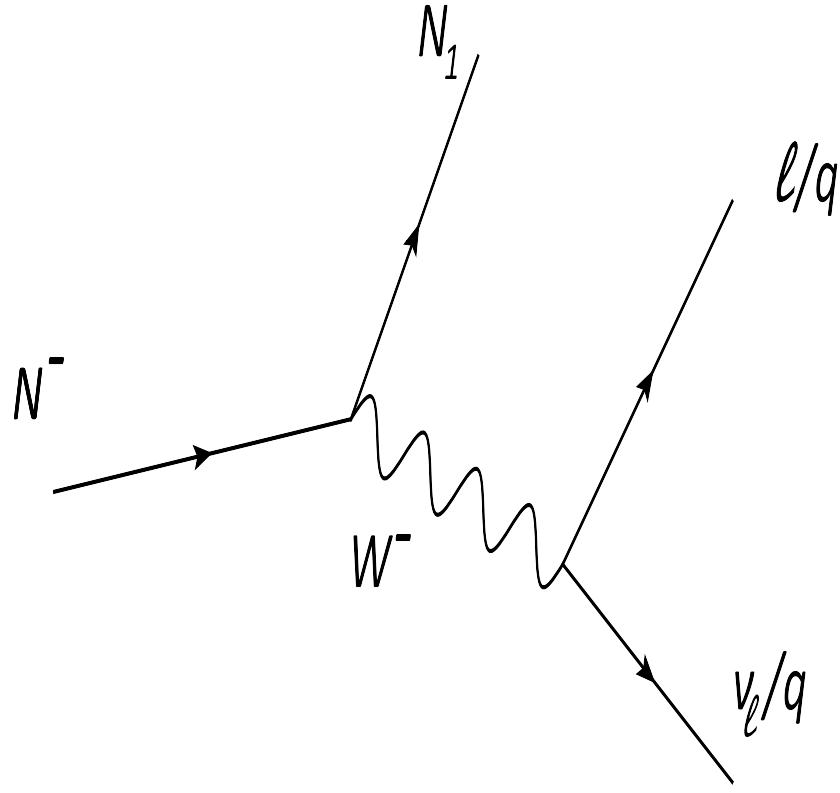


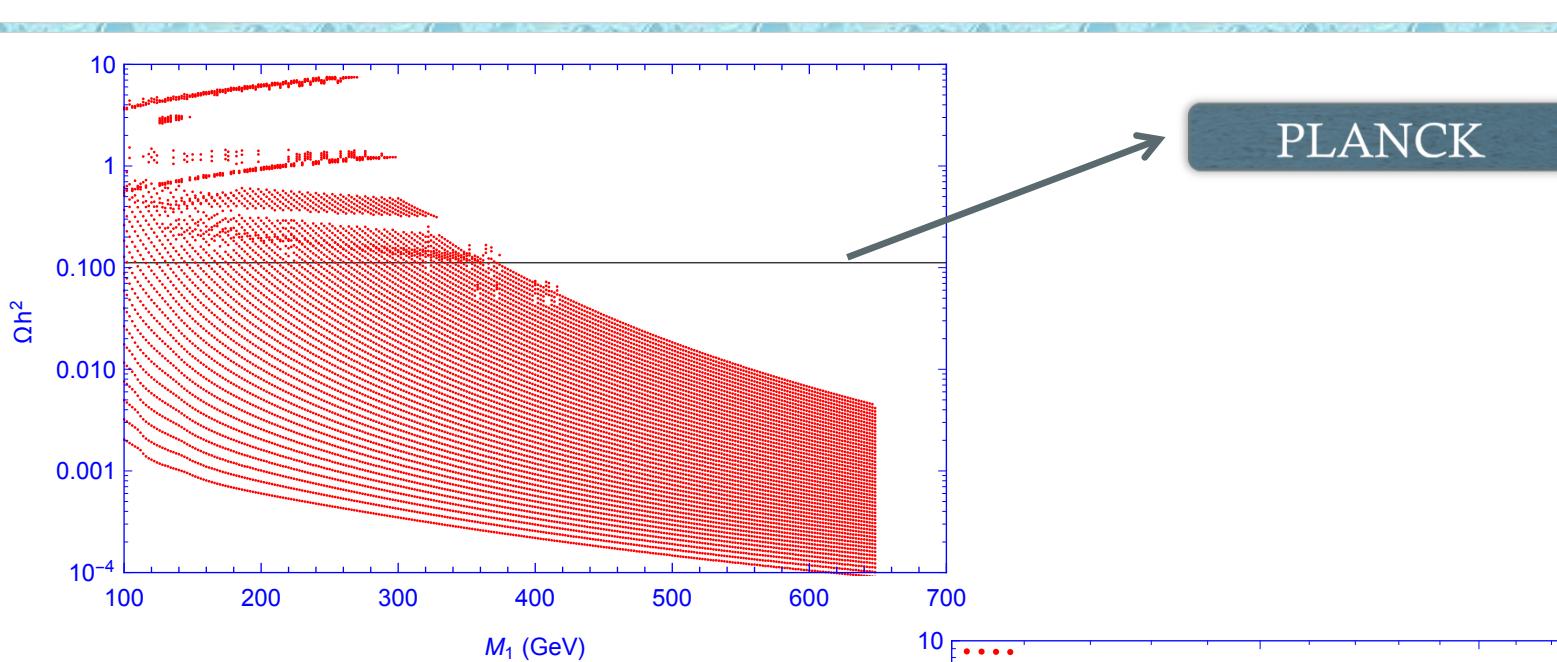
$M_\Delta = 200 \text{ GeV}$

$\Delta M = 500 \text{ GeV}$

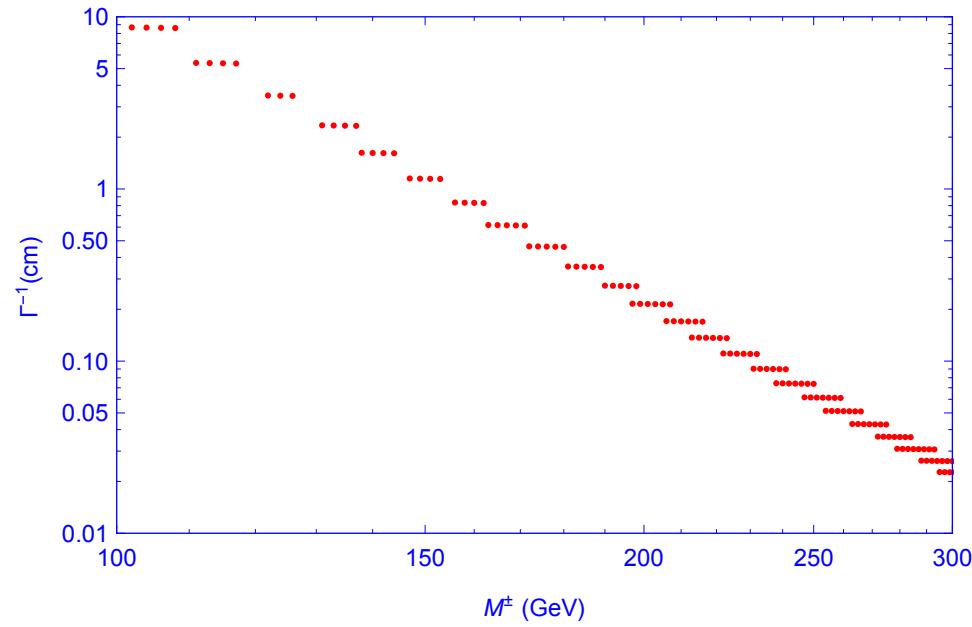


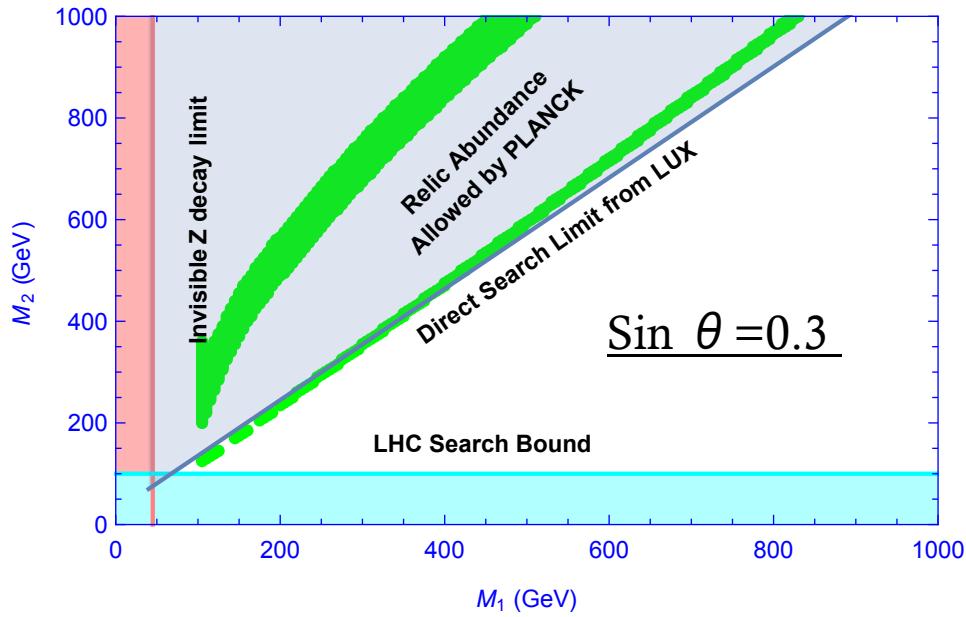
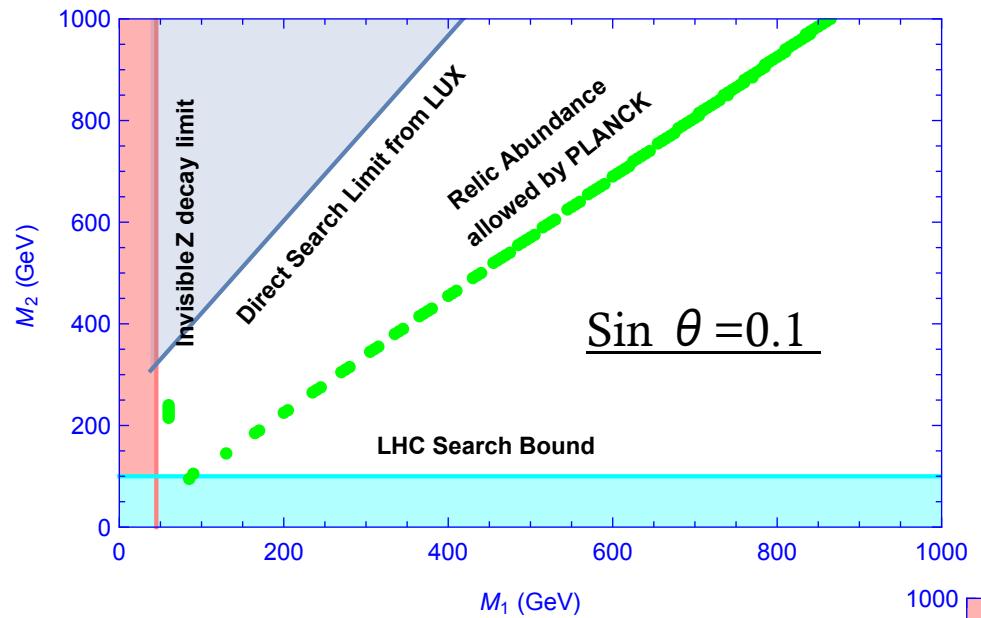
# Decay Of $N^-$ And Displaced Vertex Signature





$\Delta M < 50 \text{ GeV}$   
 $\sin \theta = 0.0003$





# Conclusion

- ★ We have studied the case of a vector like DM model where the DM is a mixed state of doublet and singlet fermion
- ★ Relic abundance is satisfied in large parameter space of the model
- ★ The relic abundance of DM largely dependent on annihilation and coannihilation of DM to the SM gauge bosons.
- ★ The scalar triplet does not contribute to the relic density except for few points near resonance
- ★ The scalar triplet after EW symmetry breaking gives a small Majorana mass to the DM as well as to the SM neutrinos.
- ★ The Majorana mass splits DM into two Majorana mass states with a mass-splitting of sub GeV order. As a result the DM can scatter inelastically from the nucleus through Z mediation.
- ★ But the elastic scattering is still be possible through scalar mediation.
- ★ The charged companion of DM decay via 3 body decay process and can show a displaced vertex.





Thank You



*Back-Up*



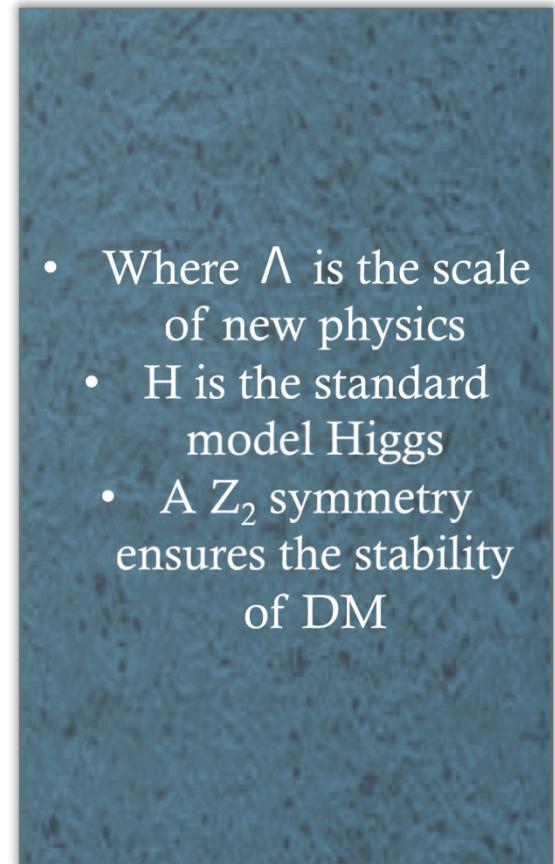
# Singlet Fermion as WIMP

$$L_{DM} = \bar{\chi} i\gamma^\mu \partial_\mu \chi + M \bar{\chi} \chi + \frac{\bar{\chi} \chi H H}{\Lambda}$$

$$\Omega h^2 = F(M_\chi, \Lambda)$$

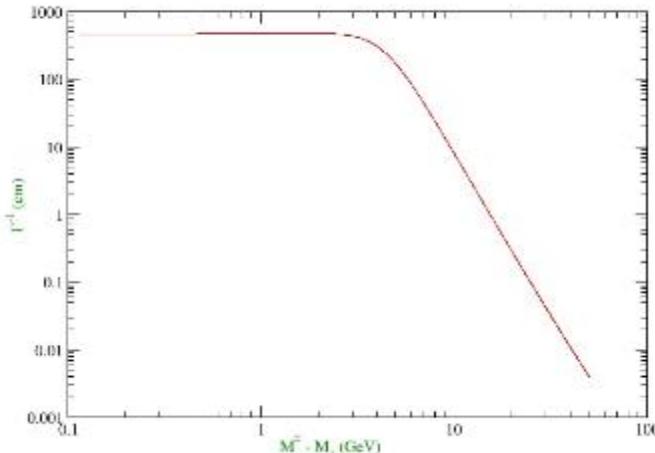
$$\sigma_{SI} = \sigma(M_\chi, \Lambda)$$

- Relic abundance is large for apart from the Higgs funnel.
- We can not do any further prediction about our weak interaction assumption of dark matter



# COLLIDER SEARCH

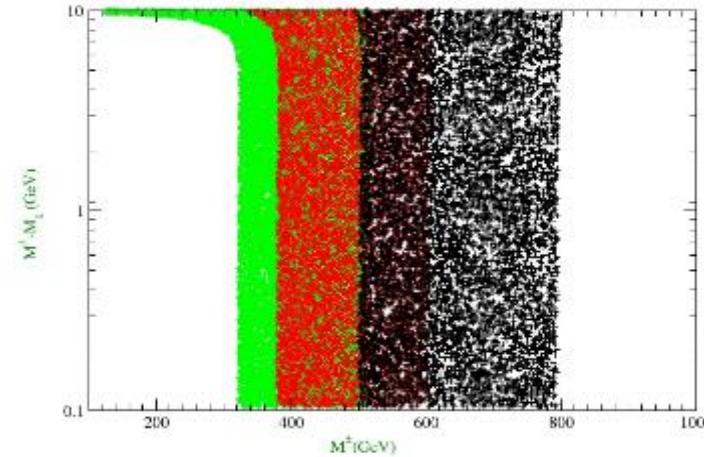
➤ Displaced vertex signature of  $N^\pm$        $N^\pm \rightarrow N_1 + l^\pm + \nu_l$



$$M^\pm = 150\text{GeV}$$

$$\sin \theta = 3 \times 10^{-4}$$

$$m_l = 105\text{MeV}$$

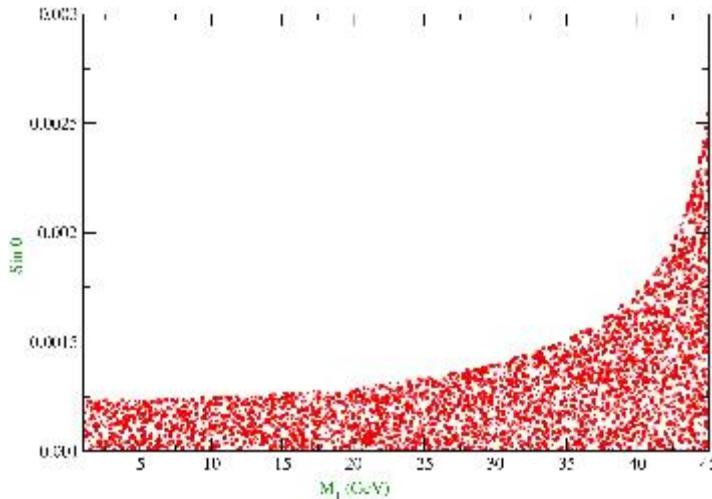


$\sin \theta = 3 \times 10^{-4}$  (Green)  
 $\sin \theta = 2 \times 10^{-4}$  (Red)  
 $\sin \theta = 10^{-4}$  (Black)

For a small mass difference we expect a large displaced vertex for the charge companion of DM



# CONSTRAINT FROM INVISIBLE HIGGS AND Z DECAY



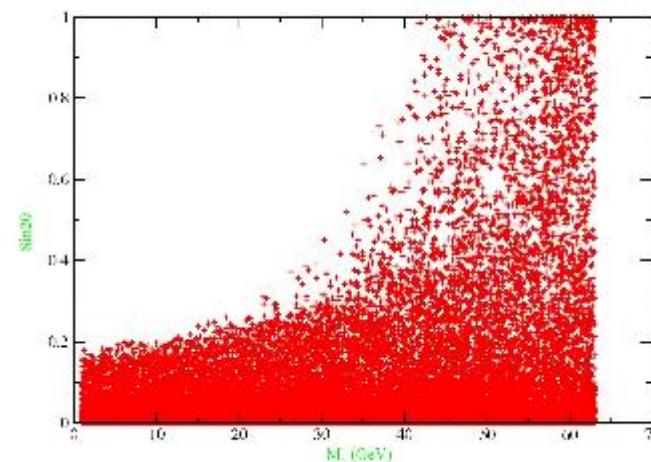
Constraint from Invisible Z Decay : LEP

$$M_1 < M_Z/2 = 45 \text{ GeV}$$
$$M_2 = 45 \text{ GeV}$$

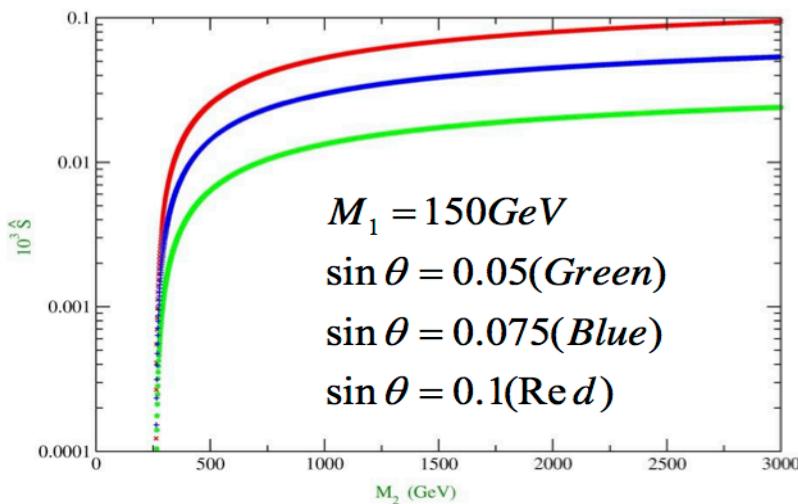
Constraint from Invisible Higgs Decay  
All possible values of  $M_2$  is used  
that opens the decay channel

$$Br_{inv} = \frac{\Gamma_h^{inv}}{\Gamma_h^{SM} + \Gamma_h^{inv}} < 0.3$$

- ATLAS COLLABORATION , [1508.07869 arXiv](https://arxiv.org/abs/1508.07869)

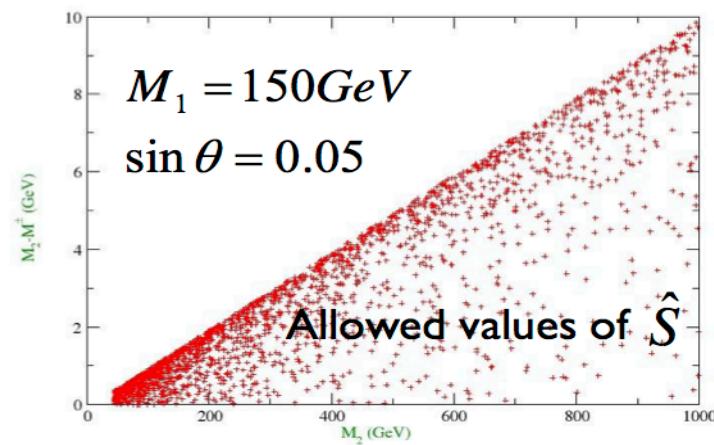


# S parameter

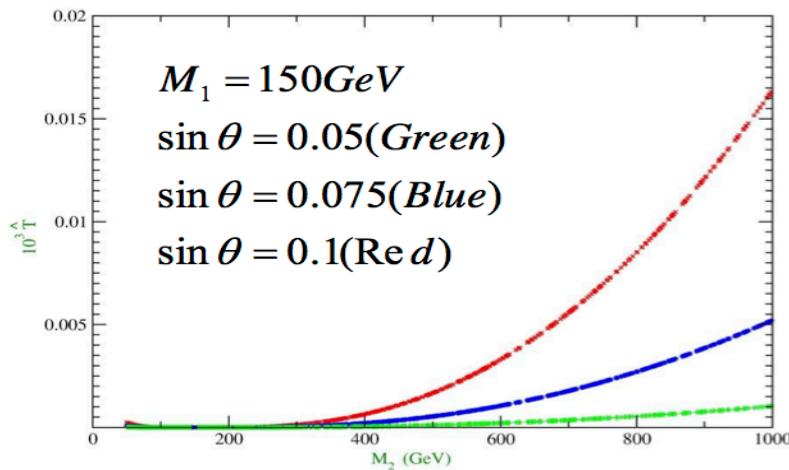


$$\hat{S} = \frac{\alpha S}{4 \sin^2 \theta_w}$$

$$1000\hat{S} = 0.0 \pm 1.3$$



# T parameter



$$\hat{T} = \alpha T$$

$$1000\hat{T} = 0.1 \pm 0.9$$

